## RESONANCE AND REVIVALS I. QUANTUM ROTOR AND INFINITE-WELL DYNAMICS



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So we thought we'd put this revival business to bed! Then...

```
Some Early History of Quantum Revivals
\begin{tabular}{ll} 
J.H. Eberly et.al. Phys. Rev. A 23,236 (1981) & Laser QuantumCavityDynamic revivals \\
R.S. McDowell, WGH, C.W. Patterson LosAlmos Sci. 3, 38(1982) & Symmetric-top revivals \\
S.I. Vetchinkin, et. al. Chem. Phys. Lett. 215,11 (1993) & 1D \(\infty\)-Square well revivals \\
Aronstein, Stroud, Berry, ..., Schleich,.. (1995-1998) & " " " " " \\
WGH, J. Mol. Spectrosc. 210, 166(2001) & Bohr-rotor revivals "
\end{tabular}
So we thought we'd put this revival business to bed! Then this...
More recent story of Quantum Revivals
Anne B. McCoy Chem. Phys. Lett. 501, 603(2011)...reminds me that Morse potential is integer-analytic.
```

Leads to cool Morse revivals in: Following Talk RJ05 by Li:
Resonance\&Revivals II. MORSE OSCILLATOR AND DOUBLE MORSE WELL DYNAMICS.

## So now we're having a revival-revival!

...and, in words by Joannie Mitchell, I find:
"I didn't really know... revivals ... at all."

What do revivals look like? (...in space-time...)


# What do revivals look like? <br> (...in space-time...) 

OK, let's try that again... with<br>quantum<br>revivals...



Observable dynamics of $N$-level-system state $|\Psi\rangle$
Depends on Fourier spectrum of probability distribution $\langle\Psi \mid \Psi\rangle$

$$
|\Psi\rangle=\sum_{n=0}^{N} e^{-i \omega_{n} t} \psi_{n}^{\text {...But individual eigenfrequenci }{ }^{\text {are not directly observable } \ldots} \text {.. }}
$$



# Observable dynamics of $N$-level-system state $|\Psi\rangle$ 

 Depends on Fourier spectrum of probability distribution $\langle\Psi \mid \Psi\rangle$$$
|\Psi\rangle=\sum_{n=0}^{N} e^{-i \omega_{n} t} \psi_{n}^{\quad . . \text { But individual eigenfrequencies }} \text { are not directly observable... }
$$



$$
\omega_{4}
$$

$$
\langle\Psi|=\sum_{m=0}^{N} e^{+i \omega_{m} t} \psi_{m}^{*} \quad \overline{\omega_{3}}
$$

$$
\frac{\frac{\omega_{2}}{\omega_{1}}}{\frac{\omega_{0}}{2}}
$$

Observable dynamics of $N$-level-system state $|\Psi\rangle$ Depends on Fourier spectrum of probability distribution $\langle\Psi \mid \Psi\rangle$

$$
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$$

$$
\frac{\frac{\omega_{2}}{\omega_{1}}}{\frac{\omega_{0}}{2}}
$$

$$
\begin{aligned}
& \left.\begin{array}{rlrl}
\langle\Psi \mid \Psi\rangle & =\sum_{n=0}^{N} e^{i\left(\omega_{m}-\omega_{n}\right) t} \underbrace{\boldsymbol{\psi}_{m}^{*} \psi_{n}} & |\Psi\rangle=\sum_{n=0}^{N} e^{-i \omega_{n} t} \psi_{n} \\
& =\sum_{m, n=0}^{N} e^{i \Delta_{n n} t} & \rho_{m n} & \omega_{0} \omega_{1} \omega_{2} \\
\omega_{3} & \omega_{4} \\
& |\mid l l l l
\end{array} \right\rvert\, \\
& { }^{N} \quad \text {...But individual eigenfrequencies }
\end{aligned}
$$

Observable dynamics of $N$-level-system state $|\Psi\rangle$
Depends on Fourier spectrum of probability distribution $\langle\Psi \mid \Psi\rangle$

$$
\begin{aligned}
& \begin{aligned}
\langle\Psi \mid \Psi\rangle & =\sum_{n=0}^{N} e^{i\left(\omega_{m}-\omega_{n}\right) t} \underbrace{N} \underbrace{*}_{m, n=0} \psi_{n}^{N} \\
& =\sum_{m}^{i \Delta_{m n} t} e^{\boldsymbol{P}_{m n}}
\end{aligned} \\
& { }^{N} \quad \text {...But individual eigenfrequencies } \\
& |\Psi\rangle=\sum_{n=0}^{N} e^{-i \omega_{n} t} \psi_{n} \text { are not directly observable... } \\
& \langle\Psi|=\sum_{m=0}^{N} e^{+i \omega_{m}{ }^{t}} \boldsymbol{\psi}_{m}^{*} \\
& \omega_{3}
\end{aligned}
$$

# Observable dynamics of $N$-level-system state $|\Psi\rangle$ 

Depends on Fourier spectrum of probability distribution $\langle\Psi \mid \Psi\rangle$
 $|\Psi\rangle=\sum_{n=0}^{N} e^{-i \omega_{n} t} \psi_{n \ldots \text { only differences } \Delta_{m n}=\left(\omega_{m}-\omega_{n}\right)}^{\begin{array}{c}\text { But individual eigenfrequencies } \\ \text { are not directly observable }\end{array}}$ $\Delta$

# Observable dynamics of $N$-level-system state $|\Psi\rangle$ 

Depends on Fourier spectrum of probability distribution $\langle\Psi \mid \Psi\rangle$
 $|\Psi\rangle=\sum_{n=0}^{N} e^{-i \omega_{n} t} \boldsymbol{\psi}_{n} \begin{gathered}\text {....only ind diffecerences observable.... } \Delta_{m n} \\ \text { are not ritual eigenfrequencies }\end{gathered}$ $\Delta_{40} \Delta_{41}$


$$
\begin{aligned}
& \mathrm{N}=5 \text { eigenfrequencies } \\
& \text { positive } \Delta_{m>n}=\omega_{m}-\omega_{n} \\
& \text { and : } N(N-1) / 2=10 \\
& \text { negative } \Delta_{m<n}=\omega_{m}-\omega_{n}
\end{aligned}
$$

# Observable dynamics of $N$-level-system state $|\Psi\rangle$ 

Depends on Fourier spectrum of probability distribution $\langle\Psi \mid \Psi\rangle$

$$
N \quad \text {...But individual eigenfrequencies }
$$

$$
\begin{aligned}
& \begin{aligned}
\langle\Psi \mid \Psi\rangle & =\sum_{n=0}^{N} e^{i\left(\omega_{m}-\omega_{n}\right) t} \underbrace{\psi_{m}^{*} \psi_{n}} \\
& =\sum_{m, n=0}^{N} e^{i \Delta_{m n} t} \rho_{m n}
\end{aligned} \\
& |\Psi\rangle=\sum e^{-i \omega_{n} t} \boldsymbol{\psi} \text { are not directly observable... } \\
& \ldots \text {...only differences } \Delta_{m n}=\left(\omega_{m}-\omega_{n}\right) \\
& \langle\Psi|=\sum_{m=0}^{N} e^{+i \omega_{m} t} \psi_{m}^{*} \quad \omega_{3} \\
& \frac{\frac{\omega_{2}}{\omega_{1}}}{\frac{\omega_{0}}{}} \\
& \begin{array}{l}
\text { Nines: } N(N-1) / 2=10 \\
\text { positive } \Delta_{m>n}=\omega_{m}-\omega_{n} \\
\text { and: } N(N-1) / 2=10 \\
\text { negative } \Delta_{m<n}=\omega_{m}-\omega_{n} \\
\hline \text { beats }\left|\Delta_{m<n}\right|=\left|\omega_{m}-\omega_{n}\right|
\end{array}
\end{aligned}
$$

Observable dynamics of 2 -level-system state $|\Psi\rangle$ Fourier spectrum of $\langle\Psi \mid \Psi\rangle$ has $O N E$ beat frequency $\Delta_{21}=-\Delta_{12}$


$$
\langle\Psi|=\sum_{m=0}^{N} e^{+i \omega_{m} t} \psi_{m}^{*}
$$

$\frac{\omega_{2}}{\omega_{1}}$

$$
|\Psi\rangle=\sum_{n=0}^{N} e^{-i \omega_{n} t} \psi_{n}
$$

$$
\Delta_{21}
$$


$\mathrm{N}=2$ eigenfrequencies gives: $N(N-1) / 2=1$ positive $\Delta_{2>1}=\omega_{2}-\omega_{1}$ and : $N(N-1) / 2=1$
$\frac{\text { negative } \Delta_{1<2}=\omega_{1}-\omega_{2}}{\text { beat }\left|\Delta_{m<n}\right|=\left|\omega_{m}-\omega_{n}\right|}$

## 2-level-system and $C_{2}$ symmetry beat dynamics


$\mathrm{C}_{2}$ Character Table describes eigenstates


## 2-level-system and $C_{2}$ symmetry beat dynamics

$\mathrm{C}_{2}$ Phasor-Character Table

$\mathrm{C}_{2}$ Character Table describes eigenstates


Phasor $\mathrm{C}_{2}$ Characters describe local state beats

Initial sum

## 2-level-system and $C_{2}$ symmetry beat dynamics

$\mathrm{C}_{2}$ Phasor-Character Table

antisymmetric $\mathrm{A}_{2}$

|  | $1=r^{0}$ | $r=r^{1}$ |
| :---: | :---: | :---: |
| $0 \bmod 2$ | 1 | 1 |
| $\pm 1 \bmod 2$ | 1 | -1 |

Phasor $\mathrm{C}_{2}$ Characters describe local state beats

Initial sum

1/4-beat

2-level-system and $C_{2}$ symmetry beat dynamics
$\mathrm{C}_{2}$ Phasor-Character Table

$\mathrm{C}_{2}$ Character Table describes eigenstates
symmetric $\mathrm{A}_{1}$
antisymmetric $\mathrm{A}_{2}$

Phasor $\mathrm{C}_{2}$ Characters describe local state beats
Initial sum

1/4-beat

1/2-beat


## 2-level-system and $C_{2}$ symmetry beat dynamics

$\mathrm{C}_{2}$ Phasor-Character Table


Coupled Optical Pendula $E(t)$ even $+45^{\circ}$
王
$\mathrm{C}_{2}$
parity

symmetric $\mathrm{A}_{1}$
VS.
antisymmetric $\mathrm{A}_{2}$

Phasor $\mathrm{C}_{2}$ Characters describe local state beats

Initial sum

1/4-beat
1/2-beat

3/4-beat


2-level-system and $C_{2}$ symmetry beat dynamics
$\mathrm{C}_{2}$ Phasor-Character Table


Coupled Optical Pendula E(t)

H>

$+\cdots$



What do revivals look like?
...in per-space-time...
(... that is:
frequency $\omega_{m}{ }^{\text {radiansece. }}$
VS
$k$-vector $\left.k_{m \text { radian/cm }}\right)$

## $N$-level-system and revival-beat wave dynamics



## $N$-level-system and revival-beat wave dynamics

Possible wave velocities for
Quadratic (Bohr-Rotor) Spectrum

$$
\begin{gathered}
\omega_{m}=B m^{2} \\
k_{m}= \pm m
\end{gathered}
$$

$$
V_{\text {phase }}=\frac{\omega_{m}}{k_{m}}=\frac{B m^{2}}{m}
$$

$$
=m B
$$



## $N$-level-system and revival-beat wave dynamics

Possible wave velocities for
Quadratic (Bohr-Rotor) Spectrum

$$
\begin{gathered}
\omega_{m}=B m^{2} \\
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\end{gathered}
$$

$$
\begin{aligned}
V_{\text {phase }}=\frac{\omega_{m}}{k_{m}}=\frac{B m^{2}}{m} \\
=m B
\end{aligned} \quad \begin{array}{r}
\text { Vgroup }= \\
=(m \pm n) B
\end{array}
$$



## $N$-level-system and revival-beat wave dynamics



Harmonic Oscillator level spectrum contains the Rotor Levels as a subset
$N$-level-system and revival-beat wave dynamics
(Just 2-levels $(0, \pm 1)$ (and some $\pm 2$ ) excited)
1/1
$3 / 4$

$1 / 2$
$1 / 4$
$N$-level-system and revival-beat wave dynamics
(Just 2-levels $(0, \pm 1)$ (and some $\pm 2$ ) excited)
(4-levels $(0, \pm 1, \pm 2, \pm 3)$ (and some $\pm 4$ ) excited)


## $N$-level-system and revival-beat wave dynamics

Zeros (clearly) and "particle-packets" (faintly) have paths labeled by fraction sequences like: $\frac{0}{7}, \frac{1}{7}, \frac{2}{7}, \frac{3}{7}, \frac{4}{7}, \frac{5}{7}, \frac{6}{7}, \frac{1}{1}$

Time $t$ e $t / 1$ (9 or10-levels $(0, \pm 1, \pm 2, \pm 3, \pm 4, \ldots, \pm 9, \pm 10, \pm 11 . .$.$) excited)$


Farey Sum algebra of revival-beat wave dynamics Label by numerators $N$ and denominators $D$ of rational fractions $N / D$


Farey Sum algebra of revival-beat wave dynamics Label by numerators $N$ and denominators $D$ of rational fractions $N / D$


# A Lesson in Rational Fractions N/D <br> (...that you can take home for your kids!) 







Farey Sum related to vector sum and Ford Circles

1/2-circle has diameter $1 / 2^{2}=1 / 4$
$1 / 3$-circles have diameter $1 / 3^{2}=1 / 9$
$\mathrm{n} / \mathrm{d}$-circles have diameter $1 / d^{2}$

## $C_{m}$ algebra of revival-phase dynamics

## Quantum rotor fractional take turns at Cn symmetry

## $C_{m}$ algebra of revival-phase dynamics

Discrete 3-State or Trigonal System (Tesla's 3-Phase AC)


Discrete 6-State or Hexagonal System (6-Phase AC)
$C_{6}$ Eigenstate Characters


## Summary

Quantum rotor revivals obey wonderfully simple geometry, number, and group theoretical analysis and
as the next talk will show...

## Summary

Quantum rotor revivals obey wonderfully simple geometry, number, and group theoretical analysis and as the next talk will show...
"I still don't really know... revivals ... at all."

## Simulation of revival-intensity dynamics




