AMOP reference links on following page 2.05.18 class 7.0: Symmetry Principles for Advanced Atomic-Molecular-Optical-Physics William G. Harter - University of Arkansas

Symmetry group  $\mathcal{G}=U(1)$  representations, 1D HO Hamiltonian  $\mathbf{H}=\hbar\omega \mathbf{a}^{\dagger}\mathbf{a}$  operators, 1D HO wave eigenfunctions  $\Psi_n$ , and coherent  $\alpha$ -states

Factoring 1D-HO Hamiltonian  $H=p^{2}+x^{2}$ Creation-Destruction at a algebra of U(1) operatorss Eigenstate creationism (and destructionism) Vacuum state |0>,  $1^{st}$  excited state |1>, |2>, ...

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#### AMOP reference links (Updated list given on 2nd page of each class presentation)

Web Resources - front page

#### <u>2014 AMOP</u>

2017 Group Theory for QM

#### UAF Physics UTube channel

2018 AMOP

Frame Transformation Relations And Multipole Transitions In Symmetric Polyatomic Molecules - RMP-1978 (Alt Scanned version)

Rotational energy surfaces and high- J eigenvalue structure of polyatomic molecules - Harter - Patterson - 1984

Galloping waves and their relativistic properties - ajp-1985-Harter

Asymptotic eigensolutions of fourth and sixth rank octahedral tensor operators - Harter-Patterson-JMP-1979

Nuclear spin weights and gas phase spectral structure of 12C60 and 13C60 buckminsterfullerene -Harter-Reimer-Cpl-1992 - (Alt1, Alt2 Erratum)

Theory of hyperfine and superfine levels in symmetric polyatomic molecules.

I) Trigonal and tetrahedral molecules: Elementary spin-1/2 cases in vibronic ground states - PRA-1979-Harter-Patterson (Alt scan)

II) <u>Elementary cases in octahedral hexafluoride molecules - Harter-PRA-1981 (Alt scan)</u>

Rotation-vibration scalar coupling zeta coefficients and spectroscopic band shapes of buckminsterfullerene - Weeks-Harter-CPL-1991 (Alt scan) Fullerene symmetry reduction and rotational level fine structure/ the Buckyball isotopomer 12C 13C59 - jcp-Reimer-Harter-1997 (HiRez) Molecular Eigensolution Symmetry Analysis and Fine Structure - IJMS-harter-mitchell-2013

Rotation-vibration spectra of icosahedral molecules.

I) Icosahedral symmetry analysis and fine structure - harter-weeks-jcp-1989

II) Icosahedral symmetry, vibrational eigenfrequencies, and normal modes of buckminsterfullerene - weeks-harter-jcp-1989

III) Half-integral angular momentum - harter-reimer-jcp-1991

QTCA Unit 10 Ch 30 - 2013

AMOP Ch 32 Molecular Symmetry and Dynamics - 2019

AMOP Ch 0 Space-Time Symmetry - 2019

#### RESONANCE AND REVIVALS

I) QUANTUM ROTOR AND INFINITE-WELL DYNAMICS - ISMSLi2012 (Talk) https://kb.osu.edu/dspace/handle/1811/52324

- II) Comparing Half-integer Spin and Integer Spin Alva-ISMS-Ohio2013-R777 (Talks)
- III) Quantum Resonant Beats and Revivals in the Morse Oscillators and Rotors (2013-Li-Diss)

Rovibrational Spectral Fine Structure Of Icosahedral Molecules - Cpl 1986 (Alt Scan)

Gas Phase Level Structure of C60 Buckyball and Derivatives Exhibiting Broken Icosahedral Symmetry - reimer-diss-1996 Resonance and Revivals in Quantum Rotors - Comparing Half-integer Spin and Integer Spin - Alva-ISMS-Ohio2013-R777 (Talk) Quantum Revivals of Morse Oscillators and Farey-Ford Geometry - Li-Harter-cpl-2013 Wave Node Dynamics and Revival Symmetry in Quantum Rotors - harter - jms - 2001 Symmetry group  $\mathcal{G}=U(1)$  representations, 1D HO Hamiltonian  $\mathbf{H}=\hbar\omega \mathbf{a}^{\dagger}\mathbf{a}$  operators, 1D HO wave eigenfunctions  $\Psi_n$ , and coherent  $\alpha$ -states

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Define Destruction operator and Creation Operator
$$\mathbf{Creation \ Operator}$$
Commutation relations between  $\mathbf{a} = (\mathbf{X} + i\mathbf{P})/2$  and  $\mathbf{a}^{\dagger} = (\mathbf{X} - i\mathbf{P})/2$  with  $\mathbf{X} \equiv \sqrt{M\omega}\mathbf{x}/\sqrt{2}$  and  $\mathbf{P} \equiv \mathbf{p}/\sqrt{2M}$ :
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Recall *commutator* [x, p] *relation*:  $[x, p] = xp - px = \hbar i 1$ 

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1D-HO Hamiltonian in terms of **a**<sup>†</sup>**a** operator

Recall:  $\mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar\omega (\mathbf{a}^{\dagger}\mathbf{a} + \mathbf{a}\mathbf{a}^{\dagger})/2$ 

Recall *commutator*  $[\mathbf{x}, \mathbf{p}]$  *relation*:  $[\mathbf{x}, \mathbf{p}] = \mathbf{x}\mathbf{p} - \mathbf{p}\mathbf{x} = \hbar i \mathbf{1}$ 



Recall *commutator* [x, p] *relation*:  $[x, p] = xp - px = \hbar i \mathbf{1}$ 



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Symmetry group  $\mathcal{G}=U(1)$  representations, 1D HO Hamiltonian  $\mathbf{H}=\hbar\omega \mathbf{a}^{\dagger}\mathbf{a}$  operators, 1D HO wave eigenfunctions  $\Psi_n$ , and coherent  $\alpha$ -states

Factoring 1D-HO Hamiltonian  $\mathbf{H}=\mathbf{p}^{2}+\mathbf{x}^{2}$ Creation-Destruction  $\mathbf{a}^{\dagger}\mathbf{a}$  algebra of U(1) operators Eigenstate creationism (and destructionism) Vacuum state |0>,  $1^{st}$  excited state |1>, |2>, ...

Normal ordering for matrix calculation (creation **a**<sup>†</sup> on left, destruction **a** on right) Commutator derivative identities, Binomial expansion identities

 $\langle a^n a^{\dagger n} \rangle$  operator calculations Number operator and Hamiltonian operator Expectation values of position, momentum, and uncertainty for eigenstate  $|n\rangle$ Harmonic oscillator beat dynamics of mixed states

*Oscillator coherent states ("Shoved" and "kicked" states) Translation operators vs. boost operators, boost-translation combinations* 

*Time evolution of a coherent state*  $|\alpha\rangle$ *Properties of coherent states and "squeezed" states*  *Eigenstate creationism (and destructionism) Given1D-HO Hamiltonian:*  $(\mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar \omega \mathbf{a}^{\dagger} \mathbf{a} + \mathbf{1}\hbar \omega/2)$  and commutation:  $([\mathbf{a},\mathbf{a}^{\dagger}] = \mathbf{1})$  or  $(\mathbf{a}\mathbf{a}^{\dagger} = \mathbf{a}^{\dagger}\mathbf{a} + \mathbf{1})$ 

Define ground state  $|0\rangle$  as the eigenstate of  $\mathbf{H}(\mathbf{x},\mathbf{p})$  with the zero point eigenvalue  $E_0 = \hbar \omega/2$ .

*Eigenstate creationism (and destructionism) Given1D-HO Hamiltonian:*  $(\mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar \omega \mathbf{a}^{\dagger} \mathbf{a} + \mathbf{1}\hbar \omega/2)$  and commutation:  $([\mathbf{a},\mathbf{a}^{\dagger}] = \mathbf{1})$  or  $(\mathbf{a}\mathbf{a}^{\dagger} = \mathbf{a}^{\dagger}\mathbf{a} + \mathbf{1})$ 

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 $\mathbf{H}(\mathbf{x},\mathbf{p}) |0\rangle = \hbar\omega/2 |0\rangle \qquad \langle 0| \mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar\omega/2 \langle 0|$ 

*Eigenstate creationism (and destructionism) Given1D-HO Hamiltonian:*  $(\mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar \omega \mathbf{a}^{\dagger} \mathbf{a} + \mathbf{1}\hbar \omega/2)$  and commutation:  $([\mathbf{a},\mathbf{a}^{\dagger}] = \mathbf{1})$  or  $(\mathbf{a}\mathbf{a}^{\dagger} = \mathbf{a}^{\dagger}\mathbf{a} + \mathbf{1})$ 

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Action by **a** on ground ket  $|0\rangle$  (or **a**<sup>†</sup> on ground bra  $\langle 0|$ ) gives *nothing* (zero vectors  $\boldsymbol{\theta}$ ).

$$\mathbf{a} |0\rangle = \boldsymbol{\theta} \qquad \qquad \langle 0| \ \mathbf{a}^{\dagger} = \boldsymbol{\theta}$$

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Eigenstate creationism (and destructionism) Given 1D-HO Hamiltonian:  $(\mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar \omega \mathbf{a}^{\dagger} \mathbf{a} + \mathbf{1}\hbar \omega/2)$  and commutation:  $([\mathbf{a},\mathbf{a}^{\dagger}] = \mathbf{1})$  or  $(\mathbf{a}\mathbf{a}^{\dagger} = \mathbf{a}^{\dagger}\mathbf{a} + \mathbf{1})$ Define ground state  $|0\rangle$  as the eigenstate of  $\mathbf{H}(\mathbf{x},\mathbf{p})$  with the zero point eigenvalue  $E_0 = \hbar \omega/2$ .  $\mathbf{H}(\mathbf{x},\mathbf{p}) |0\rangle = \hbar \omega/2 |0\rangle$   $\langle 0| \mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar \omega/2 \langle 0|$ Action by **a** on ground ket  $|0\rangle$  (or **a**<sup>†</sup> on ground bra  $\langle 0|$ ) gives *nothing* (zero vectors  $\boldsymbol{\theta}$ ).  $\mathbf{a} |0\rangle = \boldsymbol{\theta} \qquad \qquad \langle 0| \ \mathbf{a}^{\dagger} = \boldsymbol{\theta}$ But,  $\mathbf{a}^{\dagger}$  acts on ground ket to give  $|1\rangle = \mathbf{a}^{\dagger}|0\rangle$  with  $\mathbf{H}$  eigenvalue  $E_{I} = \hbar\omega + E_{0}$ . ( $|1\rangle = \mathbf{a}^{\dagger}|0\rangle$ ,  $\langle 0|\mathbf{a} = \langle 1|.\rangle$ *Proof:*  $\mathbf{H}(\mathbf{x},\mathbf{p}) \mathbf{a}^{\dagger}|0\rangle = \hbar\omega \mathbf{a}^{\dagger}\mathbf{a} \mathbf{a}^{\dagger}|0\rangle + \hbar\omega/2 \mathbf{a}^{\dagger}|0\rangle$  $\mathbf{H}(\mathbf{x},\mathbf{p}) \mathbf{a}^{\dagger}|0\rangle = \hbar\omega \mathbf{a}^{\dagger}(\mathbf{a}^{\dagger}\mathbf{a} + \mathbf{1})|0\rangle + \hbar\omega/2 \mathbf{a}^{\dagger}|0\rangle$ 

Eigenstate creationism (and destructionism) Given 1D-HO Hamiltonian:  $(\mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar \omega \mathbf{a}^{\dagger} \mathbf{a} + \mathbf{1}\hbar \omega/2)$  and commutation:  $([\mathbf{a},\mathbf{a}^{\dagger}] = \mathbf{1})$  or  $(\mathbf{a}\mathbf{a}^{\dagger} = \mathbf{a}^{\dagger}\mathbf{a} + \mathbf{1})$ Define ground state  $|0\rangle$  as the eigenstate of  $\mathbf{H}(\mathbf{x},\mathbf{p})$  with the zero point eigenvalue  $E_0 = \hbar \omega/2$ .  $\mathbf{H}(\mathbf{x},\mathbf{p}) |0\rangle = \hbar \omega/2 |0\rangle$   $\langle 0| \mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar \omega/2 \langle 0|$ Action by **a** on ground ket  $|0\rangle$  (or **a**<sup>†</sup> on ground bra  $\langle 0|$ ) gives *nothing* (zero vectors  $\boldsymbol{\theta}$ ).  $\mathbf{a} |0\rangle = \mathbf{0}$   $\langle 0| \mathbf{a}^{\dagger} = \mathbf{0}$ But,  $\mathbf{a}^{\dagger}$  acts on ground ket to give  $|1\rangle = \mathbf{a}^{\dagger}|0\rangle$  with  $\mathbf{H}$  eigenvalue  $E_1 = \hbar\omega + E_0$ .  $(|1\rangle = \mathbf{a}^{\dagger}|0\rangle, \langle 0|\mathbf{a} = \langle 1|.\rangle$ Proof:  $\mathbf{H}(\mathbf{x},\mathbf{p}) \mathbf{a}^{\dagger}|0\rangle = \hbar\omega \mathbf{a}^{\dagger}\mathbf{a} \mathbf{a}^{\dagger}|0\rangle + \hbar\omega/2 \mathbf{a}^{\dagger}|0\rangle$  $\mathbf{H}(\mathbf{x},\mathbf{p}) \mathbf{a}^{\dagger}|0\rangle = \hbar\omega \mathbf{a}^{\dagger}(\mathbf{a}^{\dagger}\mathbf{a} + \mathbf{1})|0\rangle + \hbar\omega/2 \mathbf{a}^{\dagger}|0\rangle$ =  $\hbar\omega \mathbf{a}^{\dagger}|0\rangle + \mathbf{0} + \hbar\omega/2 \mathbf{a}^{\dagger}|0\rangle$ 

Eigenstate creationism (and destructionism) Given 1D-HO Hamiltonian:  $(\mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar \omega \mathbf{a}^{\dagger} \mathbf{a} + \mathbf{1}\hbar \omega/2)$  and commutation:  $([\mathbf{a},\mathbf{a}^{\dagger}] = \mathbf{1})$  or  $(\mathbf{a}\mathbf{a}^{\dagger} = \mathbf{a}^{\dagger}\mathbf{a} + \mathbf{1})$ ..... Define ground state  $|0\rangle$  as the eigenstate of  $\mathbf{H}(\mathbf{x},\mathbf{p})$  with the zero point eigenvalue  $E_0 = \hbar \omega/2$ .  $\mathbf{H}(\mathbf{x},\mathbf{p}) |0\rangle = \hbar \omega/2 |0\rangle$   $\langle 0| \mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar \omega/2 \langle 0|$ Action by **a** on ground ket  $|0\rangle$  (or **a**<sup>†</sup> on ground bra  $\langle 0|$ ) gives *nothing* (zero vectors  $\boldsymbol{\theta}$ ).  $\mathbf{a} |0\rangle = \mathbf{0}$   $\langle 0| \mathbf{a}^{\dagger} = \mathbf{0}$ But,  $\mathbf{a}^{\dagger}$  acts on ground ket to give  $|1\rangle = \mathbf{a}^{\dagger}|0\rangle$  with  $\mathbf{H}$  eigenvalue  $E_1 = \hbar\omega + E_0$ .  $(|1\rangle = \mathbf{a}^{\dagger}|0\rangle, \langle 0|\mathbf{a} = \langle 1|.\rangle$ Proof:  $\mathbf{H}(\mathbf{x},\mathbf{p}) \mathbf{a}^{\dagger}|0\rangle = \hbar\omega \mathbf{a}^{\dagger}\mathbf{a} \mathbf{a}^{\dagger}|0\rangle + \hbar\omega/2 \mathbf{a}^{\dagger}|0\rangle$  $\mathbf{H}(\mathbf{x},\mathbf{p}) \mathbf{a}^{\dagger}|0\rangle = \hbar\omega \mathbf{a}^{\dagger}(\mathbf{a}^{\dagger}\mathbf{a} + \mathbf{1})|0\rangle + \hbar\omega/2 \mathbf{a}^{\dagger}|0\rangle$  $= \hbar\omega \mathbf{a}^{\dagger} (\mathbf{a}^{\dagger} \mathbf{a} + \mathbf{1}) |0\rangle + \hbar\omega/2 \mathbf{a}^{\dagger} |0\rangle$  $= \hbar\omega \mathbf{a}^{\dagger} |0\rangle + \mathbf{0} + \hbar\omega/2 \mathbf{a}^{\dagger} |0\rangle \qquad QED:$  $+\hbar\omega/2$ )  $|1\rangle = E_1 |1\rangle$  where:  $E_1 = \hbar\omega + E_0$  $\mathbf{H}(\mathbf{x},\mathbf{p}) |1\rangle = (\hbar\omega)$ 

Eigenstate creationism (and destructionism) Given 1D-HO Hamiltonian:  $(\mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar \omega \mathbf{a}^{\dagger} \mathbf{a} + \mathbf{1}\hbar \omega/2)$  and commutation:  $([\mathbf{a},\mathbf{a}^{\dagger}] = \mathbf{1})$  or  $(\mathbf{a}\mathbf{a}^{\dagger} = \mathbf{a}^{\dagger}\mathbf{a} + \mathbf{1})$ Define ground state  $|0\rangle$  as the eigenstate of  $\mathbf{H}(\mathbf{x},\mathbf{p})$  with the zero point eigenvalue  $E_0 = \hbar \omega/2$ .  $\mathbf{H}(\mathbf{x},\mathbf{p}) |0\rangle = \hbar \omega/2 |0\rangle$   $\langle 0| \mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar \omega/2 \langle 0|$ Action by **a** on ground ket  $|0\rangle$  (or **a**<sup>†</sup> on ground bra  $\langle 0|$ ) gives *nothing* (zero vectors  $\boldsymbol{\theta}$ ).  $\mathbf{a} |0\rangle = \mathbf{0}$   $\langle 0| \mathbf{a}^{\dagger} = \mathbf{0}$ But,  $\mathbf{a}^{\dagger}$  acts on ground ket to give  $|1\rangle = \mathbf{a}^{\dagger}|0\rangle$  with  $\mathbf{H}$  eigenvalue  $E_1 = \hbar\omega + E_0$ .  $(|1\rangle = \mathbf{a}^{\dagger}|0\rangle, \langle 0|\mathbf{a} = \langle 1|.\rangle$ Proof:  $\mathbf{H}(\mathbf{x},\mathbf{p}) \mathbf{a}^{\dagger}|0\rangle = \hbar\omega \mathbf{a}^{\dagger}\mathbf{a} \mathbf{a}^{\dagger}|0\rangle + \hbar\omega/2 \mathbf{a}^{\dagger}|0\rangle$  $\mathbf{H}(\mathbf{x},\mathbf{p}) \mathbf{a}^{\dagger}|0\rangle = \hbar\omega \mathbf{a}^{\dagger}(\mathbf{a}^{\dagger}\mathbf{a} + \mathbf{1})|0\rangle + \hbar\omega/2 \mathbf{a}^{\dagger}|0\rangle$  $= \hbar\omega \mathbf{a}' (\mathbf{a}'\mathbf{a} + \mathbf{1})|0\rangle + \hbar\omega/2 \mathbf{a}'|0\rangle$  $= \hbar\omega \mathbf{a}^{\dagger}|0\rangle + \mathbf{0} + \hbar\omega/2 \mathbf{a}^{\dagger}|0\rangle$ *QED:*  $\mathbf{H}(\mathbf{x},\mathbf{p})|1\rangle = (\hbar\omega + \hbar\omega/2)|1\rangle = E_1|1\rangle$  where:  $E_1 = \hbar\omega + E_0$ *One-quantum* or *1st excited eigenket*  $|1\rangle = \mathbf{a}^{\dagger}|0\rangle$ 

Eigenstate creationism (and destructionism) Given 1D-HO Hamiltonian:  $(\mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar \omega \mathbf{a}^{\dagger} \mathbf{a} + \mathbf{1}\hbar \omega/2)$  and commutation:  $([\mathbf{a},\mathbf{a}^{\dagger}] = \mathbf{1})$  or  $(\mathbf{a}\mathbf{a}^{\dagger} = \mathbf{a}^{\dagger}\mathbf{a} + \mathbf{1}$ Define ground state  $|0\rangle$  as the eigenstate of  $\mathbf{H}(\mathbf{x},\mathbf{p})$  with the zero point eigenvalue  $E_0 = \hbar \omega/2$ .  $\mathbf{H}(\mathbf{x},\mathbf{p}) |0\rangle = \hbar \omega/2 |0\rangle$   $\langle 0| \mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar \omega/2 \langle 0|$ Action by **a** on ground ket  $|0\rangle$  (or **a**<sup>†</sup> on ground bra  $\langle 0|$ ) gives *nothing* (zero vectors  $\boldsymbol{\theta}$ ).  $\mathbf{a} |0\rangle = \mathbf{0}$   $\langle 0| \mathbf{a}^{\dagger} = \mathbf{0}$ But,  $\mathbf{a}^{\dagger}$  acts on ground ket to give  $|1\rangle = \mathbf{a}^{\dagger}|0\rangle$  with  $\mathbf{H}$  eigenvalue  $E_1 = \hbar\omega + E_0$ .  $(|1\rangle = \mathbf{a}^{\dagger}|0\rangle, \langle 0|\mathbf{a} = \langle 1|.\rangle$ Proof:  $\mathbf{H}(\mathbf{x},\mathbf{p}) \mathbf{a}^{\dagger}|0\rangle = \hbar\omega \mathbf{a}^{\dagger}\mathbf{a} \mathbf{a}^{\dagger}|0\rangle + \hbar\omega/2 \mathbf{a}^{\dagger}|0\rangle$  $\mathbf{H}(\mathbf{x},\mathbf{p}) \mathbf{a}^{\dagger}|0\rangle = \hbar\omega \mathbf{a}^{\dagger}(\mathbf{a}^{\dagger}\mathbf{a} + \mathbf{1})|0\rangle + \hbar\omega/2 \mathbf{a}^{\dagger}|0\rangle$  $= \hbar\omega \mathbf{a}' (\mathbf{a} + \mathbf{1})|0\rangle + \hbar\omega/2 \mathbf{a}'|0\rangle$  $= \hbar\omega \mathbf{a}^{\dagger}|0\rangle + \mathbf{0} + \hbar\omega/2 \mathbf{a}^{\dagger}|0\rangle$ *QED:*  $\mathbf{H}(\mathbf{x},\mathbf{p})|1\rangle = (\hbar\omega + \hbar\omega/2)|1\rangle = E_1|1\rangle$  where:  $E_1 = \hbar\omega + E_0$ *One-quantum* or *1st excited eigenket*  $|1\rangle = \mathbf{a}^{\dagger}|0\rangle$ For kets, **a**<sup>†</sup> is *creation operator* while **a** is *destruction operator*.  $|\mathbf{a}|1\rangle = \mathbf{a}\mathbf{a}^{\dagger}|0\rangle = (\mathbf{a}^{\dagger}\mathbf{a} + \mathbf{1})|0\rangle = |0\rangle$ 

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*Time evolution of a coherent state*  $|\alpha\rangle$ *Properties of coherent states and "squeezed" states*
$$\langle x | \mathbf{a} | 0 \rangle = \frac{1}{\sqrt{2\hbar}} \left( \sqrt{M\omega} \langle x | \mathbf{x} | 0 \rangle + i \langle x | \mathbf{p} | 0 \rangle / \sqrt{M\omega} \right) = 0$$

Coordinate representation of the "nothing" equation  $\langle x | \mathbf{a} | 0 \rangle = \mathbf{0}$ 

with: 
$$\mathbf{p} = \mathbf{h}\mathbf{k} = \frac{\hbar}{i}\frac{\partial}{\partial x}$$

$$\langle x | \mathbf{a} | 0 \rangle = \frac{1}{\sqrt{2\hbar}} \left( \sqrt{M\omega} \langle x | \mathbf{x} | 0 \rangle + i \langle x | \mathbf{p} | 0 \rangle / \sqrt{M\omega} \right) = 0$$
$$\sqrt{M\omega} x \psi_0(x) + i \frac{\hbar}{i} \frac{\partial \psi_0(x)}{\partial x} / \sqrt{M\omega} = 0$$

$$\langle x | \mathbf{a} | 0 \rangle = \frac{1}{\sqrt{2\hbar}} \left( \sqrt{M\omega} \langle x | \mathbf{x} | 0 \rangle + i \langle x | \mathbf{p} | 0 \rangle / \sqrt{M\omega} \right) = 0$$

$$\sqrt{M\omega} x \psi_0(x) + i \frac{\hbar}{i} \frac{\partial \psi_0(x)}{\partial x} / \sqrt{M\omega} = 0$$

$$\psi'_0(x) = -\frac{M\omega}{\hbar} x \psi_0(x)$$

$$\begin{aligned} \langle x | \mathbf{a} | 0 \rangle &= \frac{1}{\sqrt{2\hbar}} \left( \sqrt{M\omega} \langle x | \mathbf{x} | 0 \rangle + i \langle x | \mathbf{p} | 0 \rangle / \sqrt{M\omega} \right) &= 0 \\ \sqrt{M\omega} x \psi_0 (x) + i \frac{\hbar}{i} \frac{\partial \psi_0 (x)}{\partial x} / \sqrt{M\omega} &= 0 \\ \psi'_0 (x) &= -\frac{M\omega}{\hbar} x \psi_0 (x) \end{aligned}$$
$$\begin{aligned} & \int \frac{d\psi}{\psi} &= \int \frac{-M\omega}{\hbar} x dx , \quad \ln \psi + \ln const. = \frac{-M\omega}{\hbar} \frac{x^2}{2}, \quad \psi = \frac{e^{-M\omega x^2/2\hbar}}{const.} \end{aligned}$$

$$\begin{aligned} \langle x | \mathbf{a} | 0 \rangle &= \frac{1}{\sqrt{2\hbar}} \left( \sqrt{M\omega} \langle x | \mathbf{x} | 0 \rangle + i \langle x | \mathbf{p} | 0 \rangle / \sqrt{M\omega} \right) &= 0 \\ \sqrt{M\omega} x \psi_0 (x) + i \frac{\hbar}{i} \frac{\partial \psi_0 (x)}{\partial x} / \sqrt{M\omega} &= 0 \\ \psi'_0 (x) &= -\frac{M\omega}{\hbar} x \psi_0 (x) \end{aligned}$$
$$\begin{aligned} & \int \frac{d\psi}{\psi} &= \int \frac{-M\omega}{\hbar} x dx , \quad \ln \psi + \ln const. = \frac{-M\omega}{\hbar} \frac{x^2}{2}, \quad \psi = \frac{e^{-M\omega x^2/2\hbar}}{const.} \end{aligned}$$



with:  $\mathbf{p} = \mathbf{h}\mathbf{k} = \frac{\hbar}{\cdot} \frac{\partial}{\partial}$ Coordinate representation of the "nothing" equation  $\langle x | \mathbf{a} | 0 \rangle = \mathbf{0}$  $\langle x | \mathbf{a} | 0 \rangle = \frac{1}{\sqrt{2\hbar}} \left( \sqrt{M\omega} \langle x | \mathbf{x} | 0 \rangle + i \langle x | \mathbf{p} | 0 \rangle / \sqrt{M\omega} \right) = 0$  $\sqrt{M\omega} x \psi_0(x) + i \frac{\hbar}{i} \frac{\partial \psi_0(x)}{\partial x} / \sqrt{M\omega} = 0$  $\psi_0'(x) = -\frac{M\omega}{\hbar} x \psi_0(x)$  $\int \frac{d\psi}{\psi} = \int \frac{-M\omega}{\hbar} x dx , \quad \ln \psi + \ln const. = \frac{-M\omega}{\hbar} \frac{x^2}{2}, \quad \psi = \frac{e^{-M\omega x^2/2\hbar}}{const.} = \frac{e^{-M\omega x^2/2\hbar}}{\left(\frac{\pi \hbar}{\pi}\right)^{1/4}}$ The normalization *const.* is evaluated using a standard Gaussian integral:  $\int_{-\infty}^{\infty} dx \ e^{-\alpha x^2} = \sqrt{\frac{\pi}{2}}$  $\left\langle \psi_{0} \middle| \psi_{0} \right\rangle = 1 = \int_{-\infty}^{\infty} dx \, \frac{e^{-M\omega x^{2} 2/2\hbar}}{const^{2}} = \sqrt{\frac{\pi \hbar}{M\omega}} / const.^{2} \Rightarrow const. = \left(\frac{\pi \hbar}{M\omega}\right)^{1/4}$  $\psi_0(x)$ 15.9 9.55 Zero-point 10 9.18 energy  $E_0$ ່ 🗘 ເ<sup>5</sup>  $=\hbar\omega/2$ Х Classical turning points

Symmetry group  $\mathcal{G}=U(1)$  representations, 1D HO Hamiltonian  $\mathbf{H}=\hbar\omega \mathbf{a}^{\dagger}\mathbf{a}$  operators, 1D HO wave eigenfunctions  $\Psi_n$ , and coherent  $\alpha$ -states

Factoring 1D-HO Hamiltonian H=p<sup>2</sup>+x<sup>2</sup>

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*Wavefunction creationism (1st Excited state)* 1st excited state wavefunction  $\psi_1(x) = \langle x | 1 \rangle$ ????  $\langle x \mid \mathbf{a}^{\dagger} \mid 0 \rangle = \langle x \mid 1 \rangle = \psi_1(x)$ 22.9 15.9  $\psi_0(x)$ 9.55 Zero-point 109.18 energy  $E_0$ <sup>5</sup> ا 151 1 x  $=\hbar\omega/2$ X Classical turning points

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#### Wavefunction creationism (1<sup>st</sup> Excited state)

1st excited state wavefunction  $\psi_1(x) = \langle x | 1 \rangle$  $\langle x | \mathbf{a}^{\dagger} | 0 \rangle = \langle x | 1 \rangle = \psi_1(x)$ 

Expanding the creation operator

$$\left\langle x \left| \mathbf{a}^{\dagger} \right| 0 \right\rangle = \frac{1}{\sqrt{2\hbar}} \left( \sqrt{M\omega} \left\langle x \left| \mathbf{x} \right| 0 \right\rangle - i \left\langle x \left| \mathbf{p} \right| 0 \right\rangle / \sqrt{M\omega} \right) = \left\langle x \left| 1 \right\rangle = \psi_1(x)$$



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The operator coordinate representations generate the first excited state wavefunction.

 $\langle x|1\rangle = \psi_1(x) = \frac{1}{\sqrt{2\hbar}} \left( \sqrt{M\omega} x \psi_0(x) - i\frac{\hbar}{i} \frac{\partial \psi_0(x)}{\partial x} / \sqrt{M\omega} \right)$ 



#### *Wavefunction creationism (1st Excited state)*



Expanding the creation operator

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Vacuum state  $|0\rangle$ ,  $1^{st}$  excited state  $|1\rangle$ ,  $|2\rangle$ , ...



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Creation-Destruction at a algebra of U(1) operators Eigenstate creationism (and destructionism) Vacuum state |0>, 1<sup>st</sup> excited state |1>, |2>, ...



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Binomial power expansion identities:  $aa^{\dagger n} = na^{\dagger n-1} + a^{\dagger n}a \leftarrow \cdots$  [AB, C] = - [C, AB] = -[C, A]B - A[C, B]= [A, C]B + A[B, C]

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 $= n(n-1)\mathbf{a}^{\dagger n-2} + n\mathbf{a}^{\dagger n-1}\mathbf{a} + n\mathbf{a}^{\dagger n-1}\mathbf{a} + \mathbf{a}^{\dagger n}\mathbf{a}^{2}$ 

 $= n(n-1)\mathbf{a}^{\dagger n-2} + 2n\mathbf{a}^{\dagger n-1}\mathbf{a} + \mathbf{a}^{\dagger n}\mathbf{a}^2$ 

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 $\mathbf{a}^2 \mathbf{a}^{\dagger n} = n \mathbf{a} \mathbf{a}^{\dagger n-1}$ 

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Symmetry group  $\mathcal{G}=U(1)$  representations, 1D HO Hamiltonian  $\mathbf{H}=\hbar\omega \mathbf{a}^{\dagger}\mathbf{a}$  operators, 1D HO wave eigenfunctions  $\Psi_n$ , and coherent  $\alpha$ -states

Factoring 1D-HO Hamiltonian  $\mathbf{H}=\mathbf{p}^2+\mathbf{x}^2$ Creation-Destruction  $\mathbf{a}^*\mathbf{a}$  algebra of U(1) operatorsEigenstate creationism (and destructionism)Vacuum state |0>,1st excited state |1>, |2>,...

Normal ordering for matrix calculation (creation **a**<sup>†</sup> on left, destruction **a** on right) Commutator derivative identities, Binomial expansion identities



 $\begin{array}{l} \langle \mathbf{a}^{n} \mathbf{a}^{\dagger n} \rangle \ operator \ calculations \\ Number \ operator \ and \ Hamiltonian \ operator \\ Expectation \ values \ of \ position, \ momentum, \ and \ uncertainty \ for \ eigenstate \ |n\rangle \\ Harmonic \ oscillator \ beat \ dynamics \ of \ mixed \ states \end{array}$ 

Oscillator coherent states ("Shoved" and "kicked" states) Translation operators vs. boost operators, boost-translation combinations

*Time evolution of a coherent state*  $|\alpha\rangle$ *Properties of coherent states and "squeezed" states*  *Matrix*  $\langle \mathbf{a}^{n} \mathbf{a}^{\dagger n} \rangle$  *calculation* 

Derive normalization for  $n^{th}$  state obtained by  $(\mathbf{a}^{\dagger})^n$  operator:

$$|n\rangle = \frac{\mathbf{a}^{\dagger n}|0\rangle}{const.}$$
, where:  $1 = \langle n|n\rangle = \frac{\langle 0|\mathbf{a}^n \mathbf{a}^{\dagger n}|0\rangle}{(const.)^2}$ 

 $\begin{array}{l} Matrix \left\langle \mathbf{a}^{n} \mathbf{a}^{\dagger n} \right\rangle calculation \\ \text{Derive normalization for } n^{th} \text{ state obtained by } (\mathbf{a}^{\dagger})^{n} \text{ operator:} \quad \text{Use: } \mathbf{a}^{n} \mathbf{a}^{\dagger n} = n! \left( \mathbf{1} + n \mathbf{a}^{\dagger} \mathbf{a} + \frac{n(n-1)}{2! \cdot 2!} \mathbf{a}^{\dagger 2} \mathbf{a}^{2} + \ldots \right) \\ |n\rangle = \frac{\mathbf{a}^{\dagger n} |0\rangle}{const.}, \quad \text{where:} \quad 1 = \left\langle n | n \right\rangle = \frac{\left\langle 0 | \mathbf{a}^{n} \mathbf{a}^{\dagger n} | 0 \right\rangle}{(const.)^{2}} \end{array}$ 

$$\begin{array}{l} Matrix \left\langle \mathbf{a}^{n}\mathbf{a}^{\dagger n} \right\rangle calculation \\ \text{Derive normalization for } n^{th} \text{ state obtained by } (\mathbf{a}^{\dagger})^{n} \text{ operator: } \text{Use: } \mathbf{a}^{n}\mathbf{a}^{\dagger n} = n! \left(\mathbf{1} + n\mathbf{a}^{\dagger}\mathbf{a} + \frac{n(n-1)}{2! \cdot 2!}\mathbf{a}^{\dagger 2}\mathbf{a}^{2} + \ldots\right) \\ |n\rangle = \frac{\mathbf{a}^{\dagger n}|0\rangle}{const.}, \quad \text{where: } 1 = \left\langle n|n \right\rangle = \frac{\left\langle 0|\mathbf{a}^{n}\mathbf{a}^{\dagger n}|0 \right\rangle}{(const.)^{2}} = n! \frac{\left\langle 0|\mathbf{1} + n\mathbf{a}^{\dagger}\mathbf{a} + ..|0 \right\rangle}{(const.)^{2}} = \frac{n!}{(const.)^{2}} \\ \text{So: } (const.)^{2} = n! \\ (const.) = \sqrt{n!} \end{array}$$

$$\begin{aligned} \text{Matrix } \langle \mathbf{a}^{n} \mathbf{a}^{\dagger n} \rangle \text{ calculation} \\ \text{Derive normalization for } n^{th} \text{ state obtained by } (\mathbf{a}^{\dagger})^{n} \text{ operator: } \text{Use: } \mathbf{a}^{n} \mathbf{a}^{\dagger n} = n! \left( \mathbf{1} + n \mathbf{a}^{\dagger} \mathbf{a} + \frac{n(n-1)}{2! \cdot 2!} \mathbf{a}^{\dagger 2} \mathbf{a}^{2} + \dots \right) \\ |n\rangle &= \frac{\mathbf{a}^{\dagger n} |0\rangle}{const.}, \quad \text{where: } 1 = \langle n|n\rangle = \frac{\langle 0|\mathbf{a}^{n} \mathbf{a}^{\dagger n}|0\rangle}{(const.)^{2}} = n! \frac{\langle 0|\mathbf{1} + n \mathbf{a}^{\dagger} \mathbf{a} + ..|0\rangle}{(const.)^{2}} = \frac{n!}{(const.)^{2}} \\ \text{So: } (const.)^{2} = n! \\ \left( |n\rangle = \frac{\mathbf{a}^{\dagger n} |0\rangle}{\sqrt{n!}} \quad \text{Root-factorial normalization} \right) \end{aligned}$$

$$\begin{aligned} Matrix \langle \mathbf{a}^{n} \mathbf{a}^{\dagger n} \rangle \ calculation \\ \text{Derive normalization for } n^{th} \text{ state obtained by } (\mathbf{a}^{\dagger})^{n} \text{ operator: Use: } \mathbf{a}^{n} \mathbf{a}^{\dagger n} = n! \left( \mathbf{1} + n \mathbf{a}^{\dagger} \mathbf{a} + \frac{n(n-1)}{2! \cdot 2!} \mathbf{a}^{\dagger 2} \mathbf{a}^{2} + ... \right) \\ |n\rangle &= \frac{\mathbf{a}^{\dagger n} |0\rangle}{const.}, \quad \text{where: } \mathbf{1} = \langle n|n\rangle = \frac{\langle 0|\mathbf{a}^{n} \mathbf{a}^{\dagger n}|0\rangle}{(const.)^{2}} = n! \frac{\langle 0|\mathbf{1} + n \mathbf{a}^{\dagger} \mathbf{a} + ..|0\rangle}{(const.)^{2}} = \frac{n!}{(const.)^{2}} \\ \left[ |n\rangle = \frac{\mathbf{a}^{\dagger n} |0\rangle}{\sqrt{n!}} \quad Root-factorial normalization \\ \text{Apply creation } \mathbf{a}^{\dagger}: \\ \mathbf{a}^{\dagger} |n\rangle &= \frac{\mathbf{a}^{\dagger n+1} |0\rangle}{\sqrt{n!}} \\ \mathbf{a} |n\rangle &= \frac{\mathbf{a}^{\dagger n+1} |0\rangle}{\sqrt{n!}} \end{aligned}$$

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Feynman's mnemonic rule: Larger of two quanta goes in radical factor

$$\begin{aligned} \text{Matrix } \langle \mathbf{a}^{n} \mathbf{a}^{\dagger n} \rangle \text{ calculation} \\ \text{Derive normalization for } n^{th} \text{ state obtained by } (\mathbf{a}^{\dagger})^{n} \text{ operator:} & \text{Use: } \mathbf{a}^{n} \mathbf{a}^{\dagger n} = n! \left( 1 + n \mathbf{a}^{\dagger} \mathbf{a} + \frac{n(n-1)}{2! \cdot 2!} \mathbf{a}^{\dagger 2} \mathbf{a}^{2} + \dots \right) \\ |n\rangle &= \frac{\mathbf{a}^{\dagger n} |0\rangle}{const.}, & \text{where:} & 1 = \langle n|n\rangle = \frac{\langle 0|\mathbf{a}^{n} \mathbf{a}^{\dagger n}|0\rangle}{(const.)^{2}} = n! \frac{\langle 0|\mathbf{1} + n \mathbf{a}^{\dagger} \mathbf{a} + .|0\rangle}{(const.)^{2}} = \frac{n!}{(const.)^{2}} \\ & \left( n \right) = \frac{\mathbf{a}^{\dagger n} |0\rangle}{\sqrt{n!}}, & \text{where:} & 1 = \langle n|n\rangle = \frac{\langle 0|\mathbf{a}^{n} \mathbf{a}^{\dagger n}|0\rangle}{(const.)^{2}} = n! \frac{\langle 0|\mathbf{1} + n \mathbf{a}^{\dagger} \mathbf{a} + .|0\rangle}{(const.)^{2}} = \frac{n!}{(const.)^{2}} \end{aligned} \\ & \text{Use: } \mathbf{a}^{\dagger n} = n \mathbf{a}^{\dagger n-1} + \mathbf{a}^{\dagger n} \mathbf{a} \\ & \text{Apply creation } \mathbf{a}^{\dagger}: & \text{Apply destruction } \mathbf{a}: \\ & \mathbf{a}^{\dagger} |n\rangle = \frac{\mathbf{a}^{\dagger n+1} |0\rangle}{\sqrt{n!}} = \sqrt{n+1} \frac{\mathbf{a}^{\dagger n+1} |0\rangle}{\sqrt{(n+1)!}} & \mathbf{a} |n\rangle = \frac{\mathbf{a}^{\dagger n} |0\rangle}{\sqrt{n!}} = \frac{(n \mathbf{a}^{\dagger n-1} + \mathbf{a}^{\dagger n} \mathbf{a})|0\rangle}{\sqrt{n!}} = \sqrt{n} \frac{\mathbf{a}^{\dagger n-1} |0\rangle}{\sqrt{(n-1)!}} \\ & \mathbf{a}^{\dagger} |n\rangle = \sqrt{n+1} |n+1\rangle & \mathbf{a} |n\rangle = \sqrt{n} |n-1\rangle \end{aligned}$$

Feynman's mnemonic rule: Larger of two quanta goes in radical factor

$$\left\langle \mathbf{a}^{\dagger} \right\rangle = \left( \begin{array}{cccc} \cdot & & & & \\ 1 & \cdot & & & \\ & \sqrt{2} & \cdot & & \\ & & \sqrt{3} & \cdot & \\ & & & \sqrt{4} & \cdot & \\ & & & & \ddots & \cdot \end{array} \right)$$

$$\begin{aligned} Matrix (\mathbf{a}^{\mathbf{n}}\mathbf{a}^{\dagger \mathbf{n}}) \ calculation \\ \text{Derive normalization for } n^{th} \text{ state obtained by } (\mathbf{a}^{\dagger})^{n} \text{ operator:} \quad \text{Use: } \mathbf{a}^{n}\mathbf{a}^{\dagger n} = n! \left(1 + n\mathbf{a}^{\dagger}\mathbf{a} + \frac{n(n-1)}{2!2!}\mathbf{a}^{\dagger 2}\mathbf{a}^{2} + \dots \right) \\ |n\rangle &= \frac{\mathbf{a}^{\dagger n}|0\rangle}{const.}, \quad \text{where:} \quad 1 = \langle n|n\rangle = \frac{\langle 0|\mathbf{a}^{n}\mathbf{a}^{\dagger n}|0\rangle}{(const.)^{2}} = n! \frac{\langle 0|\mathbf{1} + n\mathbf{a}^{\dagger}\mathbf{a} + .|0\rangle}{(const.)^{2}} = \frac{n!}{(const.)^{2}} \\ \hline |n\rangle &= \frac{\mathbf{a}^{\dagger n}|0\rangle}{\sqrt{n!}}, \quad \text{where:} \quad 1 = \langle n|n\rangle = \frac{\langle 0|\mathbf{a}^{n}\mathbf{a}^{\dagger n}|0\rangle}{(const.)^{2}} = n! \frac{\langle 0|\mathbf{1} + n\mathbf{a}^{\dagger}\mathbf{a} + .|0\rangle}{(const.)^{2}} = \frac{n!}{(const.)^{2}} \\ \hline |n\rangle &= \frac{\mathbf{a}^{\dagger n}|0\rangle}{\sqrt{n!}}, \quad \text{Root-factorial normalization} \\ \text{Apply creation } \mathbf{a}^{\dagger}: \\ \mathbf{a}^{\dagger}|n\rangle &= \frac{\mathbf{a}^{\dagger n+1}|0\rangle}{\sqrt{n!}} = \sqrt{n+1} \frac{\mathbf{a}^{\dagger n+1}|0\rangle}{\sqrt{(n+1)!}}, \quad \mathbf{a}|n\rangle = \frac{\mathbf{a}\mathbf{a}^{\dagger n}|0\rangle}{\sqrt{n!}} = \frac{(n\mathbf{a}^{\dagger n-1} + \mathbf{a}^{\dagger n}\mathbf{a})|0\rangle}{\sqrt{n!}} = \sqrt{n} \frac{\mathbf{a}^{\dagger n-1}|0\rangle}{\sqrt{(n-1)!}} \\ \hline \mathbf{a}^{\dagger}|n\rangle &= \sqrt{n+1} \frac{\mathbf{a}^{\dagger n+1}|0\rangle}{\sqrt{(n+1)!}}, \quad \mathbf{a}|n\rangle = \sqrt{n}|n-1\rangle \\ Feynman's mnemonic rule: Larger of two quanta goes in radical factor} \\ \langle \mathbf{a}^{\dagger} \rangle = \begin{pmatrix} 1 & \sqrt{2} & \sqrt$$

(Welcome to  $\infty$ -dimensional... quantum space!)

Symmetry group  $\mathcal{G}=U(1)$  representations, 1D HO Hamiltonian  $\mathbf{H}=\hbar\omega \mathbf{a}^{\dagger}\mathbf{a}$  operators, 1D HO wave eigenfunctions  $\Psi_n$ , and coherent  $\alpha$ -states

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$$\left\langle \mathbf{a}^{\dagger} \right\rangle = \left( \begin{array}{cccc} \cdot & & & & \\ 1 & \cdot & & & \\ & \sqrt{2} & \cdot & & \\ & & \sqrt{3} & \cdot & \\ & & & \sqrt{4} & \cdot & \\ & & & & \ddots & \cdot \end{array} \right)$$

Number operator and Hamiltonian operator Number operator  $N=a^{\dagger}a$  counts quanta.

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$$\left\langle \mathbf{a}^{\dagger} \right\rangle = \left( \begin{array}{cccc} \cdot & & & & \\ 1 & \cdot & & & \\ & \sqrt{2} & \cdot & & \\ & & \sqrt{3} & \cdot & \\ & & & \sqrt{4} & \cdot & \\ & & & & \ddots & \cdot \end{array} \right)$$

Number operator and Hamiltonian operator Number operator  $N=a^{\dagger}a$  counts quanta.

$$\mathbf{a}^{\dagger}\mathbf{a}|n\rangle = \frac{\mathbf{a}^{\dagger}\mathbf{a}\mathbf{a}^{\dagger n}|0\rangle}{\sqrt{n!}}$$

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$$\left\langle \mathbf{a}^{\dagger} \right\rangle = \left( \begin{array}{cccc} \cdot & & & & \\ 1 & \cdot & & & \\ & \sqrt{2} & \cdot & & \\ & & \sqrt{3} & \cdot & \\ & & & \sqrt{4} & \cdot & \\ & & & & \ddots & \cdot \end{array} \right)$$

$$\langle \mathbf{a} \rangle = \begin{pmatrix} \cdot & 1 \\ & \cdot & \sqrt{2} \\ & & \cdot & \sqrt{3} \\ & & & \cdot & \sqrt{4} \\ & & & & \cdot & \ddots \\ & & & & & \cdot & \ddots \\ & & & & & & \cdot & \ddots \\ & & & & & & \cdot & \ddots \\ & & & & & & & \cdot & \ddots \\ & & & & & & & \cdot & \ddots \\ & & & & & & & & \cdot & \ddots \\ & & & & & & & & \cdot & \ddots \\ & & & & & & & & \cdot & \ddots \\ \end{array}$$

$$J_{se:} aa^{\dagger n} = na^{\dagger n-1} + a^{\dagger n}a$$

Number operator and Hamiltonian operator *Number operator*  $N=a^{\dagger}a$  <u>counts quanta</u>.

$$\mathbf{a}^{\dagger}\mathbf{a}|n\rangle = \frac{\mathbf{a}^{\dagger}\mathbf{a}\mathbf{a}^{\dagger n}|0\rangle}{\sqrt{n!}}$$

rule. Laiger

$$\left\langle \mathbf{a}^{\dagger} \right\rangle = \left( \begin{array}{cccc} \cdot & & & & \\ 1 & \cdot & & & \\ & \sqrt{2} & \cdot & & \\ & & \sqrt{3} & \cdot & \\ & & & \sqrt{4} & \cdot \\ & & & & \ddots & \cdot \end{array} \right)$$

$$\langle \mathbf{a} \rangle = \begin{pmatrix} \cdot & 1 & & \\ & \cdot & \sqrt{2} & & \\ & & \cdot & \sqrt{3} & & \\ & & & \cdot & \sqrt{4} & \\ & & & & \cdot & \ddots \end{pmatrix} \qquad \mathbf{Use:} \ \mathbf{aa}^{\dagger n} = n\mathbf{a}^{\dagger n-1} + \mathbf{a}^{\dagger n}\mathbf{a}$$

Number operator and Hamiltonian operator *Number operator*  $N=a^{\dagger}a$  <u>counts quanta</u>.

$$\mathbf{a}^{\dagger}\mathbf{a}|n\rangle = \frac{\mathbf{a}^{\dagger}\mathbf{a}\mathbf{a}^{\dagger n}|0\rangle}{\sqrt{n!}} = n\frac{\mathbf{a}^{\dagger}\mathbf{a}^{\dagger n-1}|0\rangle}{\sqrt{n!}}$$

$$\begin{aligned} \text{Matrix } \langle \mathbf{a}^{n} \mathbf{a}^{\dagger n} \rangle \text{ calculation Number operator} \\ \text{Derive normalization for } n^{th} \text{ state obtained by } (\mathbf{a}^{\dagger})^{n} \text{ operator: Use: } \mathbf{a}^{n} \mathbf{a}^{\dagger n} = n! \left( 1 + n \mathbf{a}^{\dagger} \mathbf{a} + \frac{n(n-1)}{2! \cdot 2!} \mathbf{a}^{\dagger 2} \mathbf{a}^{2} + \dots \right) \\ |n\rangle &= \frac{\mathbf{a}^{\dagger n} |0\rangle}{const.}, \quad \text{where: } 1 = \langle n|n\rangle = \frac{\langle 0|\mathbf{a}^{n} \mathbf{a}^{\dagger n}|0\rangle}{(const.)^{2}} = n! \frac{\langle 0|\mathbf{1} + n \mathbf{a}^{\dagger} \mathbf{a} + ..|0\rangle}{(const.)^{2}} = \frac{n!}{(const.)^{2}} \\ &= \frac{n!}{(const.)^{2}} \\ \text{Use: } \mathbf{a}^{\dagger n} = n \mathbf{a}^{\dagger n-1} + \mathbf{a}^{\dagger n} \mathbf{a} \\ \text{Apply creation } \mathbf{a}^{\dagger}: & \text{Apply destruction } \mathbf{a}: \\ \mathbf{a}^{\dagger}|n\rangle &= \frac{\mathbf{a}^{\dagger n+1}|0\rangle}{\sqrt{n!}} = \sqrt{n+1} \frac{\mathbf{a}^{\dagger n+1}|0\rangle}{\sqrt{(n+1)!}} \\ \mathbf{a}|n\rangle &= \frac{\mathbf{a}^{\dagger n}|0\rangle}{\sqrt{n!}} = \frac{(n \mathbf{a}^{\dagger n-1} + \mathbf{a}^{\dagger n} \mathbf{a})|0\rangle}{\sqrt{n!}} = \sqrt{n} \frac{\mathbf{a}^{\dagger n-1}|0\rangle}{\sqrt{(n-1)!}} \\ \mathbf{a}^{\dagger}|n\rangle &= \sqrt{n+1}|n+1\rangle \\ \text{Feynman's mnemonic rule: Larger of two quanta goes in radical factor} \end{aligned}$$

$$\left\langle \mathbf{a}^{\dagger} \right\rangle = \left( \begin{array}{cccc} \cdot & & & & \\ 1 & \cdot & & & \\ & \sqrt{2} & \cdot & & \\ & & \sqrt{3} & \cdot & \\ & & & \sqrt{4} & \cdot & \\ & & & & \ddots & \cdot \end{array} \right)$$

$$\langle \mathbf{a} \rangle = \begin{pmatrix} \cdot & 1 & & \\ & \cdot & \sqrt{2} & & \\ & & \cdot & \sqrt{3} & & \\ & & & \cdot & \sqrt{4} & \\ & & & & \cdot & \ddots \end{pmatrix} \qquad \mathbf{Use:} \ \mathbf{aa}^{\dagger n} = n\mathbf{a}^{\dagger n-1} + \mathbf{a}^{\dagger n}\mathbf{a}$$

*Number operator and Hamiltonian operator Number operator* N=a<sup>†</sup>a <u>counts quanta</u>.

$$\mathbf{a}^{\dagger}\mathbf{a}|n\rangle = \frac{\mathbf{a}^{\dagger}\mathbf{a}\mathbf{a}^{\dagger n}|0\rangle}{\sqrt{n!}} = n\frac{\mathbf{a}^{\dagger}\mathbf{a}^{\dagger n-1}|0\rangle}{\sqrt{n!}} = n\frac{\mathbf{a}^{\dagger n}|0\rangle}{\sqrt{n!}} = n|n\rangle$$

 $\langle \mathbf{a}^{*} \rangle = \left( \begin{array}{ccc} \sqrt{3} & \cdot & \\ & \sqrt{4} & \cdot \\ & & \ddots & \cdot \end{array} \right)$ 

Number operator and Hamiltonian operator Number operator  $N=a^{\dagger}a$  counts quanta.

$$\mathbf{a}^{\dagger}\mathbf{a}|n\rangle = \frac{\mathbf{a}^{\dagger}\mathbf{a}\mathbf{a}^{\dagger n}|0\rangle}{\sqrt{n!}} = n\frac{\mathbf{a}^{\dagger}\mathbf{a}^{\dagger n-1}|0\rangle}{\sqrt{n!}} = n\frac{\mathbf{a}^{\dagger n}|0\rangle}{\sqrt{n!}} = n|n\rangle$$

Hamiltonian operator

 $\mathbf{H} |n\rangle = \hbar \omega \mathbf{a}^{\dagger} \mathbf{a} |n\rangle + \hbar \omega/2 \mathbf{1} |n\rangle = \hbar \omega (n+1/2) |n\rangle$ 

$$\begin{array}{l} \text{Matrix} \left\langle \mathbf{a}^{n} \mathbf{a}^{\dagger n} \right\rangle calculation \quad Hamiltonian \ operator\\ \text{Derive normalization for } n^{th} \text{ state obtained by } (\mathbf{a}^{\dagger})^{n} \text{ operator: } \text{Use: } \mathbf{a}^{n} \mathbf{a}^{\dagger n} = n! \left(1 + n \mathbf{a}^{\dagger} \mathbf{a} + \frac{n(n-1)}{2!2!} \mathbf{a}^{\dagger 2} \mathbf{a}^{2} + ...\right) \\ |n\rangle = \frac{\mathbf{a}^{\dagger n} |0\rangle}{const.}, \quad \text{where: } 1 = \left\langle n|n \right\rangle = \frac{\left\langle 0|\mathbf{a}^{n} \mathbf{a}^{\dagger n}|0 \right\rangle}{(const.)^{2}} = n! \frac{\left\langle 0|\mathbf{1} + n \mathbf{a}^{\dagger} \mathbf{a} + ..|0 \right\rangle}{(const.)^{2}} = \frac{n!}{(const.)^{2}} \\ \hline \left( |n\rangle = \frac{\mathbf{a}^{\dagger n} |0\rangle}{\sqrt{n!}}, \quad \text{where: } 1 = \left\langle n|n \right\rangle = \frac{\left\langle 0|\mathbf{a}^{n} \mathbf{a}^{\dagger n}|0 \right\rangle}{(const.)^{2}} = n! \frac{\left\langle 0|\mathbf{1} + n \mathbf{a}^{\dagger} \mathbf{a} + ..|0 \right\rangle}{(const.)^{2}} = \frac{n!}{(const.)^{2}} \\ \hline \left( |n\rangle = \frac{\mathbf{a}^{\dagger n} |0\rangle}{\sqrt{n!}}, \quad \text{where: } 1 = \left\langle n|n\rangle = \frac{\left\langle 0|\mathbf{a}^{n} \mathbf{a}^{\dagger n}|0 \right\rangle}{\sqrt{n!}} = n! \frac{n!}{(const.)^{2}} \\ \hline \left( |n\rangle = \frac{\mathbf{a}^{\dagger n+1} |0\rangle}{\sqrt{n!}}, \quad \text{where: } 1 = \left\langle n|n\rangle = \frac{\mathbf{a}^{\dagger n} |0\rangle}{\sqrt{n!}}, \quad \text{Root-factorial normalization} \\ \text{Apply creation } \mathbf{a}^{\dagger}: \\ \mathbf{a}^{\dagger}|n\rangle = \frac{\mathbf{a}^{\dagger n+1} |0\rangle}{\sqrt{n!}} = \sqrt{n+1} \frac{\mathbf{a}^{\dagger n+1} |0\rangle}{\sqrt{(n+1)!}}, \quad \mathbf{a}|n\rangle = \frac{\mathbf{a}^{\dagger n} |0\rangle}{\sqrt{n!}} = \left\langle n \mathbf{a}^{\dagger n-1} + \mathbf{a}^{\dagger n} \mathbf{a} \right\rangle \\ \text{Apply destruction } \mathbf{a}: \\ \mathbf{a}^{\dagger}|n\rangle = \sqrt{n+1} |\mathbf{a}^{\dagger}|n\rangle = \frac{\mathbf{a}^{\dagger n+1} |\mathbf{a}^{\dagger}|}{\sqrt{(n+1)!}}, \quad \mathbf{a}|n\rangle = \frac{\mathbf{a}^{\dagger n} |\mathbf{a}^{\dagger}|}{\sqrt{n!}} = \sqrt{n} \frac{\mathbf{a}^{\dagger n-1} + \mathbf{a}^{\dagger n} \mathbf{a}}{\sqrt{(n-1)!}} \\ \text{Feynman's mnemonic rule: Larger of two quanta goes in radical factor} \\ \left\langle \mathbf{a}^{\circ} - \left( \begin{array}{c} 1 & \sqrt{2} & \sqrt{2} \\ \sqrt{3} & \sqrt{3} & \sqrt{3} \\ \sqrt{n!} & \sqrt{n!} & = n \frac{\mathbf{a}^{\dagger n} |0\rangle}{\sqrt{n!}} = n \frac{\mathbf{a}^{\dagger n} |0\rangle}{\sqrt{n!}} \\ \text{Hamiltonian operator} \\ \text{Number operator N=} \mathbf{a}^{\dagger \mathbf{a}} (n) + \hbar m(n(n+1/2)|n\rangle \quad (h) = hm(\mathbf{a}^{\dagger \mathbf{a}}; n) \in m \left( \begin{array}{c} 1 & 2 \\ 3 & \sqrt{n} \\ \sqrt{n!} & \frac{1}{2} \\ \frac$$

.

Symmetry group  $\mathcal{G}=U(1)$  representations, 1D HO Hamiltonian  $\mathbf{H}=\hbar\omega \mathbf{a}^{\dagger}\mathbf{a}$  operators, 1D HO wave eigenfunctions  $\Psi_n$ , and coherent  $\alpha$ -states

Factoring 1D-HO Hamiltonian  $\mathbf{H}=\mathbf{p}^{2}+\mathbf{x}^{2}$ Creation-Destruction  $\mathbf{a}^{\dagger}\mathbf{a}$  algebra of U(1) operatorsEigenstate creationism (and destructionism)Vacuum state |0>,1st excited state |1>, |2>,

Normal ordering for matrix calculation (creation **a**<sup>†</sup> on left, destruction **a** on right) Commutator derivative identities, Binomial expansion identities

 $\langle a^n a^{\dagger n} \rangle$  operator calculations Number operator and Hamiltonian operator Expectation values of position, momentum, and uncertainty for eigenstate  $|n \rangle$ Harmonic oscillator beat dynamics of mixed states

*Oscillator coherent states ("Shoved" and "kicked" states) Translation operators vs. boost operators, boost-translation combinations* 

*Time evolution of a coherent state*  $|\alpha\rangle$ *Properties of coherent states and "squeezed" states* 

Operator for position **X**:  $\sqrt{\frac{M\omega}{2\hbar}}$ **x** =  $\frac{\mathbf{a} + \mathbf{a}^{\dagger}}{2}$ 

Operator for momentum **p**:  $\sqrt{\frac{1}{2\hbar M\omega}}$  **p** =  $\frac{\mathbf{a} - \mathbf{a}^{\dagger}}{2i}$ 

Operator for position 
$$\mathbf{X}: \sqrt{\frac{M\omega}{2\hbar}} \mathbf{x} = \frac{\mathbf{a} + \mathbf{a}^{\dagger}}{2}$$
  
expectation for position  $\langle \mathbf{X} \rangle$ :  
 $\overline{\mathbf{x}}|_{n} = \langle n|\mathbf{x}|n \rangle = \sqrt{\frac{\hbar}{2M\omega}} \langle n|(\mathbf{a} + \mathbf{a}^{\dagger})|n \rangle = 0$ 

Operator for momentum  $\mathbf{p}: \sqrt{\frac{1}{2\hbar M\omega}} \mathbf{p} = \frac{\mathbf{a} - \mathbf{a}^{\dagger}}{2i}$ expectation for momentum  $\langle \mathbf{p} \rangle$ :  $\overline{\mathbf{p}}|_{n} = \langle n|\mathbf{p}|n \rangle = i \sqrt{\frac{\hbar M\omega}{2}} \langle n|(\mathbf{a}^{\dagger} - \mathbf{a})|n \rangle = 0$ 

Operator for position 
$$\mathbf{X}: \sqrt{\frac{M\omega}{2\hbar}} \mathbf{x} = \frac{\mathbf{a} + \mathbf{a}^{\dagger}}{2}$$
  
expectation for position  $\langle \mathbf{X} \rangle$ :  
 $\overline{\mathbf{x}}|_{n} = \langle n|\mathbf{x}|n \rangle = \sqrt{\frac{\hbar}{2M\omega}} \langle n|(\mathbf{a} + \mathbf{a}^{\dagger})|n \rangle = 0$   
expectation for (position)<sup>2</sup>  $\langle \mathbf{X}^{2} \rangle$ :  
 $\overline{\mathbf{x}^{2}}|_{n} = \langle n|\mathbf{x}^{2}|n \rangle = \frac{\hbar}{2M\omega} \langle n|(\mathbf{a} + \mathbf{a}^{\dagger})^{2}|n \rangle$ 

Operator for momentum  $\mathbf{p}: \sqrt{\frac{1}{2\hbar M\omega}} \mathbf{p} = \frac{\mathbf{a} - \mathbf{a}^{\dagger}}{2i}$ expectation for momentum  $\langle \mathbf{p} \rangle$ :  $\overline{\mathbf{p}}|_{n} = \langle n|\mathbf{p}|n \rangle = i\sqrt{\frac{\hbar M\omega}{2}} \langle n|(\mathbf{a}^{\dagger} - \mathbf{a})|n \rangle = 0$ expectation for (momentum)<sup>2</sup>  $\langle \mathbf{p}^{2} \rangle$ :  $\overline{\mathbf{p}^{2}}|_{n} = \langle n|\mathbf{p}^{2}|n \rangle = i^{2}\frac{\hbar M\omega}{2} \langle n|(\mathbf{a}^{\dagger} - \mathbf{a})^{2}|n \rangle$ 

Operator for position 
$$\mathbf{x}: \sqrt{\frac{M\omega}{2\hbar}} \mathbf{x} = \frac{\mathbf{a} + \mathbf{a}^{\dagger}}{2}$$
  
expectation for position  $\langle \mathbf{x} \rangle$ :  
 $\overline{\mathbf{x}}|_{n} = \langle n|\mathbf{x}|n \rangle = \sqrt{\frac{\hbar}{2M\omega}} \langle n|(\mathbf{a} + \mathbf{a}^{\dagger})|n \rangle = 0$   
expectation for (position)<sup>2</sup>  $\langle \mathbf{x}^{2} \rangle$ :  
 $\overline{\mathbf{x}^{2}}|_{n} = \langle n|\mathbf{x}^{2}|n \rangle = \frac{\hbar}{2M\omega} \langle n|(\mathbf{a} + \mathbf{a}^{\dagger})^{2}|n \rangle$   
 $= \frac{\hbar}{2M\omega} \langle n|(\mathbf{a}^{2} + \mathbf{a}^{\dagger}\mathbf{a} + \mathbf{a}\mathbf{a}^{\dagger} + \mathbf{a}^{\dagger 2})|n \rangle$   
 $use:$   
 $\mathbf{a}\mathbf{a}^{\dagger} = \mathbf{1} + \mathbf{a}^{\dagger}\mathbf{a}$ 

Operator for momentum  $\mathbf{p}: \sqrt{\frac{1}{2\hbar M\omega}} \mathbf{p} = \frac{\mathbf{a} - \mathbf{a}^{\dagger}}{2i}$ expectation for momentum  $\langle \mathbf{p} \rangle$ :  $\overline{\mathbf{p}}|_{n} = \langle n|\mathbf{p}|n \rangle = i\sqrt{\frac{\hbar M\omega}{2}} \langle n|(\mathbf{a}^{\dagger} - \mathbf{a})|n \rangle = 0$ expectation for (momentum)<sup>2</sup>  $\langle \mathbf{p}^{2} \rangle$ :  $\overline{\mathbf{p}^{2}}|_{n} = \langle n|\mathbf{p}^{2}|n \rangle = i^{2}\frac{\hbar M\omega}{2} \langle n|(\mathbf{a}^{\dagger} - \mathbf{a})^{2}|n \rangle$  $= -\frac{\hbar M\omega}{2} \langle n|(\mathbf{a}^{\dagger 2} - \mathbf{a}^{\dagger}\mathbf{a} - \mathbf{a}\mathbf{a}^{\dagger} + \mathbf{a}^{2})|n \rangle$ 

Operator for position 
$$\mathbf{x}: \sqrt{\frac{M\omega}{2\hbar}} \mathbf{x} = \frac{\mathbf{a} + \mathbf{a}^{\dagger}}{2}$$
  
expectation for position  $\langle \mathbf{X} \rangle$ :  
 $\overline{\mathbf{x}}|_{n} = \langle n|\mathbf{x}|n \rangle = \sqrt{\frac{\hbar}{2M\omega}} \langle n|(\mathbf{a} + \mathbf{a}^{\dagger})|n \rangle = 0$   
expectation for (position)<sup>2</sup>  $\langle \mathbf{X}^{2} \rangle$ :  
 $\overline{\mathbf{x}^{2}}|_{n} = \langle n|\mathbf{x}^{2}|n \rangle = \frac{\hbar}{2M\omega} \langle n|(\mathbf{a} + \mathbf{a}^{\dagger})^{2}|n \rangle$   
 $= \frac{\hbar}{2M\omega} \langle n|(\mathbf{a}^{2} + \mathbf{a}^{\dagger}\mathbf{a} + \mathbf{a}\mathbf{a}^{\dagger} + \mathbf{a}^{\dagger 2})|n \rangle$   $\mathbf{a}\mathbf{a}^{\dagger} = \mathbf{1} + \mathbf{a}^{\dagger}\mathbf{a}$   
 $= \frac{\hbar}{2M\omega}$  (2n+1)

Operator for momentum  $\mathbf{p}: \sqrt{\frac{1}{2\hbar M\omega}} \mathbf{p} = \frac{\mathbf{a} - \mathbf{a}^{\dagger}}{2i}$ expectation for momentum  $\langle \mathbf{p} \rangle$ :  $\overline{\mathbf{p}}|_{n} = \langle n|\mathbf{p}|n \rangle = i\sqrt{\frac{\hbar M\omega}{2}} \langle n|(\mathbf{a}^{\dagger} - \mathbf{a})|n \rangle = 0$ expectation for (momentum)<sup>2</sup>  $\langle \mathbf{p}^{2} \rangle$ :  $\overline{\mathbf{p}^{2}}|_{n} = \langle n|\mathbf{p}^{2}|n \rangle = i^{2}\frac{\hbar M\omega}{2} \langle n|(\mathbf{a}^{\dagger} - \mathbf{a})^{2}|n \rangle$   $= -\frac{\hbar M\omega}{2} \langle n|(\mathbf{a}^{\dagger 2} - \mathbf{a}^{\dagger}\mathbf{a} - \mathbf{a}\mathbf{a}^{\dagger} + \mathbf{a}^{2})|n \rangle$  $= \frac{\hbar M\omega}{2} (2n+1)$ 

Operator for position **x**: 
$$\sqrt{\frac{M\omega}{2\hbar}} \mathbf{x} = \frac{\mathbf{a} + \mathbf{a}^{\dagger}}{2}$$
  
expectation for position  $\langle \mathbf{x} \rangle$ :  
 $\overline{\mathbf{x}}|_{n} = \langle n|\mathbf{x}|n \rangle = \sqrt{\frac{\hbar}{2M\omega}} \langle n|(\mathbf{a} + \mathbf{a}^{\dagger})|n \rangle = 0$   
expectation for (position)<sup>2</sup>  $\langle \mathbf{x}^{2} \rangle$ :  
 $\overline{\mathbf{x}^{2}}|_{n} = \langle n|\mathbf{x}^{2}|n \rangle = \frac{\hbar}{2M\omega} \langle n|(\mathbf{a} + \mathbf{a}^{\dagger})^{2}|n \rangle$   
 $= \frac{\hbar}{2M\omega} \langle n|(\mathbf{a}^{2} + \mathbf{a}^{\dagger}\mathbf{a} + \mathbf{a}\mathbf{a}^{\dagger} + \mathbf{a}^{\dagger 2})|n \rangle$   $\mathbf{a}\mathbf{a}^{\dagger} = \mathbf{1} + \mathbf{a}^{\dagger}\mathbf{a}$   
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Operator for momentum  $\mathbf{p}: \sqrt{\frac{1}{2\hbar M\omega}} \mathbf{p} = \frac{\mathbf{a} - \mathbf{a}^{\dagger}}{2i}$ expectation for momentum  $\langle \mathbf{p} \rangle$ :  $\overline{\mathbf{p}}|_{n} = \langle n|\mathbf{p}|n \rangle = i\sqrt{\frac{\hbar M\omega}{2}} \langle n|(\mathbf{a}^{\dagger} - \mathbf{a})|n \rangle = 0$ expectation for (momentum)<sup>2</sup>  $\langle \mathbf{p}^{2} \rangle$ :  $\overline{\mathbf{p}^{2}}|_{n} = \langle n|\mathbf{p}^{2}|n \rangle = i^{2}\frac{\hbar M\omega}{2} \langle n|(\mathbf{a}^{\dagger} - \mathbf{a})^{2}|n \rangle$   $= -\frac{\hbar M\omega}{2} \langle n|(\mathbf{a}^{\dagger 2} - \mathbf{a}^{\dagger}\mathbf{a} - \mathbf{a}\mathbf{a}^{\dagger} + \mathbf{a}^{2})|n \rangle$  $= \frac{\hbar M\omega}{2} (2n+1)$ 

Uncertainty or standard deviation  $\Delta q$  of a statistical quantity q is its root mean-square difference.

$$(\Delta q)^2 = \overline{(q - \overline{q})^2}$$
 or:  $\Delta q = \sqrt{(q - \overline{q})^2}$ 

Operator for position **x**: 
$$\sqrt{\frac{M\omega}{2\hbar}} \mathbf{x} = \frac{\mathbf{a} + \mathbf{a}^{\dagger}}{2}$$
  
expectation for position  $\langle \mathbf{X} \rangle$ :  
 $\overline{\mathbf{x}}|_{n} = \langle n|\mathbf{x}|n \rangle = \sqrt{\frac{\hbar}{2M\omega}} \langle n|(\mathbf{a} + \mathbf{a}^{\dagger})|n \rangle = 0$   
expectation for (position)<sup>2</sup>  $\langle \mathbf{X}^{2} \rangle$ :  
 $\overline{\mathbf{x}^{2}}|_{n} = \langle n|\mathbf{x}^{2}|n \rangle = \frac{\hbar}{2M\omega} \langle n|(\mathbf{a} + \mathbf{a}^{\dagger})^{2}|n \rangle$   
 $= \frac{\hbar}{2M\omega} \langle n|(\mathbf{a}^{2} + \mathbf{a}^{\dagger}\mathbf{a} + \mathbf{a}\mathbf{a}^{\dagger} + \mathbf{a}^{\dagger 2})|n \rangle$   $\mathbf{a}\mathbf{a}^{\dagger} = \mathbf{1} + \mathbf{a}^{\dagger}\mathbf{a}$   
 $= \frac{\hbar}{2M\omega}$  (2n+1)

Operator for momentum  $\mathbf{p}: \sqrt{\frac{1}{2\hbar M\omega}} \mathbf{p} = \frac{\mathbf{a} - \mathbf{a}^{\dagger}}{2i}$ expectation for momentum  $\langle \mathbf{p} \rangle$ :  $\overline{\mathbf{p}}|_n = \langle n | \mathbf{p} | n \rangle = i \sqrt{\frac{\hbar M\omega}{2}} \langle n | (\mathbf{a}^{\dagger} - \mathbf{a}) | n \rangle = 0$ expectation for (momentum)<sup>2</sup>  $\langle \mathbf{p}^2 \rangle$ :  $\overline{\mathbf{p}^2}|_n = \langle n | \mathbf{p}^2 | n \rangle = i^2 \frac{\hbar M\omega}{2} \langle n | (\mathbf{a}^{\dagger} - \mathbf{a})^2 | n \rangle$   $= -\frac{\hbar M\omega}{2} \langle n | (\mathbf{a}^{\dagger 2} - \mathbf{a}^{\dagger} \mathbf{a} - \mathbf{a} \mathbf{a}^{\dagger} + \mathbf{a}^2) | n \rangle$  $= \frac{\hbar M\omega}{2}$  (2n+1)

Uncertainty or standard deviation  $\Delta q$  of a statistical quantity q is its root mean-square difference.

$$\Delta x|_{n} = \sqrt{|\mathbf{x}^{2}|_{n}} = \sqrt{\frac{\hbar(2n+1)}{2M\omega}} \qquad (\Delta q)^{2} = \overline{(q-\overline{q})^{2}} \quad \text{or:} \quad \Delta q = \sqrt{(q-\overline{q})^{2}} \\ \Delta p|_{n} = \sqrt{\frac{\hbar(2n+1)}{2M\omega}} \qquad \Delta p|_{n} = \sqrt{\frac{\hbar(2n+1)}{2M\omega}}$$

Operator for position 
$$\mathbf{x}: \sqrt{\frac{M\omega}{2\hbar}} \mathbf{x} = \frac{\mathbf{a} + \mathbf{a}^{\dagger}}{2}$$
  
expectation for position  $\langle \mathbf{x} \rangle$ :  
 $\overline{\mathbf{x}} |_{n} = \langle n | \mathbf{x} | n \rangle = \sqrt{\frac{\hbar}{2M\omega}} \langle n | (\mathbf{a} + \mathbf{a}^{\dagger}) | n \rangle = 0$   
expectation for (position)<sup>2</sup>  $\langle \mathbf{x}^{2} \rangle$ :  
 $\overline{\mathbf{x}^{2}} |_{n} = \langle n | \mathbf{x}^{2} | n \rangle = \frac{\hbar}{2M\omega} \langle n | (\mathbf{a} + \mathbf{a}^{\dagger})^{2} | n \rangle$   
 $= \frac{\hbar}{2M\omega} \langle n | (\mathbf{a}^{2} + \mathbf{a}^{\dagger}\mathbf{a} + \mathbf{a}\mathbf{a}^{\dagger} + \mathbf{a}^{\dagger 2}) | n \rangle$   $\mathbf{a}\mathbf{a}^{\dagger} = \mathbf{1} + \mathbf{a}^{\dagger}\mathbf{a}$   
 $= \frac{\hbar}{2M\omega} (2n+1)$ 

Operator for momentum  $\mathbf{p}: \sqrt{\frac{1}{2\hbar M\omega}} \mathbf{p} = \frac{\mathbf{a} - \mathbf{a}^{\dagger}}{2i}$ expectation for momentum  $\langle \mathbf{p} \rangle$ :  $\bar{\mathbf{p}} |_n = \langle n | \mathbf{p} | n \rangle = i \sqrt{\frac{\hbar M\omega}{2}} \langle n | (\mathbf{a}^{\dagger} - \mathbf{a}) | n \rangle = 0$ expectation for (momentum)<sup>2</sup>  $\langle \mathbf{p}^2 \rangle$ :  $\bar{\mathbf{p}}^2 |_n = \langle n | \mathbf{p}^2 | n \rangle = i^2 \frac{\hbar M\omega}{2} \langle n | (\mathbf{a}^{\dagger} - \mathbf{a})^2 | n \rangle$   $= -\frac{\hbar M\omega}{2} \langle n | (\mathbf{a}^{\dagger 2} - \mathbf{a}^{\dagger} \mathbf{a} - \mathbf{a} \mathbf{a}^{\dagger} + \mathbf{a}^2) | n \rangle$  $= \frac{\hbar M\omega}{2} (2n+1)$ 

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$$\Delta x|_{n} = \sqrt{\mathbf{x}^{2}} = \sqrt{\frac{\hbar(2n+1)}{2M\omega}} \qquad (\Delta q)^{2} = \overline{(q-\overline{q})^{2}} \quad \text{or:} \quad \Delta q = \sqrt{(q-\overline{q})^{2}} \\ \Delta p|_{n} = \sqrt{\mathbf{p}^{2}} = \sqrt{\frac{\hbar M\omega(2n+1)}{2}}$$

*Heisenberg uncertainty product* for the *n*-quantum eigenstate  $|n\rangle$ 

$$(\Delta x \cdot \Delta p)\Big|_{n} = \sqrt{\mathbf{x}^{2}} \sqrt{\mathbf{p}^{2}} = \sqrt{\frac{\hbar(2n+1)}{2M\omega}} \sqrt{\frac{\hbar M\omega(2n+1)}{2}}$$

Operator for position 
$$\mathbf{x}: \sqrt{\frac{M\omega}{2\hbar}} \mathbf{x} = \frac{\mathbf{a} + \mathbf{a}^{\dagger}}{2}$$
  
expectation for position  $\langle \mathbf{x} \rangle$ :  
 $\overline{\mathbf{x}} |_{n} = \langle n | \mathbf{x} | n \rangle = \sqrt{\frac{\hbar}{2M\omega}} \langle n | (\mathbf{a} + \mathbf{a}^{\dagger}) | n \rangle = 0$   
expectation for (position)<sup>2</sup>  $\langle \mathbf{x}^{2} \rangle$ :  
 $\overline{\mathbf{x}^{2}} |_{n} = \langle n | \mathbf{x}^{2} | n \rangle = \frac{\hbar}{2M\omega} \langle n | (\mathbf{a} + \mathbf{a}^{\dagger})^{2} | n \rangle$   
 $= \frac{\hbar}{2M\omega} \langle n | (\mathbf{a}^{2} + \mathbf{a}^{\dagger}\mathbf{a} + \mathbf{a}\mathbf{a}^{\dagger} + \mathbf{a}^{\dagger 2}) | n \rangle$   $\mathbf{a}\mathbf{a}^{\dagger} = \mathbf{1} + \mathbf{a}^{\dagger}\mathbf{a}$   
 $= \frac{\hbar}{2M\omega} (2n+1)$ 

Operator for momentum  $\mathbf{p}: \sqrt{\frac{1}{2\hbar M\omega}} \mathbf{p} = \frac{\mathbf{a} - \mathbf{a}^{\dagger}}{2i}$ expectation for momentum  $\langle \mathbf{p} \rangle$ :  $\mathbf{\bar{p}}|_n = \langle n|\mathbf{p}|n \rangle = i\sqrt{\frac{\hbar M\omega}{2}} \langle n|(\mathbf{a}^{\dagger} - \mathbf{a})|n \rangle = 0$ expectation for (momentum)<sup>2</sup>  $\langle \mathbf{p}^2 \rangle$ :  $\mathbf{\bar{p}}^2|_n = \langle n|\mathbf{p}^2|n \rangle = i^2 \frac{\hbar M\omega}{2} \langle n|(\mathbf{a}^{\dagger} - \mathbf{a})^2|n \rangle$   $= -\frac{\hbar M\omega}{2} \langle n|(\mathbf{a}^{\dagger 2} - \mathbf{a}^{\dagger}\mathbf{a} - \mathbf{a}\mathbf{a}^{\dagger} + \mathbf{a}^2)|n \rangle$  $= \frac{\hbar M\omega}{2} (2n+1)$ 

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$$\Delta x|_{n} = \sqrt{\mathbf{x}^{2}} = \sqrt{\frac{\hbar(2n+1)}{2M\omega}} \qquad (\Delta q)^{2} = \overline{(q-\overline{q})^{2}} \quad \text{or:} \quad \Delta q = \sqrt{(q-\overline{q})^{2}} \\ \Delta p|_{n} = \sqrt{\mathbf{p}^{2}} = \sqrt{\frac{\hbar M\omega(2n+1)}{2}}$$

*Heisenberg uncertainty product* for the *n*-quantum eigenstate  $|n\rangle$ 

$$(\Delta x \cdot \Delta p) \Big|_{n} = \sqrt{\mathbf{x}^{2}} \sqrt{\mathbf{p}^{2}} = \sqrt{\frac{\hbar(2n+1)}{2M\omega}} \sqrt{\frac{\hbar M\omega(2n+1)}{2}}$$
$$(\Delta x \cdot \Delta p) \Big|_{n} = \hbar \left( n + \frac{1}{2} \right)$$

Operator for position 
$$\mathbf{x}: \sqrt{\frac{M\omega}{2\hbar}} \mathbf{x} = \frac{\mathbf{a} + \mathbf{a}^{\dagger}}{2}$$
  
expectation for position  $\langle \mathbf{x} \rangle$ :  
 $\overline{\mathbf{x}} |_{n} = \langle n | \mathbf{x} | n \rangle = \sqrt{\frac{\hbar}{2M\omega}} \langle n | (\mathbf{a} + \mathbf{a}^{\dagger}) | n \rangle = 0$   
expectation for (position)<sup>2</sup>  $\langle \mathbf{x}^{2} \rangle$ :  
 $\overline{\mathbf{x}^{2}} |_{n} = \langle n | \mathbf{x}^{2} | n \rangle = \frac{\hbar}{2M\omega} \langle n | (\mathbf{a} + \mathbf{a}^{\dagger})^{2} | n \rangle$   
 $= \frac{\hbar}{2M\omega} \langle n | (\mathbf{a}^{2} + \mathbf{a}^{\dagger}\mathbf{a} + \mathbf{a}\mathbf{a}^{\dagger} + \mathbf{a}^{\dagger 2}) | n \rangle$   $\mathbf{a}\mathbf{a}^{\dagger} = \mathbf{1} + \mathbf{a}^{\dagger}\mathbf{a}$   
 $= \frac{\hbar}{2M\omega} (2n+1)$ 

Operator for momentum  $\mathbf{p}: \sqrt{\frac{1}{2\hbar M\omega}} \mathbf{p} = \frac{\mathbf{a} - \mathbf{a}^{\dagger}}{2i}$ expectation for momentum  $\langle \mathbf{p} \rangle$ :  $\bar{\mathbf{p}} |_n = \langle n | \mathbf{p} | n \rangle = i \sqrt{\frac{\hbar M\omega}{2}} \langle n | (\mathbf{a}^{\dagger} - \mathbf{a}) | n \rangle = 0$ expectation for (momentum)<sup>2</sup>  $\langle \mathbf{p}^2 \rangle$ :  $\bar{\mathbf{p}}^2 |_n = \langle n | \mathbf{p}^2 | n \rangle = i^2 \frac{\hbar M\omega}{2} \langle n | (\mathbf{a}^{\dagger} - \mathbf{a})^2 | n \rangle$   $= -\frac{\hbar M\omega}{2} \langle n | (\mathbf{a}^{\dagger 2} - \mathbf{a}^{\dagger} \mathbf{a} - \mathbf{a} \mathbf{a}^{\dagger} + \mathbf{a}^2) | n \rangle$  $= \frac{\hbar M\omega}{2} (2n+1)$ 

Uncertainty or standard deviation  $\Delta q$  of a statistical quantity q is its root mean-square difference.

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Heisenberg minimum uncertainty product occurs for the 0-quantum (ground) eigenstate.

$$(\Delta x \cdot \Delta p) \Big|_0 = \frac{\hbar}{2}$$

We pause for sobering considerations of the quantum world *vs*. the classical one. Consider a "high"-quantum (n=20) eigenstate wavefunction:



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Symmetry group  $\mathcal{G}=U(1)$  representations, 1D HO Hamiltonian  $\mathbf{H}=\hbar\omega \mathbf{a}^{\dagger}\mathbf{a}$  operators, 1D HO wave eigenfunctions  $\Psi_n$ , and coherent  $\alpha$ -states

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*Time evolution of a coherent state*  $|\alpha\rangle$ *Properties of coherent states and "squeezed" states* 

 $\Psi(x) = \langle x | \Psi \rangle = \langle x | 0 \rangle \langle 0 | \Psi \rangle + \langle x | 1 \rangle \langle 1 | \Psi \rangle = \psi_0(x) \Psi 0 + \psi_1(x) \Psi 1$ 

The time dependence  $\Psi(x,t)$  of the mixed wave is then

 $\Psi(x,t) = \psi_0(x) \ e^{-i\omega_0 t} \ \Psi_0 + \psi_1(x) \ e^{-i\omega_1 t} \ \Psi_1 = (\psi_0(x) \ e^{-i\omega_0 t} + \psi_1(x) \ e^{-i\omega_1 t})/\sqrt{2}$ 

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$$\Psi(x,t) = \sqrt{\Psi^*\Psi} = \sqrt{\left(e^{-i\omega_0 t}\psi_0(x) + e^{-i\omega_1 t}\psi_1(x)\right)^* \left(e^{-i\omega_0 t}\psi_0(x) + e^{-i\omega_1 t}\psi_1(x)\right)/2}$$
$$= \sqrt{\left(\left|\psi_0(x)\right|^2 + \left|\psi_1(x)\right|^2 + \psi_0(x)\psi_1(x)\left(e^{i(\omega_1 - \omega_0)t} + e^{-i(\omega_1 - \omega_0)t}\right)\right)/2}$$

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$$\begin{split} \Psi(x,t) &| = \sqrt{\Psi^* \Psi} = \sqrt{\left( e^{-i\omega_0 t} \psi_0(x) + e^{-i\omega_1 t} \psi_1(x) \right)^* \left( e^{-i\omega_0 t} \psi_0(x) + e^{-i\omega_1 t} \psi_1(x) \right) / 2} \\ &= \sqrt{\left( \left| \psi_0(x) \right|^2 + \left| \psi_1(x) \right|^2 + \psi_0(x) \psi_1(x) \left( e^{i(\omega_1 - \omega_0)t} + e^{-i(\omega_1 - \omega_0)t} \right) \right) / 2} \\ &= \sqrt{\left( \left| \psi_0(x) \right|^2 + \left| \psi_1(x) \right|^2 + 2\psi_0(x) \psi_1(x) \cos(\omega_1 - \omega_0)t \right) / 2} \end{split}$$

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$$= \sqrt{\left(|\psi_0(x)|^2 + |\psi_1(x)|^2 + \psi_0(x)\psi_1(x)\left(e^{i(\omega_1 - \omega_0)t} + e^{-i(\omega_1 - \omega_0)t}\right)\right) / 2}$$

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$$t = 0$$

$$t = \tau/4$$

$$E_1^{str}$$

$$E_0^{str}$$

$$E_0^{str}$$

$$t = \tau/2$$

$$t = 3\tau/4$$

$$E_0^{str}$$

$$E_0^{str}$$








![](_page_109_Figure_0.jpeg)

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Boost operators and generators: (A "kick") Boost operator **B**(b) boosts p-wavefunctions **B**(b). $\psi(p) = \psi(p-b) = \langle x | \mathbf{B}(b) | \psi \rangle = \langle p-b | \psi \rangle$ 

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Shoves  $\psi$  *a*-units to right or *x*-space *a*-units left

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Boost operators and generators: (A "kick") *Boost operator* **B**(*b*) boosts *p*-wavefunctions  $\mathbf{B}(b) \cdot \psi(p) = \psi(p-b) = \langle x | \mathbf{B}(b) | \psi \rangle = \langle p-b | \psi \rangle$ Increases momentum of ket-state by *b* units  $\langle p | \mathbf{B}(b) = \langle p - b |$ , or:  $\mathbf{B}^{\dagger}(b) | p \rangle = | p - b \rangle$ Tiny boost  $b \rightarrow db$  is identity **1** plus  $\mathbf{K} \cdot db$  $\mathbf{B}(db) = \mathbf{1} + \mathbf{K} \cdot db \quad \text{where: } \mathbf{K} = \frac{\partial \mathbf{B}}{\partial b}\Big|_{b=0}$ is generator **K** of boosts  $\mathbf{B}(b) = \left(\mathbf{B}(\frac{b}{N})\right)^{N} = \lim_{N \to \infty} \left(1 + \frac{b}{N}\mathbf{K}\right)^{N} = e^{b\mathbf{K}}$  $\mathbf{B}(b) \cdot \psi(p) = e^{b\mathbf{K}} \cdot \psi(p) = e^{-b\frac{\partial}{\partial p}} \cdot \psi(p)$  $=\psi(p)-b\frac{\partial\psi(p)}{\partial p}+\frac{b^2}{2!}\frac{\partial^2\psi(p)}{\partial p^2}-\frac{b^3}{2!}\frac{\partial^3\psi(p)}{\partial p^3}+\dots$ **K** relates to position  $\mathbf{X} \rightarrow \hbar i \frac{\partial}{\partial p} = i \frac{\partial}{\partial k}$  **K** =  $\frac{i}{\hbar} \mathbf{X} \rightarrow -\frac{\partial}{\partial p} = \frac{-1}{\hbar} \frac{\partial}{\partial k}$  $\mathbf{B}(b) = e^{b\frac{i}{\hbar}\mathbf{X}} = e^{ib(\mathbf{a}^{\dagger} + \mathbf{a})/\sqrt{2\hbar M\omega}}$ 

$$\mathbf{T}(a)\cdot\psi(x) = \psi(x-a) = \langle x|\mathbf{T}(a)|\psi\rangle = \langle x-a|\psi\rangle$$

Shoves  $\psi$  *a*-units to right or *x*-space *a*-units left

$$\langle x | \mathbf{T}(a) = \langle x - a | \text{ or: } \mathbf{T}^{\dagger}(a) | x \rangle = | x - a \rangle$$

Tiny translation  $a \rightarrow da$  is identity 1 plus  $\mathbf{G} \cdot da$   $\mathbf{T}(da) = \mathbf{1} + \mathbf{G} \cdot da$  where:  $\mathbf{G} = \frac{\partial \mathbf{T}}{\partial a}\Big|_{a=0}$ is generator **G** of translations

$$\mathbf{T}(a) = \left(\mathbf{T}\left(\frac{a}{N}\right)\right)^{N} = \lim_{N \to \infty} \left(1 + \frac{a}{N}\mathbf{G}\right)^{N} = e^{a\mathbf{G}}$$
$$\mathbf{T}(a) \cdot \psi(x) = e^{a\mathbf{G}} \cdot \psi(x) = e^{-a\frac{\partial}{\partial x}} \cdot \psi(x)$$
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**Bottom Line** 

(Move up)

$$\mathbf{T}(a) = e^{-a\frac{\iota}{\hbar}\mathbf{p}} = e^{a\left(\mathbf{a}^{\dagger} - \mathbf{a}\right)\sqrt{M\omega/2\hbar}}$$

Check T(a) on plane-wave with  $p=\hbar k$  (More T(a) $e^{ikx} = e^{-ia\mathbf{p}/\hbar}e^{ikx} = e^{-iak}e^{ikx} = e^{ik(x-a)}$ 

Boost operators and generators: (A "kick") *Boost operator* **B**(*b*) boosts *p*-wavefunctions  $\mathbf{B}(b) \cdot \psi(p) = \psi(p-b) = \langle x | \mathbf{B}(b) | \psi \rangle = \langle p-b | \psi \rangle$ Increases momentum of ket-state by *b* units  $\langle p | \mathbf{B}(b) = \langle p - b |$ , or:  $\mathbf{B}^{\dagger}(b) | p \rangle = | p - b \rangle$ Tiny boost  $b \rightarrow db$  is identity 1 plus  $\mathbf{K} \cdot db$  $\mathbf{B}(db) = \mathbf{1} + \mathbf{K} \cdot db \quad \text{where: } \mathbf{K} = \frac{\partial \mathbf{B}}{\partial b} \Big|_{b=0}$ is generator **K** of boosts  $\mathbf{B}(b) = \left(\mathbf{B}(\frac{b}{N})\right)^{N} = \lim_{N \to \infty} \left(1 + \frac{b}{N}\mathbf{K}\right)^{N} = e^{b\mathbf{K}}$  $\mathbf{B}(b) \cdot \psi(p) = e^{b\mathbf{K}} \cdot \psi(p) = e^{-b\frac{\partial}{\partial p}} \cdot \psi(p)$  $=\psi(p)-b\frac{\partial\psi(p)}{\partial p}+\frac{b^2}{2!}\frac{\partial^2\psi(p)}{\partial p^2}-\frac{b^3}{2!}\frac{\partial^3\psi(p)}{\partial p^3}+\dots$ **K** relates to position  $\mathbf{X} \rightarrow \hbar i \frac{\partial}{\partial p} = i \frac{\partial}{\partial k}$  **K** =  $\frac{i}{\hbar} \mathbf{X} \rightarrow -\frac{\partial}{\partial p} = \frac{-1}{\hbar} \frac{\partial}{\partial k}$  $\mathbf{B}(b) = e^{b\frac{i}{\hbar}\mathbf{x}} = e^{ib(\mathbf{a}^{\dagger} + \mathbf{a})/\sqrt{2\hbar M\omega}}$ Check **B**(*b*) on plane-wave with  $p=\hbar k$  $\mathbf{B}(b)e^{ikx} = e^{ib\mathbf{X}/\hbar}e^{ikx} = e^{ibx/\hbar}e^{ikx} = e^{i(k+b/\hbar)x}$ 

$$\mathbf{T}(a) \cdot \psi(x) = e^{a\mathbf{G}} \cdot \psi(x) = e^{-a\frac{\partial}{\partial x}} \cdot \psi(x)$$

$$= \psi(x) - a\frac{\partial\psi(x)}{\partial x} + \frac{a^2}{2!}\frac{\partial^2\psi(x)}{\partial x^2} - \frac{a^3}{2!}\frac{\partial^3\psi(x)}{\partial x^3} + \dots$$

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$$\frac{\mathbf{Check } \mathbf{T}(a) \text{ on plane-wave with } p = \hbar k}{\mathbf{T}(a)e^{ikx}} = e^{-ia\mathbf{p}/\hbar}e^{ikx} = e^{-iak}e^{ikx} = e^{ik(x-a)}$$

$$\begin{split} \mathbf{B}(b) \cdot \psi(p) &= e^{b\mathbf{K}} \cdot \psi(p) = e^{-b\frac{\partial}{\partial p}} \cdot \psi(p) \\ &= \psi(p) - b\frac{\partial \psi(p)}{\partial p} + \frac{b^2}{2!} \frac{\partial^2 \psi(p)}{\partial p^2} - \frac{b^3}{2!} \frac{\partial^3 \psi(p)}{\partial p^3} + \dots \\ \mathbf{K} \text{ relates to position } \mathbf{X} \rightarrow \hbar i \frac{\partial}{\partial p} = i \frac{\partial}{\partial k} \\ \mathbf{K} &= \frac{i}{\hbar} \mathbf{X} \rightarrow -\frac{\partial}{\partial p} = \frac{-1}{\hbar} \frac{\partial}{\partial k} \\ \mathbf{B}(b) &= e^{b\frac{i}{\hbar} \mathbf{X}} = e^{ib\left(\mathbf{a}^{\dagger} + \mathbf{a}\right)/\sqrt{2\hbar M\omega}} \\ \frac{Check \mathbf{B}(b) \text{ on plane-wave with } p = \hbar k}{\mathbf{B}(b)e^{ikx}} = e^{ib\mathbf{x}/\hbar}e^{ikx} = e^{ibx/\hbar}e^{ikx} = e^{i\left(k+b/\hbar\right)x} \end{split}$$

Symmetry group  $\mathcal{G}=U(1)$  representations, 1D HO Hamiltonian  $\mathbf{H}=\hbar\omega \mathbf{a}^{\dagger}\mathbf{a}$  operators, 1D HO wave eigenfunctions  $\Psi_n$ , and coherent  $\alpha$ -states

Factoring 1D-HO Hamiltonian  $H=p^{2}+x^{2}$ *Creation-Destruction* **a**<sup>*†*</sup>**a** *algebra of U*(1) *operators Eigenstate creationism (and destructionism) Vacuum state* |0>, *1st excited state* |1>, |2>, ...

Normal ordering for matrix calculation (creation **a**<sup>†</sup> on left, destruction **a** on right) Commutator derivative identities, Binomial expansion identities

 $\langle a^n a^{\dagger n} \rangle$  operator calculations Number operator and Hamiltonian operator Expectation values of position, momentum, and uncertainty for eigenstate  $|n\rangle$ Harmonic oscillator beat dynamics of mixed states

Oscillator coherent states ("Shoved" and "kicked" states) *Translation operators vs. boost operators,* **boost-translation combinations** 

*Time evolution of a coherent state*  $|\alpha>$ Properties of coherent states and "squeezed" states Applying boost-translation combinationsT(a) and B(b) operations do not commute.Q. Which should come first?

 $\mathbf{T}(a)$  and  $\mathbf{B}(b)$  operations do not commute. Q. Which should come first?  $\mathbf{T}(a) = e^{-ia\mathbf{p}/\hbar}$  or  $\mathbf{B}(b) = e^{ib\mathbf{x}/\hbar}$ ??

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May evaluate with *Baker-Campbell-Hausdorf identity* since  $[\mathbf{x},\mathbf{p}]=i\hbar\mathbf{1}$  and  $[[\mathbf{x},\mathbf{p}],\mathbf{x}]=[[\mathbf{x},\mathbf{p}],\mathbf{p}]=\mathbf{0}$ .  $e^{\mathbf{A}+\mathbf{B}} = e^{\mathbf{A}}e^{\mathbf{B}}e^{-[\mathbf{A},\mathbf{B}]/2} = e^{\mathbf{B}}e^{\mathbf{A}}e^{[\mathbf{A},\mathbf{B}]/2}$ , where:  $[\mathbf{A},[\mathbf{A},\mathbf{B}]] = \mathbf{0} = [\mathbf{B},[\mathbf{A},\mathbf{B}]]$  (left as an exercise)  $\mathbf{C}(a,b) = e^{i(b\mathbf{x}-a\mathbf{p})/\hbar} = e^{ib\mathbf{x}/\hbar}e^{-ia\mathbf{p}/\hbar}e^{-ab[\mathbf{x},\mathbf{p}]/2\hbar^2} = e^{ib\mathbf{x}/\hbar}e^{-ia\mathbf{p}/\hbar}e^{-iab/2\hbar}$  $= \mathbf{B}(b)\mathbf{T}(a)e^{-iab/2\hbar} = \mathbf{T}(a)\mathbf{B}(b)e^{iab/2\hbar}$ 

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Reordering only affects the overall phase.

$$\mathbf{C}(a,b) = e^{i(b\mathbf{x}-a\mathbf{p})/\hbar} = e^{ib(\mathbf{a}^{\dagger}+\mathbf{a})/\sqrt{2\hbar M\omega} + a(\mathbf{a}^{\dagger}-\mathbf{a})\sqrt{M\omega/2\hbar}}$$
$$= e^{\alpha \mathbf{a}^{\dagger}-\alpha^{*}\mathbf{a}} = e^{-|\alpha|^{2}/2}e^{\alpha \mathbf{a}^{\dagger}}e^{-\alpha^{*}\mathbf{a}} = e^{|\alpha|^{2}/2}e^{-\alpha^{*}\mathbf{a}}e^{\alpha \mathbf{a}^{\dagger}}$$

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Complex *phasor coordinate*  $\alpha(a,b)$  is defined by: Reordering only affects the overall phase.  $\mathbf{C}(a,b) = e^{i(b\mathbf{x}-a\mathbf{p})/\hbar} = e^{ib(\mathbf{a}^{\dagger}+\mathbf{a})/\sqrt{2\hbar M\omega} + a(\mathbf{a}^{\dagger}-\mathbf{a})\sqrt{M\omega/2\hbar}}$  $\alpha(a,b)$  $=a\sqrt{M\omega/2\hbar}+ib/\sqrt{2\hbar M\omega}$  $=e^{\alpha \mathbf{a}^{\dagger}-\alpha * \mathbf{a}} = e^{-|\alpha|^2/2}e^{\alpha \mathbf{a}^{\dagger}}e^{-\alpha * \mathbf{a}} = e^{|\alpha|^2/2}e^{-\alpha * \mathbf{a}}e^{\alpha \mathbf{a}^{\dagger}}e^{-\alpha * \mathbf{a}} = e^{|\alpha|^2/2}e^{-\alpha * \mathbf{a}}e^{\alpha \mathbf{a}^{\dagger}}e^{-\alpha * \mathbf{a}}e^{\alpha \mathbf{a}}e$  $= \left[a + i\frac{b}{M\omega}\right]\sqrt{M\omega/2\hbar}$ Coherent wavepacket state  $|\alpha(x_0, p_0)\rangle$ :  $|\alpha_0(x_0, p_0)\rangle = \mathbf{C}(x_0, p_0)|0\rangle = e^{i(x_0 \mathbf{x} - p_0 \mathbf{p})/\hbar}|0\rangle$  $=e^{-|\boldsymbol{\alpha}_0|^2/2}e^{\boldsymbol{\alpha}_0\mathbf{a}^{\dagger}}e^{-\boldsymbol{\alpha}_0^*\mathbf{a}}|0\rangle$  $=e^{-|\boldsymbol{\alpha}_0|^2/2}e^{\boldsymbol{\alpha}_0\mathbf{a}^{\dagger}}|0\rangle$  $=e^{-|\boldsymbol{\alpha}_0|^2/2}\sum_{n=0}^{\infty}\left(\boldsymbol{\alpha}_0\,\mathbf{a}^{\dagger}\right)^n |0\rangle/n!$  $=e^{-|\boldsymbol{\alpha}_0|^2/2}\sum_{n=0}^{\infty}\frac{(\boldsymbol{\alpha}_0)^n}{\sqrt{n!}}|n\rangle , \quad \text{where: } |n\rangle = \frac{\mathbf{a}^{\dagger n}|0\rangle}{\sqrt{n!}}$ 

Symmetry group  $\mathcal{G}=U(1)$  representations, 1D HO Hamiltonian  $\mathbf{H}=\hbar\omega \mathbf{a}^{\dagger}\mathbf{a}$  operators, 1D HO wave eigenfunctions  $\Psi_n$ , and coherent  $\alpha$ -states

Factoring 1D-HO Hamiltonian  $\mathbf{H}=\mathbf{p}^{2}+\mathbf{x}^{2}$ Creation-Destruction  $\mathbf{a}^{\dagger}\mathbf{a}$  algebra of U(1) operatorsEigenstate creationism (and destructionism)Vacuum state |0>,1st excited state |1>, |2>,

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Time evolution operator for constant **H** has general form :  $\mathbf{U}(t,0) = e^{-|\boldsymbol{\alpha}_0|^2/2} \sum_{n=0}^{\infty} \frac{(\boldsymbol{\alpha}_0)^n}{\sqrt{n!}} |n\rangle$ 

# *Time evolution of coherent state:*

$$\left| \frac{\alpha_0}{\alpha_0} (x_0, p_0) \right\rangle = e^{-\left| \frac{\alpha_0}{\alpha_0} \right|^2 / 2} \sum_{n=0}^{\infty} \frac{\left( \frac{\alpha_0}{\sqrt{n!}} \right)^n}{\sqrt{n!}} \left| n \right\rangle$$

Time evolution operator for constant **H** has general form :**U**(t,0)=e<sup>-i**H** $t/\hbar$ </sup>

Oscillator eigenstate time evolution is simply determined by harmonic phases.

$$\mathbf{U}(t,0)|n\rangle = e^{-i\mathbf{H}t/\hbar}|n\rangle = e^{-i(n+1/2)\omega t}|n\rangle$$

## *Time evolution of coherent state:*

 $\left| \frac{\boldsymbol{\alpha}_{0}(x_{0}, p_{0})}{\boldsymbol{\alpha}_{0}(x_{0}, p_{0})} \right\rangle = e^{-\left| \boldsymbol{\alpha}_{0} \right|^{2}/2} \sum_{n=0}^{\infty} \frac{\left( \boldsymbol{\alpha}_{0} \right)^{n}}{\sqrt{n!}} \left| n \right\rangle$   $\mathbf{H}_{t}/\hbar$ 

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Coherent state evolution results.

$$\begin{aligned} \mathbf{U}(t,0) \Big| \boldsymbol{\alpha}_{0}(x_{0},p_{0}) \Big\rangle &= e^{-|\boldsymbol{\alpha}_{0}|^{2}/2} \sum_{n=0}^{\infty} \frac{(\boldsymbol{\alpha}_{0})^{n}}{\sqrt{n!}} \mathbf{U}(t,0) \Big| n \Big\rangle = e^{-|\boldsymbol{\alpha}_{0}|^{2}/2} \sum_{n=0}^{\infty} \frac{(\boldsymbol{\alpha}_{0})^{n}}{\sqrt{n!}} e^{-i(n+1/2)\omega t} \Big| n \Big\rangle \\ &= e^{-i\omega t/2} e^{-|\boldsymbol{\alpha}_{0}|^{2}/2} \sum_{n=0}^{\infty} \frac{(\boldsymbol{\alpha}_{0}e^{-i\omega t})^{n}}{\sqrt{n!}} \Big| n \Big\rangle \end{aligned}$$

## *Time evolution of coherent state:*

 $\left| \frac{\alpha_0}{\alpha_0} (x_0, p_0) \right\rangle = e^{-\left| \alpha_0 \right|^2 / 2} \sum_{n=0}^{\infty} \frac{\left( \alpha_0 \right)^n}{\sqrt{n!}} \left| n \right\rangle$ Time evolution operator for constant **H** has general form :**U**(t,0)=e<sup>-i**H** $t/\hbar$ </sup>

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$$\begin{aligned} \mathsf{J}(t,0) \Big| \alpha_0(x_0,p_0) \Big\rangle &= e^{-|\alpha_0|^2/2} \sum_{n=0}^{\infty} \frac{(\alpha_0)^n}{\sqrt{n!}} \mathsf{U}(t,0) |n\rangle \\ &= e^{-i\omega t/2} e^{-|\alpha_0|^2/2} \sum_{n=0}^{\infty} \frac{(\alpha_0 e^{-i\omega t})^n}{\sqrt{n!}} |n\rangle \end{aligned}$$

Evolution simplifies to a variable- $\alpha_0$  coherent state with a *time dependent phasor coordinate*  $\alpha_t$ :

$$\mathbf{U}(t,0)|\boldsymbol{\alpha}_{0}(x_{0},p_{0})\rangle = e^{-i\omega t/2}|\boldsymbol{\alpha}_{t}(x_{t},p_{t})\rangle \text{ where:}$$

 $\boldsymbol{\alpha}_{t}(x_{t}, p_{t}) = e^{-i\omega t} \boldsymbol{\alpha}_{0}(x_{0}, p_{0})$  $\begin{bmatrix} x_{t} + i\frac{p_{t}}{M\omega} \end{bmatrix} = e^{-i\omega t} \begin{bmatrix} x_{0} + i\frac{p_{0}}{M\omega} \end{bmatrix}$ 

 $\left| \alpha_{0}(x_{0},p_{0}) \right\rangle = e^{-\left| \alpha_{0} \right|^{2}/2} \sum_{n=0}^{\infty} \frac{\left( \alpha_{0} \right)^{n}}{\sqrt{n!}} \left| n \right\rangle$ Time evolution operator for constant **H** has general form :**U**(*t*,0)=e^{-i\mathbf{H}t/\hbar} Oscillator eigenstate time evolution is interval.

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Coherent state evolution results.

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$$(x_t, p_t)$$
 mimics classical oscillato

$$x_t = x_0 \cos \omega t + \frac{p_0}{M\omega} \sin \omega t$$

$$\frac{p_t}{M\omega} = -x_0 \sin \omega t + \frac{p_0}{M\omega} \cos \omega t$$

Real and imaginary parts ( $x_t$  and  $p_t/M\omega$ ) of  $\alpha_t$  go clockwise on phasor circle

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Coherent ket  $|\alpha(x_0, p_0)\rangle$  is eigenvector of destruct-op. **a**.  $\mathbf{a} |\alpha_0(x_0, p_0)\rangle = e^{-|\alpha_0|^2/2} \sum_{n=0}^{\infty} \frac{(\alpha_0)^n}{\sqrt{n!}} \mathbf{a} |n\rangle$ 



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$$\mathbf{a} |\alpha_0(x_0, p_0)\rangle = e^{-|\alpha_0|^2/2} \sum_{n=0}^{\infty} \frac{(\alpha_0)^n}{\sqrt{n!}} \mathbf{a} |n\rangle$$
$$= e^{-|\alpha_0|^2/2} \sum_{n=0}^{\infty} \frac{(\alpha_0)^n}{\sqrt{n!}} \sqrt{n} |n-1\rangle$$

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$$\mathbf{a} \left| \boldsymbol{\alpha}_{0} \left( x_{0}, p_{0} \right) \right\rangle = e^{-|\boldsymbol{\alpha}_{0}|^{2}/2} \sum_{n=0}^{\infty} \frac{\left( \boldsymbol{\alpha}_{0} \right)^{n}}{\sqrt{n!}} \mathbf{a} | n \rangle$$
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$$\mathbf{a} | \boldsymbol{\alpha}_{0} (x_{0}, p_{0}) \rangle = e^{-|\boldsymbol{\alpha}_{0}|^{2}/2} \sum_{n=0}^{\infty} \frac{(\boldsymbol{\alpha}_{0})^{n}}{\sqrt{n!}} \mathbf{a} | n \rangle$$
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 $||\alpha_t\rangle$  Coherent bra  $\langle \alpha(x_0, p_0) |$  is eigenvector of create-op. **a**<sup>†</sup>.

$$\langle \boldsymbol{\alpha}_0(x_0, p_0) | \mathbf{a}^{\dagger} = \langle \boldsymbol{\alpha}_0(x_0, p_0) | \boldsymbol{\alpha}_0^*$$

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Yay! Cosine trajectory! (That is "fat")



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what happens if you apply operators with non-linear "tensor" exponents  $exp(s\mathbf{x}^2)$ ,  $exp(f\mathbf{p}^2)$ , etc.



Coherent wave packet uncertainty relation:  $\Delta n \cdot \Delta \phi > \pi/n$ 



???? Some uncertainty remains about this uncertainty ????