

Lecture 30.

Relativity of interfering and galloping waves: SWR and SWQ II.

(Ch. 4-6 of Unit 2 4.12.12)

Unmatched amplitudes giving galloping waves

Standing Wave Ratio (SWR) and Standing Wave Quotient (SWQ)

Analogy with group and phase

Analogy between wave galloping, Keplarlian IHO orbits, and optical polarization

Waves that go back in time - The Feynman-Wheeler Switchback

1st Quantization: Quantizing phase variables ω and k

Understanding how quantum transitions require “mixed-up” states

Closed cavity vs Ring cavity

Lecture 30 ended here

Galloping waves due to unmatched amplitudes

2-CW dynamics has two 1-CW amplitudes A_{\rightarrow} and A_{\leftarrow} that we now allow to be *unmatched*. ($A_{\rightarrow} \neq A_{\leftarrow}$)

$$A_{\rightarrow} e^{i(k_{\rightarrow} x - \omega_{\rightarrow} t)} + A_{\leftarrow} e^{i(k_{\leftarrow} x - \omega_{\leftarrow} t)} = e^{i(k_{\Sigma} x - \omega_{\Sigma} t)} [A_{\rightarrow} e^{i(k_{\Delta} x - \omega_{\Delta} t)} + A_{\leftarrow} e^{-i(k_{\Delta} x - \omega_{\Delta} t)}]$$

Waves have half-sum mean-phase rates $(k_{\Sigma}, \omega_{\Sigma})$ and half-difference group rates $(k_{\Delta}, \omega_{\Delta})$.

$$k_{\Sigma} = (k_{\rightarrow} + k_{\leftarrow}) / 2$$

$$\omega_{\Sigma} = (\omega_{\rightarrow} + \omega_{\leftarrow}) / 2$$

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Also important is amplitude mean $A_{\Sigma} = (A_{\rightarrow} + A_{\leftarrow}) / 2$ and half-difference $A_{\Delta} = (A_{\rightarrow} - A_{\leftarrow}) / 2$.

Detailed wave motion depends on **standing-wave-ratio *SWR*** or the inverse **standing-wave-quotient *SWQ***.

$$SWR = \frac{(A_{\rightarrow} - A_{\leftarrow})}{(A_{\rightarrow} + A_{\leftarrow})}$$

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These are analogous to frequency ratios for *group velocity* $V_{group} < c$ and its inverse that is *phase velocity* $V_{phase} > c$.

$$V_{group} = \frac{\omega_{\Delta}}{k_{\Delta}} = \frac{(\omega_{\rightarrow} - \omega_{\leftarrow})}{(k_{\rightarrow} - k_{\leftarrow})} = c \frac{(\omega_{\rightarrow} - \omega_{\leftarrow})}{(\omega_{\rightarrow} + \omega_{\leftarrow})}$$

$$V_{phase} = \frac{\omega_{\Sigma}}{k_{\Sigma}} = \frac{(\omega_{\rightarrow} + \omega_{\leftarrow})}{(k_{\rightarrow} + k_{\leftarrow})} = c \frac{(\omega_{\rightarrow} + \omega_{\leftarrow})}{(\omega_{\rightarrow} - \omega_{\leftarrow})}$$

$$\frac{V_{group}}{c} = \frac{c}{V_{phase}}$$

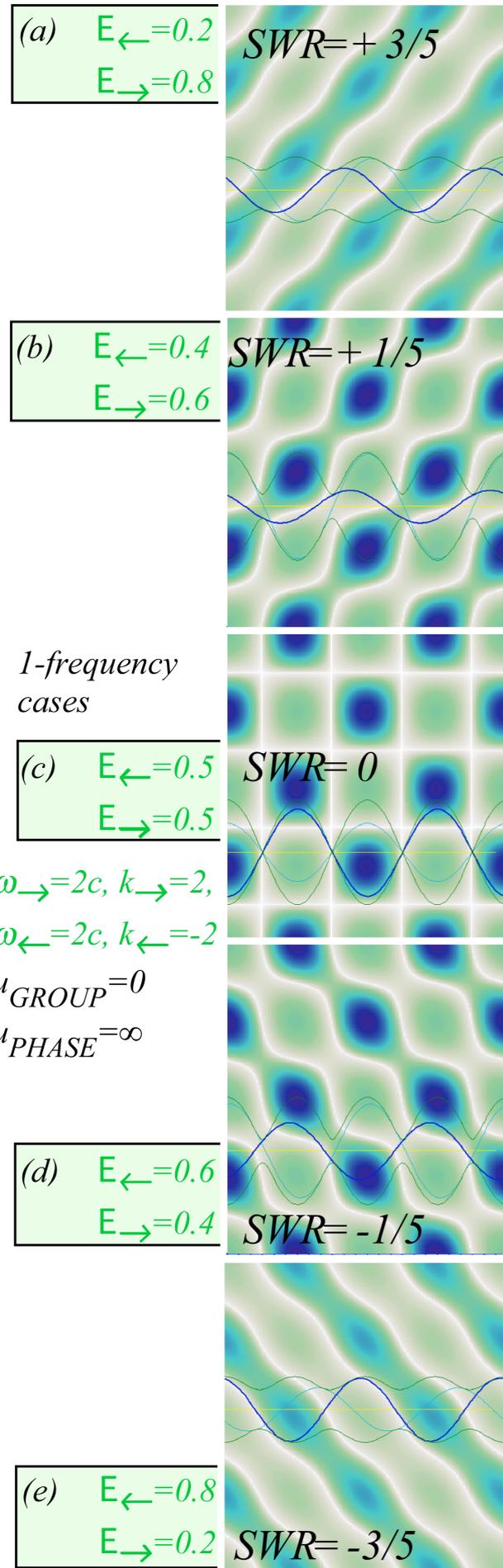


Fig. 6.1 Monochromatic (1-frequency) 2-CW wave space-time patterns.



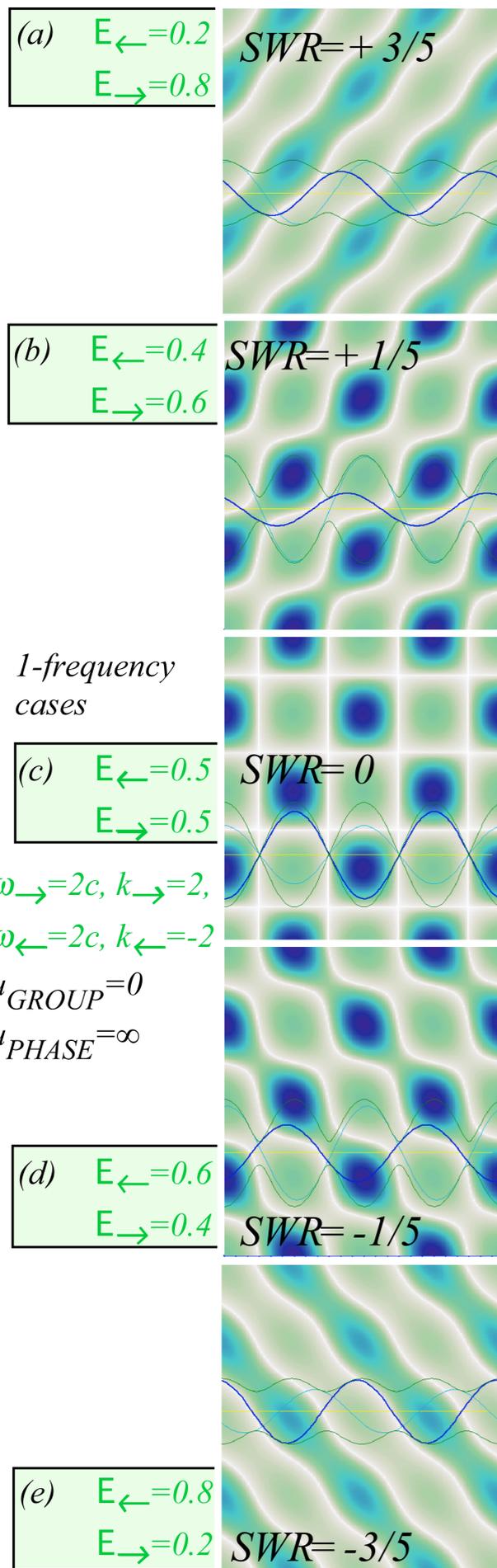


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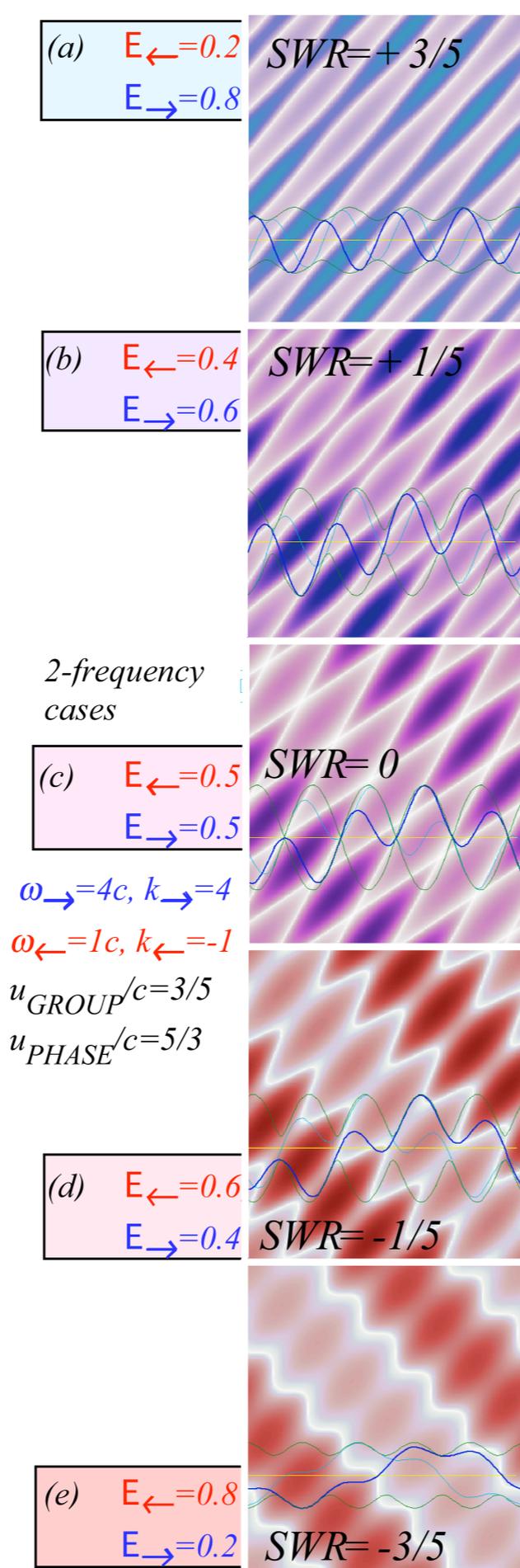


Fig. 6.2 Dichromatic (2-frequency) 2-CW wave space-time patterns.

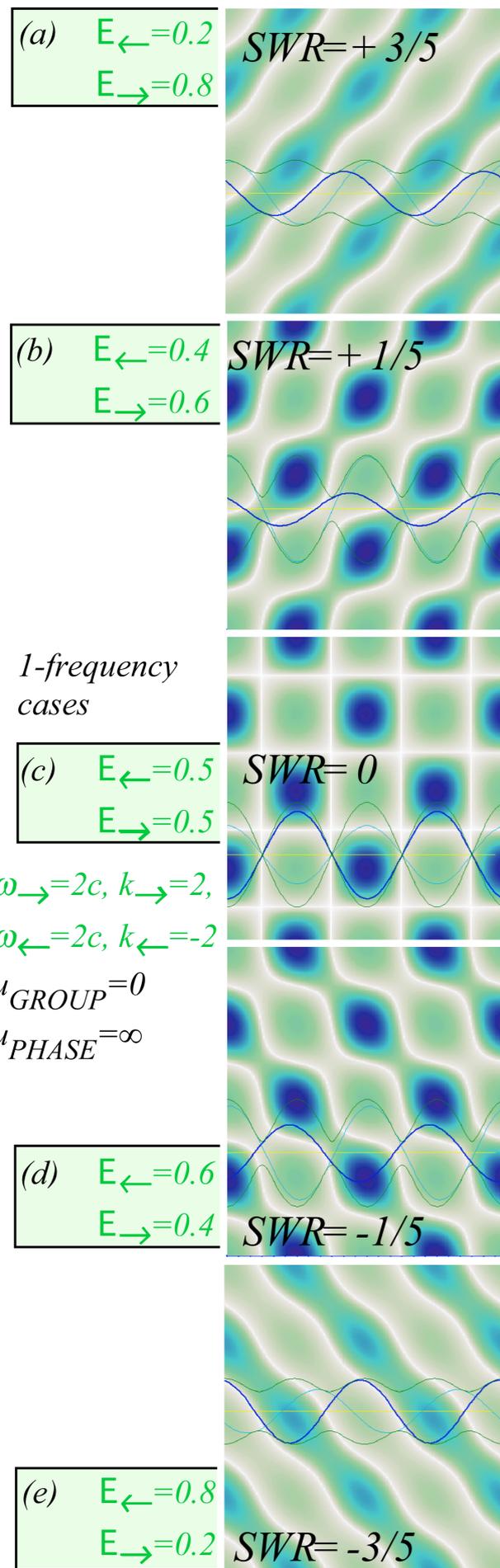


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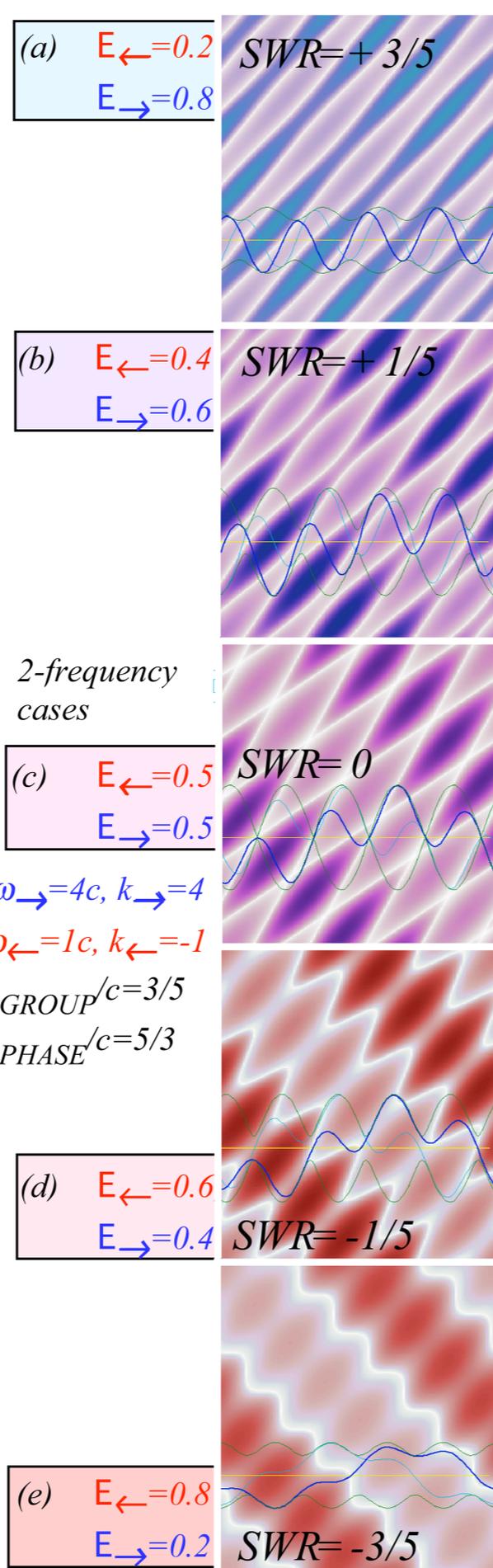


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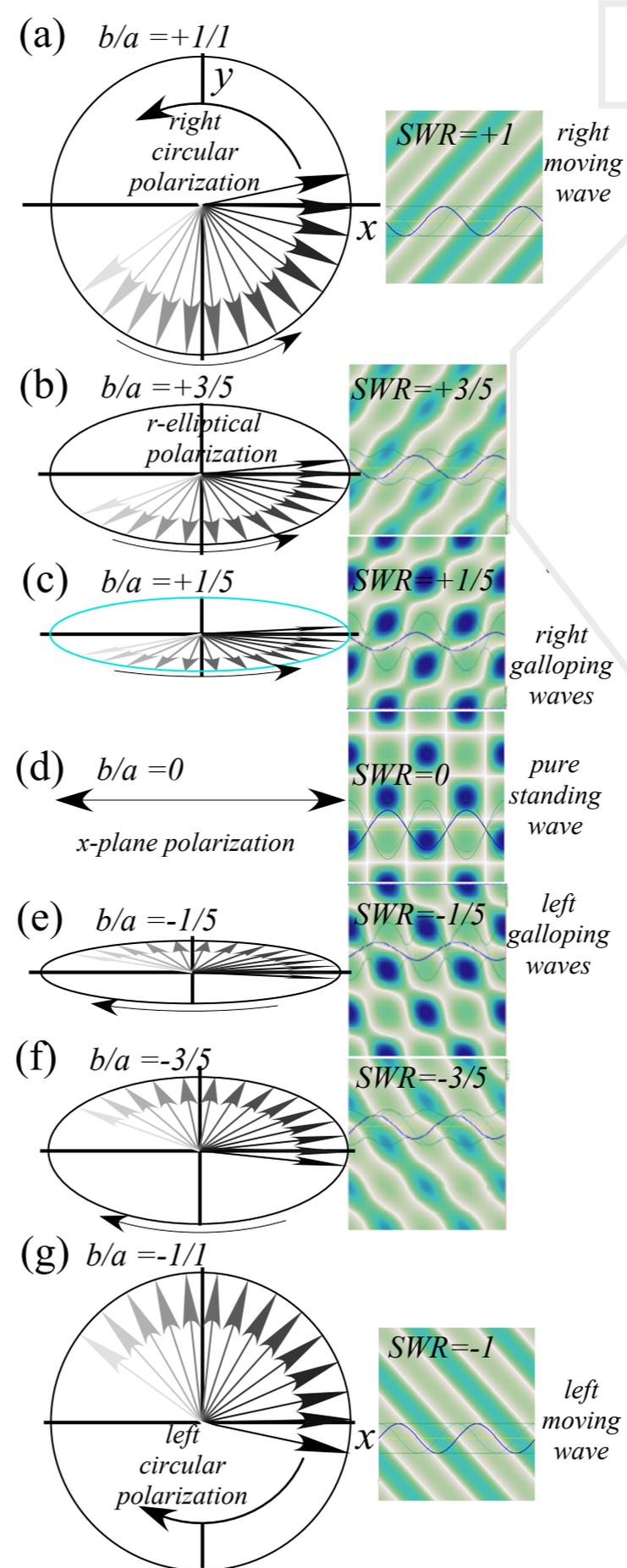


Fig. 6.3 (a-g) Elliptic polarization ellipses relate to galloping waves in Fig. 6.1. (h-i) Kepler anomalies.

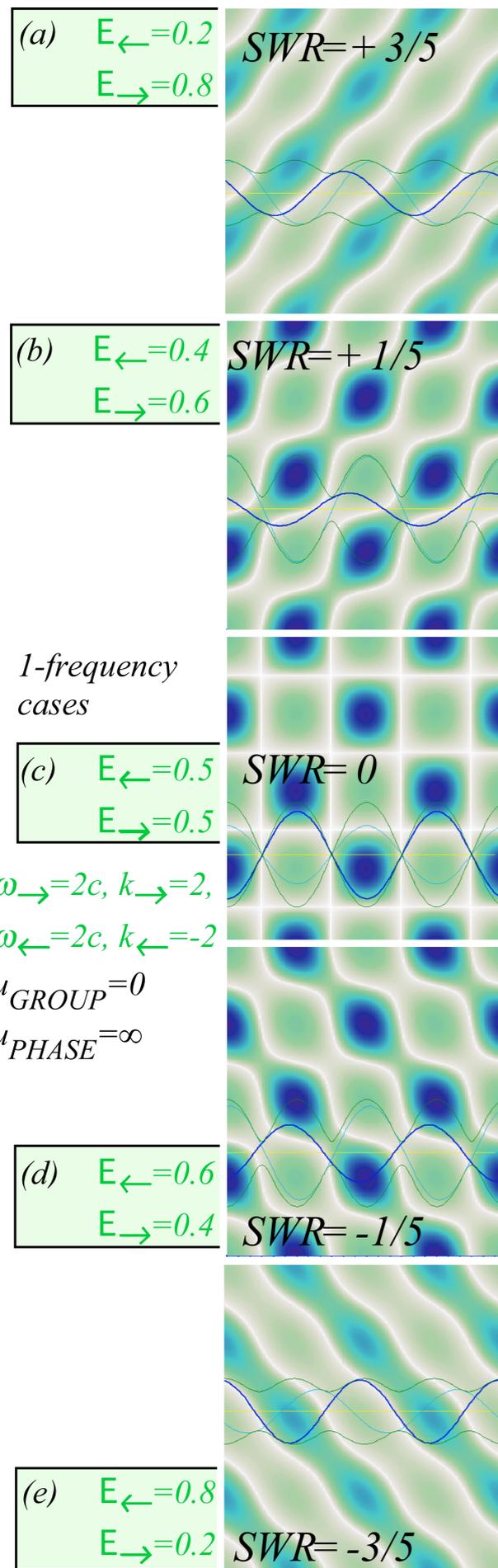


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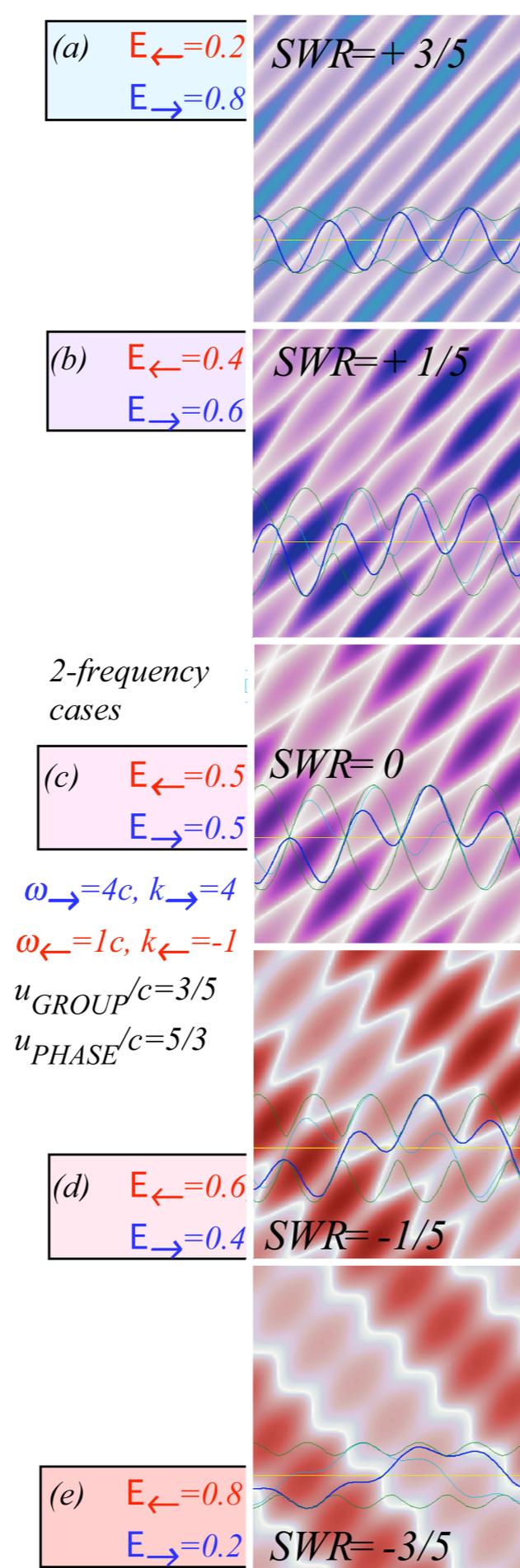


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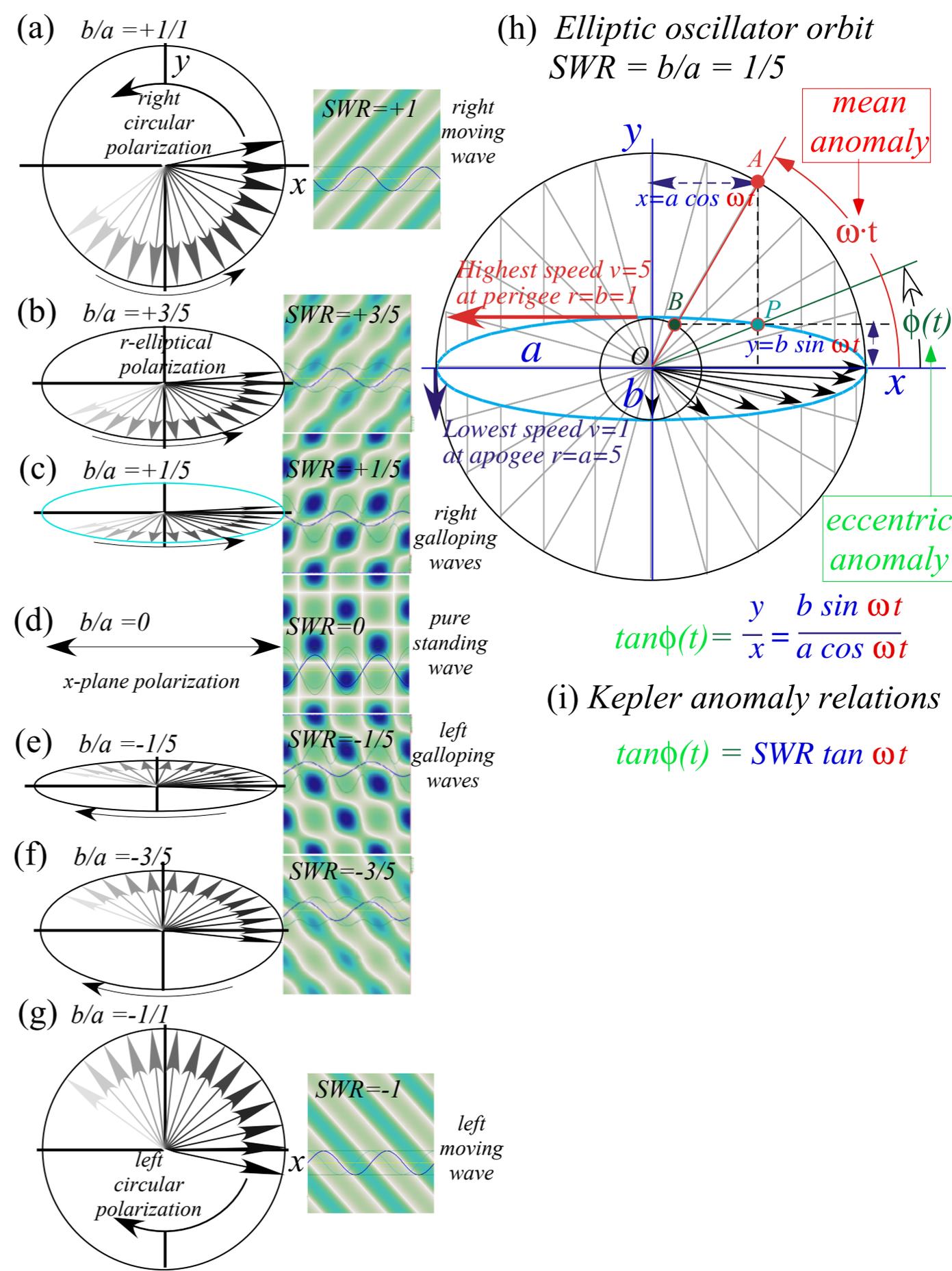
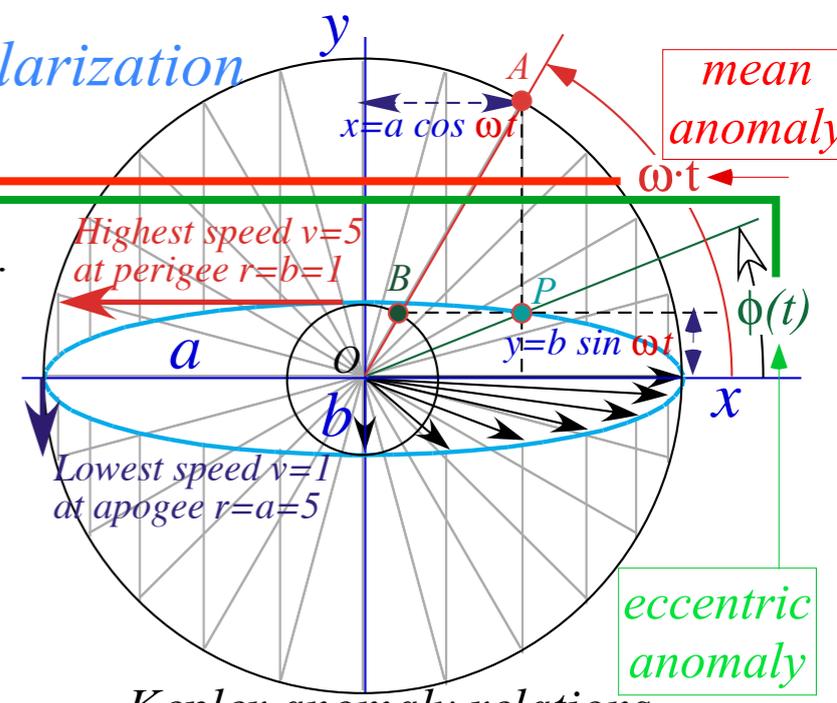


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Analogy between wave galloping, Keplerian IHO orbits, and optical polarization

We'll show wave galloping is analogous to Keplerian orbital motion of angles $\omega \cdot t$ and ϕ of orbits.

$$\tan \phi(t) = \frac{b}{a} \tan \omega \cdot t$$



Kepler anomaly relations

$$\tan \phi(t) = \frac{y}{x} = \frac{b \sin \omega t}{a \cos \omega t} = SWR \cdot \tan \omega t$$

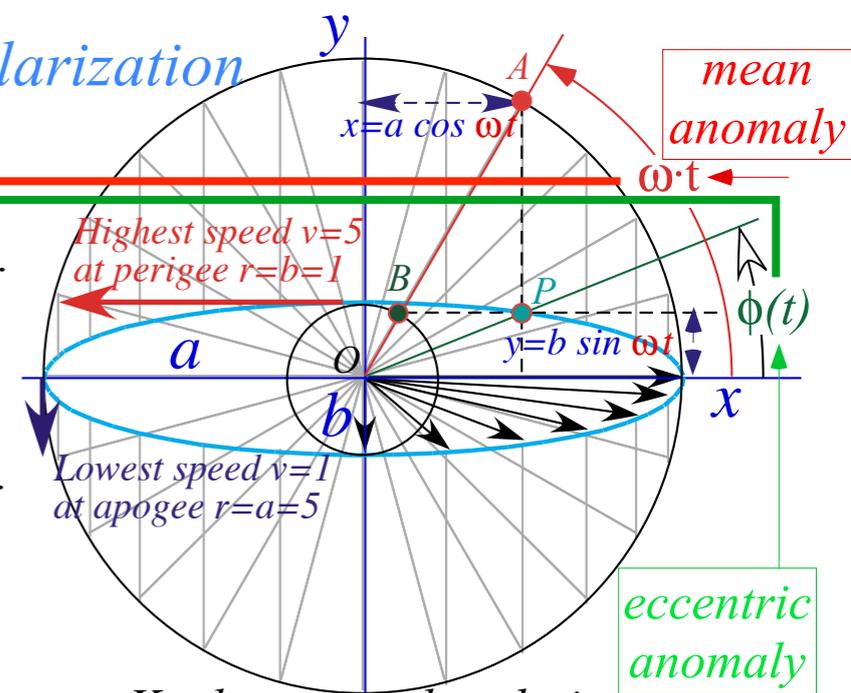
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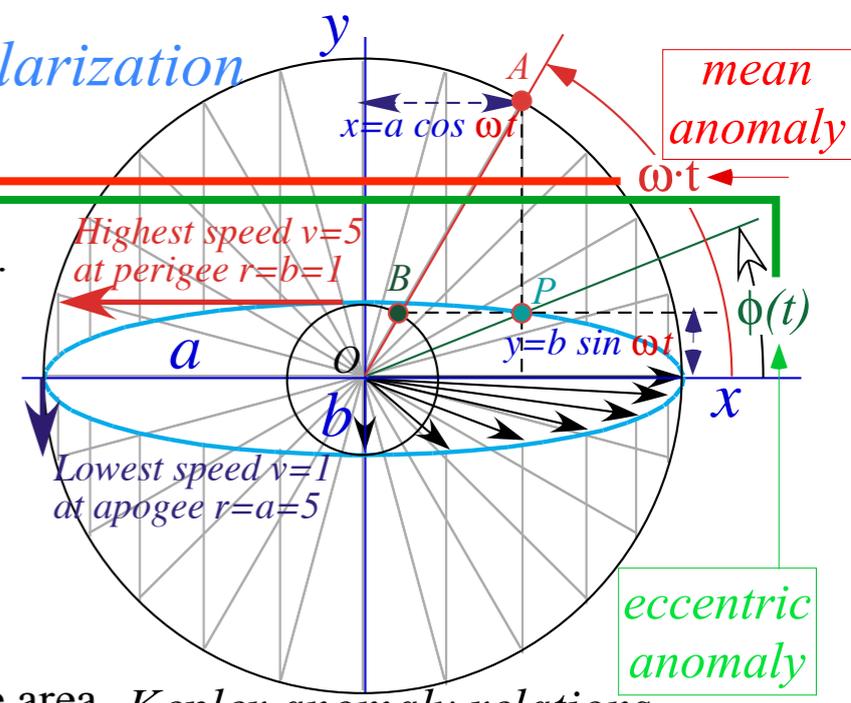
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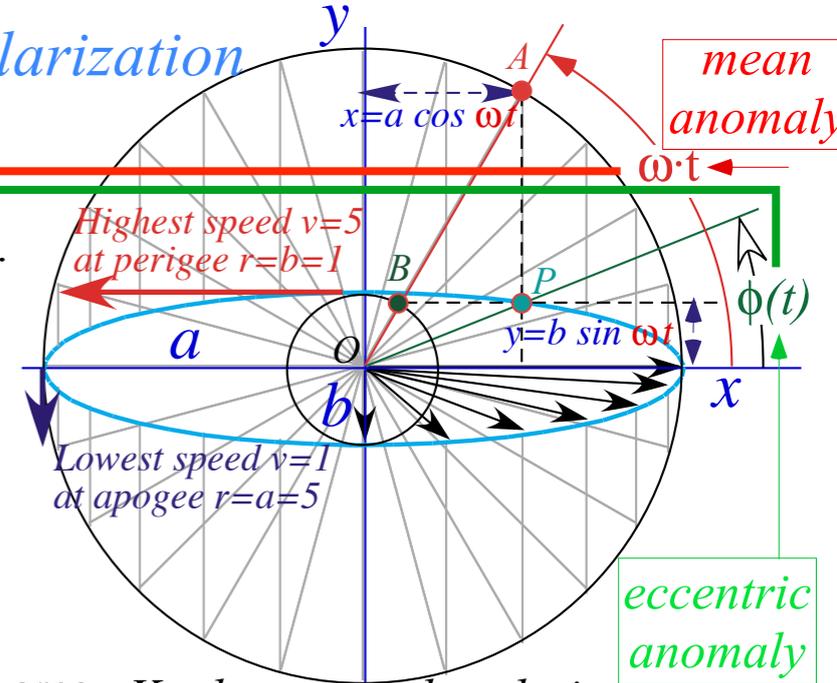
The product of angular momentum r^2 and $\dot{\phi}$ is orbital momentum, a constant proportional to ellipse area. *Kepler anomaly relations*

$$r^2 \frac{d\phi}{dt} = \text{constant} = (a^2 \cos^2 \omega t + b^2 \cdot \sin^2 \omega t) \frac{d\phi}{dt} = \omega \cdot ab$$



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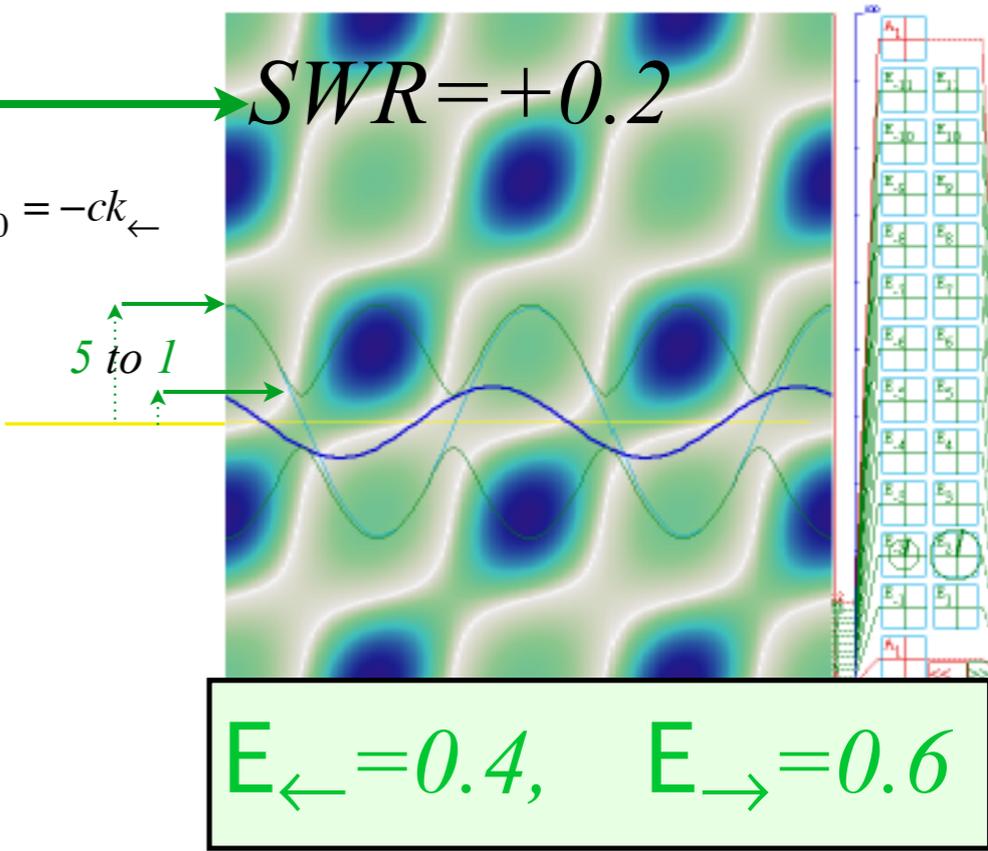
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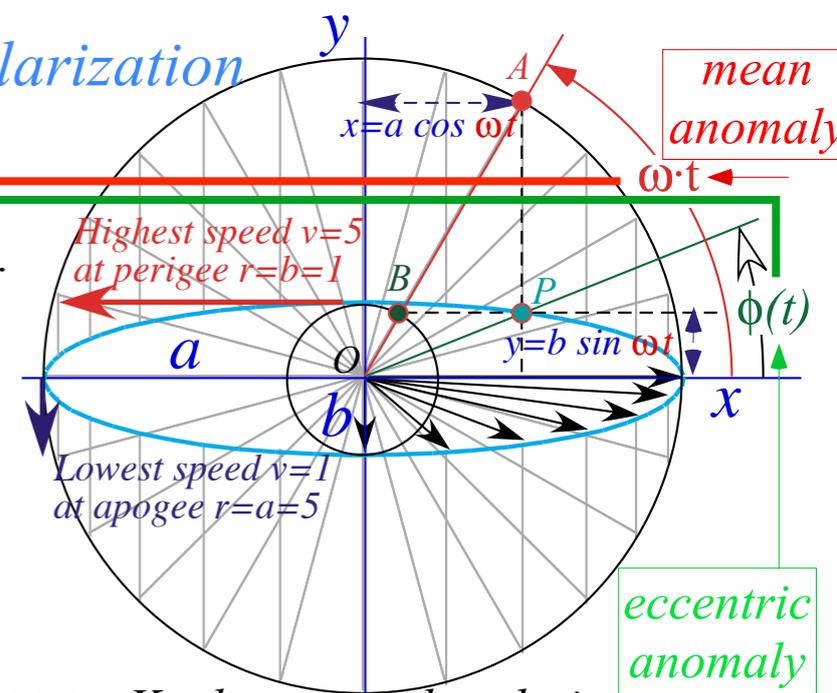
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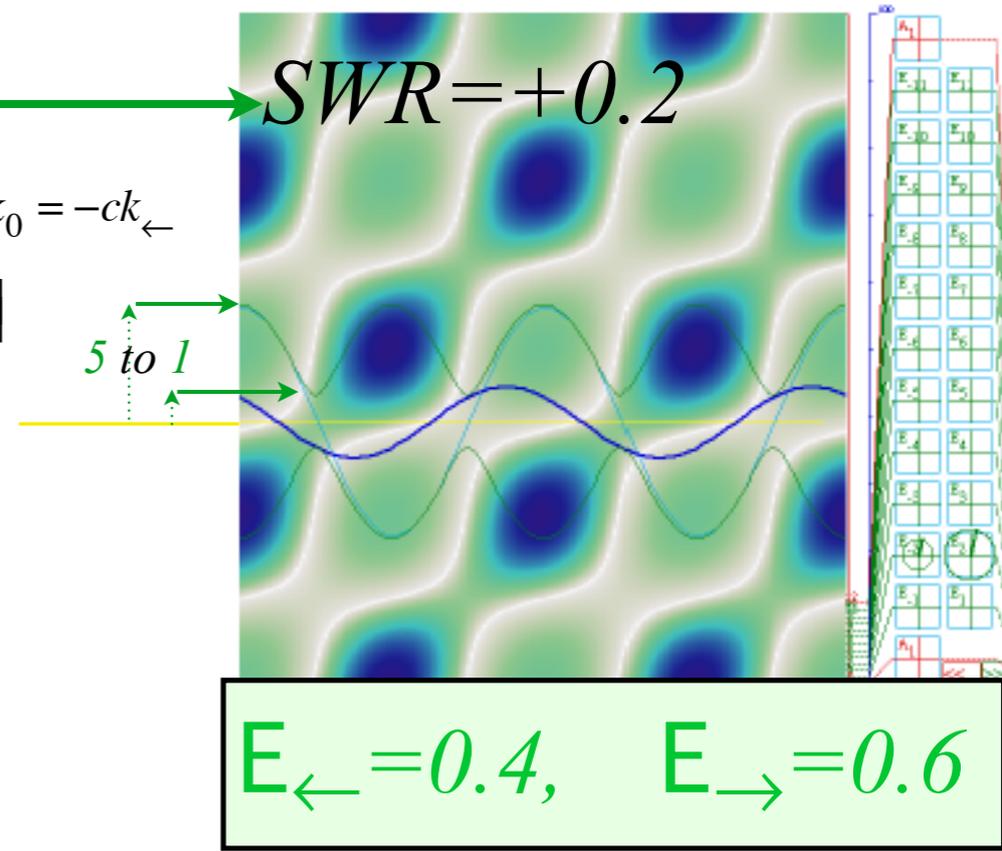
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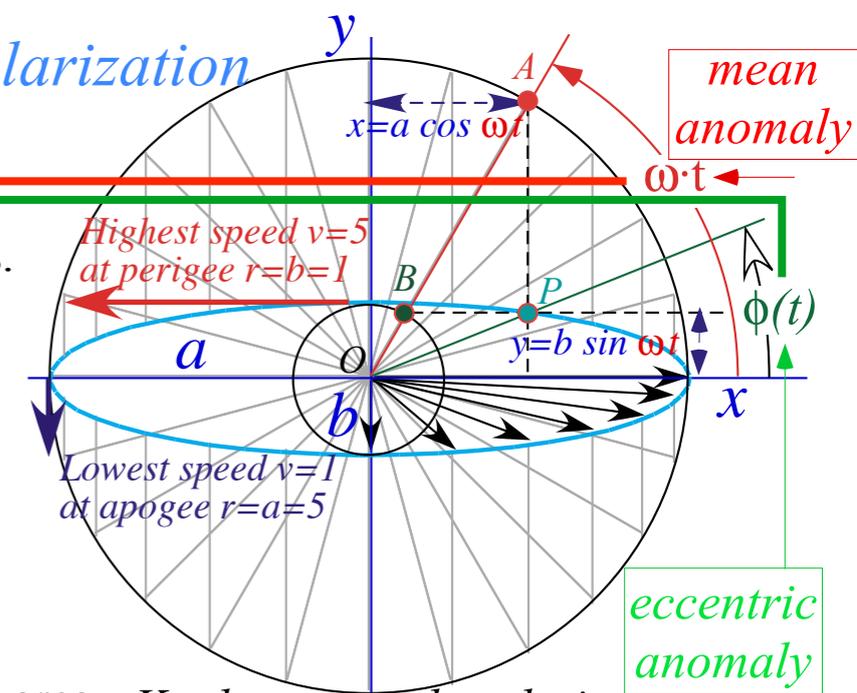
$$(A_{\rightarrow} + A_{\leftarrow}) [\cos k_0 x \cos \omega_0 t] = -(A_{\rightarrow} - A_{\leftarrow}) [\sin k_0 x \sin \omega_0 t]$$

Space $k_0 x$ varies with time $\omega_0 t$ in the same way that eccentric anomaly ϕ varies with $\omega \cdot t$.

$$\tan k_0 x = -SWR \cdot \cot \omega_0 t = SWR \cdot \tan \omega_0 \bar{t} \text{ where: } \omega_0 \bar{t} = \omega_0 t - \pi/2$$



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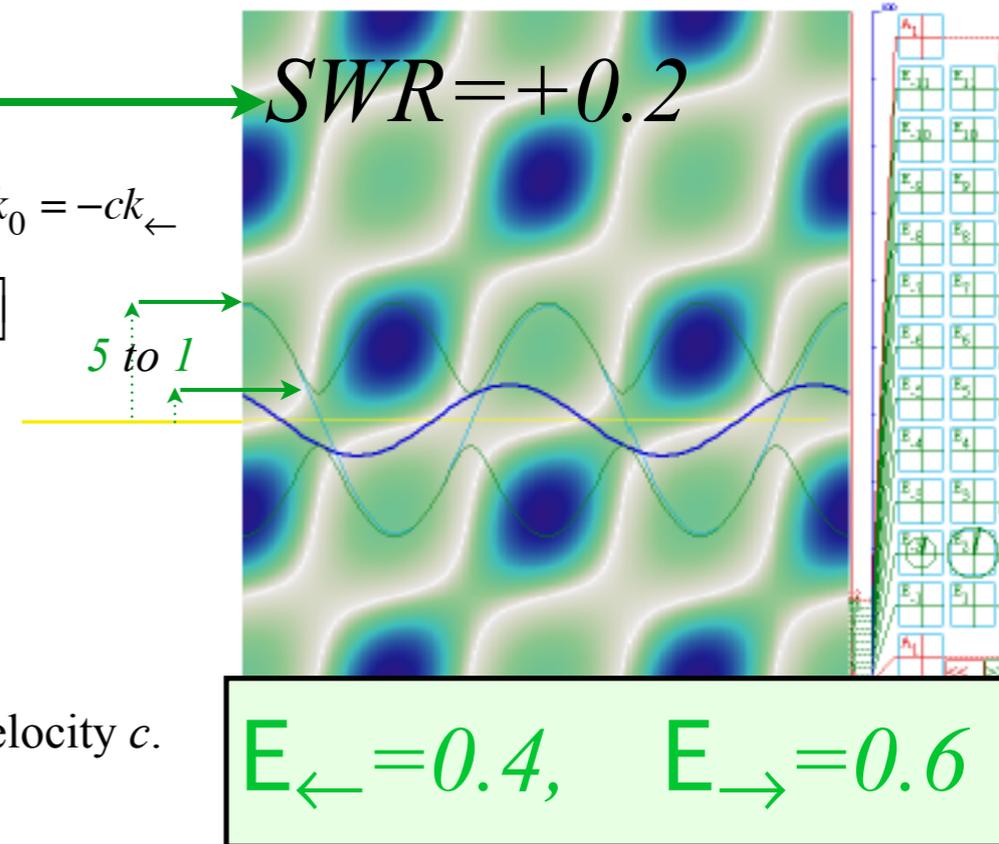
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Speed of galloping wave zeros is the time derivative of root location x in units of light velocity c .

$$\frac{dx}{dt} = c \cdot SWR \frac{\sec^2 \omega_0 \bar{t}}{\sec^2 k_0 x} = \frac{c \cdot SWR}{\cos^2 \omega_0 \bar{t} + SWR^2 \cdot \sin^2 \omega_0 \bar{t}} = \begin{cases} c \cdot SWR & \text{for: } \bar{t} = 0, \pi, 2\pi, \dots \\ c \cdot SWQ & \bar{t} = \pi/2, 3\pi/2, \dots \end{cases}$$



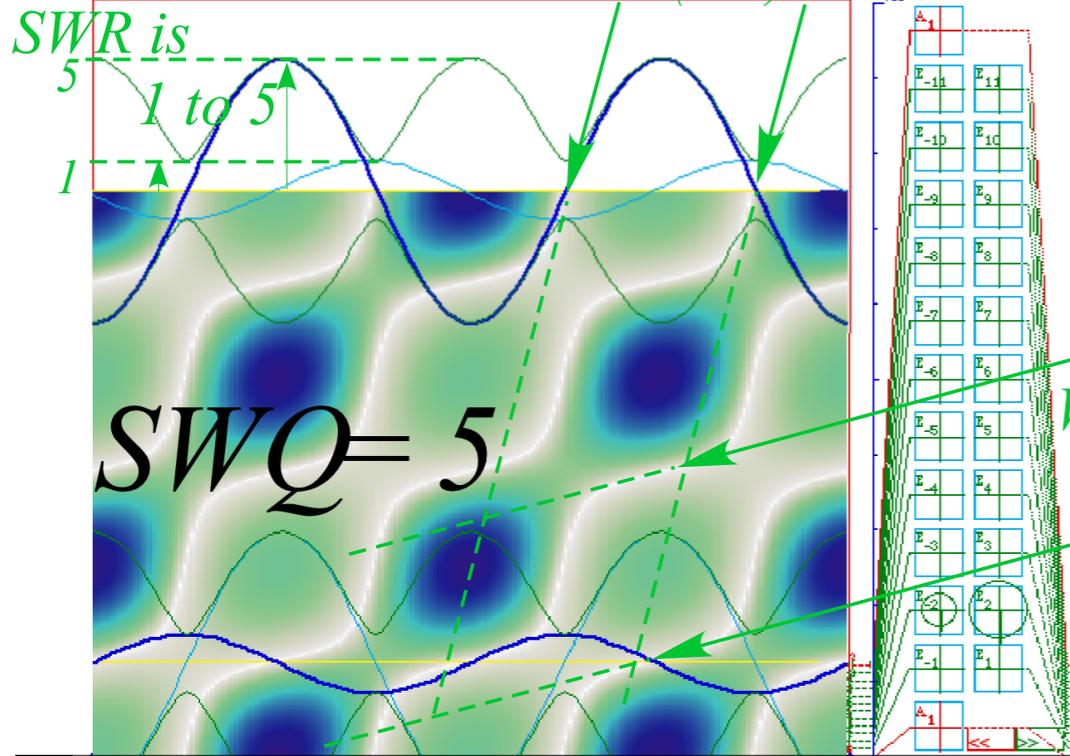
Wave-Zero Speed-Limits

Standing Wave Ratio SWR and Quotient SWQ

$$SWR = (E_{\rightarrow} - E_{\leftarrow}) / (E_{\rightarrow} + E_{\leftarrow}) = 1 / SWQ$$

$$SWR = 1/5$$

Wave zeros
"resting"
at $(1/5)c$



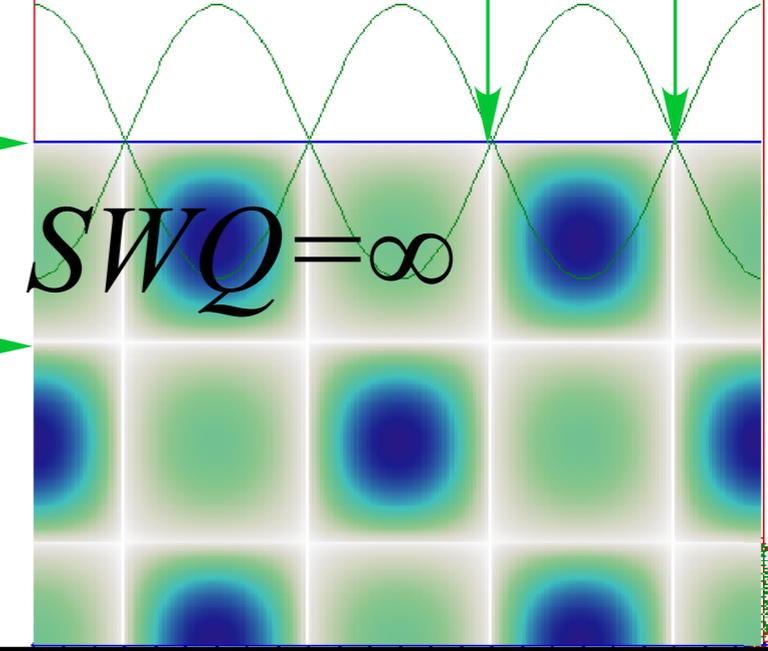
$$\begin{aligned} \omega_{\rightarrow} &= 2c & \omega_{\leftarrow} &= -2c \\ k_{\rightarrow} &= 2 & k_{\leftarrow} &= -2 \\ u_{GROUP} &= 0 & u_{PHASE} &= \infty \end{aligned}$$

Wave zeros
"galloping"
at $5c$

$$E_{\leftarrow} = 0.4, E_{\rightarrow} = 0.6$$

$$SWR = 0$$

Wave zeros
"standing"
at 0-speed

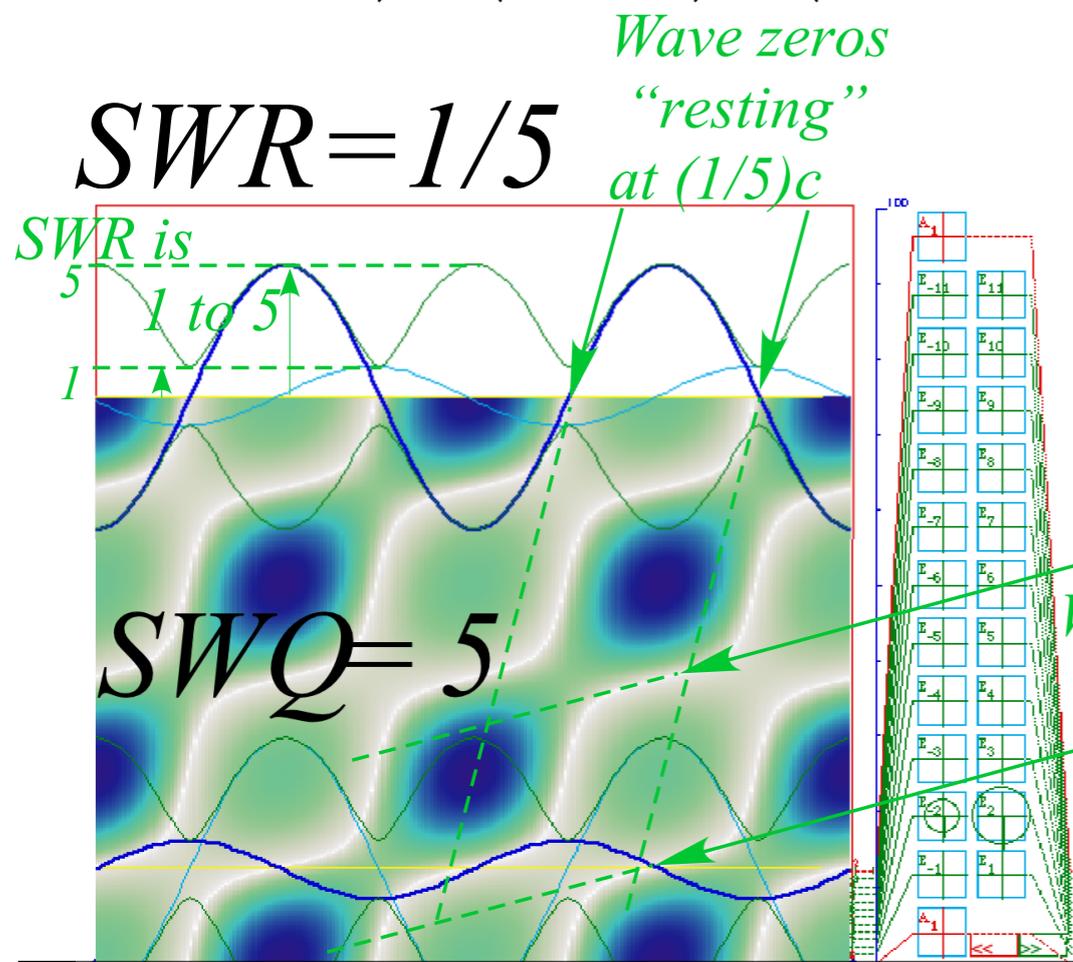


$$E_{\leftarrow} = 0.5, E_{\rightarrow} = 0.5$$

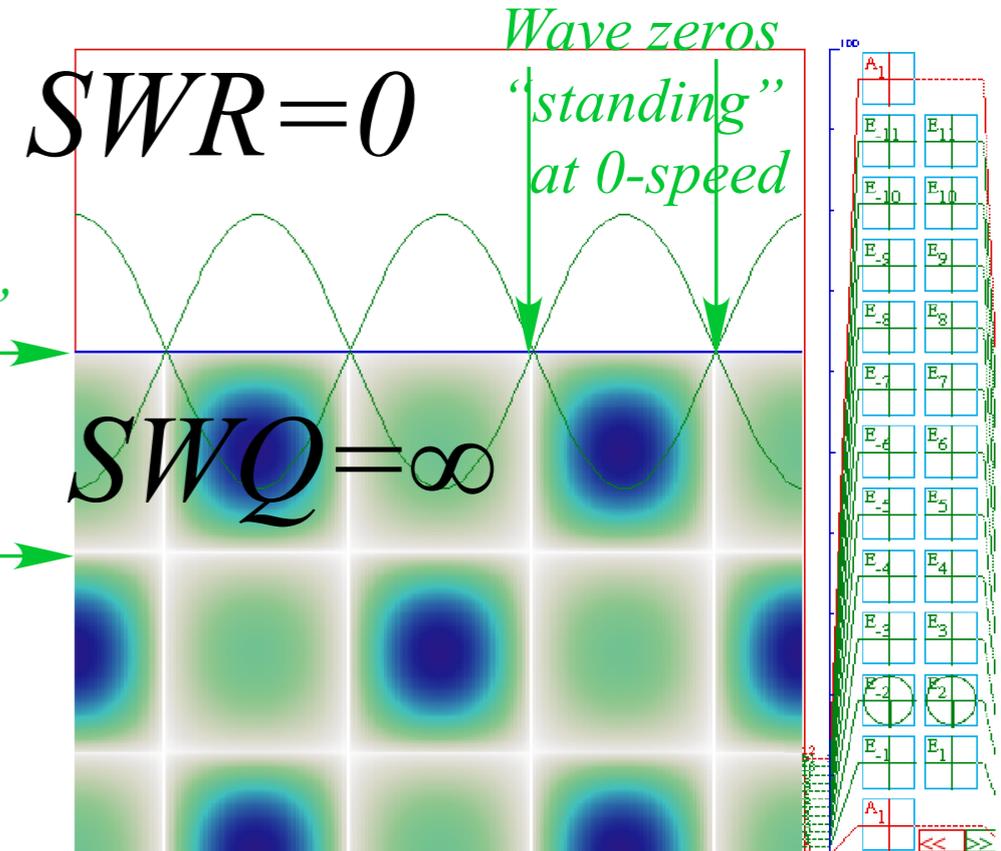
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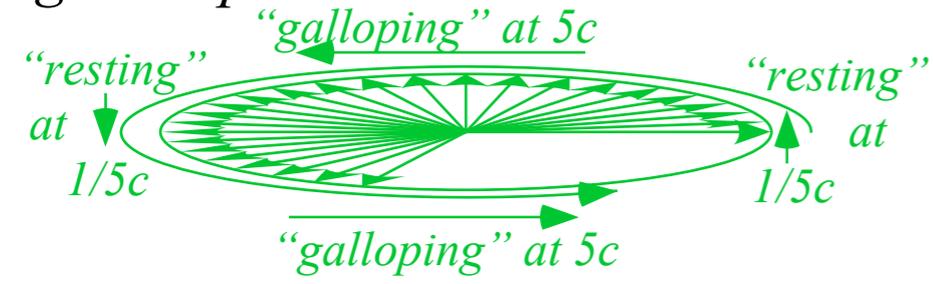
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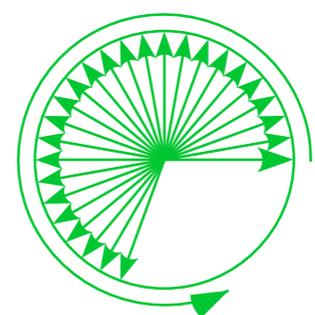
$$E_{\leftarrow} = 0.5, E_{\rightarrow} = 0.5$$

$$E_{\leftarrow} = 0.4, E_{\rightarrow} = 0.6$$

$SWR=1/5$ is analogous to (5-to-1) Right Elliptic Polarization



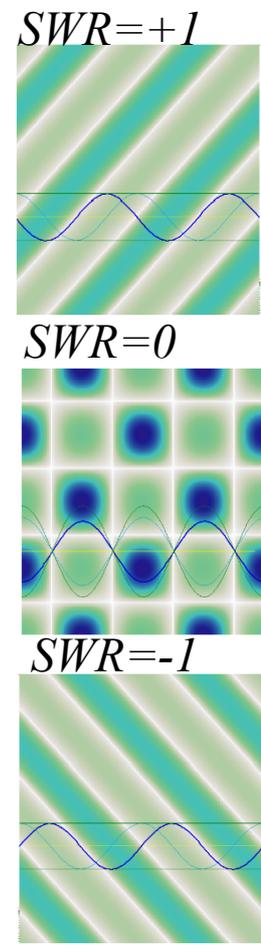
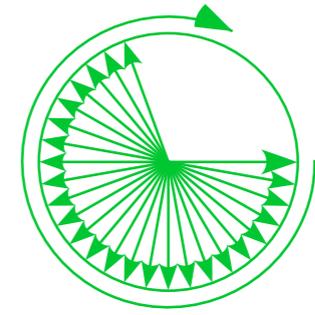
$SWR=1$ is analogous to (1,i) Right Circular Polarization



$SWR=0$ is analogous to (1,0) x-Plane Linear Polarization



$SWR=-1$ is analogous to (1,-i) Left Circular Polarization



Waves that go back in time - The Feynman-Wheeler Switchback

Minkowski Zero-Grids are Spacetime Switchbacks for $-u_{GROUP} < SWR < 0$

$\omega_{\rightarrow} = 4c$	$\omega_{\leftarrow} = 1c$
$k_{\rightarrow} = 4$	$k_{\leftarrow} = -1$
$u_{GROUP} = c3/5$	$u_{PHASE} = c5/3$

Group zero speed limit

$$\frac{u_{GROUP} + SWR}{1 + u_{GROUP} \cdot \frac{SWR}{c^2}} = 5c/11$$

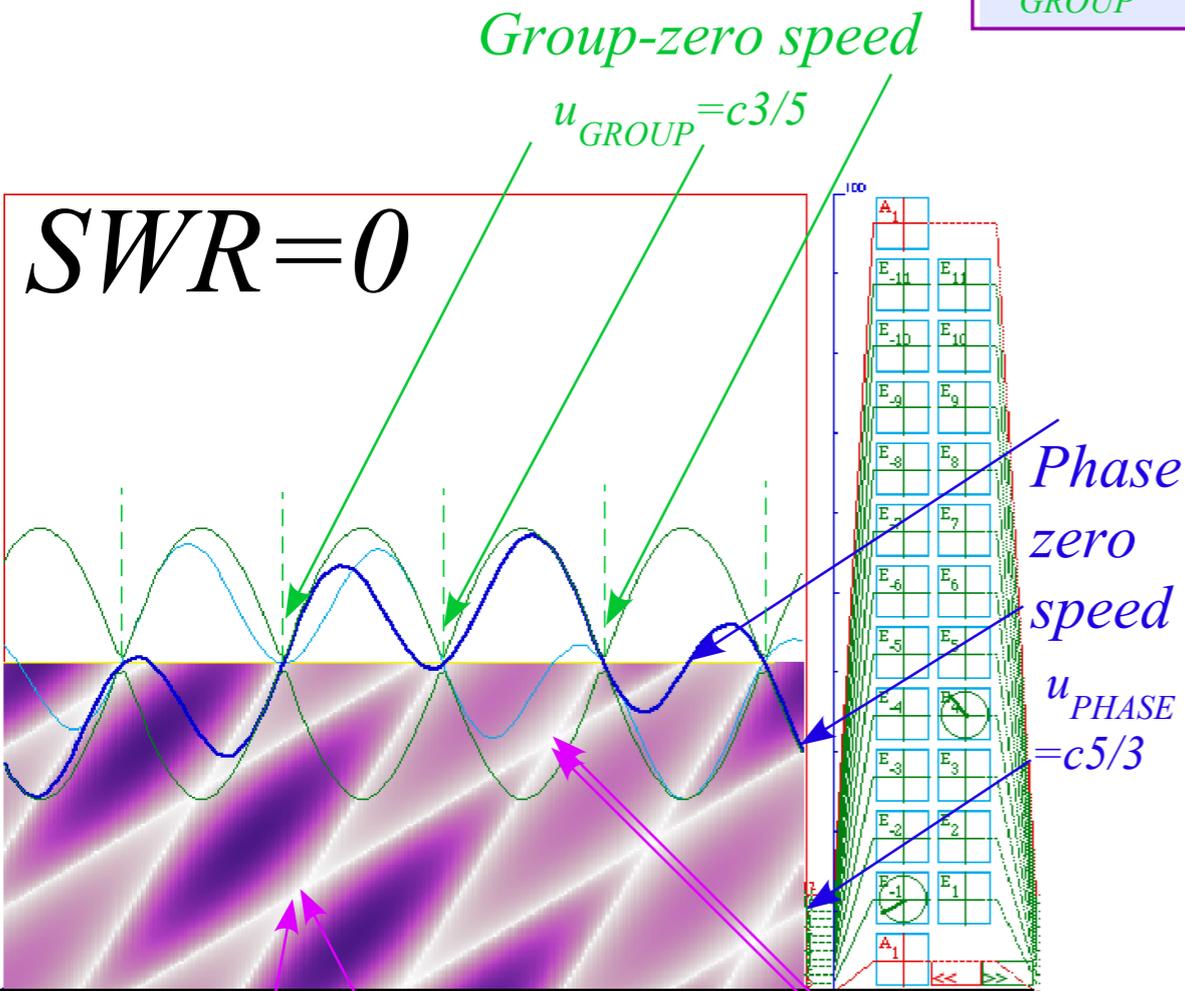
$$\frac{\frac{3}{5} + \frac{-1}{5}}{1 + \frac{3 \cdot -1}{5 \cdot 5}} = \frac{\frac{2}{5}}{\frac{22}{25}} = \frac{5}{11}$$

Phase "anti-zero" going "back-in-time"

Phase zero speed limit

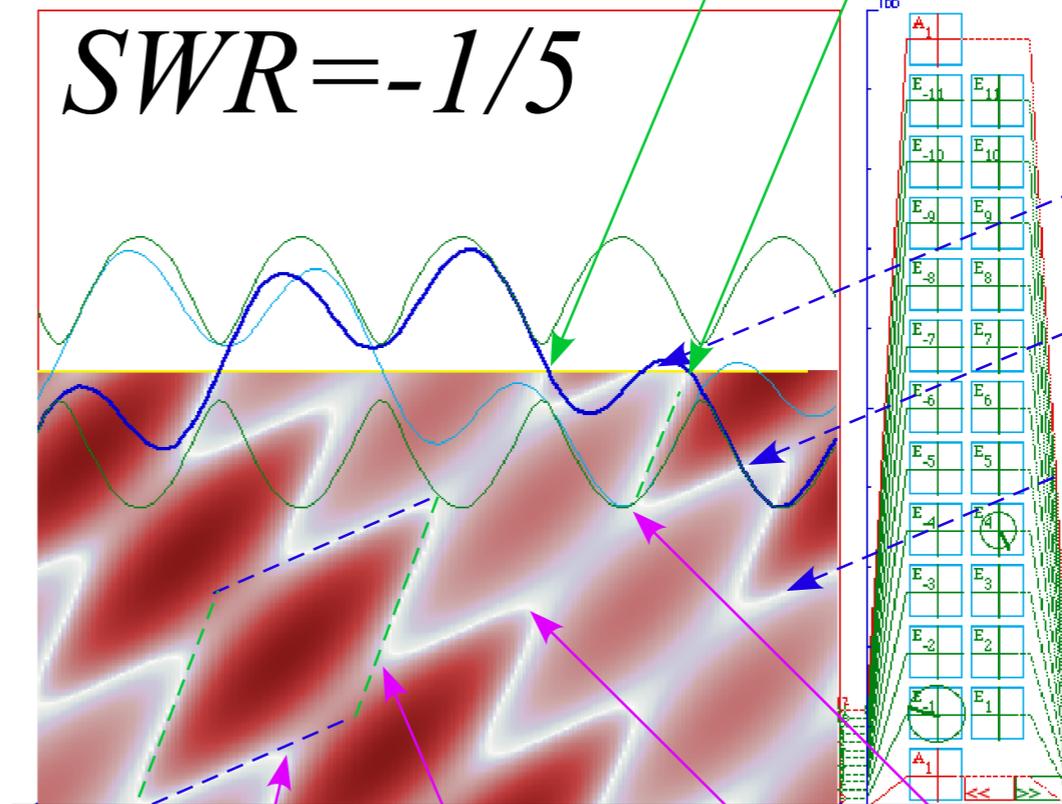
$$\frac{u_{PHASE} + SWR}{1 + u_{PHASE} \cdot \frac{SWR}{c^2}} = 11c/5$$

$$\frac{\frac{5}{3} + \frac{-1}{5}}{1 + \frac{5 \cdot -1}{3 \cdot 5}} = \frac{\frac{22}{15}}{\frac{10}{25}} = \frac{11}{5}$$



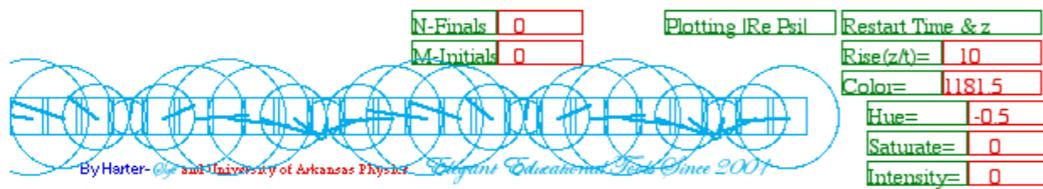
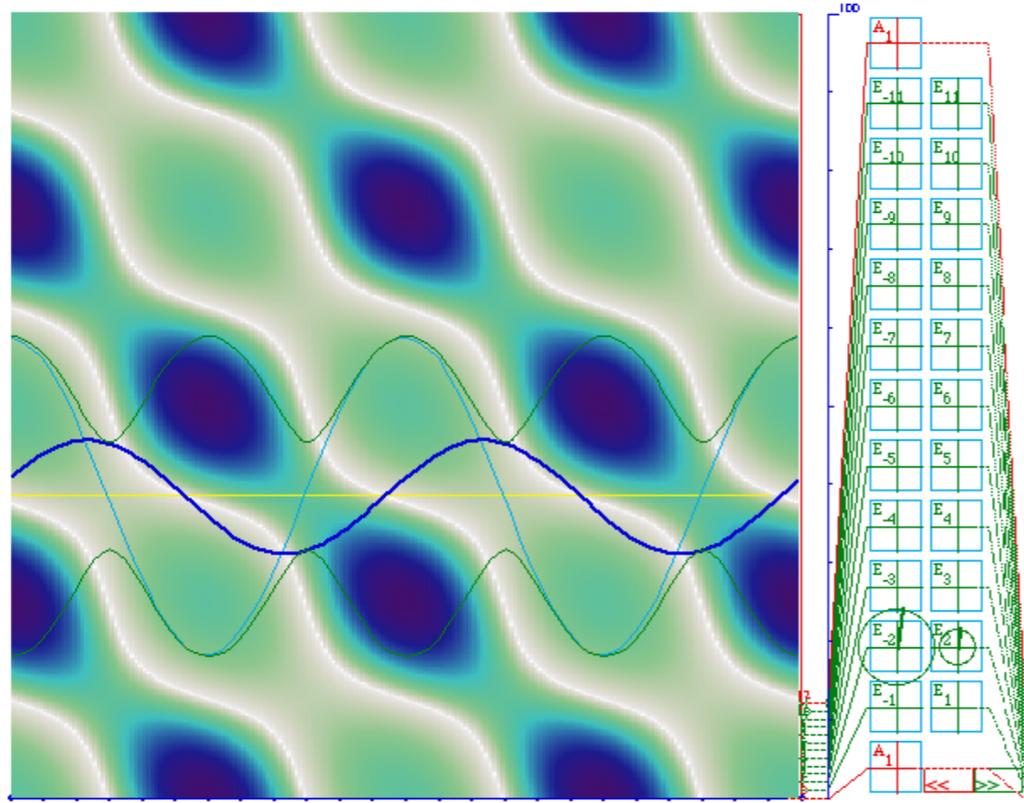
$E_{\leftarrow} = 0.5, E_{\rightarrow} = 0.5$

Wave zero-anti-zero annihilation and creation occur together at the same spacetime point for $SWR=0$

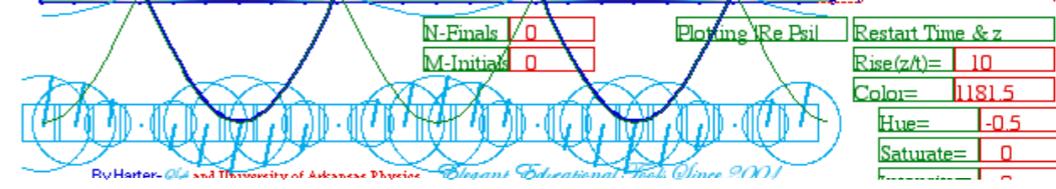
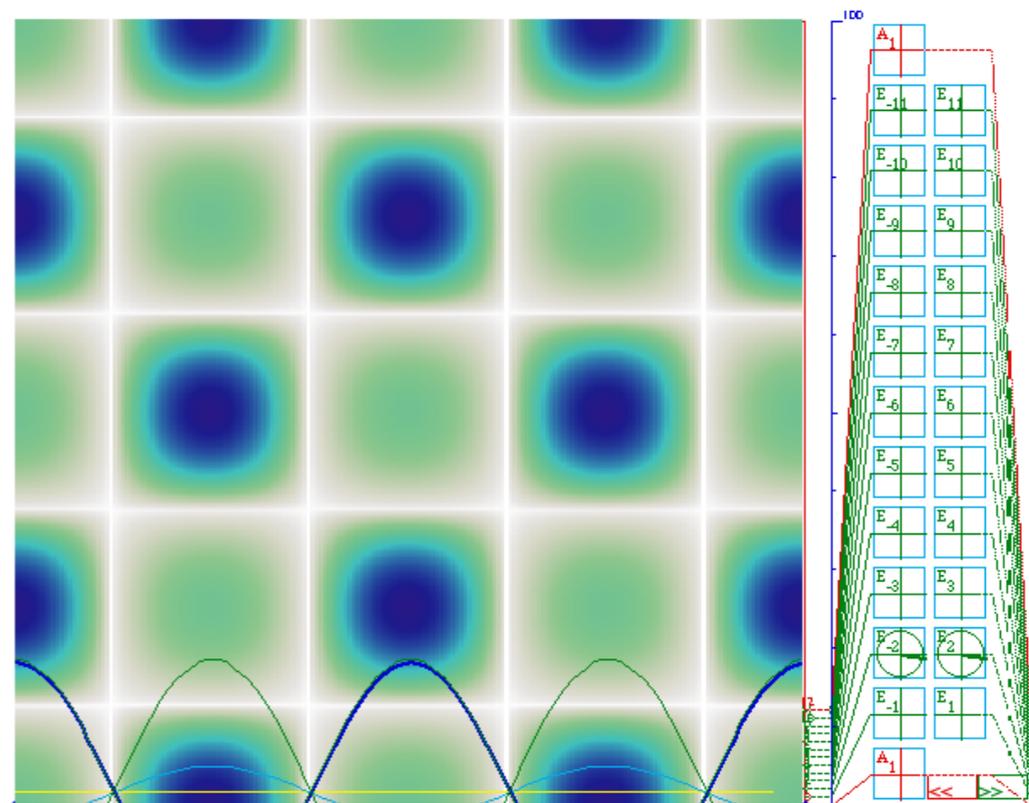


$E_{\leftarrow} = 0.6, E_{\rightarrow} = 0.4$

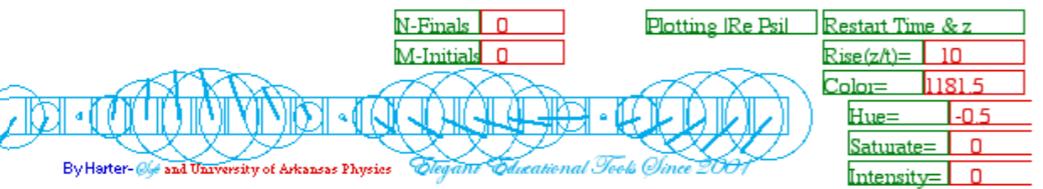
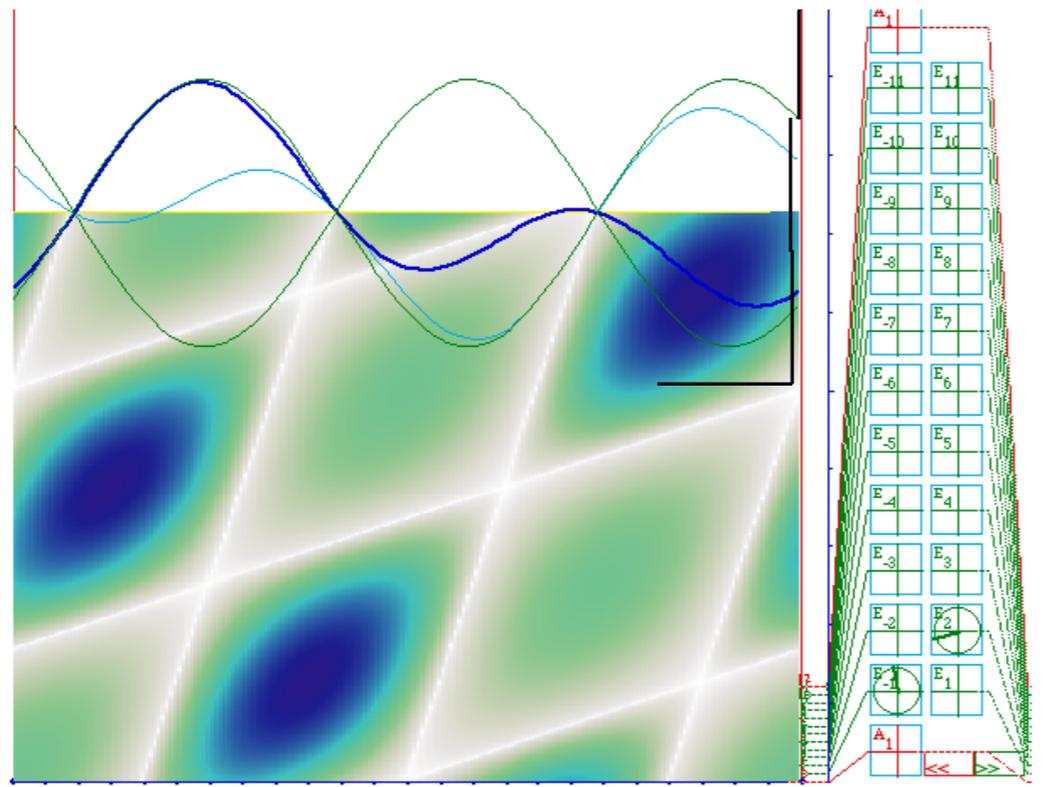
Wave zero-anti-zero annihilation and creation occur separately at different spacetime points for $-u_{GROUP} < SWR < 0$



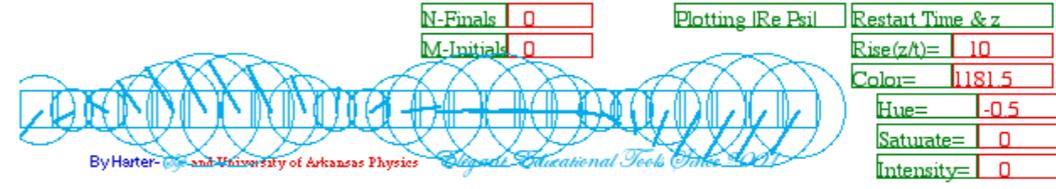
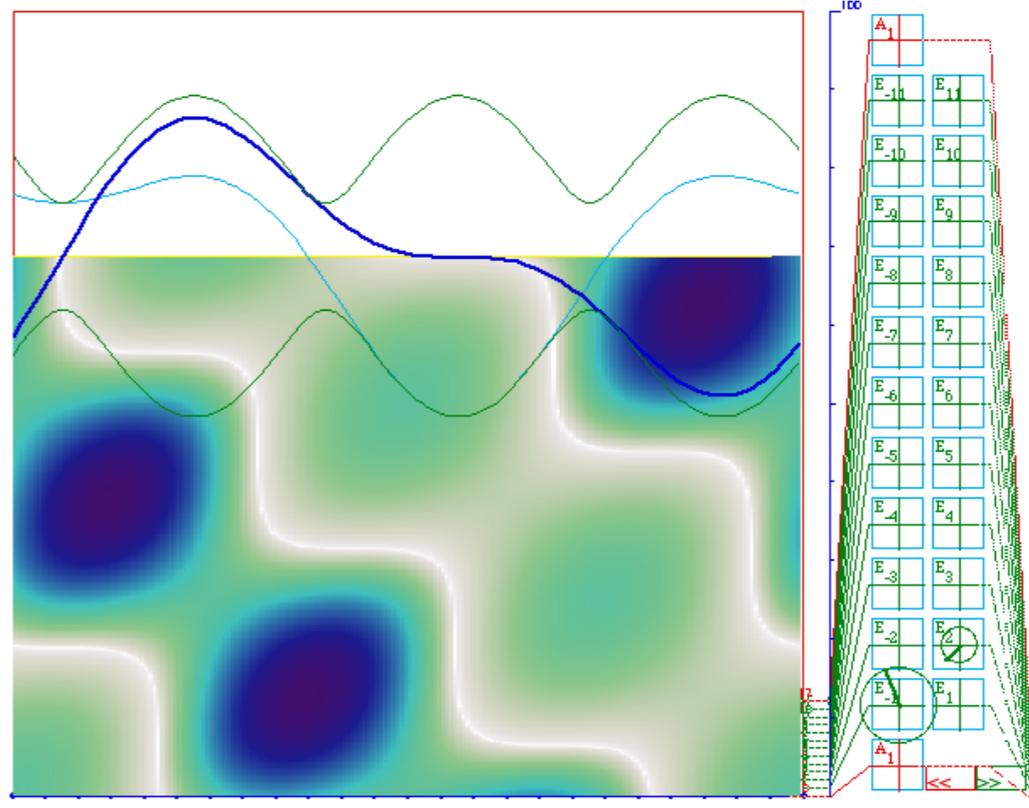
N-Finals	0	Plotting IRe Psi	Restart Time & z
M-Initials	0	Rise(z/t)=	10
		Color=	1181.5
		Hue=	-0.5
		Saturate=	0
		Intensity=	0



N-Finals	0	Plotting IRe Psi	Restart Time & z
M-Initials	0	Rise(z/t)=	10
		Color=	1181.5
		Hue=	-0.5
		Saturate=	0
		Intensity=	0



N-Finals	0	Plotting IRe Psi	Restart Time & z
M-Initials	0	Rise(z/t)=	10
		Color=	1181.5
		Hue=	-0.5
		Saturate=	0
		Intensity=	0



N-Finals	0	Plotting IRe Psi	Restart Time & z
M-Initials	0	Rise(z/t)=	10
		Color=	1181.5
		Hue=	-0.5
		Saturate=	0
		Intensity=	0

At High Speed 2-CW Modes Look More Like 1-CW Beams

$$\psi = E \sqrt{\frac{\epsilon_0}{h\nu}}$$

Various combinations of *opposite-k* 1-CW beams occur with open boundaries.

E-wave: $\mathbf{E} = \mathbf{E}_{\rightarrow} e^{i(k_{\rightarrow}x - \omega_{\rightarrow}t)} + \mathbf{E}_{\leftarrow} e^{i(k_{\leftarrow}x - \omega_{\leftarrow}t)}$ is related to Ψ -wave: $\Psi = \psi_{\rightarrow} e^{i(k_{\rightarrow}x - \omega_{\rightarrow}t)} + \psi_{\leftarrow} e^{i(k_{\leftarrow}x - \omega_{\leftarrow}t)}$

Standing Wave Ratio (or Quotient)

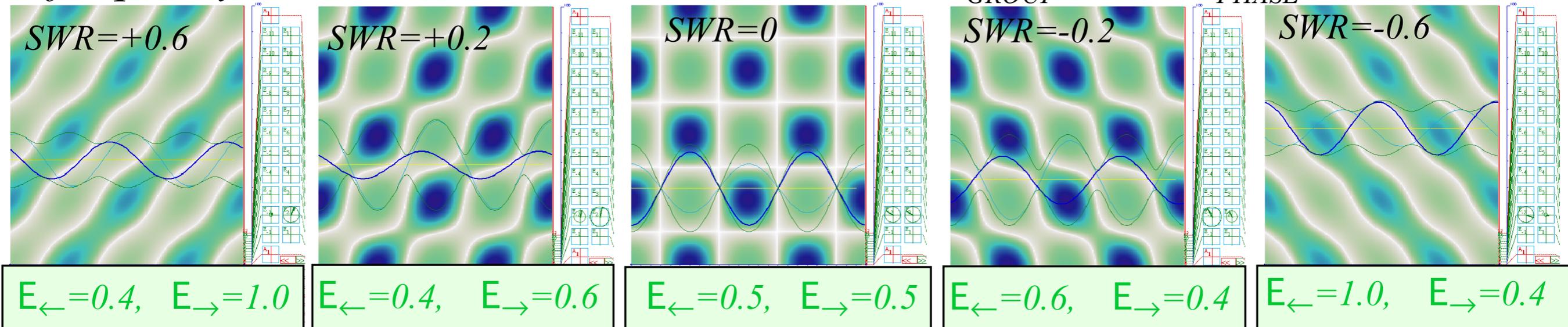
$$SWR = (E_{\rightarrow} - E_{\leftarrow}) / (E_{\rightarrow} + E_{\leftarrow}) = 1/SWQ$$

key
numbers

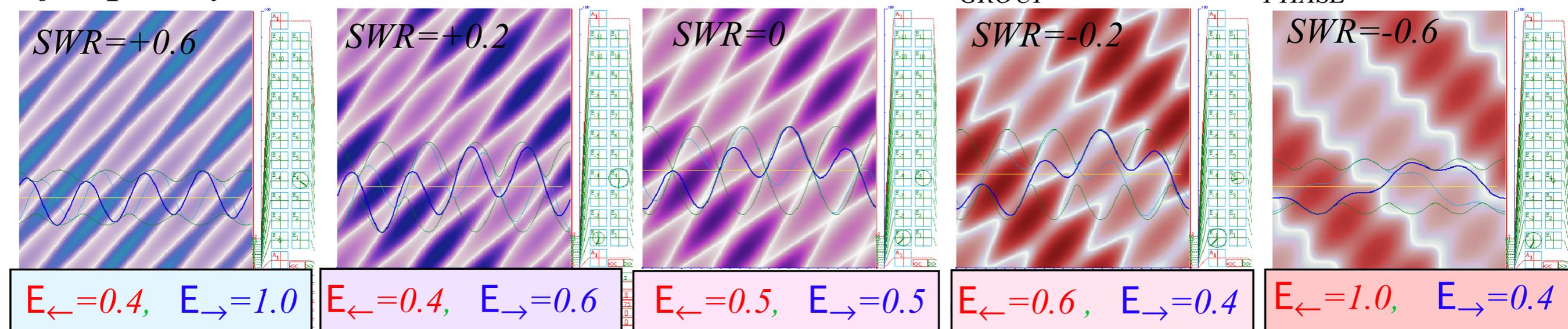
Wave Group (or Phase) Velocity

$$u_{GROUP}/c = (\omega_{\rightarrow} - \omega_{\leftarrow}) / (\omega_{\rightarrow} + \omega_{\leftarrow}) = c/u_{PHASE}$$

1-frequency case : $\omega_{\rightarrow} = 2c, k_{\rightarrow} = 2, \omega_{\leftarrow} = 2c, k_{\leftarrow} = -2$ gives: $u_{GROUP} = 0$ and $u_{PHASE} = \infty$



2-frequency case : $\omega_{\rightarrow} = 4c, k_{\rightarrow} = 4, \omega_{\leftarrow} = 1c, k_{\leftarrow} = -1$ gives: $u_{GROUP}/c = 3/5$ and $u_{PHASE}/c = 5/3$



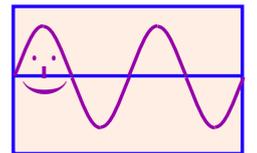
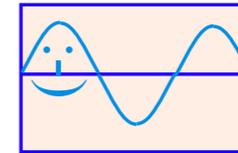
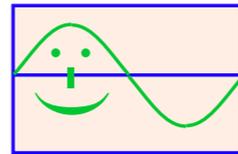
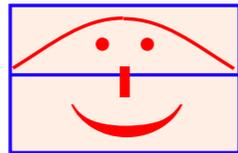
1st Quantization: Quantizing phase variables ω and k
Understanding how quantum transitions require “mixed-up” states
Closed cavity vs Ring cavity

Quantized ω and k Counting wave kink numbers

If everything is made of waves then we expect *quantization* of everything because waves only thrive if *integral* numbers n of their “kinks” fit into whatever structure (box, ring, etc.) they’re supposed to live. The numbers n are called *quantum numbers*.

OK box quantum numbers: $n=1$ $n=2$ $n=3$ $n=4$

(+ integers only)



Some

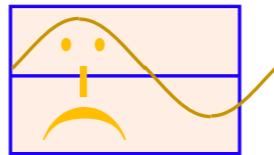
NOT OK numbers: $n=0.67$

too fat!



$n=1.7$

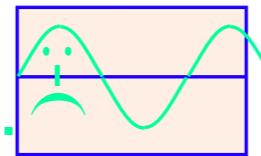
too thin!



$n=2.59$

wrong color again!

misfits...



$n=4$

...not tolerated!

NOTE: We’re using “false-color” here.

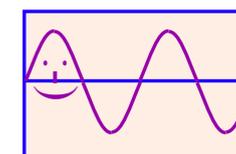
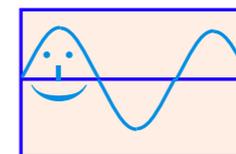
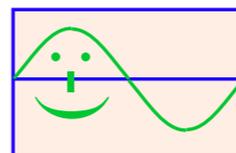
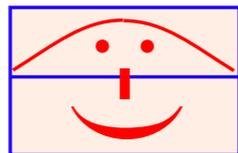
This doesn’t mean a system’s energy can’t vary continuously between “OK” values $E_1, E_2, E_3, E_4, \dots$

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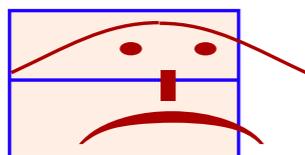
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Some

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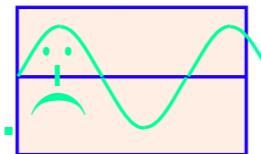
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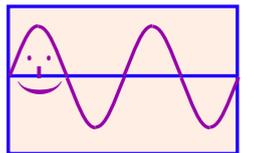
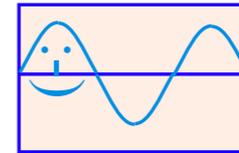
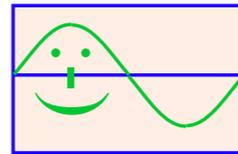
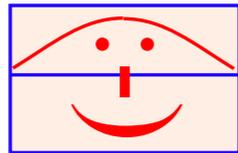
In fact its state can be a linear combination of any of the “OK” waves $|E_1\rangle, |E_2\rangle, |E_3\rangle, |E_4\rangle, \dots$

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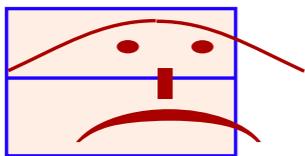
(+ integers only)



Some

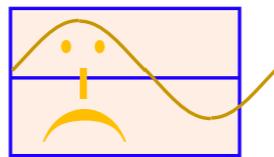
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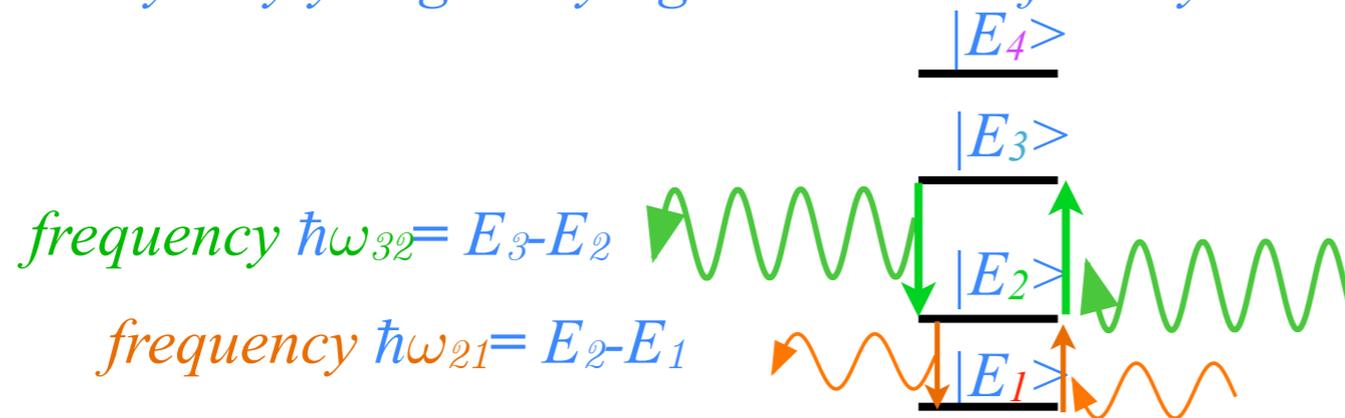
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That's the only way you get any light in or out of the system to “see” it.

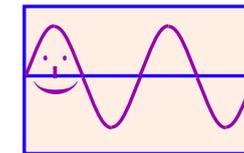
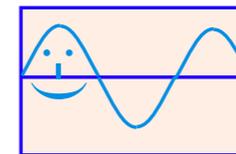
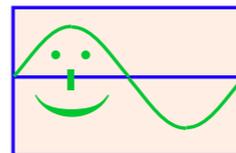
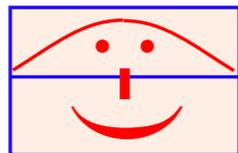


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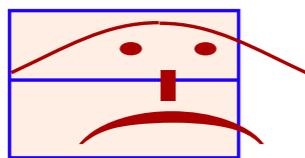
(+ integers only)



Some

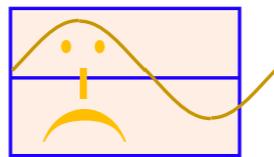
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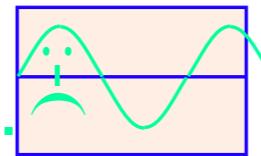
too thin!



$n=2.59$

wrong color again!

misfits...



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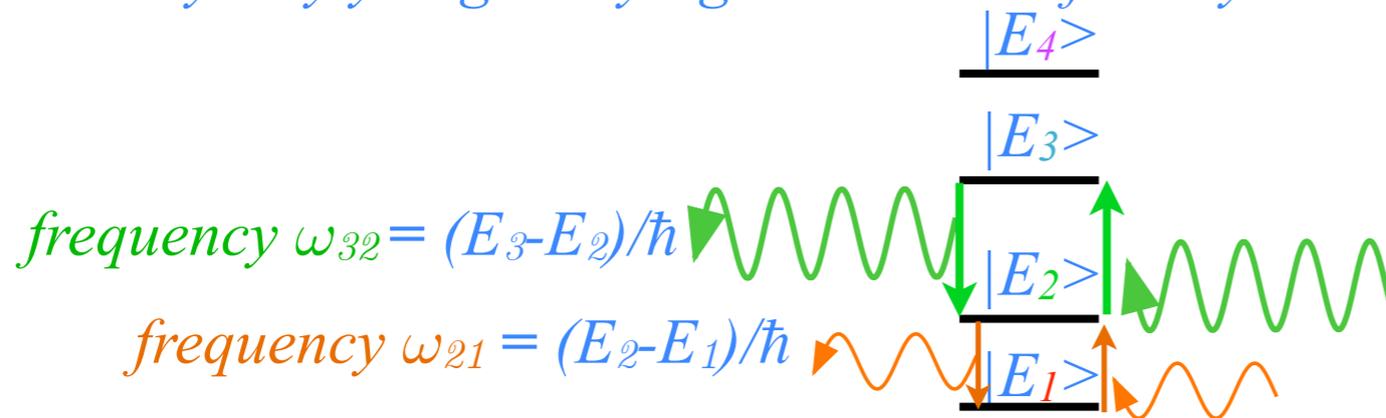
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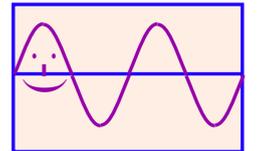
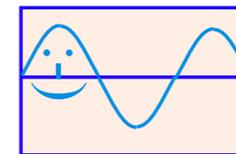
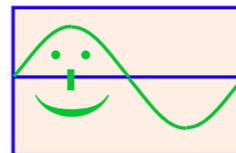
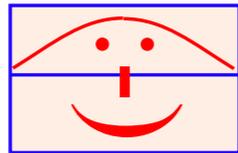
These *eigenstates* are the only ways the system can “play dead” ...
... “sleep with the fishes”...

Quantized ω and k Counting wave kink numbers

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OK box quantum numbers: $n=1$ $n=2$ $n=3$ $n=4$

(+ integers only)



Some

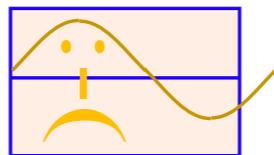
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too fat!



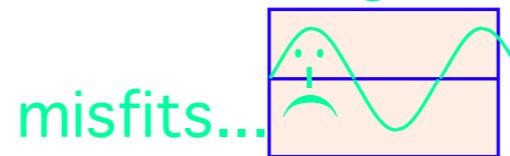
$n=1.7$

too thin!



$n=2.59$

wrong color again!



$n=4$

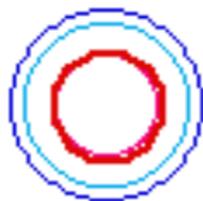
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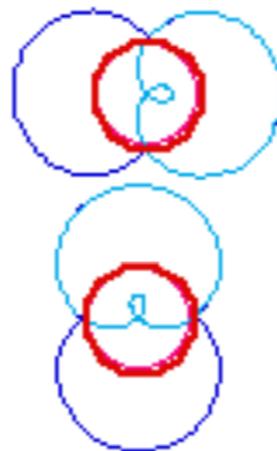
Rings tolerate a *zero* (kinkless) quantum wave but require \pm integral wave number.

OK ring quantum numbers: $m=0$

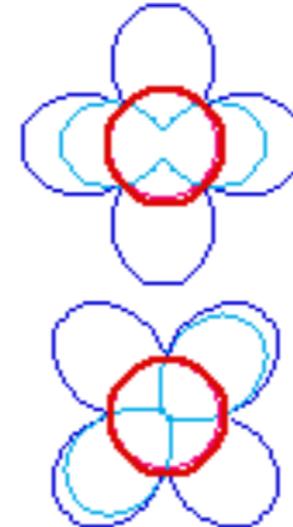
(\pm integral number of wavelengths)



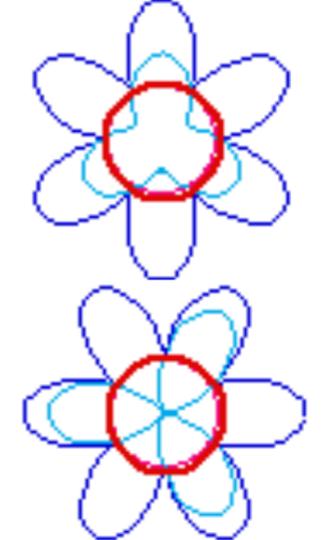
$m=\pm 1$



$m=\pm 2$



$m=3$



Bohr’s models of *atomic spectra* (1913-1923) are beginnings of *quantum wave mechanics* built on *Planck-Einstein* (1900-1905) relation $E=h\nu$. *DeBroglie* relation $p=h/\lambda$ comes around 1923.

Lecture 30 ended here

2nd Quantization: Quantizing amplitudes (“photons”, “vibrons”, and “what-ever-ons”)

Introducing coherent states (What lasers use)

Analogy with (ω, k) wave packets

Wave coordinates need coherence

Quantized *Amplitude* Counting “photon” number

Planck’s relation $E=Nh\nu$ began as a tentative axiom to explain low-T light. Then he tried to disavow it! Einstein picked it up in his 1905 paper. Since then its use has grown enormously and continues to amaze, amuse (or bewilder) all who study it. A current view is that it represents the *quantization* of optical field *amplitude*. We picture this below as *N-photon* wave states for each box-mode of *m* wave kinks.

