

# Quantum Revivals of Morse Oscillators and Farey-Ford Geometry

Alvason Zhenhua Li<sup>a,b</sup>, William G. Harter<sup>a</sup>

<sup>a</sup>*Microelectronics-Photonics Program, Department of Physics, University of Arkansas, Fayetteville, AR 72701, USA*

<sup>b</sup>*Present Address: Fred Hutchinson Cancer Research Center, Seattle, WA 98109, USA*

---

## Abstract

Analytical eigensolutions for Morse oscillators are used to investigate quantum resonance and revivals and show how Morse anharmonicity affects revival times. A minimum semi-classical Morse revival time  $T_{min-rev}$  found by Heller is related to a complete quantum revival time  $T_{rev}$  using a quantum deviation  $\delta_N$  parameter that in turn relates  $T_{rev}$  to the maximum quantum beat period  $T_{max-beat}$ . Also, number theory of Farey and Thales-circle geometry of Ford is shown to elegantly analyze and display fractional revivals. Such quantum dynamical analysis may have applications for spectroscopy or quantum information processing and computing.

---

## 1. Introduction

Wavepacket dynamics has a long history that has more recently been accelerated by graphics that exhibit spacetime behavior. Such studies began with revivals in cavity QED simulations by Eberly [1] and later simulations of molecular rovibronic dynamics [2, 3]. Ultrafast laser spectroscopy made it possible to observe wavepacket resonance and localized periodic motion in experimental situations [4, 5, 6] involving AMOP dynamics [6, 7]. This helped reveal new physics and chemistry of ultrafast spectroscopy [7, 8].

Some of this involves symmetry and number theoretic properties of wavepacket space-time structure, a still largely unexplored field. The following development is based upon earlier  $Cn$ -group and Farey-sum-tree [9] analysis of

---

*Email addresses:* [alvali@fhcrc.org](mailto:alvali@fhcrc.org) (Alvason Zhenhua Li), [wharter@uark.edu](mailto:wharter@uark.edu) (William G. Harter)

quantum rotors [10, 11] as cited by Schleich et al. [12, 13] for possible numeric factorizing applications. That work treated only R(2) rings or 1D infinite-wells but nevertheless revealed general symmetry properties.

Here Morse oscillators are shown to share Farey-sum revival structure of R(2) rings or 1D infinite-wells. Moreover, Morse revivals reveal concise ways to find complete revival times  $T_{rev}$  along with new ways to quantify quantum wavepacket dynamics using Ford circles [14][15].

The Morse oscillator potential Equation 1a is an anharmonic potential [16] used to describe covalent molecular bonding. Some dynamics of Morse states have been studied [17, 18, 19, 20, 21] as a model of vibrational anharmonicity.

$$V_M(x) = D(1 - e^{-\alpha x})^2 \quad (1a)$$

Coordinate  $x$  is variation of bond from equilibrium where the potential has its minimum and zero value at  $x = 0$ . Coefficient  $D$  is bond dissociation energy and its maximum inflection value at infinite  $x$ .  $D$  relates harmonic frequency  $\omega_e$  in Equation 1b and anharmonic frequency  $\omega_\chi$  in Equation 1c that gives width parameter  $\alpha$ . The latter is related to reduced mass  $\mu$  and anharmonic frequency  $\omega_\chi$ .

$$D = \frac{\omega_e^2}{4\omega_\chi} \hbar \quad (1b)$$

$$\alpha = \sqrt{\frac{2\omega_\chi\mu}{\hbar}} = \sqrt{\frac{\omega_e^2\mu}{2D}} \quad (1c)$$

McCoy [22] revived interest in exact eigenfunctions and eigenvalues [23] of Morse oscillator used in Equation 2 and Equation 3a below and allows analysis of its quantum dynamics that may be relevant to anharmonic dynamics in general.

The Morse oscillator, being anharmonic, has varying spacing of its energy levels in contrast to uniform (harmonic) spacing. At high quanta  $n$ , energy levels  $E_n = \hbar\omega_n$  have low- $n$  spacing  $\Delta E = \hbar\omega_e$  compressed for positive anharmonic frequency  $\omega_\chi$  in Equation 2.

$$E_n = \hbar\omega_n = \hbar\omega_e\left(n + \frac{1}{2}\right) - \hbar\omega_\chi\left(n + \frac{1}{2}\right)^2 \quad (2)$$

The corresponding Morse eigenfunctions of the eigenvalues are given by Equation 3a where  $L_n^{2s}$  represents a generalized associated Laguerre polynomial

[22].

$$\phi_n(x) = e^{\frac{-y(x)}{2}} y(x)^{s(n)} \sqrt{\frac{\alpha(\nu - 2n - 1)n!}{\Gamma(\nu - n)}} L_n^{2s(n)}(y(x)) \quad (3a)$$

Exponentially scaled  $y(x)$  has exponent  $s(n)$  as given.

$$y(x) = \nu e^{-\alpha x} \quad (3b)$$

$$s(n) = \frac{1}{2}(\nu - 2n - 1) \quad (3c)$$

The scaling parameter  $\nu$  is as follows.

$$\nu = \frac{4D}{\hbar\omega_e} \quad (3d)$$

Dynamic waves are combinations of eigenfunctions.

$$\psi(x, t) = \sum_{n=0}^{n_{max}} c_n \phi_n(x) e^{-i\frac{E_n t}{\hbar}} \quad (4)$$

Here  $n_{max}$  is the highest bound state. Its eigenvalue is nearest to dissociative limit  $D$ . To get maximum beating we assume equal Fourier coefficients  $c_n = 1$ . (We do not consider shorter revivals had by zeroing select  $c_n$ .)

A sample Morse oscillator potential shown in Figure 1(a) has a total of nine stationary bound states (from  $n = 0$  to  $n_{max} = 8$ ). The initial wave packet (Equation 4 at  $t = 0$ ) is a sum of these stationary bound states and evolves as shown in Figure 1(b) ending in its lowest  $\psi(x, T)^* \psi(x, T)$  trace as the initial shape fully revived.

Space-time plots of the norm  $|\psi(x, t)|$  in Figure 1(c) show resonant beat nodes and anti-nodes that outline semi-classical trajectories  $x(t)$  corresponding to energy values  $E_n$  ranging from the lowest ground state  $E_0$  up to the highest bound state  $E_{n_{max}}$ .

## 2. Analysis

An essential part of wave packet dynamics analysis of anharmonic systems is to predict if and when exact wave packet revival might occur. If  $T_{rev}$  is a time for a Morse oscillator revival, then Wang and Heller [21] have shown

$$T_{rev} = \frac{\pi}{\omega_\chi} \mathbb{M} \quad (5a)$$

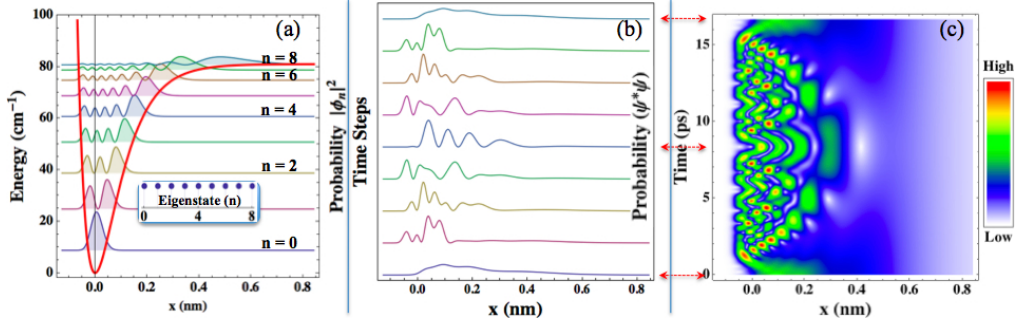


Figure 1: The Morse oscillator with a harmonic frequency  $\omega_e/2\pi c = 18(\text{cm}^{-1})$  and an anharmonic frequency  $\omega_\chi/2\pi c = 1(\text{cm}^{-1})$ . (a) Each of its stationary eigenstate  $|\phi_n|^2$  was list-plotted on a energy level of eigenvalue  $E_n$  in the potential well (red-color-line), these wave functions are normalized (indicated by the same-height dotted-line). (b) The wave packet  $\psi^*\psi$  is propagated along the time steps. (c) The probability density map of the wave packet  $|\psi|$  as a function of space and time. The double arrows connecting (b)-(c) indicate the corresponding time events.

where  $\mathbb{M}$  is an integer. This revealed two facts about Morse oscillator dynamics. First, there may be minimum or fundamental revival period at

$$T_{min-rev} = \frac{\pi}{\omega_\chi} \quad (5b)$$

This is the shortest revival time for Morse oscillator found by Wang and Heller [21]. Second, any complete revival period is made of integer numbers of the fundamental period. That is, any complete quantum revival must contain integer numbers of semiclassical-trajectory-profile periods (minimum revival period) which is approximately outlined by a classical particle oscillating with a frequency of  $2\omega_\chi$  in the Morse potential.

To illustrate relations between quantum periods and semiclassical-trajectory-profile periods, consider three cases of classical particles with corresponding quantum eigenvalue energies orbiting in a Morse potential as shown in Figure 2(a). Here the rainbow-shape trajectory of a classical particle with energy  $E_2$  has a classical oscillating period  $T$  close to the fundamental period of  $\pi/\omega_\chi$ , while a classical trajectory with energy  $E_3 = D$  is of a particle barely escaping from its Morse potential well.

The preceding case has a simple revival period formula. More analysis is required to determine a specific integer  $\mathbb{M}$  of Equation 5a for Morse revivals for a given  $(\omega_e, \omega_\chi)$ .

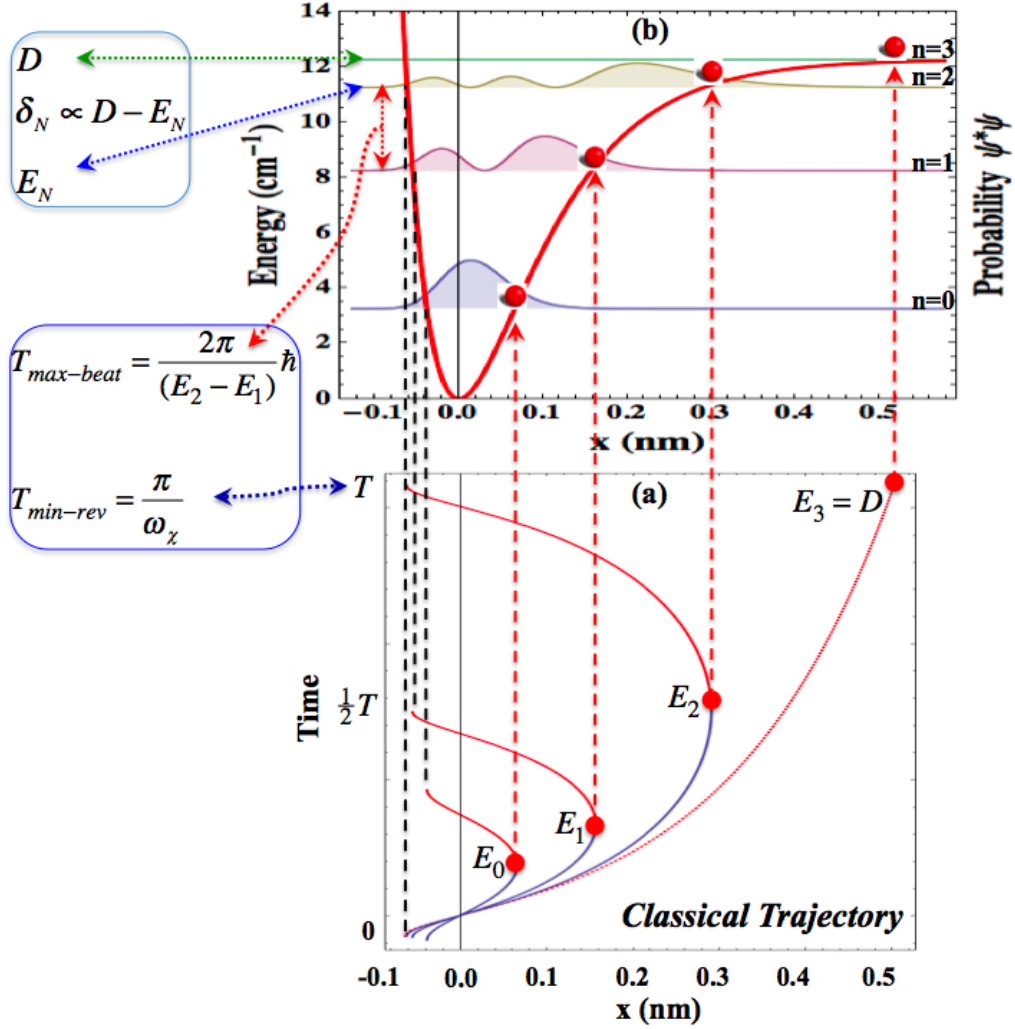


Figure 2: Relating the maximum beat period and semiclassical-trajectory-profile period (the minimum revival period). (a) 3 classical trajectories of particles oscillating in a Morse potential are plotted in one period time, and one additional classical trajectory of particle with dissociation energy  $D$  is also plotted. The red-dots in (a) and (b) indicate that these classical particles have the same energies as the corresponding quantum eigenvalue energies. (b) The probability amplitudes of 3 bound quantum eigenfunctions are listed along energy level in a Morse potential (red-thick-line).

Beating of waves with nearby frequency plays a key role in quantum dynamics. The maximum beat period  $T_{max-beat}$  due to the closest bound energy level pair in the Morse well is one key to finding its revival period. A complete revival of  $|\Psi(x, t)|^2$  at time  $T_{rev}$  must contain integer numbers of all beat periods including at least one fundamental time period  $T_{max-beat}$  for the slowest beat frequency. This relates it to revival period.

$$T_{rev} = T_{max-beat}\mathbb{N} \quad (6)$$

Here  $\mathbb{N}$  is an integer. The Morse energy level Equation 2 gives a beat-gap between neighboring energy.

$$\Delta E = E_n - E_{n-1} = \hbar(\omega_e - 2\omega_\chi n) \quad (7)$$

The  $\Delta E$  is the minimum for maximum  $n$  occurring between the highest bound quantum numbers  $n_{max}$  and  $n_{max-1}$ . Planck's relation  $E = \hbar\omega$  gives maximum beat period.

$$\begin{aligned} T_{max-beat} &= \frac{2\pi}{(\Delta\omega)_{min}} = \frac{2\pi}{E_{n_{max}} - E_{n_{max-1}}} \hbar \\ &= \frac{2\pi}{\omega_e - 2\omega_\chi n_{max}} \end{aligned} \quad (8)$$

To estimate upper bound quantum  $n_{max}$  in Equation 8, we suppose  $n_{max}$  is the integer part of a real number  $n_{real}$  and substitute  $n_{real}$  into energy Equation 2 to give  $E_{n_{real}}$  that equals dissociative limit  $D$  in Equation 1b. This equivalent relation is expressed as

$$\begin{aligned} E_{n_{real}} &= \hbar\omega_e(n_{real} + \frac{1}{2}) - \hbar\omega_\chi(n_{real} + \frac{1}{2})^2 \\ &= D = \frac{\omega_e^2}{4\omega_\chi} \hbar \end{aligned} \quad (9a)$$

A perfect square equation gives one root.

$$n_{real} = \frac{\omega_e}{2\omega_\chi} - \frac{1}{2} \quad (9b)$$

The integer part or floor of  $n_{real}$  is the highest Morse quantum number  $n_{max}$  (For Figure 1, this is  $n_{max} = 8$ ).

$$n_{max} = Floor[n_{real}] = Floor[\frac{\omega_e}{2\omega_\chi} - \frac{1}{2}] \quad (9c)$$

The fractional part  $\delta_N$  of  $n_{real}$  is quantum defect of dissociative level  $D$  and highest allowed bound energy level.

$$\delta_N = n_{real} - n_{max} \quad (9d)$$

As illustrated in Figure 2(b),  $\delta_N$  is proportional to energy gap between  $D$  and the highest bound energy level.

Then, the fundamental period  $T_{max-beat}$  in Equation 8 is expressed in term of  $\delta_N$ .

$$\begin{aligned} T_{max-beat} &= \frac{2\pi}{\omega_e - 2\omega_\chi n_{max}} = \frac{2\pi}{\omega_e - 2\omega_\chi(n_{real} - \delta_N)} \\ &= \frac{2\pi}{\omega_e - 2\omega_\chi(\frac{\omega_e}{2\omega_\chi} - \frac{1}{2} - \delta_N)} \\ &= \frac{\pi}{\omega_\chi(\delta_N + \frac{1}{2})} \end{aligned} \quad (10a)$$

Equation 10a is rewritten by substituting  $T_{min-rev} = \pi/\omega_\chi$  given by Equation 5b.

$$T_{max-beat} = \frac{\pi}{\omega_\chi(\delta_N + \frac{1}{2})} = T_{min-rev} \frac{1}{(\delta_N + \frac{1}{2})} \quad (10b)$$

This relates two fundamental building blocks of a complete Morse revival period.

$$\frac{T_{min-rev}}{T_{max-beat}} = \delta_N + \frac{1}{2} \quad (11a)$$

As discussed for Equation 5a and Equation 6, a perfect quantum revival period of the Morse oscillator  $T_{rev}$  is composed of integer numbers of the fundamental periods as follows.

$$T_{rev} = T_{min-rev}\mathbb{M} = T_{max-beat}\mathbb{N} \quad (11b)$$

Then Equation 11a and Equation 11b relate  $\mathbb{N}$  and  $\mathbb{M}$  integers.

$$\frac{\mathbb{N}}{\mathbb{M}} = \delta_N + \frac{1}{2} \quad (11c)$$

The quantum beat-period approach gives Morse revival time  $T_{rev}$  in terms of  $T_{max-beat}$  and  $\delta_N$  as follows.

$$T_{rev} = T_{max-beat}\mathbb{N} = T_{max-beat} \text{Numerator}[\delta_N + \frac{1}{2}] \quad (12a)$$

The semiclassical-trajectory-profile approach gives  $T_{rev}$  in terms of  $T_{min-rev}$  and  $\delta_N$  as follows.

$$T_{rev} = T_{min-rev}\mathbb{M} = T_{min-rev}Denominator[\delta_N + \frac{1}{2}] \quad (12b)$$

Both  $T_{min-rev}$  or  $T_{max-beat}$  serve as a fundamental building blocks of  $T_{rev}$ . Examples of this follow.

### 2.1. Fibonacci Sequence and Exchange Rate of $\frac{T_{min-rev}}{T_{max-beat}}$

Interplay of harmonicity and anharmonicity of Morse oscillators affects revival period  $T_{rev}$ . Consider Figure 3(a)-(c) where the value of  $T_{rev}$  is increased from the minimum revival period  $T_{min-rev}$  to multiples thereof with fixed anharmonic frequency  $w_\chi/2\pi c = 1(cm^{-1})$ .

In Figure 3(a) with  $\omega_e/2\pi c = 18(cm^{-1})$  one perfect revival occurs in the minimum revival time:  $T_{rev} = T_{max-beat} = T_{min-rev}$  giving a unit ratio  $T_{min-rev}/T_{max-beat} = 1/1$ .

Then in Figure 3(b) with  $\omega_e/2\pi c = 17(cm^{-1})$  and the same  $\omega_\chi$ , is seen a double time for perfect revival of  $T_{rev} = T_{max-beat} = 2T_{min-rev}$  with a half ratio  $T_{min-rev}/T_{max-beat} = 1/2$ . We note that this double revival time  $T_{rev} = 2T_{min-rev} = 2\pi/\omega_\chi$  exactly equals  $T_{approx}$  in Equation 13 given by a semiclassical treatment of general anharmonic oscillators [24, 25, 26, 27] that assumes large quantum numbers  $n$  around their average  $\bar{n}$ .

$$T_{approx} = \frac{2\pi}{\frac{1}{2} \left| \frac{d^2 E_n}{dn^2} \right|_{n=\bar{n}}} = \frac{2\pi}{w_\chi} \quad (13)$$

In Figure 3(c) with  $w_e/2\pi c = 17 + \frac{1}{3}(cm^{-1})$  is a perfect revival time  $T_{rev} = 2T_{max-beat} = 3T_{min-rev}$  with ratio of minimum revival period to maximum beat period of  $T_{min-rev}/T_{max-beat} = 2/3$ .

All revival periods  $T_{rev}$  are composed of an integer number of fundamental period  $T_{min-rev}$  (or  $T_{max-beat}$ ) but have a differing ratios  $T_{min-rev}/T_{max-beat}$  that range between 1/2 and 3/2. The Fibonacci sequence  $\{1/1, 1/2, 2/3, 3/5, 5/8, \dots\}$  is a subset of the possible rational ratios  $T_{min-rev}/T_{max-beat}$ .

### 2.2. Farey-sum and Ford geometry of fractional revivals

Fractional or intra-revival structure of Morse vibrators is quite like that of rotor revivals. In Figure 4(a) is a Morse revival of higher frequency  $\omega_e/2\pi c = 42(cm^{-1})$  and more states ( $n_{max} = 20$ ) than the one in Figure 3(a), but



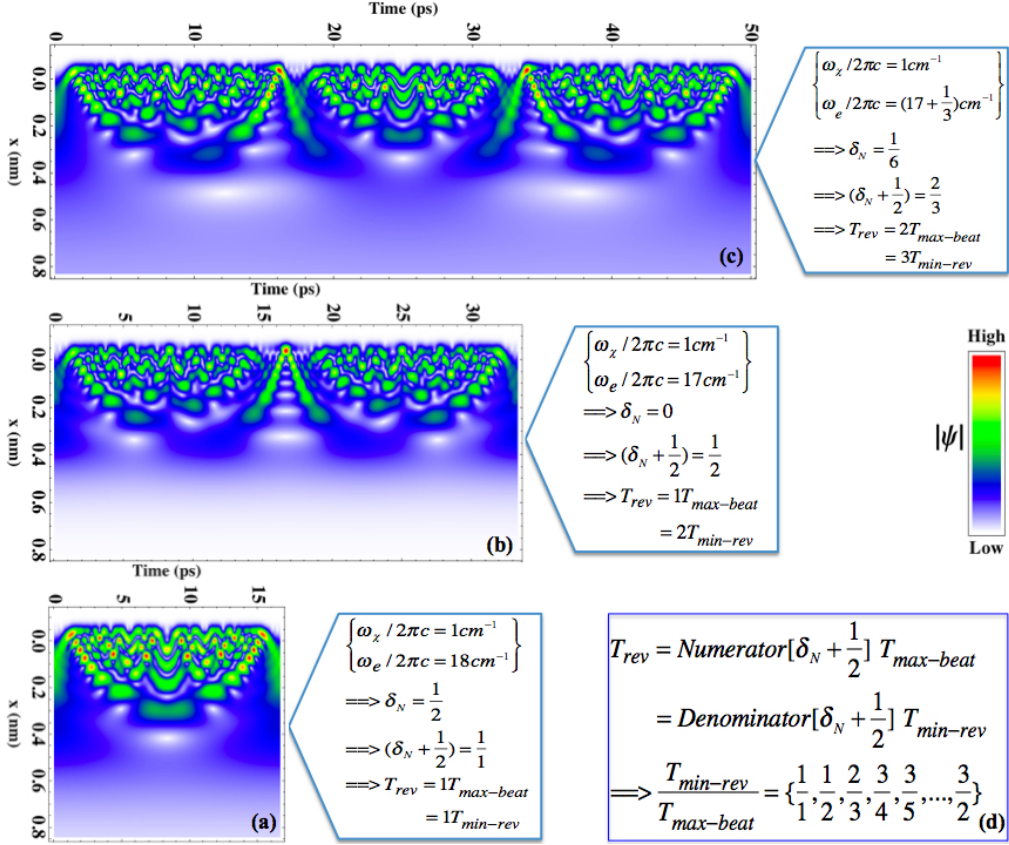


Figure 3: Fibonacci sequence and exchange rate. (a) When  $\delta_N = 1/2$ , the  $T_{rev}$  is composed of one  $T_{max-beat}$  and one  $T_{min-rev}$ . (b) When  $\delta_N = 0$ , the  $T_{rev}$  is composed of  $1T_{max-beat}$  and  $2T_{min-rev}$ . (c) When  $\delta_N = 1/6$ , the  $T_{rev}$  is composed of  $2T_{max-beat}$  and  $3T_{min-rev}$ . (d) Allowed ratios of  $T_{min-rev}$  to  $T_{max-beat}$ .

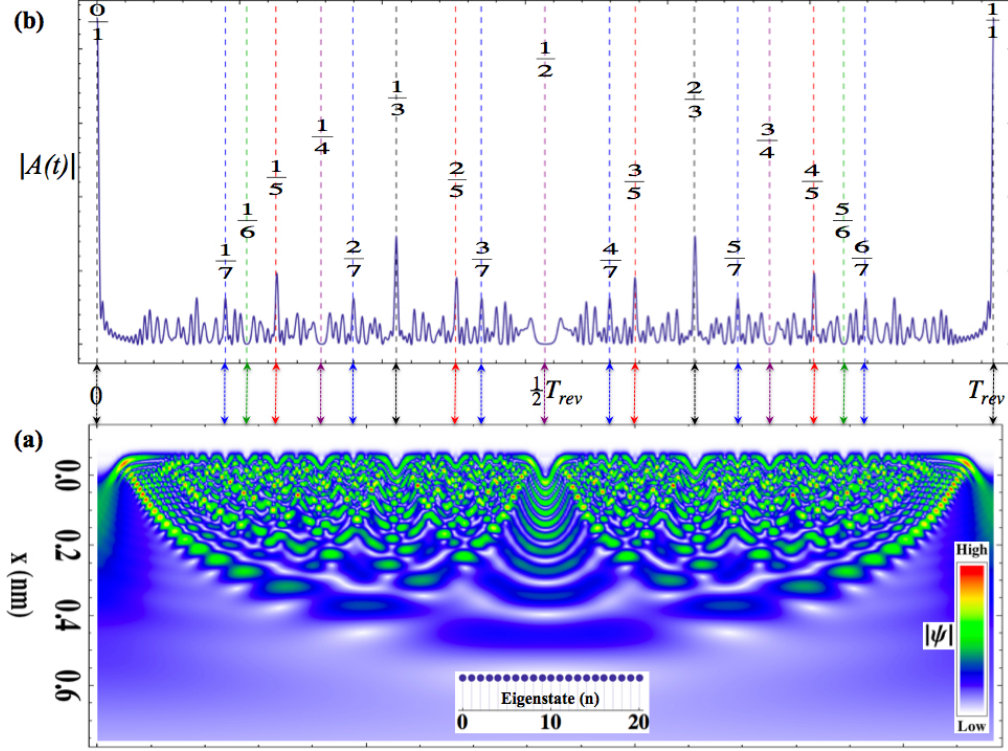


Figure 4: The Farey-sum sequence structure appears in Morse oscillator space-time pattern for  $w_e/2\pi c = 42(\text{cm}^{-1})$  and  $w_\chi/2\pi c = 1(\text{cm}^{-1})$  with  $n_{max} = 20$ . (a) One complete revival period plot of the wave packet has color denote magnitude  $|\Psi(x,t)|$ . (b) The norm of autocorrelation function ( $|A(t)|$ ) with  $n_{max} = 20$  is plotted in one complete revival period  $T_{rev}$  whose fractions  $\{\frac{0}{1}, \frac{1}{7}, \frac{1}{6}, \frac{1}{5}, \frac{1}{4}, \frac{2}{7}, \frac{1}{3}, \frac{2}{5}, \frac{3}{7}, \frac{1}{2}, \frac{4}{7}, \frac{3}{5}, \frac{2}{3}, \frac{5}{7}, \frac{3}{4}, \frac{4}{5}, \frac{5}{6}, \frac{6}{7}, \frac{1}{1}\}$  are denoted by the vertical dashed lines. The double arrows connecting (a)-(b) indicate the corresponding time events having peaks (or nodes) for time fraction  $\frac{n}{d}$  of odd (or even) depth  $d$ .

with the same anharmonicity  $\omega_\chi/2\pi c = 1(cm^{-1})$  and revival period:  $T_{rev} = 1T_{min-rev} = 1T_{max-beat} = 1/(2c(cm)^{-1}) \approx 16.7(\text{pico} - \text{second})$ .

Fractional revival structure is visible as a series of dips on top of Figure 4(a) and in Fourier amplitude frequency  $\omega_n$  sum or autocorrelation  $A(t)$  spectra [25] in Figure 4(b).

$$A(t) = \sum_{n=0}^{n_{max}} e^{-i\frac{E_n t}{\hbar}} = \sum_{n=0}^{n_{max}} e^{-i\omega_n t} \quad (14)$$

$|A(t)|$  spectra match Farey-sum sequence used in 1815 by geologist John Farey [9][28] to analyze tidal beats.

A Farey sequence, starting with fraction 0/1 and ending with fraction 1/1, builds hierarchies of irreducible rational fractions on a real line between 0.0 and 1.0 [28]. A Farey-sum  $\frac{a}{b} + \frac{c}{d} = \frac{a+c}{b+d}$  is a curious process to locate significant fractions  $\frac{n}{d}$  between  $\frac{a}{b}$  and  $\frac{c}{d}$  or overtone ( $n : d$ ) resonances in between an ( $a : b$ ) and a ( $c : d$ ) resonance.

In 1938, Leslie Ford [14] found a geometric description that helps elucidate Farey-sums. Ford geometry views each fraction as a vector and the Farey sum as a vector sum in Denominator( $y$ )-vs-Numerator( $x$ ) space such as is plotted for  $0 \leq (x, y) \leq 1$  in Figure 5. A fraction  $\frac{a}{b}$  is drawn as a vector with tail at origin and head at the point ( $x = a, y = b$ ) as shown by examples  $\mathbf{V}_{\frac{0}{1}}$  and  $\mathbf{V}_{\frac{1}{1}}$ , the black and red arrows in Figure 5 (lower left).

$$\mathbf{V}_{\frac{0}{1}} + \mathbf{V}_{\frac{1}{1}} = \begin{pmatrix} 0 \\ 1 \end{pmatrix} + \begin{pmatrix} 1 \\ 1 \end{pmatrix} = \begin{pmatrix} 1 \\ 2 \end{pmatrix} = \mathbf{V}_{\frac{1}{2}} \quad (15a)$$

The green arrow in Figure 5 is  $\mathbf{V}_{\frac{1}{2}}$ . Each vector  $\mathbf{V}_{\frac{n}{d}}$  points to or intersects a real value  $x = n/d$  on the top ( $y = 1$ )-line of Figure 5. Ford [14] discovered that each  $x = n/d$  is a tangent point for a circle having diameter  $1/d^2$  hanging below the top ( $y = 1$ )-line that is itself tangent to infinite sequences of smaller such circles, each tangent to the next and converging on  $x = n/d$ . The ( $d = 1$ )-Ford-circle is a unit-diameter circle sliced to fit the unit ( $x, y$ )-area with a pair of tangent semi circles belonging to unit Ford base vectors  $\mathbf{V}_{\frac{0}{1}}$  and  $\mathbf{V}_{\frac{1}{1}}$ . Fractions  $\frac{0}{1}$  and  $\frac{1}{1}$  make a second ( $d=2$ )-Ford circle of diameter  $1/2^2$  in the upper center of Figure 5 pointed out by sum  $\mathbf{V}_{\frac{1}{2}} = \mathbf{V}_{\frac{0}{1}} + \mathbf{V}_{\frac{1}{1}}$  in Equation 15a.

It is tangent to “parent” Ford circles for  $\frac{0}{1}$  and  $\frac{1}{1}$ . Also shown is ( $d=3$ )-Ford circle for vector  $\mathbf{V}_{\frac{1}{3}} = \mathbf{V}_{\frac{0}{1}} + \mathbf{V}_{\frac{1}{2}}$  that is tangent to circles of its parent

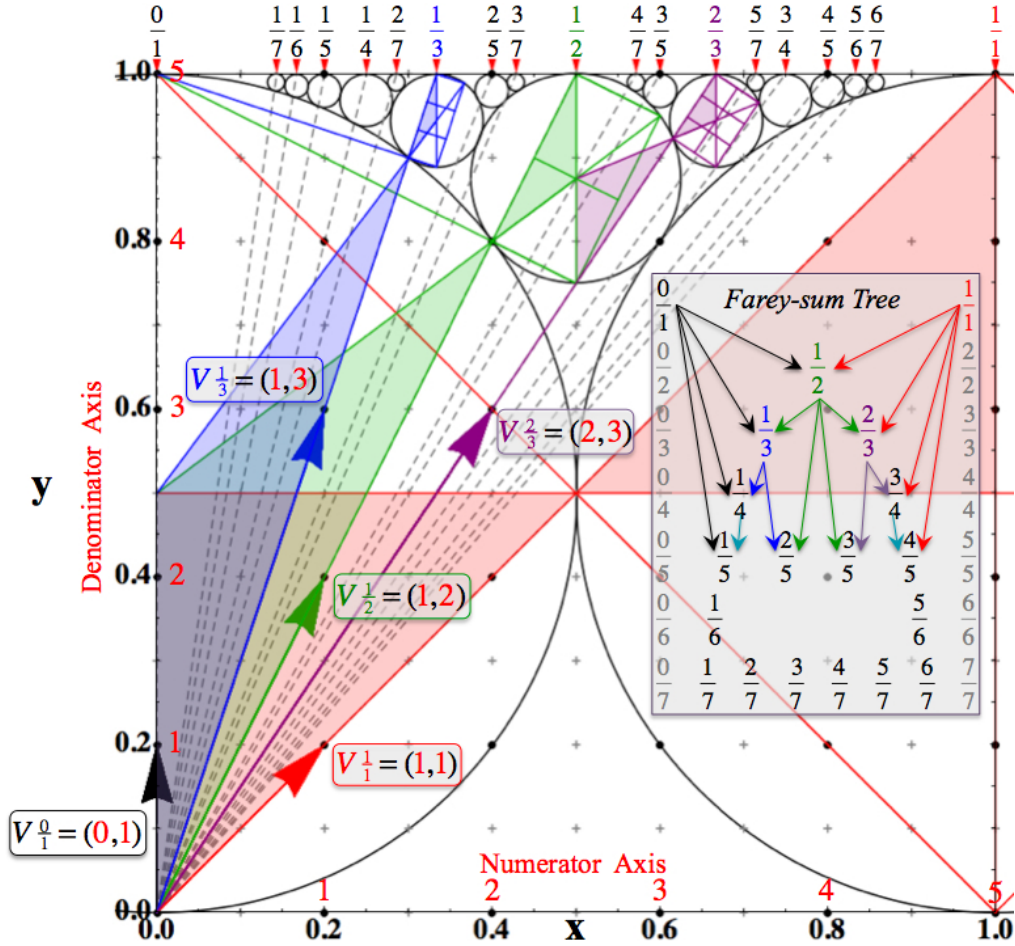


Figure 5: Ford circles and vectors with Farey-sum sequence.

fractions  $\frac{0}{1}$  and  $\frac{1}{2}$ .

$$\mathbf{V}_{\frac{0}{1}} + \mathbf{V}_{\frac{1}{2}} = \begin{pmatrix} 0 \\ 1 \end{pmatrix} + \begin{pmatrix} 1 \\ 2 \end{pmatrix} = \begin{pmatrix} 1 \\ 3 \end{pmatrix} = \mathbf{V}_{\frac{1}{3}} \quad (15b)$$

Thales 600-BCE rectangle-in-circle geometry is sufficient to derive Ford geometry. Tangent Ford circles like the  $\frac{0}{1}$ ,  $\frac{1}{1}$ , and  $\frac{1}{2}$  circles in Figure 5 meet at corners of similar Thales rectangles whose vertical diagonals are circle diameters hanging below their respective fraction points. Circle diameters subtend  $90^\circ$  corners by Thales theorem. The  $\frac{0}{1}$  and  $\frac{1}{2}$  corners meet where the  $\mathbf{V}_{\frac{1}{2}}$ -vector line crosses the  $\frac{0}{1}$ -circle. This is the  $\frac{0}{1}$ - $\frac{1}{2}$ -circle-tangent point.

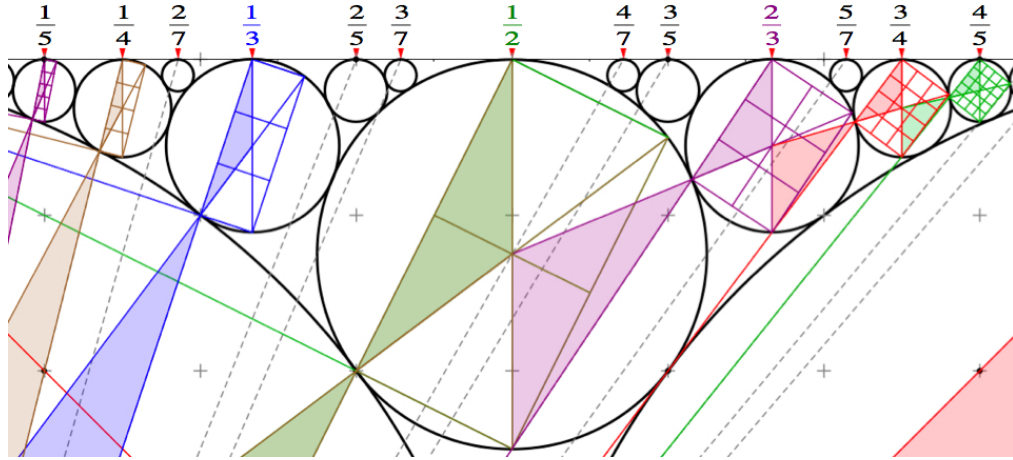


Figure 6: “Quantization” inside Ford circles: pixel lattices of  $\{(1 \times 5), (1 \times 4), (1 \times 3), (1 \times 2), (2 \times 3), (3 \times 4), (4 \times 5)\}$  rectangles lie circumscribed by circles of fractions  $\{\frac{1}{5}, \frac{1}{4}, \frac{1}{3}, \frac{1}{2}, \frac{2}{3}, \frac{3}{4}, \frac{4}{5}\}$ .

A  $\frac{0}{1}$ -diameter line through that point intersects the vertical diameter of the  $\frac{1}{2}$ -circle at its center thus defining it. Similar geometry (not drawn) applies to the  $\frac{1}{2}$ - $\frac{1}{1}$ -circle-tangent. A Farey-sum of a circle-tangent pair is a new Ford circle and fraction as shown by examples  $\frac{0}{1} + \frac{1}{2} = \frac{1}{3}$  in Equation 15b or  $\frac{1}{2} + \frac{1}{1} = \frac{2}{3}$  listed in level-3 of the Farey-sum-tree on the righthand side of Figure 5. Farey sums that give reducible fractions  $\frac{N}{D} = \frac{n \cdot f}{d \cdot j}$  are labeled by their reduced form  $\frac{n}{d}$  with the shortest allowed Ford vector, least depth or denominator  $d$ , and largest possible Ford circle.

Continued Farey-sums of Ford vectors give sequences of circles each belonging to an irreducible fraction  $\frac{n}{d}$  and its vector  $\mathbf{V}_{\frac{n}{d}}$ . In Figure 5, these circles nest in the area between their original Farey “grandparent” circles  $\mathbf{V}_{\frac{0}{1}}$  and  $\mathbf{V}_{\frac{1}{1}}$ .

A Farey-sum-tree of fractions of depth  $d \leq 7$   $\{\frac{0}{1}, \frac{1}{7}, \frac{1}{6}, \frac{1}{5}, \frac{1}{4}, \frac{2}{7}, \frac{1}{3}, \frac{2}{5}, \frac{3}{7}, \frac{1}{2}, \frac{4}{7}, \frac{3}{5}, \frac{2}{3}, \frac{5}{7}, \frac{3}{4}, \frac{4}{5}, \frac{6}{7}, \frac{1}{1}\}$  is shown in box on the right of Figure 5 and represented by a total of 19 sequentially and mutually tangent circles.

A revealing portrait emerges of quantum “fractal” structure filling the area below the top line with ever tinier  $\frac{1}{d^2}$ -diameter circles as spectral depth  $d$  increases.

By construction all Ford-vector and Thales-rectangle slopes are rational, but surprisingly so are their aspect ratios “quantized” into  $n$ -by- $d$  pixel arrays. For example,  $(1 \times 3), (1 \times 2)$ , and  $(2 \times 3)$  pixel arrays lie inside  $\{\frac{1}{3}, \frac{1}{2}, \frac{2}{3}\}$  circles in Figure 5, and Figure 6 shows finer  $(n \times d)$  arrays of pixel rectangles

circumscribed by  $\frac{n}{d}$ -circles.

### 3. Conclusion

In conclusion, exact Morse oscillator eigensolutions allow more detailed analysis of their quantum dynamics. A key top-level-to-dissociation gap parameter  $\delta_N$  provides a concise revival time formula in terms of two fundamental periods, a semiclassical  $T_{min-rev}$  found by Wang and Heller and a longest quantum beat period  $T_{max-beat}$ . This shows that complete revival periods may be composed of integer numbers of the two. Finally, fractional revivals seen in rotor pulse evolution is also shown to be present in Morse wave dynamics in the form of Farey-sum spectral sub-structure. A Ford-circles geometry relating rational fractions to real numbers may be developed to visualize these phenomena and may eventually have application to quantum information processing and computing.

### References

- [1] J. H. Eberly, N. B. Narozhny, J. J. Sanchez-Mondragon, Periodic spontaneous collapse and revival in a simple quantum model, Phys. Rev. Lett. 44 (1980) 1323.
- [2] E. J. Heller, Time dependent approach to semiclassical dynamics, J. Chem. Phys. 62 (1975) 1544.
- [3] R. S. McDowell, C. W. Patterson, W. G. Harter, The modern revolution in infrared spectroscopy, Los Alamos Science 3 (1982) 38.
- [4] A. H. Zewail, Laser femtochemistry, Science 242 (1988) 1645.
- [5] B. Feuerstein, U. Thumm, Mapping of coherent and decohering nuclear wave-packet dynamics in d2 with ultrashort laser pulses, Phys. Rev. A 67 (063408).
- [6] A. Rudenko, T. Ergler, B. Feuerstein, K. Zrost, C. D. Schroter, R. Moshhammer, J. Ullrich, Real-time observation of vibrational revival in the fastest molecular system, Chem. Phys. 329 (2006) 193–202.
- [7] T. Niederhausen, U. Thumm, Controlled vibrational quenching of nuclear wave packets in d2, Phys. Rev. A 77 (013407).

- [8] K. Ohmori, Wave-packet and coherent control dynamics, *Annu. Rev. Phys. Chem.* 60 (2009) 487–511.
- [9] J. Farey, *Philos. Mag.* 47 (1816) 385.
- [10] W. G. Harter, Quantum-fractal revival structure in  $cn$  quadratic spectra: Base- $n$  quantum computer registers, *Phys. Rev. A* 64 (012312).
- [11] W. G. Harter, Wave node dynamics and revival symmetry in quantum rotors, *J. Mol. Spectrosc.* 210 (2001) 166.
- [12] H. Mack, M. Bienert, F. Haug, M. Freyberger, W. Schleich, Wave packets can factorize numbers, *phys. stat. sol. (b)* 233 (3) (2002) 408–415.
- [13] M. Gilowski, T. Wendrich, T. Müller, C. Jentsch, W. Ertmer, E. M. Rasel, W. P. Schleich, Gauss sum factorization with cold atoms, *Phys. Rev. Lett.* 100 (030201).
- [14] L. R. Ford, Fractions, *The American Mathematical Monthly* 45 (9) (1938) 586–601.
- [15] A. Z. Li, Quantum resonant beats and revivals in the morse oscillators and rotors, Ph.D. thesis, University of Arkansas (2013).
- [16] P. M. Morse, Diatomic molecules according to the wave mechanics. ii. vibrational levels, *Phys. Rev.* 34 (1929) 57.
- [17] V. P. Gutschick, M. M. Nieto, Coherent states for general potentials. v. time evolution, *Phys. Rev. D* 22 (1980) 403.
- [18] J. P. Dahl, M. Springborg, The morse oscillator in position space, momentum space, and phase space, *J. Chem. Phys.* 88 (1988) 4535.
- [19] S. Kais, R. D. Levine, Coherent states for the morse oscillator, *Phys. Rev. A* 41 (1990) 2301.
- [20] M. Angelova, V. Hussin, Generalized and gaussian coherent states for the morse potential, *J. Phys. A: Math. Theor.* 41 (304016).
- [21] Z. Wang, E. J. Heller, Semiclassical investigation of the revival phenomena in a one-dimensional system, *J. Phys. A: Math. Theor.* 42 (285304).

- [22] A. B. McCoy, Curious properties of the morse oscillator, *Chem. Phys. Lett.* 501 (2011) 603–607.
- [23] S. Dong, R. Lemus, A. Frank, Ladder operators for the morse potential, *Int. J. Quantum Chem.* 86 (2002) 433.
- [24] I. S. Averbukh, N. F. Perelman, Fractional revivals universality in the long-term evolution of quantum wave packets beyond the correspondence principle dynamics, *Phys. Lett. A* 139 (1989) 449.
- [25] R. W. Robinett, Quantum wave packet revivals, *Phys. Rep.* 392 (2004) 1–119.
- [26] T. Lohmuller, V. Engel, J. Beswick, C. Meier, Fractional revivals in the rovibrational motion of  $i_2$ , *J. Chem. Phys.* 120 (22) (2004) 10442.
- [27] V. V. Eryomin, S. I. Vetchinkin, I. M. Umanskii, Manifestations of wave packet fractional revivals in a morselike anharmonic system, *J. Chem. Phys.* 101 (1994) 10730.
- [28] G. H. Hardy, E. M. Wright, *An Introduction to the Theory of Numbers*, 5th Edition, Oxford University Press, New York, 1979.