

Lecture 31 *Relativity*-Dynamics

Tuesday 5.03.2016

Relativity: Quantizing wave variables of phase and amplitude (Unit 3 p.45-64)

- ➔ Review of wave parameters used to develop relativistic quantum theory
 - Bohr-Schrodinger (BS) approximation throws out Mc^2 (*Is frequency really relative?*)
 - Effect on group velocity (*None*) and phase velocity (*Absurd*)
- 1st Quantization: Quantizing phase variables k_m and $\omega(k_m)$
 - Understanding how quantum dynamics and transitions involve “mixed” states
 - Square well example of mixing unequal frequencies
 - Circle well or ring example of mixing equal or unequal frequencies
- Mixing unequal amplitudes makes “Gallop” wave: Analogy of (*SWR, SWQ*) to (V_{group}, V_{phase})
 - Analogy with optical polarization geometry and Kepler orbits
 - Super-luminal speed and Feynman-Wheeler pair-creation switchbacks
- 2nd Quantization: Quantizing wave amplitudes A_N and invariance of photon number
 - Analogy 1: Many CW (Continuous Waves) add up to make PW (Pulse Waves)
 - Analogy 2: Many Photon-Number-Modes add up to make Coherent-Laser-Modes
 - Heisenberg $\Delta v \cdot \Delta t \sim 1 \sim \Delta \kappa \cdot \Delta x$ analogous to $\Delta N \cdot \Delta phase \sim 1$ uncertainty relations
- Electromagnetic wave mode energy: Maxwell vs. Planck-Einstein
 - 1st quantization for wave phase variables and classical energy of **E**, **B**, and **A** fields
 - 2nd quantization for wave and Planck quantum energy of **E**, **B**, and **A** fields
 - Scaling **E**-waves to mime quantum Ψ -waves and ψ -waves
- Relativistic effects on charge, current, and Maxwell Fields



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Electromagnetic wave mode energy: Maxwell vs. Planck-Einstein

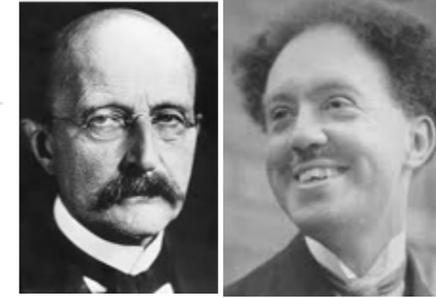
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Relativistic effects on charge, current, and Maxwell Fields

Using (some) wave parameters to develop relativistic quantum theory



Max Planck 1858-1947 Louis DeBroglie 1892-1987

$B = v_A$
 $B = v_A = cK_A$

$v_{phase} = B \cosh \rho \approx B + \frac{1}{2} B \rho^2$ (for $u \ll c$)

$cK_{phase} = B \sinh \rho \approx B \rho$ (for $u \ll c$)

$\frac{u}{c} = \tanh \rho \approx \rho$ (for $u \ll c$)

$\cosh \rho \approx 1 + \frac{1}{2} \rho^2 \approx 1 + \frac{1}{2} \frac{u^2}{c^2}$
 $\sinh \rho \approx \rho \approx \frac{u}{c}$

Lect. 30
 p. 3 to 29

At low speeds:

$v_{phase} \approx B + \frac{1}{2} \frac{B}{c^2} u^2$

\Leftarrow for $(u \ll c) \Rightarrow$

$K_{phase} \approx \frac{B}{c^2} u$

v_{phase} and K_{phase} resemble formulae for Newton's kinetic energy $\frac{1}{2} Mu^2$ and momentum Mu .

Rescale v_{phase} by h so: $M = \frac{hB}{c^2}$

or: $hB = Mc^2$ (The famous Mc^2 shows up here!)

$h v_{phase} \approx hB + \frac{1}{2} \frac{hB}{c^2} u^2$

\Leftarrow for $(u \ll c) \Rightarrow$

$h K_{phase} \approx \frac{hB}{c^2} u$

So attach scale factor h (or hN) to match units.

$h v_{phase} \approx Mc^2 + \frac{1}{2} Mu^2$

\Leftarrow for $(u \ll c) \Rightarrow$

$h K_{phase} \approx Mu$

Natural wave conspiracy? Lucky coincidences? Expensive Cheap trick? ...Try exact v_{phase} and K_{phase} ...

Energy density, photon number N , and normalization discussed p.74-81

Need to replace h with hN to match e.m. energy density $\epsilon_0 \mathbf{E} \cdot \mathbf{E} = hN v_{phase}$

$h v_{phase} = hB \cosh \rho = Mc^2 \cosh \rho$
 Planck (1900)
 Total Energy: $E = \frac{Mc^2}{\sqrt{1-u^2/c^2}}$
 Einstein (1905)

This motivates the "particle" normalization $\int \Psi^* \Psi dV = N \quad \Psi = \sqrt{\frac{\epsilon_0}{h v}} E$

$h c K_{phase} = hB \sinh \rho = Mc^2 \sinh \rho$

$\frac{1}{\sqrt{\beta^2 - 1}} = \frac{c}{\sqrt{1 - \frac{u^2}{c^2}}}$ (old-fashioned notation)

$cp = \frac{Mu c}{\sqrt{1 - u^2/c^2}}$

Momentum: $h c K_{phase} = p = \frac{Mu}{\sqrt{1 - u^2/c^2}}$
 DeBroglie (1921)

group	$b_{Doppler RED}$	$\frac{v_{group}}{c}$	$\frac{v_{group}}{v_A}$	$\frac{\lambda_{group}}{\lambda_A}$	$\frac{K_{group}}{K_A}$	$\frac{\tau_{group}}{\tau_A}$	$\frac{v_{phase}}{c}$
phase	$\frac{1}{b_{Doppler BLUE}}$	$\frac{c}{v_{phase}}$	$\frac{K_{phase}}{K_A}$	$\frac{\tau_{phase}}{\tau_A}$	$\frac{v_{phase}}{v_A}$		
rapidity ρ	$e^{-\rho}$	$\tanh \rho$	$\sinh \rho$	$\text{sech } \rho$	$\cosh \rho$		
stellar angle σ	$1/e^{+\rho}$	$\sin \sigma$	$\tan \sigma$	$\cos \sigma$	$\sec \sigma$	$\cot \sigma$	$\csc \sigma$
$\beta \equiv \frac{u}{c}$	$\sqrt{\frac{1-\beta}{1+\beta}}$	$\frac{\beta}{1}$	$\frac{1}{\sqrt{\beta^2-1}}$	$\frac{\sqrt{1-\beta^2}}{1}$	$\frac{1}{\sqrt{1-\beta^2}}$	$\frac{\sqrt{\beta^2-1}}{1}$	$\frac{1}{\beta}$
value for $\beta=3/5$	$\frac{1}{2} = 0.5$	$\frac{3}{5} = 0.6$	$\frac{3}{4} = 0.75$	$\frac{4}{5} = 0.80$	$\frac{5}{4} = 1.25$	$\frac{4}{3} = 1.33$	$\frac{5}{3} = 1.67$

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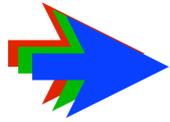
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$$E = \hbar\omega \quad \text{ENERGY}$$

$$= \frac{Mc^2}{\sqrt{1-u^2/c^2}} = Mc^2 \cosh \rho$$

given:

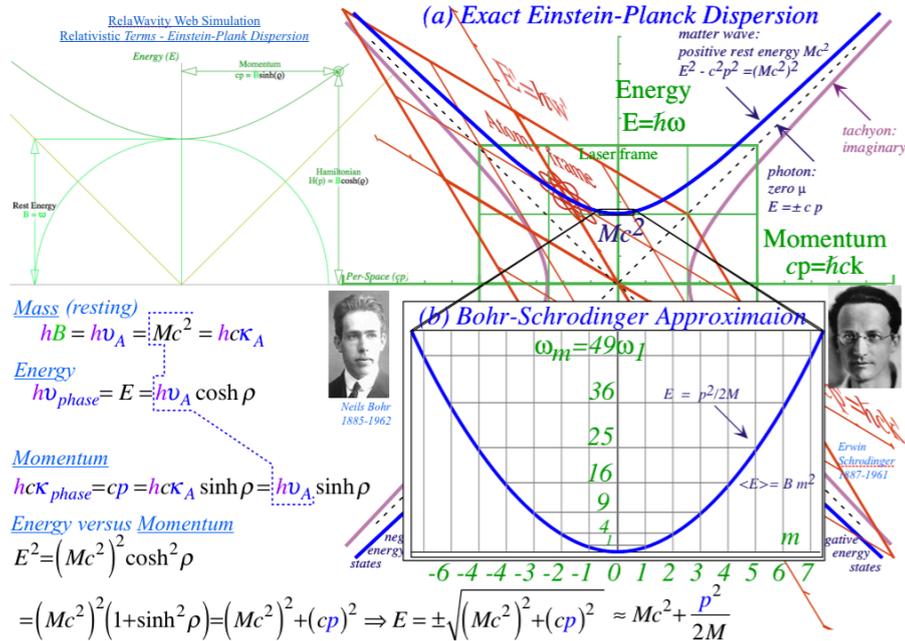
$$cp = \hbar ck \quad \text{MOMENTUM}$$

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[Relativistic Terms - Einstein-Planck Dispersion](#)

Some details concerning Lect. 30 - slide 31

Using (some) wave coordinates for relativistic quantum theory



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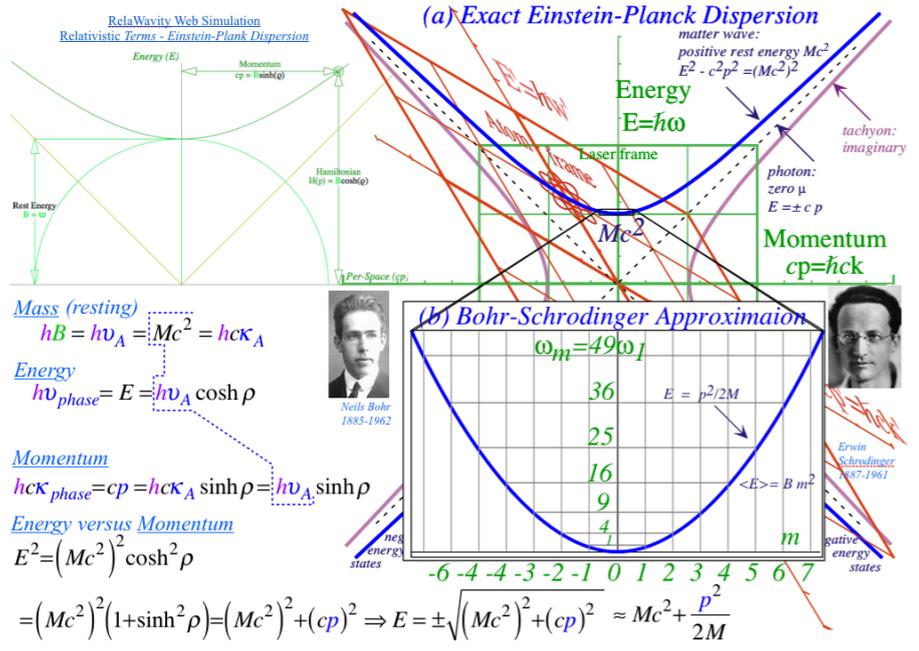
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$$E = \left[(Mc^2)^2 + (cp)^2 \right]^{1/2} \approx \cancel{Mc^2} + \frac{1}{2M} p^2 \xrightarrow{BS\text{-}approx} \frac{1}{2M} p^2$$

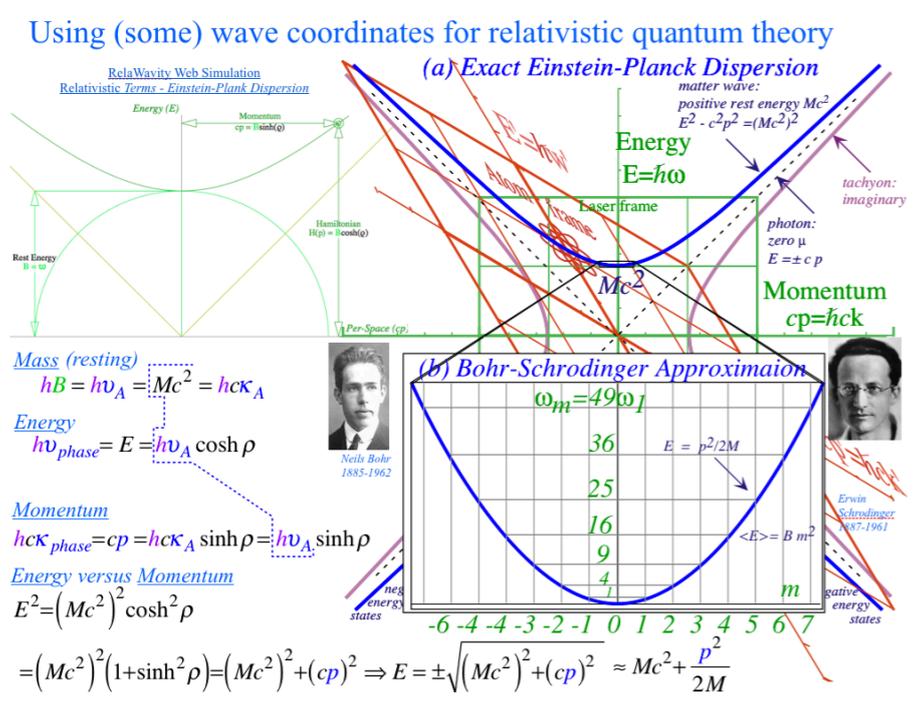
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BS- binomial approximation

$$(a+b)^n = a^n + na^{n-1}b + \dots$$

replaces hyperbola: $E = \left[(Mc^2)^2 + (cp)^2 \right]^{1/2} = Mc^2 + \frac{1}{2} \frac{c^2 p^2}{Mc^2} + \dots$

with parabola: $E = Mc^2 + \frac{1}{2M} p^2$



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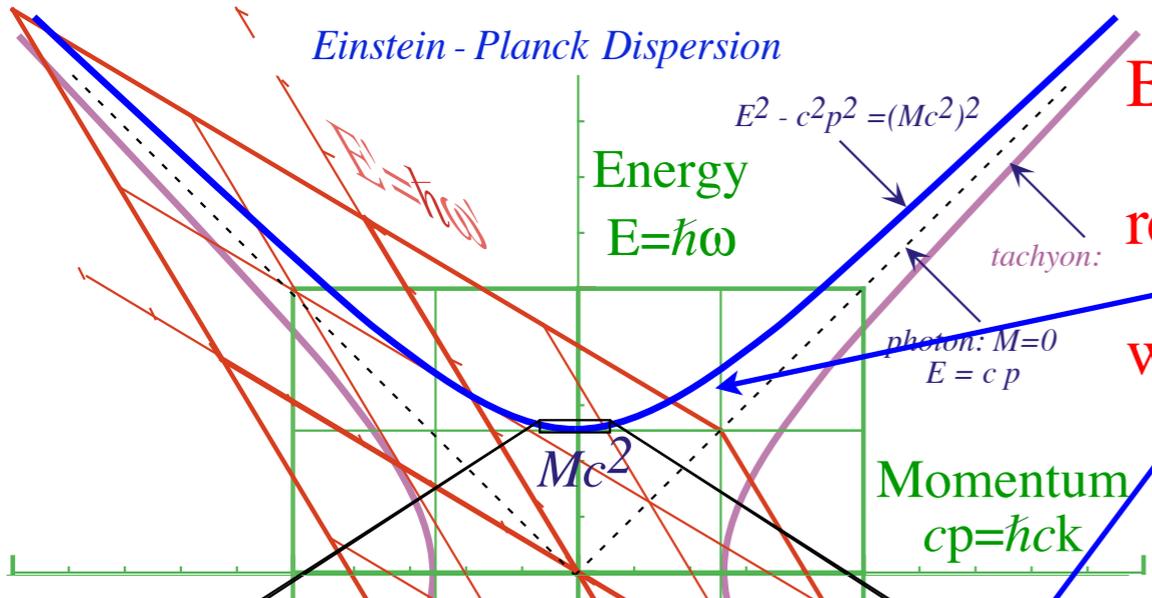
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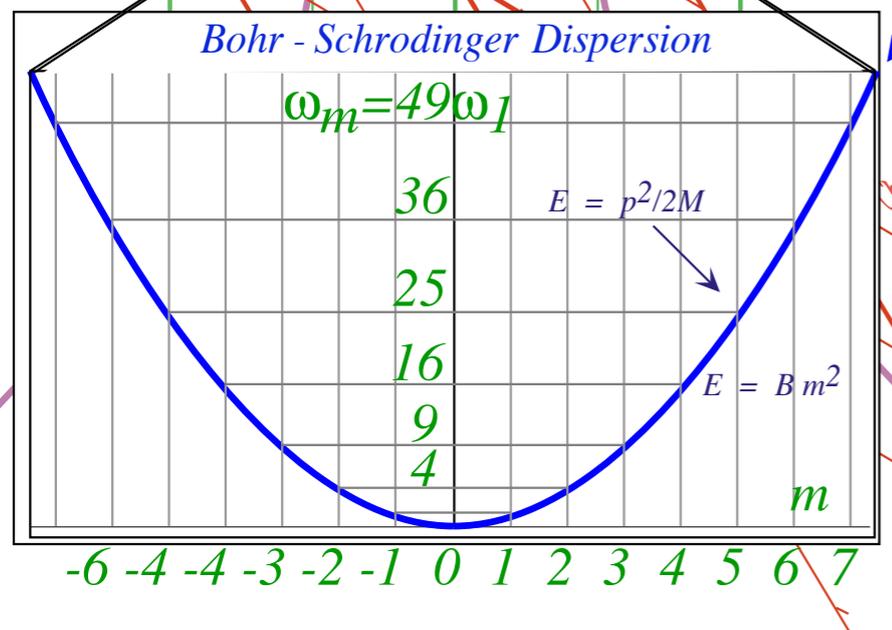
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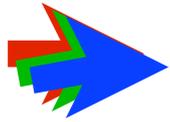
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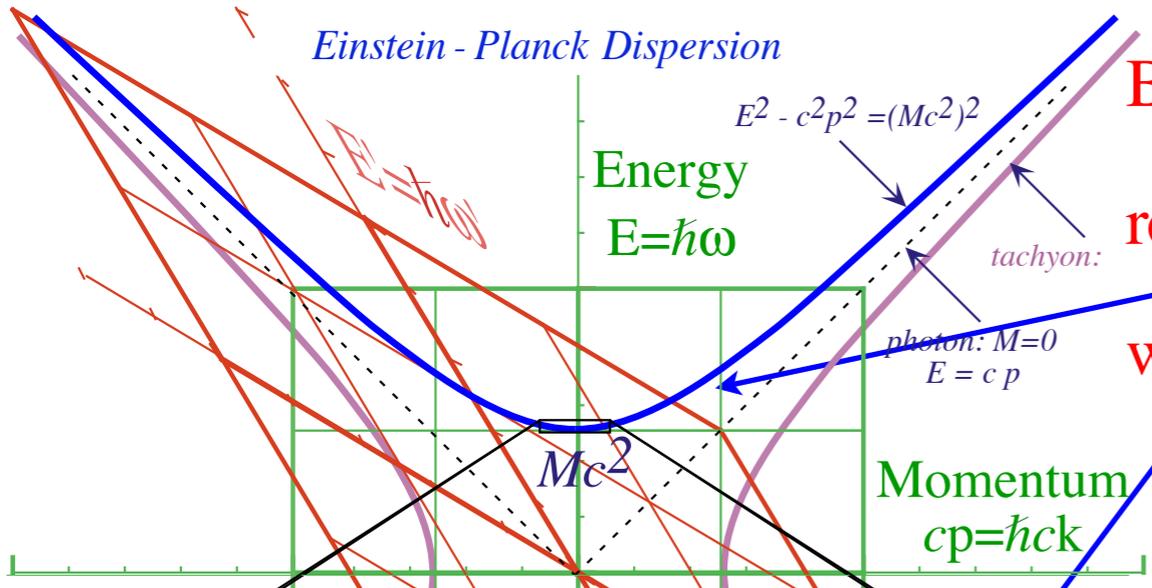
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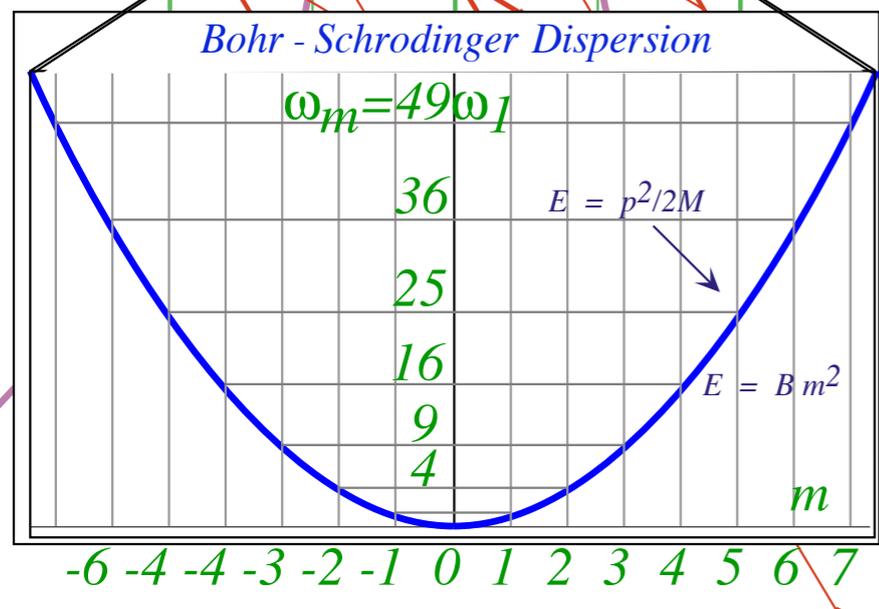
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This affects wave phase dynamics only:



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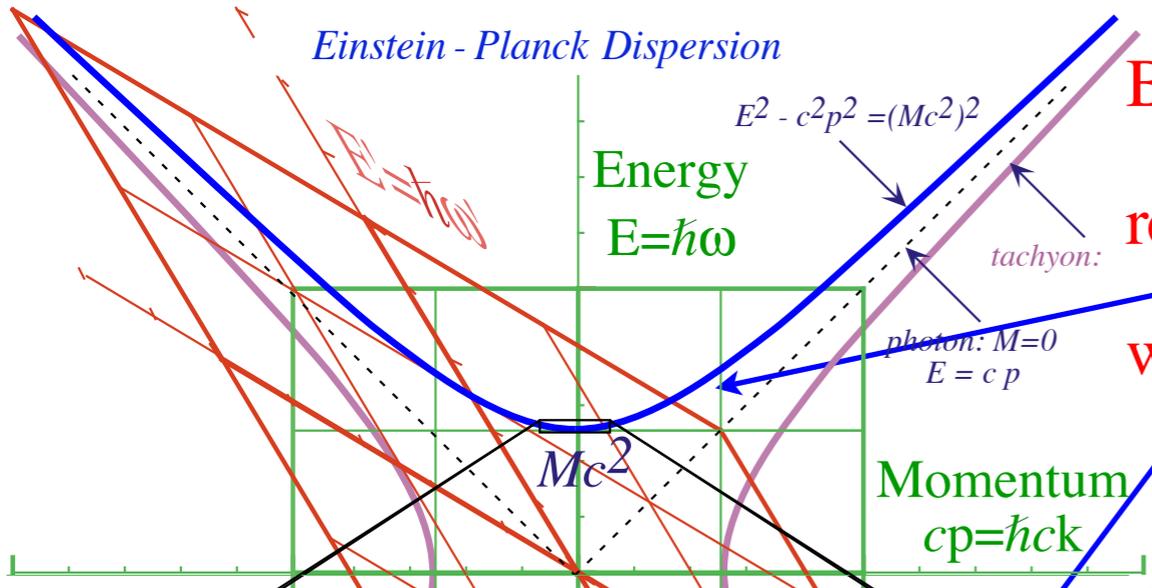
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replaces hyperbola:

$$E = \left[(Mc^2)^2 + (cp)^2 \right]^{1/2} = Mc^2 + \frac{1}{2} \frac{c^2 p^2}{Mc^2} + \dots$$

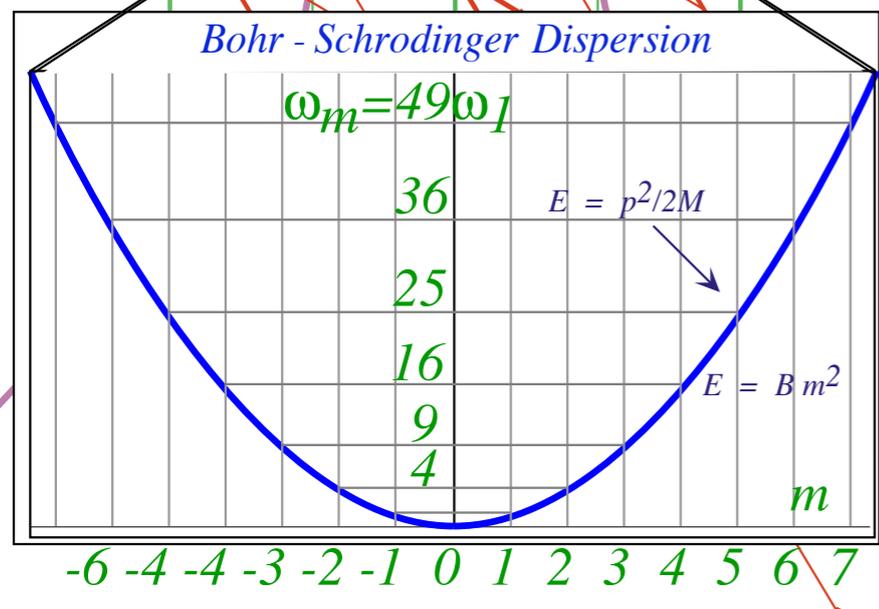
with parabola:

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This affects wave phase dynamics only:

Group velocity $u = V_{group} = \frac{d\omega}{dk}$ is a differential quantity unaffected by origin shift.

But, Phase velocity $\frac{\omega}{k} = V_{phase}$ is greatly reduced by deleting Mc^2 from $E = \hbar\omega$.



Bohr-Schrodinger (BS) approximation throws out Mc^2 (Is frequency really relative?)

$E = \hbar\omega$ ENERGY

given:

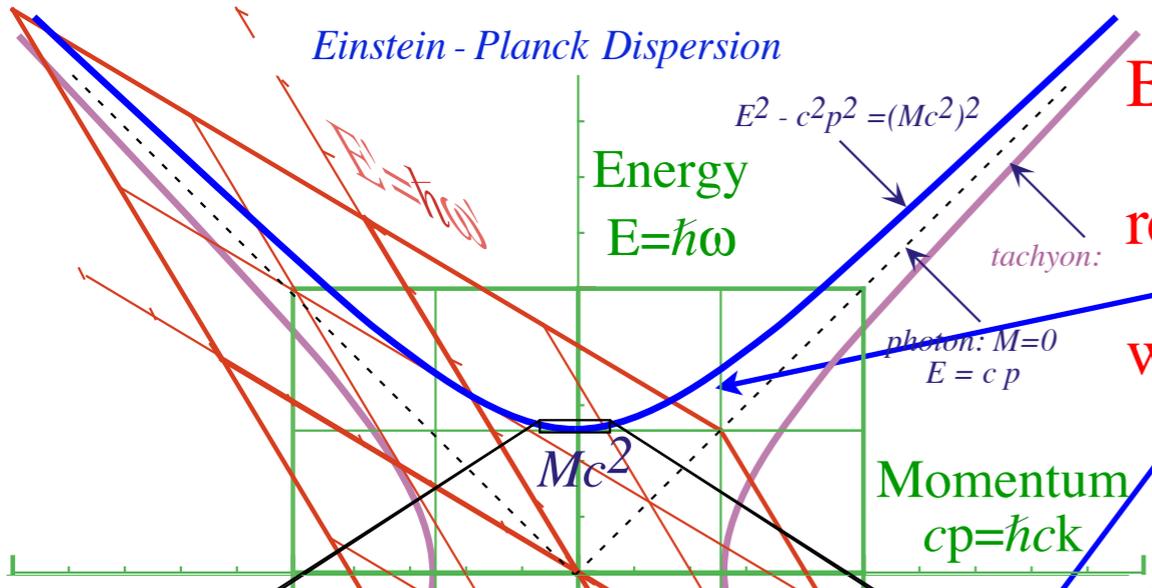
$cp = \hbar ck$ MOMENTUM
 $= \frac{Mc u}{\sqrt{1 - u^2/c^2}} = Mc^2 \sinh \rho$

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[RelaWavity Web Simulation](#)
[Relativistic Terms - Einstein-Planck Dispersion](#)

$$E = \left[(Mc^2)^2 + (cp)^2 \right]^{1/2} \approx \cancel{Mc^2} + \frac{1}{2M} p^2 \xrightarrow{\text{BS-approx}} \frac{1}{2M} p^2$$

The **BS claim**: may shift energy origin ($E=Mc^2, cp=0$) to ($E=0, cp=0$). (Frequency is relative!)



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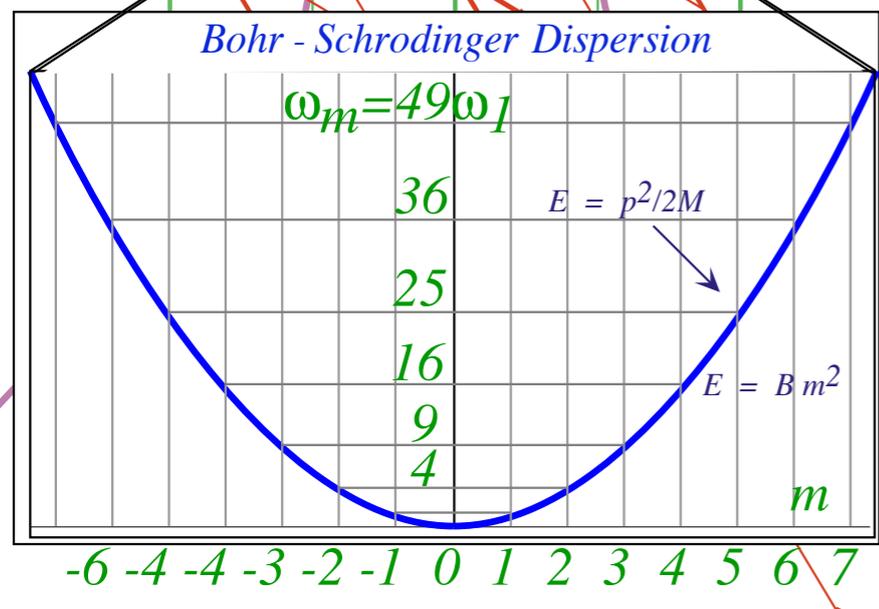
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It slows from super-luminal $V_{phase} = c^2/u$ to a sedate sub-luminal speed of $V_{group}/2$.

$$\omega_{BS}(k) = \frac{k^2}{2M} \quad \text{gives:} \quad V_{phase} = \frac{\omega_{BS}}{k} = \frac{k}{2M}$$

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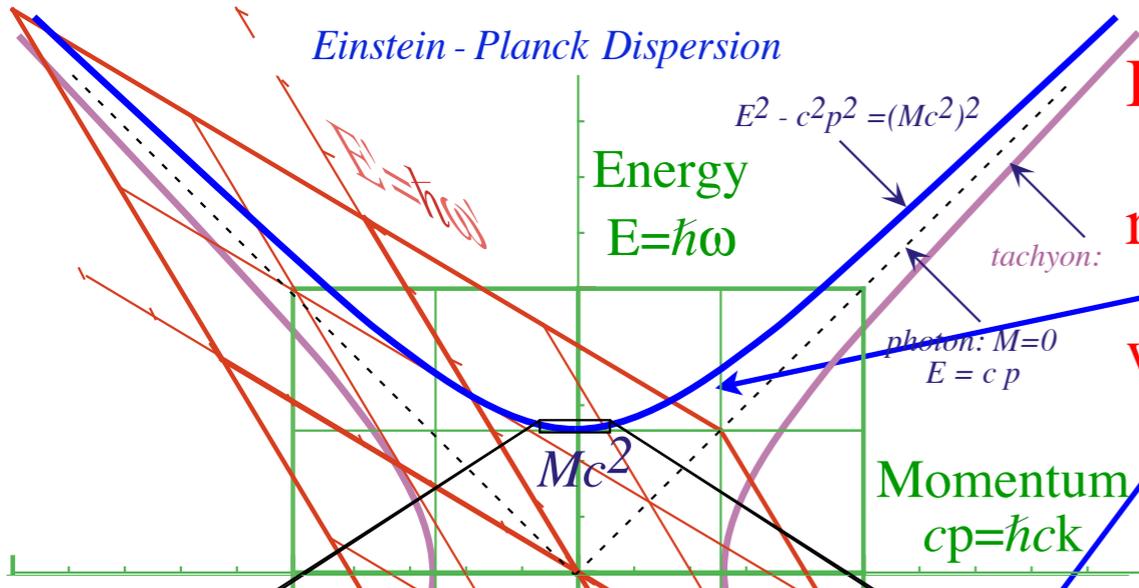
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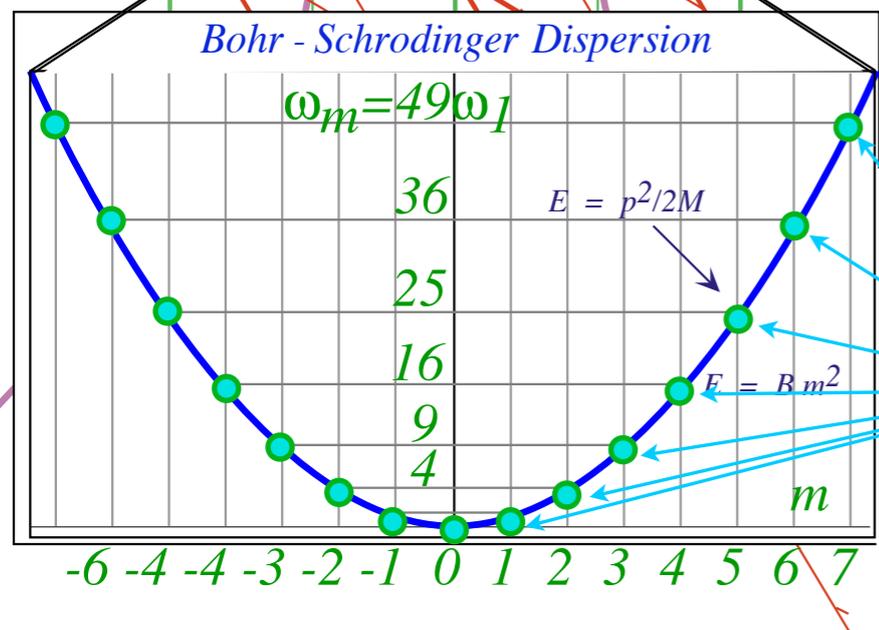
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1st quantization:

Restrict wave-number k to integer or $1/2$ -integer quanta (to fit in cavity)



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Here is a rare but important case where $\frac{d\omega(k)}{dk}$ equals $\frac{\Delta\omega}{\Delta k}$. (Usually not so unless limit $\Delta k \rightarrow 0$ exists.)

Standard formula for classical group velocity is $V_{group} = \frac{d\omega(k)}{dk}$

But this may fail if $\omega(k)$ is quantized and thus discrete $\omega(k_m)$.

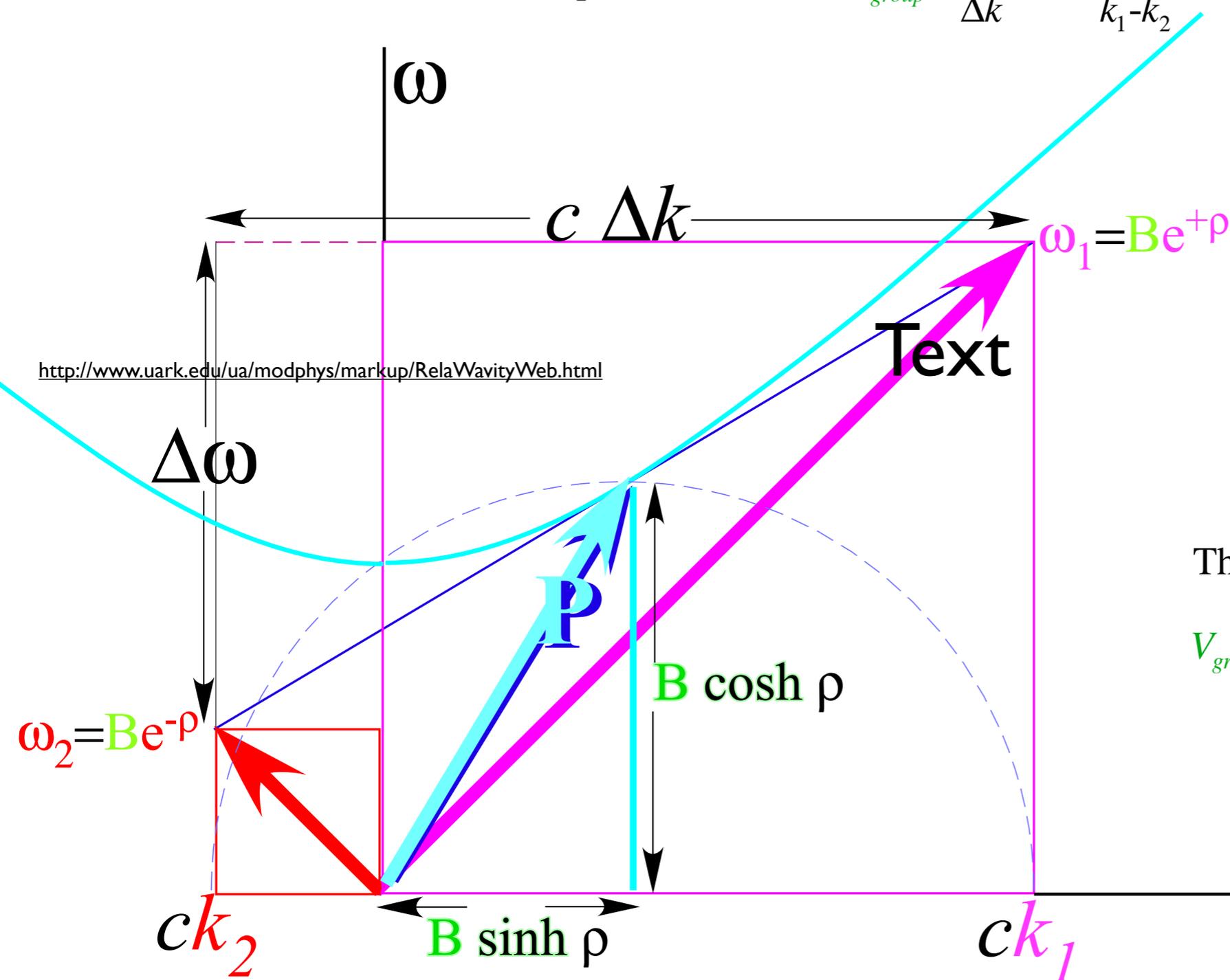
Then we need to use an exact quantum form: $V_{group} = \frac{\Delta\omega}{\Delta k} = \frac{\omega(k_1) - \omega(k_2)}{k_1 - k_2}$

Here the exact discrete value is:

$$\begin{aligned} V_{group} &= \frac{\Delta\omega}{\Delta k} = \frac{\omega(k_1) - \omega(k_2)}{k_1 - k_2} \\ &= \frac{Be^{+\rho} - Be^{-\rho}}{Be^{+\rho} - (-Be^{-\rho})} \\ &= \frac{e^{+\rho} - e^{-\rho}}{e^{+\rho} + Be^{-\rho}} = \frac{\sinh \rho}{\cosh \rho} = \tanh \rho \end{aligned}$$

This time it matches calculus value:

$$\begin{aligned} V_{group} &= \frac{d\omega(k)}{dk} = \frac{d(\cosh \rho)}{d(\sinh \rho)} \\ &= \frac{\sinh \rho}{\cosh \rho} = \tanh \rho \end{aligned}$$



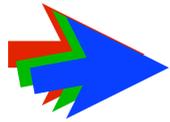
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[Review per-space-time](#)

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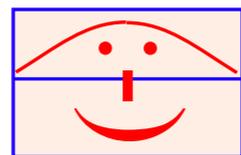
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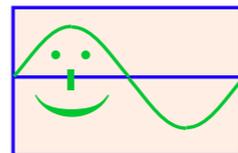
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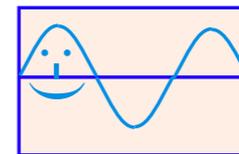
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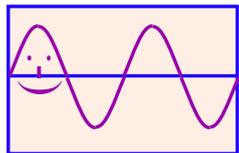
1 half-wave



2 half-waves



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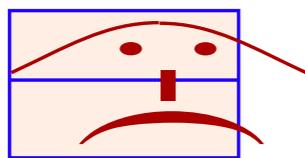


4 half-waves

Some

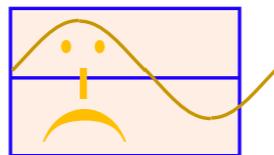
NOT OK numbers: $n=0.67$

too fat!



$n=1.7$

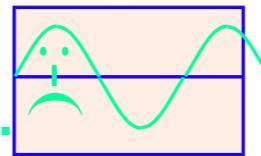
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$n=2.59$

wrong color again!

misfits...



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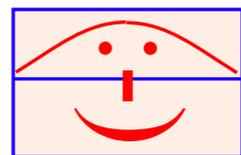
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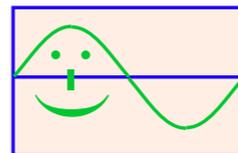
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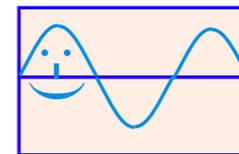
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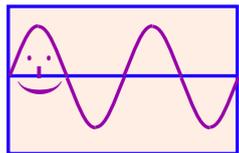
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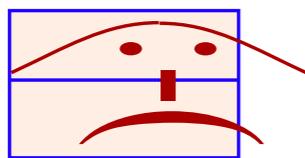


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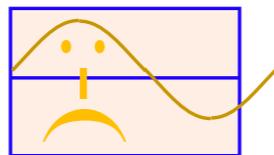
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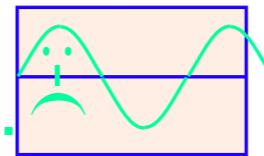
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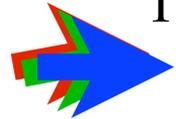
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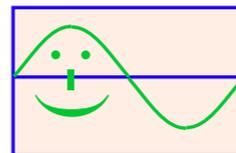
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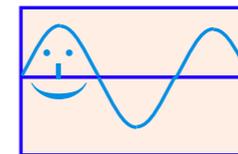
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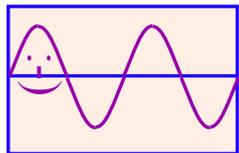
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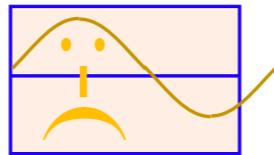
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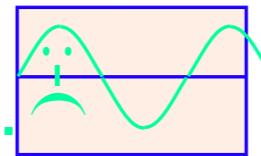
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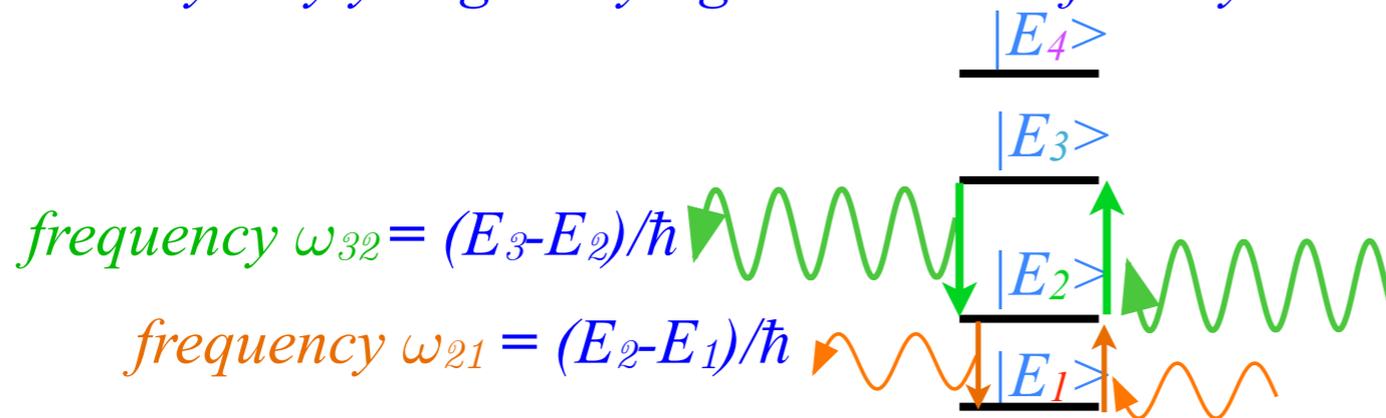
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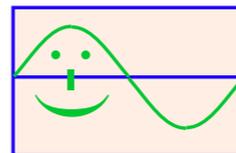
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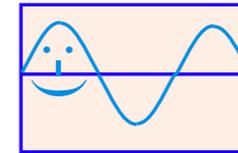
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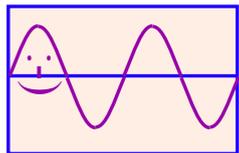
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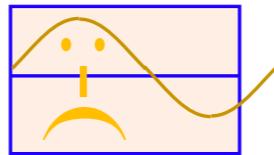
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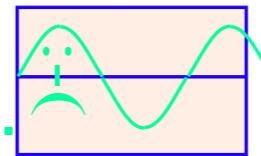
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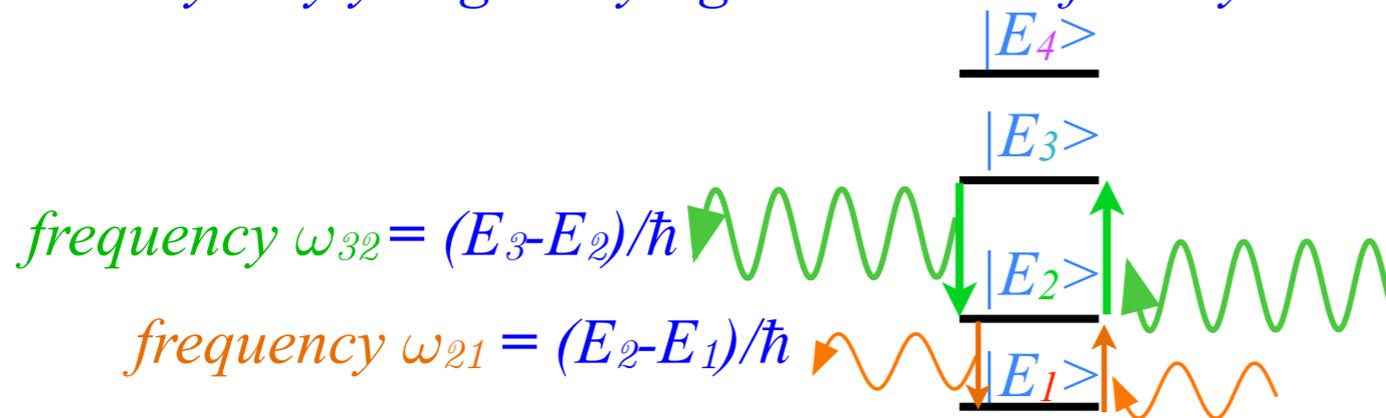
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These *eigenstates* are just the ways the wavy system can “play dead” ...

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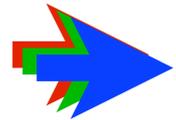
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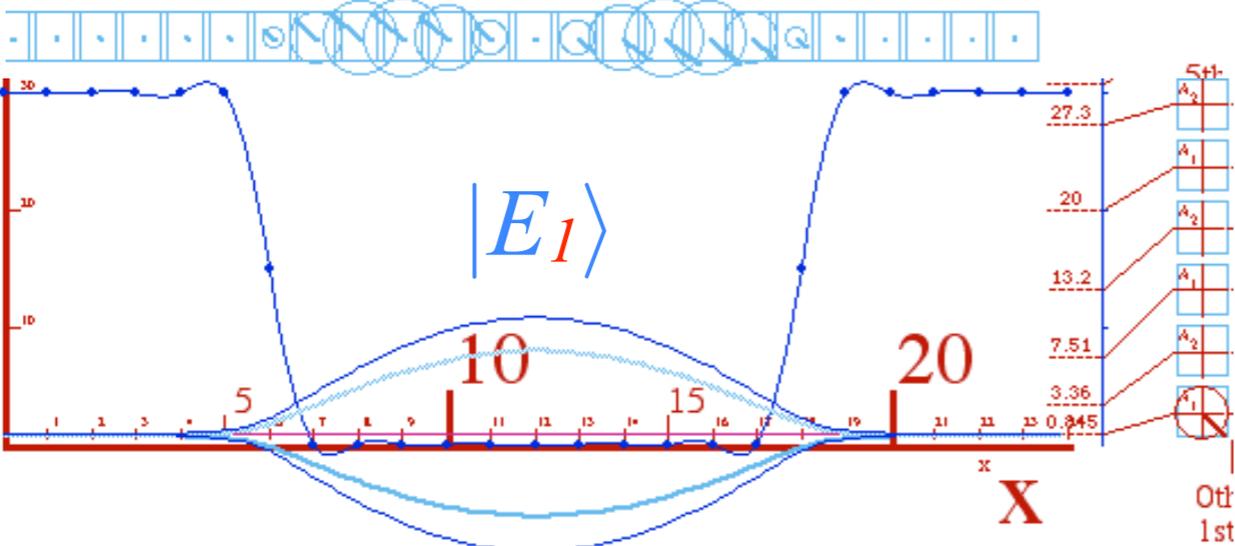
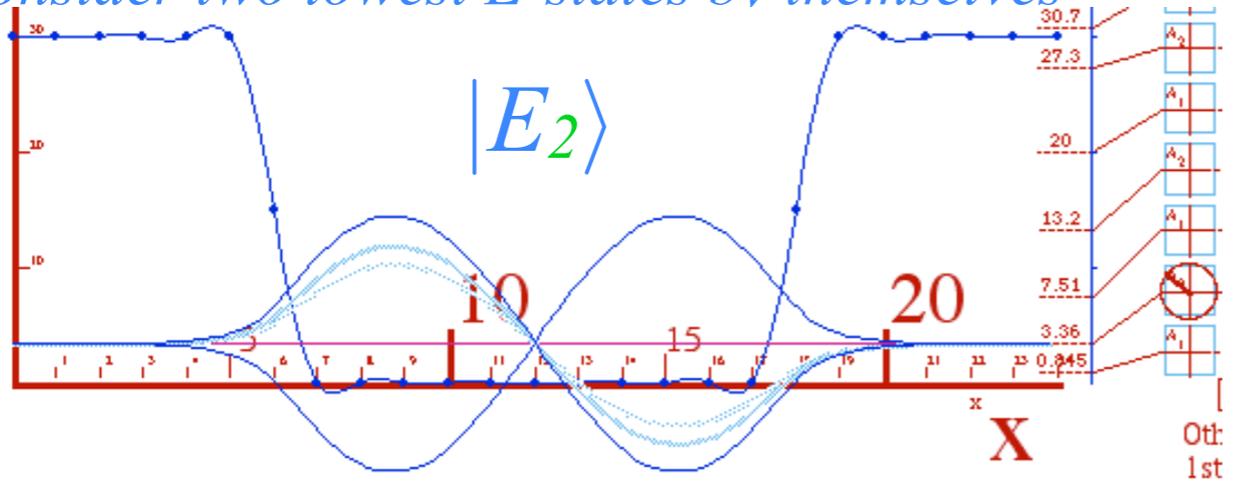
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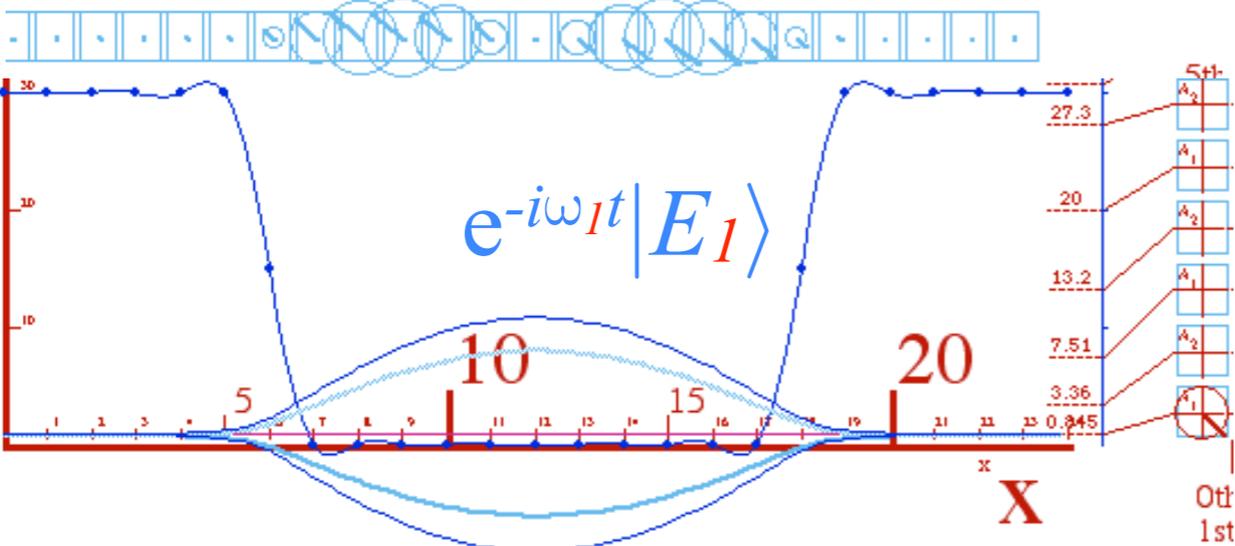
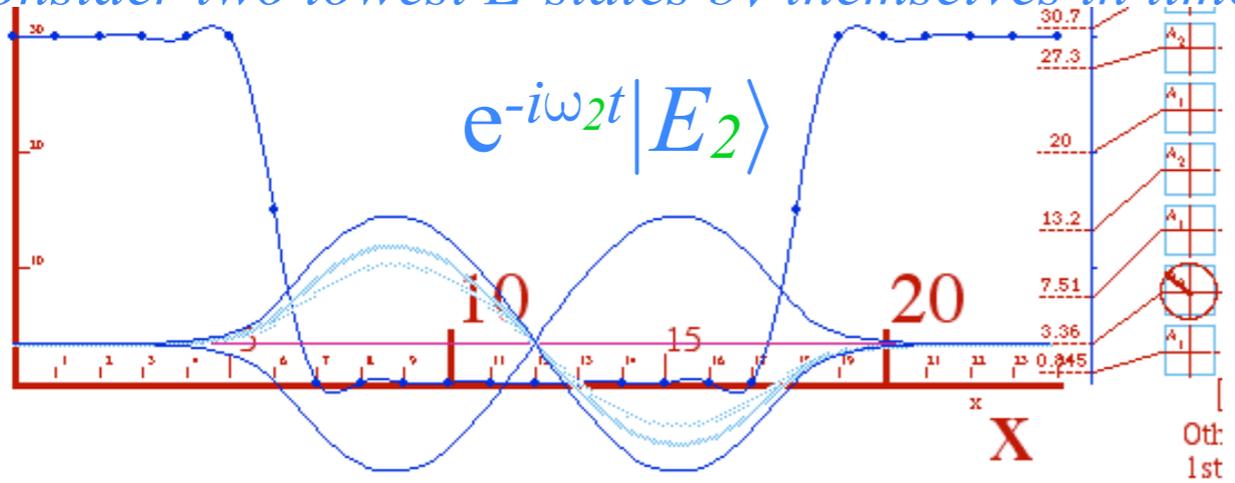


Consider two lowest E -states by themselves



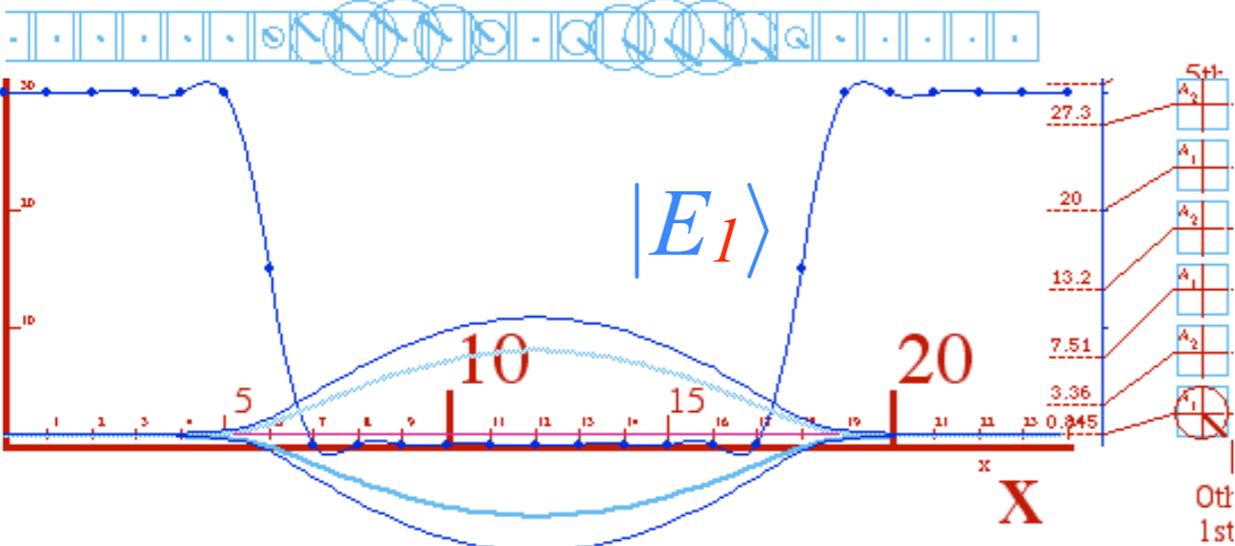
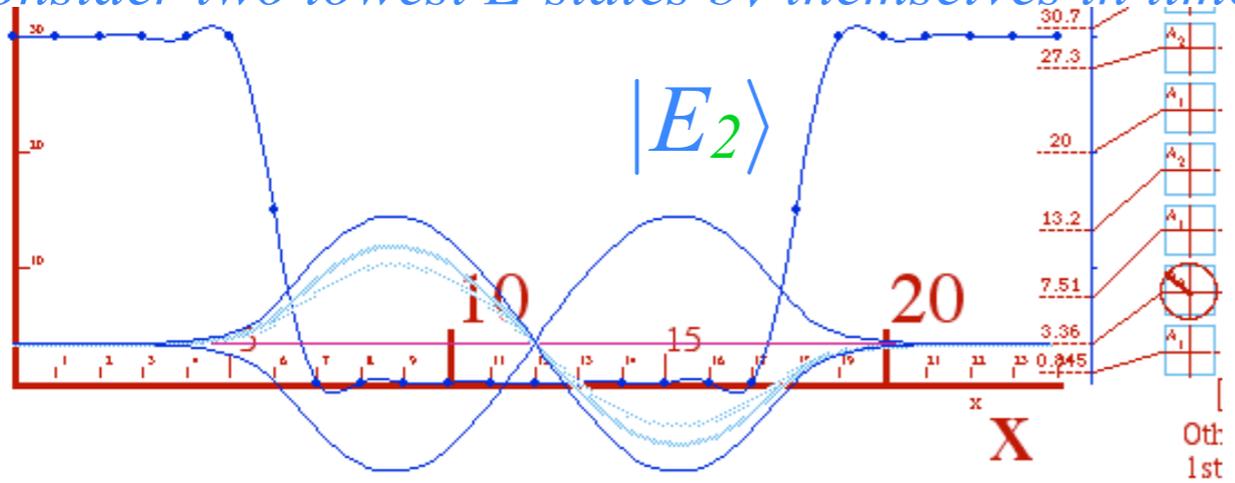
By Harter-*et al* and University of Arkansas Physics *Elegant Educational Tools Since 2001*

Consider two lowest E-states by themselves in time



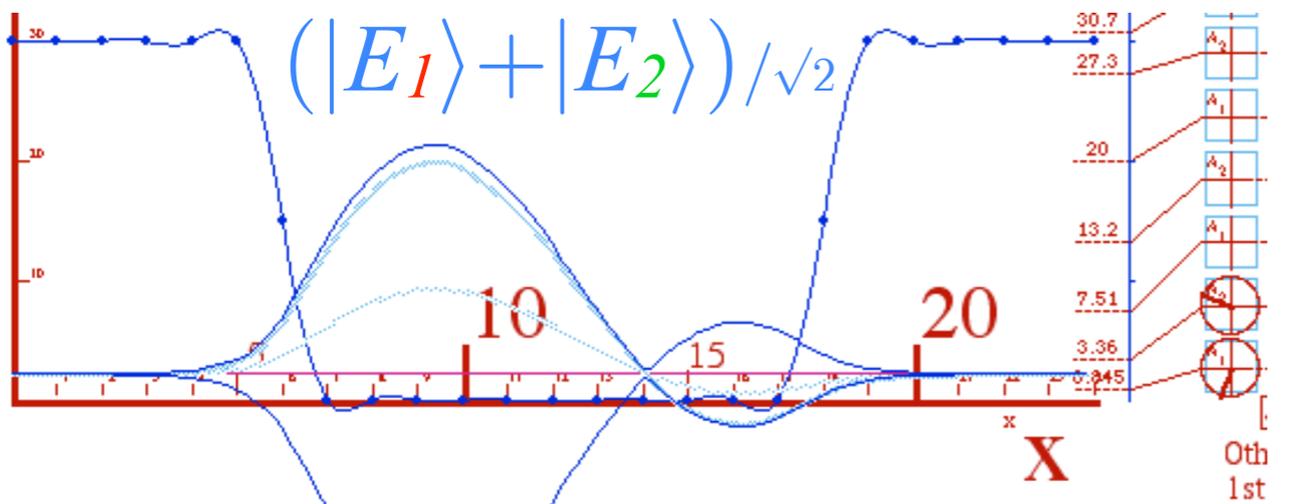
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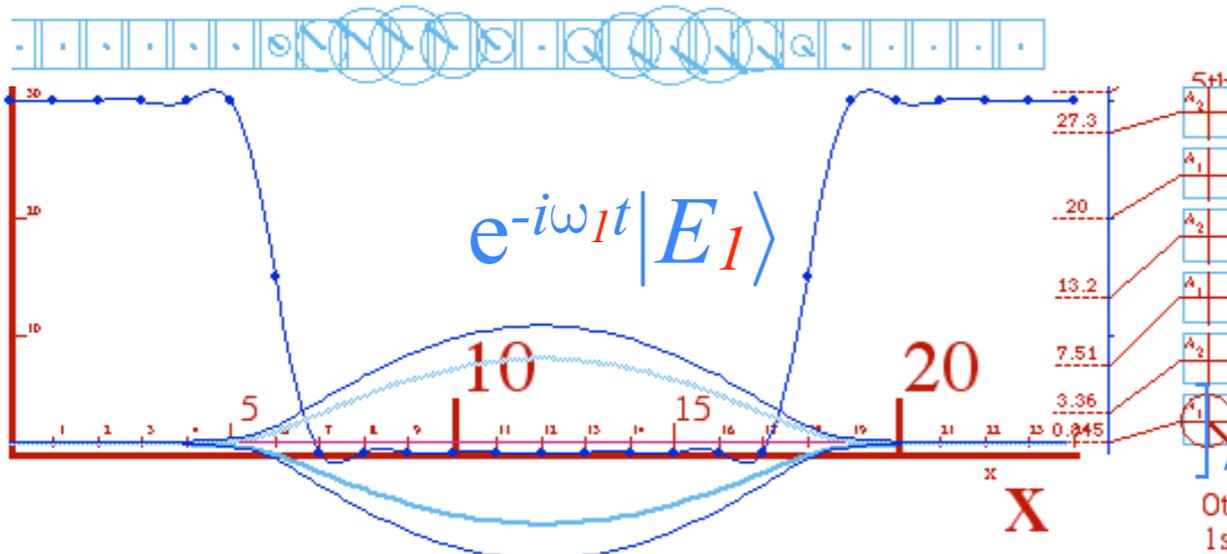
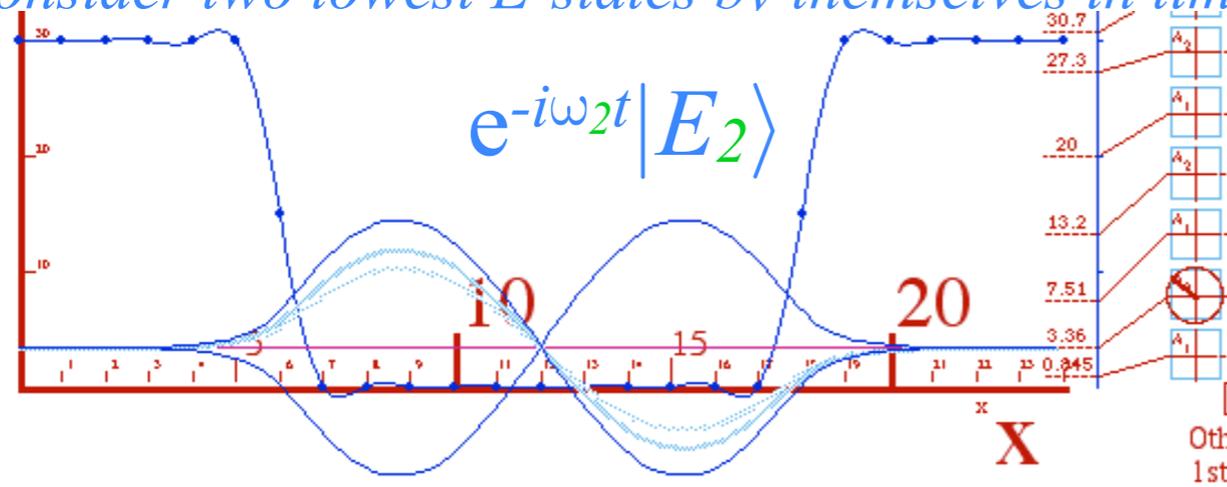
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Now combine (add) them



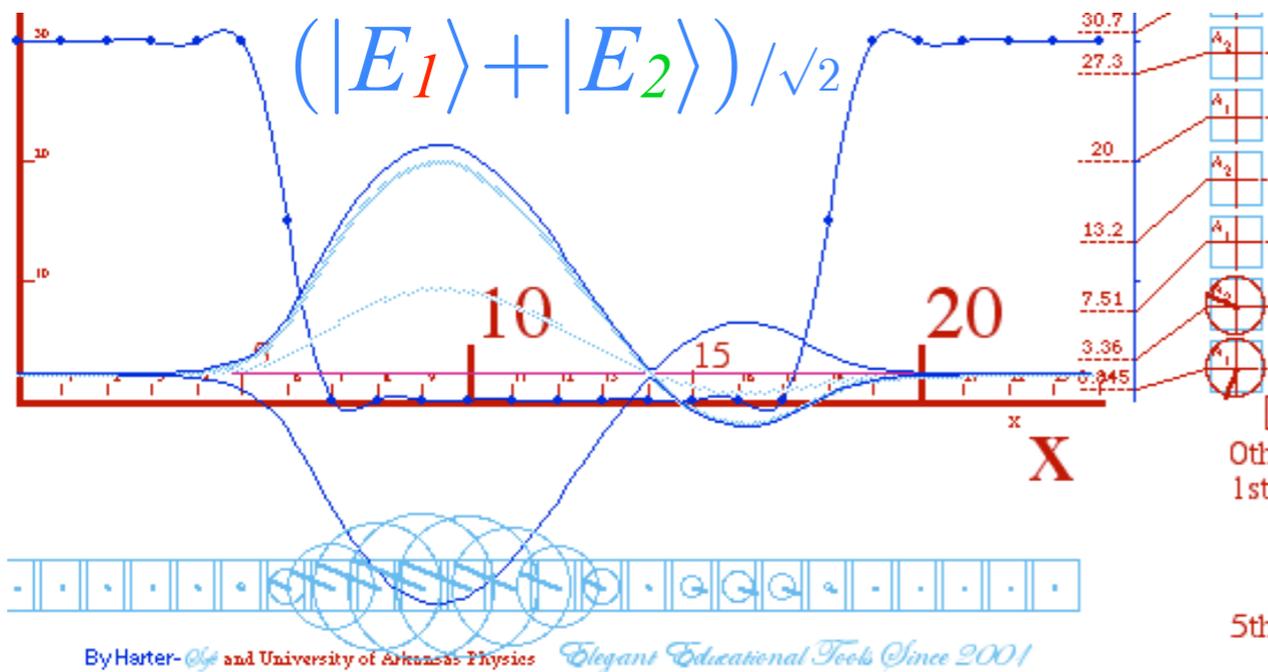
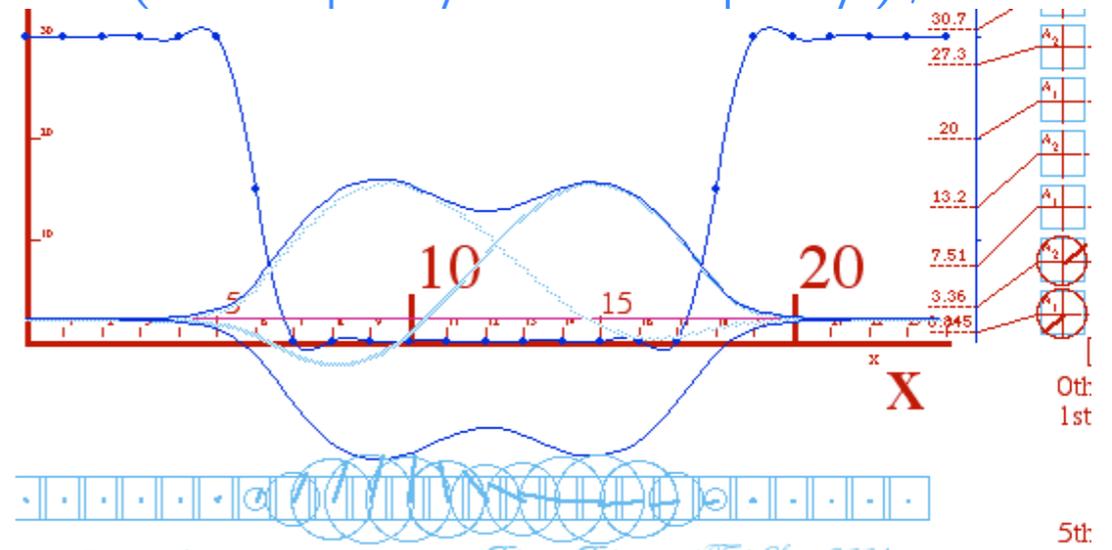
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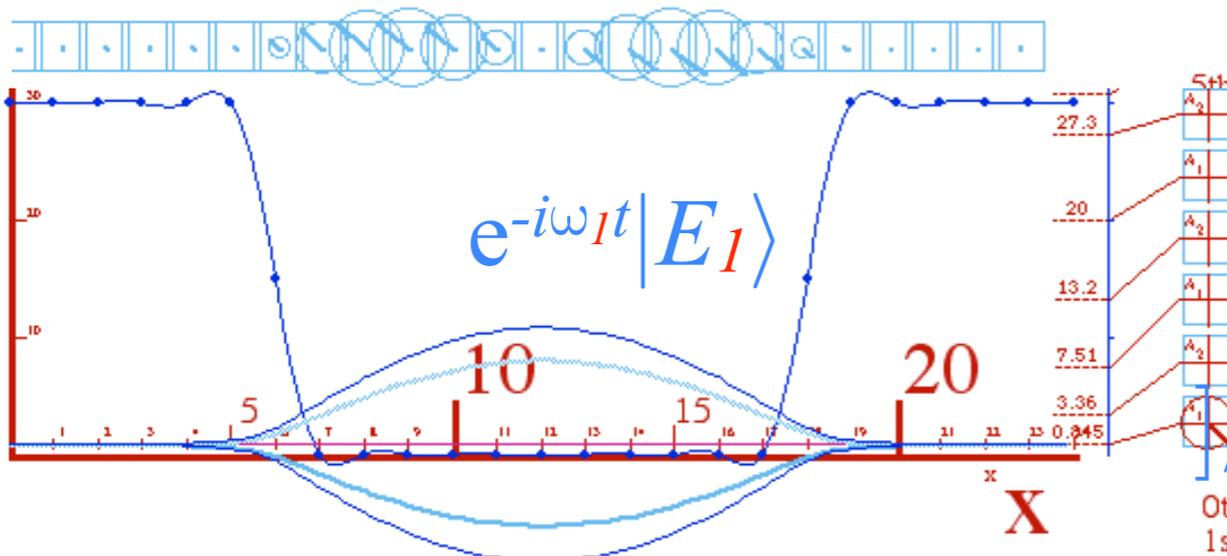
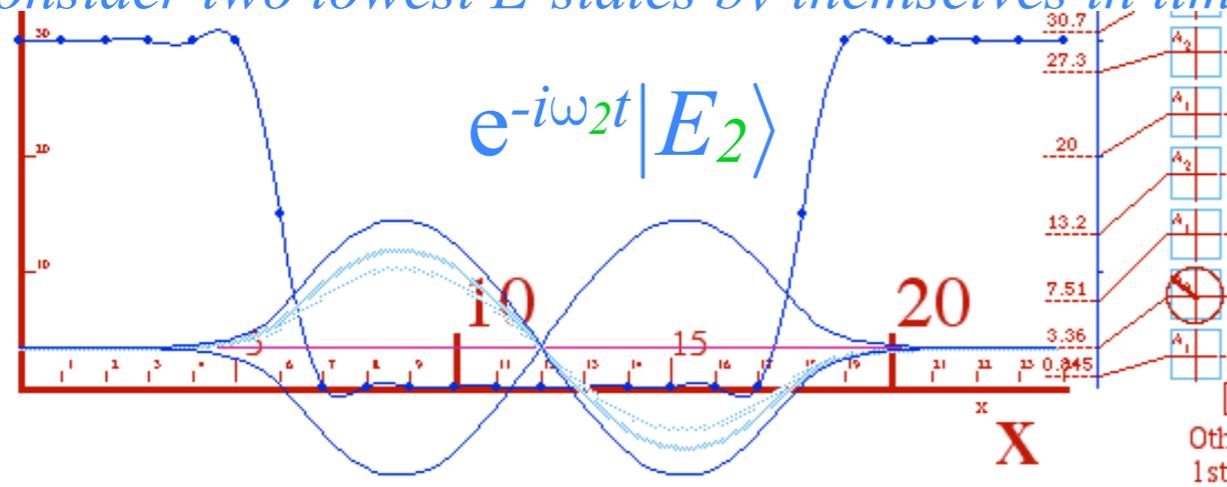


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$$\left(e^{-i\omega_1 t} |E_1\rangle + e^{-i\omega_2 t} |E_2\rangle \right) / \sqrt{2}$$



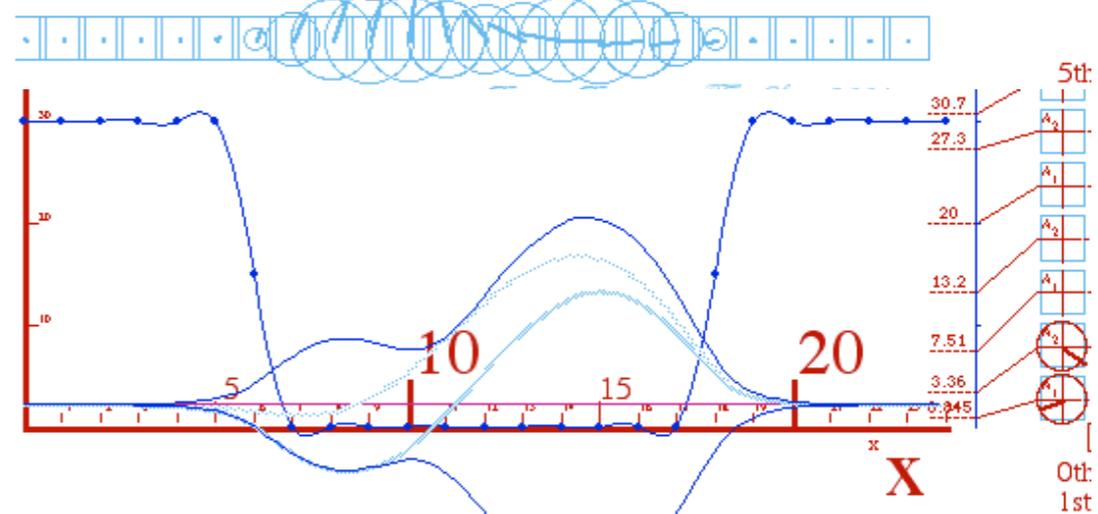
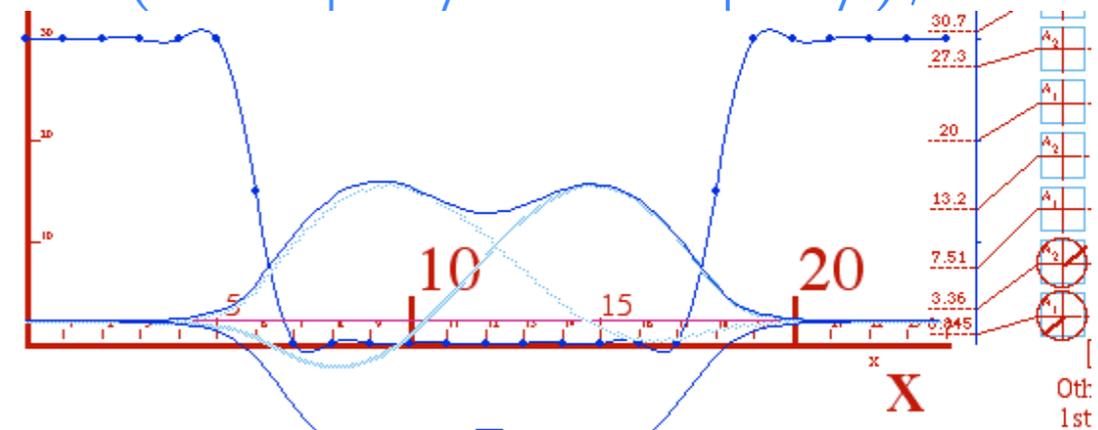
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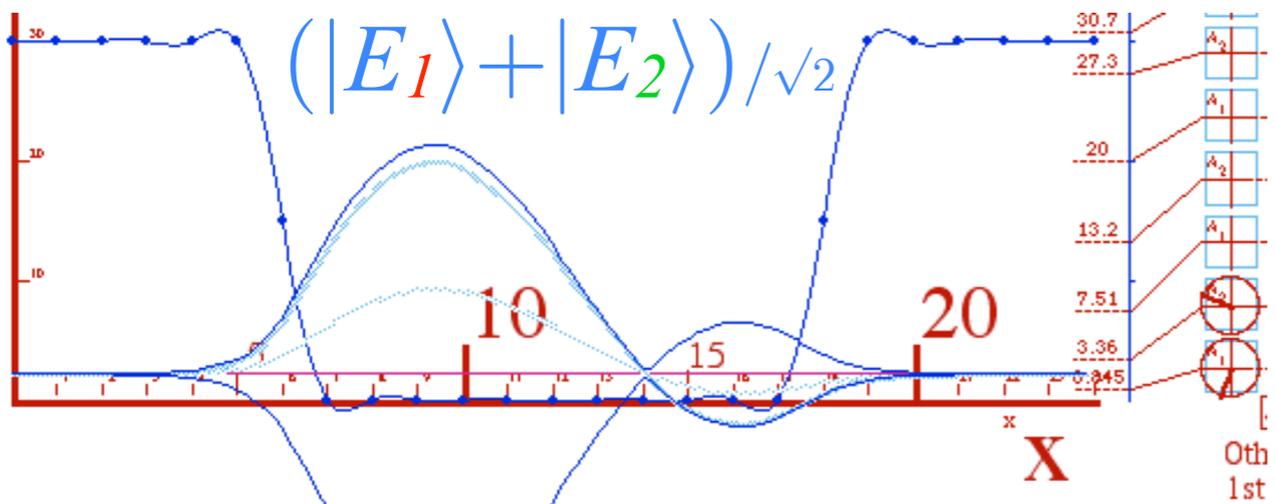
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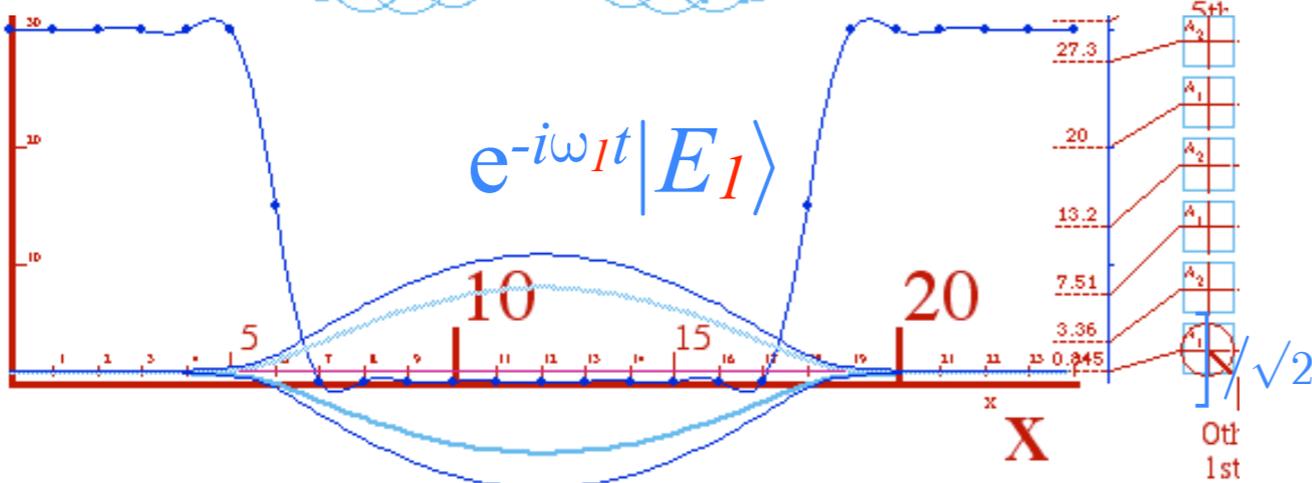
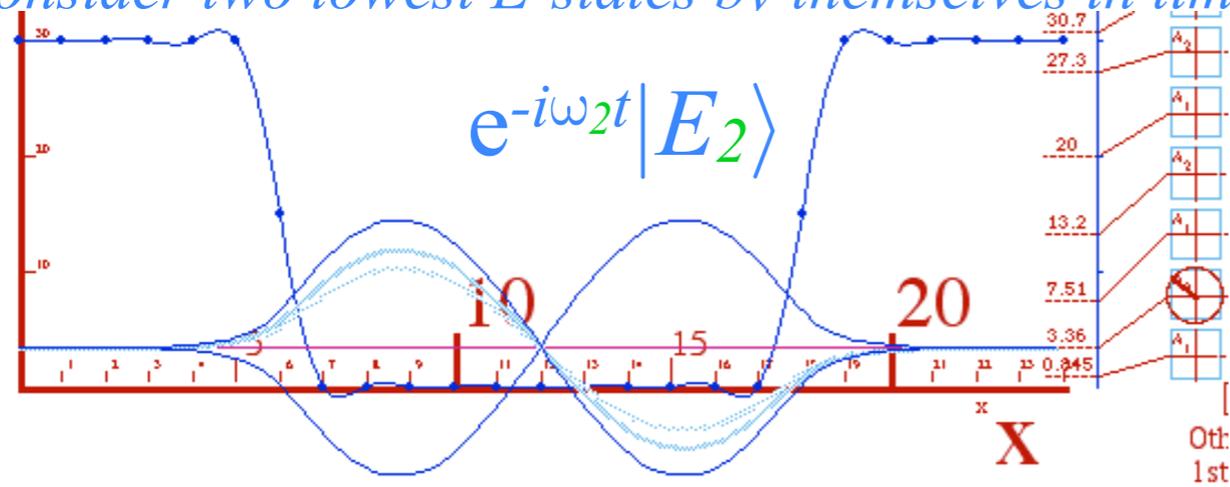


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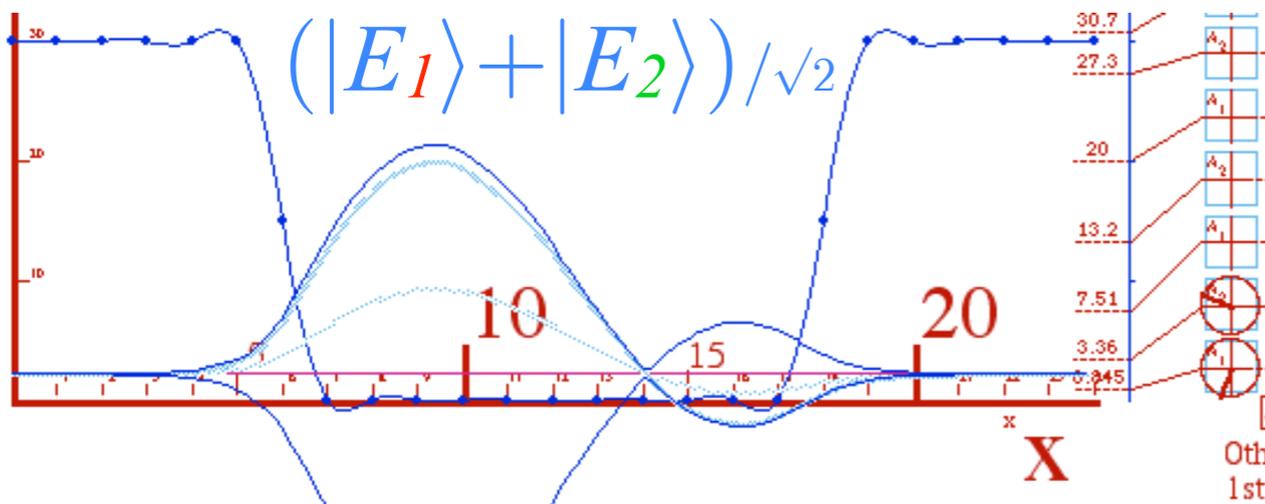


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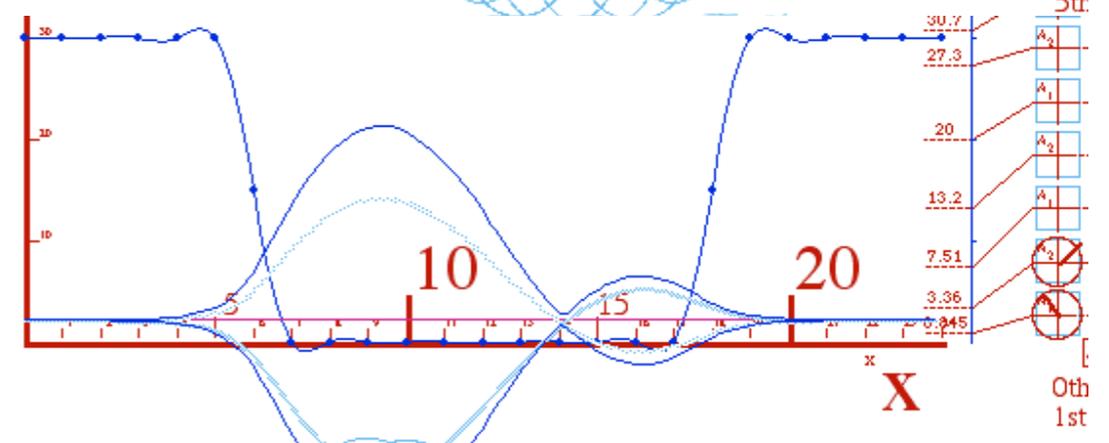
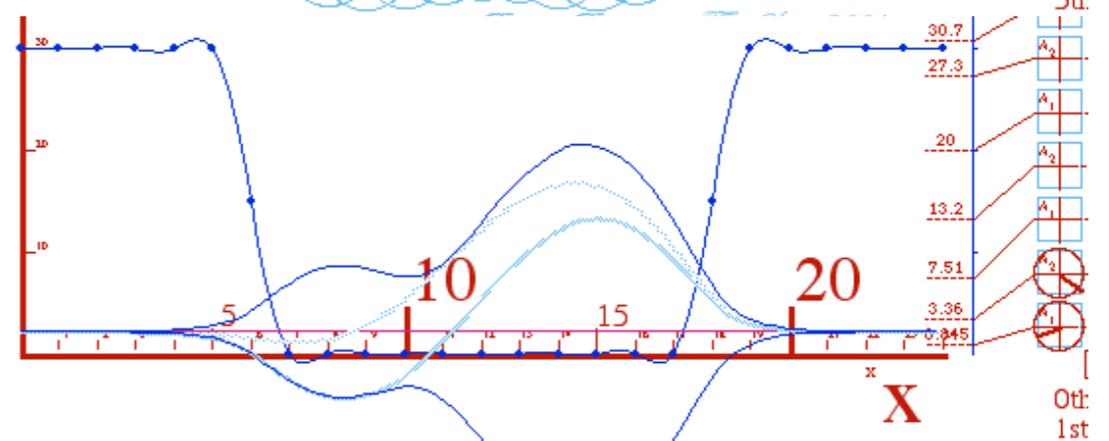
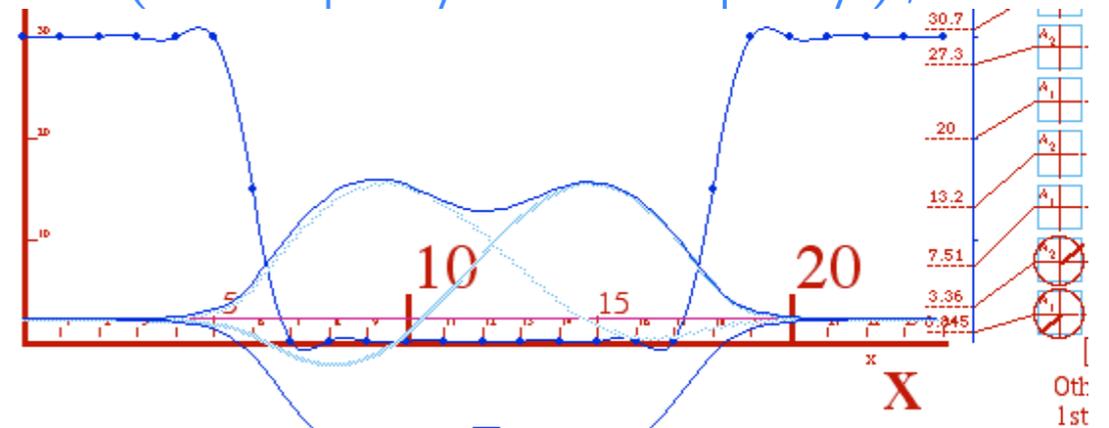
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$$(e^{-i\omega_1 t} |E_1\rangle + e^{-i\omega_2 t} |E_2\rangle) / \sqrt{2}$$



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5th

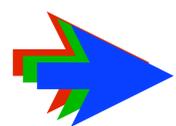
Review of wave parameters used to develop relativistic quantum theory

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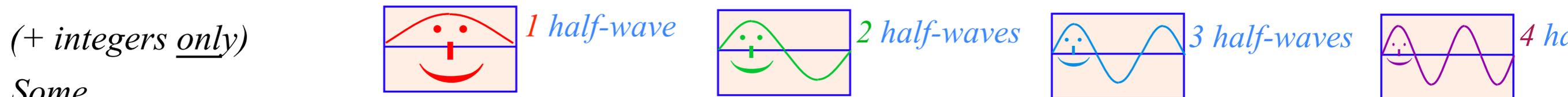
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Relativistic effects on charge, current, and Maxwell Fields

Quantized ω and k Counting wave kink numbers

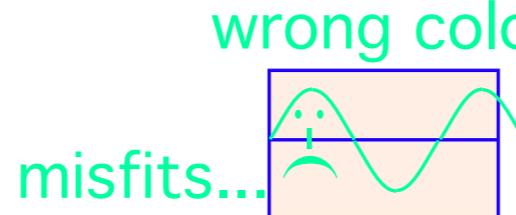
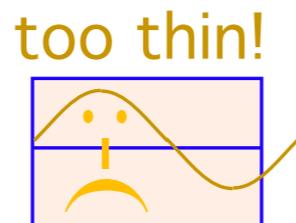
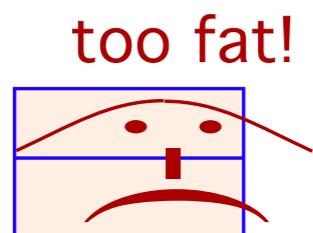
If everything is made of waves then we expect *quantization* of everything because waves only thrive if *integral* numbers n of their “kinks” fit into whatever structure (box, ring, etc.) they’re supposed to live. The numbers n are called *quantum numbers*.

OK box quantum numbers: $n=1$ $n=2$ $n=3$ $n=4$



Some

NOT OK numbers: $n=0.67$ $n=1.7$ $n=2.59$ $n=4$



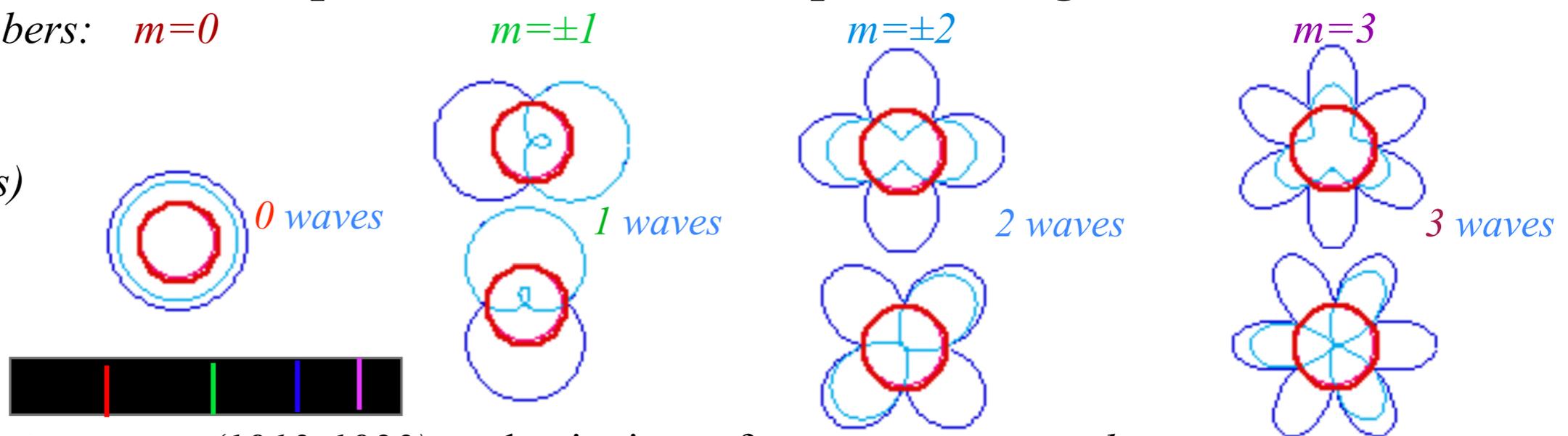
...not tolerated!

NOTE: We're using "false-color" here.

Rings tolerate a *zero* (kinkless) quantum wave but require \pm integral wave number.

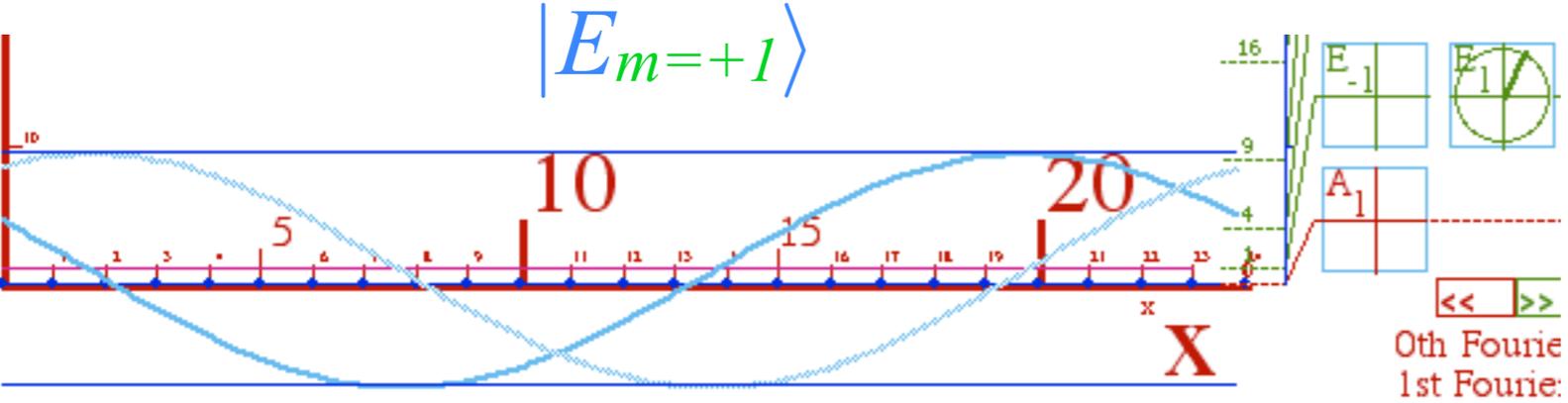
OK ring quantum numbers: $m=0$ $m=\pm 1$ $m=\pm 2$ $m=3$

(\pm integral number of wavelengths)

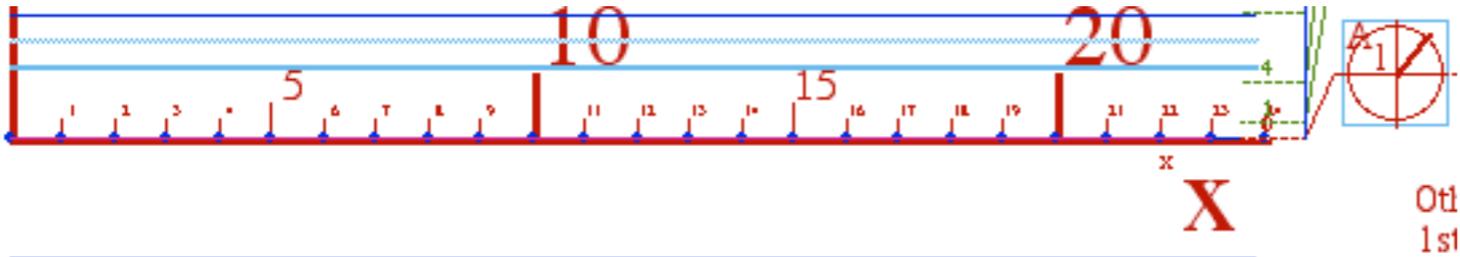


Bohr's models of *atomic spectra* (1913-1923) are beginnings of *quantum wave mechanics* built on *Planck-Einstein* (1900-1905) relation $E=h\nu$. *DeBroglie* relation $p=h/\lambda$ comes around 1923.

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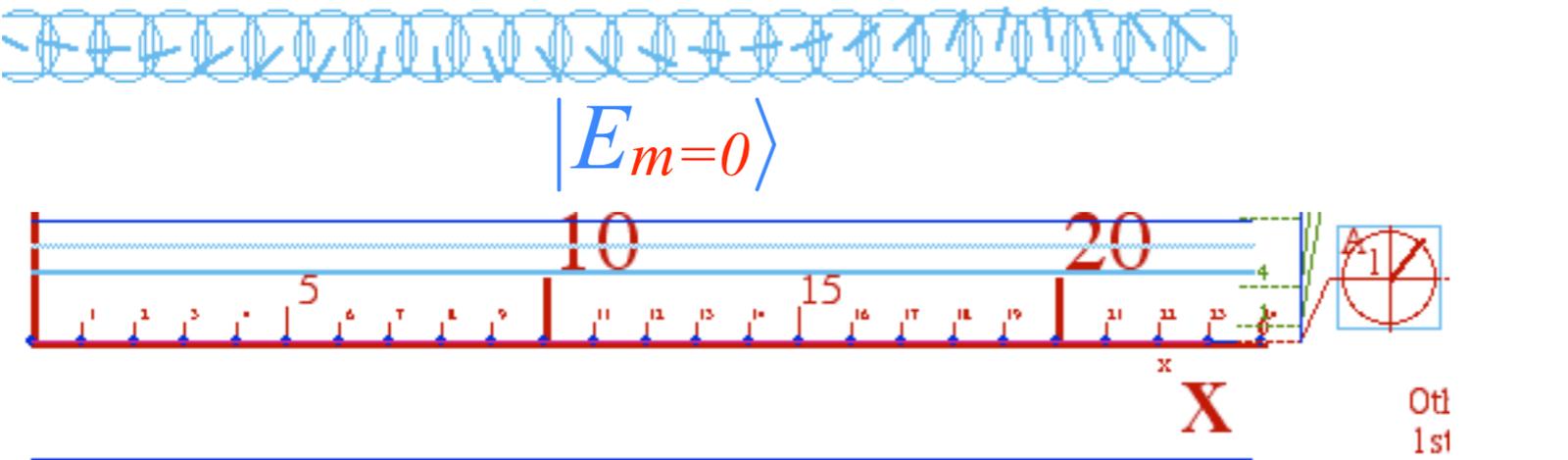
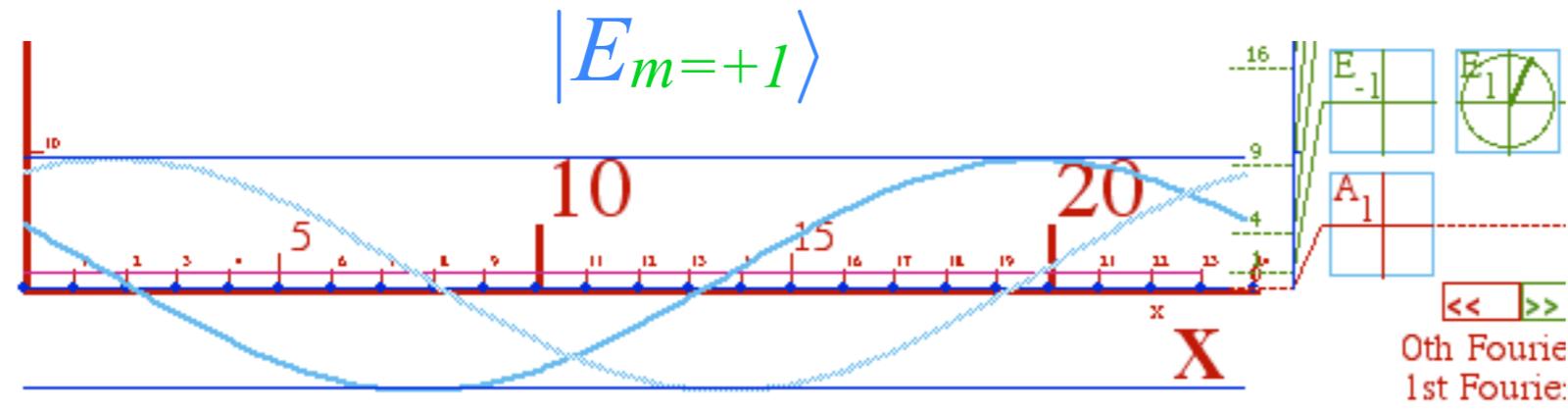
$|E_{m=0}\rangle$



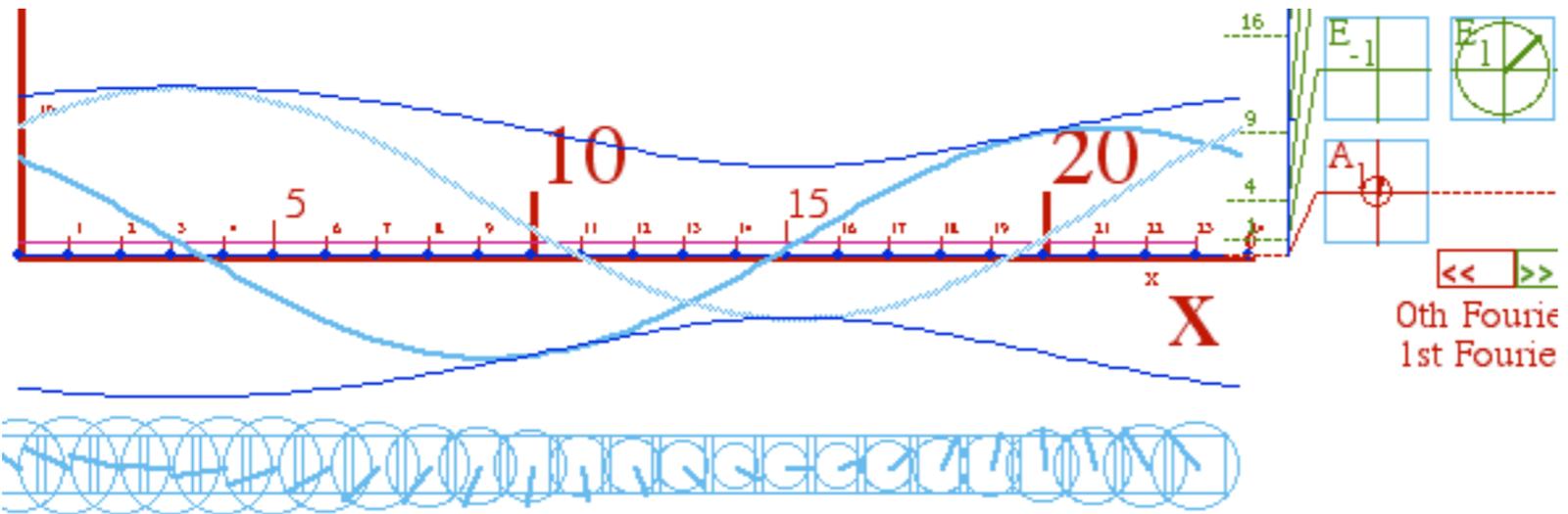
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Now combine (add) them and let time roll!

$$\left(e^{-i\omega_0 t} |E_0\rangle + e^{-i\omega_{+1} t} |E_{+1}\rangle \right) / \sqrt{2}$$



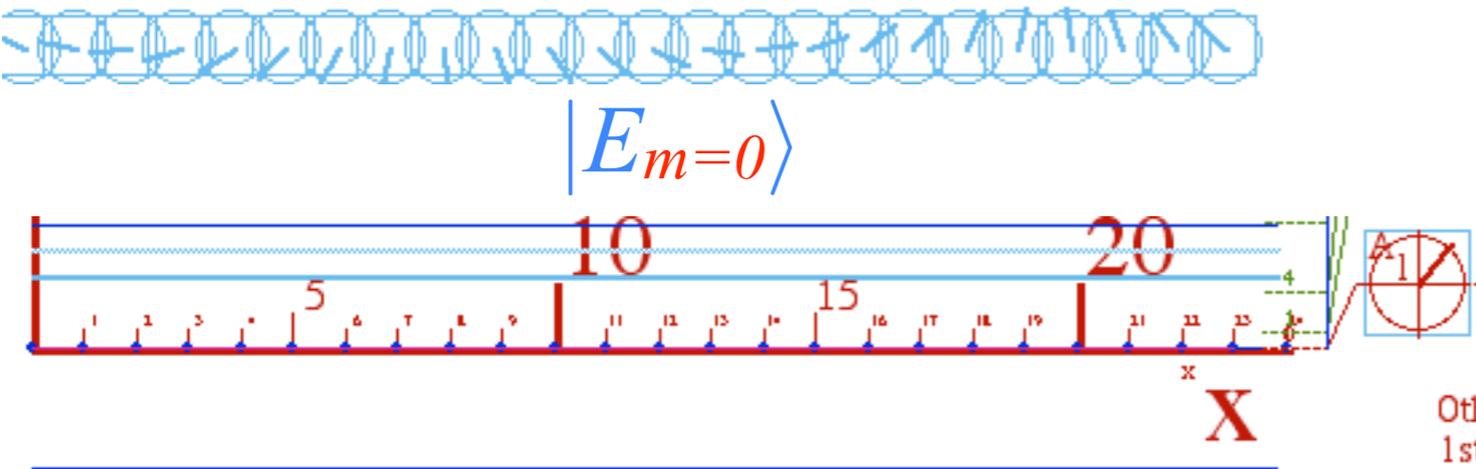
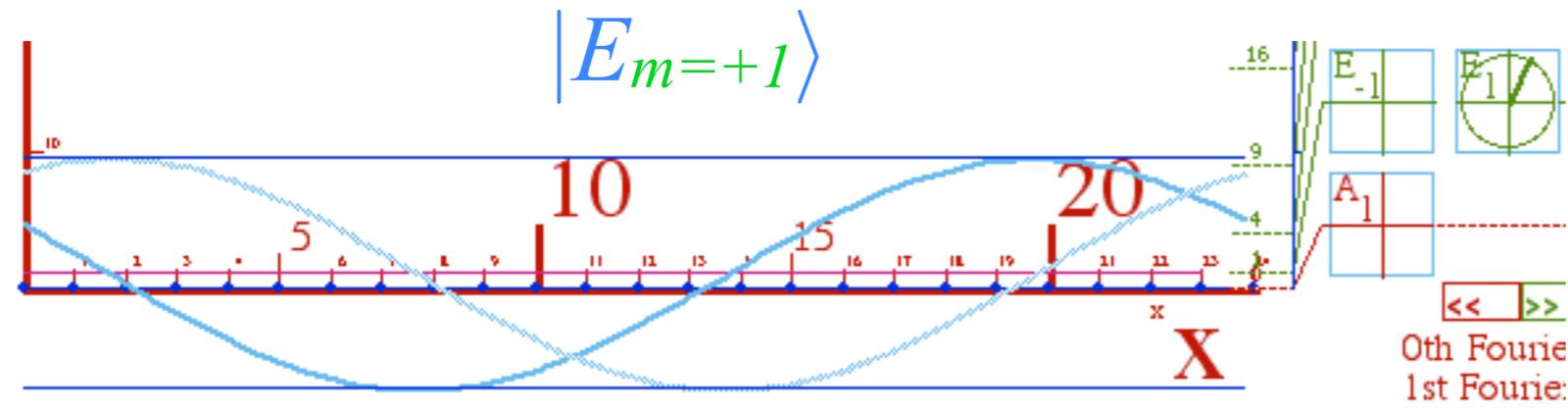
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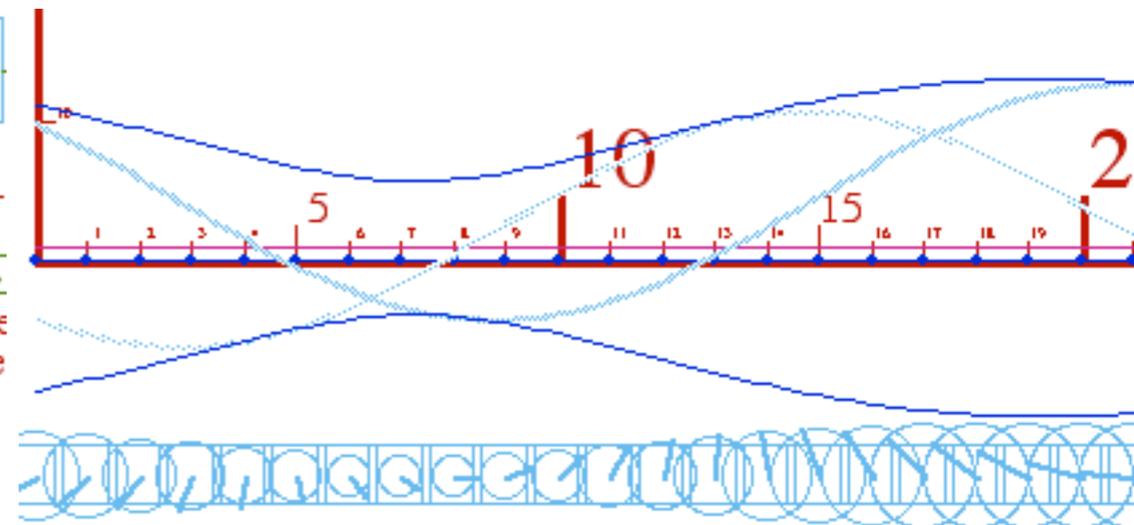
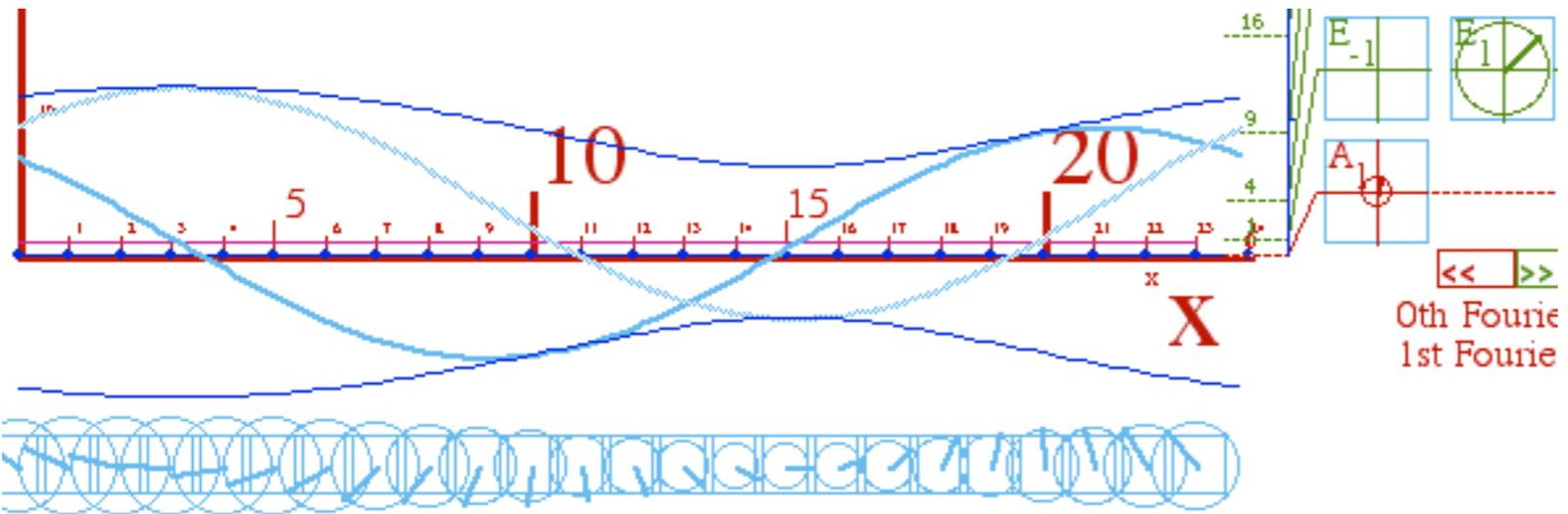
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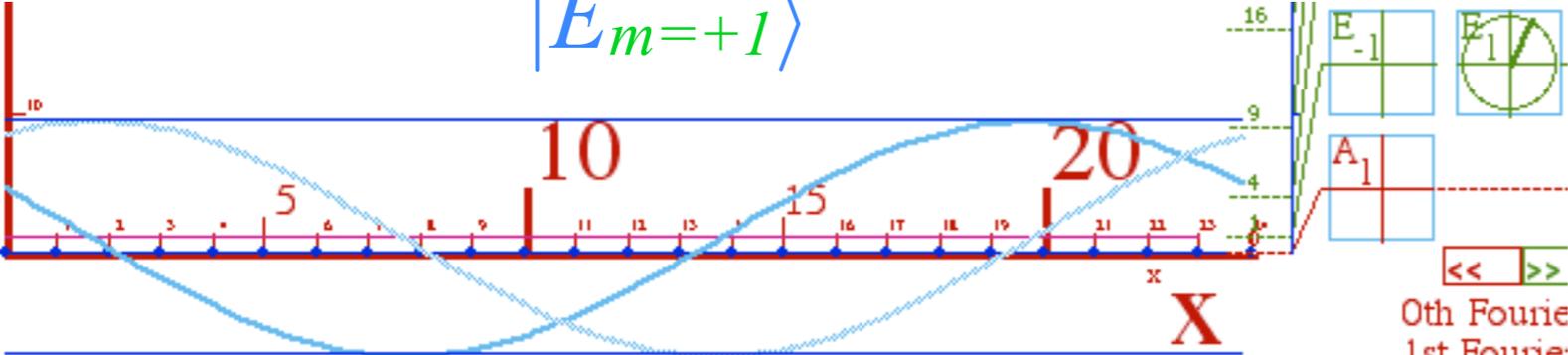
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(Just moves forward rigidly)

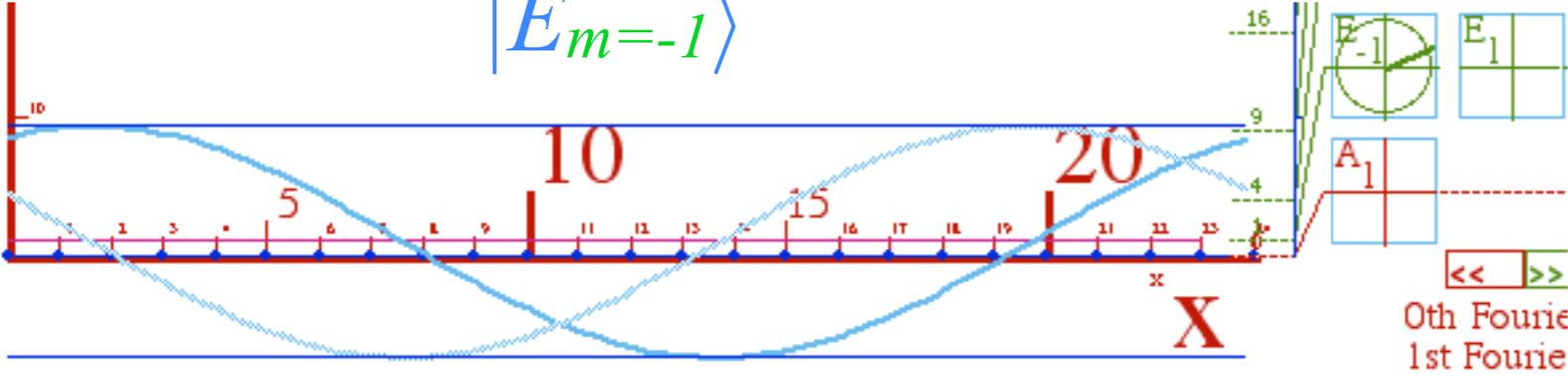


Consider two degenerate E-states by themselves

$$|E_{m=+1}\rangle$$

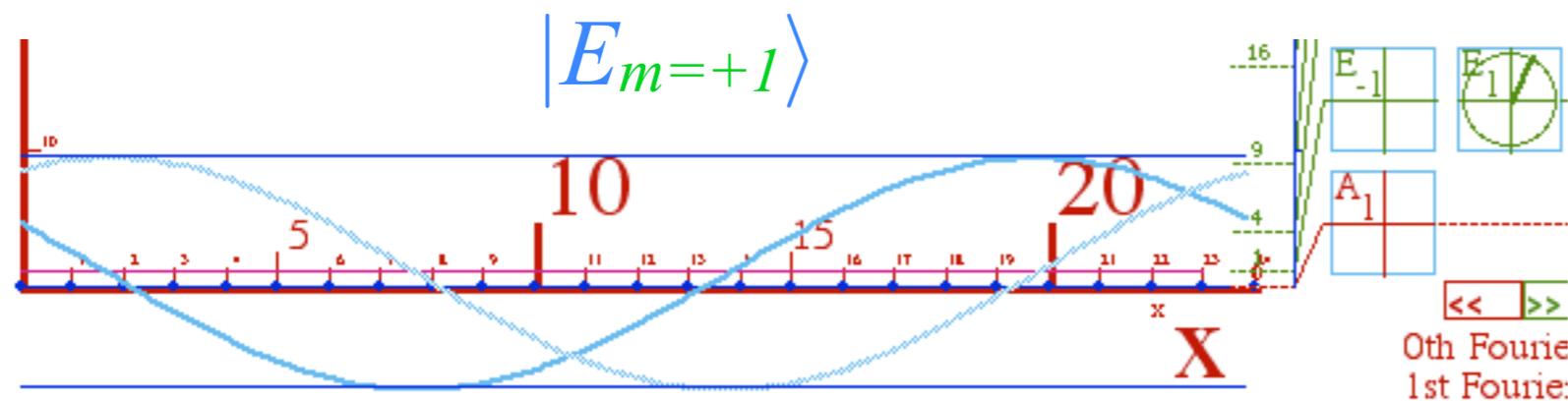


$$|E_{m=-1}\rangle$$

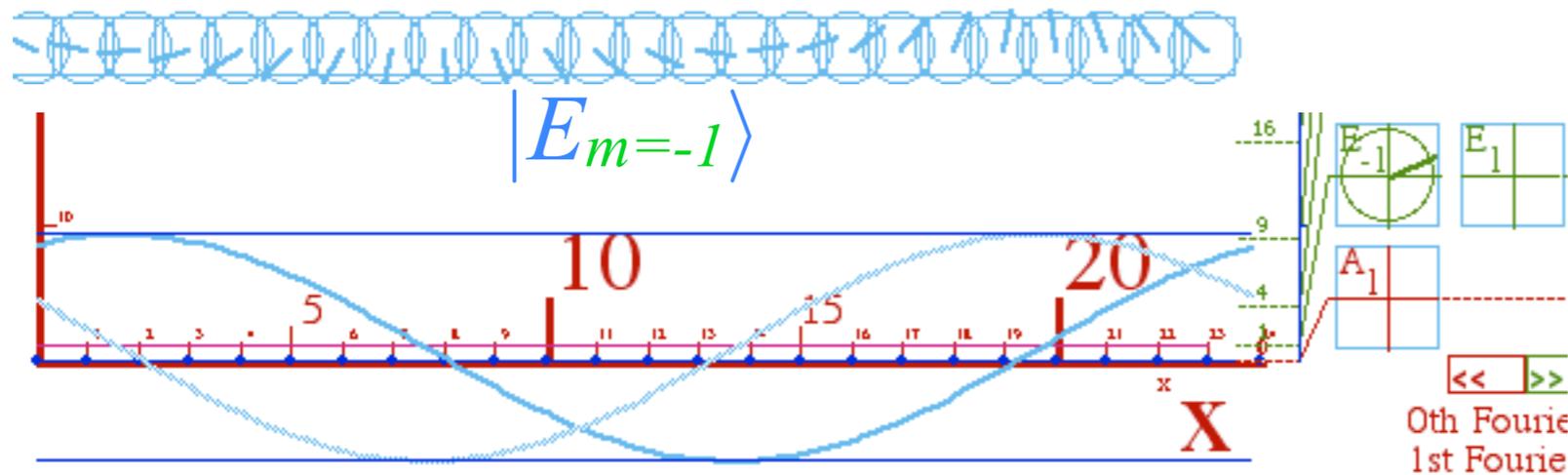


Consider two *degenerate* E-states by themselves

Now combine (add) them and let time roll!



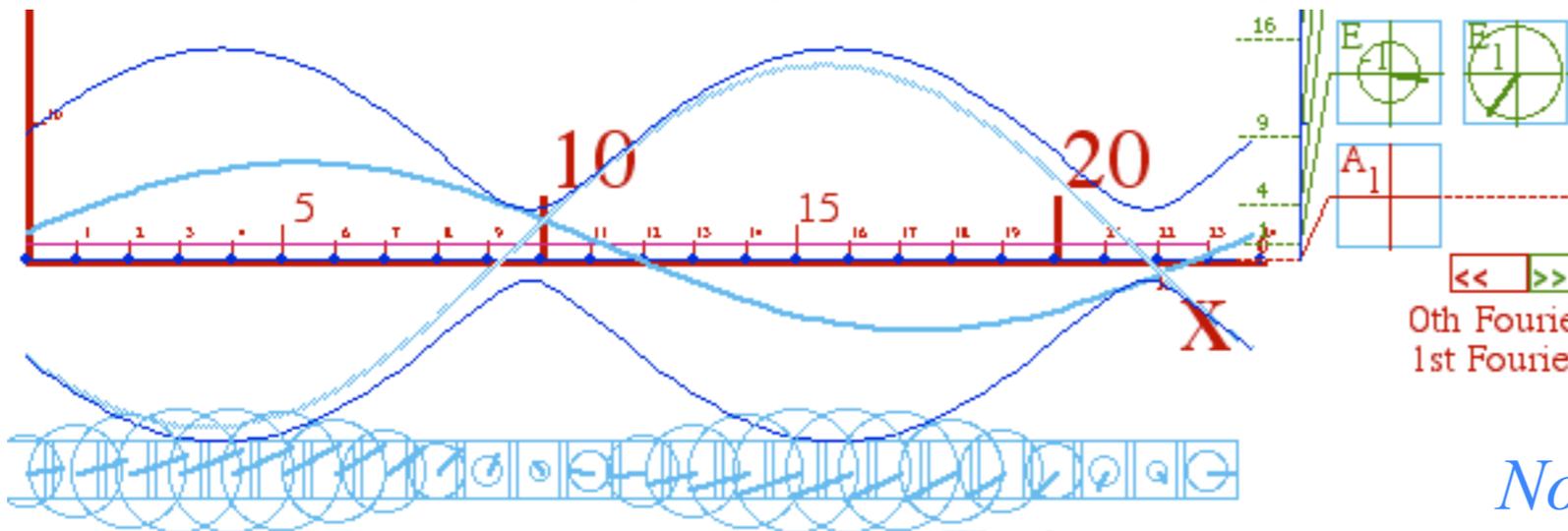
$$\frac{(e^{-i\omega_{-1}t}|E_{-1}\rangle + e^{-i\omega_{+1}t}|E_{+1}\rangle)}{\sqrt{2}}$$



Group wave is stationary if $\omega_{-1} = \omega_{+1}$ but phase can move or “gallop” faster than light!

If $\omega_{-1} < \omega_{+1}$ then $V_{group} < 0$

If $\omega_{-1} > \omega_{+1}$ then $V_{group} > 0$



Nothing CAN go faster than light

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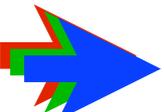
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2-CW dynamics has two 1-CW amplitudes A_{\rightarrow} and A_{\leftarrow} that may be *unmatched*. ($A_{\rightarrow} \neq A_{\leftarrow}$)

$$A_{\rightarrow} e^{i(k_{\rightarrow} x - \omega_{\rightarrow} t)} + A_{\leftarrow} e^{i(k_{\leftarrow} x - \omega_{\leftarrow} t)} = e^{i(k_{\Sigma} x - \omega_{\Sigma} t)} [A_{\rightarrow} e^{i(k_{\Delta} x - \omega_{\Delta} t)} + A_{\leftarrow} e^{-i(k_{\Delta} x - \omega_{\Delta} t)}]$$

Waves have half-sum mean-phase rates $(k_{\Sigma}, \omega_{\Sigma})$ and half-difference group rates $(k_{\Delta}, \omega_{\Delta})$.

$$k_{\Sigma} = (k_{\rightarrow} + k_{\leftarrow}) / 2$$

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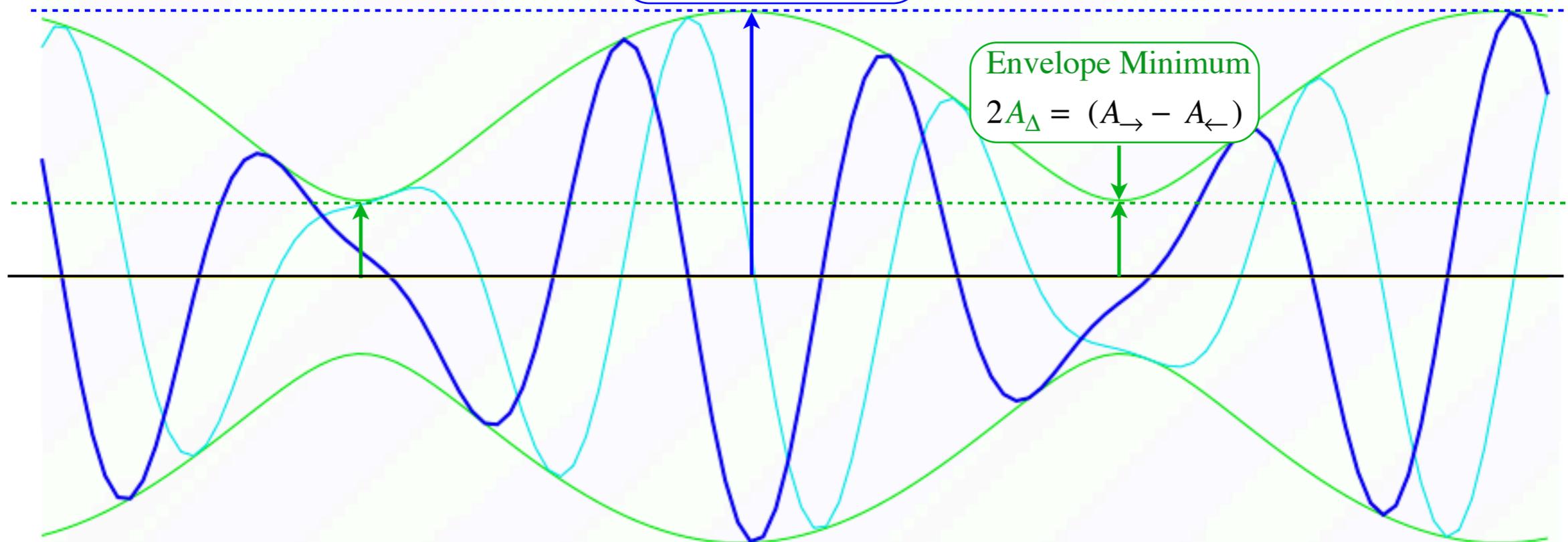
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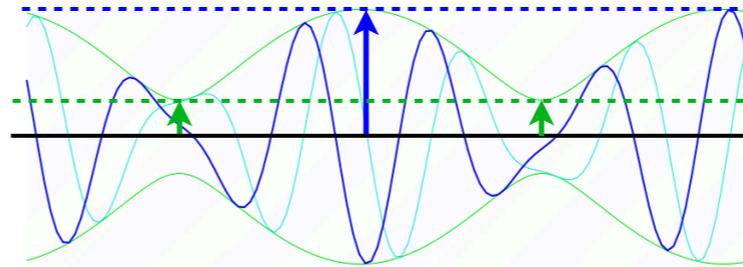
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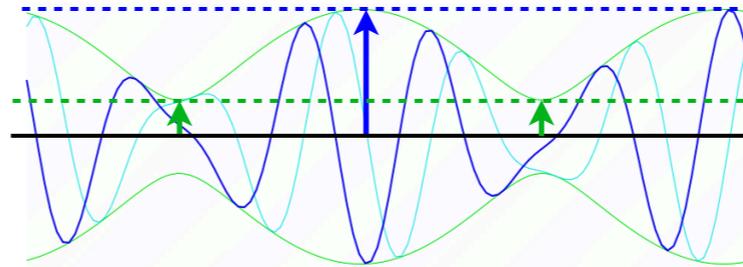
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$$V_{\text{phase}} = \frac{\omega_{\Sigma}}{k_{\Sigma}} = \frac{(\omega_{\rightarrow} + \omega_{\leftarrow})}{(k_{\rightarrow} + k_{\leftarrow})} = c \frac{(\omega_{\rightarrow} + \omega_{\leftarrow})}{(\omega_{\rightarrow} - \omega_{\leftarrow})}$$

$$\frac{V_{\text{group}}}{c} = \frac{\omega_{\Delta}}{ck_{\Delta}} = \frac{(\omega_{\rightarrow} - \omega_{\leftarrow})}{c(k_{\rightarrow} - k_{\leftarrow})} = \frac{(\omega_{\rightarrow} - \omega_{\leftarrow})}{(\omega_{\rightarrow} + \omega_{\leftarrow})}$$

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Galloping waves due to unmatched 2-CW amplitudes

2-CW dynamics has two 1-CW amplitudes A_{\rightarrow} and A_{\leftarrow} that may be *unmatched*. ($A_{\rightarrow} \neq A_{\leftarrow}$)

$$A_{\rightarrow} e^{i(k_{\rightarrow}x - \omega_{\rightarrow}t)} + A_{\leftarrow} e^{i(k_{\leftarrow}x - \omega_{\leftarrow}t)} = e^{i(k_{\Sigma}x - \omega_{\Sigma}t)} [A_{\rightarrow} e^{i(k_{\Delta}x - \omega_{\Delta}t)} + A_{\leftarrow} e^{-i(k_{\Delta}x - \omega_{\Delta}t)}]$$

Waves have half-sum mean-phase rates $(k_{\Sigma}, \omega_{\Sigma})$ and half-difference group rates $(k_{\Delta}, \omega_{\Delta})$.

$$k_{\Sigma} = (k_{\rightarrow} + k_{\leftarrow})/2$$

$$k_{\Delta} = (k_{\rightarrow} - k_{\leftarrow})/2$$

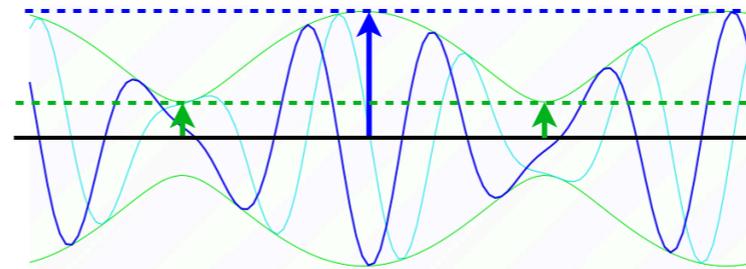
$$\omega_{\Sigma} = (\omega_{\rightarrow} + \omega_{\leftarrow})/2$$

$$\omega_{\Delta} = (\omega_{\rightarrow} - \omega_{\leftarrow})/2$$

Now consider amplitude mean $A_{\Sigma} = (A_{\rightarrow} + A_{\leftarrow})/2$ and amplitude half-difference $A_{\Delta} = (A_{\rightarrow} - A_{\leftarrow})/2$.

Detailed wave motion depends on standing-wave-ratio **SWR** or the inverse standing-wave-quotient **SWQ**.

$$\frac{\text{Envelope-Min.}}{\text{Envelope-Max.}} = \text{SWR} = \frac{(A_{\rightarrow} - A_{\leftarrow})}{(A_{\rightarrow} + A_{\leftarrow})}$$



$$\text{SWQ} = \frac{(A_{\rightarrow} + A_{\leftarrow})}{(A_{\rightarrow} - A_{\leftarrow})} = \frac{1}{\text{SWR}}$$

They're analogous to group velocity $V_{\text{group}} < c$ frequency ratios and inverse phase velocity $V_{\text{phase}} > c$ ratios.

$$V_{\text{group}} = \frac{\omega_{\Delta}}{k_{\Delta}} = \frac{(\omega_{\rightarrow} - \omega_{\leftarrow})}{(k_{\rightarrow} - k_{\leftarrow})} = c \frac{(\omega_{\rightarrow} - \omega_{\leftarrow})}{(\omega_{\rightarrow} + \omega_{\leftarrow})}$$

$$V_{\text{phase}} = \frac{\omega_{\Sigma}}{k_{\Sigma}} = \frac{(\omega_{\rightarrow} + \omega_{\leftarrow})}{(k_{\rightarrow} + k_{\leftarrow})} = c \frac{(\omega_{\rightarrow} + \omega_{\leftarrow})}{(\omega_{\rightarrow} - \omega_{\leftarrow})}$$

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$$\frac{V_{\text{phase}}}{c} = \frac{\omega_{\Sigma}}{ck_{\Sigma}} = \frac{(\omega_{\rightarrow} + \omega_{\leftarrow})}{c(k_{\rightarrow} + k_{\leftarrow})} = \frac{(\omega_{\rightarrow} + \omega_{\leftarrow})}{(\omega_{\rightarrow} - \omega_{\leftarrow})} = \frac{c}{V_{\text{group}}}$$

$$\frac{V_{\text{group}}}{c} = \frac{c}{V_{\text{phase}}} \quad \text{is analogous to:} \quad \text{SWR} = \frac{1}{\text{SWQ}}$$

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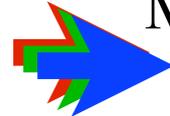
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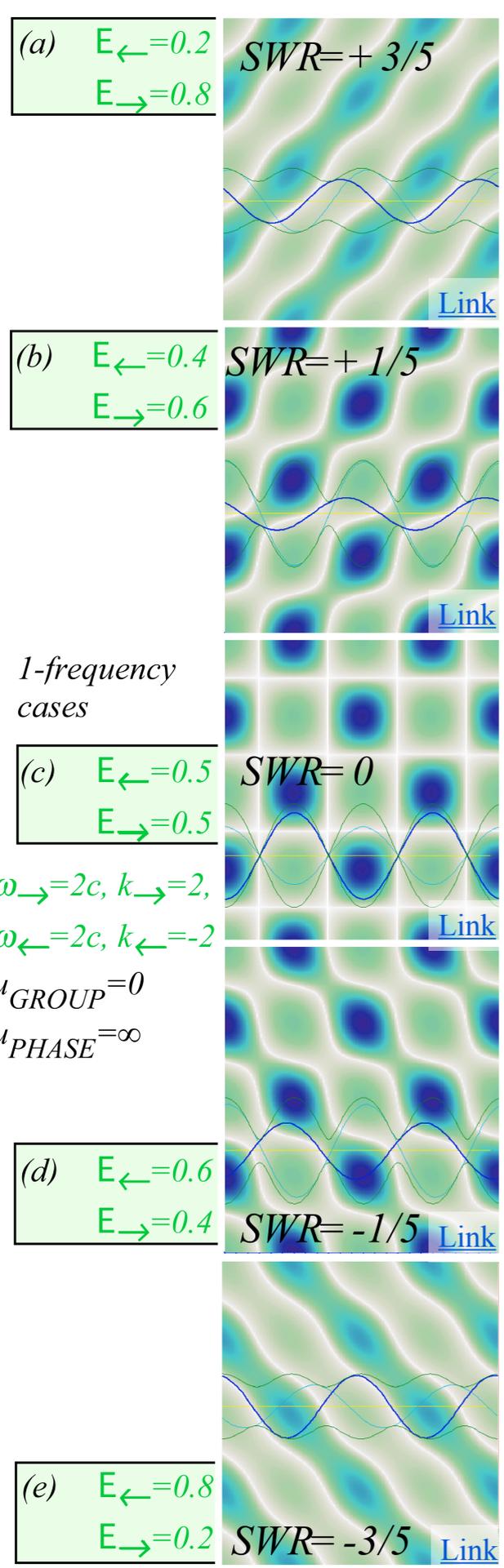
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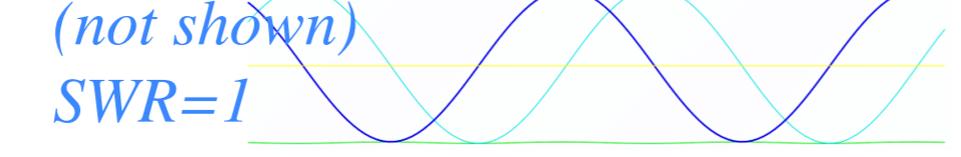
1-frequency cases

$$\omega_{\rightarrow}=2c, k_{\rightarrow}=2,$$

$$\omega_{\leftarrow}=2c, k_{\leftarrow}=-2$$

$$u_{GROUP}=0$$

$$u_{PHASE}=\infty$$



SWR=+3/5

SWR=+1/5

SWR=0

SWR=-1/5

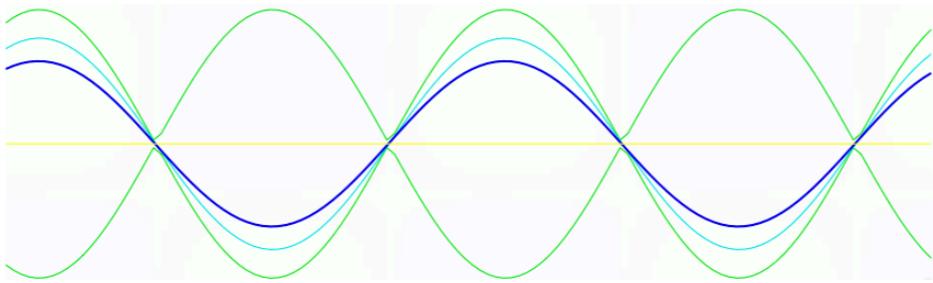
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(not shown in (x,ct) plots)

SWR=1

Two extremes for Standing Wave Ratio

SWR=0

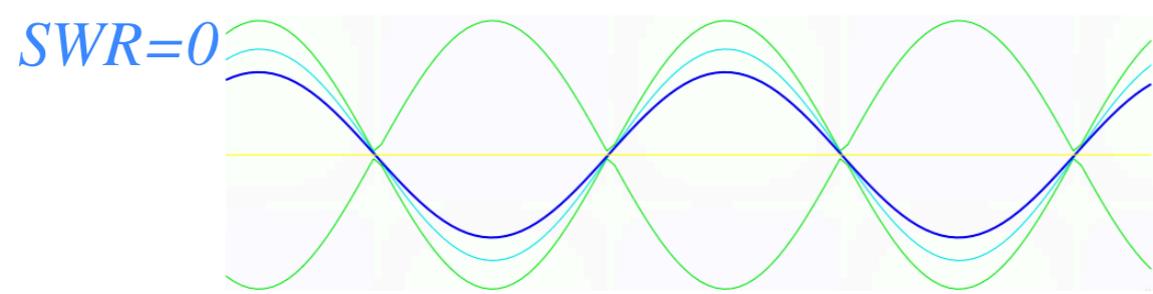
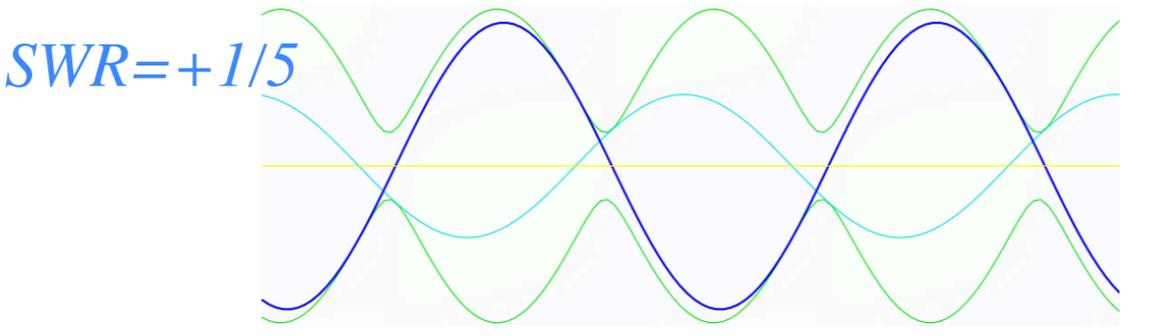
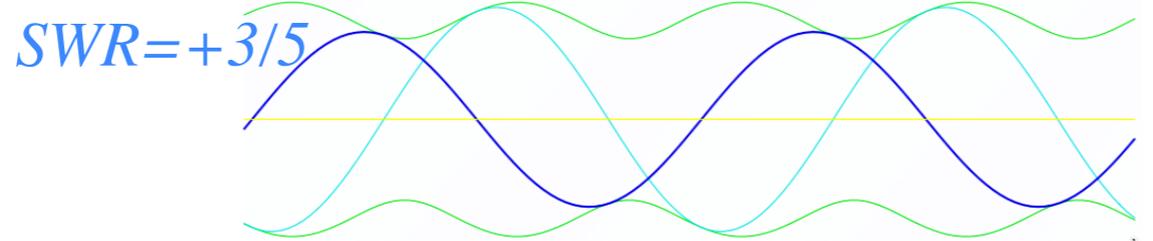
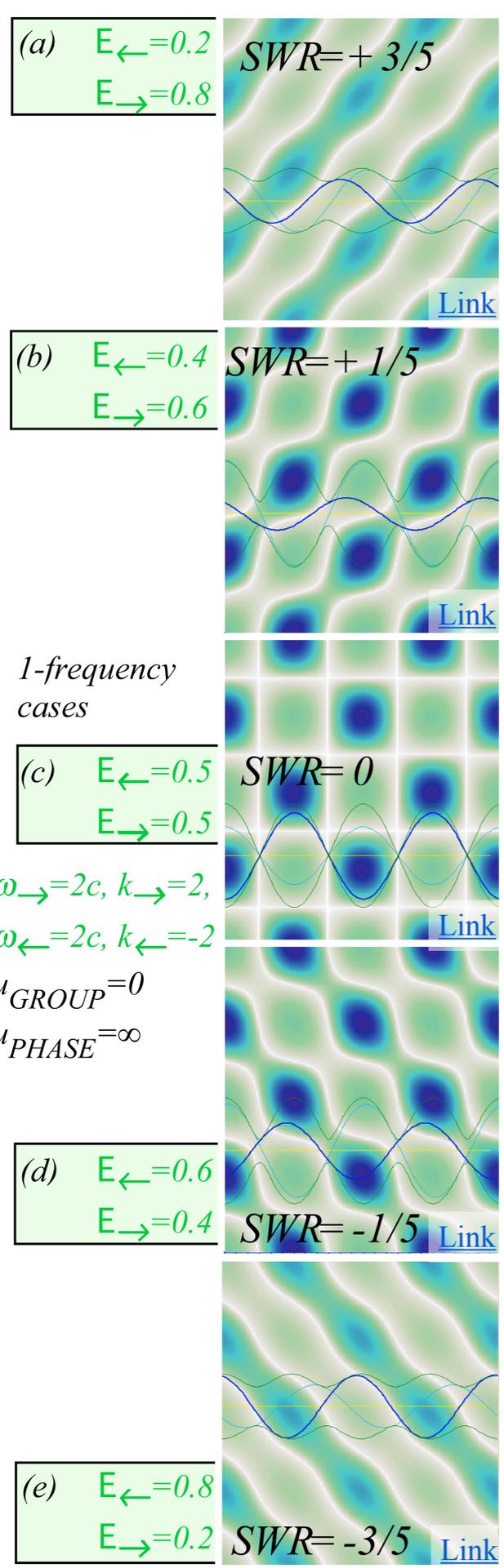


BohrIt Web Simulations Links
(embedded in corners)

from: Fig. 4.5.2
QTforCA
Unit 2 Ch.4

from: Fig. 8.6.3
CMwBang!
Unit 8 Ch.6

Fig. 6.1 Monochromatic (1-frequency) 2-CW wave space-time patterns.



(not shown in (x,ct) plots)

SWR=1
Two extremes for Standing
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SWR=0

1-frequency
cases

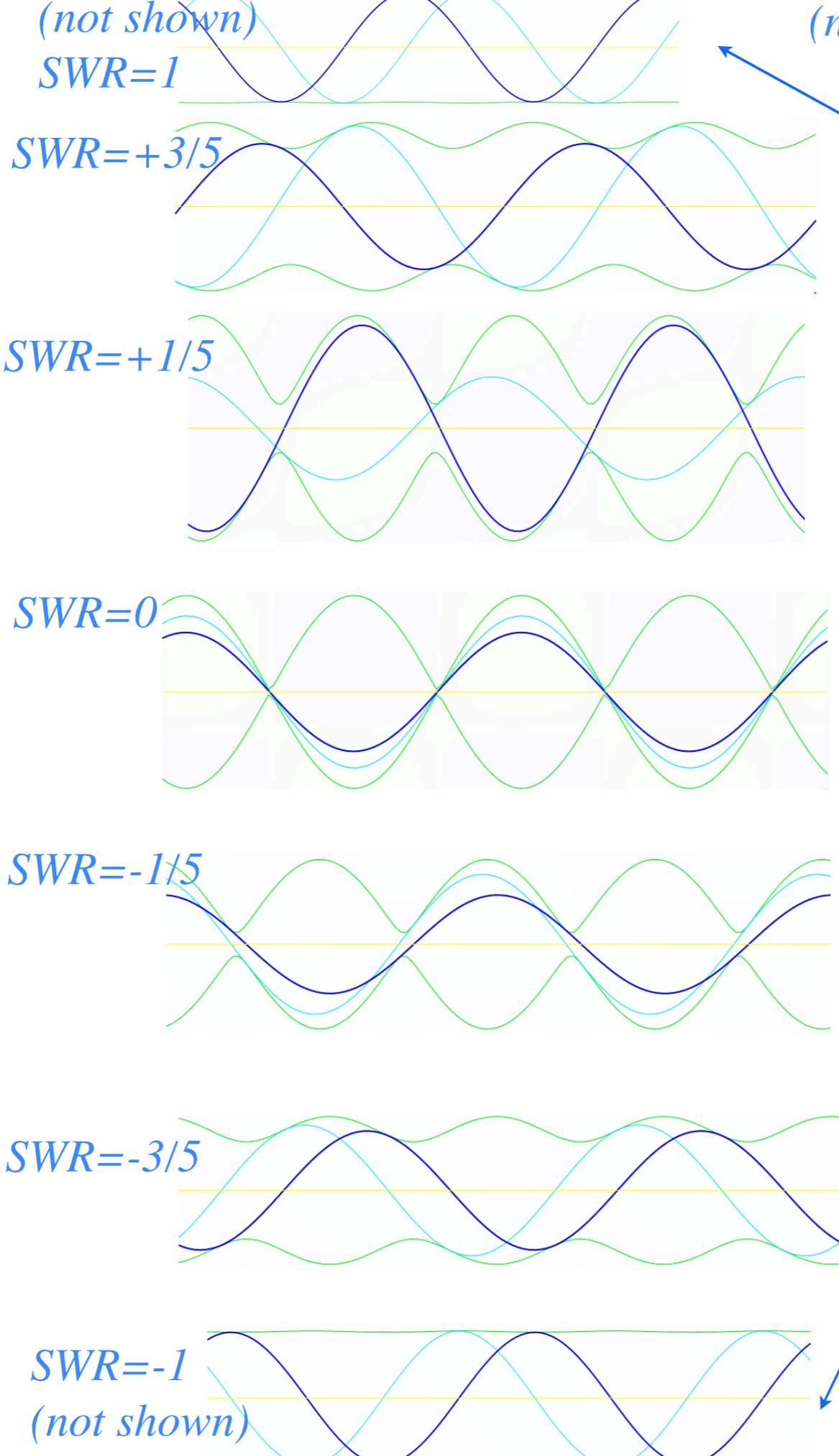
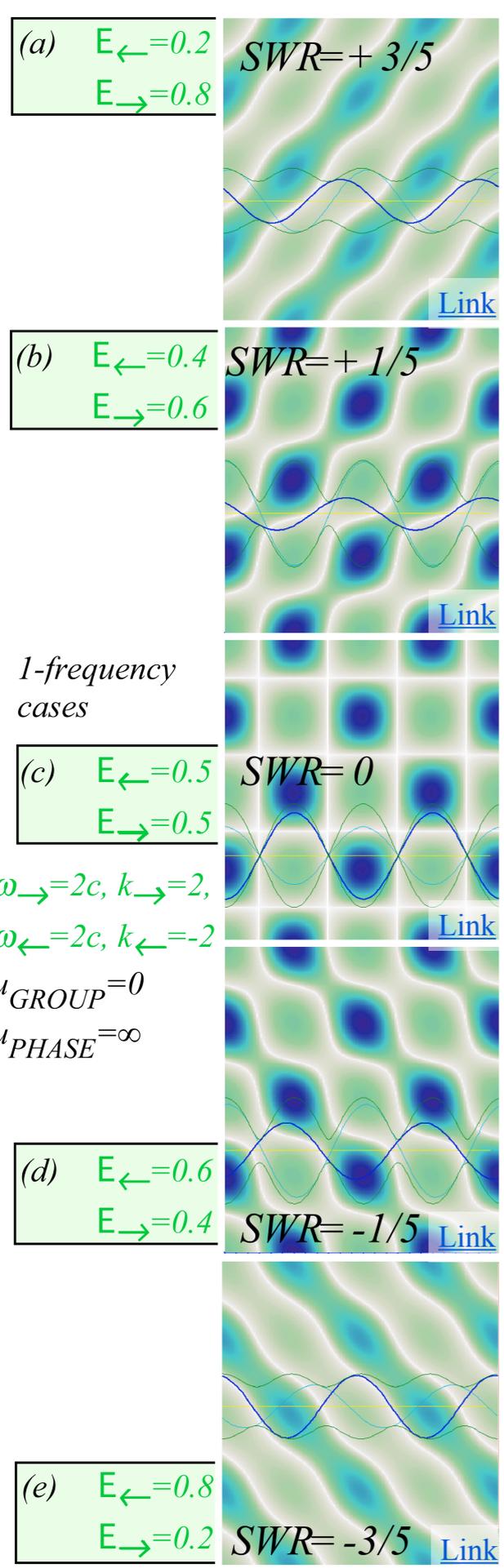
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(not shown in (x,ct) plots)

$SWR=1$
Two extremes for Standing Wave Ratio
 $SWR=0$...and
 $SWR=-1$

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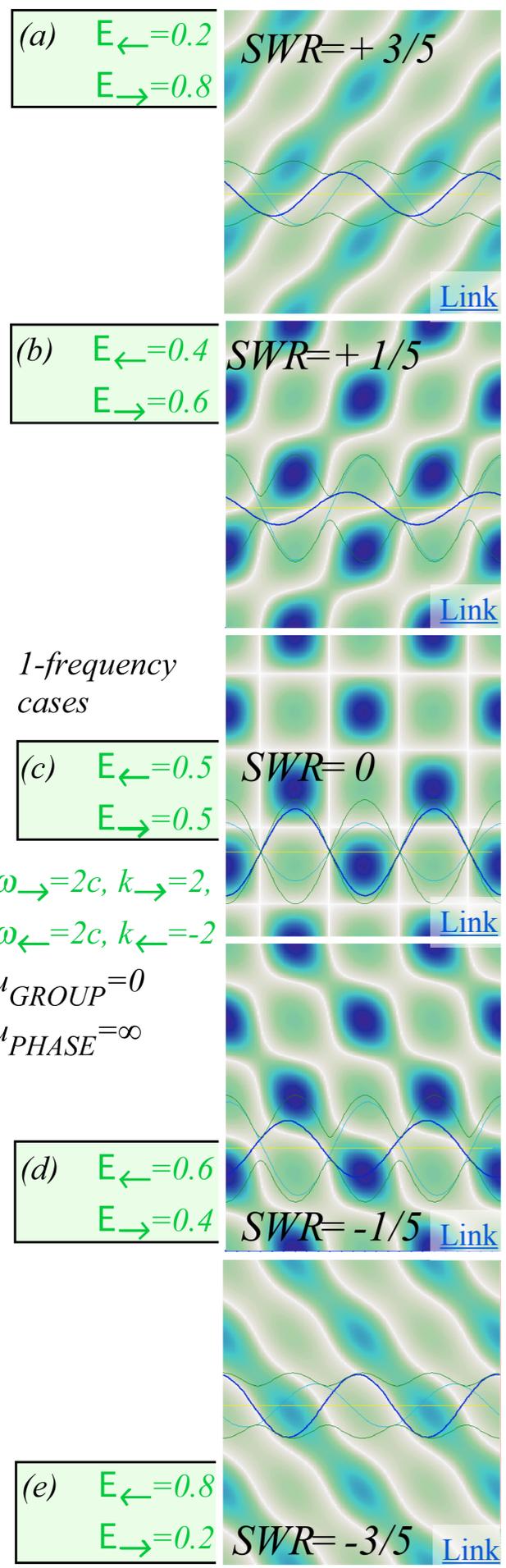


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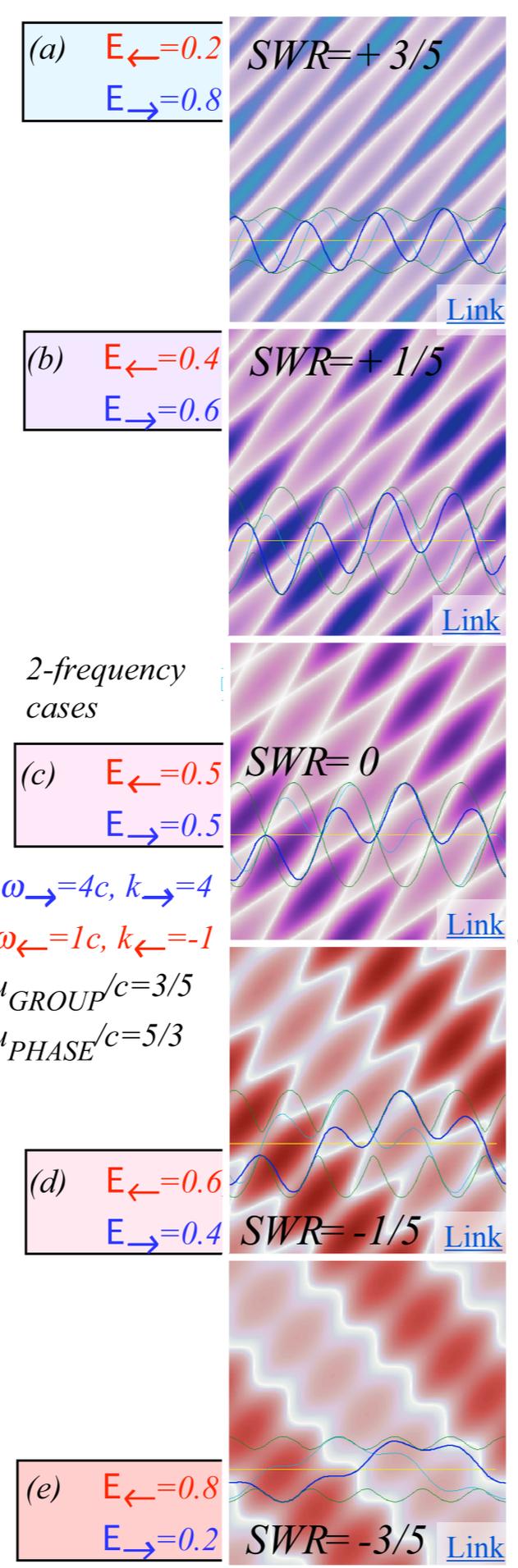
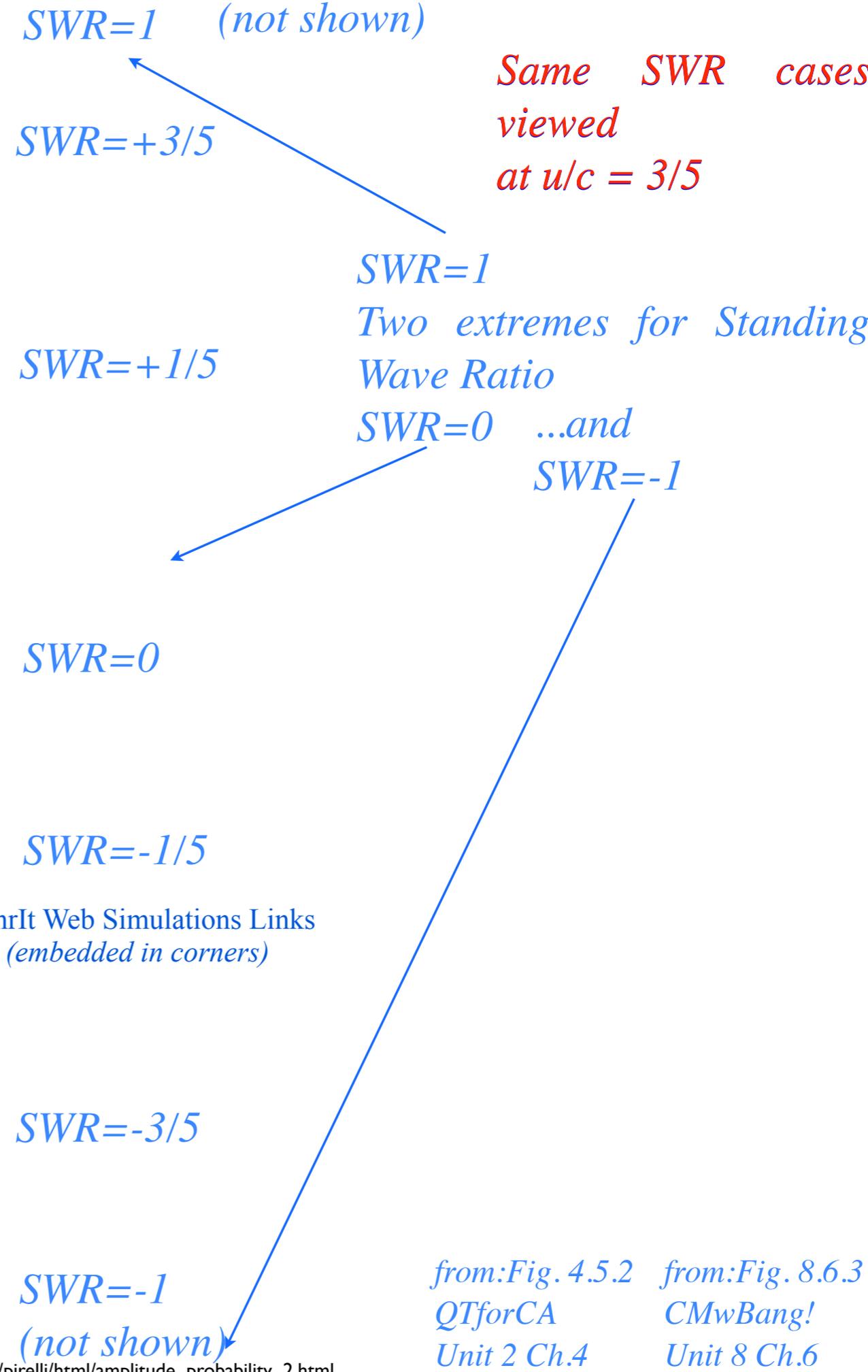


Fig. 6.2 Dichromatic (2-frequency) 2-CW wave space-time patterns.

www.uark.edu/ua/pirelli/html/amplitude_probability_2.html



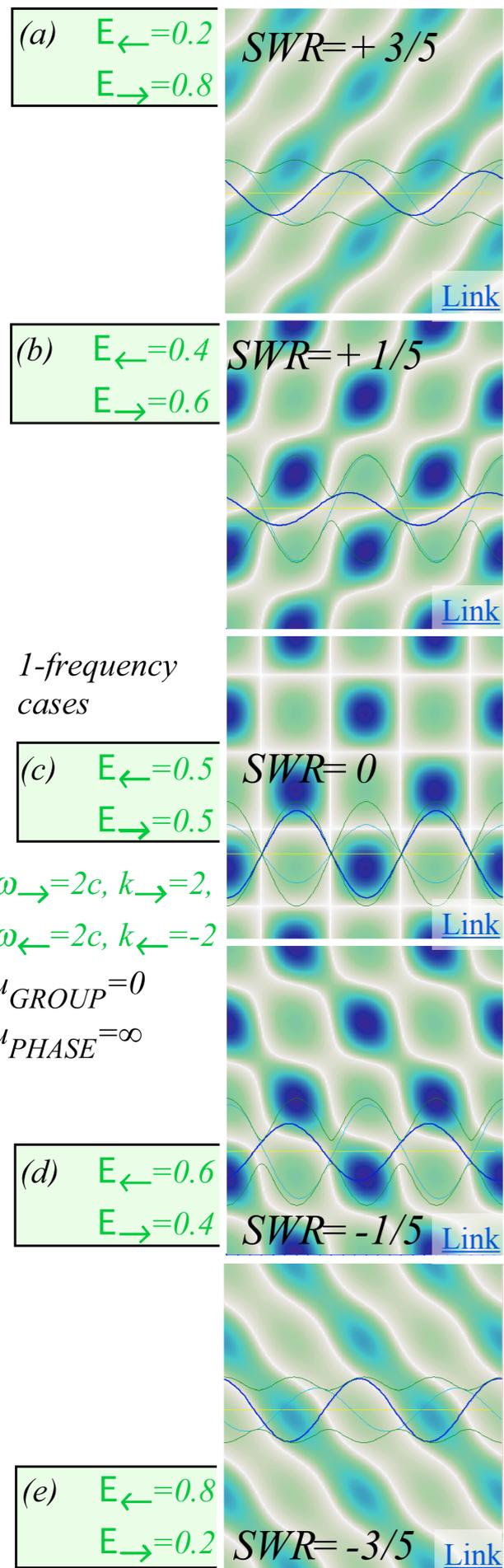


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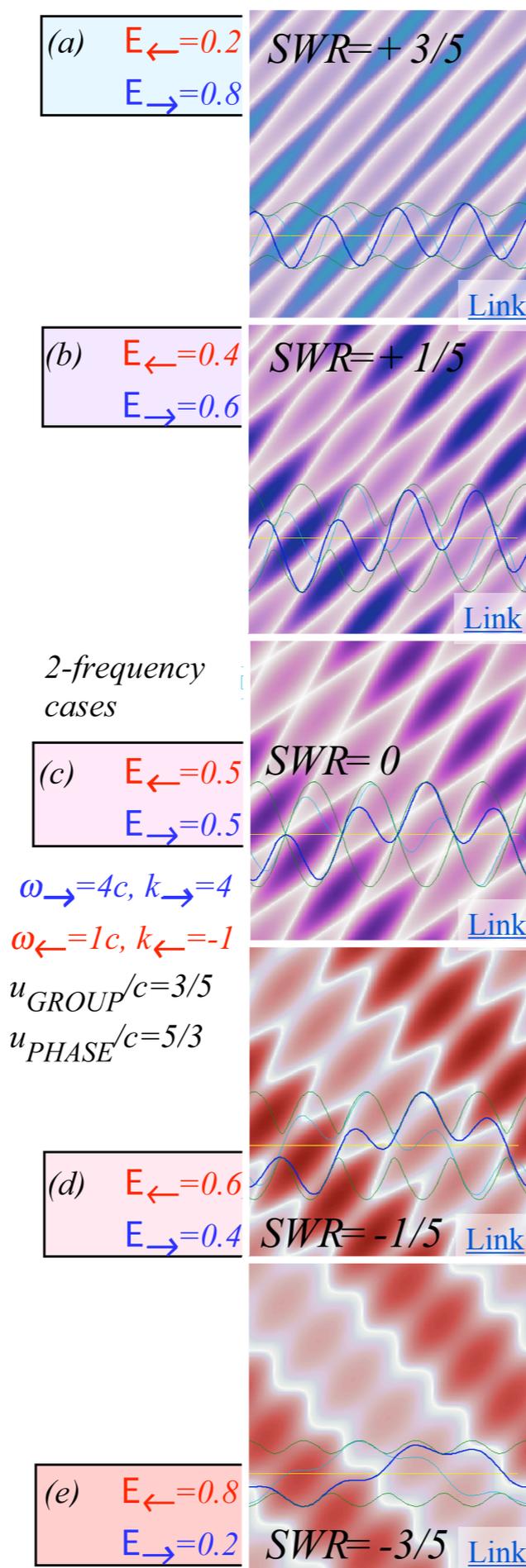


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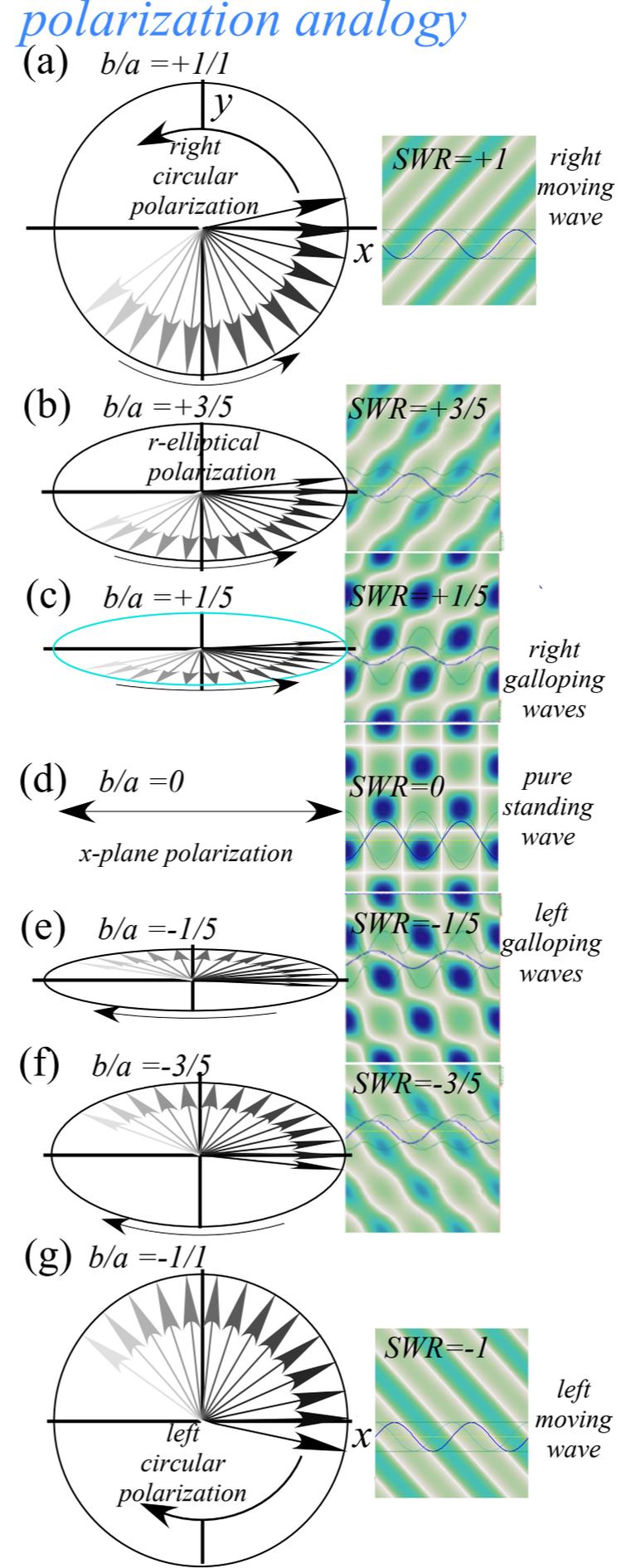


Fig. 6.3 (a-g) Elliptic polarization ellipses relate to galloping waves in Fig. 6.1. (h-i) Kepler anomalies.

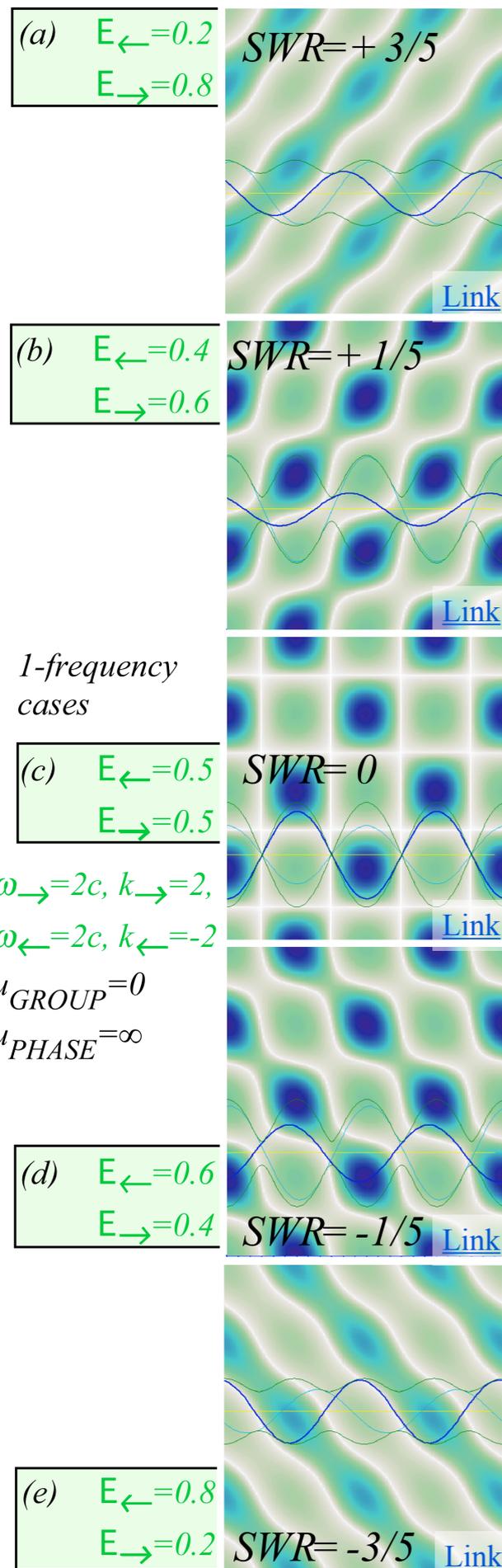


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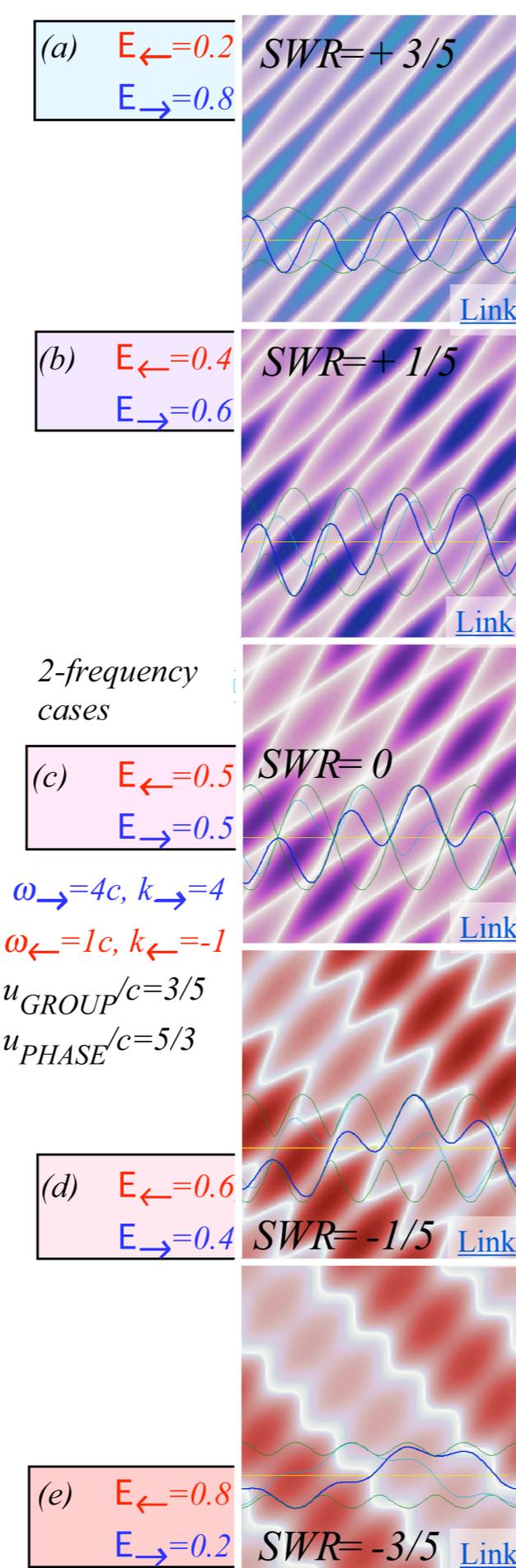


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<http://www.uark.edu/ua/modphys/markup/BoxItWeb.html?AU2=1.0&BU2=0.0&CU2=0.0&DU2=1.0&xInitial=1.0&yInitial=0.0&pxInitial=0.0&pyInitial=0.4&wantBoxLines=1>

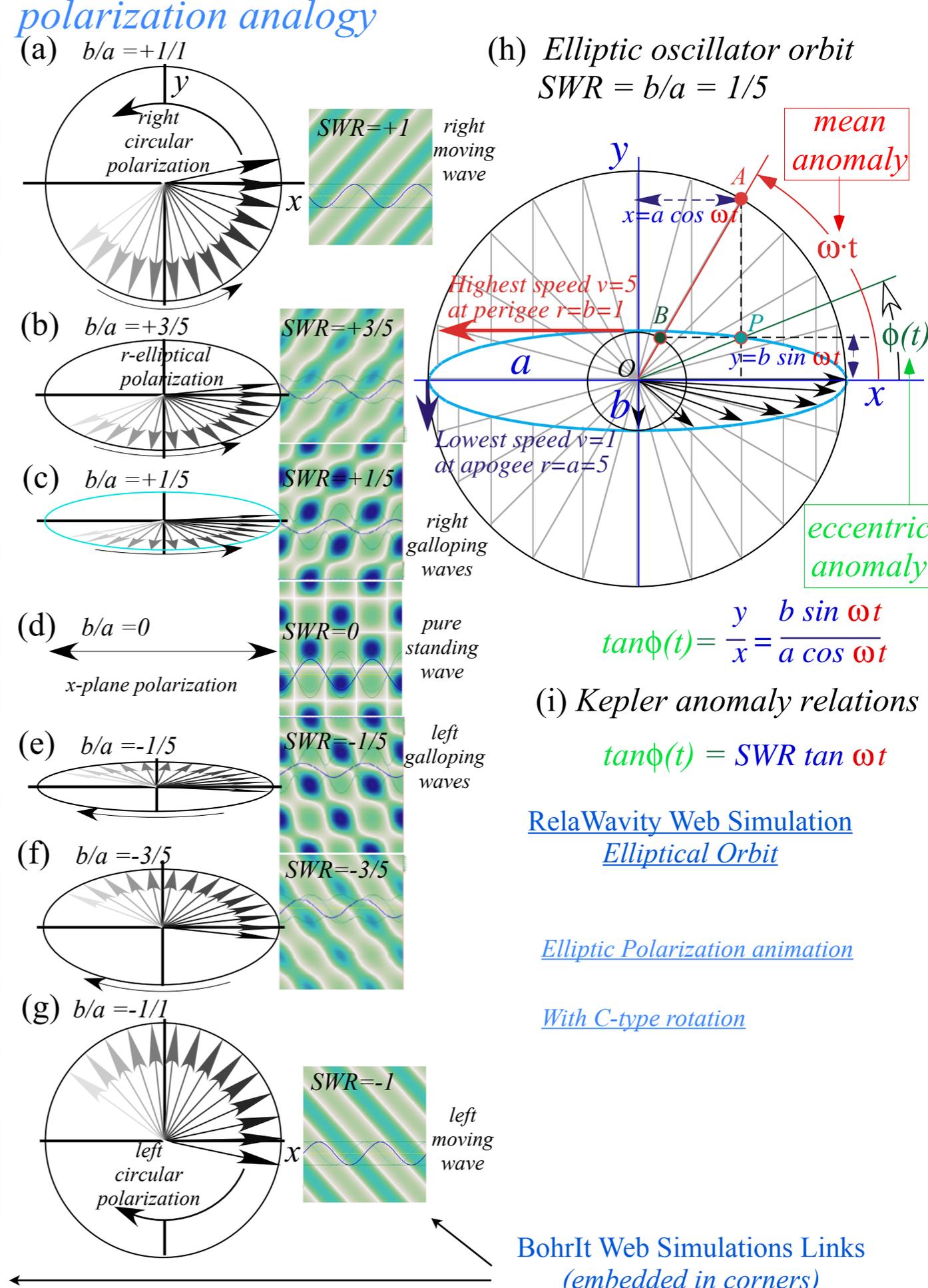


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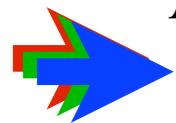
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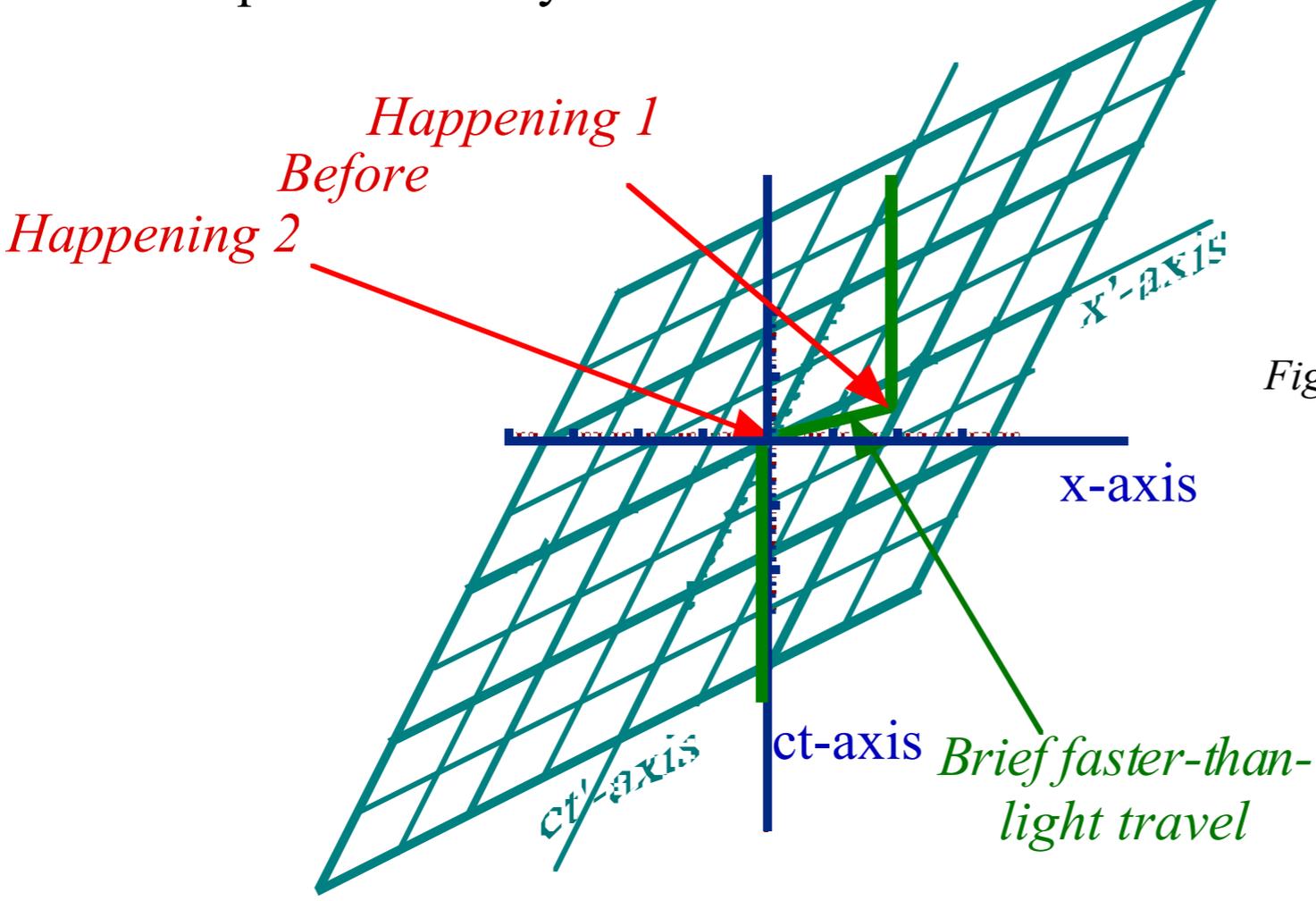


Fig. 2.B.10 Lighthouse plot of two Happenings

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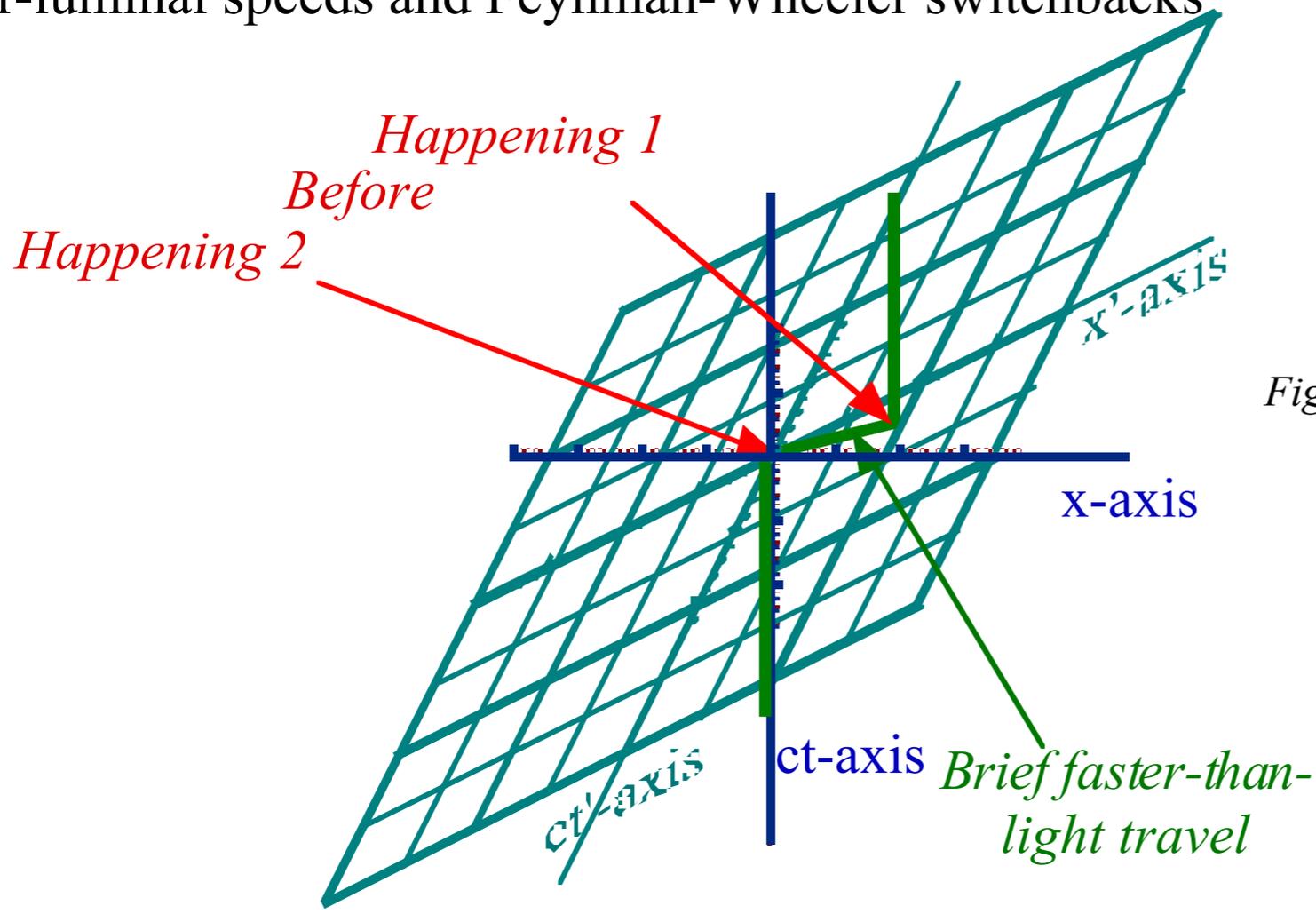


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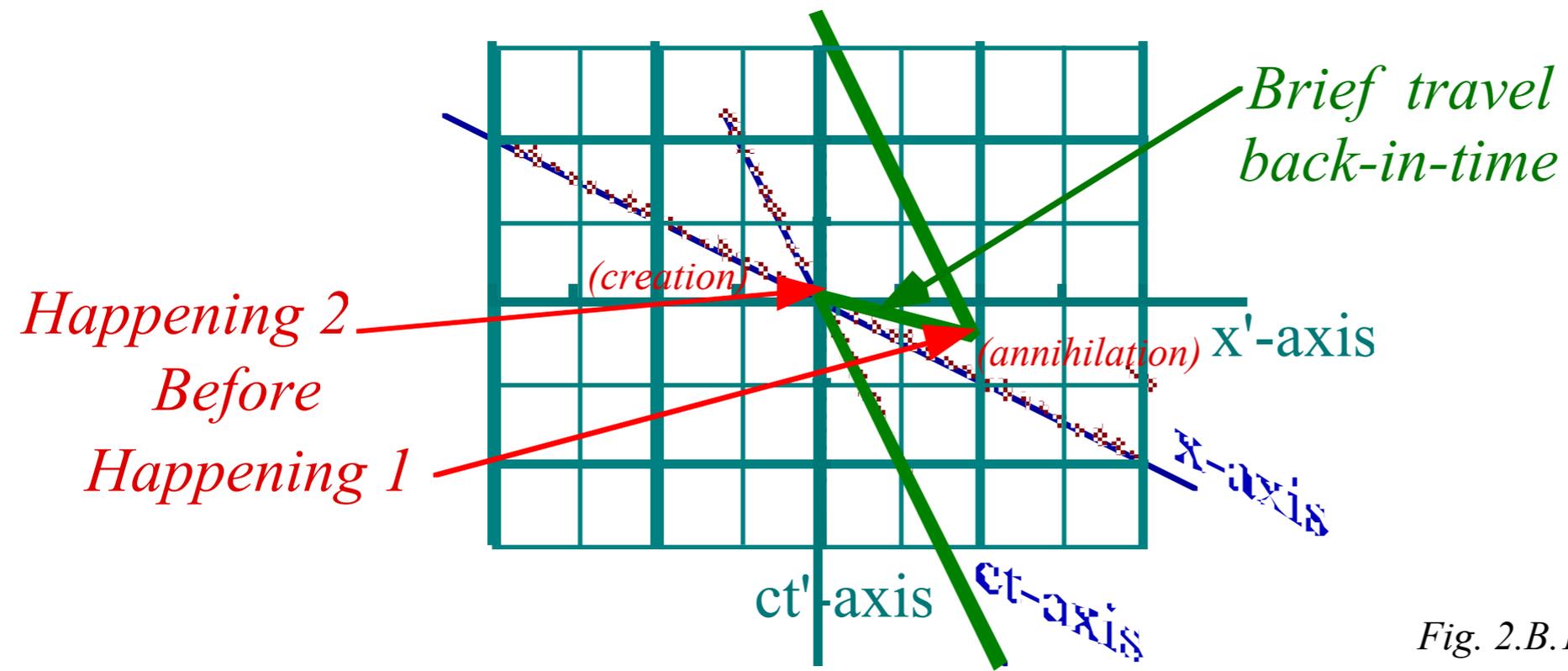


Fig. 2.B.11 Ship plot of two Happenings

Waves that go back in time - The Feynman-Wheeler Switchback

Minkowski Zero-Grids are Spacetime Switchbacks for $-u_{GROUP} < SWR < 0$

$\omega_{\rightarrow} = 4c$	$\omega_{\leftarrow} = 1c$
$k_{\rightarrow} = 4$	$k_{\leftarrow} = -1$
$u_{GROUP} = c3/5$	$u_{PHASE} = c5/3$

Group zero speed limit

$$\frac{u_{GROUP} + SWR}{1 + u_{GROUP} \cdot \frac{SWR}{c^2}} = 5c/11$$

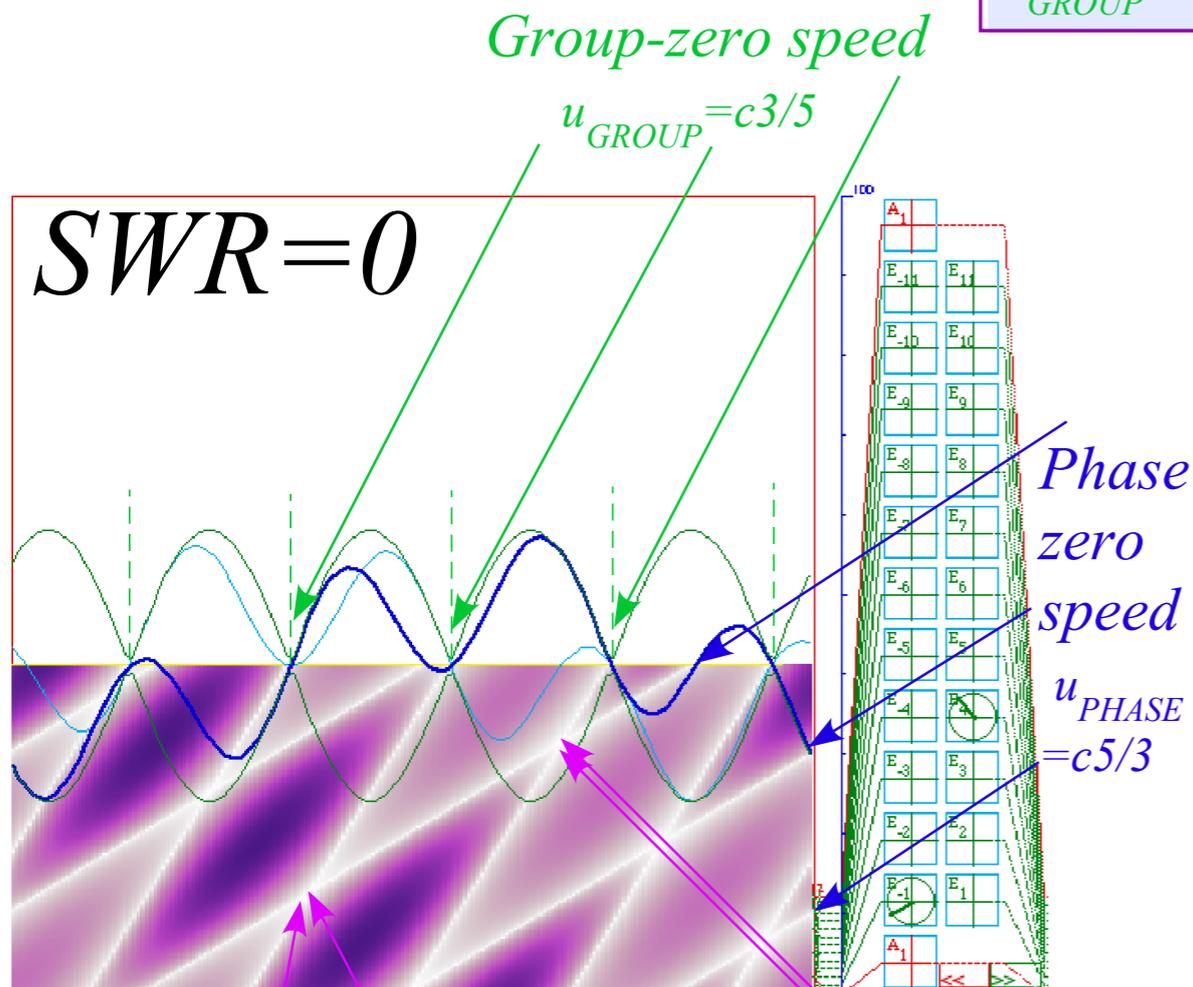
$$\frac{\frac{3}{5} + \frac{-1}{5}}{1 + \frac{3 \cdot -1}{5 \cdot 5}} = \frac{\frac{2}{5}}{\frac{22}{25}} = \frac{5}{11}$$

Phase "anti-zero" going "back-in-time"

Phase zero speed limit

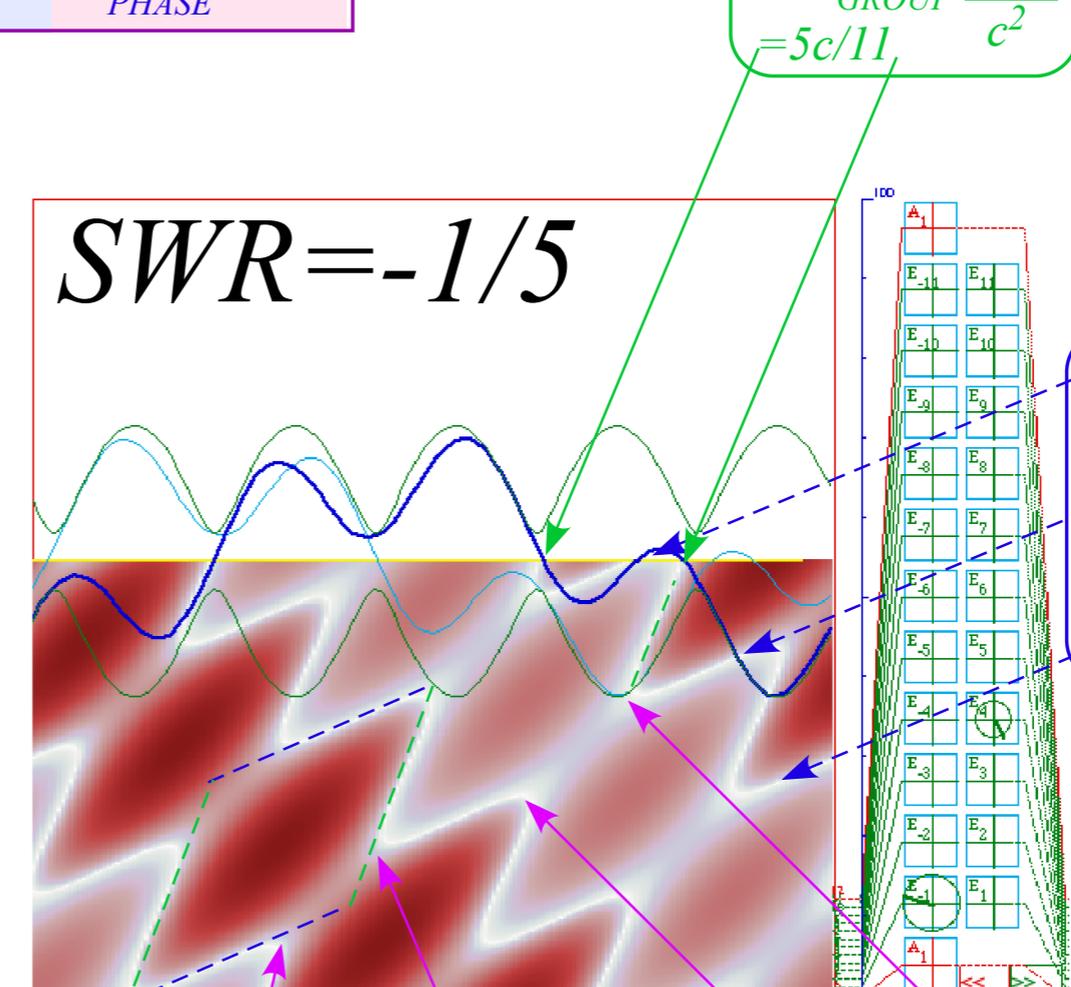
$$\frac{u_{PHASE} + SWR}{1 + u_{PHASE} \cdot \frac{SWR}{c^2}} = 11c/5$$

$$\frac{\frac{5}{3} + \frac{-1}{5}}{1 + \frac{5 \cdot -1}{3 \cdot 5}} = \frac{\frac{22}{15}}{\frac{10}{25}} = \frac{11}{5}$$



$E_{\leftarrow} = 0.5, E_{\rightarrow} = 0.5$

Wave zero-anti-zero annihilation and creation occur together at the same spacetime point for $SWR=0$



$E_{\leftarrow} = 0.6, E_{\rightarrow} = 0.4$

Wave zero-anti-zero annihilation and creation occur separately at different spacetime points for $-u_{GROUP} < SWR < 0$

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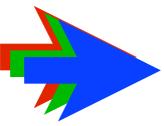
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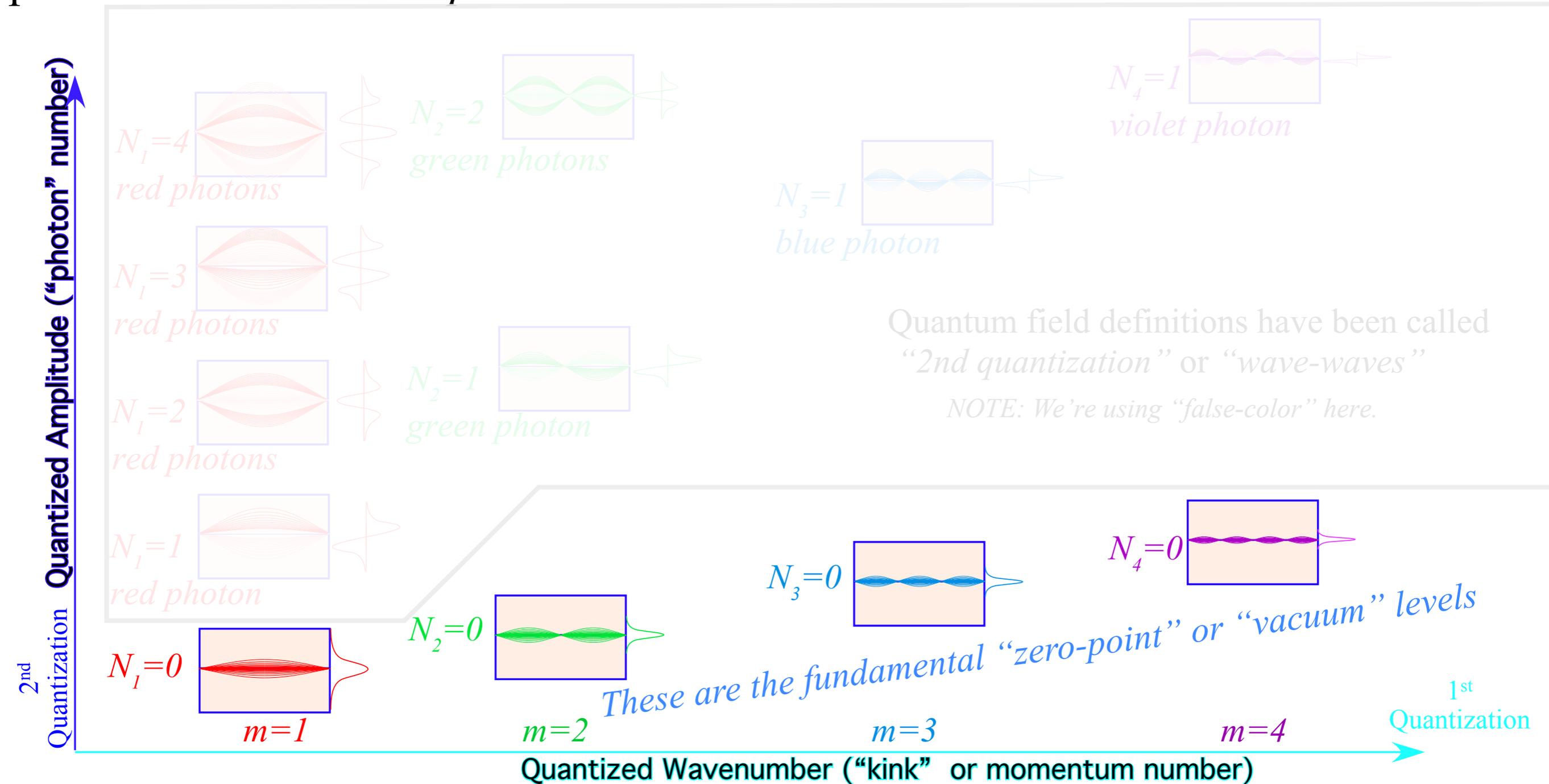
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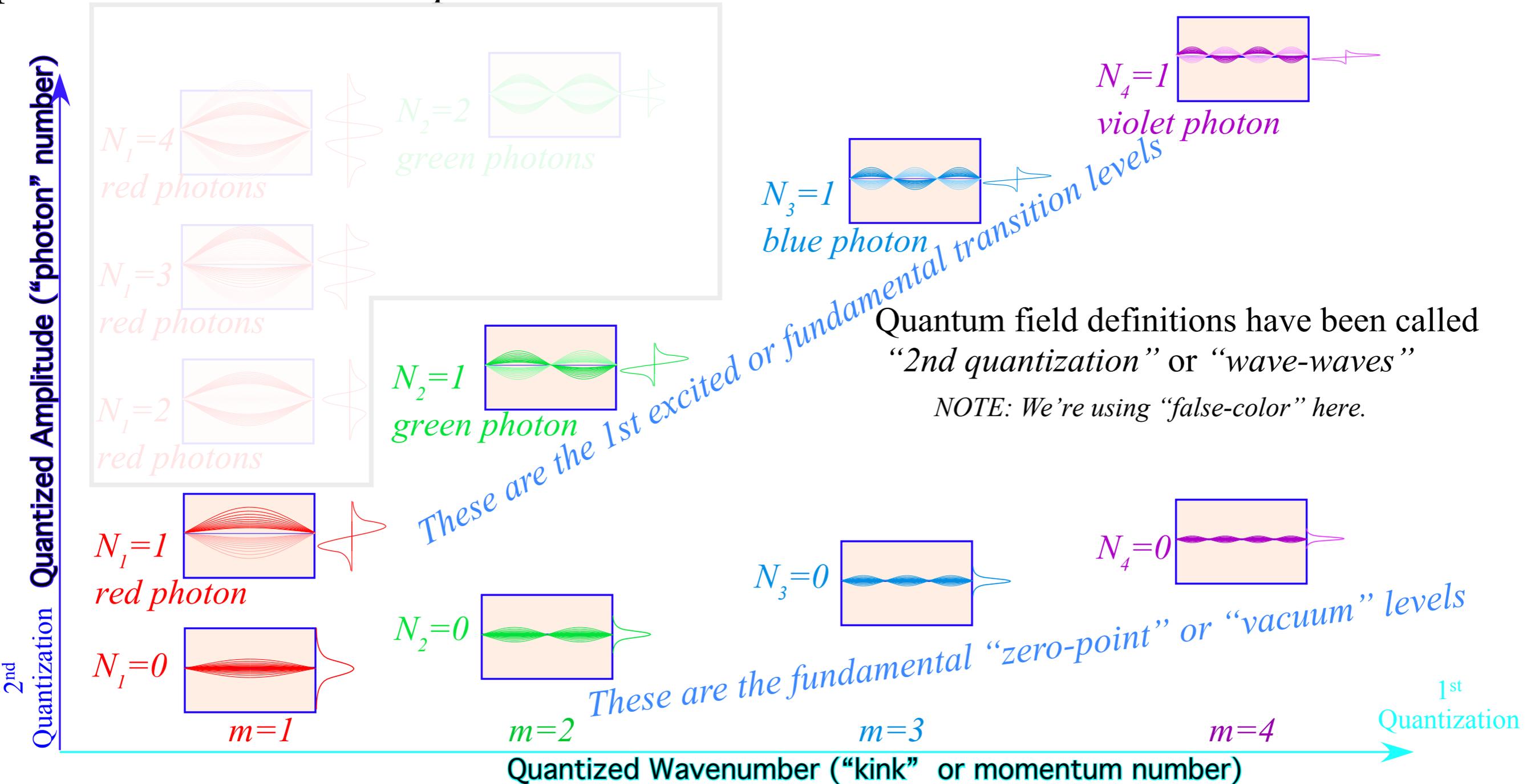
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Planck’s relation $E=Nh\nu$ began as a tentative axiom to explain low-T light. Then he tried to disavow it! Einstein picked it up in his 1905 paper. Since then its use has grown enormously and continues to amaze, amuse (or bewilder) all who study it. A current view is that it represents the *quantization* of optical field *amplitude*. We picture this below as *N-photon* wave states for each box-mode of *m* wave kinks.



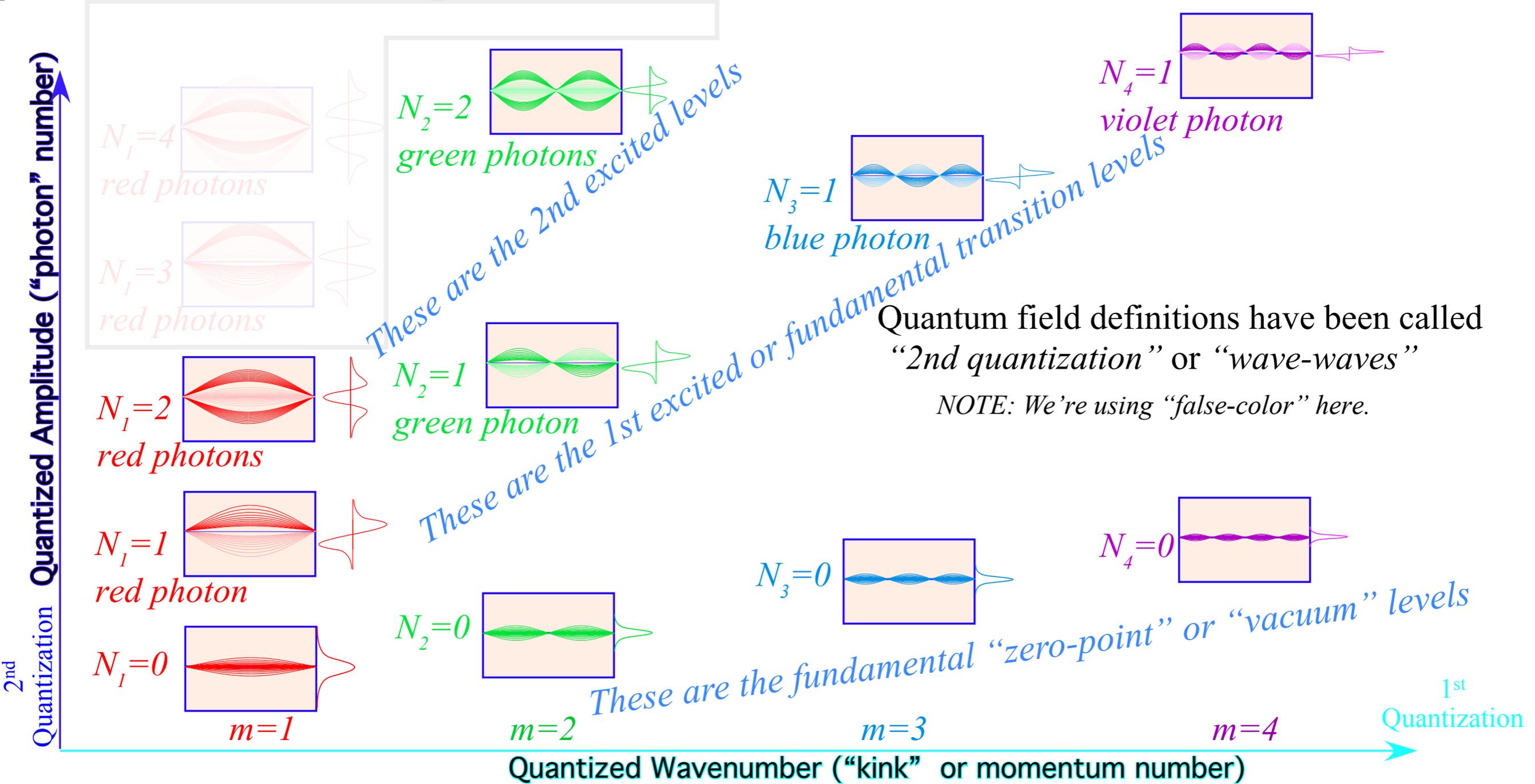
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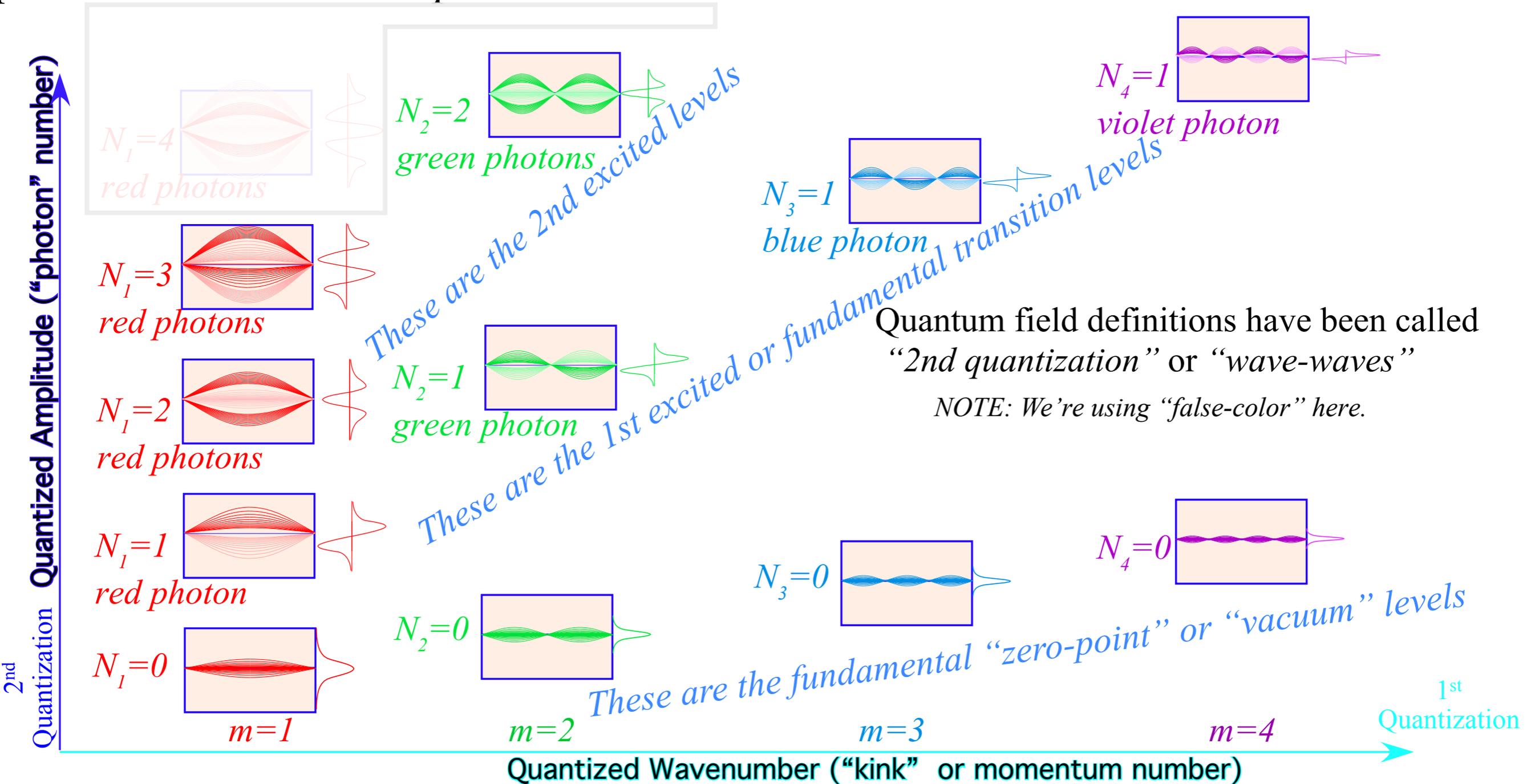
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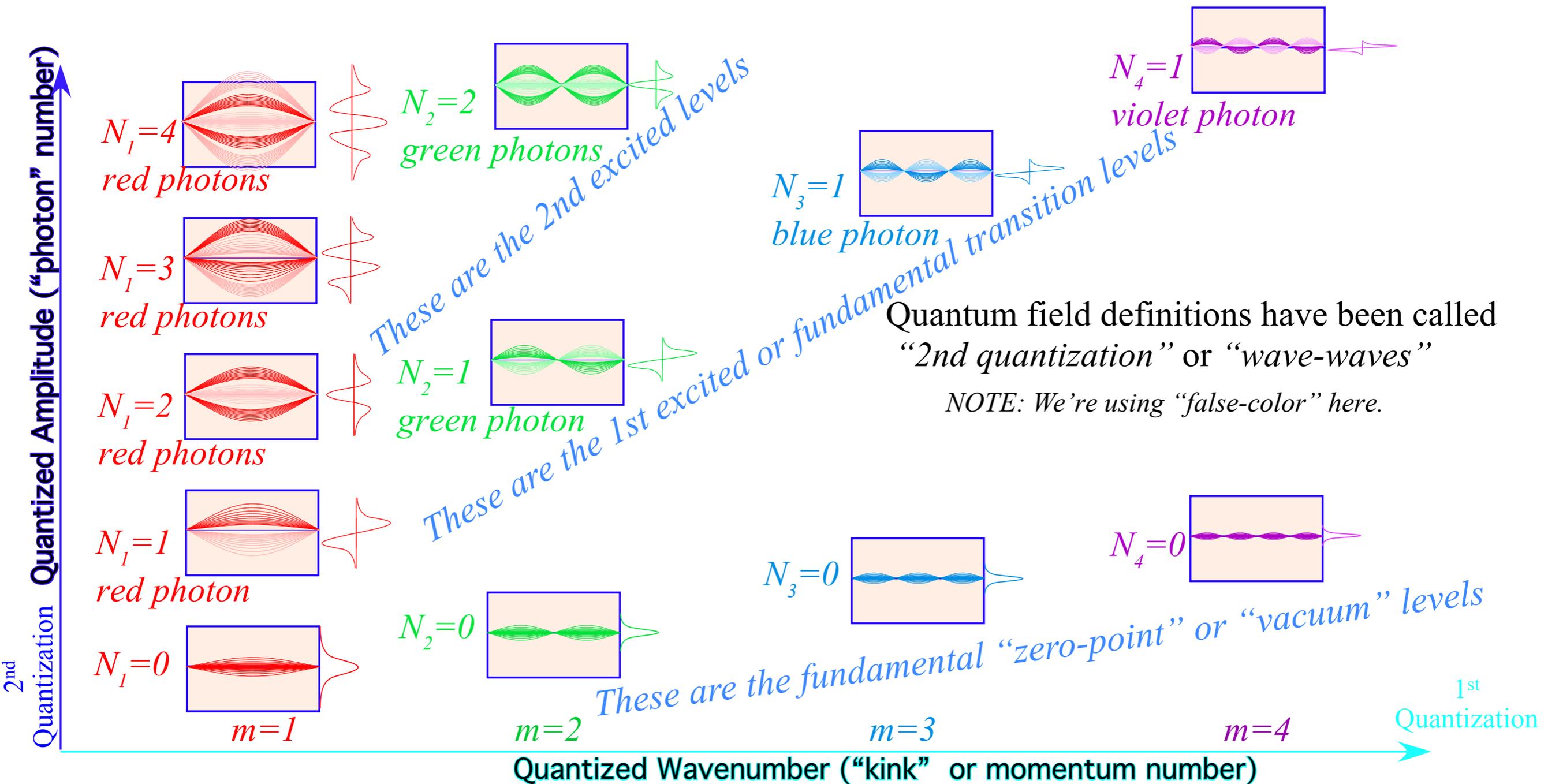
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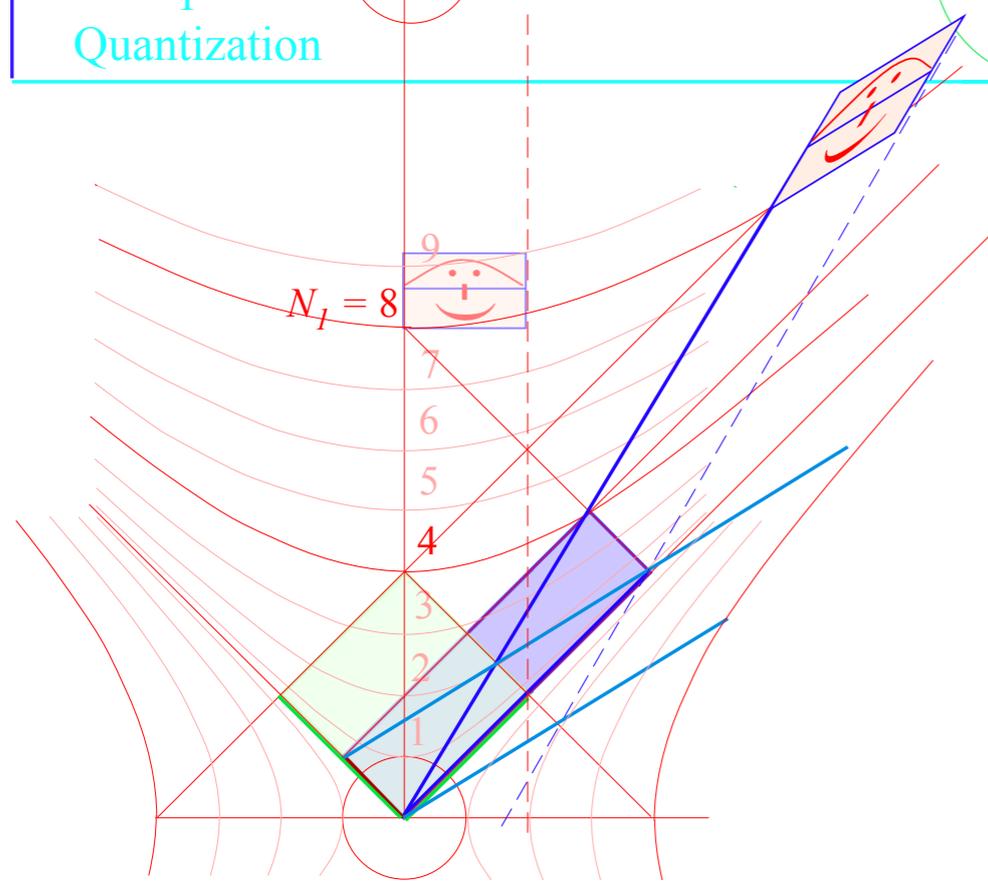
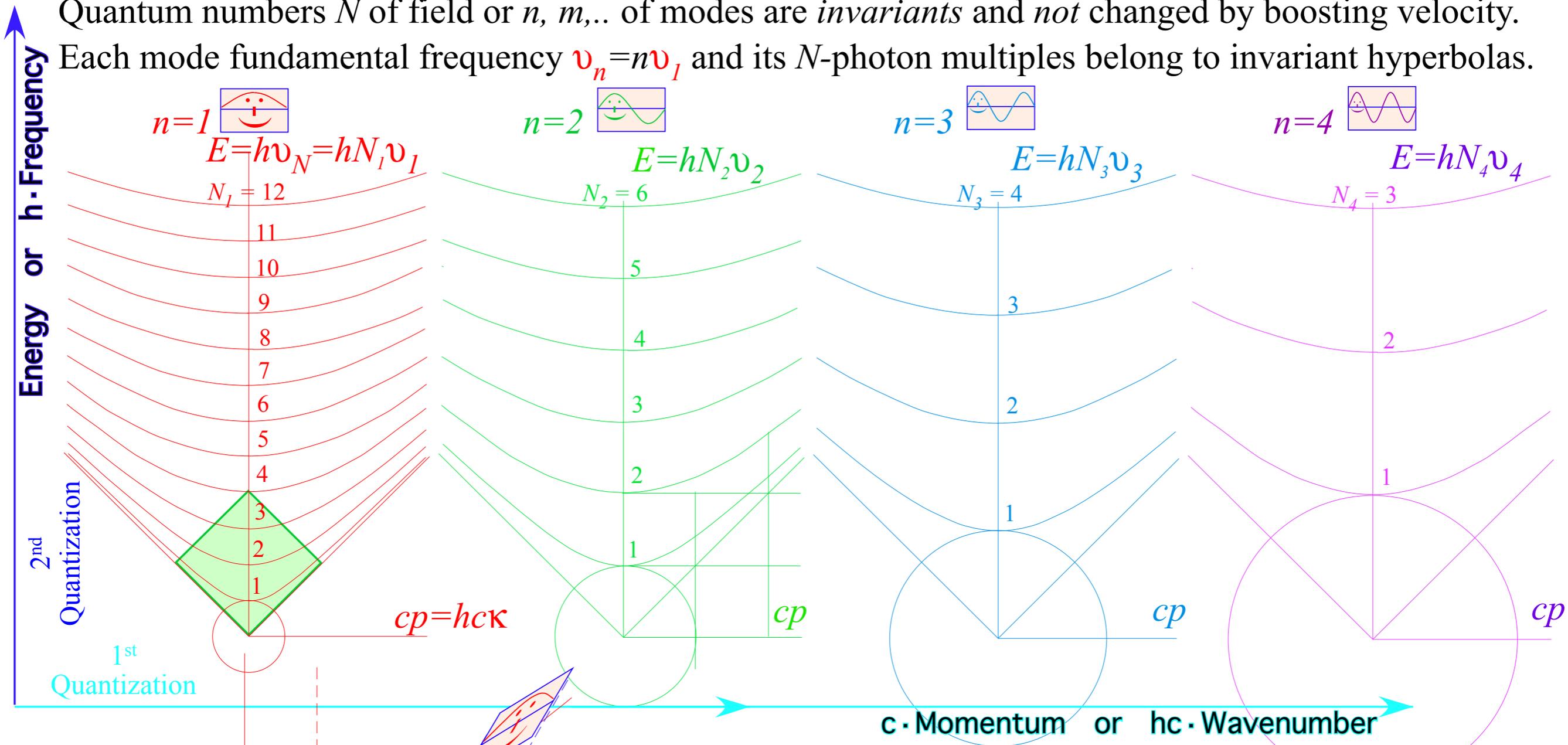


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Quantum numbers N of field or n, m, \dots of modes are *invariants* and *not* changed by boosting velocity. Each mode fundamental frequency $\nu_n = n\nu_1$ and its N -photon multiples belong to invariant hyperbolas.



Boosted observers see distorted frequencies and lengths, but will agree on the *numbers* n and N of mode *nodes* and *photons*.

This is how light waves can “fake” some of the properties of classical “things” such as *invariance* or *object permanence*.

It takes at least *TWO CW*'s to achieve such invariance. One CW is not enough and cannot have non-zero invariant N . Invariance is an *interference* effect that needs at least *two-to-tango*!

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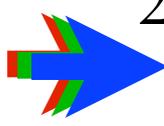
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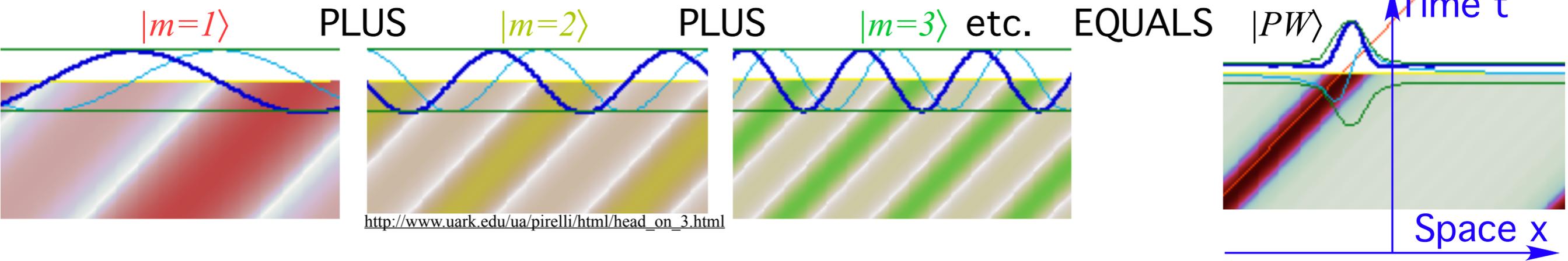
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Coherent States: Oscillator Amplitude Packets analogous to Wave Packets

We saw how adding *CW*'s (*Continuous Waves* $m=1,2,3\dots$) can make *PW* (*Pulse Wave*) or *WP* (*Wave Packet*) that is more like a classical "thing" with more *localization* in space x and time t .



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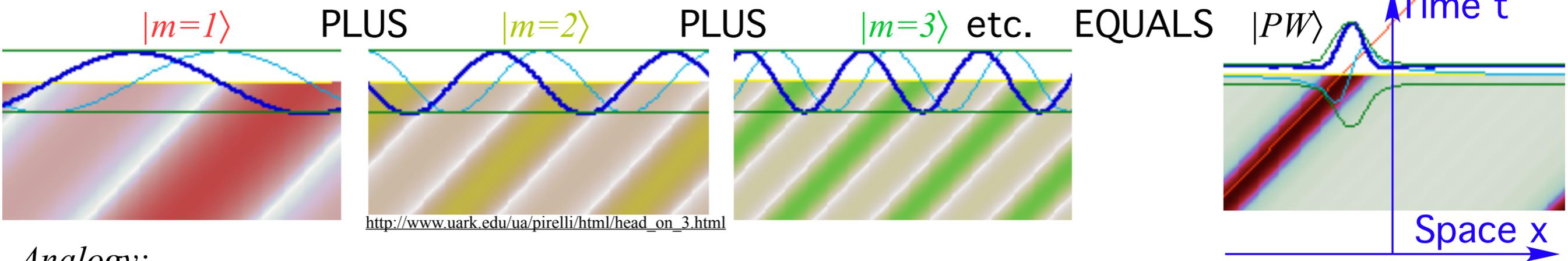
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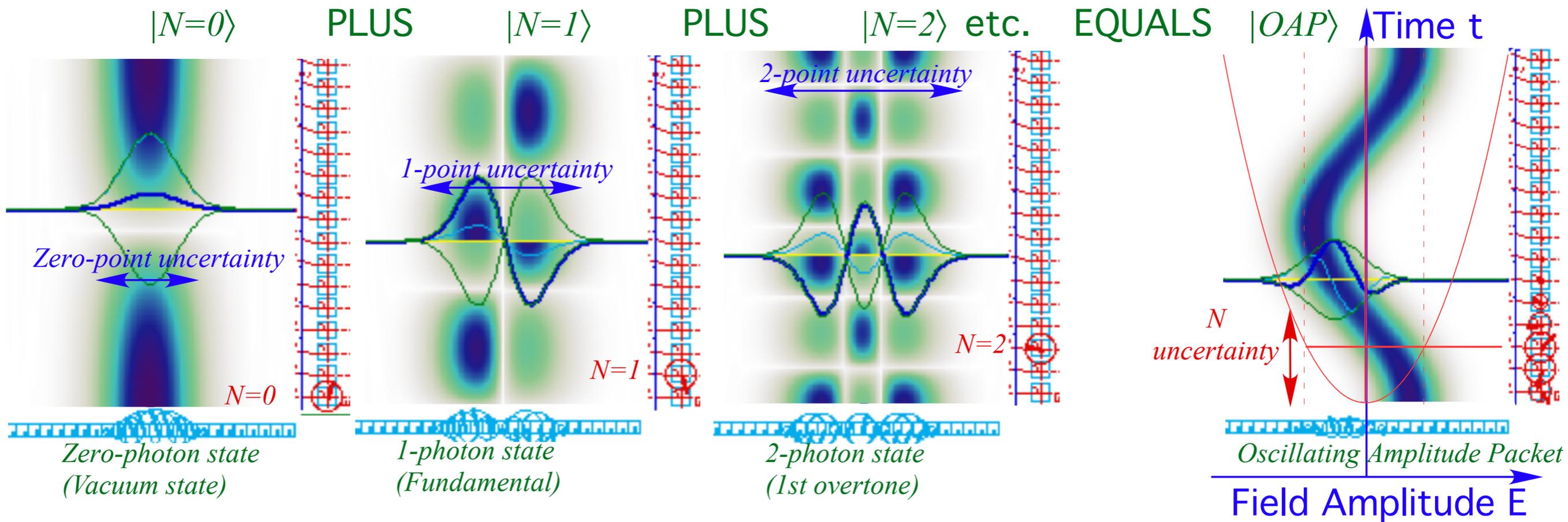
Coherent States: Oscillator Amplitude Packets analogous to Wave Packets

We saw how adding CW's (Continuous Waves $m=1,2,3\dots$) can make PW (Pulse Wave) or WP (Wave Packet) that is more like a classical "thing" with more localization in space x and time t .



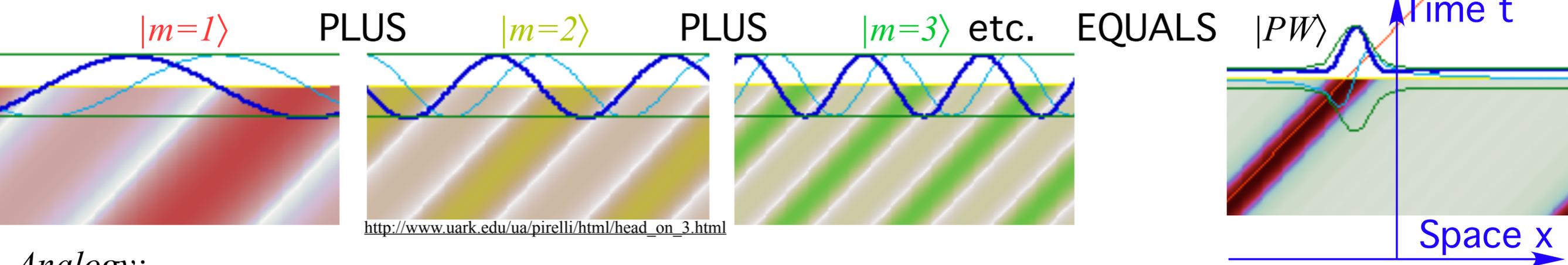
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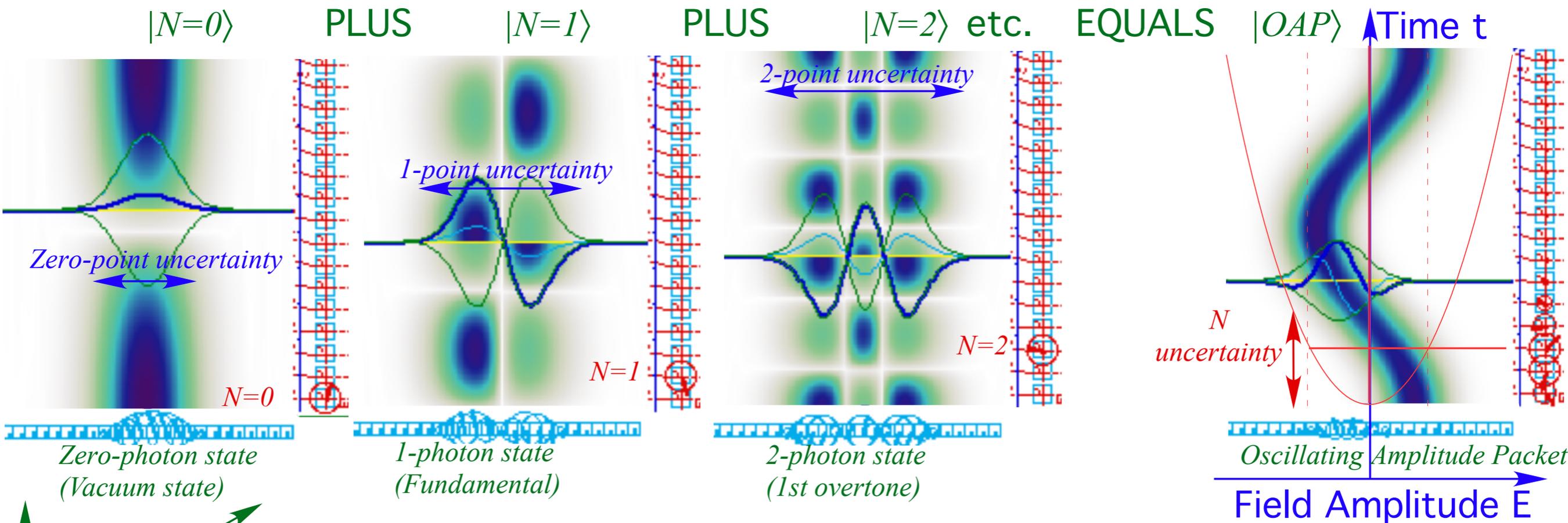
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Pure photon states have localized (certain) N but delocalized (uncertain) amplitude and phase.
 OAP states have delocalized (uncertain) N but more localized (certain) amplitude and phase.

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Effect on group velocity (*None*) and phase velocity (*Absurd*)

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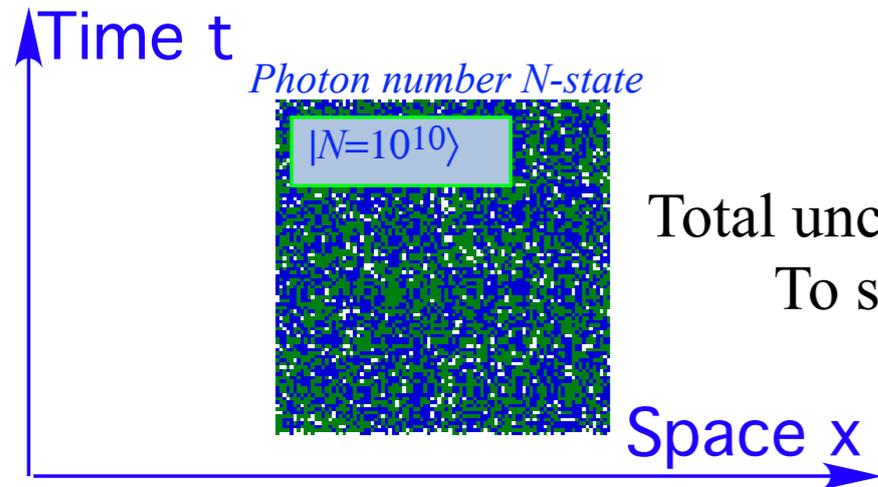
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Coherent States(contd.) *Spacetime wave grid is impossible without coherent states*

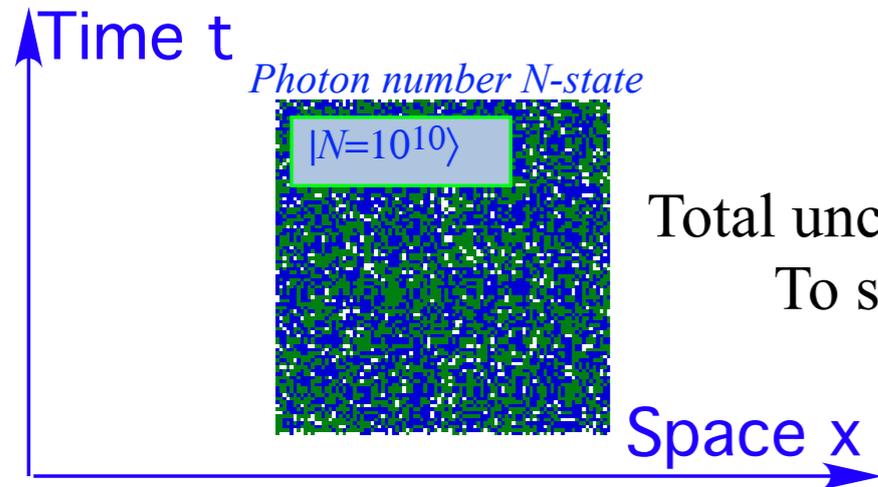
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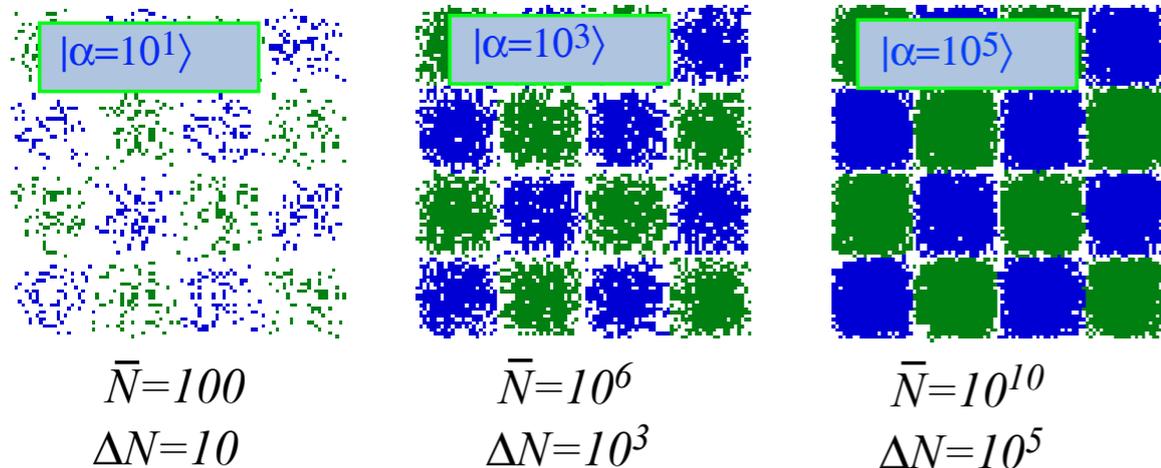
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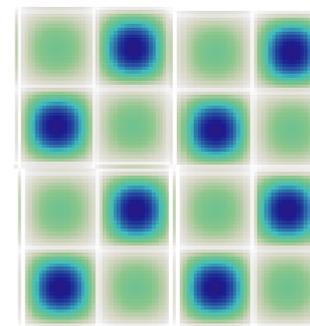
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Coherent- α -states are defined by continuous amplitude-packet parameter α whose square is average photon number $\bar{N}=|\alpha|^2$. Coherent-states make better spacetime coordinates for larger $\bar{N}=|\alpha|^2$.

Quantum field coherent α -states



Classical limit



Coherent-state uncertainty in photon number (and mass) varies with amplitude parameter $\Delta N \sim \alpha \sim \sqrt{N}$ so a coherent state with $\bar{N}=|\alpha|^2 = 10^6$ only has a 1-in-1000 uncertainty $\Delta N \sim \alpha \sim \sqrt{N} = 1000$.

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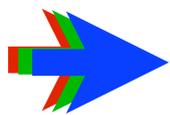
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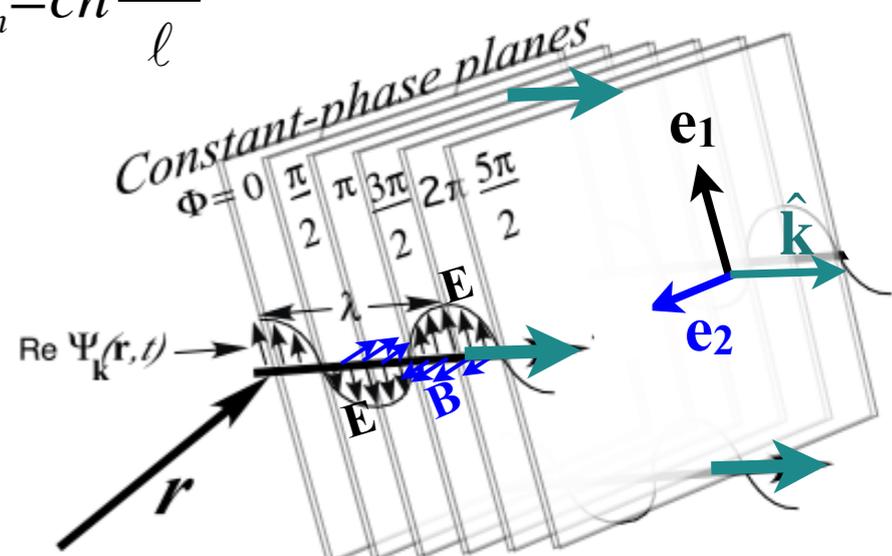
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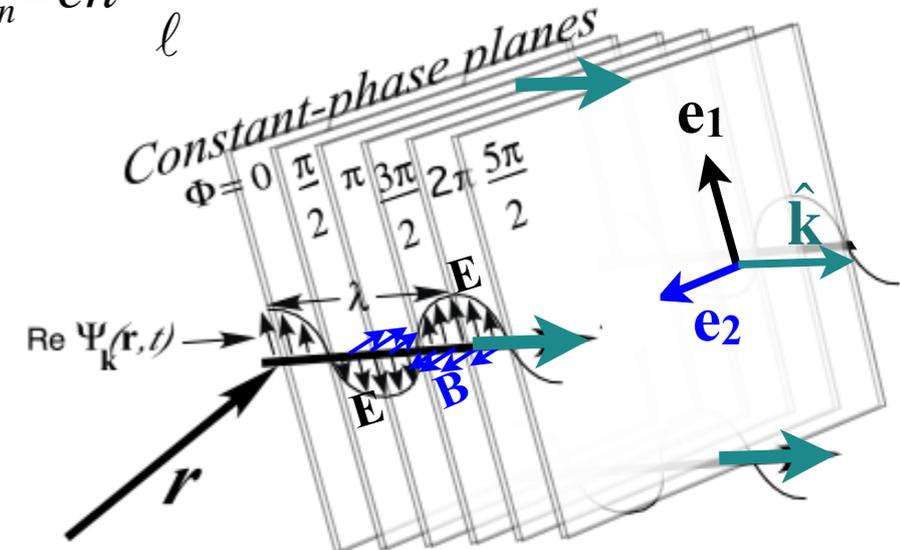
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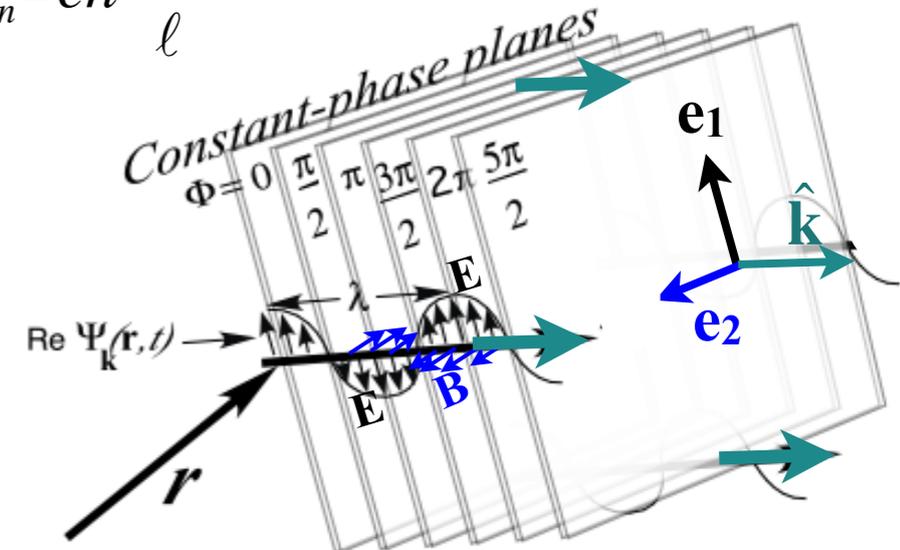
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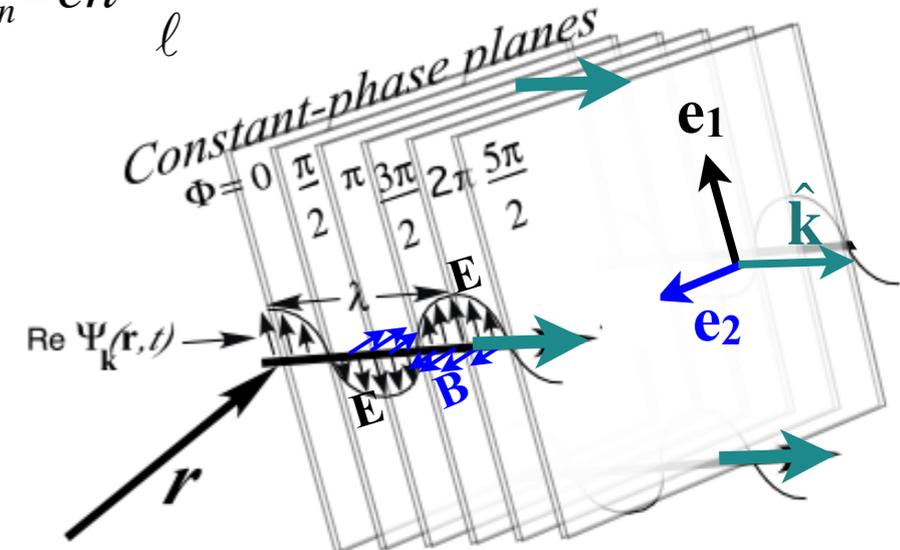
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QUANTITY
photon counts per sec.

QUALITY
cycles per sec.

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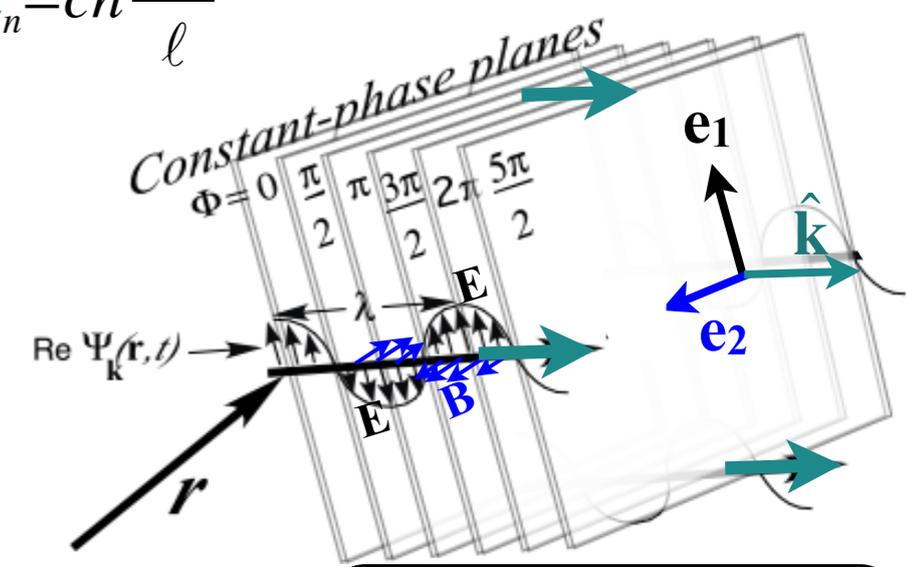
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QUANTITY
photon counts per sec.

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cycles per sec.

N and ω are both frequencies for quantum wave so \mathbf{E} -field has Doppler $e^{\pm\rho}$ -shifts just like N and ω

Now we see how Planck's $E_N(\omega) = \hbar N \omega$ axiom has the classical quadratic $\omega^2 |\mathbf{A}|^2$ oscillator energy

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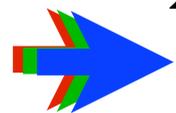
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$$\langle U \rangle = \frac{\hbar N \omega}{V} = \frac{\epsilon_0}{2} \omega^2 |\mathbf{A}|^2 = \frac{\epsilon_0}{2} |\mathbf{E}|^2$$

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$$\langle U \rangle = \frac{\hbar N \omega}{V} = \frac{\epsilon_0}{2} \omega^2 |\mathbf{A}|^2 = \frac{\epsilon_0}{2} |\mathbf{E}|^2$$

$$\frac{\langle U \rangle}{\hbar \omega} = \frac{N}{V} = \frac{\epsilon_0}{2\hbar} \omega |\mathbf{A}|^2 = \frac{\epsilon_0}{2\hbar \omega} |\mathbf{E}|^2$$

Rescale \mathbf{E} by $s = \sqrt{\frac{\epsilon_0}{2\hbar\omega}}$ to get x and y component wave function

$$\vec{\Psi} = \begin{pmatrix} \Psi_x \\ \Psi_y \end{pmatrix} = s \mathbf{E} = \sqrt{\frac{\epsilon_0}{2\hbar\omega}} \mathbf{E}$$

whose volume integral $\iiint_V dV \vec{\Psi}^* \cdot \vec{\Psi} = \iiint_V dV (|\Psi_x|^2 + |\Psi_y|^2) \propto \iiint_V dV \frac{N}{V} = N$

is $\iiint_V dV (|\Psi_x|^2 + |\Psi_y|^2) = N$ (It is normalized to particle number N .)

Poynting flux S is scaled to get counts per area·second.

$$S = cU = c\epsilon_0 |\mathbf{E}|^2 = \hbar n \omega \quad \text{where:} \quad n = Nc/V \text{ (per } m^2 \text{ per sec.)}$$

Review of wave parameters used to develop relativistic quantum theory

Bohr-Schrodinger (BS) approximation throws out Mc^2 (*Is frequency really relative?*)

Effect on group velocity (*None*) and phase velocity (*Absurd*)

1st Quantization: Quantizing phase variables k_m and $\omega(k_m)$

Understanding how quantum dynamics and transitions involve “mixed” states

Square well example of mixing unequal frequencies

Circle well or ring example of mixing equal or unequal frequencies

Mixing unequal amplitudes makes “Gallop” wave: Analogy of (*SWR, SWQ*) to (V_{group}, V_{phase})

Analogy with optical polarization geometry and Kepler orbits

Super-luminal speed and Feynman-Wheeler pair-creation switchbacks

2nd Quantization: Quantizing wave amplitudes A_N and invariance of photon number

Analogy 1: Many CW (Continuous Waves) add up to make PW (Pulse Waves)

Analogy 2: Many Photon-Number-Modes add up to make Coherent-Laser-Modes

Heisenberg $\Delta\nu\Delta t \sim 1 \sim \Delta\kappa\Delta x$ analogous to $\Delta N\Delta phase \sim 1$ uncertainty relations

Electromagnetic wave mode energy: Maxwell vs. Planck-Einstein

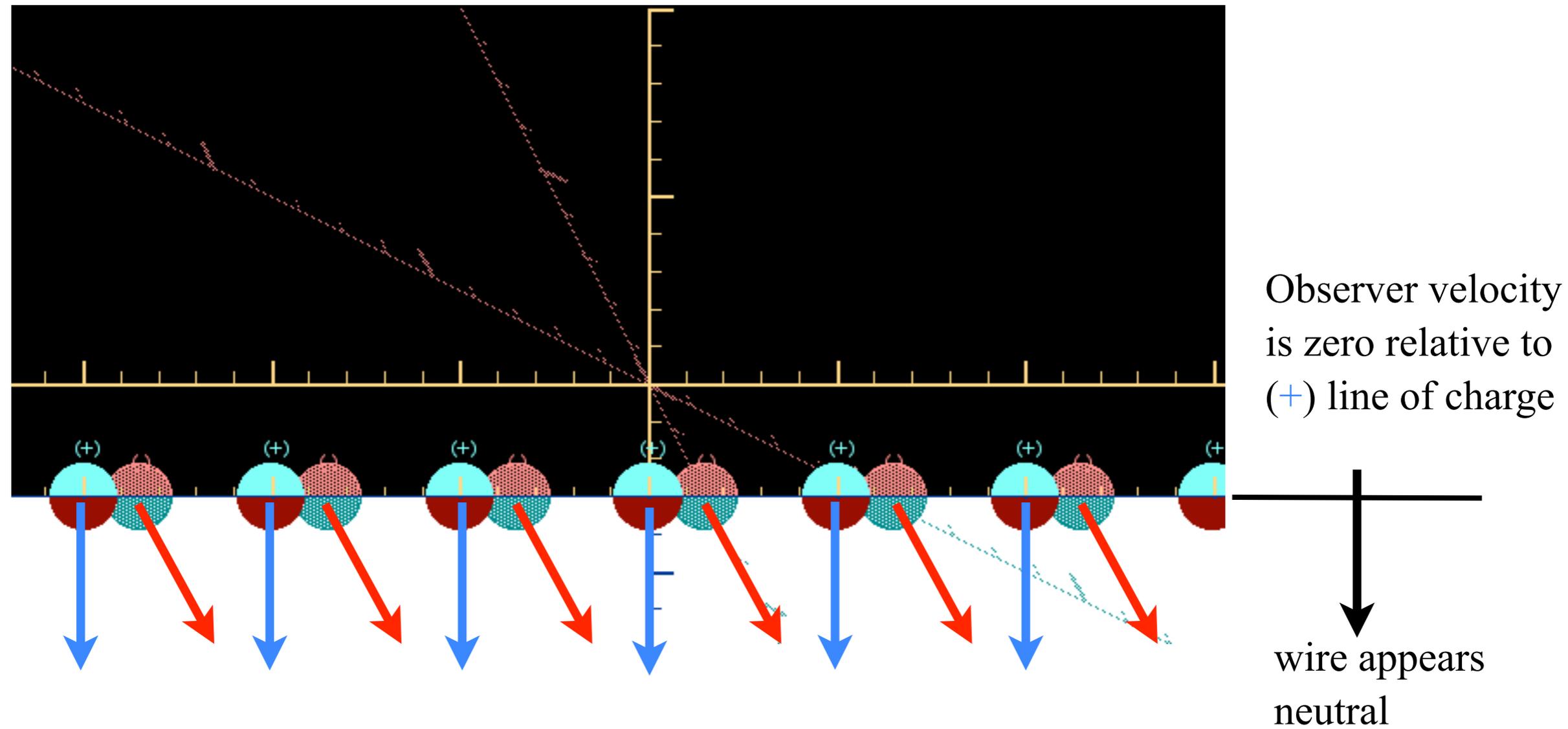
1st quantization for wave phase variables and classical energy of **E**, **B**, and **A** fields

2nd quantization for wave and Planck quantum energy of **E**, **B**, and **A** fields

Scaling **E**-waves to mime quantum Ψ -waves and ψ -waves

 Relativistic effects on charge, current, and Maxwell Fields

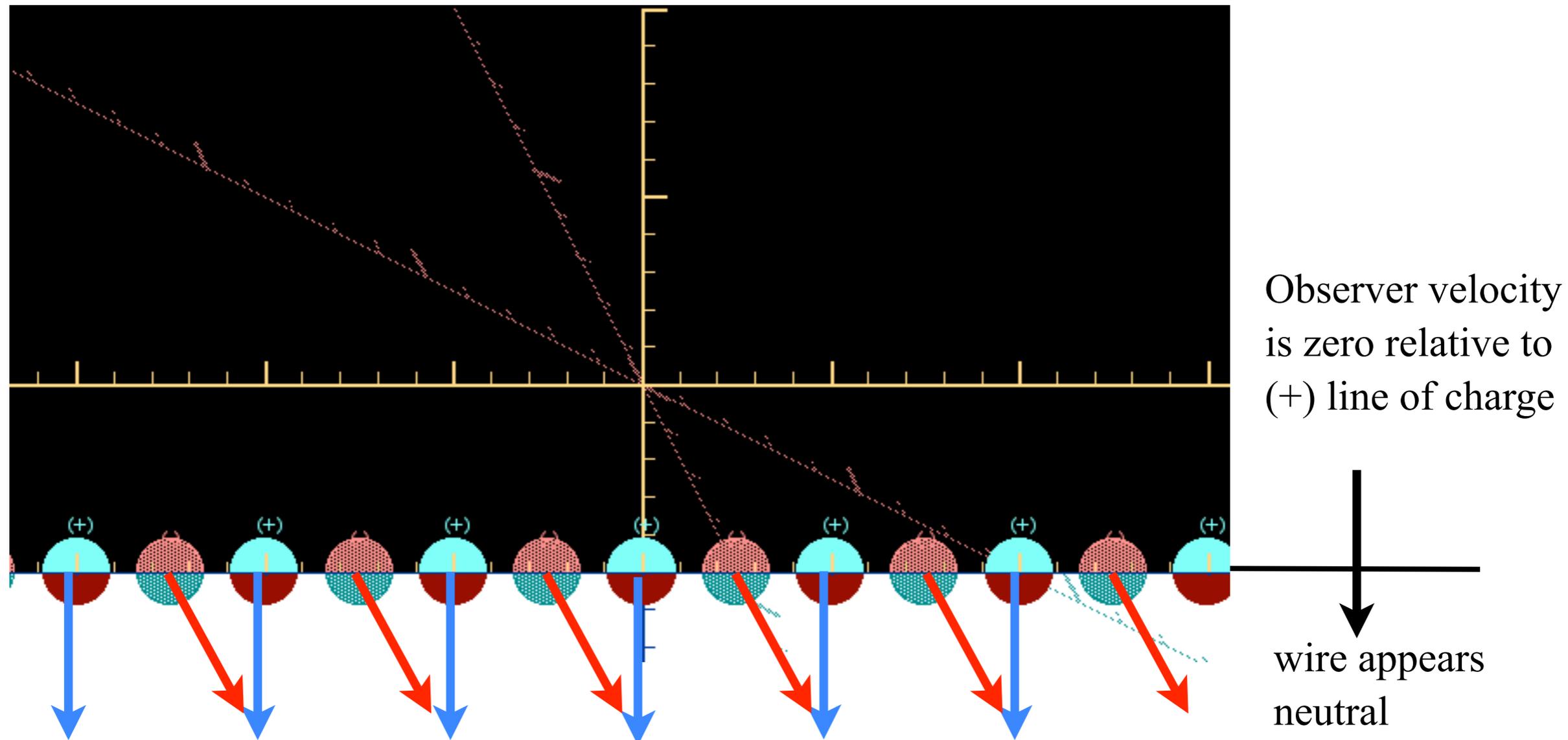
Relativistic effects on charge, current, and Maxwell Fields



(+) Charge fixed (-) Charge moving to right (*Negative current density*)

(+) Charge density is Equal to the (-) Charge density

Relativistic effects on charge, current, and Maxwell Fields



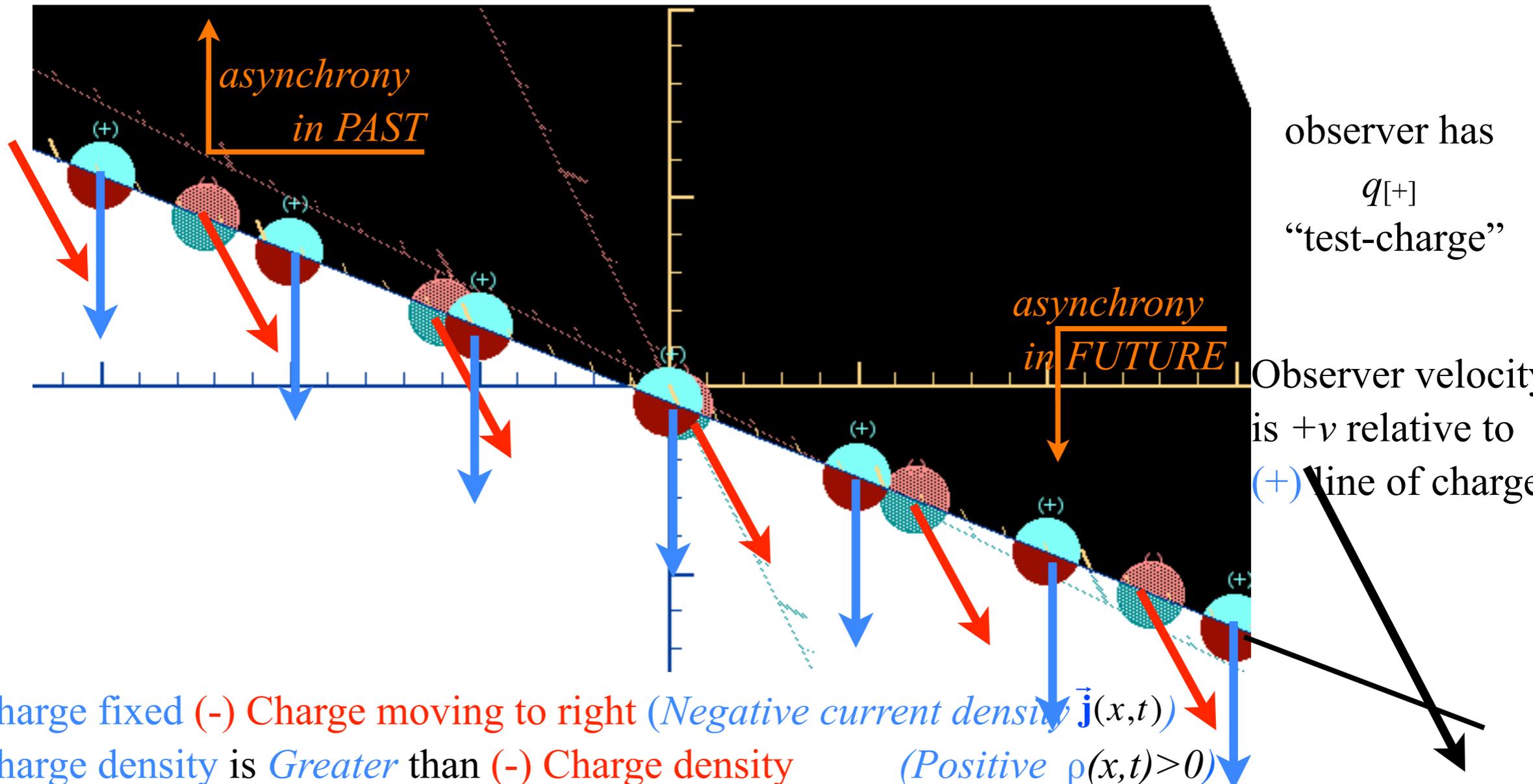
- (+) Charge fixed (-) Charge moving to right (*Negative current density* $\vec{j}(x,t)$)
- (+) Charge density is Equal to the (-) Charge density (*Zero* $\rho(x,t)=0$)

Relativistic effects on charge, current, and Maxwell Fields

Current density changes by Lorentz *asynchrony*

Asynchrony due to off-diagonal $\sinh \rho$ (a 1st-order effect)

in Lorentz transform:
$$\begin{pmatrix} \cosh \rho & \sinh \rho \\ \sinh \rho & \cosh \rho \end{pmatrix} \sim \begin{pmatrix} 1 & v/c \\ v/c & 1 \end{pmatrix}$$



(+) Charge fixed (-) Charge moving to right (Negative current density $\vec{j}(x,t)$)
 (+) Charge density is Greater than (-) Charge density (Positive $\rho(x,t) > 0$)

observer has $q_{[+]}$ "test-charge"

Observer velocity is $+v$ relative to (+) line of charge

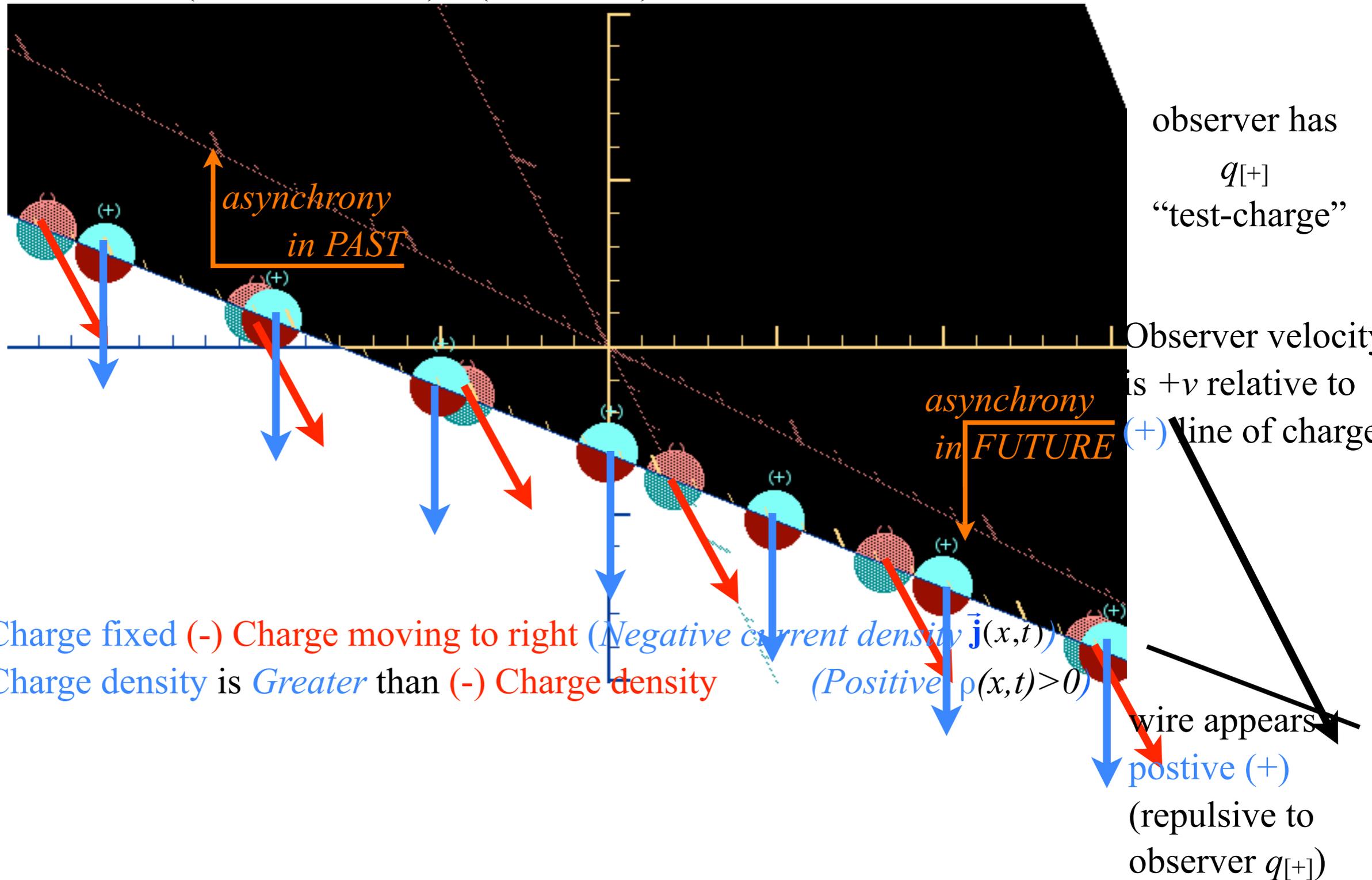
wire appears positive (+) (repulsive to observer $q_{[+]}$)

Relativistic effects on charge, current, and Maxwell Fields

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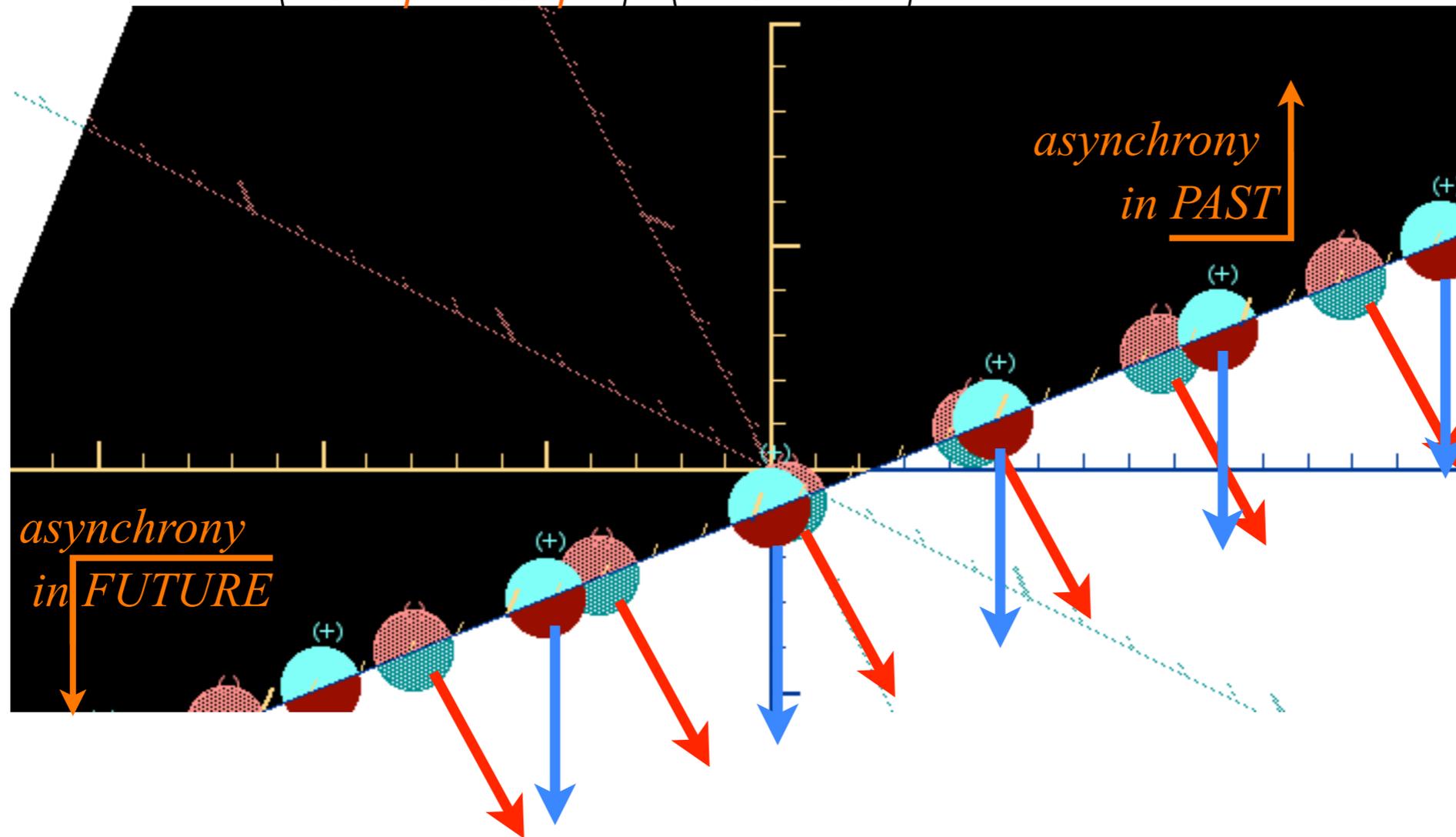
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observer has

$q_{[+]}$

“test-charge”

Observer velocity is $-v$ relative to $(+)$ line of charge



wire appears **negative (-)** (attractive to observer $q_{[+]}$)

(+) Charge fixed (-) Charge moving to right (Negative current density $\vec{j}(x,t)$)

(+) Charge density is *Less* than (-) Charge density (Negative $\rho(x,t) < 0$)

Relativistic effects on charge, current, and Maxwell Fields

Current density changes by Lorentz *asynchrony*

Asynchrony due to off-diagonal $\sinh \rho$ (a 1st-order effect)

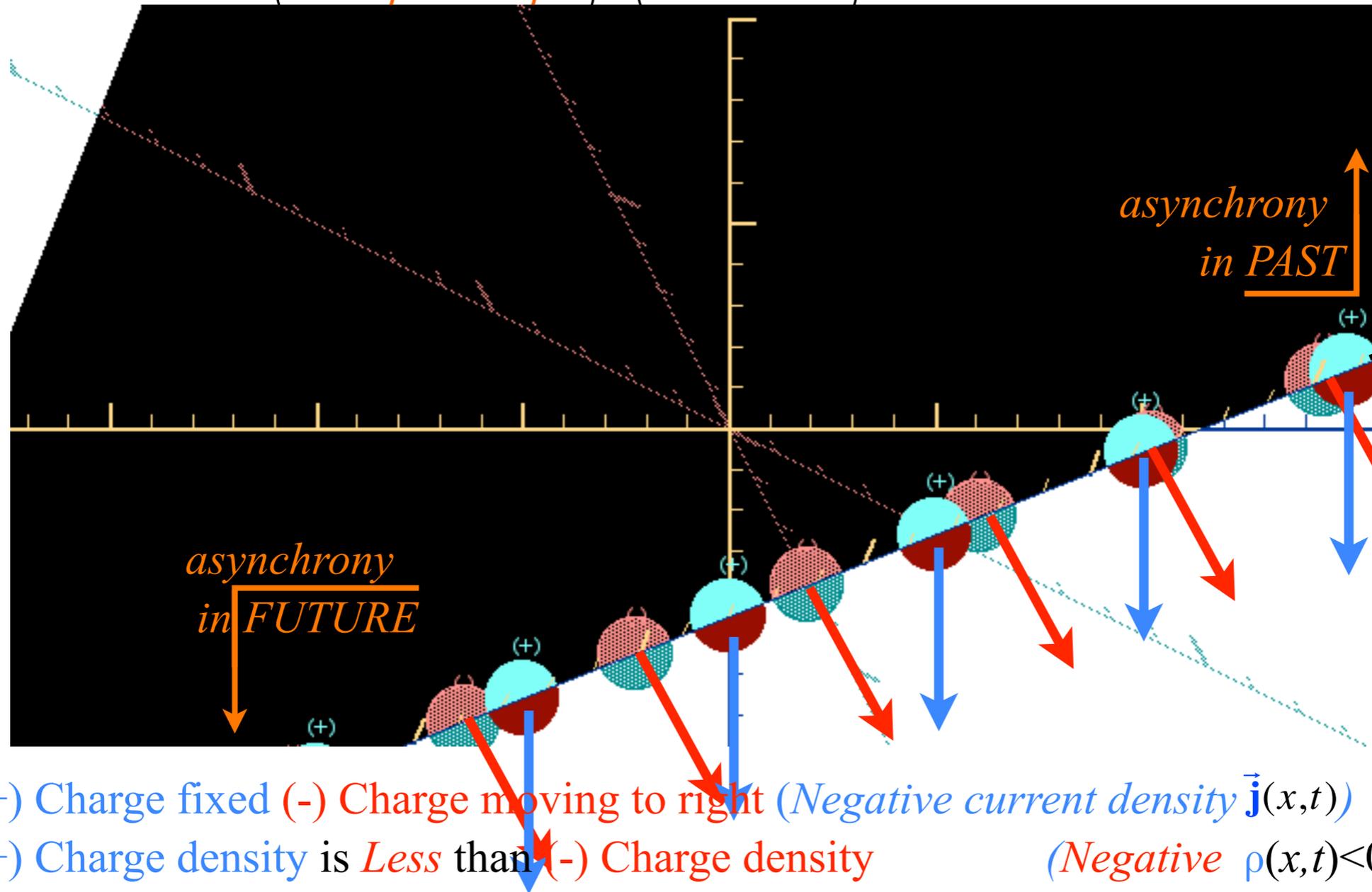
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$q_{[+]}$

“test-charge”

Observer velocity is $-v$ relative to (+) line of charge



asynchrony in PAST

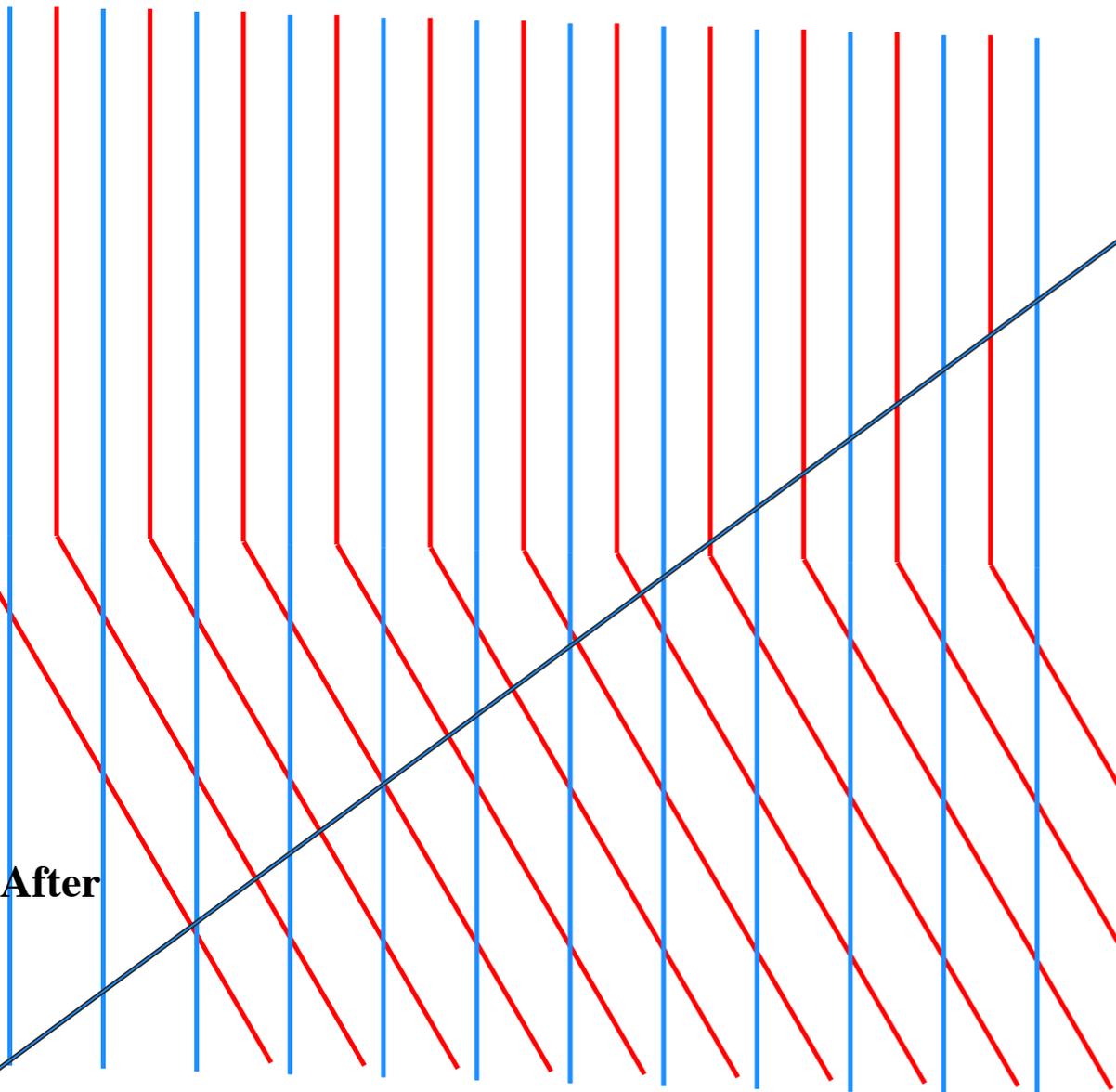
asynchrony in FUTURE

wire appears negative (-) (attractive to observer $q_{[+]}$)

- (+) Charge fixed (-) Charge moving to right (Negative current density $\vec{j}(x,t)$)
- (+) Charge density is *Less* than (-) Charge density (Negative $\rho(x,t) < 0$)

Simple 1st-order relativistic geometry of magnetism

Before



After

If Black is moving to Left

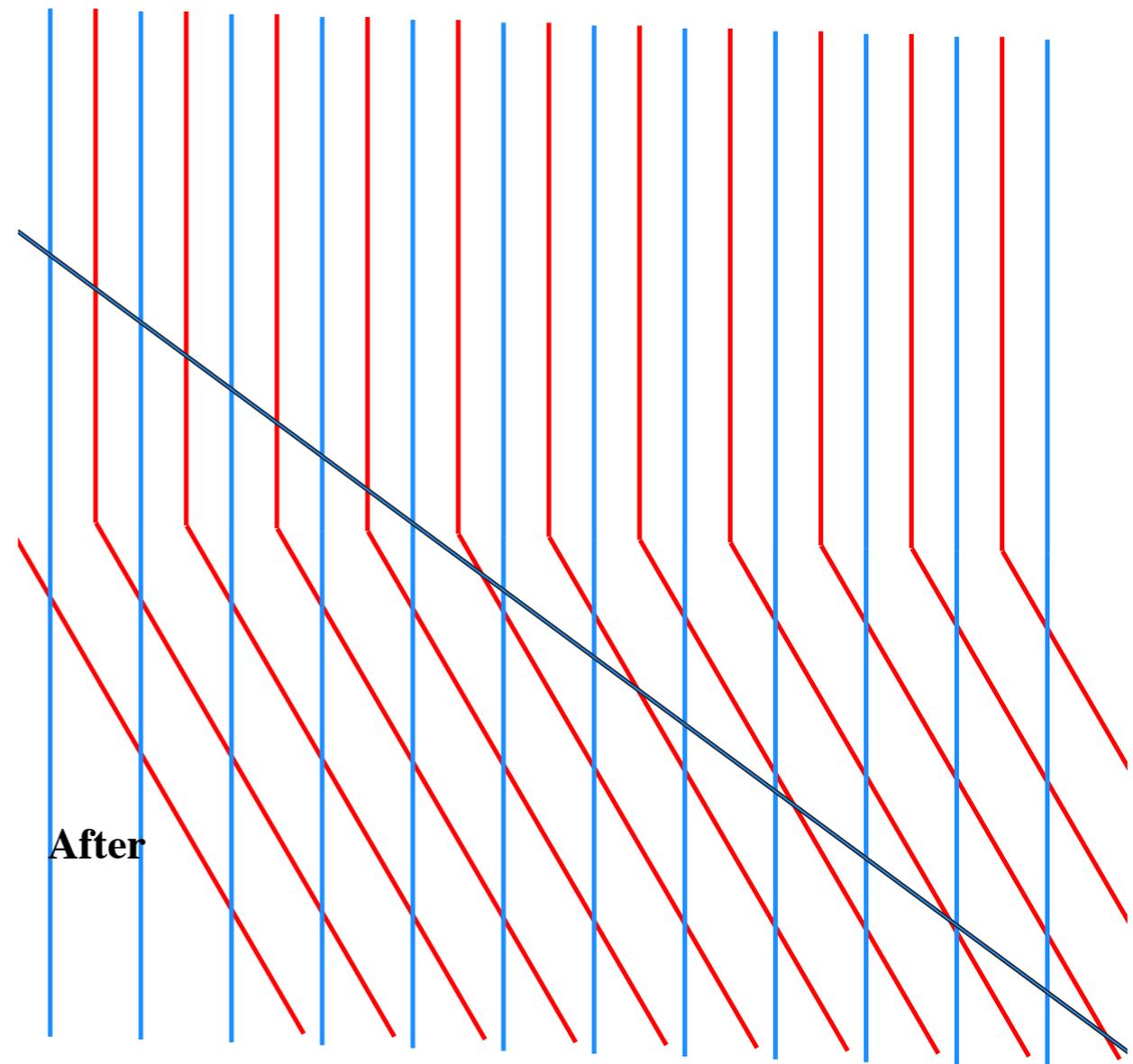
Before red starts moving to right

Black sees same number of red and blue

After red starts moving to right

Black sees more red than blue

Before



After

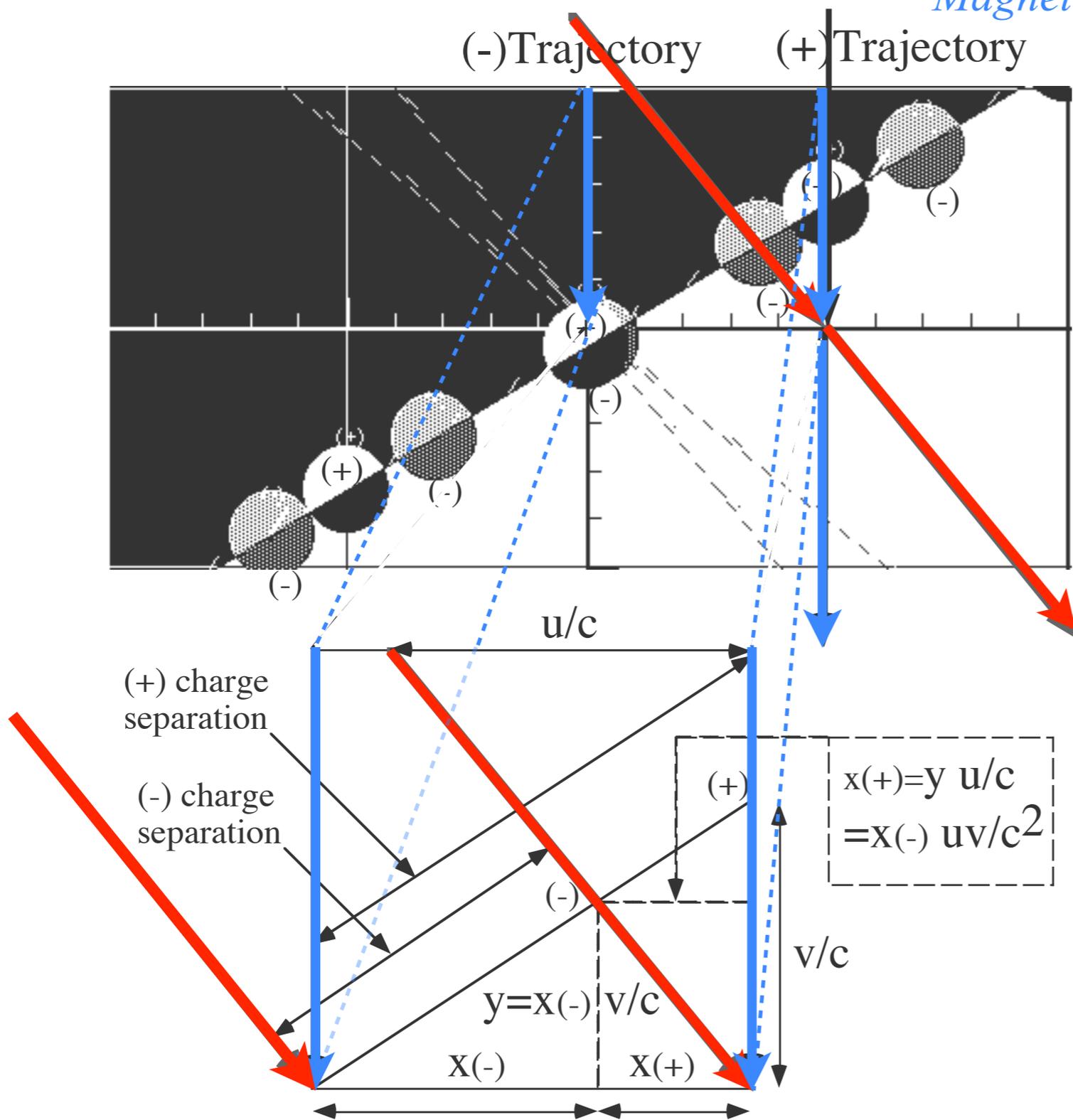
If Black is moving to Right

Before red starts moving to right

Black sees same number of red and blue

After red starts moving to right

Black sees more blue than red

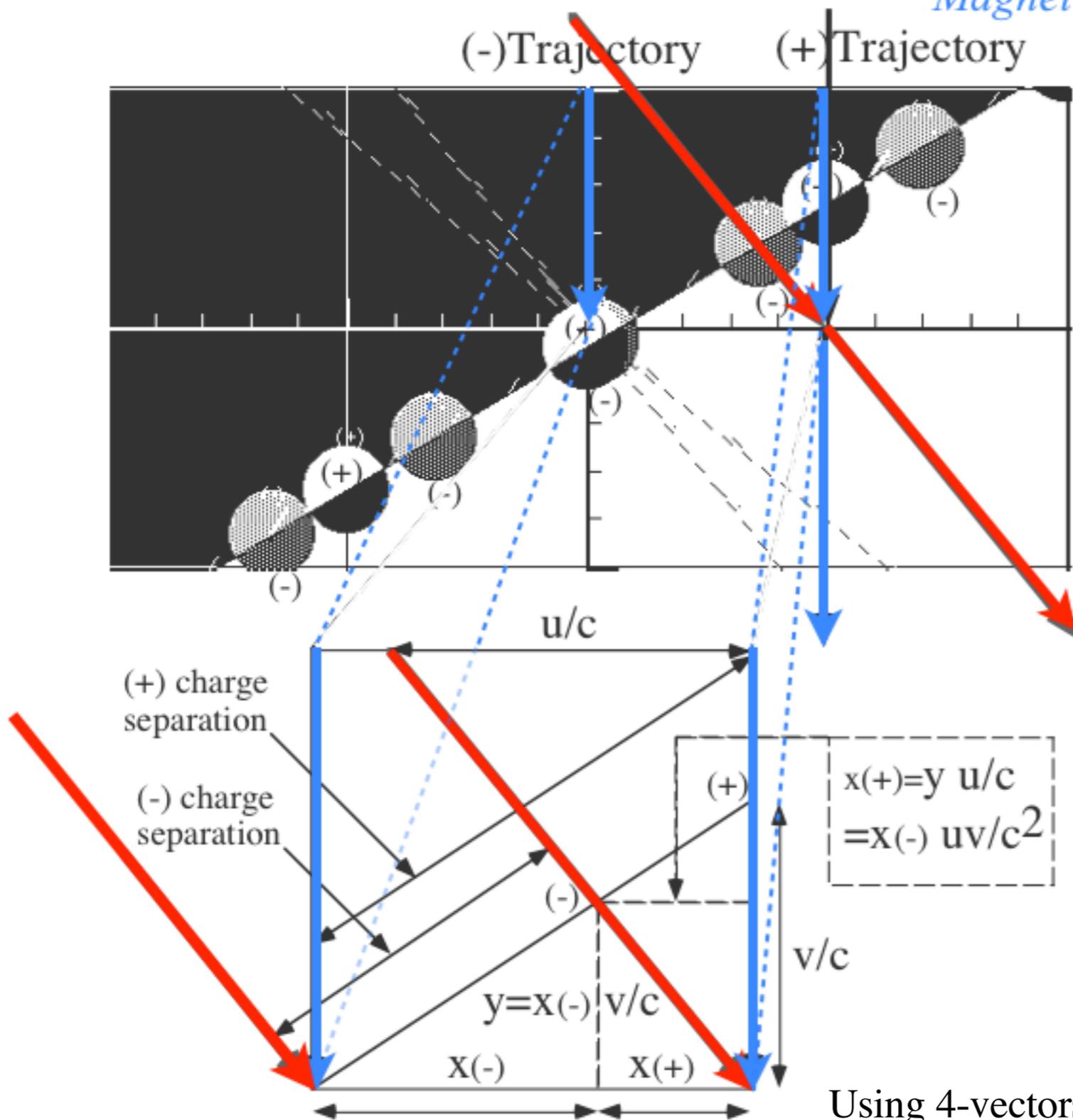


$$\frac{\rho(-)}{\rho(+)} = \frac{(+)\text{ charge separation}}{(-)\text{ charge separation}} = \frac{x(+)+x(-)}{x(-)}$$

$$\frac{\rho(-)}{\rho(+)} = \frac{x(+)}{x(-)} + 1 = \frac{uv}{c^2} + 1$$

$$\rho(+)-\rho(-) = \rho(+)\left(1 - \frac{\rho(-)}{\rho(+)}\right) = -\frac{uv}{c^2}\rho(+)$$

Unit square: $(u/c) / 1 = x(+)/y$
 $(v/c) / 1 = y/x(-)$



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Unit square: $(u/c) / 1 = x(+)/y$
 $(v/c) / 1 = y/x(-)$

Using 4-vectors to EL Transform (charge-current)=($c\rho, \mathbf{j}$)

$$\begin{pmatrix} c\rho' \\ j_{x'} \\ j_{y'} \\ j_{z'} \end{pmatrix} = \begin{pmatrix} \cosh \rho & \sinh \rho & \cdot & \cdot \\ \sinh \rho & \cosh \rho & \cdot & \cdot \\ \cdot & \cdot & 1 & \cdot \\ \cdot & \cdot & \cdot & 1 \end{pmatrix} \begin{pmatrix} c\rho \\ j_x \\ j_y \\ j_z \end{pmatrix}$$

The electric force field \mathbf{E} of a charged line varies inversely with radius. The Gauss formula for force in mks units :

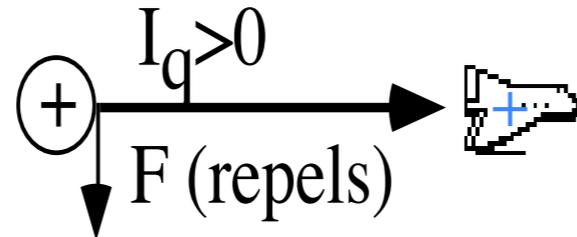
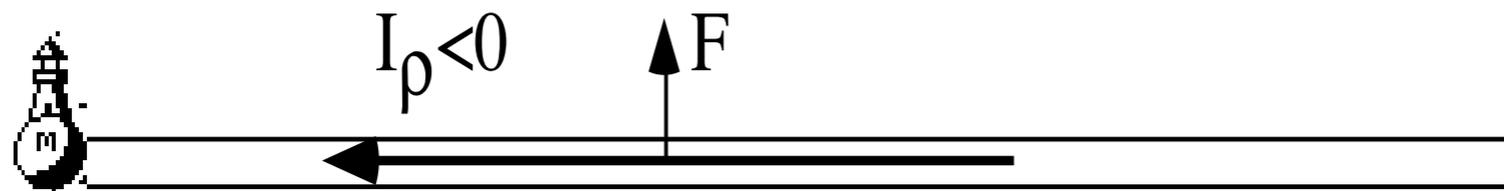
$$F = qE = q \left[\frac{1}{4\pi\epsilon_0} \frac{2\rho}{r} \right], \quad \text{where: } \frac{1}{4\pi\epsilon_0} = 9 \times 10^9 \frac{N \cdot m^2}{\text{Coul.}}$$

$$F = qE = q \left[\frac{1}{4\pi\epsilon_0} \frac{2}{r} \left(-\frac{uv}{c^2} \rho(+)\right) \right] = -\frac{2qv\rho(+)}{4\pi\epsilon_0 c^2 r} = -2 \times 10^{-7} \frac{I_q I_\rho}{r}$$

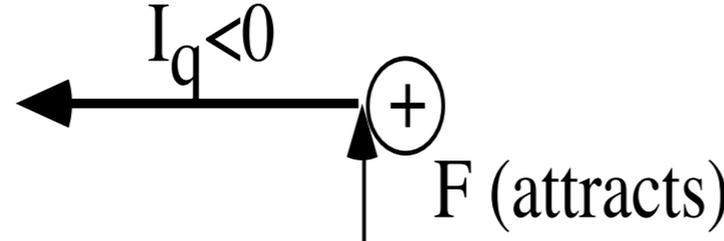
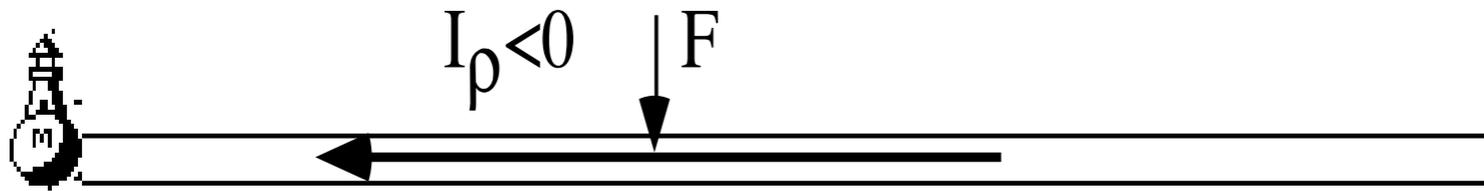
$$1/4\pi\epsilon_0 = 9 \cdot 10^9$$

$$c^2 = 9 \cdot 10^{16}$$

$$1/(4\pi\epsilon_0 c^2) = 10^{-7}$$



*I see excess (+)
charge up there. Yuk!*



*I see excess (-)
charge up there. Yum!*

The electric force field \mathbf{E} of a charged line varies inversely with radius. The Gauss formula for force in mks units :

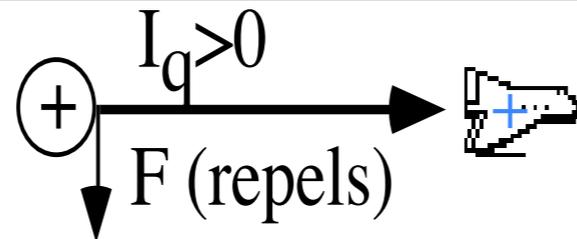
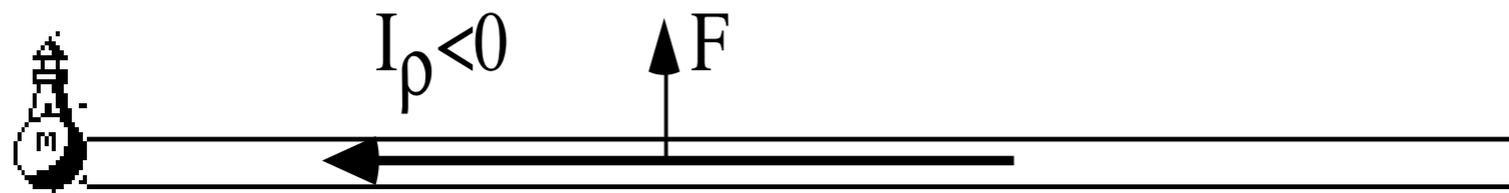
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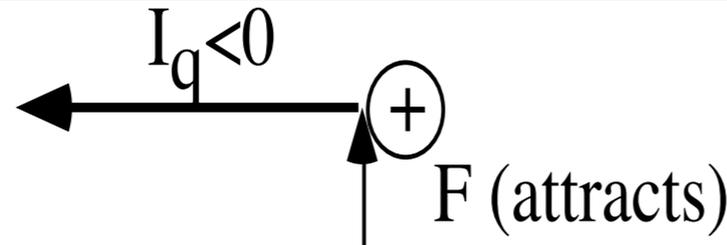
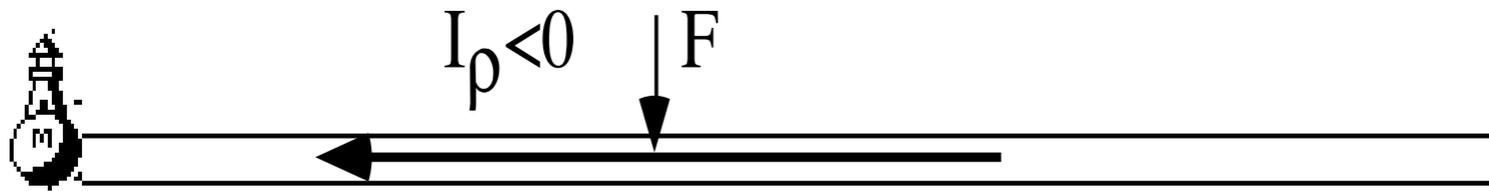
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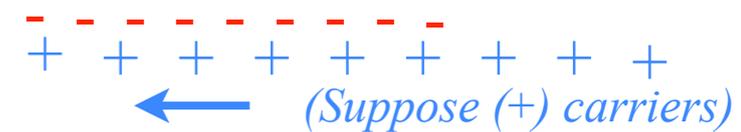
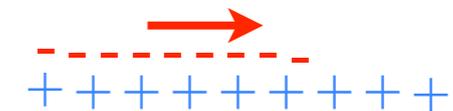
$$1/(4\pi\epsilon_0 c^2) = 10^{-7}$$



I see excess (+) charge up there. Yuk!



I see excess (-) charge up there. Yum!



2005 Pirelli treatments

Relating photons to Maxwell energy density and Poynting flux

Relativistic variation and invariance of frequency (ω, k) and amplitudes

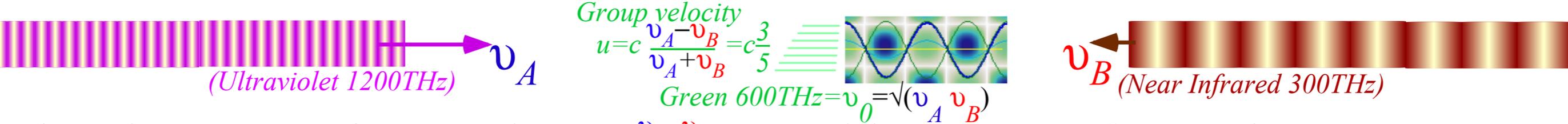
How probability ψ -waves and flux ψ -waves evolved

Properties of amplitude ψ^ψ -squares*

More on unmatched amplitudes AND unmatched frequencies AND unmatched quanta

Light Energy and Flux 2-CW vs 1-CW-light

What if head-on CW's $\nu_A=1200\text{THz}$ and $\nu_B=300\text{THz}$ pair-up in a 2-CW-light beam?

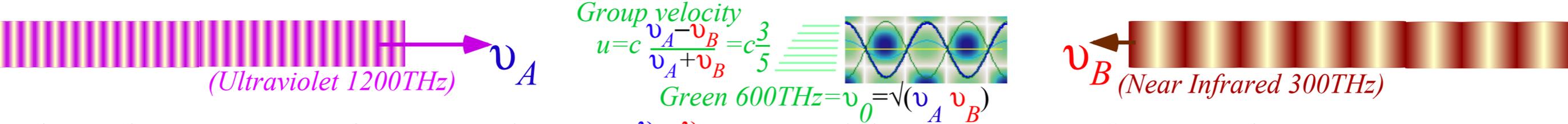


They form a *rest frame* going $u = c \frac{\nu_A - \nu_B}{\nu_A + \nu_B} = 3c/5$ with a *mean* or *base* color $\nu_0 = \sqrt{\nu_A \nu_B}$

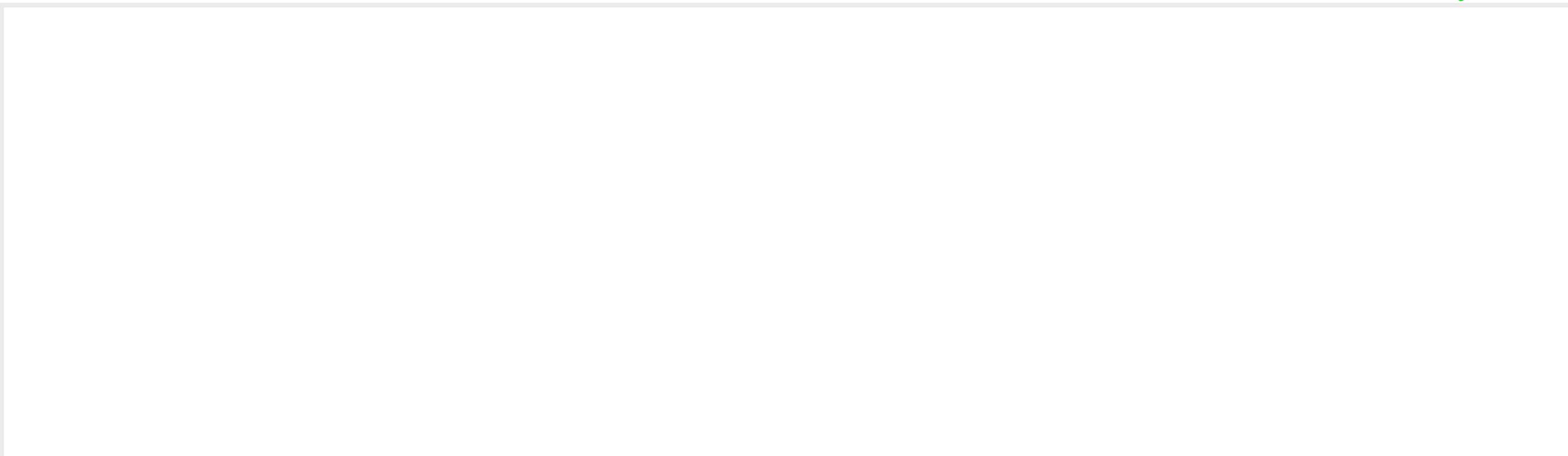
($\nu_0 = B = 600\text{THz}$ is *green* here. Neither has this singly.)

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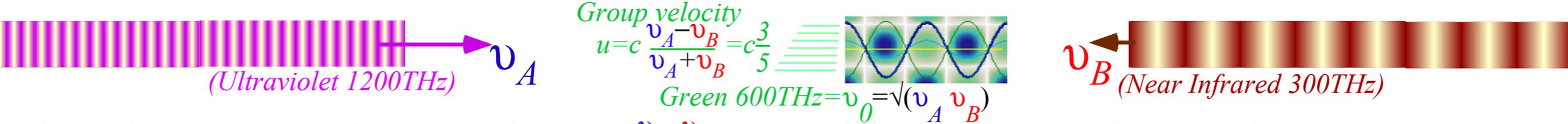


They form a *rest frame* going $u=c\frac{\nu_A-\nu_B}{\nu_A+\nu_B}=3c/5$ with a *mean* or *base* color $\nu_0=\sqrt{\nu_A\nu_B}$ ($\nu_0=B=600\text{THz}$ is *green* here. Neither has this singly.) *All* observers agree on ν_0 since *all* shift-products $b\nu_A r\nu_B$ equal $(\nu_0)^2$ due to Doppler-time-symmetry ($b=1/r$). Single CW's get *invariant* properties if they pair-up. The $\nu_A-\nu_B$ pairing above makes a number \bar{N} of *invariant mass quanta* $M_1=h\nu_0/c^2=4.42\cdot 10^{-36}\text{kg}$ where the number \bar{N} is invariant, too. \bar{N} is Planck's *photon number* for the cavity rest energy $E=\bar{N}h\nu_0$.



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Relating Planck's E to Maxwell's Density U=E/V

Maxwell field energy E , a product of mean-square electric field $\langle E^2 \rangle$, volume of cavity V , and constant $\epsilon_0=8.854\cdot 10^{-12}\text{C}^2/\text{N}\cdot\text{m}^2$, approximates Planck's energy $\bar{N}h\nu_0$.

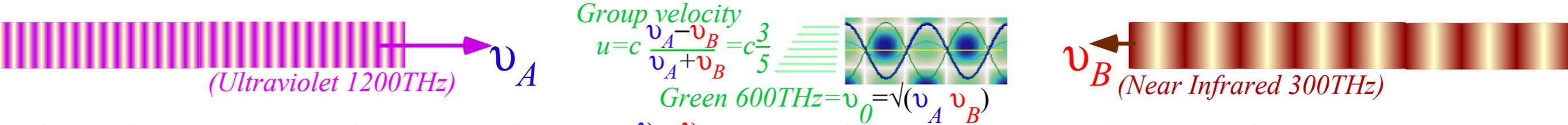
$$E = \langle E^2 \rangle V \epsilon_0 = \bar{N} h \nu_0 \quad \text{Maxwell-Planck Energy}$$

$$U = \langle E^2 \rangle \epsilon_0 = \bar{N} h \nu_0 / V \quad \text{Maxwell-Planck Density}$$

$$\text{Field Energy} = |\mathbf{E}|^2 \epsilon_0 \quad 1/4\pi\epsilon_0 = 9\cdot 10^9$$

Light Energy and Flux 2-CW vs 1-CW-light

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$$E = \langle E^2 \rangle V \epsilon_0 = \bar{N} h \nu_0 \quad \text{Maxwell-Planck Energy}$$

$$U = \langle E^2 \rangle \epsilon_0 = \bar{N} h \nu_0 / V \quad \text{Maxwell-Planck Density}$$

Example: Let a $\frac{1}{4}\mu\text{m}$ -cube cavity (*Half-wave at 600Thz*) have $\bar{N}=10^{10}$ photons in volume $V=(\frac{1}{4}10^{-6}\text{m})^3$.

$$\text{Energy per photon: } h\nu_0 = 4\cdot 10^{-19}\text{J} = 2.5\text{ eV}$$

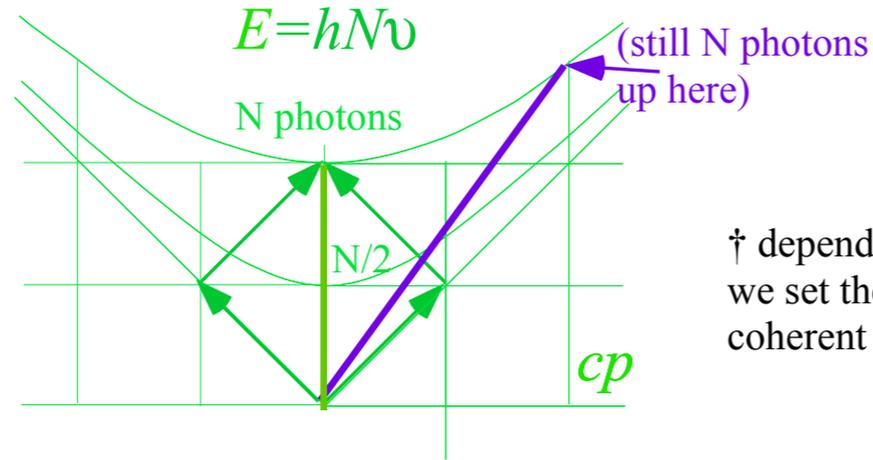
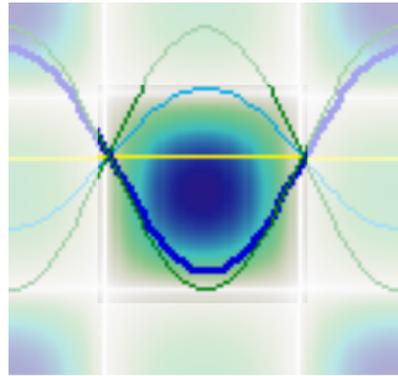
$$\text{Energy of } \bar{N} \text{ photons: } \bar{N} h \nu_0 = 4\cdot 10^{-9}\text{J} = 25\text{GeV}$$

$$\text{E-field per photon: } E_1 = \sqrt{(h\nu_0 / V \epsilon_0)} = 7.6\cdot 10^3\text{V/m}$$

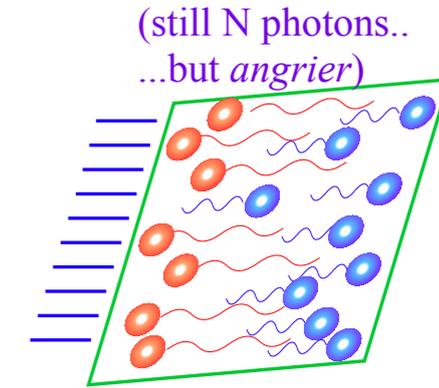
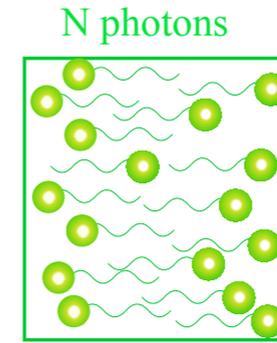
$$\text{E-field of } \bar{N} \text{ photons: } E_N = 7.6\cdot 10^{13}\text{V/m}$$

Energy and Flux (contd) 2-CW- vs 1-CW-light

Planck $E=Nh\nu$ relation allows us to interpret our N -quantized 2-CW mode as a box or *cavity* of $N_{(\text{more-or-less}\dagger)}$ photons where N is invariant to speed u of box.

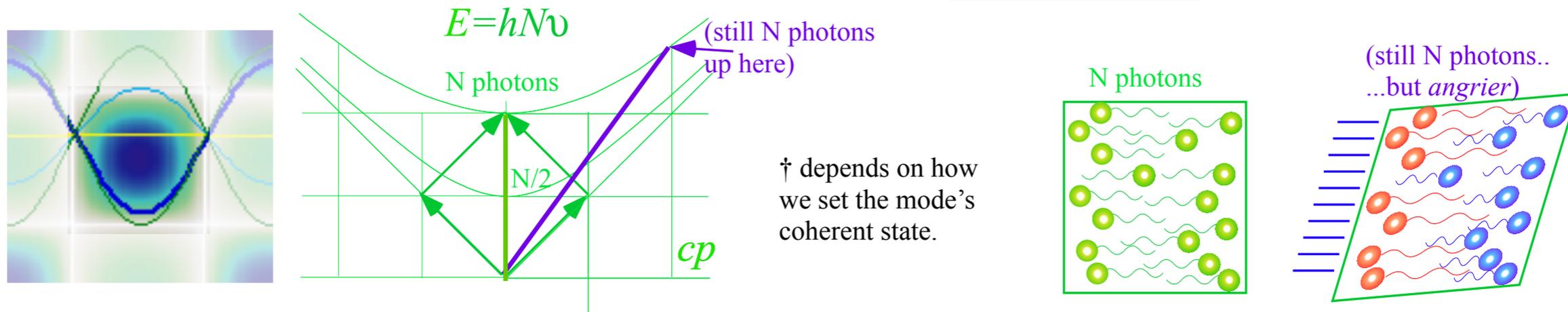


† depends on how we set the mode's coherent state.



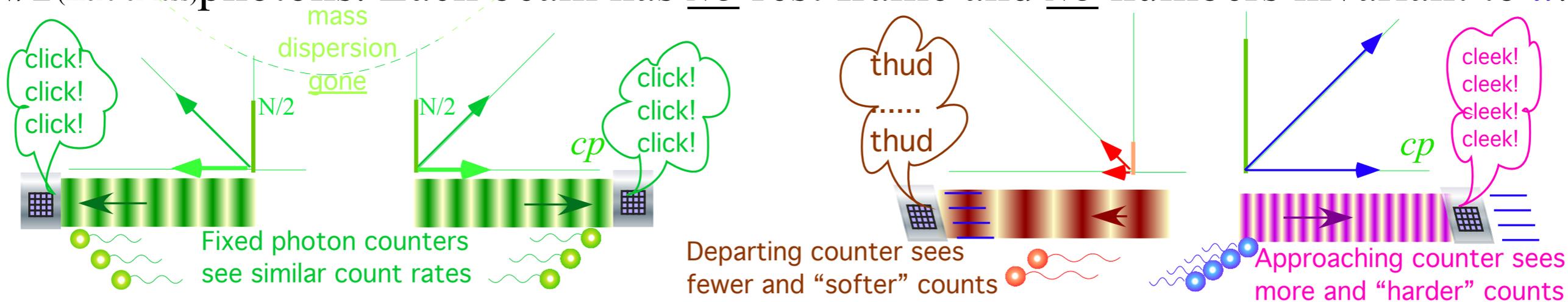
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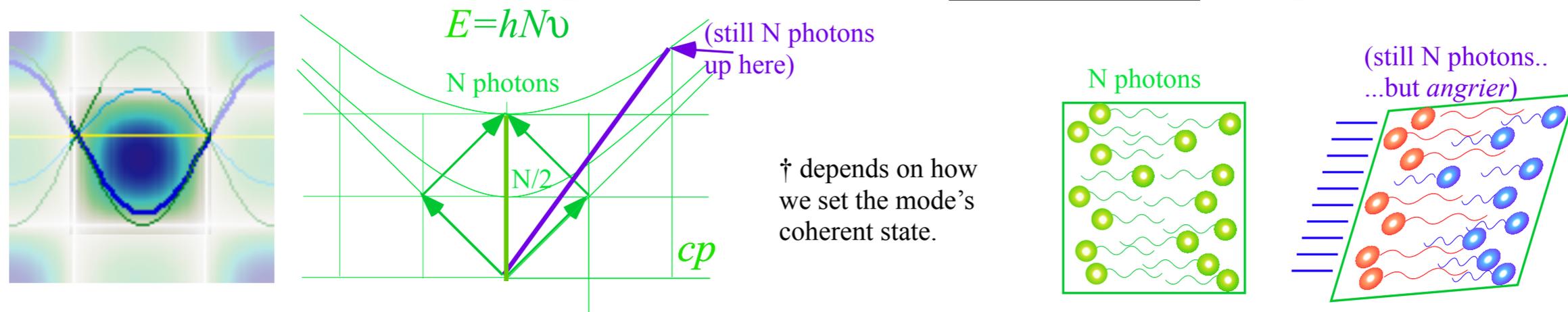
† depends on how we set the mode's coherent state.

If we open the box our 2-CW mode “divorces” into two separate 1-CW beams of $N/2_{(\text{more-or-less})}$ photons. Each beam has NO rest frame and NO numbers invariant to u .

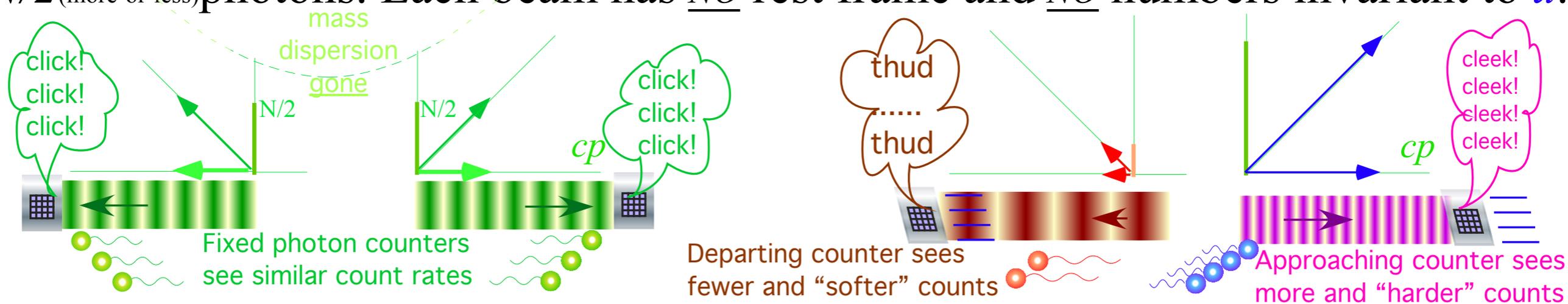


Energy and Flux (contd) 2-CW- vs 1-CW-light

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Relating Poynting's Intensity $S=cU$ to Planck's Flux

Poynting intensity S is a product of $c=2.99792458m/s$ and density U . It approximates Planck's energy $E=Nh\nu$ times c and divided by cavity volume V .

$$S=cU=(Nc/V)h\nu = n h\nu \quad \text{Poynting-Planck Flux (Watts per square meter)}$$

The *photon-count rate* is $n=Nc/V$ (per square meter per second) and $h\nu$ is energy (per count).

Relating photons to Maxwell energy density and Poynting flux

➔ *Relativistic variation and invariance of frequency (ω, k) and amplitudes*

How probability ψ -waves and flux ψ -waves evolved

Properties of amplitude ψ^ψ -squares*

More on unmatched amplitudes AND unmatched frequencies AND unmatched quanta

Frequency and Amplitude Variance 2-CW-light vs 1-CW-light

2-CW modes have invariance

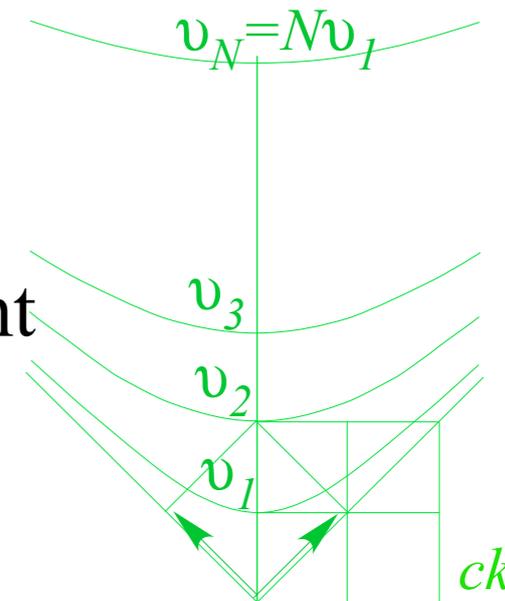
Maxwell-Planck energy E is photon number $N(m^{-3})$ times 2-CW-frequency ν_1 .

Invariant to ρ

Each is ρ -invariant

$$E = \langle U \rangle \cdot V = \epsilon_0 \langle E^2 \rangle \cdot V = \epsilon_0 \langle E_{2-CW}^* E_{2-CW} \rangle \cdot V = hN\nu_1 = h\nu_N$$

Photon number N and rest-frame frequencies $\nu_1 \dots \nu_N$ are invariant to rapidity ρ and occupy (ω, ck) -hyperbolas in per-spacetime.



Frequency and Amplitude Variance 2-CW-light vs 1-CW-light

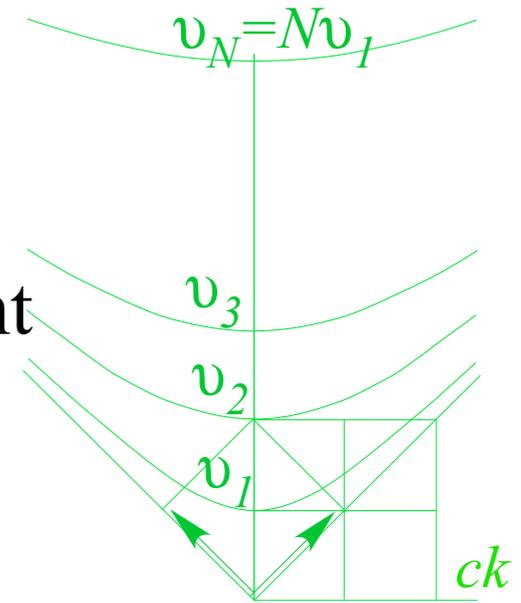
2-CW modes have invariance

Maxwell-Planck energy E is *photon number* $N(m^{-3})$ times *2-CW-frequency* ν_1 .

Invariant to ρ

Each is ρ -invariant

$$E = \langle U \rangle \cdot V = \epsilon_0 \langle E^2 \rangle \cdot V = \epsilon_0 \langle E_{2-CW}^* E_{2-CW} \rangle \cdot V = hN\nu_1 = h\nu_N$$



Photon number N and rest-frame frequencies $\nu_1 \dots \nu_N$ are invariant to rapidity ρ and occupy (ω, ck) -hyperbolas in per-spacetime.

1-CW beams lack invariance (have “variance” ala’ Doppler)

Planck-Poynting flux S is *count rate* $n = Nc/V(m^{-2}s^{-1})$ times *1-CW-frequency* ν_{\leftarrow} or ν_{\rightarrow} .

Count rate n and frequency ν Doppler shift by $b = e^{\pm\rho}$ factors and occupy $(\omega = \pm ck)$ -baselines.

Shifts by $b = e^{+2\rho}$

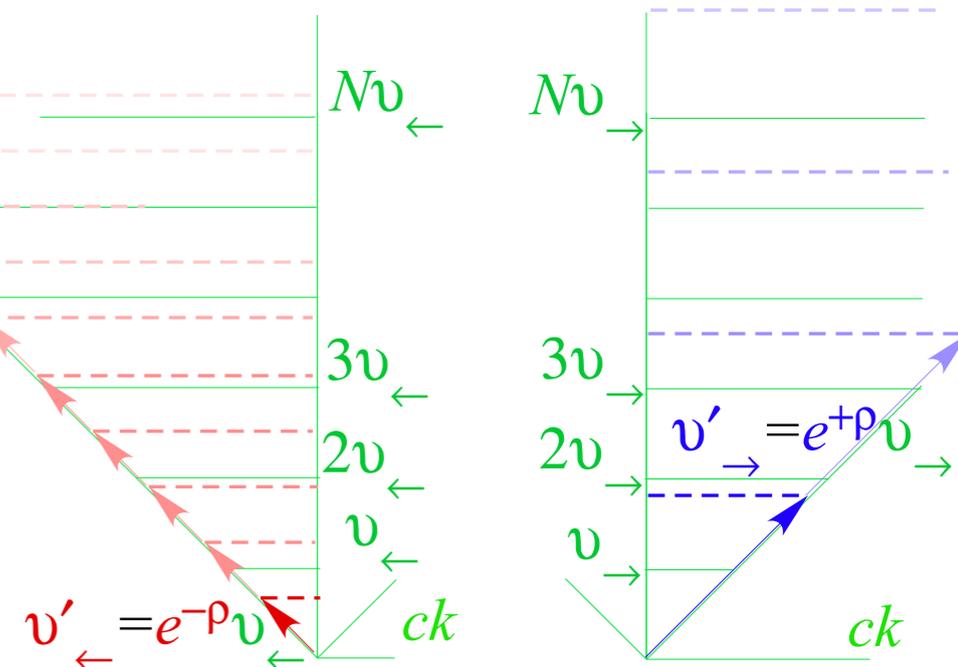
Each blue shifts by $b = e^{+\rho}$

$$S_{\rightarrow} = cU_{\rightarrow} = c\epsilon_0 \langle E^2 \rangle = c\epsilon_0 \langle E_{1-CW}^{\rightarrow} * E_{1-CW}^{\rightarrow} \rangle = hn_{\rightarrow} \nu_{\rightarrow}$$

$$S_{\leftarrow} = cU_{\leftarrow} = c\epsilon_0 \langle E^2 \rangle = c\epsilon_0 \langle E_{1-CW}^{\leftarrow} * E_{1-CW}^{\leftarrow} \rangle = hn_{\leftarrow} \nu_{\leftarrow}$$

Shifts by $r = e^{-2\rho}$

Each red shifts by $r = e^{-\rho}$



Note: $E_{1-CW}^{\leftrightarrow} \sqrt{(c\epsilon_0/h)} = \sqrt{(n_{\leftrightarrow} \nu_{\leftrightarrow})}$ is geometric mean of *amplitude frequency* n_{\leftrightarrow} and *phase frequency* ν_{\leftrightarrow} .

Important result below:

*Amplitudes of 1-CW “exponentiate” just like frequency,
and intensity does at twice the rate
(A double-double whammy!)*

1-CW beams lack invariance (have “variance” ala’ Doppler)

Planck-Poynting flux S is *count rate* $n=Nc/V(m^{-2}s^{-1})$ times *1-CW-frequency* ν_{\leftarrow} or ν_{\rightarrow} .

Count rate n and frequency ν Doppler shift

by $b=e^{\pm\rho}$ factors and occupy $(\omega=\pm ck)$ -*baselines*.

Shifts by $b=e^{+2\rho}$

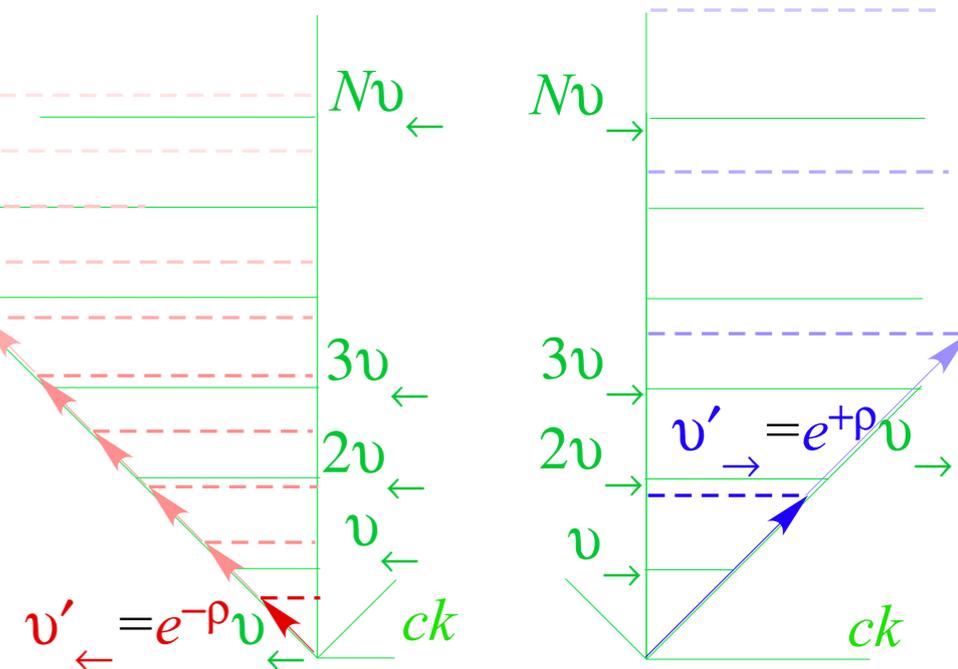
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Shifts by $r=e^{-2\rho}$

Each blue shifts by $b=e^{+\rho}$

Each red shifts by $r=e^{-\rho}$



Note: $E_{1-CW}^{\leftrightarrow} \sqrt{(c\epsilon_0/h)} = \sqrt{(n_{\leftrightarrow} \nu_{\leftrightarrow})}$ is geometric mean of *amplitude frequency* n_{\leftrightarrow} and *phase frequency* ν_{\leftrightarrow} .

Relating photons to Maxwell energy density and Poynting flux

Relativistic variation and invariance of frequency (ω, k) and amplitudes

 *How probability ψ -waves and flux ψ -waves evolved*

Properties of amplitude ψ^ψ -squares*

More on unmatched amplitudes AND unmatched frequencies AND unmatched quanta

How Probability Amplitudes ψ or Ψ Come About (An optical view)

Maxwell-Planck-Poynting flux $S=cU=c\epsilon_0|E|^2=c\epsilon_0E^*E=n\hbar\nu$ has count rate $n=Nc/V(m^{-2}s^{-1})$

If each E-field amplitude factor is scaled by a factor $\sqrt{\frac{c\epsilon_0}{\hbar\nu}} = \sqrt{\frac{\epsilon_0}{\hbar\kappa}}$ the result is a flux probability amplitude $\psi = E\sqrt{\frac{c\epsilon_0}{\hbar\nu}}$ whose square equals flux count rate $n(m^{-2}s^{-1})$.

$$\psi^*\psi = n \quad (m^{-2}s^{-1})$$

A fixed probability amplitude $\Psi = E\sqrt{\frac{\epsilon_0}{\hbar\nu}}$ has square equal to N/V (particles per volume).

$$\Psi^*\Psi = N/V \quad (m^{-3})$$

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Here's how to answer Planck's worry about photons

Q: How can classical oscillator energy (Amplitude)²(frequency)² give with linear Planck law $S=n\hbar\nu$?

A: Let amplitude ψ or Ψ contain inverse square root of frequency: $\psi = E\sqrt{\frac{c\epsilon_0}{\hbar\nu}}$ the "quantum amplitude"

Energy $\sim |A|^2 \nu^2$ where vector potential \mathbf{A} defines electric field: $\mathbf{E} = \frac{\partial \mathbf{A}}{\partial t} = i\omega \mathbf{A} = 2\pi i\nu \mathbf{A}$

$$\text{Energy} \sim |A|^2 \nu^2 = |A\sqrt{\nu}|^2 \nu = \left| \frac{E}{2\pi\nu} \sqrt{\nu} \right|^2 \nu = \left| \frac{E}{2\pi\sqrt{\nu}} \right|^2 \nu \sim \left| E\sqrt{\frac{c\epsilon_0}{\hbar\nu}} \right|^2 = n\hbar\nu$$

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Probability Waves $\psi(x,t)$ (More optical views)

Optical E-field amplitudes like $E(x,t)=E_0e^{i(kx-\omega t)}$ vary with space x and time t . So do scaled $\psi(x,t)$ amplitudes whose sum- Σ (integral- \int) over cells ΔV (or dV) must be particle number N . For 1-particle systems ($N=1$) this is the *unit norm* rule.

$$\sum_j \psi(x_j, t)^* \psi(x_j, t) \Delta V_j = N \quad \text{or:} \quad \int \psi(x, t)^* \psi(x, t) dV = N$$

How Probability Amplitudes ψ or Ψ Come About (An optical view)

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Born interpreted $\psi(x,t)^*\psi(x,t)$ as *probable expectation* of particle count. Schrodinger objected to the *probability wave* interpretation that is now accepted and called the Schrodinger theory. A relativistic wave view lends merit to his objections.

Doppler Transformation of 2-CW Modes

Doppler shift of *opposite-k 1-CW* beams. As derived before phases are invariant: $(k'x' - \omega't' = kx - \omega t)$

E-wave: $\mathbf{E} = \mathbf{E}_{\rightarrow} e^{i(k_{\rightarrow}x - \omega_{\rightarrow}t)} + \mathbf{E}_{\leftarrow} e^{i(k_{\leftarrow}x - \omega_{\leftarrow}t)}$

Ψ -wave: $\Psi = \psi_{\rightarrow} e^{i(k_{\rightarrow}x - \omega_{\rightarrow}t)} + \psi_{\leftarrow} e^{i(k_{\leftarrow}x - \omega_{\leftarrow}t)}$

blue shift

$$\begin{aligned} \mathbf{E}'_{\rightarrow} &= \mathbf{b} \mathbf{E}_{\rightarrow} \\ &= e^{+\rho} \mathbf{E}_{\rightarrow} \end{aligned}$$

red shift

$$\begin{aligned} \mathbf{E}'_{\leftarrow} &= \mathbf{r} \mathbf{E}_{\leftarrow} \\ &= e^{-\rho} \mathbf{E}_{\leftarrow} \end{aligned}$$

$$\psi = E \sqrt{\frac{\epsilon_0}{h\nu}}$$

scaled blue shift

$$\begin{aligned} \psi'_{\rightarrow} &= \sqrt{\mathbf{b}} \psi_{\rightarrow} \\ &= e^{+\rho/2} \psi_{\rightarrow} \end{aligned}$$

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Parameters related to *relative velocity* u :

$$\beta = u/c = \tanh \rho = \frac{\sinh \rho}{\cosh \rho} = \frac{e^{+\rho} - e^{-\rho}}{e^{+\rho} + e^{-\rho}} = \frac{b^2 - 1}{b^2 + 1}$$

$$b^2 = \frac{1 + \beta}{1 - \beta} = \frac{1 + \tanh \rho}{1 - \tanh \rho}$$

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Transformation of SWR (or SWQ) and u_{GROUP} (or u_{PHASE}) is a *non-linear* transformation

$$SWR' = \frac{\mathbf{E}'_{\rightarrow} - \mathbf{E}'_{\leftarrow}}{\mathbf{E}'_{\rightarrow} + \mathbf{E}'_{\leftarrow}} = \frac{b^2 \mathbf{E}_{\rightarrow} - \mathbf{E}_{\leftarrow}}{b^2 \mathbf{E}_{\rightarrow} + \mathbf{E}_{\leftarrow}} = \frac{(1 + \beta) \mathbf{E}_{\rightarrow} - (1 - \beta) \mathbf{E}_{\leftarrow}}{(1 + \beta) \mathbf{E}_{\rightarrow} + (1 - \beta) \mathbf{E}_{\leftarrow}} = \frac{(\mathbf{E}_{\rightarrow} - \mathbf{E}_{\leftarrow}) + \beta(\mathbf{E}_{\rightarrow} + \mathbf{E}_{\leftarrow})}{(\mathbf{E}_{\rightarrow} + \mathbf{E}_{\leftarrow}) + \beta(\mathbf{E}_{\rightarrow} - \mathbf{E}_{\leftarrow})} = \frac{SWR + \beta}{1 + \beta \cdot SWR}$$

SWR (or SWQ) Transformation

$$SWR' = \frac{SWR + \beta}{1 + SWR \cdot \beta} = \frac{SWR + u/c}{1 + SWR \cdot u/c}$$

u_{GROUP} (or u_{PHASE}) Transformation

$$u'_{GROUP}/c = \frac{u_{GROUP}/c + \beta}{1 + u_{GROUP} \cdot \beta/c} = \frac{(u_{GROUP} + u)/c}{1 + u_{GROUP} \cdot u/c^2}$$

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$$u'_{GROUP}/c = \frac{u_{GROUP}/c + \beta}{1 + u_{GROUP} \cdot \beta/c} = \frac{(u_{GROUP} + u)/c}{1 + u_{GROUP} \cdot u/c^2}$$

Both are restatements of hyperbolic trig identity: $\tanh(a+b) = \frac{\tanh(a) + \tanh(b)}{1 + \tanh(a) \cdot \tanh(b)}$ last term is ignorable if both a and b are small

Velocity addition is *non-linear* but *rapidity* addition is always *linear*: $\rho_{a+b} = \rho_a + \rho_b$

Special Relativity and Quantum Mechanics by Ruler and Compass I.

The simplest "molecule": 2 CW Lasers form Minkowski Space-time (and Reciprocally-related) Frame Coordinates



A Making sense of light-wave axiom(s).
Speed of light $c = 299,792,458$ m/s

Einstein Pulse Wave (PW) Axiom: PW speed seen by all observers is c

Using Occam's Razor: Joseph Fourier (1768-1830)

Evenson Continuous Wave (CW) axiom: CW speed for all colors is c

1CW is affected by 1st-order Doppler and Red shifts $\lambda = c/\nu$ of frequency ν and wavenumber k

B How does space-time and per-space-per-time carry light-waves?

$(\text{wavelength } \lambda = \text{period } \tau)$ and/or $(\text{wavenumber } k = \text{frequency } \nu)$

$(\lambda = 1/k \text{ and } \tau = 1/\nu)$ and $(\nu = 1/\tau \text{ and } k = 1/\lambda)$

$(\lambda = \text{meters per wave and } \tau = \text{seconds per wave})$ and $(\nu = \text{waves per meter and } k = \text{waves per second})$

Greek "k" for wavenumber for Kayser (or "kinks")

The "Keyboard of the gods" Or per-space-per-time graphs vs. space-time graphs

Wave-velocity formula: $v_{\text{wave}} = \frac{\text{distance}}{\text{time}} = \frac{\text{wavelength}}{\text{period}} = \frac{\nu}{k}$

Light wave-velocity c (near main topic): $v_{\text{wave}} = c = \frac{\lambda}{\tau} = \frac{\nu}{k} = 299,792,458 \text{ m/s}$

Dimensionless Light wave-velocity $c/c=1$: $\frac{v_{\text{wave}}}{c} = \frac{\lambda}{c\tau} = \frac{\nu}{ck} = 1$

C Doppler Shift in per-space-per-time

Atom traveling along wave sees less wave "hits"/sec. (that is: Doppler red-shift)

Atom traveling against wave sees more wave "hits"/sec. (that is: Doppler blue-shift)

Other Doppler versions: $\lambda'/\lambda = \nu/\nu'$ must match this phasor clock/clock-array, too. That's gauge invariance! $c\lambda c' = c'\lambda' c$

Special relativity and quantum mechanics are very much a story of the geometry of light-wave motion

D Clarify Evenson's CW Axiom (All colors go c) by Doppler effects

Alice tries to fool Bob that she's shining a 600THz laser. (Bob's unaware she's moving really fast...)

Q1: Can Bob tell it's a "phony" 600THz by measuring his received wavelength?

Q2: If so, what "phony" λ does Bob see?

Answer to Q2 is C: the one with $\lambda/\nu = c/c = 1$. If he sees Green 600THz then he measures $\lambda = 0.5 \mu\text{m}$. If he sees Red 300THz then he measures $\lambda = 1.0 \mu\text{m}$.

Answer to Q1 is NO! CW light carries no birth-certificate!

All colors go $c = \lambda\nu = c/\lambda \cdot \lambda = c$

Then Evenson's axiom holds: $v = ck$

E Easy Doppler-shift and Rapidity calculation

Doppler ratio: $(R/S) = \frac{v_{\text{observer}}}{v_{\text{source}}}$

rapidity: $\rho_{RS} = \log_e(R/S)$ or $(R/S) = e^{\rho_{RS}}$

Definition of Rapidity: $(R/S) = \frac{v_{\text{observer}}}{v_{\text{source}}}$ is time-reversed $(A/B) = \frac{v_A}{v_B}$

Bob-Alice Doppler ratio: $(R/A) = \frac{v_B}{v_A} = \frac{c}{v_A} = \frac{1}{\beta_A}$

Carla-Alice Doppler ratio: $(C/A) = \frac{v_B}{v_A} = \frac{c}{v_A} = \frac{1}{\beta_A}$

Carla-Bob Doppler ratio: $(C/B) = \frac{v_B}{v_A} = \frac{c}{v_A} = \frac{1}{\beta_A}$

Carla-Alice rapidity: $\rho_{CA} = \log_e(C/A) = \log_e \frac{1}{\beta_A}$

Carla-Bob rapidity: $\rho_{CB} = \log_e(C/B) = \log_e \frac{1}{\beta_A}$

Galileo's Revenge (part 1): Rapidity adds just like Galilean velocity

Galileo's Revenge (part 2): Phasor angular velocity adds just like Galilean velocity

F 1 CW Laser-phasor Wave Function

Dimensionless Light wave-velocity $c/c=1$: $\frac{v_{\text{wave}}}{c} = \frac{\lambda}{c\tau} = \frac{\nu}{ck} = 1$

laser-phasors $\psi(x,t)$: $\psi = A e^{i(kx - \omega t)} = A \cos(kx - \omega t) + i A \sin(kx - \omega t)$

Real $\psi = \text{Re } \psi$ and **Imaginary $\psi = \text{Im } \psi$**

Wavelength $\lambda = 2\pi/k = 1/k$ and **Period $T = 2\pi/\omega = 1/\nu$**

Wavenumber $k = 1/\lambda = 2\pi/\lambda$ and **Angular frequency $\omega = 2\pi\nu$**

How do we measure space and time with light waves? Use 1CW laser-phasors for a phase-based theory

G 2 CWs Interfering in Space-Time

Right-moving CW $e^{i(kx - \omega t)}$ and **Left-moving CW $e^{i(-kx - \omega t)}$**

Sum of Wave Phasor Array: $\psi(x,t) = e^{i(kx - \omega t)} + e^{i(-kx - \omega t)} = 2 \cos(kx) e^{-i\omega t}$

Group Beat: $\cos(kx) \cos(\omega t)$

Minkowski Space-Time Grid: Shows the relationship between space and time axes for two frames.

H 2 Doppler shifted CWs Interfering in Space-Time

Right-directed CW $e^{i(kx - \omega t)}$ and **Left-directed CW $e^{i(-kx - \omega t)}$**

Doppler blue shifted to 1200THz and **Doppler red shifted to 300THz**

2CW per-Spacetime Plot: Shows the interference pattern in the (x, ct) plane.

2CW Minkowski-Spacetime Grid: Shows the grid for the two frames.

Relativity - Using light's own wave-like nature to better understand special relativity and quantum mechanics

I Thales Mean Geometry (600BCE) helps "Relativity"

Thales showed a circle diameter subtends a right angle with any circle point P

Per-Time ω axis and **Per-Space ck axis**

Transformed/Per-Time ω' axis and **Transformed/Per-Space ck' axis**

Arithmetic Mean Bobshp = $(1+4)/2 = 5/2$

Difference Mean Bobshp = $(4-1)/2 = 3/2$

J Comparing Longitudinal relativity parameter: Rapidity $\rho = \log_e(\text{Doppler Shift})$ to a Transverse relativity parameter: Stellar aberration angle σ

Observer fixed below star sees it directly overhead. **Observer going u sees star at angle σ in u direction.**

Stellar aberration angle σ : $c \tanh \rho = u = c \sin \sigma$

Epstein's trick is to turn a hyperbolic form $ct = \sqrt{(ct')^2 - (x')^2}$ into a circular form: $\sqrt{(ct')^2 + (x')^2} = (ct')$

Then everything (and everybody) always goes speed c through (x', ct') space!

K Table of 12 wave parameters (includes inverses) for relativity

Table of 12 wave parameters: A table with columns for parameter, symbol, and value.

Ocean's Sword ($\lambda/c=3/5$): A diagram showing the relationship between wave parameters and the speed of light.

...and values for $u/c=3/5$

Geometry applies to (x, y) space-space to (k_x, k_y) per-space-per-space to (x, ct) space-time

L Optical wave guide relativistic geometry aided by Occam's Sword

Relativistic mode with near- c $V_{\text{group}}=c/2$ and $V_{\text{phase}}=2c$ (Low dispersion): Shows wave propagation in a waveguide.

Example of near-cut-off mode with low $V_{\text{group}}=c/3$ and high $V_{\text{phase}}=2c$ (High dispersion): Shows wave propagation in a waveguide.

M Per-space Per-space Geometry

Per-space Per-space Geometry: A diagram showing the geometry of space and space-time.

[Link to pdf version of Part I online](#)

Note: When printed at their optimal resolution, each poster is 7 feet across!

