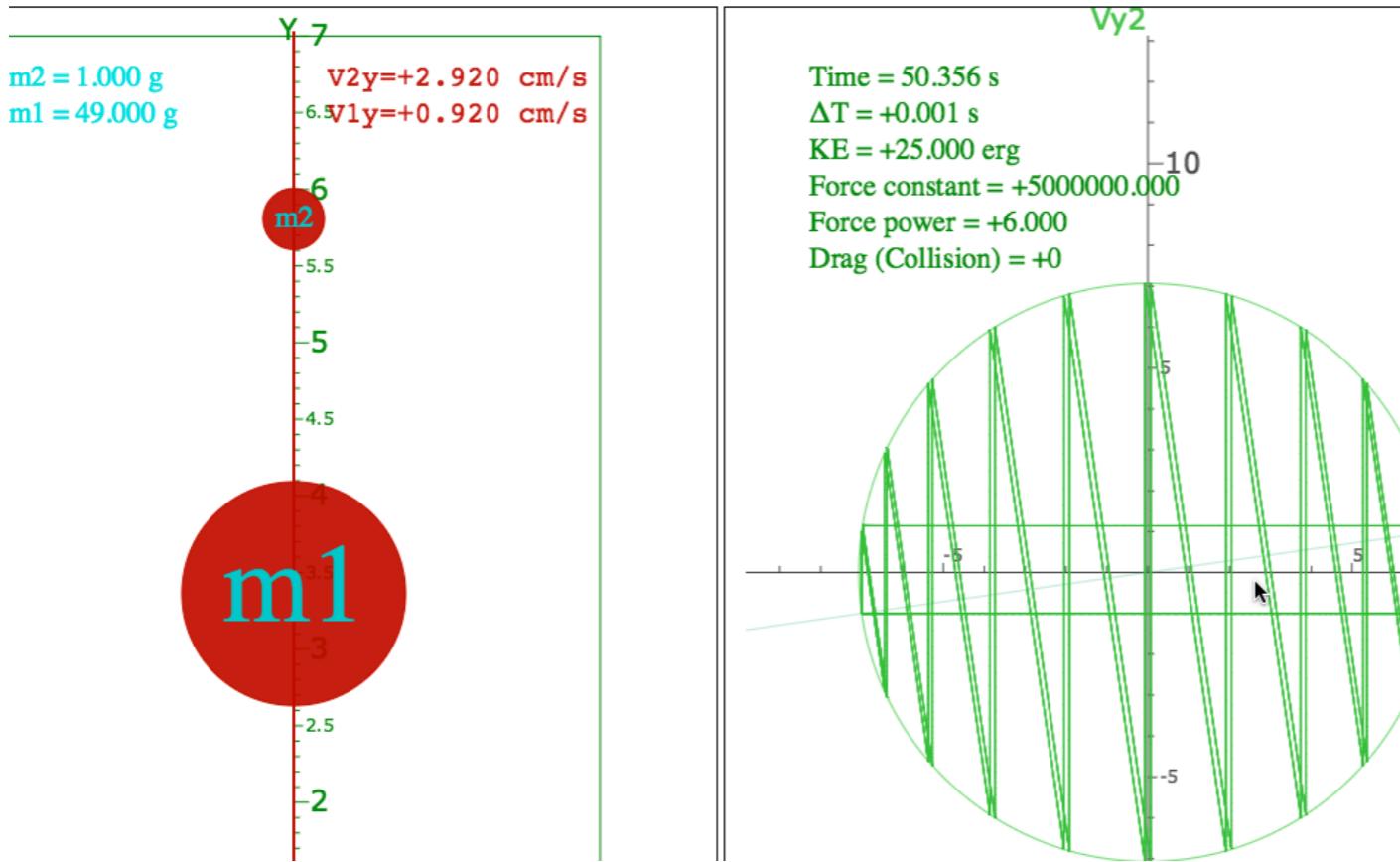


Pseudo-Rotations

Exercise 4.1 Estrangian plot in text Fig. 4.2 (Details on p.64 of Lect. 3to4x. See p.131, too.) has mass ratio $M_1/m_2 = 49/1$ and has nearly (but not quite) periodic path plot. (Let the pen-mass be $m_2=1$ here.)

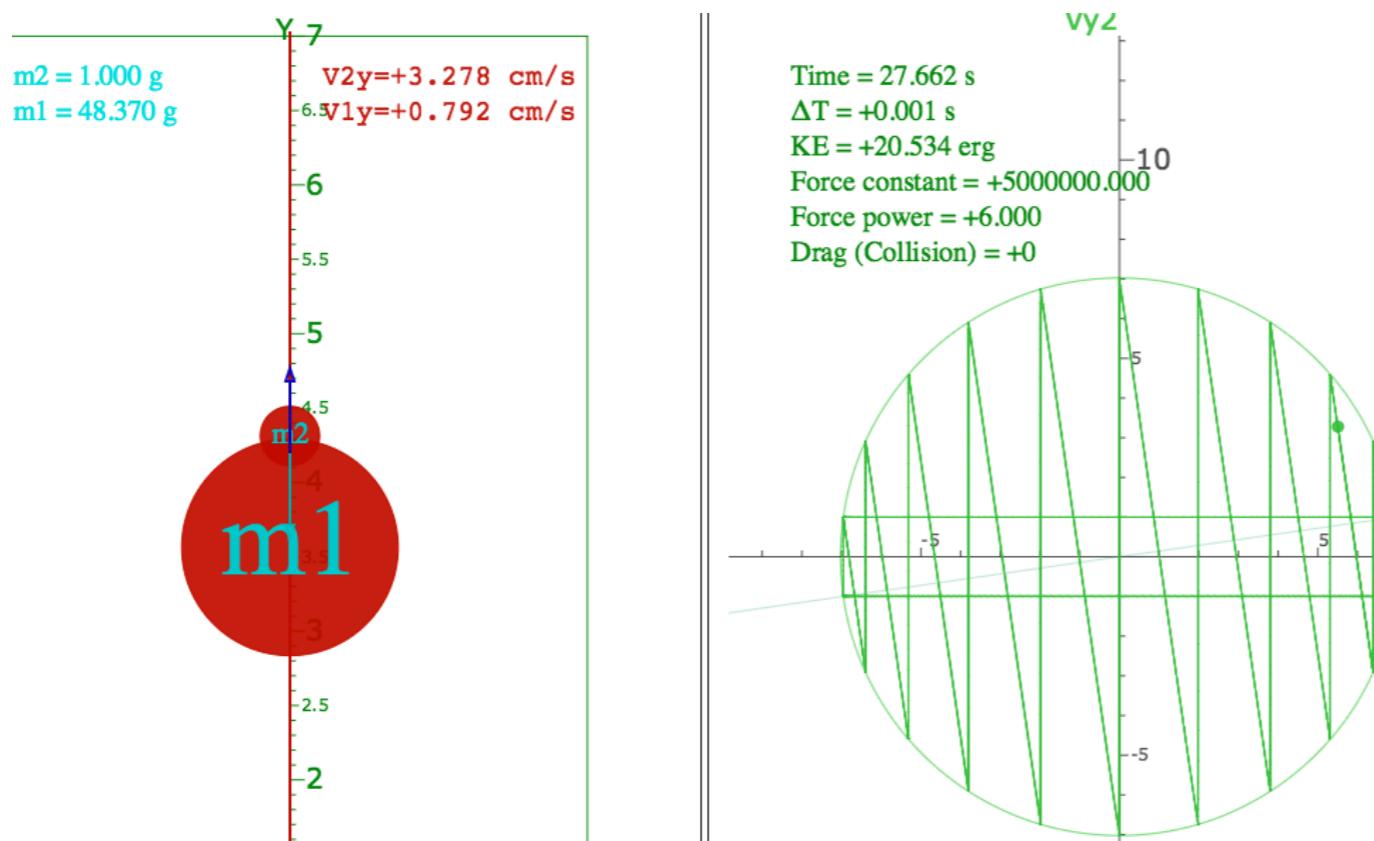


Derive a closed formula for value of $M_1 = 48.37\ldots$ (to at least 7 figures) having *exactly* periodic behavior.

Simplest formula should relate to tangent of a desired Estrangian rotation half-angle $\theta/2$ for mass M_1 .

Changing to $M_1 = 48.37$ gives more nearly periodic paths shown below. (Seems perfect but it's not.)

(Experiment using BounceIt on web. <http://www.uark.edu/ua/modphys/markup/BounceItWeb.php>)



Solution to Pseudo-Rotations

Solving the mass ratio equation (5.10b) for m_1/m_2 in terms of angle θ :

$$\cos\theta \equiv \left(\frac{m_1 - m_2}{M}\right) \quad \text{and} : \quad \sin\theta \equiv \left(\frac{2\sqrt{m_1 m_2}}{M}\right)$$

This will simplify to:

$$\cos\theta = \frac{m_1 - m_2}{m_1 + m_2} = \frac{\frac{m_1}{m_2} - 1}{\frac{m_1}{m_2} + 1} \quad \text{and} : \quad \frac{m_1}{m_2} = \frac{1 + \cos\theta}{1 - \cos\theta} = \frac{\cos^2 \frac{\theta}{2}}{\sin^2 \frac{\theta}{2}} = \cot^2 \frac{\theta}{2} \quad \text{or} : \quad \cot \frac{\theta}{2} = \sqrt{\frac{m_1}{m_2}}$$

Given angle $\theta = \pi/11$ or $\theta/2 = \pi/22$ and $m_2=1$ predicts

$$m_1 = \tan^2 \frac{\pi}{22} = 48.3741500787$$

Pseudo-Vibrations

Exercise 4.2 On p.50-55 of *Lect. 5to6*, is shown pseudo-harmonic motion of the large mass $M=50\text{kg}$ attacked on either side by a pair of tiny masses $m=0.1\text{kg}$ each traveling back and forth in a range of $Y_0=3.5\text{m}$ at an average speed of 20m/sec . The calculation seems to come up a period about $\sqrt{3}$ times too big for mass M . Explain what was overlooked and derive an improved formula for the period.

Example of oscillator with opposing Isothermal potentials

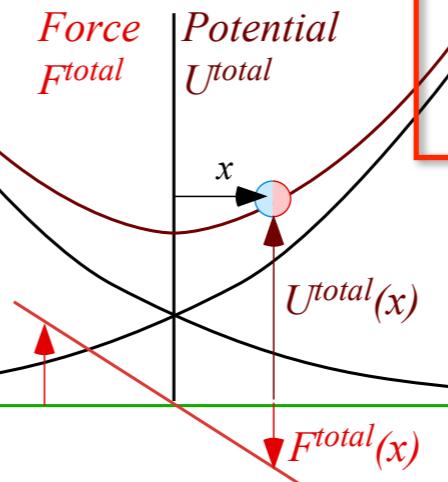
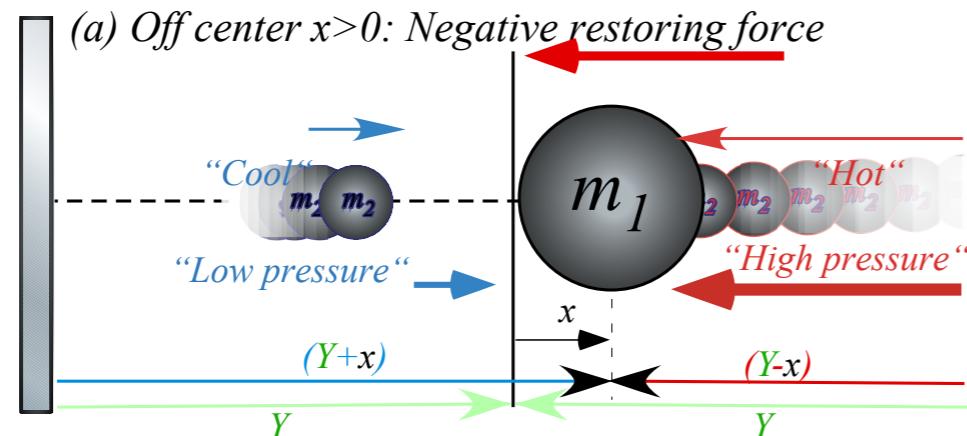
1D-Isothermal Force Law (assume v_2 is constant for all Y):

$$F = \frac{m_2 v_2^2}{Y} = \frac{\text{const.}}{Y}$$

$$F^{phys} = \frac{m_2 v_2^2}{Y} = -\frac{\Delta U}{\Delta Y}$$

implies :

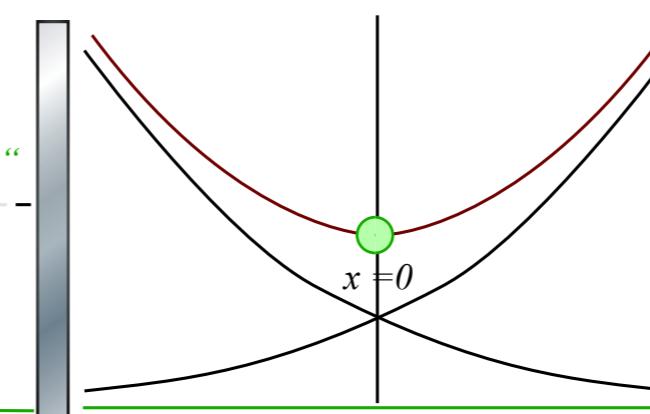
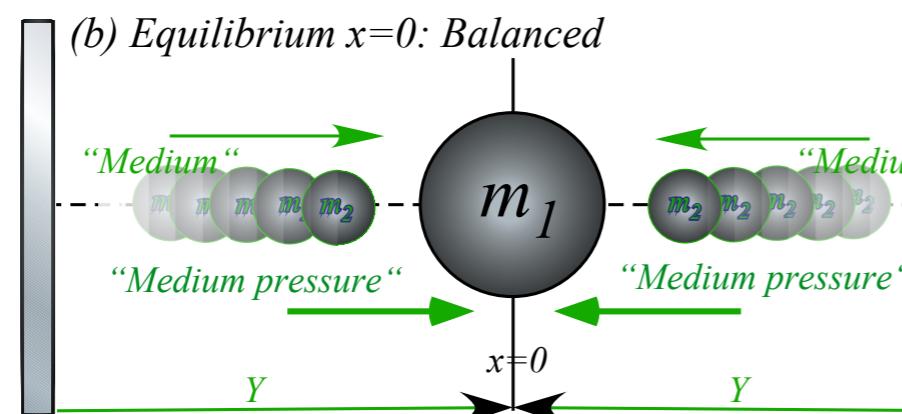
$$U = -m_2 v_2^2 \ln(Y)$$



$$F = \frac{m_2 (\bar{v}_2^{IN} Y_0)^2}{Y^3} = \frac{\text{const.}}{Y^3}$$

1D-Adiabatic Force Law

Fig. 5.3
Unit 1



Two opposing 1D-Adiabatic Force Law fields

$$F^{total} = \frac{f}{(Y_0 + x)^3} - \frac{f}{(Y_0 - x)^3} = f \left[\frac{1}{Y_0^3} - \frac{3x}{Y_0^4} + \frac{6x^2}{Y_0^5} - \frac{10x^3}{Y_0^6} + \dots \right] - f \left[\frac{1}{Y_0^3} + \frac{3x}{Y_0^4} + \frac{6x^2}{Y_0^5} + \frac{10x^3}{Y_0^6} + \dots \right]$$

$$(Y_0 + x)^{-3} = Y_0^{-3} - 3xY_0^{-4} + \frac{3 \cdot 4}{2} x^2 Y_0^{-5} - \frac{3 \cdot 4 \cdot 5}{2 \cdot 3} x^3 Y_0^{-6} \dots \quad (Y_0 - x)^{-3} = Y_0^{-3} + 3xY_0^{-4} + \frac{3 \cdot 4}{2} x^2 Y_0^{-5} + \frac{3 \cdot 4 \cdot 5}{2 \cdot 3} x^3 Y_0^{-6} \dots$$

$$(Y_0 + x)^n = Y_0^n + nY_0^{n-1}x + \frac{n(n-1)}{2} Y_0^{n-2}x^2 + \frac{n(n-1)(n-2)}{2 \cdot 3} Y_0^{n-3}x^3 + \frac{n(n-1)(n-2)(n-3)}{2 \cdot 3 \cdot 4} Y_0^{n-4}x^4 \dots$$

Example of oscillator with opposing Isothermal potentials

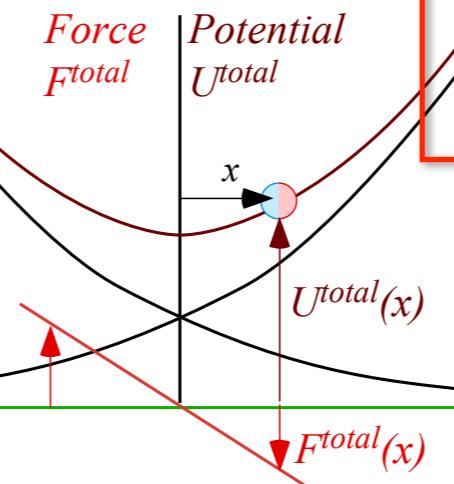
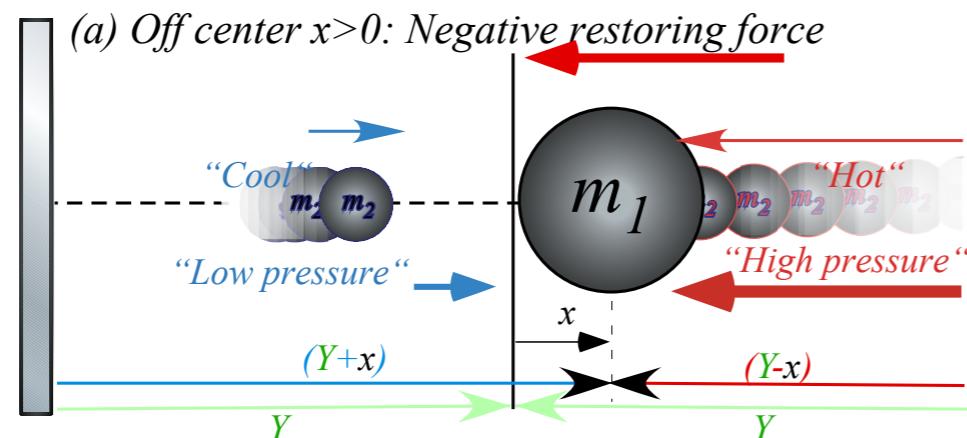
1D-Isothermal Force Law (assume v_2 is constant for all Y):

$$F = \frac{m_2 v_2^2}{Y} = \frac{\text{const.}}{Y}$$

$$F^{phys} = \frac{m_2 v_2^2}{Y} = -\frac{\Delta U}{\Delta Y}$$

implies :

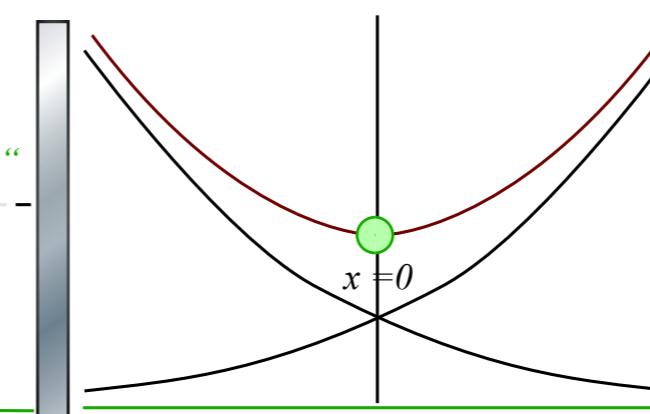
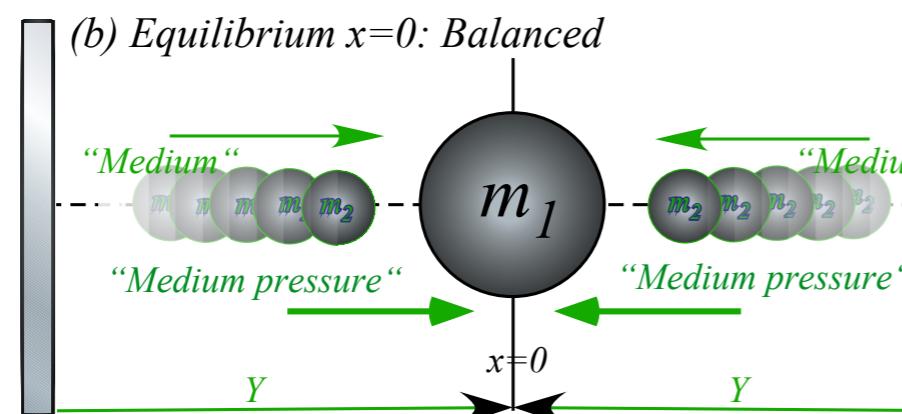
$$U = -m_2 v_2^2 \ln(Y)$$



$$F = \frac{m_2 (\bar{v}_2^{IN} Y_0)^2}{Y^3} = \frac{\text{const.}}{Y^3}$$

1D-Adiabatic Force Law

Fig. 5.3
Unit 1



Two opposing 1D-Adiabatic Force Law fields

$$F^{total} = \frac{f}{(Y_0 + x)^3} - \frac{f}{(Y_0 - x)^3} = f \left[\cancel{\frac{1}{Y_0^3}} - \frac{3x}{Y_0^4} + \cancel{\frac{6x^2}{Y_0^5}} - \frac{10x^3}{Y_0^6} + \dots \right] - f \left[\cancel{\frac{1}{Y_0^3}} + \frac{3x}{Y_0^4} + \cancel{\frac{6x^2}{Y_0^5}} + \frac{10x^3}{Y_0^6} + \dots \right]$$

$$(Y_0 + x)^{-3} = Y_0^{-3} - 3xY_0^{-4} + \frac{3 \cdot 4}{2} x^2 Y_0^{-5} - \frac{3 \cdot 4 \cdot 5}{2 \cdot 3} x^3 Y_0^{-6} \dots \quad (Y_0 - x)^{-3} = Y_0^{-3} + 3xY_0^{-4} + \frac{3 \cdot 4}{2} x^2 Y_0^{-5} + \frac{3 \cdot 4 \cdot 5}{2 \cdot 3} x^3 Y_0^{-6} \dots$$

$$(Y_0 + x)^n = Y_0^n + n Y_0^{n-1} x + \frac{n(n-1)}{2} Y_0^{n-2} x^2 + \frac{n(n-1)(n-2)}{2 \cdot 3} Y_0^{n-3} x^3 + \frac{n(n-1)(n-2)(n-3)}{2 \cdot 3 \cdot 4} Y_0^{n-4} x^4 \dots$$

Example of oscillator with opposing Isothermal potentials

1D-Isothermal Force Law (assume v_2 is constant for all Y):

$$F = \frac{m_2 v_2^2}{Y} = \frac{\text{const.}}{Y}$$

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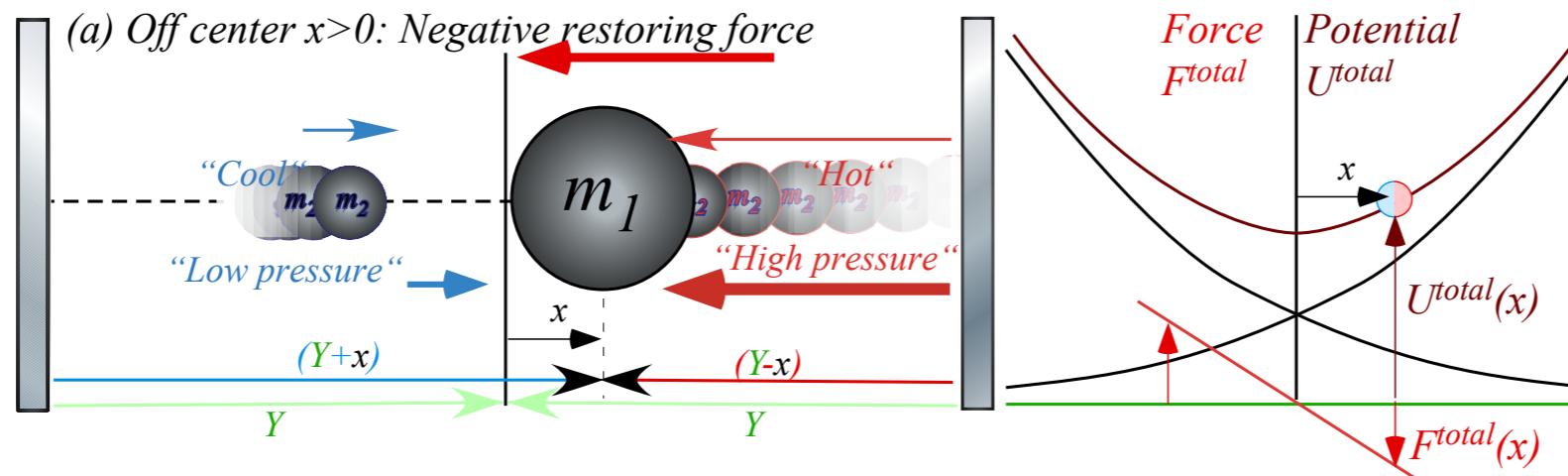
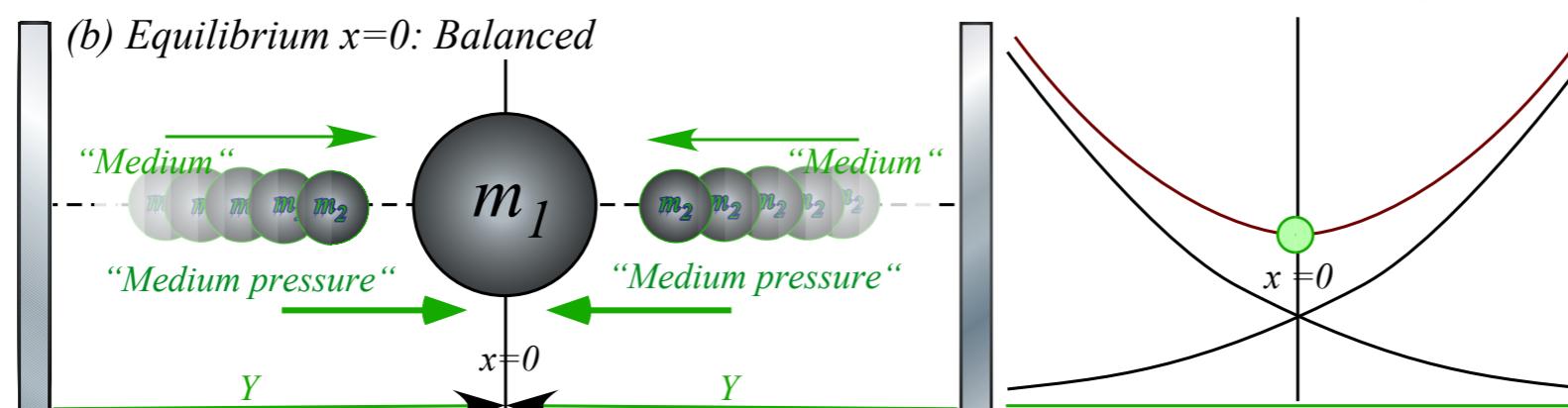


Fig. 5.3
Unit 1



Anharmonic oscillator terms...

Harmonic oscillator term

Two opposing 1D-Adiabatic Force Law fields

$$F^{total} = \frac{f}{(Y_0 + x)^3} - \frac{f}{(Y_0 - x)^3} = f \left[-\frac{3x}{Y_0^4} - \frac{10x^3}{Y_0^6} \dots \right] + f \left[-\frac{3x}{Y_0^4} - \frac{10x^3}{Y_0^6} + \dots \right] = \underbrace{-2f \frac{3x}{Y_0^4}}_{\text{Harmonic oscillator term}} - \underbrace{2f \frac{10x^3}{Y_0^6}}_{\text{Anharmonic oscillator terms...}} - \dots$$

where:

Adiabatic harmonic force constant : $k = 2 \cdot \underbrace{3f}_{\text{Harmonic oscillator term}} / Y_0^4 = 2 \cdot 3m_2 (\bar{v}_2^{IN})^2 / Y_0^2$.

$$f = m_2 (\bar{v}_2^{IN} Y_0)^2 = \text{const.}$$

$$F = \frac{m_2 (\bar{v}_2^{IN} Y_0)^2}{Y^3} = \frac{f}{Y^3}$$

1D-Adiabatic Force Law

Example of oscillator with opposing Isothermal potentials

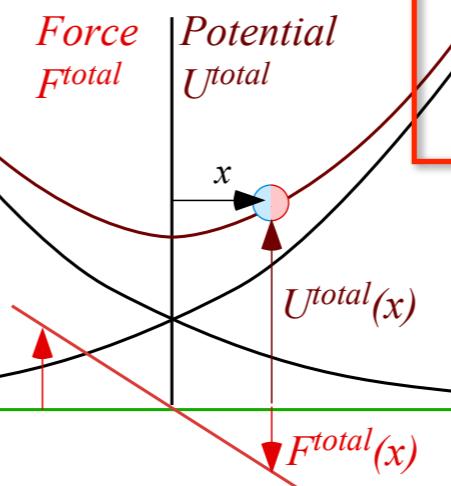
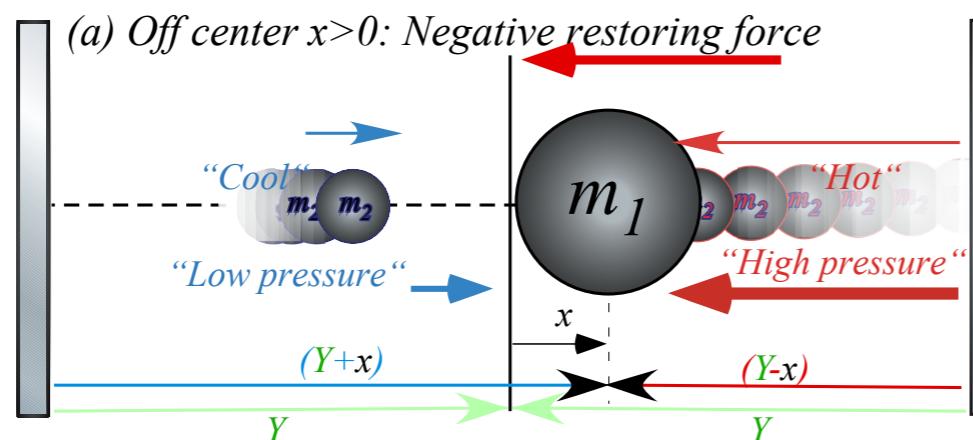
1D-Isothermal Force Law (assume v_2 is constant for all Y):

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implies :

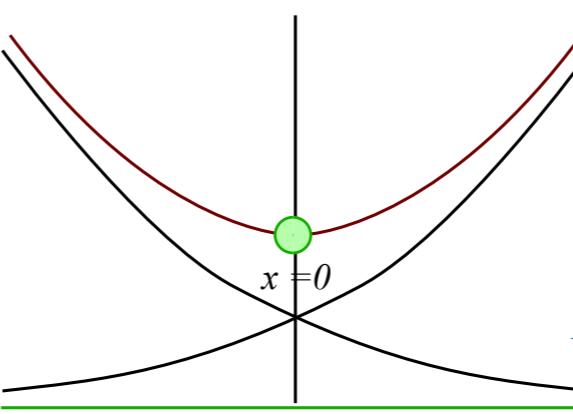
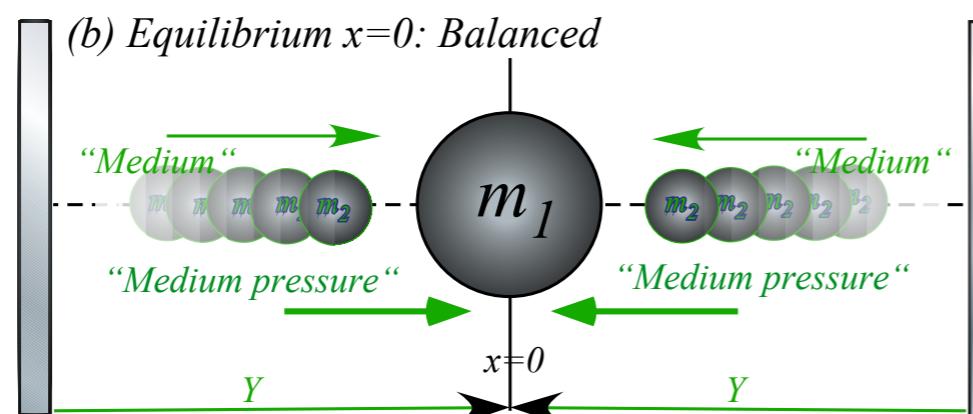
$$U = -m_2 v_2^2 \ln(Y)$$



$$F = \frac{m_2 (\bar{v}_2^{IN} Y_0)^2}{Y^3} = \frac{\text{const.}}{Y^3}$$

1D-Adiabatic Force Law

Fig. 5.3
Unit 1



Anharmonic oscillator terms...

Two opposing 1D-Adiabatic Force Law fields

$$F^{total} = \frac{f}{(Y_0 + x)^3} - \frac{f}{(Y_0 - x)^3} = f \left[-\frac{3x}{Y_0^4} - \frac{10x^3}{Y_0^6} \dots \right] + f \left[-\frac{3x}{Y_0^4} - \frac{10x^3}{Y_0^6} + \dots \right] = -2f \underbrace{\frac{3x}{Y_0^4}}_{\text{Harmonic oscillator term}} - 2f \underbrace{\frac{10x^3}{Y_0^6}}_{\text{Anharmonic oscillator terms...}} - \dots$$

where:

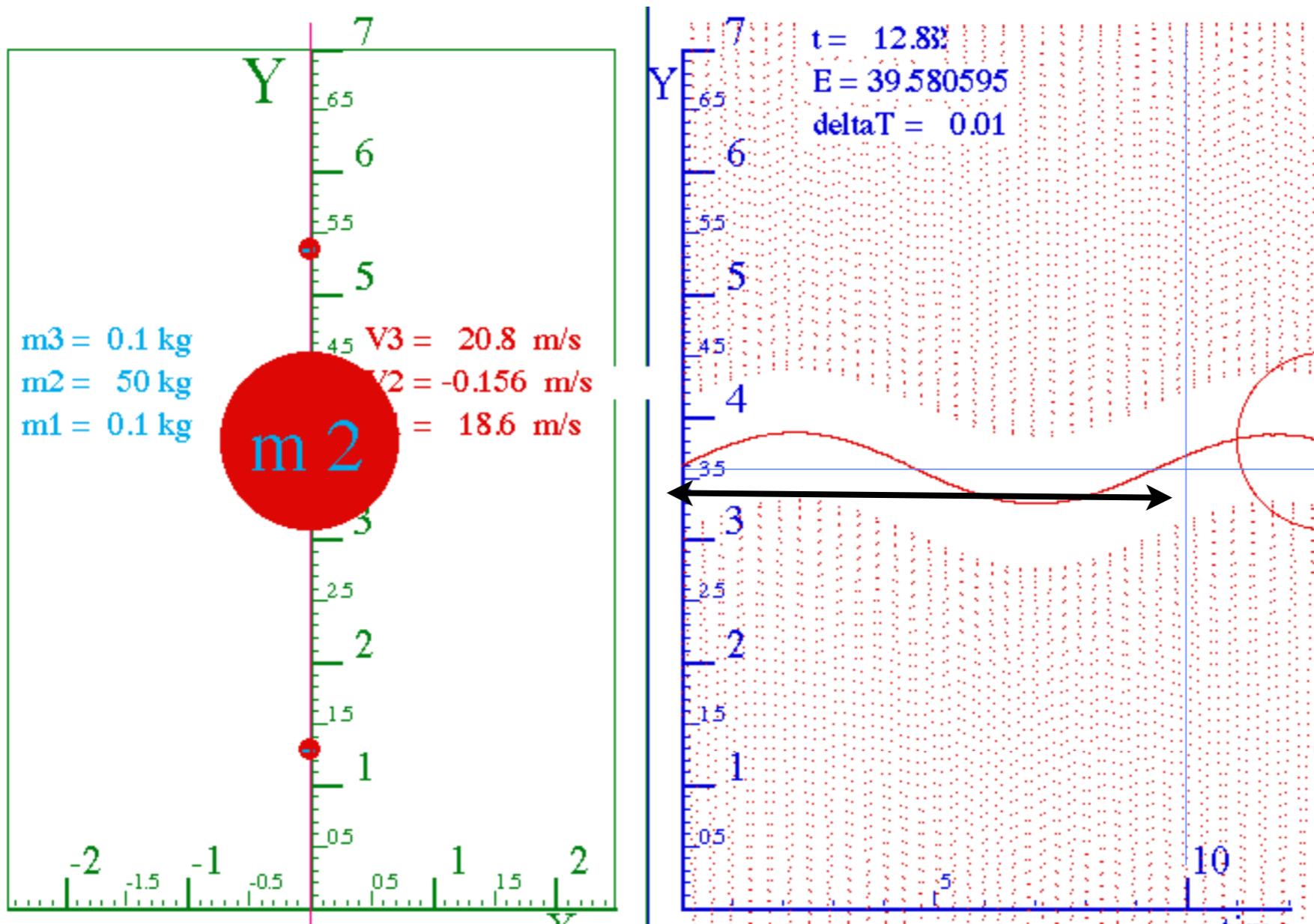
Adiabatic harmonic force constant : $k = 2 \cdot 3f/Y_0^4 = 2 \cdot 3m_2(\bar{v}_2^{IN})^2/Y_0^2$.

$$f = m_2(\bar{v}_2^{IN} Y_0)^2 = \text{const.}$$

$$\text{HO } \not\propto \text{frequency: } \omega = \sqrt{\frac{k}{m_1}} = \sqrt{\frac{2 \cdot 3m_2}{m_1} \frac{\bar{v}_2^{IN}}{Y_0}} = 2\pi\nu$$

Adiabatic frequency is $\sqrt{3}$ times faster than isothermal frequency

Switch
 $m_1=m_3$
 with
 m_2
 to match
 formula



BounceIt Superball Collision Web Simulator: 1:500:1 mass ratios (Small Amplitude)

Sample problem: Compute period given $m_1=50$, $m_2=0.1=m_3$, $v_2=20$, $Y_0=3.5$

Adiabatic Period is $\sqrt{3}$ times shorter:

$$\tau = 2\pi \sqrt{\frac{m_1}{2 \cdot 3 m_2} \frac{Y_0}{v_2}} = 6.28 \sqrt{\frac{50}{2 \cdot 3 \cdot (0.1)} \frac{3.5}{20}} \\ = 10.03 \quad \text{That's only a little too big!}$$

$$\text{Adiabatic Period : } \tau = \frac{1}{v} = 2\pi \sqrt{\frac{m_1}{k}} = 2\pi \sqrt{\frac{m_1}{2 \cdot 3 m_2} \frac{Y_0}{v_2}}$$

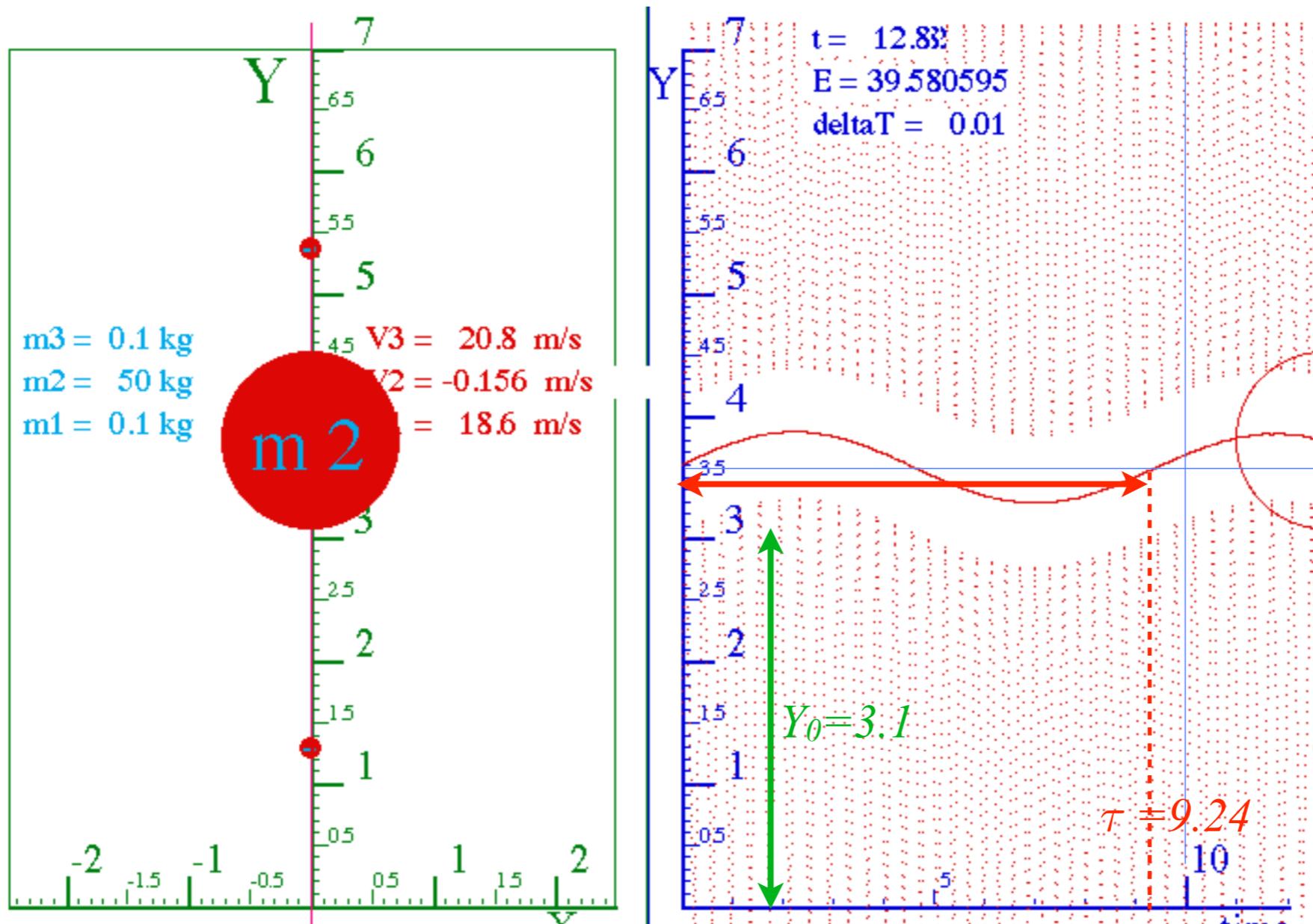
Adiabatic Frequency

$$\text{HO } \angle \text{frequency: } \omega = \sqrt{\frac{k}{m_1}} = \sqrt{\frac{2 \cdot 3 m_2}{m_1} \frac{v_2}{Y_0}} = 2\pi v$$

Fig. 5.3
 Unit 1

Simulation of
 an adiabatic case

*Switch
 $m_1=m_3$
with
 m_2
to match
formula*



BounceIt Superball Collision Web Simulator: 1:500:1 mass ratios (Small Amplitude)

Sample problem: Compute period given $m_1=50$, $m_2=0.1=m_3$, $v_2=20$, $Y_0=3.5$

Perhaps we
overestimated
 Y_0 range....
Let's try $Y_0=3.1$

Adiabatic Period is $\sqrt{3}$ times shorter:

$$\tau = 2\pi \sqrt{\frac{m_1}{2 \cdot 3 m_2} \frac{Y_0}{v_2}} = 6.28 \sqrt{\frac{50}{2 \cdot 3 \cdot (0.1)} \frac{3.1}{20}} = 9.24$$

That's better!

Adiabatic Period: $\tau = \frac{1}{v} = 2\pi \sqrt{\frac{m_1}{k}} = 2\pi \sqrt{\frac{m_1}{2 \cdot 3 m_2} \frac{Y_0}{v_2}}$

Adiabatic Frequency

HO ↘ frequency: $\omega = \sqrt{\frac{k}{m_1}} = \sqrt{\frac{2 \cdot 3 m_2}{m_1} \frac{v_2}{Y_0}} = 2\pi v$

Fig. 5.3
Unit 1

Simulation of
an adiabatic case