

# *Group Theory in Quantum Mechanics*

## *Lecture 2 (1.15.15)*

### *Quantum amplitudes, analyzers, and axioms*

*(Quantum Theory for Computer Age - Ch. 1 of Unit 1 )*

*(Principles of Symmetry, Dynamics, and Spectroscopy - Sec. 1-2 of Ch. 1 )*

*Review: “Abstraction” of bra and ket vectors from a Transformation Matrix*

*Introducing scalar and matrix products*

*Planck's energy and N-quanta (Cavity/Beam wave mode)*

*Did Max Planck Goof? What's 1-photon worth?*

*Feynman amplitude axiom 1*

*What comes out of a beam sorter channel or branch-b?*

*Sample calculations*

*Feynman amplitude axioms 2-3*

*Beam analyzers: Sorter-unsorters*

*The “Do-Nothing” analyzer*

*Feynman amplitude axiom 4*

*Some “Do-Something” analyzers*

*Sorter-counter, Filter, 1/2-wave plate, 1/4-wave plate*

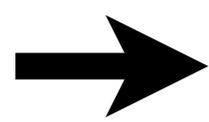
[http://www.uark.edu/ua/modphys/pdfs/QTCA\\_Pdfs/QTCA\\_Text\\_2013/QTCA\\_Unit\\_1\\_Ch\\_1\\_2013.pdf](http://www.uark.edu/ua/modphys/pdfs/QTCA_Pdfs/QTCA_Text_2013/QTCA_Unit_1_Ch_1_2013.pdf)

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*See also 2005 site:*

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# “Abstraction” of bra and ket vectors from a Transformation Matrix

*Ket or column vectors*

*Bra or row vectors*

Given  
Transformation  
Matrix  $T_{m,n'}$  :

$$\begin{pmatrix} \langle x|x' \rangle & \langle x|y' \rangle \\ \langle y|x' \rangle & \langle y|y' \rangle \end{pmatrix} = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix}$$

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$$|x'\rangle = \begin{pmatrix} \langle x|x' \rangle \\ \langle y|x' \rangle \end{pmatrix} = \begin{pmatrix} \cos\theta \\ \sin\theta \end{pmatrix}$$

*Abstracting ket  $|n'\rangle$  state vectors  
from  
Transformation Matrix*

$$T_{m,n'} = \langle m | n' \rangle$$

# “Abstraction” of bra and ket vectors from a Transformation Matrix

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Given Transformation Matrix  $T_{m,n'}$  :

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$$\langle x| = \begin{pmatrix} \langle x|x' \rangle & \langle x|y' \rangle \end{pmatrix} = \begin{pmatrix} \cos\theta & -\sin\theta \end{pmatrix}$$

$$\langle y| = \begin{pmatrix} \langle y|x' \rangle & \langle y|y' \rangle \end{pmatrix} = \begin{pmatrix} \sin\theta & \cos\theta \end{pmatrix}$$

*Abstracting bra  $\langle m|$  state vectors*

*from*

*Transformation Matrix*

$$T_{m,n'} = \langle m|n' \rangle$$

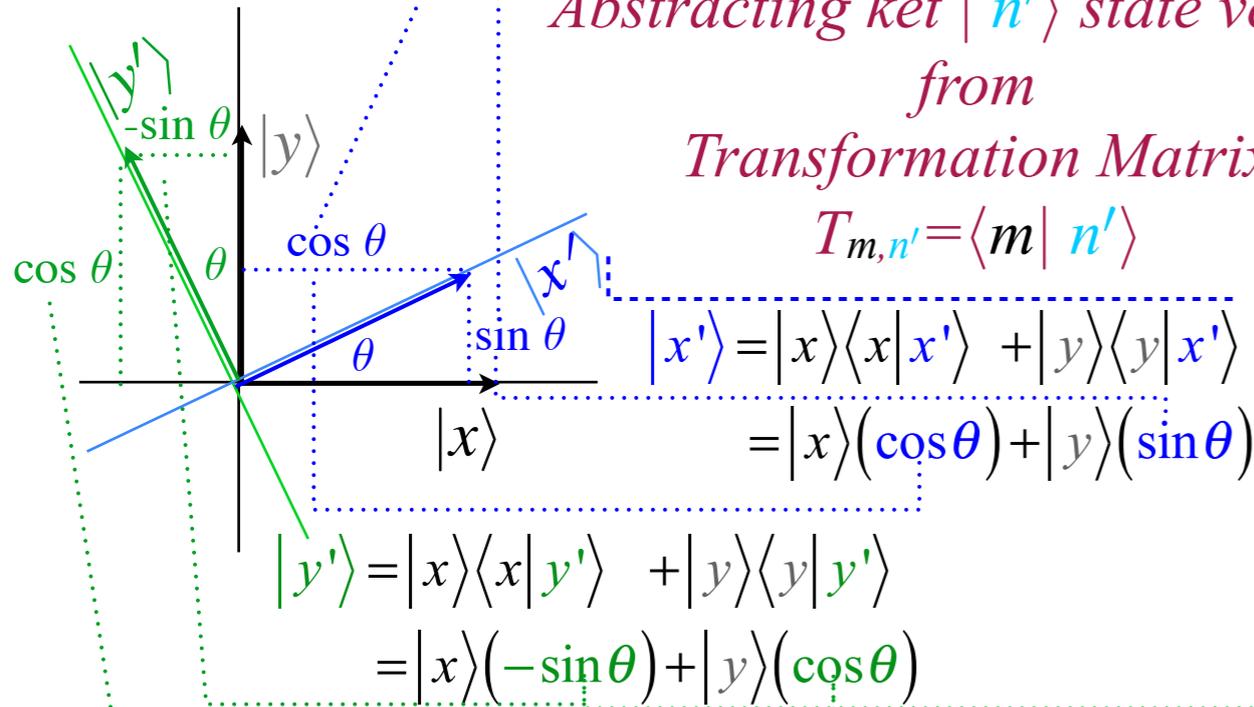
$$|x'\rangle = \begin{pmatrix} \langle x|x' \rangle \\ \langle y|x' \rangle \end{pmatrix} = \begin{pmatrix} \cos\theta \\ \sin\theta \end{pmatrix}, |y'\rangle = \begin{pmatrix} \langle x|y' \rangle \\ \langle y|y' \rangle \end{pmatrix} = \begin{pmatrix} -\sin\theta \\ \cos\theta \end{pmatrix}$$

*Abstracting ket  $|n'\rangle$  state vectors*

*from*

*Transformation Matrix*

$$T_{m,n'} = \langle m|n' \rangle$$



$$|x'\rangle = |x\rangle\langle x|x'\rangle + |y\rangle\langle y|x'\rangle$$

$$= |x\rangle(\cos\theta) + |y\rangle(\sin\theta)$$

$$|y'\rangle = |x\rangle\langle x|y'\rangle + |y\rangle\langle y|y'\rangle$$

$$= |x\rangle(-\sin\theta) + |y\rangle(\cos\theta)$$

$$\langle y| = \langle y|x'\rangle\langle x'| + \langle y|y'\rangle\langle y'|$$

$$= (\sin\theta)\langle x'| + (\cos\theta)\langle y'|$$

$$\langle x| = \langle x|x'\rangle\langle x'| + \langle x|y'\rangle\langle y'|$$

$$= (\cos\theta)\langle x'| + (-\sin\theta)\langle y'|$$

$(\theta=+30^\circ)$ -Rotated kets  $\{|x'\rangle, |y'\rangle\}$  or  $\{\mathbf{x}', \mathbf{y}'\}$  represented in page-aligned  $\{|x\rangle, |y\rangle\}$  basis.

$(\theta=-30^\circ)$ -Rotated bras  $\{\langle x|, \langle y|\}$  or  $\{\mathbf{x}, \mathbf{y}\}$  represented in page-aligned  $\{\langle x'|, \langle y'|\}$  basis.

*Ket vector algebra has the order of  $T_{m,n'}$  transposed*

*Bra vector algebra has the same order as  $T_{m,n'}$*

$$|x'\rangle = |x\rangle\langle x|x'\rangle + |y\rangle\langle y|x'\rangle = |x\rangle(\cos\theta) + |y\rangle(\sin\theta)$$

$$|y'\rangle = |x\rangle\langle x|y'\rangle + |y\rangle\langle y|y'\rangle = |x\rangle(-\sin\theta) + |y\rangle(\cos\theta)$$

$$\langle x| = \langle x|x'\rangle\langle x'| + \langle x|y'\rangle\langle y'| = (\cos\theta)\langle x'| + (-\sin\theta)\langle y'|$$

$$\langle y| = \langle y|x'\rangle\langle x'| + \langle y|y'\rangle\langle y'| = (\sin\theta)\langle x'| + (\cos\theta)\langle y'|$$

➔ *Review: “Abstraction” of bra and ket vectors from a Transformation Matrix*  
*Introducing scalar and matrix products*

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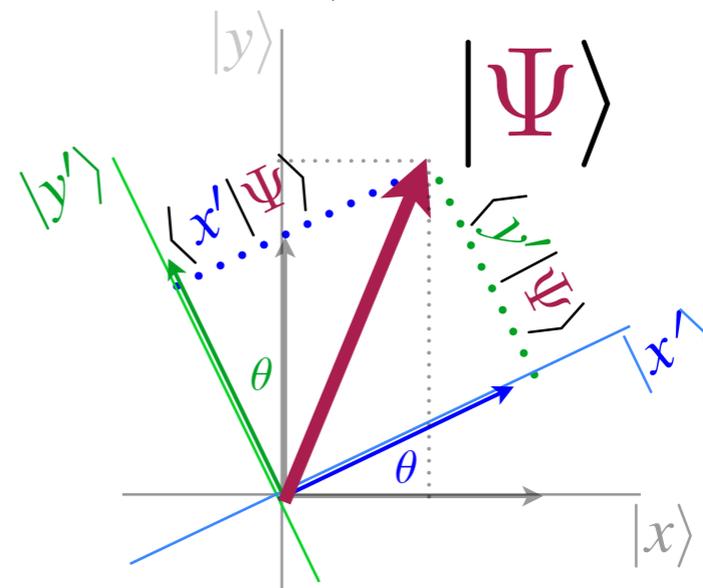
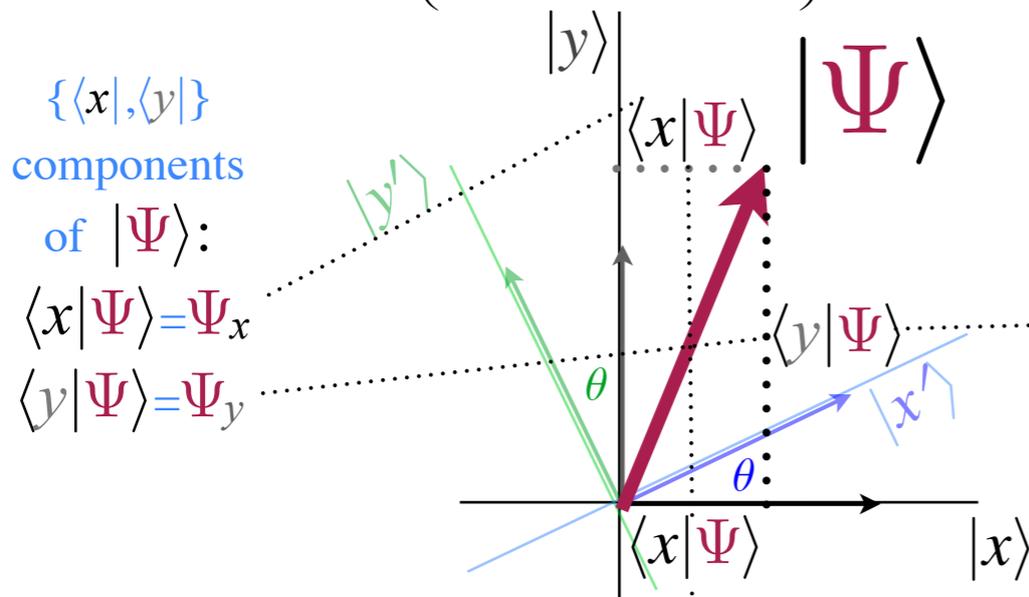
*Feynman amplitude axiom 4*

*Some “Do-Something” analyzers*

*Sorter-counter, Filter, 1/2-wave plate, 1/4-wave plate*

Transformation matrix  $T_{m,n'} = \langle m | n' \rangle$  is array of dot or *scalar products* (dot products) of unit vectors or *direction cosines*.

$$\begin{pmatrix} \langle x | x' \rangle & \langle x | y' \rangle \\ \langle y | x' \rangle & \langle y | y' \rangle \end{pmatrix} = \begin{pmatrix} \cos\theta & -\sin\theta \\ \sin\theta & \cos\theta \end{pmatrix} = \begin{pmatrix} (\mathbf{x} \cdot \mathbf{x}') & (\mathbf{x} \cdot \mathbf{y}') \\ (\mathbf{y} \cdot \mathbf{x}') & (\mathbf{y} \cdot \mathbf{y}') \end{pmatrix}$$



Any state  $|\Psi\rangle$  can be expanded in any basis  $\{ \langle x |, \langle y | \}$ , or  $\{ \langle x' |, \langle y' | \}$ , ...etc.

$$|\Psi\rangle = |x\rangle \langle x | \Psi \rangle + |y\rangle \langle y | \Psi \rangle = |x'\rangle \langle x' | \Psi \rangle + |y'\rangle \langle y' | \Psi \rangle$$

Transformation matrix  $T_{m,n'}$  relates  $\{ \langle x | \Psi \rangle, \langle y | \Psi \rangle \}$  amplitudes to  $\{ \langle x' | \Psi \rangle, \langle y' | \Psi \rangle \}$ .

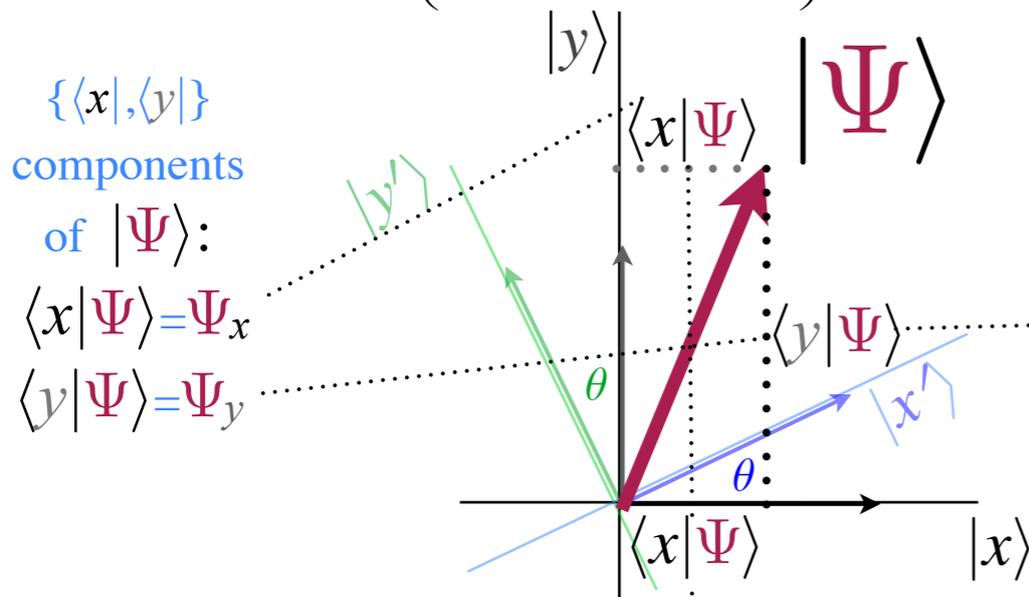
$$\begin{pmatrix} \langle x | \Psi \rangle \\ \langle y | \Psi \rangle \end{pmatrix} = \begin{pmatrix} \langle x | x' \rangle & \langle x | y' \rangle \\ \langle y | x' \rangle & \langle y | y' \rangle \end{pmatrix} \begin{pmatrix} \langle x' | \Psi \rangle \\ \langle y' | \Psi \rangle \end{pmatrix} \quad \text{or:} \quad \begin{pmatrix} \Psi_x \\ \Psi_y \end{pmatrix} = \begin{pmatrix} \langle x | x' \rangle & \langle x | y' \rangle \\ \langle y | x' \rangle & \langle y | y' \rangle \end{pmatrix} \begin{pmatrix} \Psi_{x'} \\ \Psi_{y'} \end{pmatrix}$$

Hybrid  
Gibbs-Dirac  
notation  
(Ug-ly!)

**Proof:**  $\langle x | = \langle x | x' \rangle \langle x' | + \langle x | y' \rangle \langle y' |$  implies:  $\langle x | \Psi \rangle = \langle x | x' \rangle \langle x' | \Psi \rangle + \langle x | y' \rangle \langle y' | \Psi \rangle$   
 $\langle y | = \langle y | x' \rangle \langle x' | + \langle y | y' \rangle \langle y' |$  implies:  $\langle y | \Psi \rangle = \langle y | x' \rangle \langle x' | \Psi \rangle + \langle y | y' \rangle \langle y' | \Psi \rangle$

Transformation matrix  $T_{m,n'} = \langle m | n' \rangle$  is array of dot or *scalar products* (dot products) of unit vectors or *direction cosines*.

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$\{\langle x |, \langle y | \}$   
components  
of  $|\Psi\rangle$ :

$$\langle x | \Psi \rangle = \Psi_x$$

$$\langle y | \Psi \rangle = \Psi_y$$

$\{\langle x' |, \langle y' | \}$   
components  
of  $|\Psi\rangle$ :

$$\langle x' | \Psi \rangle = \Psi_{x'}$$

$$\langle y' | \Psi \rangle = \Psi_{y'}$$

Any state  $|\Psi\rangle$  can be expanded in any basis  $\{\langle x |, \langle y | \}$ , or  $\{\langle x' |, \langle y' | \}$ , ...etc.

$$|\Psi\rangle = |x\rangle \langle x | \Psi \rangle + |y\rangle \langle y | \Psi \rangle = |x'\rangle \langle x' | \Psi \rangle + |y'\rangle \langle y' | \Psi \rangle$$

Transformation matrix  $T_{m,n'}$  relates  $\{\langle x | \Psi \rangle, \langle y | \Psi \rangle\}$  amplitudes to  $\{\langle x' | \Psi \rangle, \langle y' | \Psi \rangle\}$ .

$$\begin{pmatrix} \langle x | \Psi \rangle \\ \langle y | \Psi \rangle \end{pmatrix} = \begin{pmatrix} \langle x | x' \rangle & \langle x | y' \rangle \\ \langle y | x' \rangle & \langle y | y' \rangle \end{pmatrix} \begin{pmatrix} \langle x' | \Psi \rangle \\ \langle y' | \Psi \rangle \end{pmatrix} \quad \text{or:} \quad \begin{pmatrix} \Psi_x \\ \Psi_y \end{pmatrix} = \begin{pmatrix} \langle x | x' \rangle & \langle x | y' \rangle \\ \langle y | x' \rangle & \langle y | y' \rangle \end{pmatrix} \begin{pmatrix} \Psi_{x'} \\ \Psi_{y'} \end{pmatrix}$$

Hybrid  
Gibbs-Dirac  
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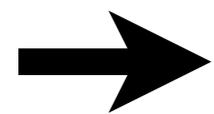
Inverse ( $\dagger = T^* = -1$ ) matrix  $T_{n',m}$  relates  $\{\langle x' | \Psi \rangle, \langle y' | \Psi \rangle\}$  amplitudes to  $\{\langle x | \Psi \rangle, \langle y | \Psi \rangle\}$ .

$$\begin{pmatrix} \langle x' | \Psi \rangle \\ \langle y' | \Psi \rangle \end{pmatrix} = \begin{pmatrix} \langle x' | x \rangle & \langle x' | y \rangle \\ \langle y' | x \rangle & \langle y' | y \rangle \end{pmatrix} \begin{pmatrix} \langle x | \Psi \rangle \\ \langle y | \Psi \rangle \end{pmatrix} \quad \text{or:} \quad \begin{pmatrix} \Psi_{x'} \\ \Psi_{y'} \end{pmatrix} = \begin{pmatrix} \langle x' | x \rangle & \langle x' | y \rangle \\ \langle y' | x \rangle & \langle y' | y \rangle \end{pmatrix} \begin{pmatrix} \Psi_x \\ \Psi_y \end{pmatrix}$$

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(Still Ug-ly!)

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## Planck's energy and $N$ -quanta (*Cavity/Beam of volume $V$ with wave mode of frequency $\omega=2\pi\nu$* )

*Planck axiom:*  $E$ -field energy density  $U$  in cavity/beam mode- $\omega$  is:  $U=N\hbar\omega/V =Nh\nu/V$  ( $N$  “photons”)

$$h=2\pi\hbar=6.6310^{-34}\text{Js}$$

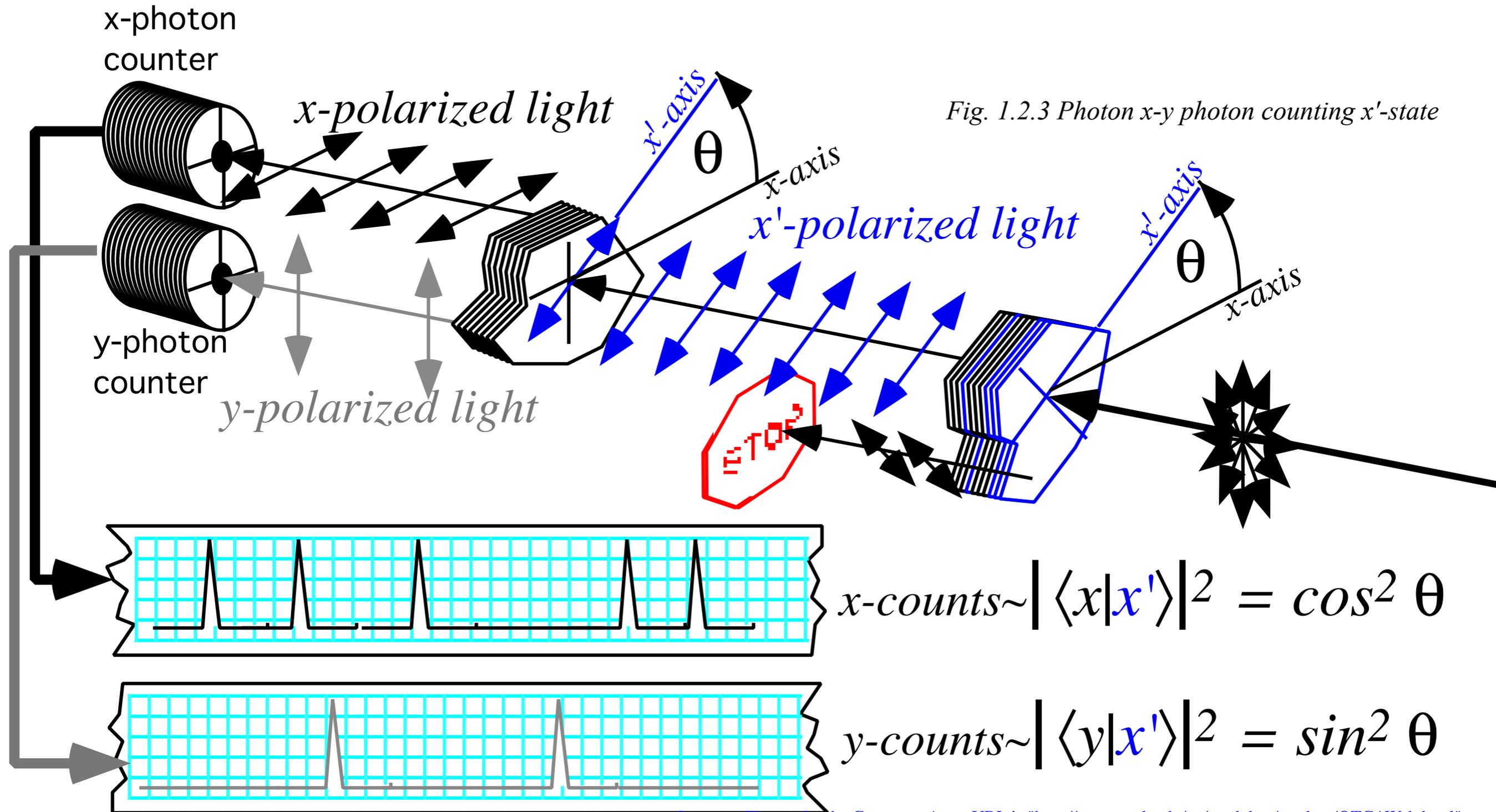
*Planck constant*

**Planck's energy and  $N$ -quanta** (Cavity/Beam of volume  $V$  with wave mode of frequency  $\omega=2\pi\nu$ )

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$E$ -field vector  $(E_x, E_y) = f \cdot (\Psi_x, \Psi_y)$  where *quantum field proportionality constant* is  $f$ .

$$h=2\pi\hbar=6.63 \cdot 10^{-34} \text{Js}$$

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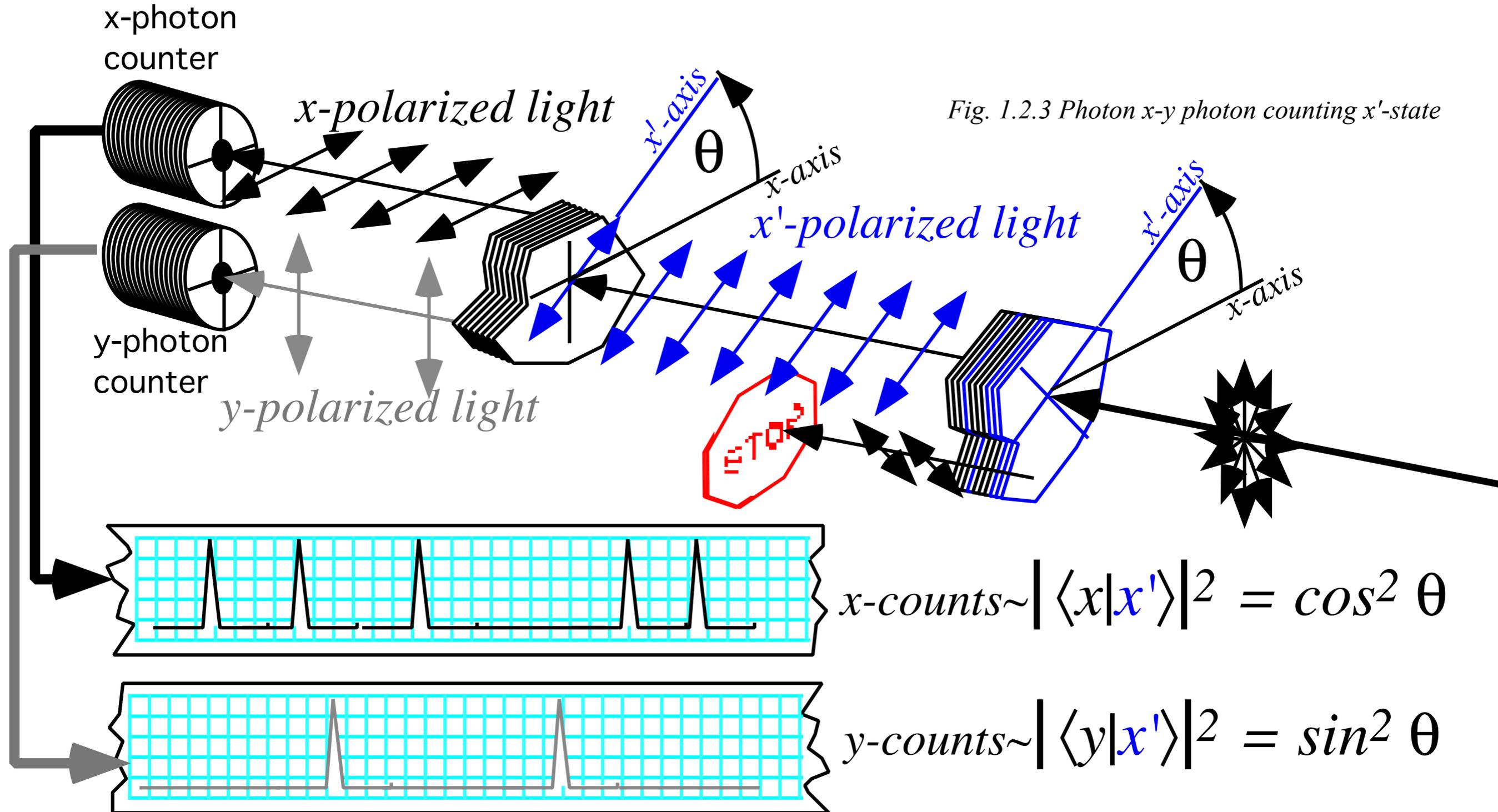


Fig. 1.2.3 Photon x-y photon counting  $x'$ -state

**Planck's energy and  $N$ -quanta** (Cavity/Beam of volume  $V$  with wave mode of frequency  $\omega=2\pi\nu$ )

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$$\begin{pmatrix} E_x \\ E_y \end{pmatrix} = \begin{pmatrix} E_x(0)e^{ikz-i\omega t} \\ E_y(0)e^{ikz-i\omega t} \end{pmatrix} \cong \begin{pmatrix} E_x(0)e^{-i\omega t} \\ E_y(0)e^{-i\omega t} \end{pmatrix} = f \begin{pmatrix} \Psi_x \\ \Psi_y \end{pmatrix}$$

$h=2\pi\hbar=6.63 \cdot 10^{-34} \text{Js}$

*Planck constant*

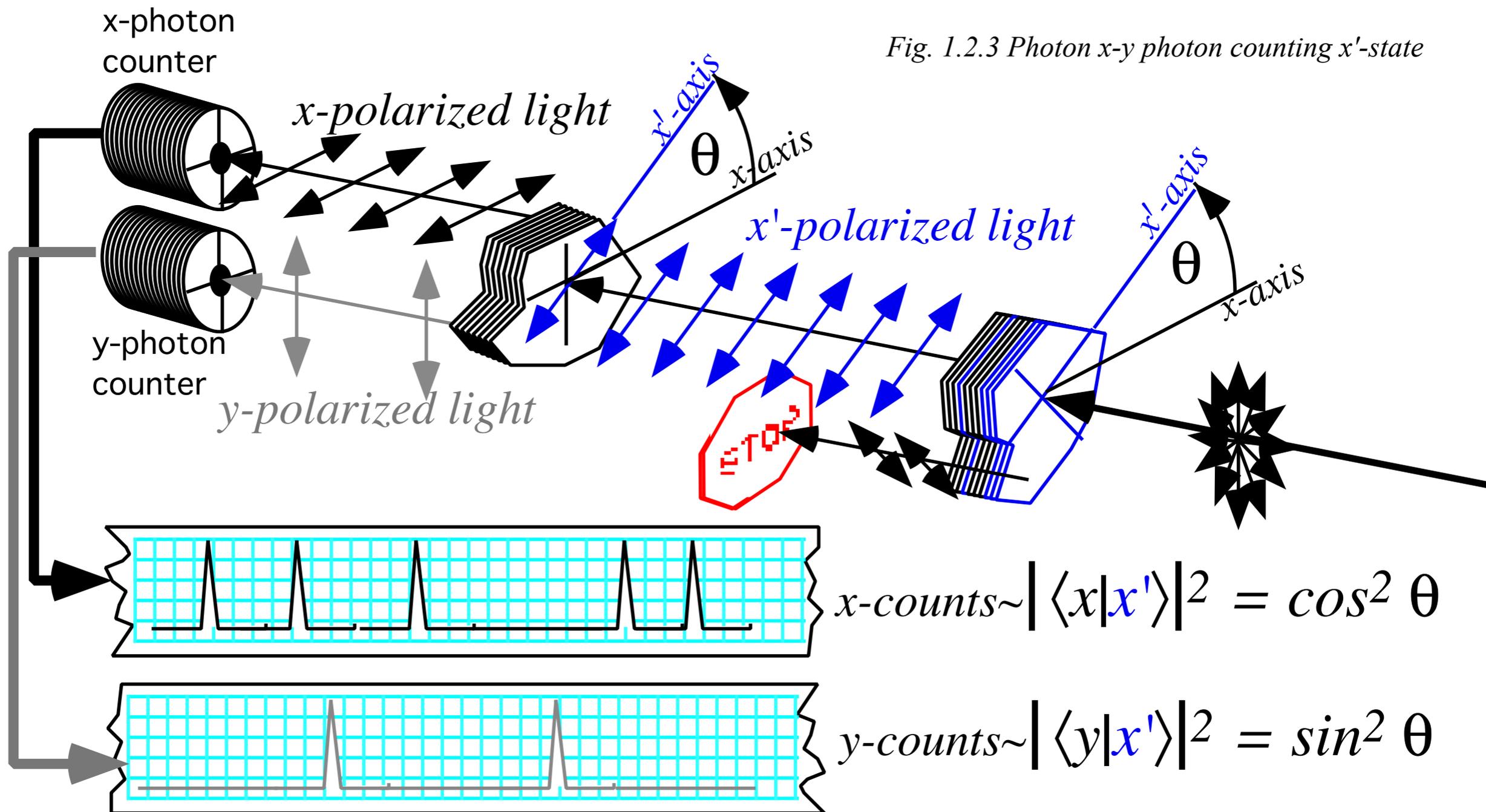


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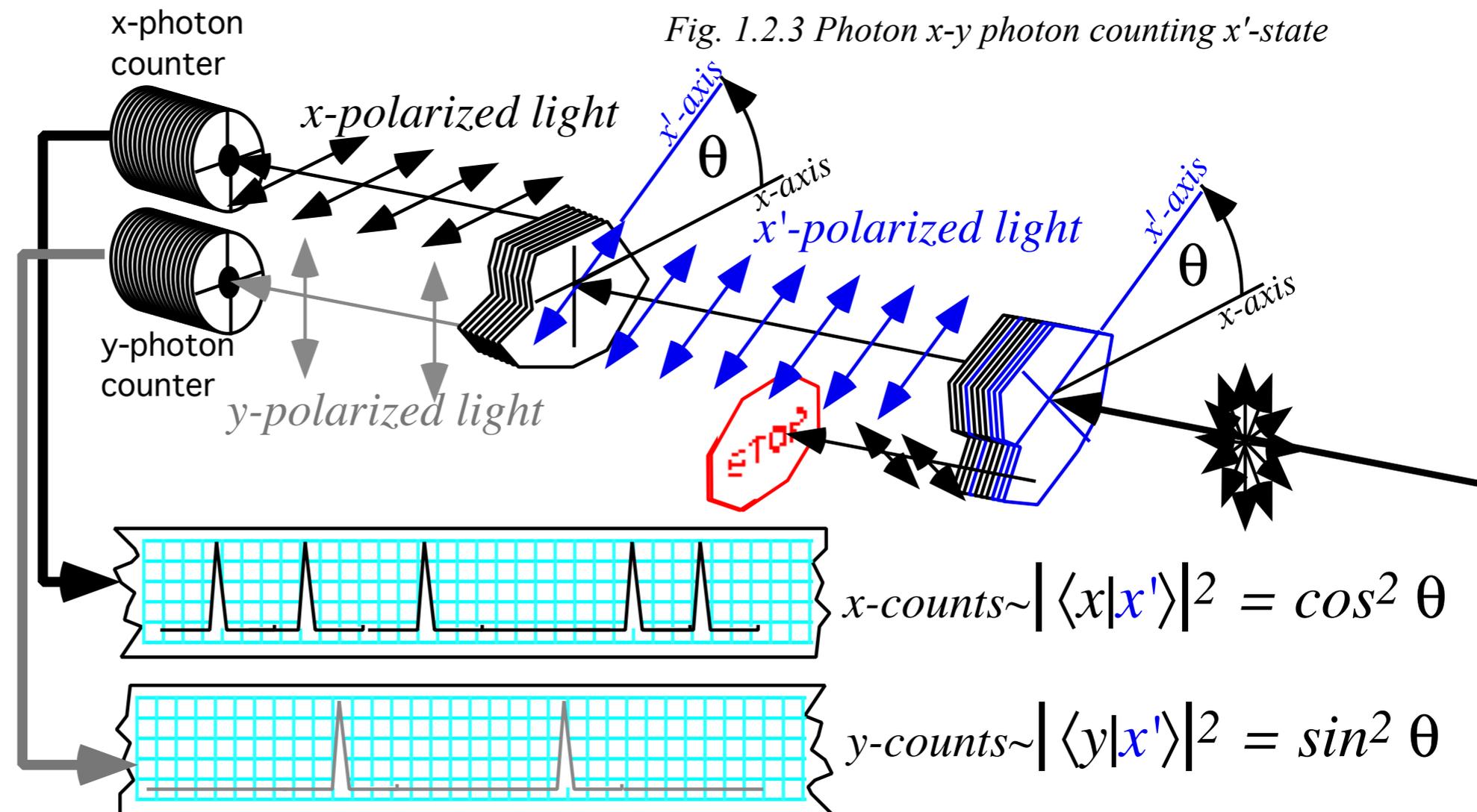
$h=2\pi\hbar=6.626\ 075 \cdot 10^{-34} \text{ Js}$  Planck constant  
 $c=2.997\ 924\ 58 \cdot 10^8 \text{ ms}^{-1}$  Light speed  
 $\epsilon_0=8.854 \cdot 10^{-12} \text{ C}^2 \text{ N}^{-1} \text{ m}^{-2}$  Electrostatic constant

Coulomb constant:  $k=1/4\pi\epsilon_0$   
 $=9 \cdot 10^9 \text{ J/C}$

$\Psi$ -amplitude-squares sum to *exciton-number*  $N$ . (...or *photon-number*  $N$ )

$$N = |\Psi_x|^2 + |\Psi_y|^2 = \langle x | \Psi \rangle^* \langle x | \Psi \rangle + \langle y | \Psi \rangle^* \langle y | \Psi \rangle$$

Quantum  
em wave  
theory



**Planck's energy and  $N$ -quanta** (Cavity/Beam of volume  $V$  with wave mode of frequency  $\omega=2\pi\nu$ )

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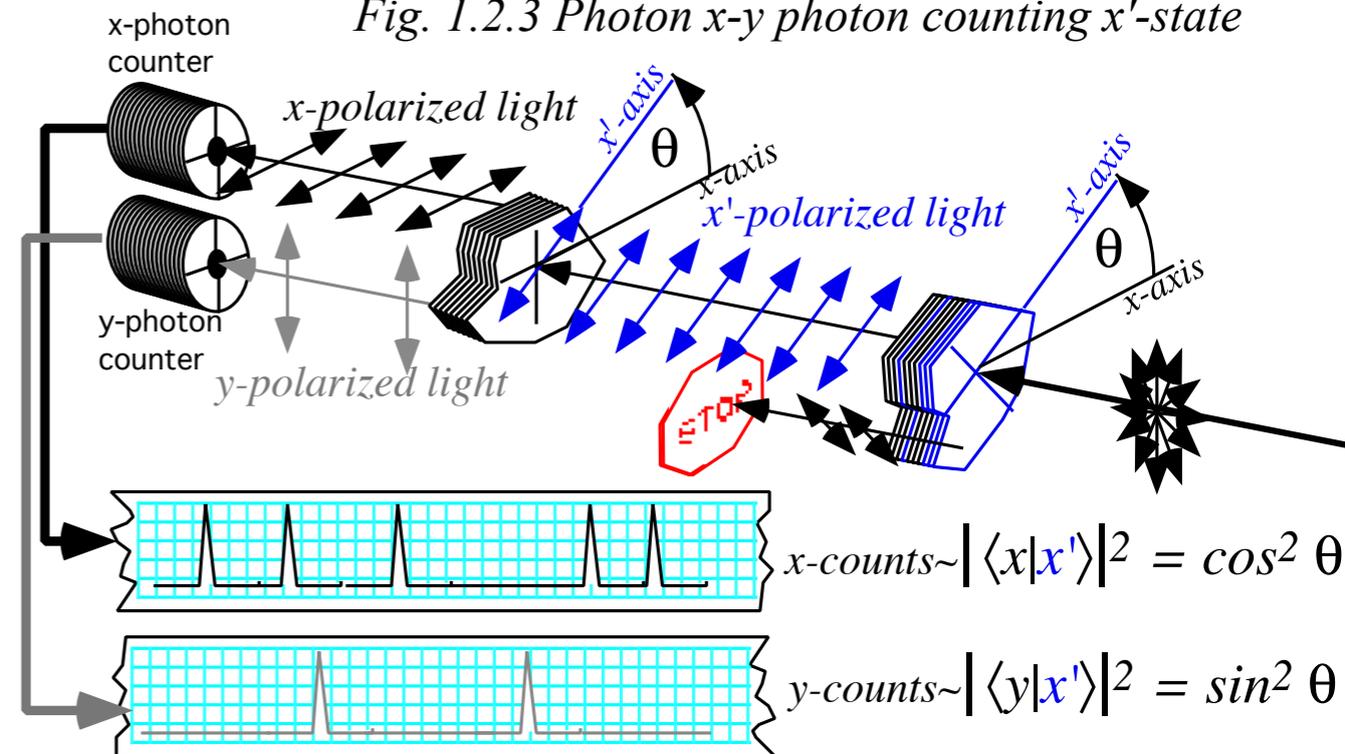
Quantum em wave theory

*Poynting energy flux  $S(\text{J/m}^2\text{s})$  or energy density  $U(\text{J/m}^3)$  of light beam.*

$$S = cU, \text{ where: } U = \epsilon_0 \left( |E_x|^2 + |E_y|^2 \right) = \epsilon_0 \left( E_x^* E_x + E_y^* E_y \right) = \epsilon_0 \left( E_x(0)^2 + E_y(0)^2 \right)$$

Classical em wave theory

Fig. 1.2.3 Photon  $x$ - $y$  photon counting  $x'$ -state



See also 2005 site:

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# Planck's energy and $N$ -quanta (Cavity/Beam of volume $V$ with wave mode of frequency $\omega=2\pi\nu$ )

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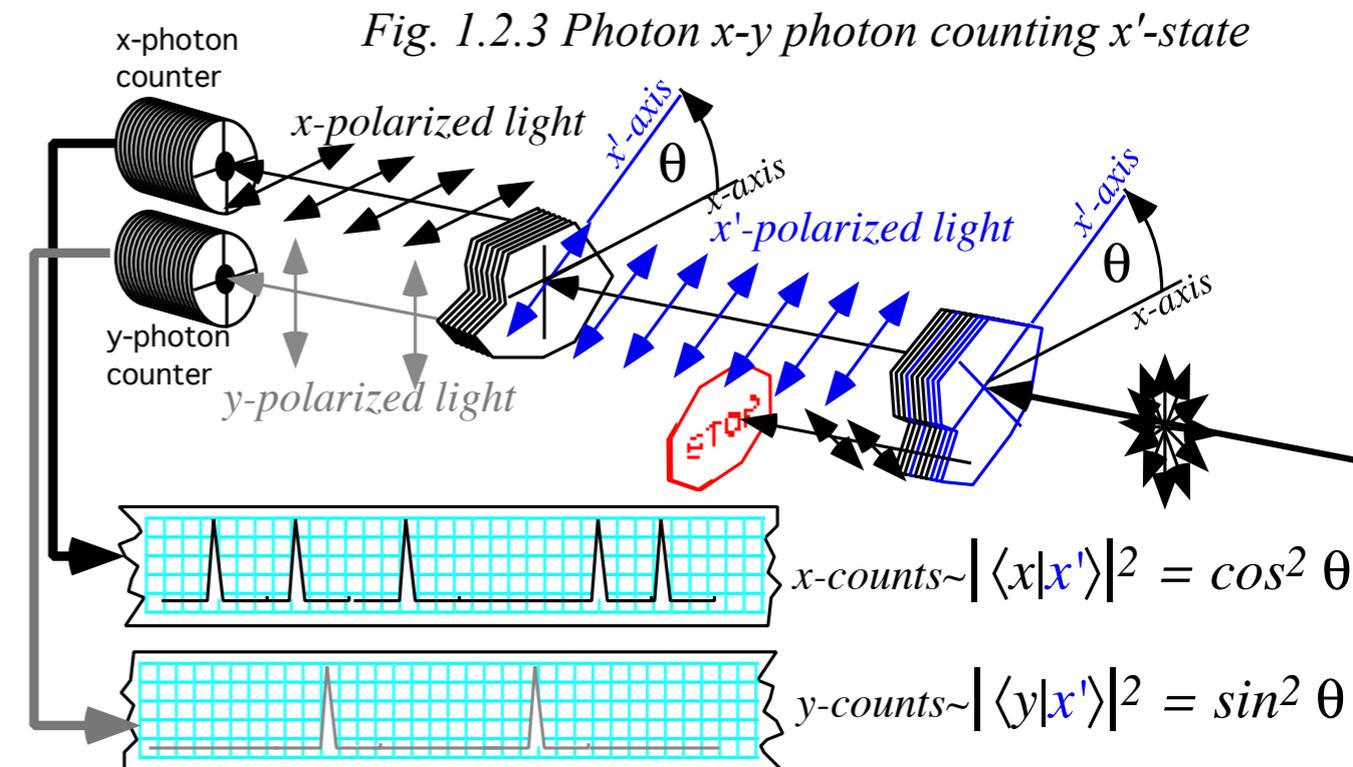
Quantum  
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*Poynting energy flux*  $S(\text{J/m}^2\text{s})$  or *energy density*  $U(\text{J/m}^3)$  of light beam.

$$S = cU, \text{ where: } U = \epsilon_0 \left( |E_x|^2 + |E_y|^2 \right) = \epsilon_0 \left( E_x^* E_x + E_y^* E_y \right) = \epsilon_0 \left( E_x(0)^2 + E_y(0)^2 \right)$$

Classical  
em wave  
theory

Equate  $U$  to Planck's  $N$ -*"photon"* quantum energy density  $N\hbar\omega/V$



# Planck's energy and $N$ -quanta (Cavity/Beam of volume $V$ with wave mode of frequency $\omega=2\pi\nu$ )

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$h=2\pi\hbar=6.626\ 075 \cdot 10^{-34} \text{ Js}$  Planck constant  
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Quantum  
em wave  
theory

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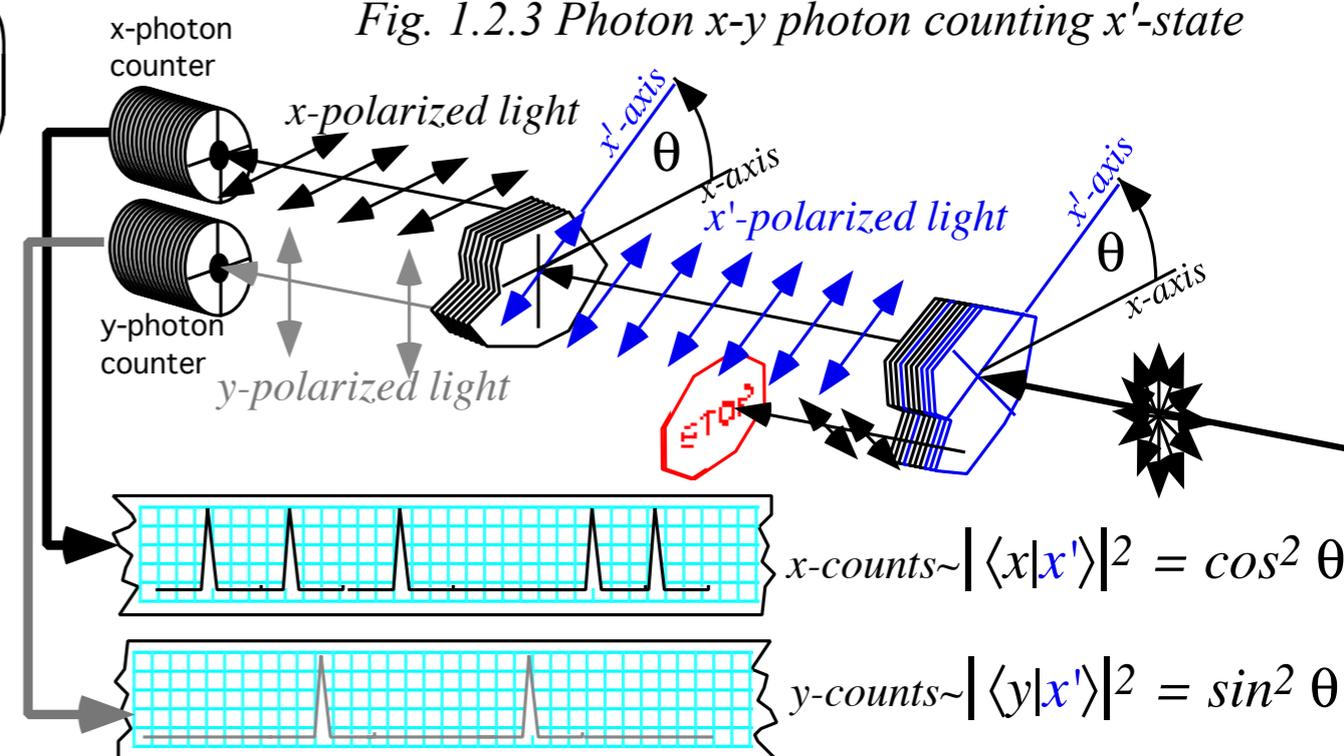
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Classical  
em wave  
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Quantum  
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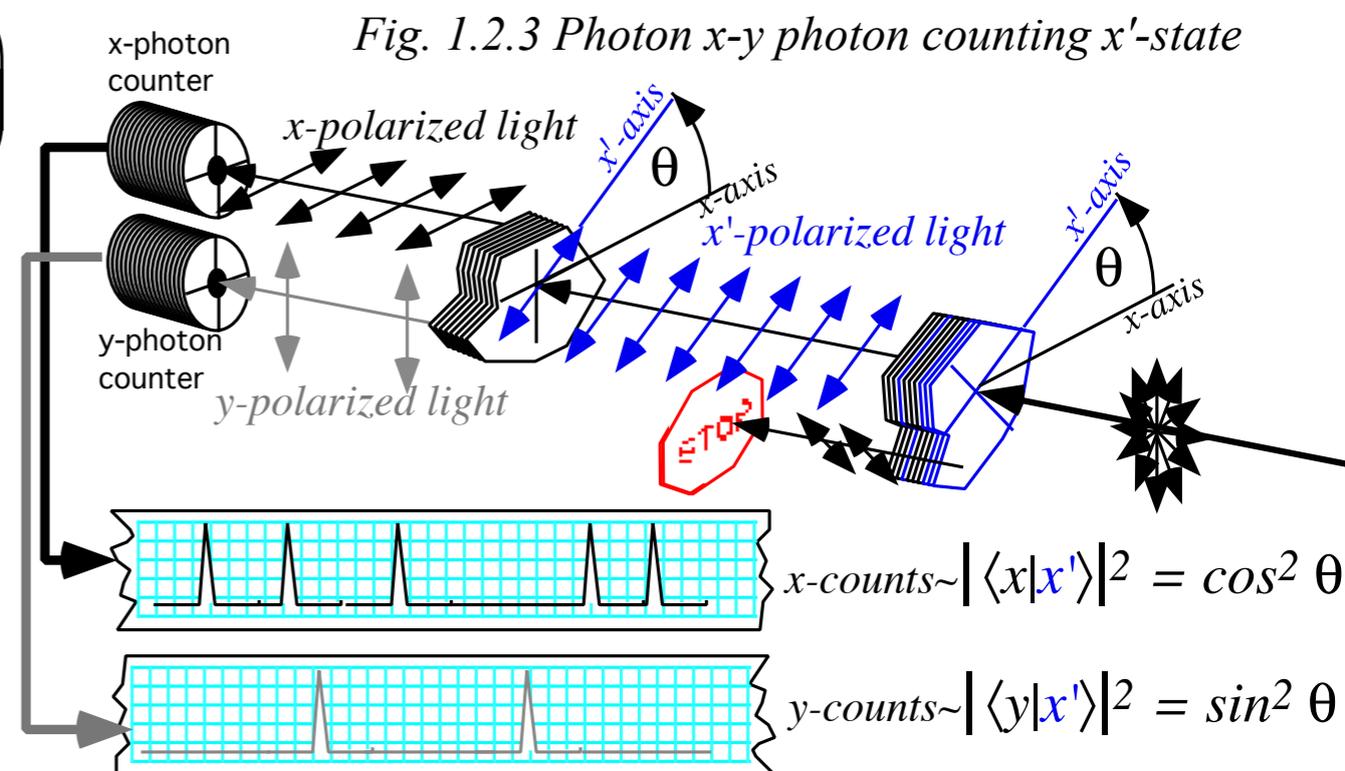
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Quantum em wave theory

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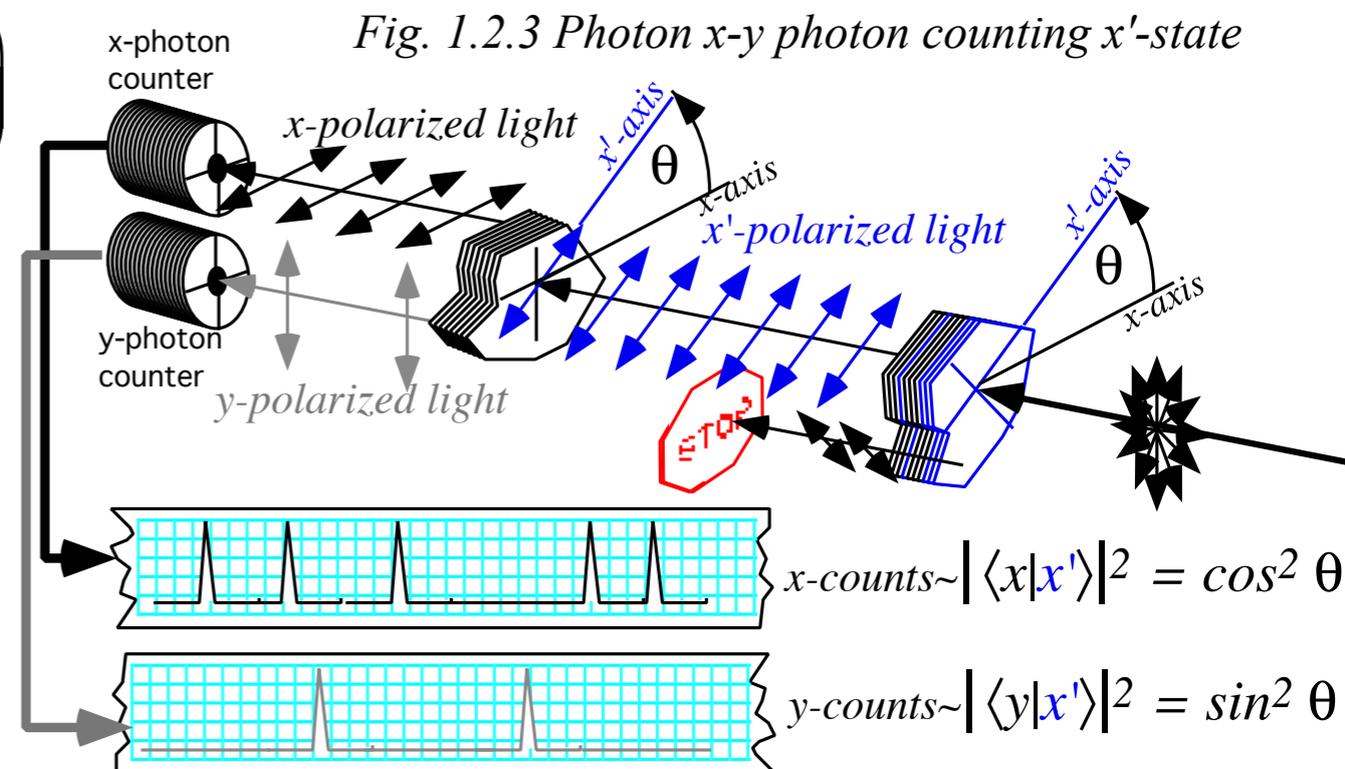
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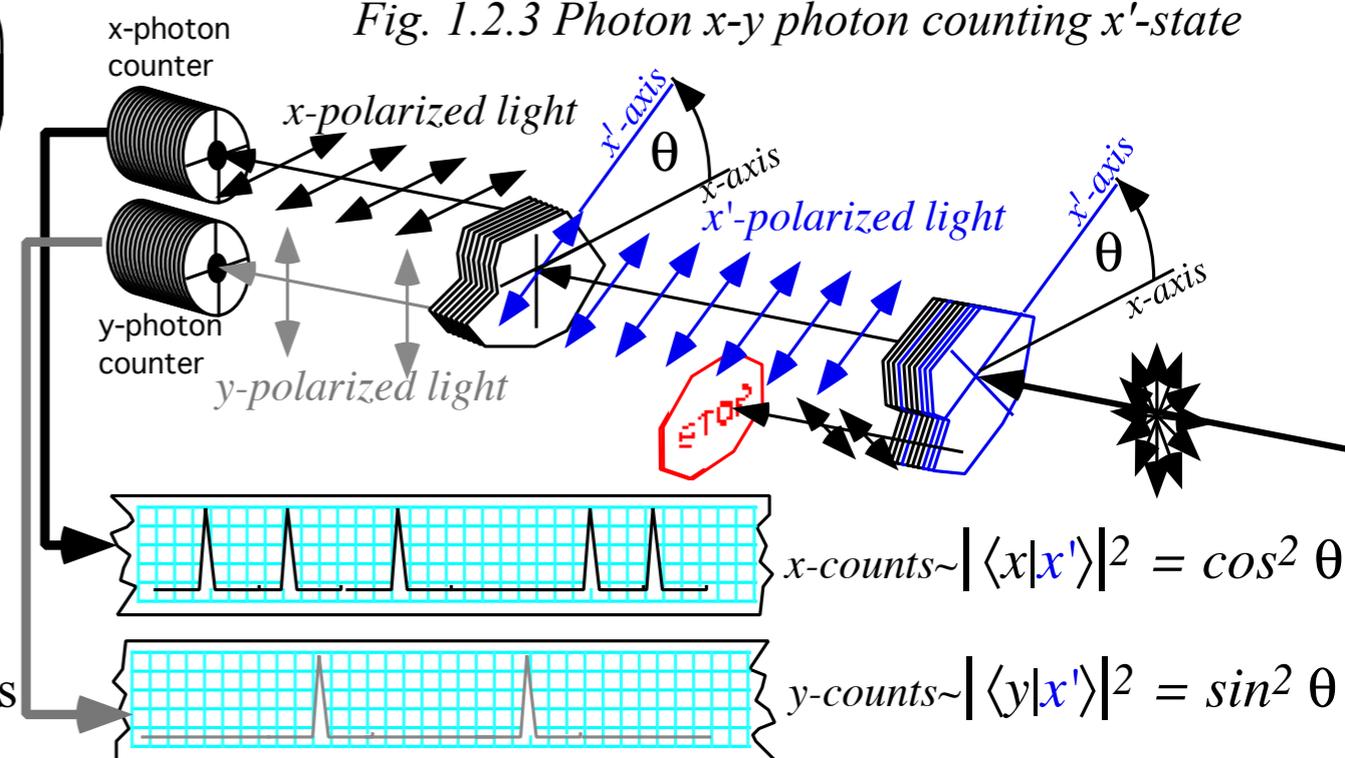
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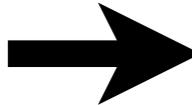
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*Review: “Abstraction” of bra and ket vectors from a Transformation Matrix*

*Introducing scalar and matrix products*

*Planck's energy and N-quanta (Cavity/Beam wave mode)*

 *Did Max Planck Goof? What's 1-photon worth?*

*Feynman amplitude axiom 1*

*What comes out of a beam sorter channel or branch-b?*

*Sample calculations*

*Feynman amplitude axioms 2-3*

*Beam analyzers: Sorter-unsorters*

*The “Do-Nothing” analyzer*

*Feynman amplitude axiom 4*

*Some “Do-Something” analyzers*

*Sorter-counter, Filter, 1/2-wave plate, 1/4-wave plate*

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Harmonic Oscillator Hamiltonian:

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$$\begin{aligned} \text{factors into: } H &= \frac{\hbar\omega}{2} \frac{(\sqrt{\omega}\mathbf{x} - i\mathbf{p}/\sqrt{\omega})}{\sqrt{2\hbar}} \frac{(\sqrt{\omega}\mathbf{x} + i\mathbf{p}/\sqrt{\omega})}{\sqrt{2\hbar}} + \frac{\hbar\omega}{2} \frac{(\sqrt{\omega}\mathbf{x} + i\mathbf{p}/\sqrt{\omega})}{\sqrt{2\hbar}} \frac{(\sqrt{\omega}\mathbf{x} - i\mathbf{p}/\sqrt{\omega})}{\sqrt{2\hbar}} \quad (\mathbf{a}\mathbf{a}^\dagger = \mathbf{a}^\dagger\mathbf{a} + 1) \\ H &= \frac{\hbar\omega}{2} \left( \underbrace{\quad}_{\mathbf{a}^\dagger} \cdot \underbrace{\quad}_{\mathbf{a}} + \underbrace{\quad}_{\mathbf{a}} \cdot \underbrace{\quad}_{\mathbf{a}^\dagger} \right) = \hbar\omega \left( \mathbf{a}^\dagger\mathbf{a} + \frac{1}{2} \cdot \mathbf{1} \right) \end{aligned}$$

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*Was Planck's linear-in-frequency- $\nu$  energy axiom a Goof?*

*Not at all. It's really not linear-in-frequency after all.*

*The  $\epsilon_0 |\mathbf{E}|^2 = h N \cdot \nu$  axiom IS in fact a product of TWO frequencies!*

*The 2nd "frequency" is COUNT RATE  $N$ .*

*As shown later: Both frequencies transform by Relativistic Doppler factor  $e^{\pm\phi}$ :*

*Coherent frequency is light quality:  $\nu' = e^{\pm\phi} \nu$ . Incoherent frequency is light quantity:  $N' = e^{\pm\phi} N$ .*

*This would imply that the  $|\mathbf{E}|$ -field also transforms like a frequency:  $|\mathbf{E}'| = e^{\pm\phi} |\mathbf{E}|$ .*

*Indeed, E-field amplitude is a frequency, too!*

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# Feynman amplitude axiom 1

## (1) The probability axiom

The first axiom deals with physical interpretation of amplitudes  $\langle j|k' \rangle$ .

*Axiom 1: The absolute square  $|\langle j|k' \rangle|^2 = \langle j|k' \rangle^* \langle j|k' \rangle$  gives probability for occurrence in state- $j$  of a system that started in state- $k'=1',2',\dots,$  or  $n'$  from one sorter and then was forced to choose between states  $j=1,2,\dots,n$  by another sorter.*

*Feynman-Dirac  
Interpretation of*

$$\langle j|k' \rangle$$

*=Amplitude of state- $j$  after  
state- $k'$  is forced to choose  
from available  $m$ -type states*

# Amplitude axioms apply to all intensity-<sup>probability</sup>conserving systems

This includes, first of all, spin-1/2 electron, proton, ..., <sup>13</sup>C, ... particles (Fermions)

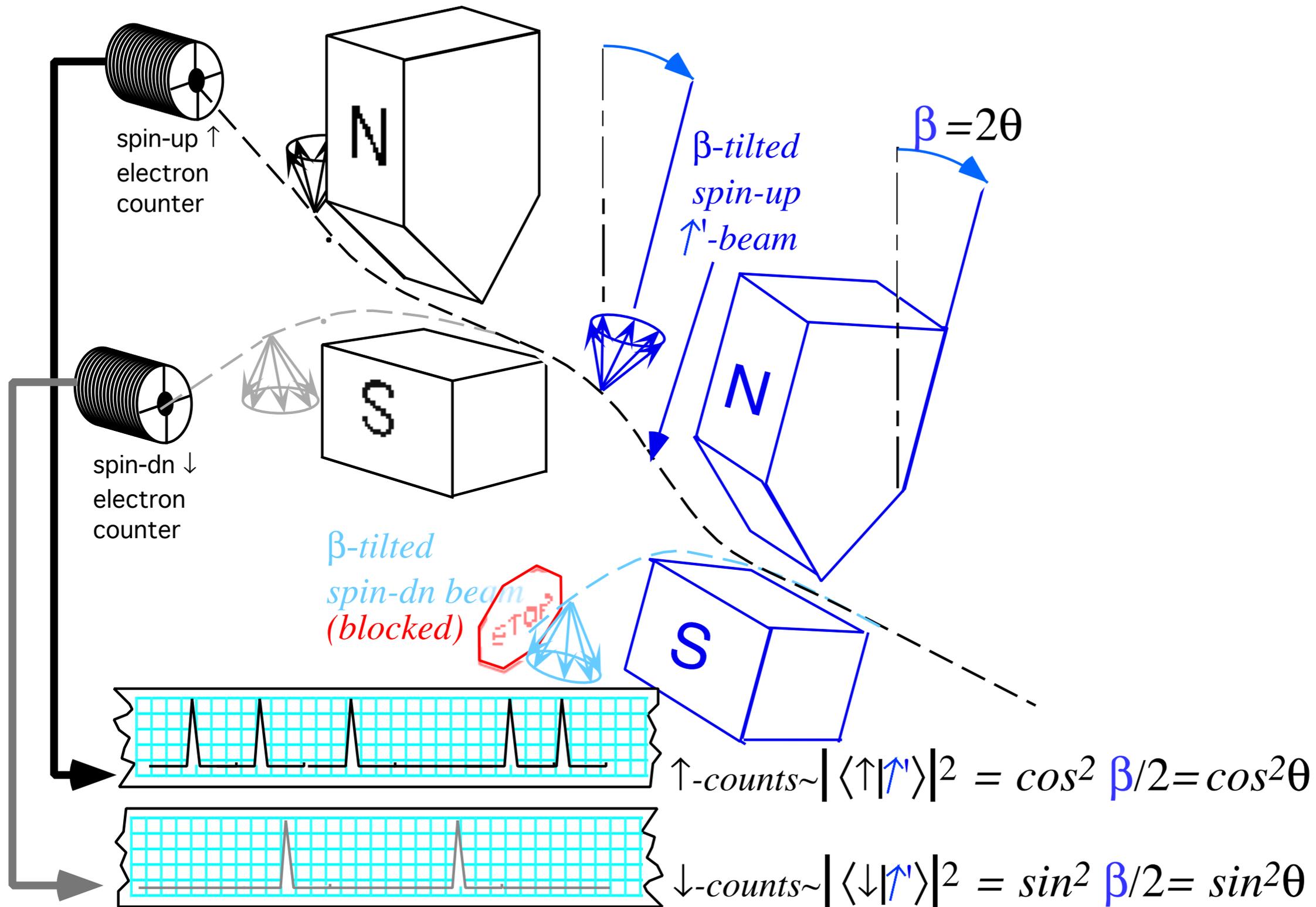


Fig. 1.2.4 Electron up-dn-spin counting of a tilted spin-up ( $\uparrow'$ )-state

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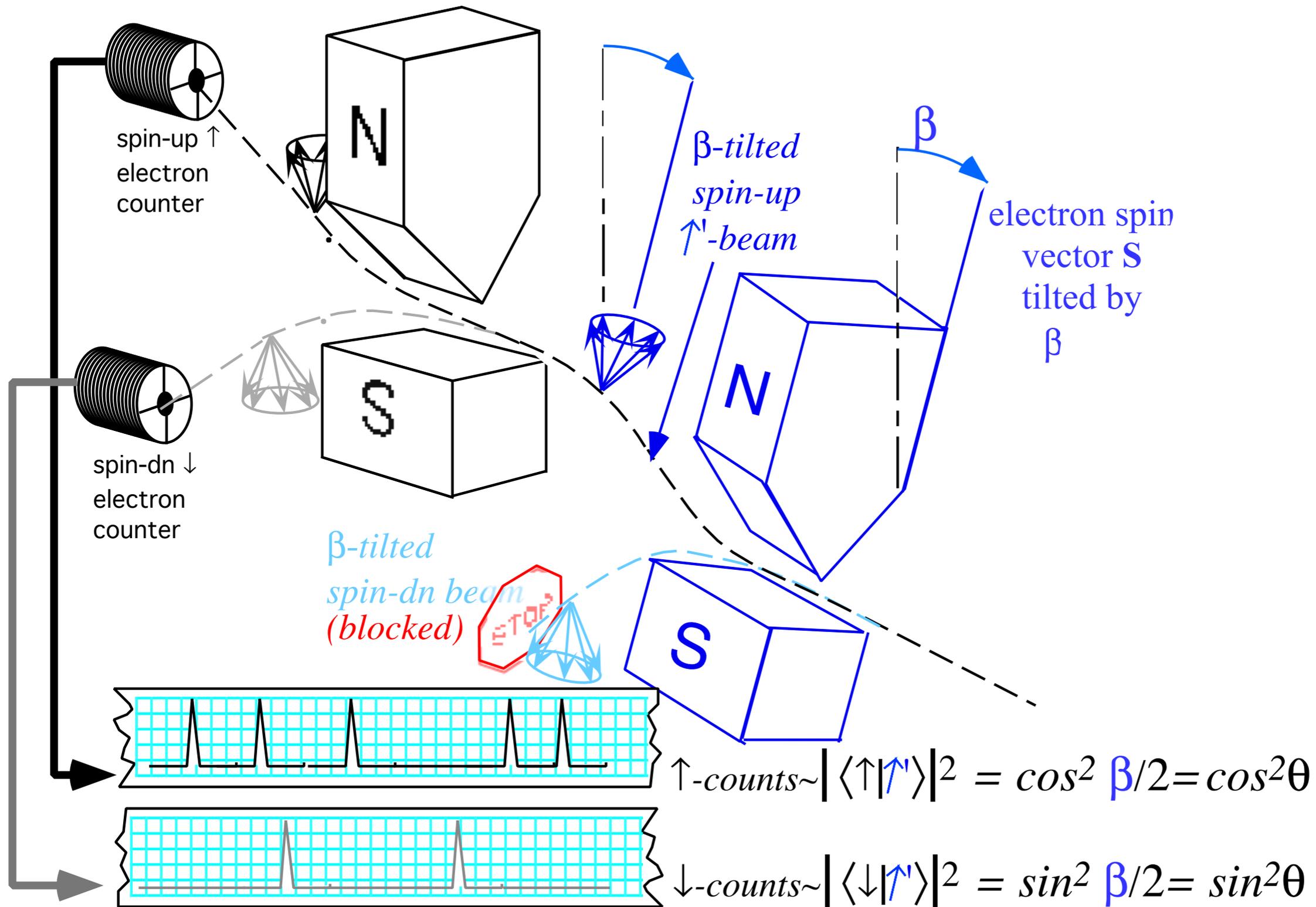


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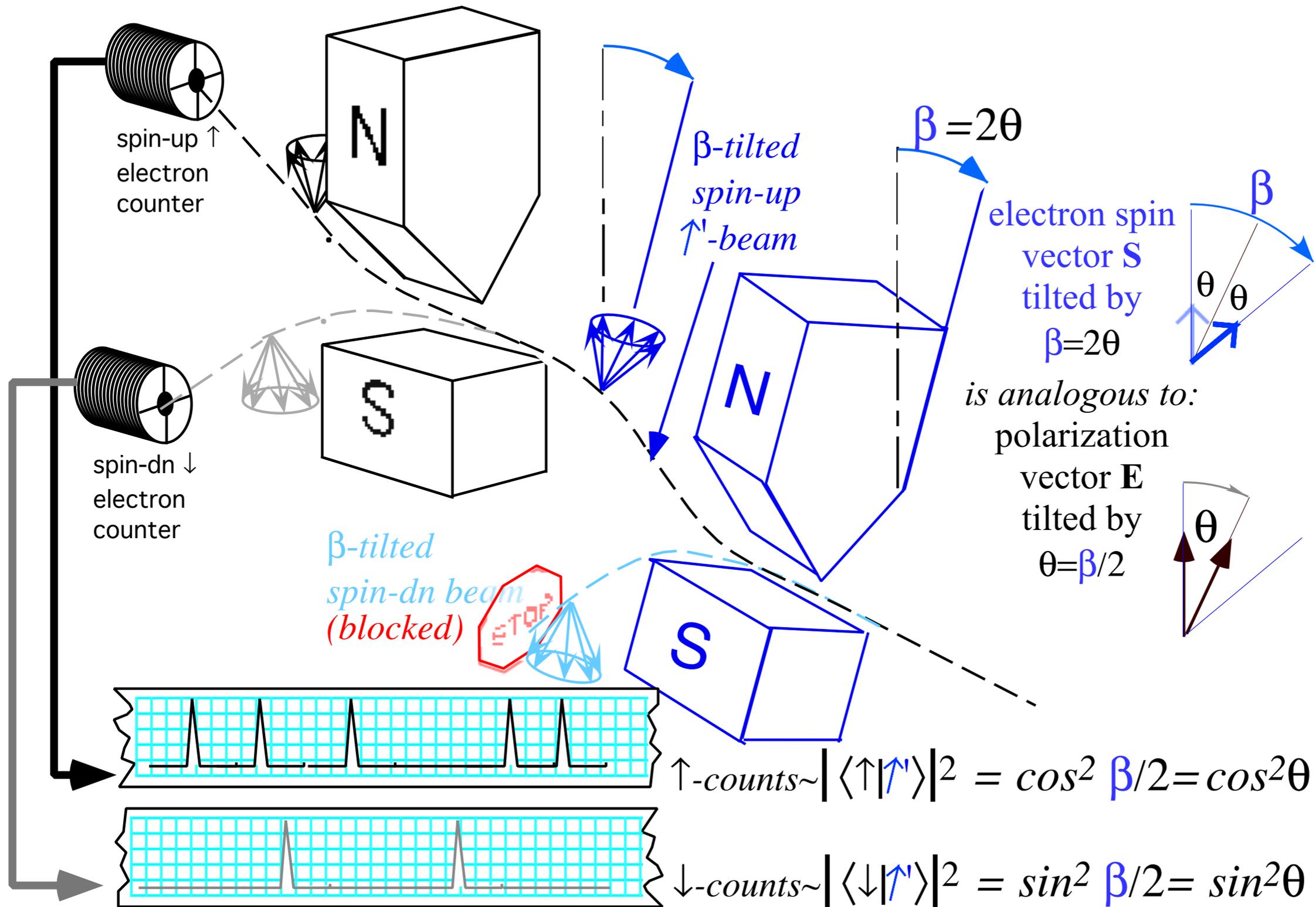


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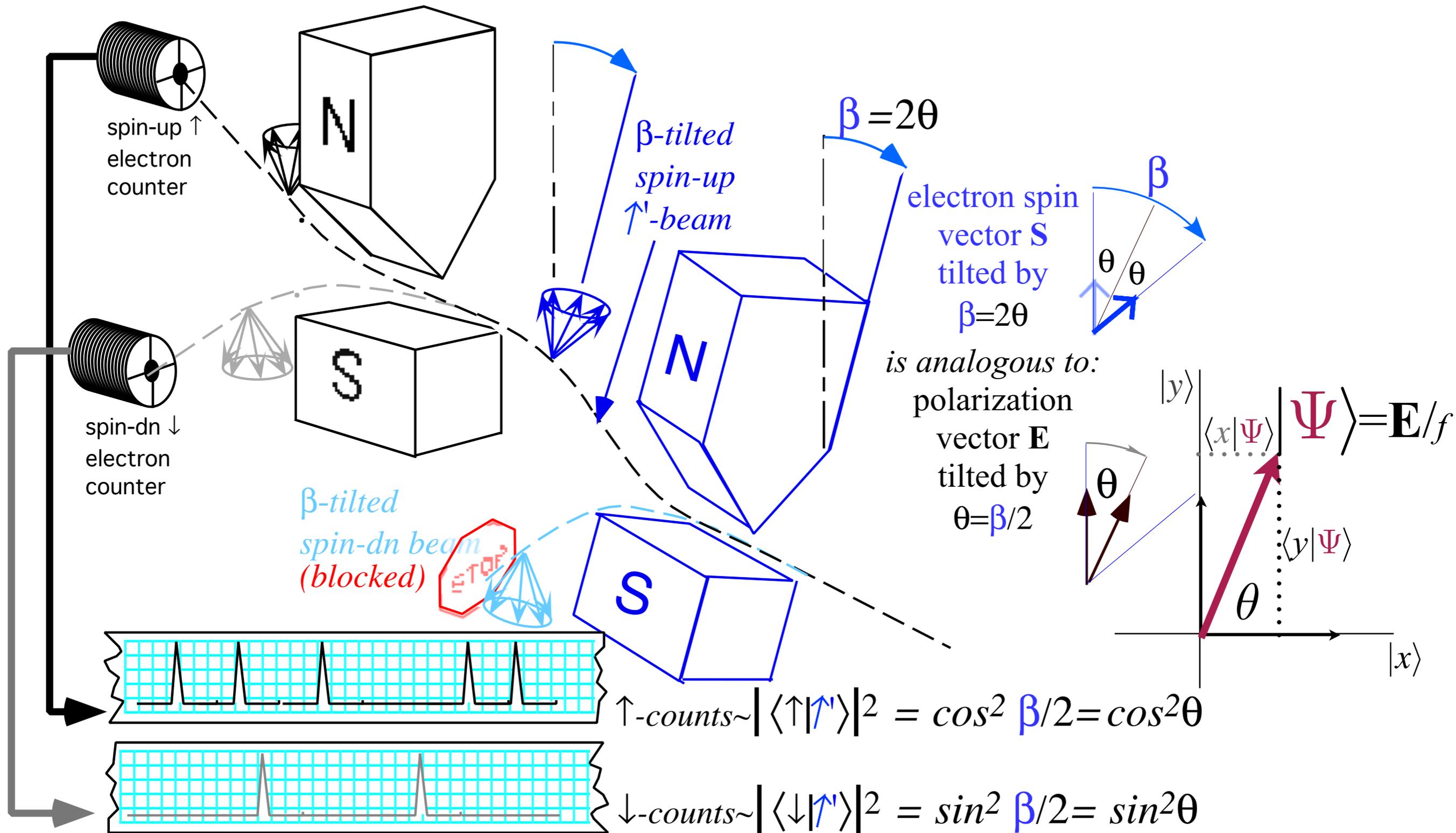


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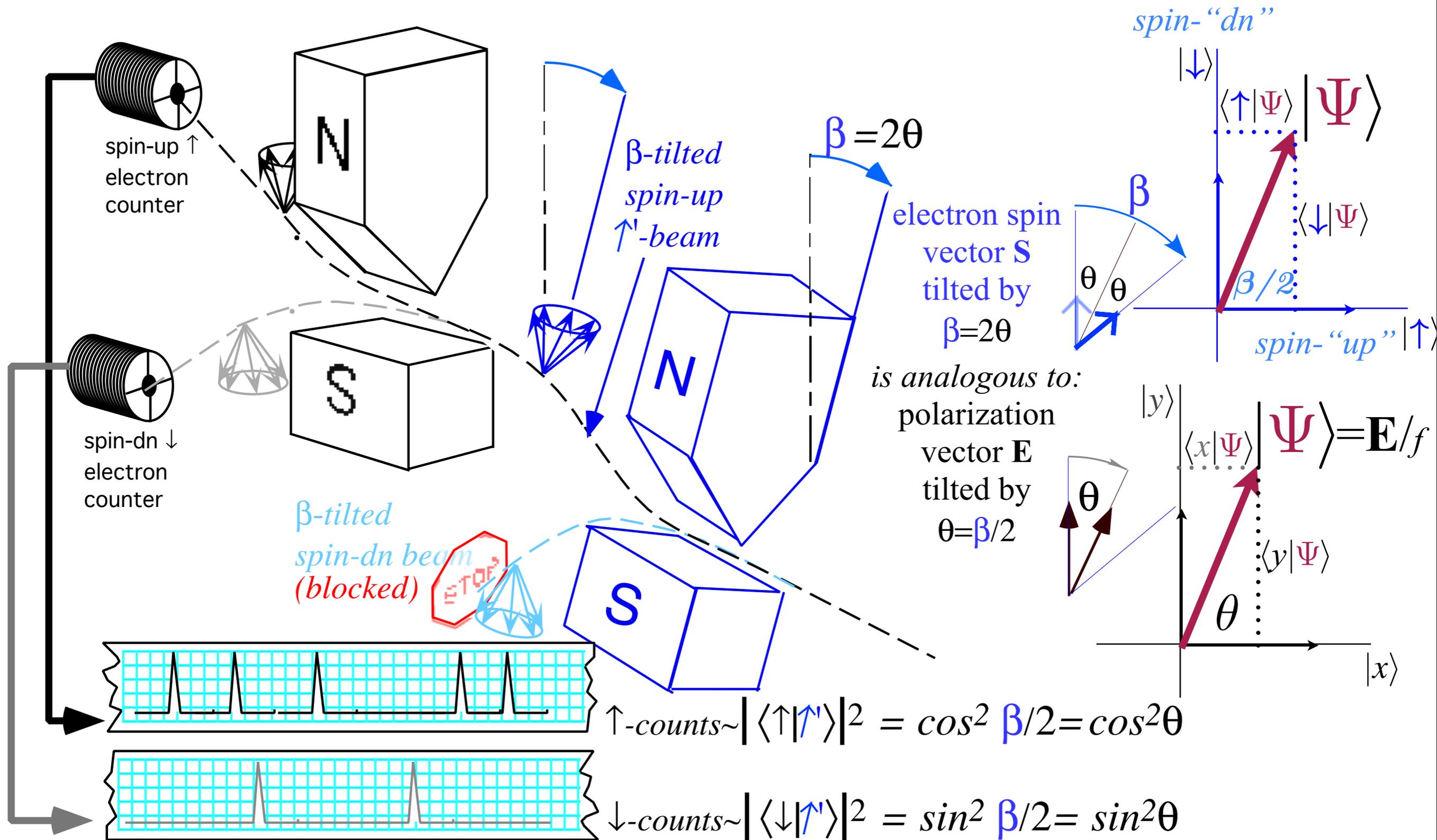


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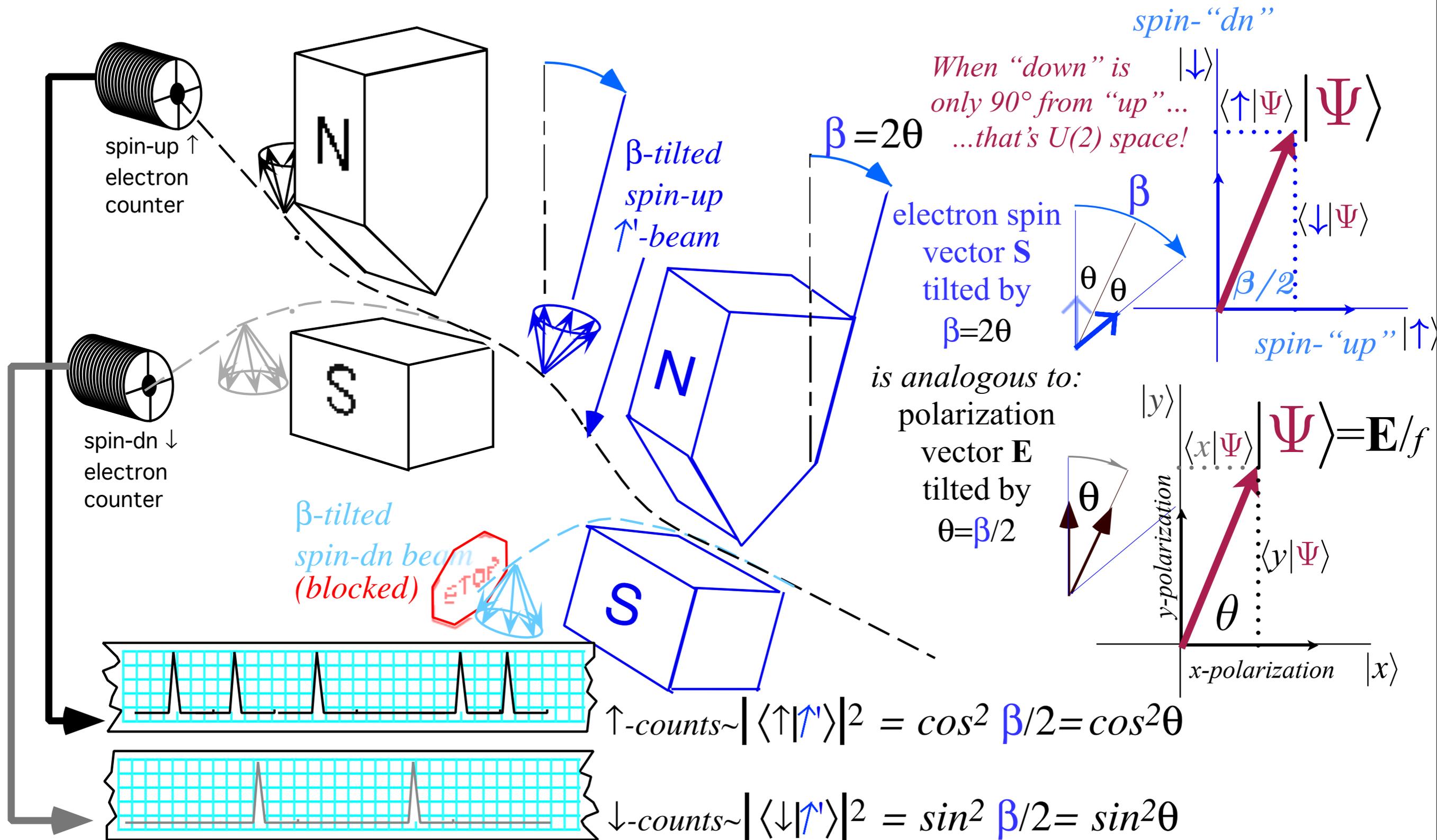


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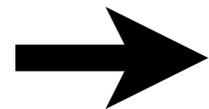
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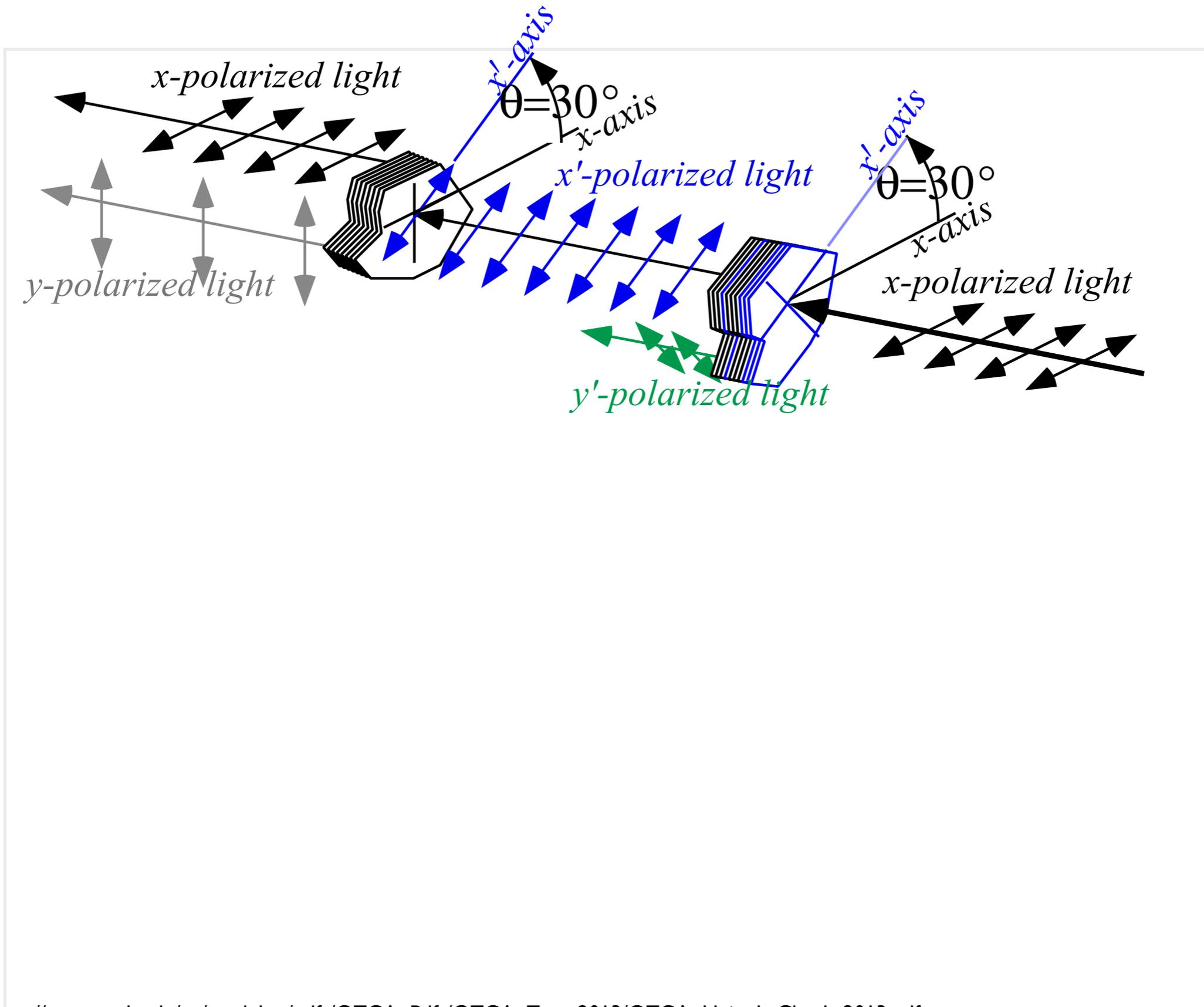
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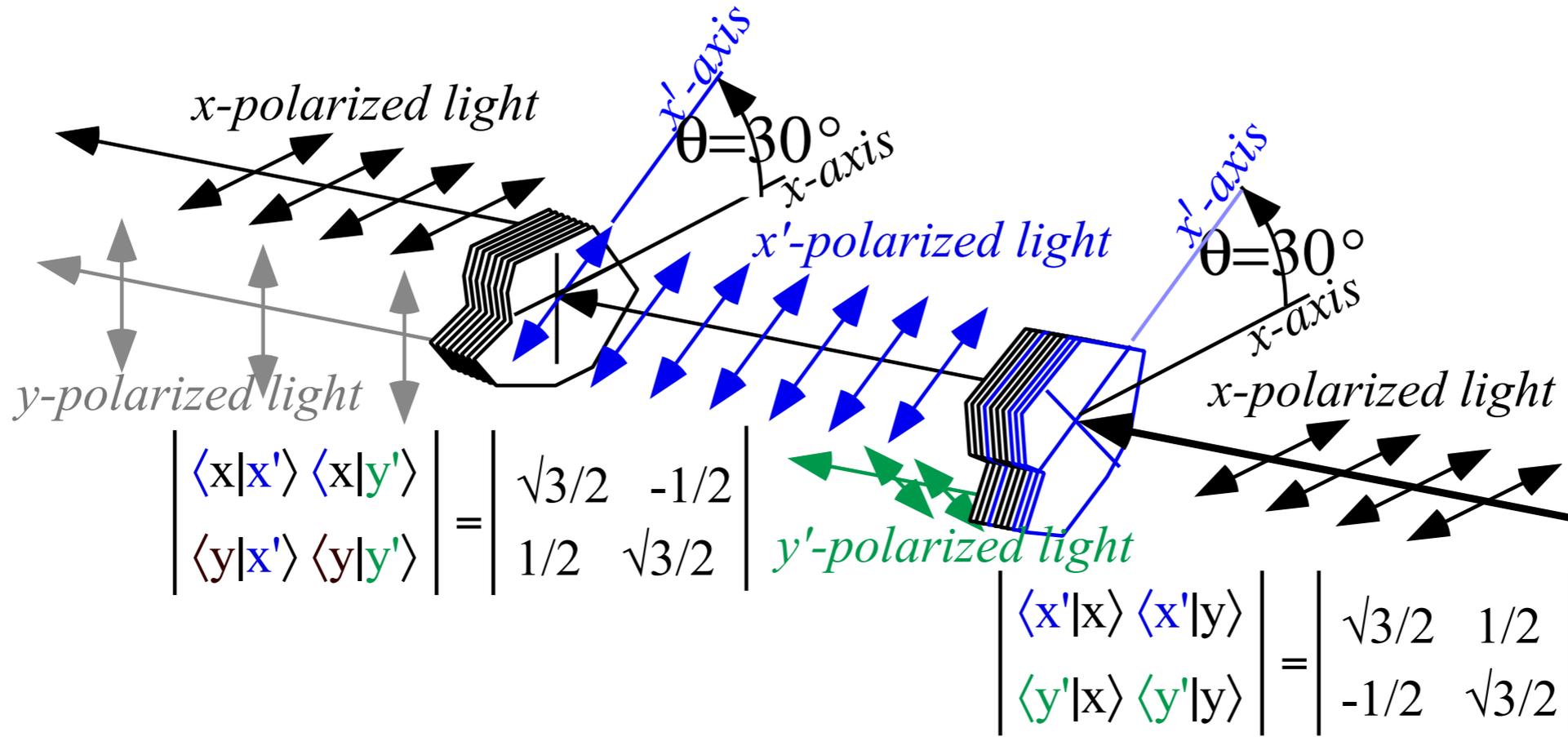
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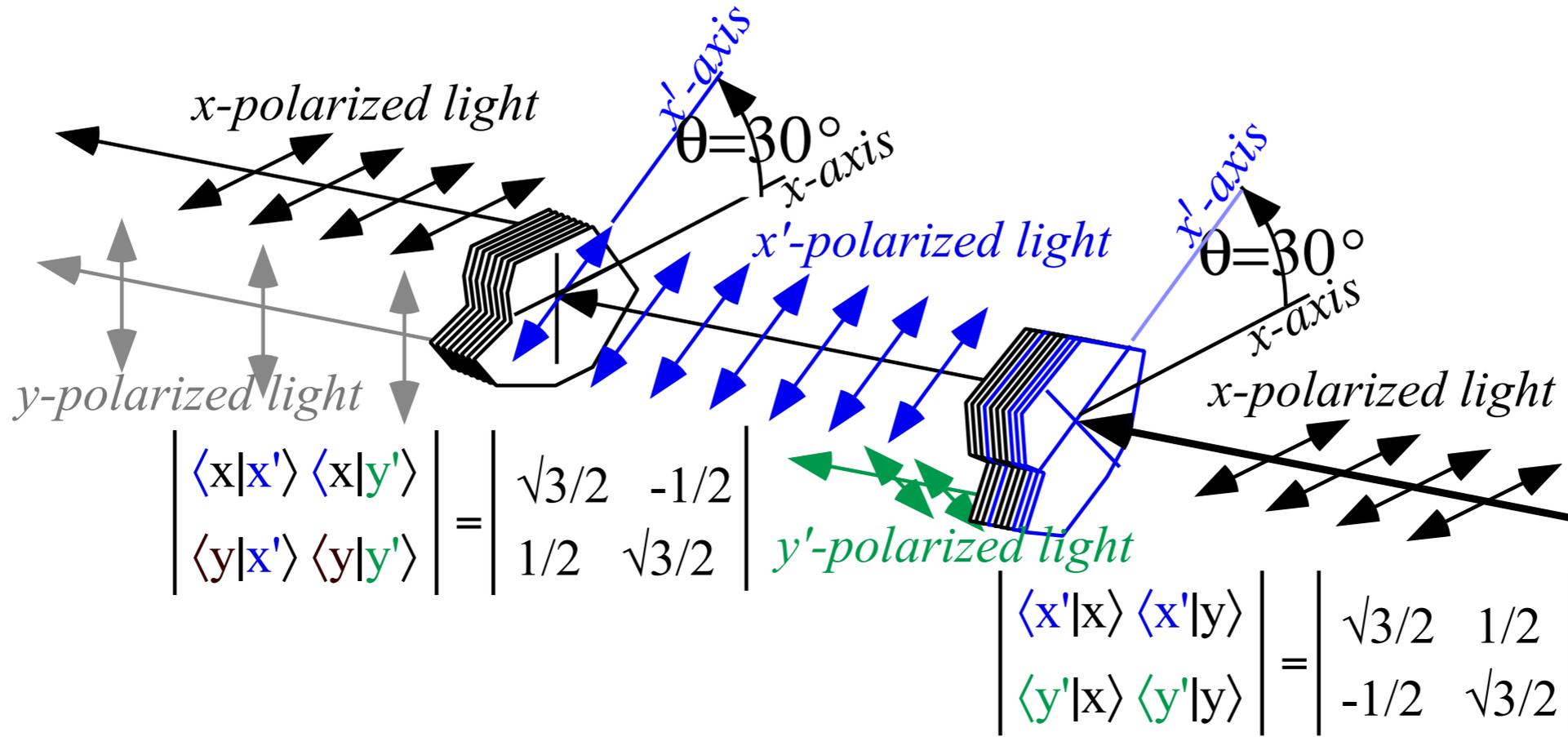
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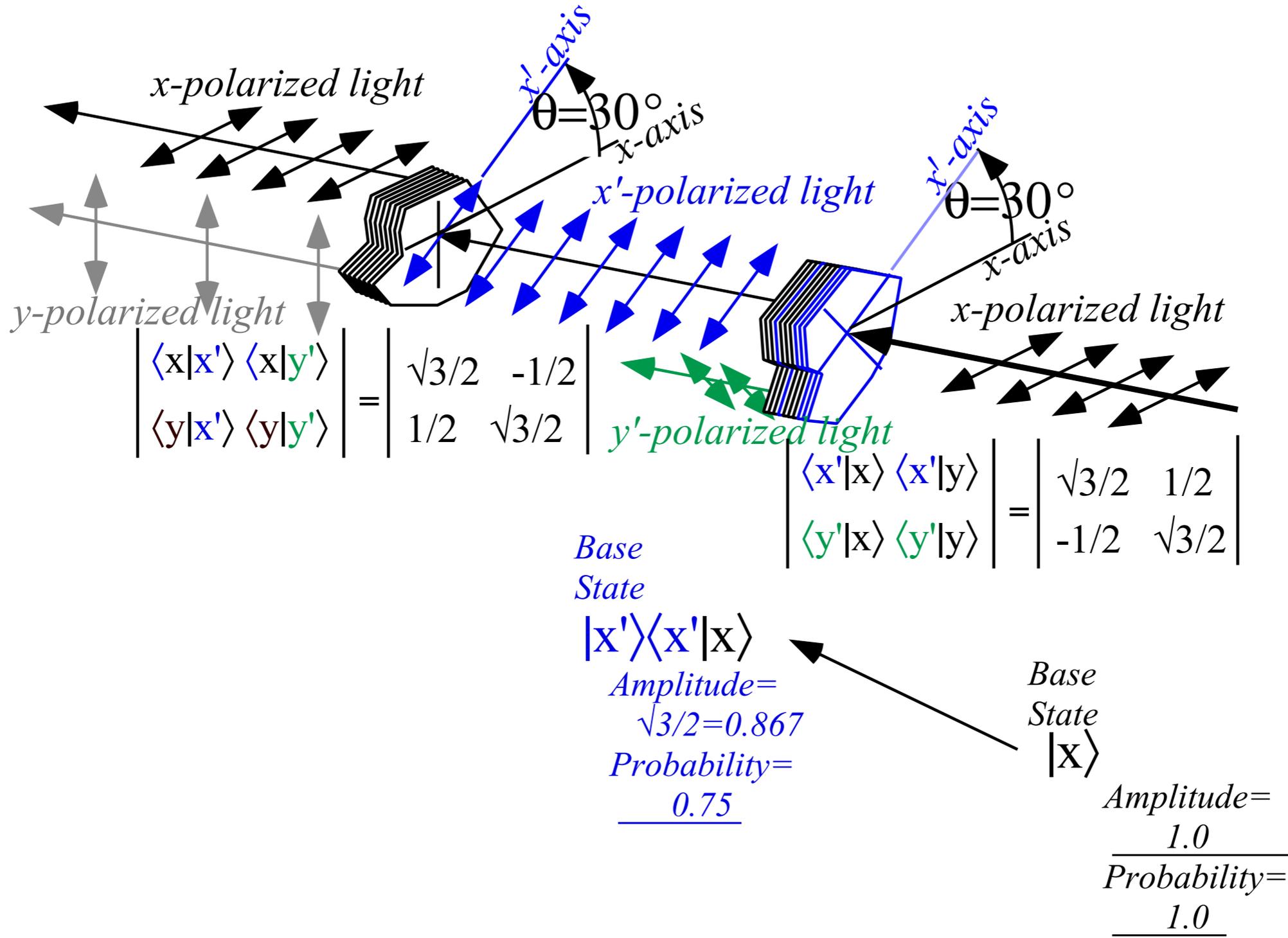
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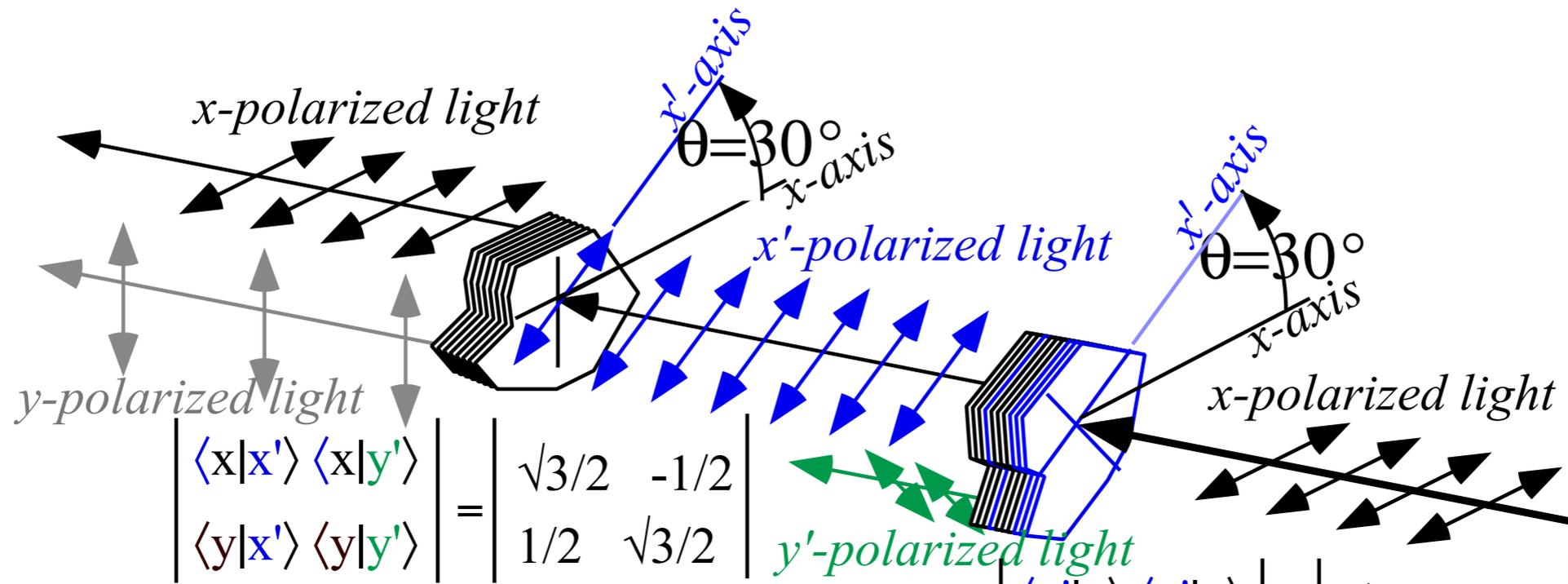
its *base state*  $|b\rangle$ , its *amplitude*  $\langle b|\Psi\rangle$ , and its *probability*  $|\langle b|\Psi\rangle|^2$  using *T-matrices*



## Q: What comes out of an analyzer channel or branch-b?

A: Want to determine or calculate:

its *base state*  $|b\rangle$ , its *amplitude*  $\langle b|\Psi\rangle$ , and its *probability*  $|\langle b|\Psi\rangle|^2$  using *T-matrices*



$$\begin{vmatrix} \langle x|x'\rangle & \langle x|y'\rangle \\ \langle y|x'\rangle & \langle y|y'\rangle \end{vmatrix} = \begin{vmatrix} \sqrt{3}/2 & -1/2 \\ 1/2 & \sqrt{3}/2 \end{vmatrix}$$

$$\begin{vmatrix} \langle x'|x\rangle & \langle x'|y\rangle \\ \langle y'|x\rangle & \langle y'|y\rangle \end{vmatrix} = \begin{vmatrix} \sqrt{3}/2 & 1/2 \\ -1/2 & \sqrt{3}/2 \end{vmatrix}$$

Base State

$$|x'\rangle\langle x'|x\rangle$$

Amplitude =  $\sqrt{3}/2 = 0.867$   
 Probability = 0.75

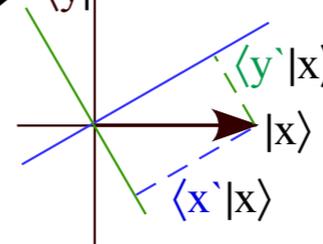
Base State

$$|y'\rangle\langle y'|x\rangle$$

Amplitude =  $-1/2 = -0.500$   
 Probability = 0.25

Base State  $|x\rangle$

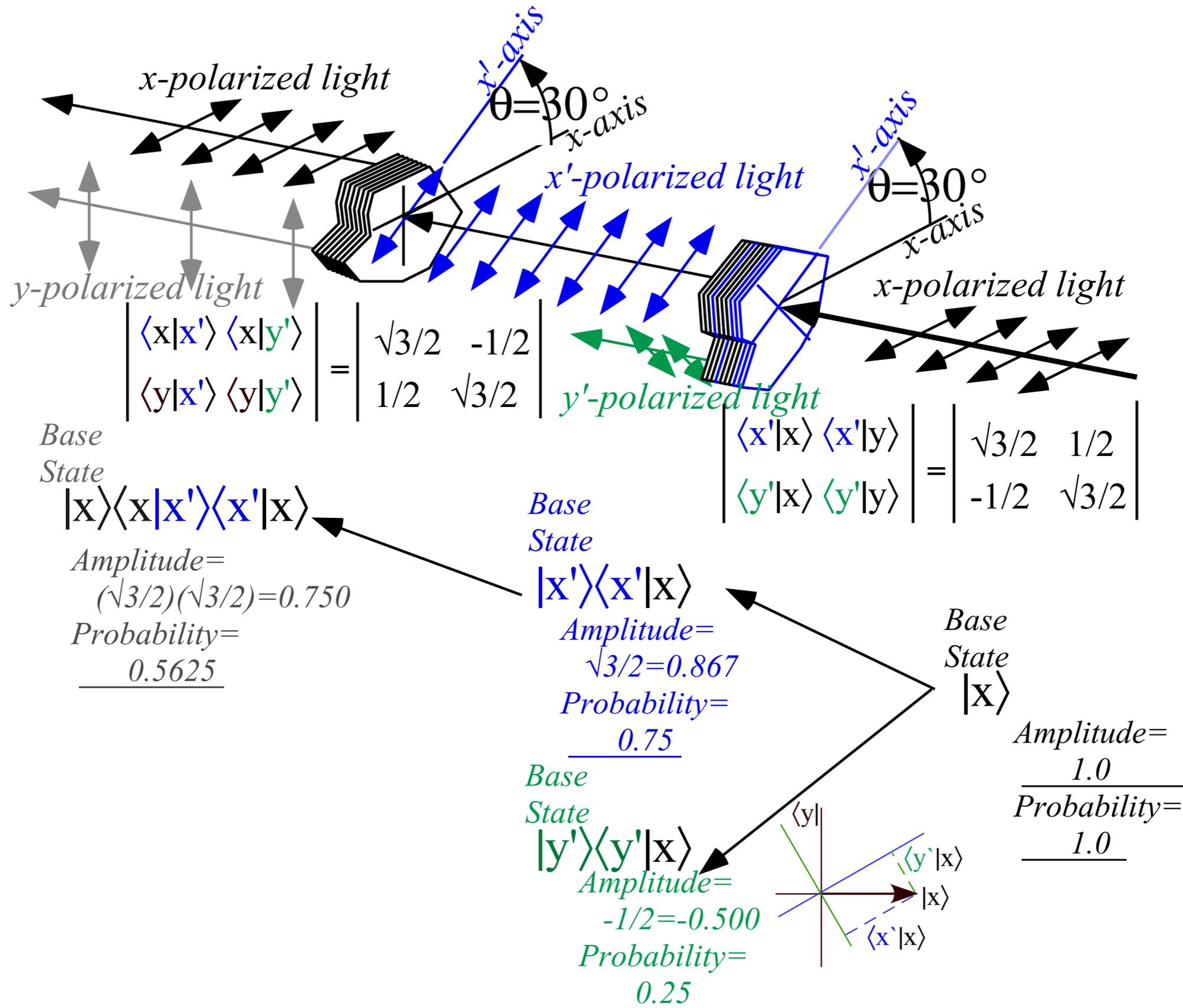
Amplitude = 1.0  
 Probability = 1.0



## Q: What comes out of an analyzer channel or branch-b?

A: Want to determine or calculate:

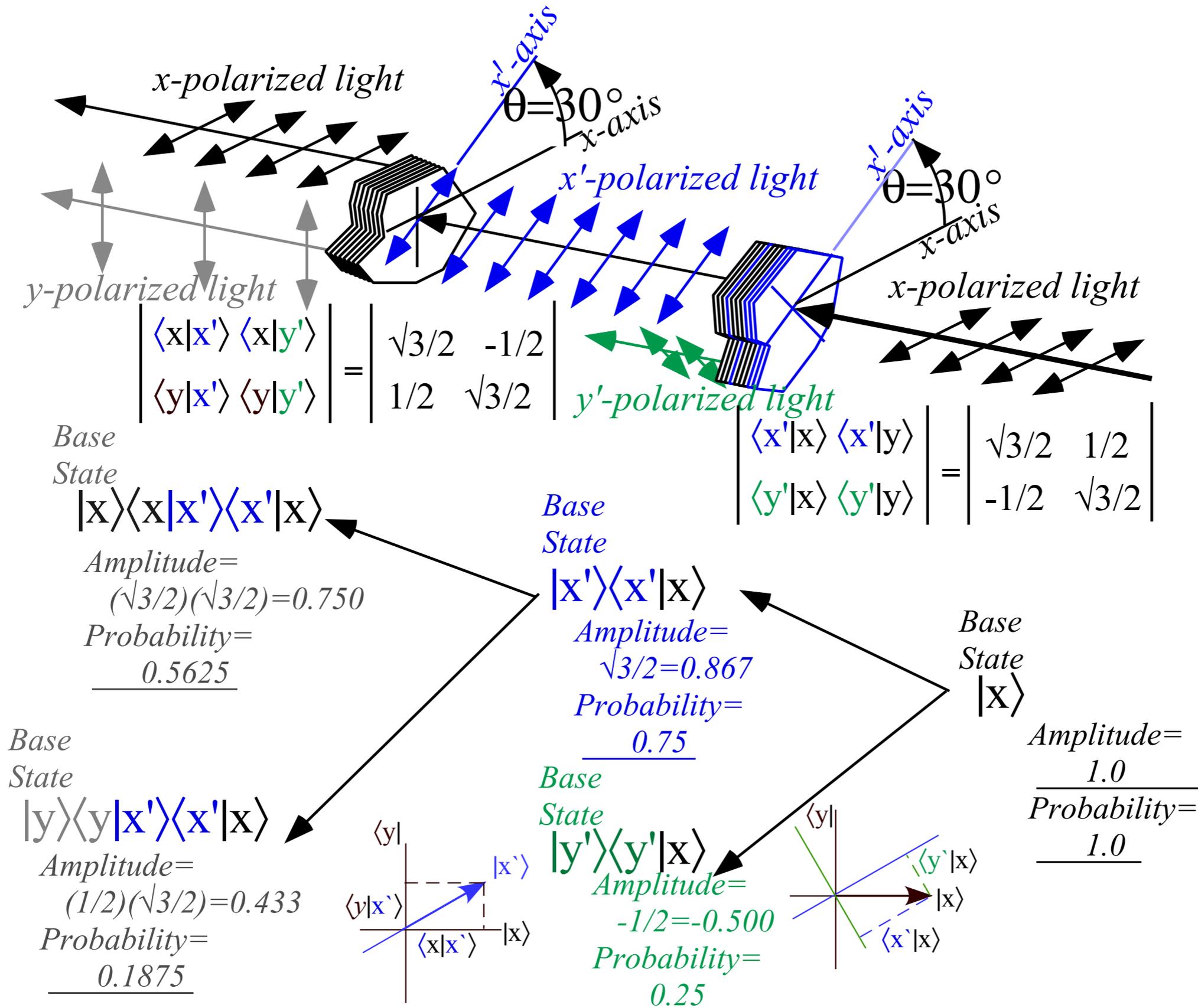
its *base state*  $|b\rangle$ , its *amplitude*  $\langle b|\Psi\rangle$ , and its *probability*  $|\langle b|\Psi\rangle|^2$  using *T-matrices*



## Q: What comes out of an analyzer channel or branch-b?

A: Want to determine or calculate:

its *base state*  $|b\rangle$ , its *amplitude*  $\langle b|\Psi\rangle$ , and its *probability*  $|\langle b|\Psi\rangle|^2$  using *T-matrices*



$$\langle e_p | D_o C_n B_m | \Psi \rangle = \langle e_p | d_o \rangle \langle d_o | c_n \rangle \langle c_n | b_m \rangle \langle b_m | \Psi \rangle$$

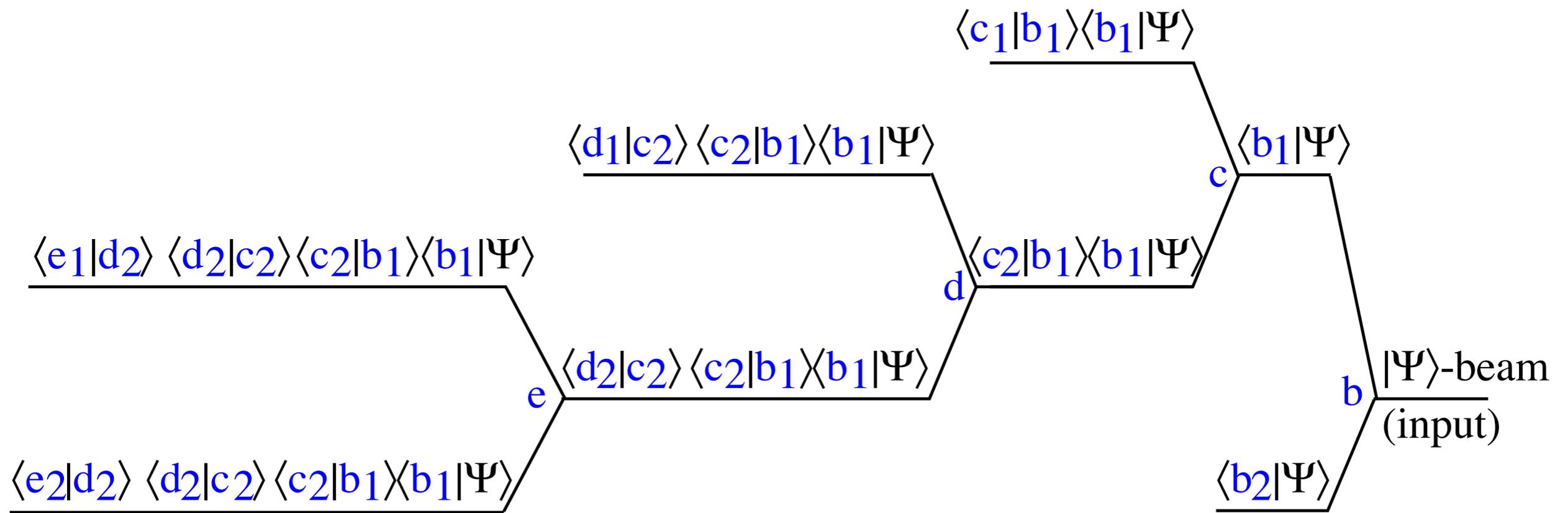


Fig. 1.3.10 Beams-amplitude products for successive beam sorting

*Review: “Abstraction” of bra and ket vectors from a Transformation Matrix*

*Introducing scalar and matrix products*

*Planck's energy and N-quanta (Cavity/Beam wave mode)*

*Did Max Planck Goof? What's 1-photon worth?*

*Feynman amplitude axiom 1*

*What comes out of a beam sorter channel or branch-b?*

*Sample calculations*

 *Feynman amplitude axioms 2-3*

*Beam analyzers: Sorter-unsorters*

*The “Do-Nothing” analyzer*

*Feynman amplitude axiom 4*

*Some “Do-Something” analyzers*

*Sorter-counter, Filter, 1/2-wave plate, 1/4-wave plate*

# *Feynman amplitude axiom 1* (Given above p.35)

## **(1) The probability axiom**

The first axiom deals with physical interpretation of amplitudes  $\langle j|k' \rangle$ .

*Axiom 1: The absolute square  $|\langle j|k' \rangle|^2 = \langle j|k' \rangle^* \langle j|k' \rangle$  gives probability for occurrence in state- $j$  of a system that started in state- $k'=1',2',\dots,$  or  $n'$  from one sorter and then was forced to choose between states  $j=1,2,\dots,n$  by another sorter.*

*Feynman-Dirac  
Interpretation of*

$$\langle j|k' \rangle$$

*=Amplitude of state- $j$  after  
state- $k'$  is forced to choose  
from available  $m$ -type states*

# Feynman amplitude axioms 1-2

## (1) The probability axiom

The first axiom deals with physical interpretation of amplitudes  $\langle j|k' \rangle$ .

*Axiom 1: The absolute square  $|\langle j|k' \rangle|^2 = \langle j|k' \rangle^* \langle j|k' \rangle$  gives probability for occurrence in state- $j$  of a system that started in state- $k'=1',2',\dots,$  or  $n'$  from one sorter and then was forced to choose between states  $j=1,2,\dots,n$  by another sorter.*

Feynman-Dirac  
Interpretation of

$$\langle j|k' \rangle$$

=Amplitude of state- $j$  after state- $k'$  is forced to choose from available  $m$ -type states

## (2) The conjugation or inversion axiom (time reversal symmetry)

The second axiom concerns going backwards through a sorter or the reversal of amplitudes.

*Axiom 2: The complex conjugate  $\langle j|k' \rangle^*$  of an amplitude  $\langle j|k' \rangle$  equals its reverse:  $\langle j|k' \rangle^* = \langle k'|j \rangle$*

# Feynman amplitude axioms 1-3

Feynman-Dirac  
Interpretation of

$$\langle j | k' \rangle$$

= Amplitude of state- $j$  after  
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## (1) The probability axiom

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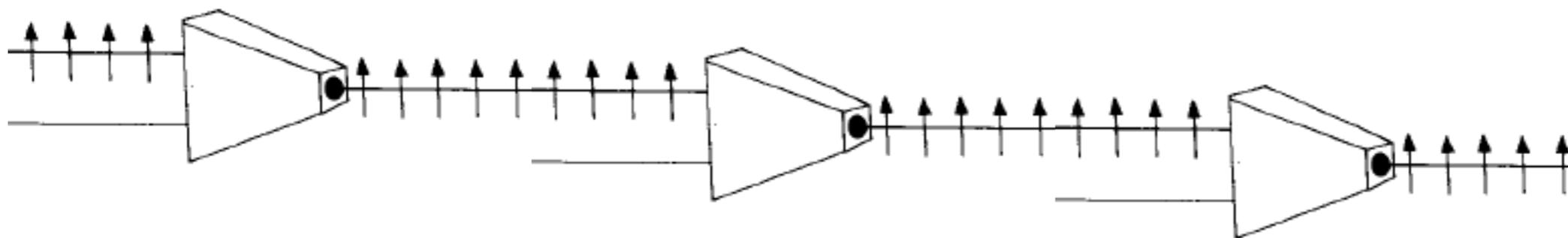
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## (3) The orthonormality or identity axiom

The third axiom concerns the amplitude for "re measurement" by the same analyzer.

*Axiom 3: If identical analyzers are used twice or more the amplitude for a passed state- $k$  is one, and for all others it is zero:*

$$\langle j | k \rangle = \delta_{jk} = \begin{cases} 1 & \text{if: } j=k \\ 0 & \text{if: } j \neq k \end{cases} = \langle j' | k' \rangle$$



*Review: “Abstraction” of bra and ket vectors from a Transformation Matrix*

*Introducing scalar and matrix products*

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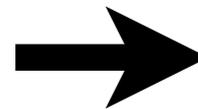
*Did Max Planck Goof? What's 1-photon worth?*

*Feynman amplitude axiom 1*

*What comes out of a beam sorter channel or branch-b?*

*Sample calculations*

*Feynman amplitude axioms 2-3*

 *Beam analyzers: Sorter-unsorters*

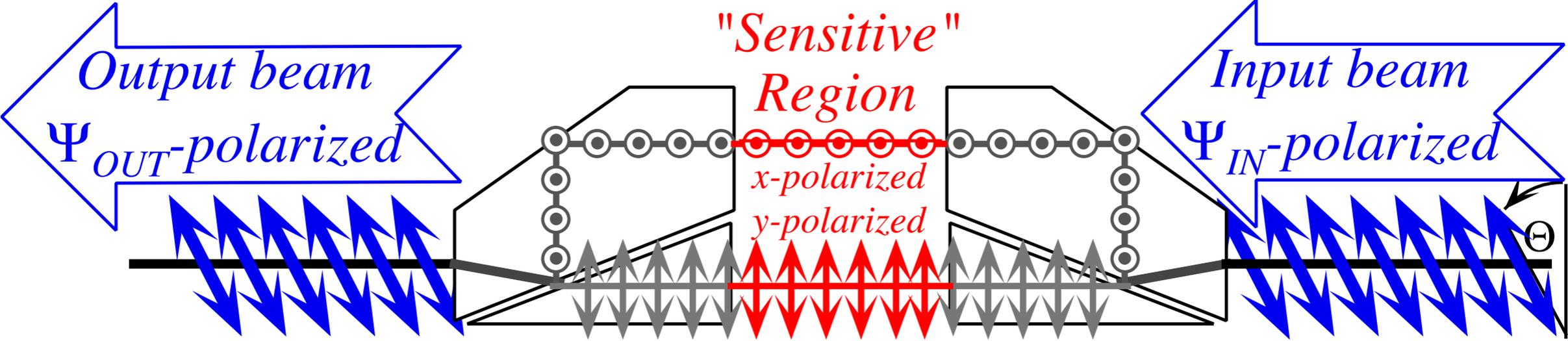
*The “Do-Nothing” analyzer*

*Feynman amplitude axiom 4*

*Some “Do-Something” analyzers*

*Sorter-counter, Filter, 1/2-wave plate, 1/4-wave plate*

*Beam Analyzers: Fundamental Quantum Processes (Sorter-Unsorter pairs)*



*Fig. 1.3.1 Anatomy of ideal optical polarization analyzer in its "Do nothing" mode*

# Beam Analyzers: Fundamental Quantum Processes (Sorter-Unsorter pairs)

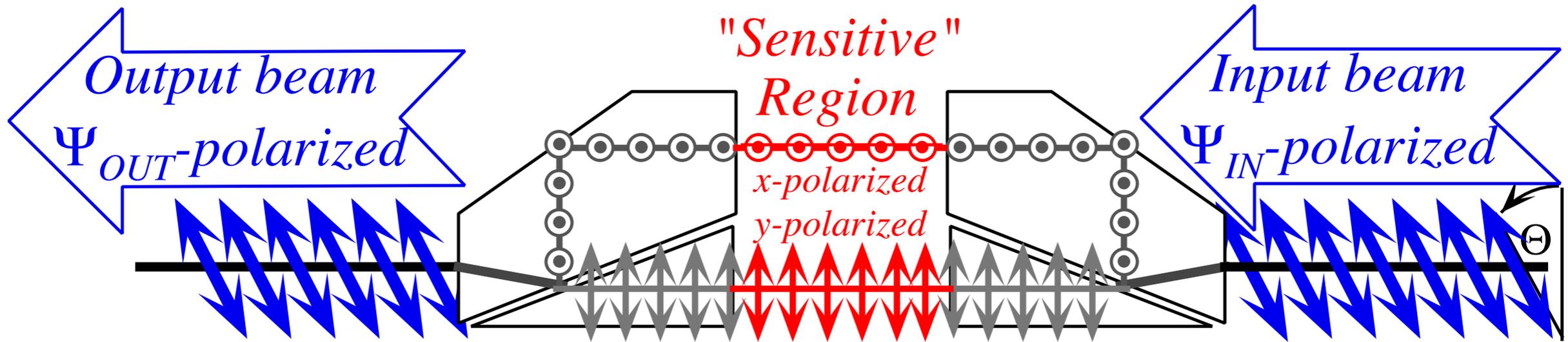


Fig. 1.3.1 Anatomy of ideal optical polarization analyzer in its "Do nothing" mode

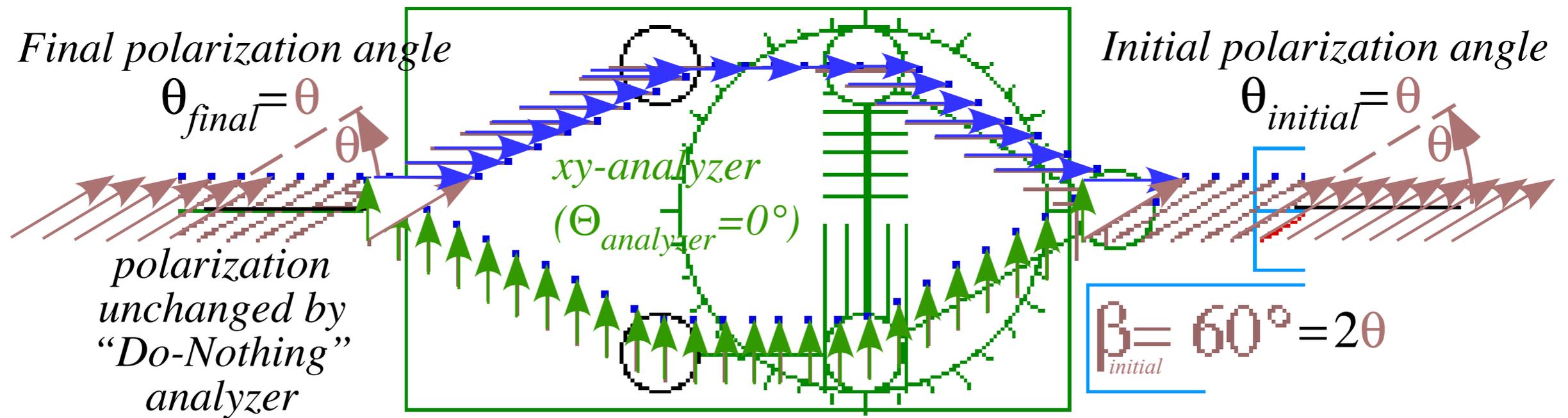


Fig. 1.3.2 Computer sketch of simulated polarization analyzer in "do-nothing" mode

*Review: “Abstraction” of bra and ket vectors from a Transformation Matrix*  
*Introducing scalar and matrix products*

*Planck's energy and N-quanta (Cavity/Beam wave mode)*

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# Beam Analyzers: Fundamental Quantum Processes (Sorter-Unsorter pairs)

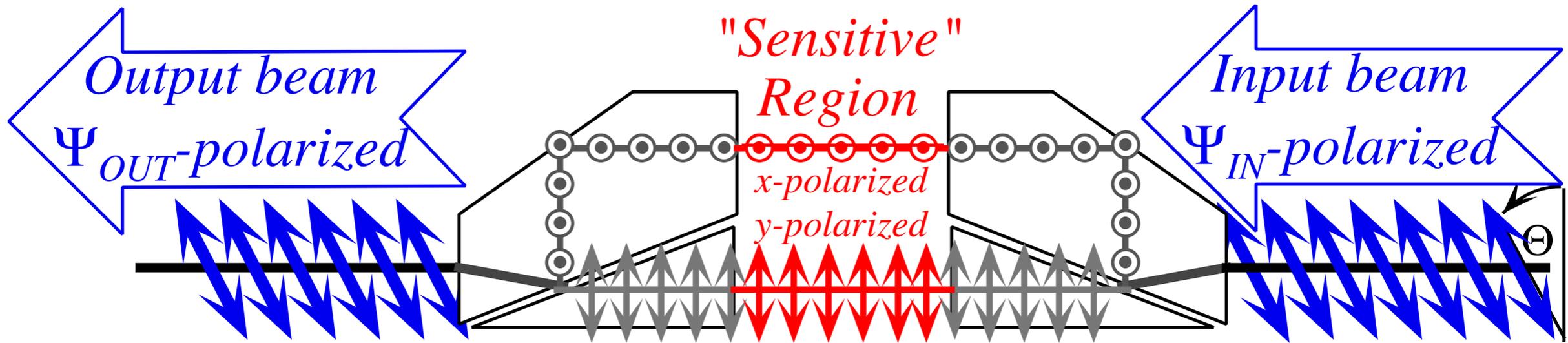


Fig. 1.3.1 Anatomy of ideal optical polarization analyzer in its "Do nothing" mode

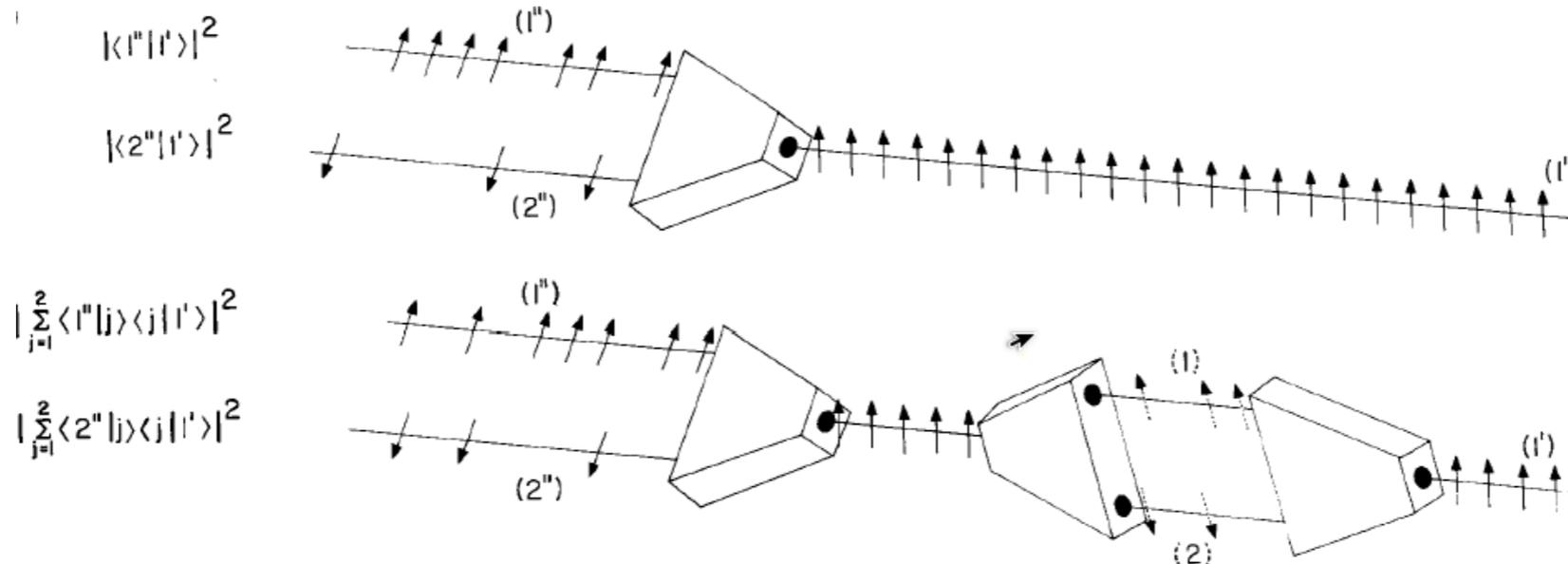
## Feynman amplitude axiom 4

### (4) The completeness or closure axiom

The fourth axiom concerns the "Do-nothing" property of an ideal analyzer, that is, a sorter followed by an "unsorter" or "put-back-togetherer" as sketched above.

*Axiom 4. Ideal sorting followed by ideal recombination of amplitudes has no effect:*

$$\langle j'' | m' \rangle = \sum_{k=1}^n \langle j'' | k \rangle \langle k | m' \rangle$$



# Beam Analyzers: Fundamental Quantum Processes (Sorter-Unsorter pairs)

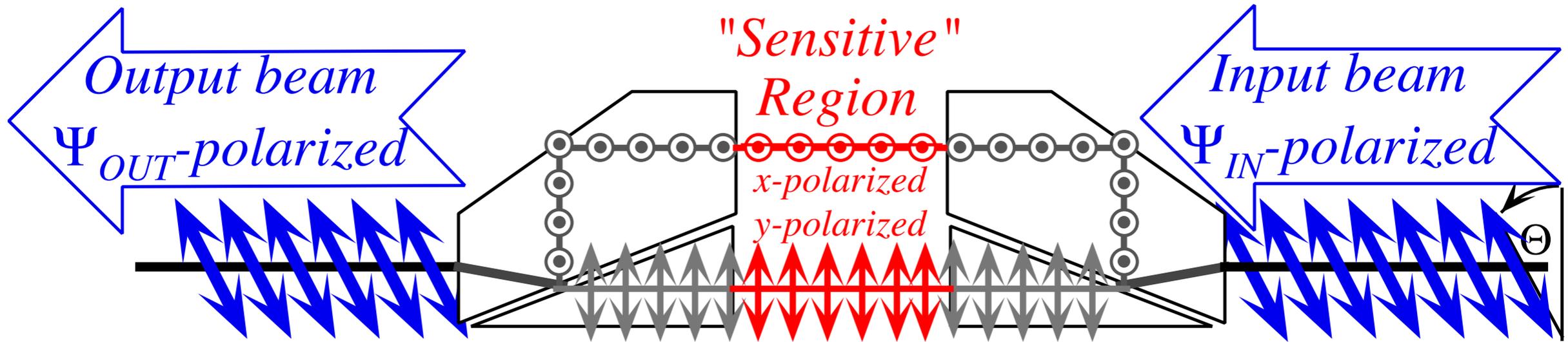


Fig. 1.3.1 Anatomy of ideal optical polarization analyzer in its "Do nothing" mode

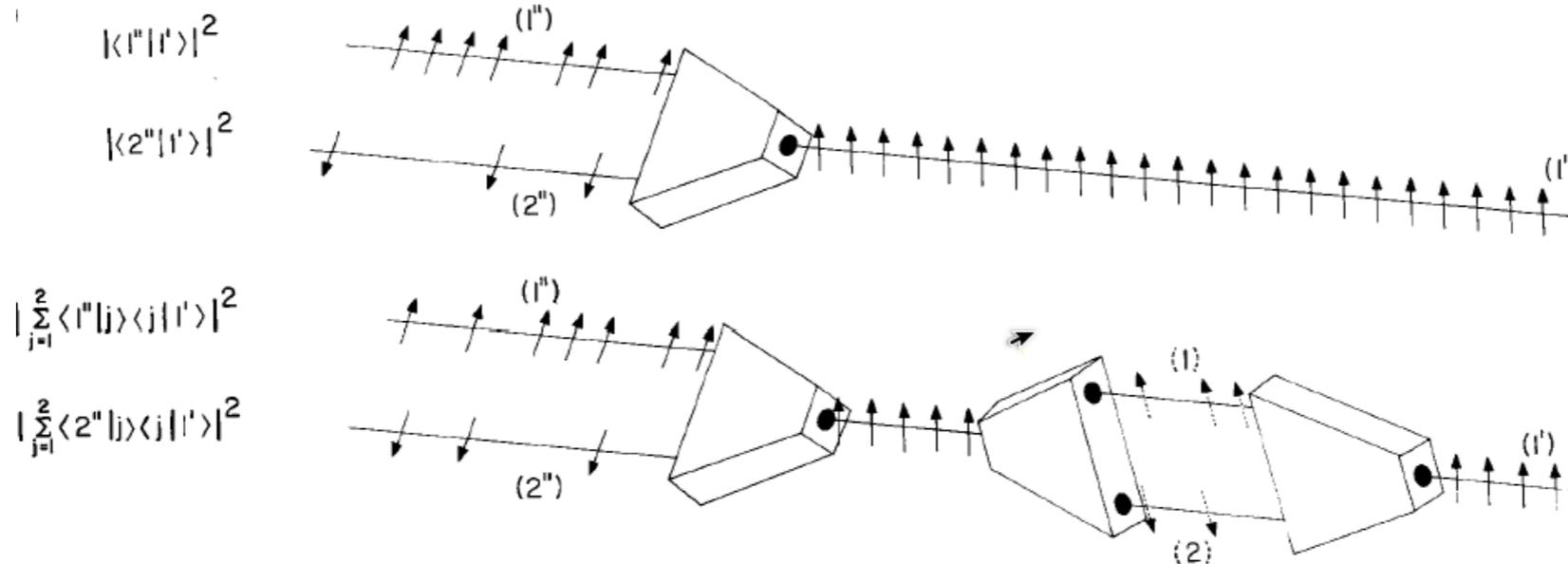
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$$\langle j'' | m' \rangle = \sum_{k=1}^n \langle j'' | k \rangle \langle k | m' \rangle$$



May use axioms 1-3 to prove special case:

$$1 = \langle m' | m' \rangle \text{ by 3}$$

$\Downarrow_{\text{by 1}}$

$$1 = \sum_{k=1}^n |\langle k | m' \rangle|^2$$

$\Downarrow_{\text{by 1}}$

$$1 = \sum_{k=1}^n \langle k | m' \rangle^* \langle k | m' \rangle$$

$\Downarrow_{\text{by 2}}$

$$\sum_{k=1}^n \langle m' | k \rangle \langle k | m' \rangle = \langle m' | m' \rangle$$

# Feynman amplitude axioms 1-4

Feynman-Dirac  
Interpretation of

$$\langle j | k' \rangle$$

= Amplitude of state- $j$  after  
state- $k'$  is forced to choose  
from available  $m$ -type states

## (1) The probability axiom

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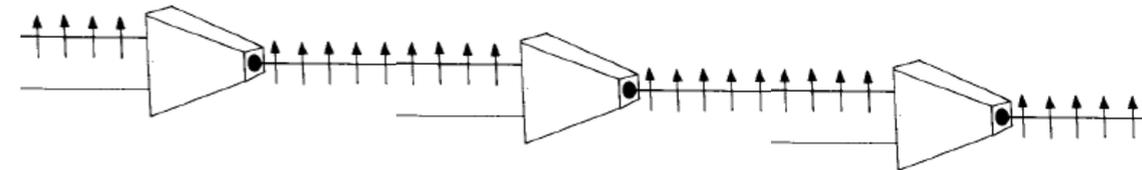
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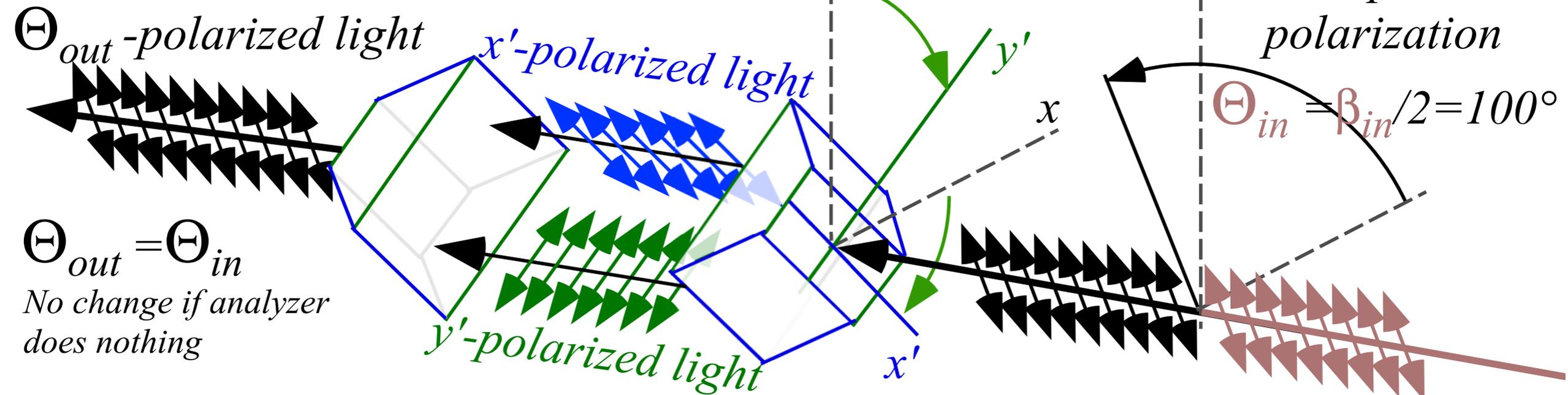
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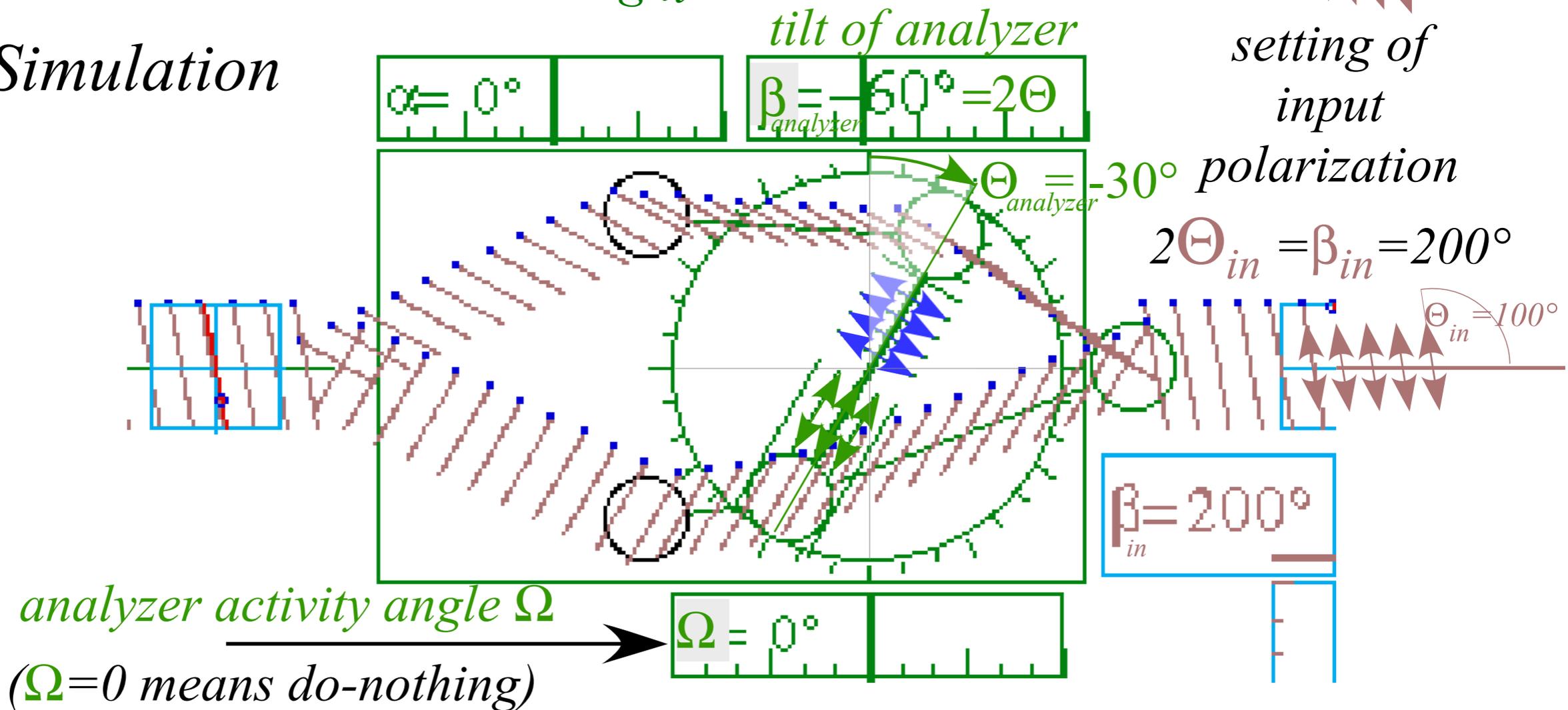
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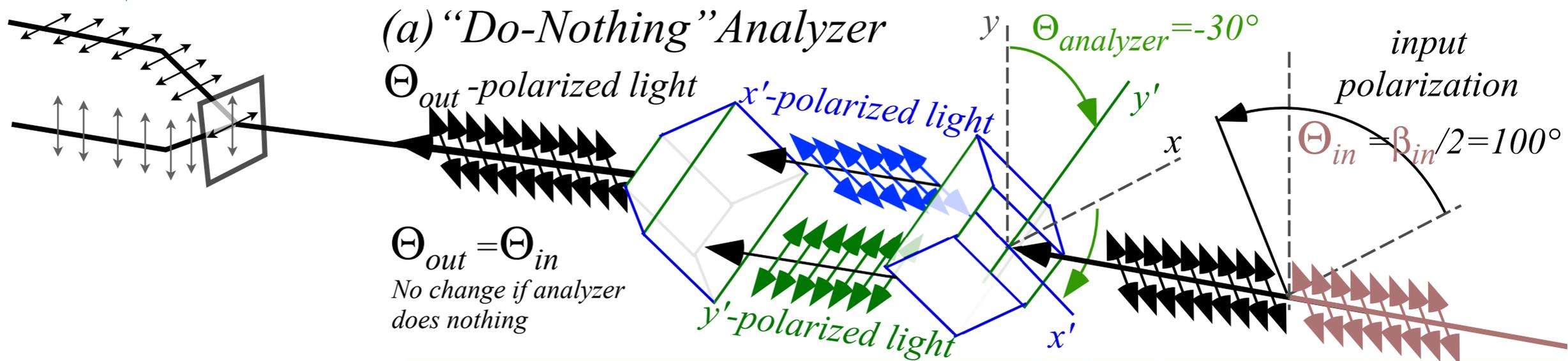
(a) "Do-Nothing" Analyzer



(b) Simulation



Imagine final  $xy$ -sorter analyzes output beam into  $x$  and  $y$ -components.



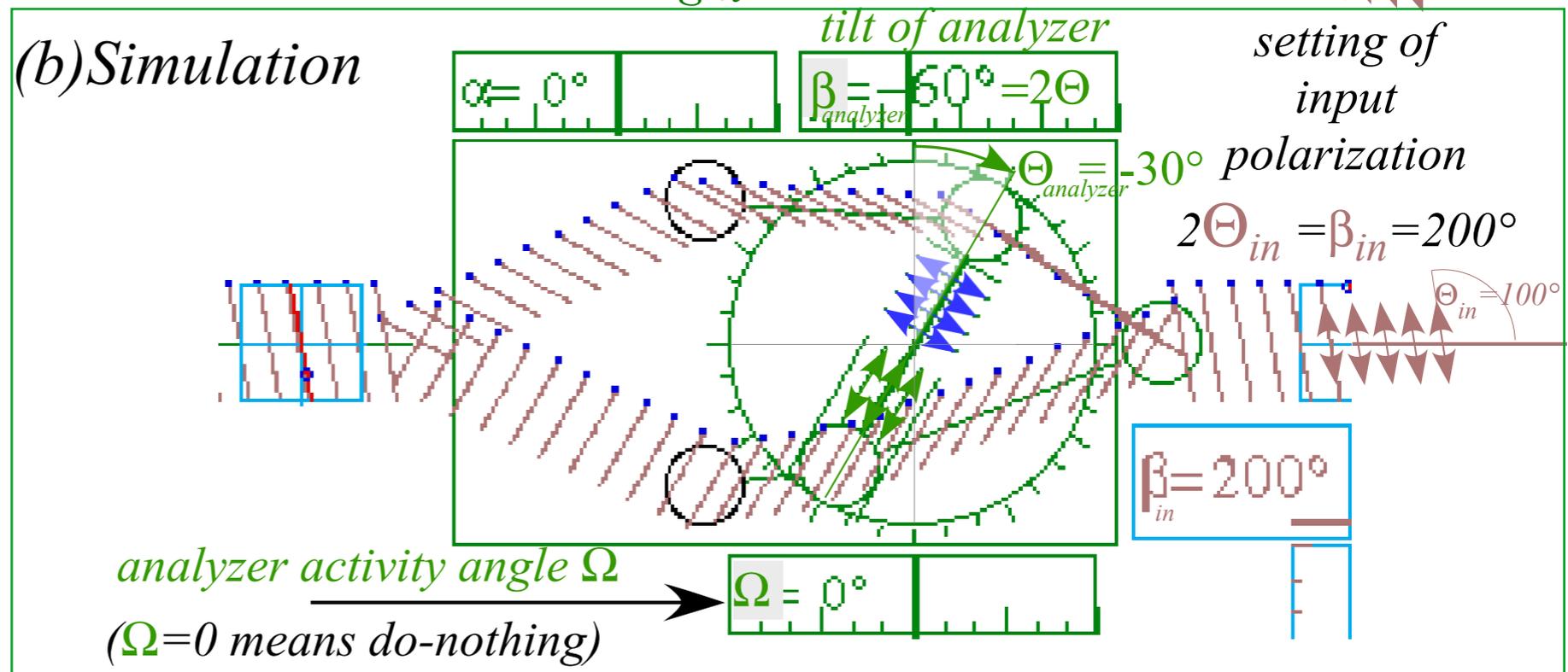
Amplitude in  $x$  or  $y$ -channel is sum over  $x'$  and  $y'$ -amplitudes

$\langle x' | \Theta_{in} \rangle = \cos(\Theta_{in} - \Theta)$

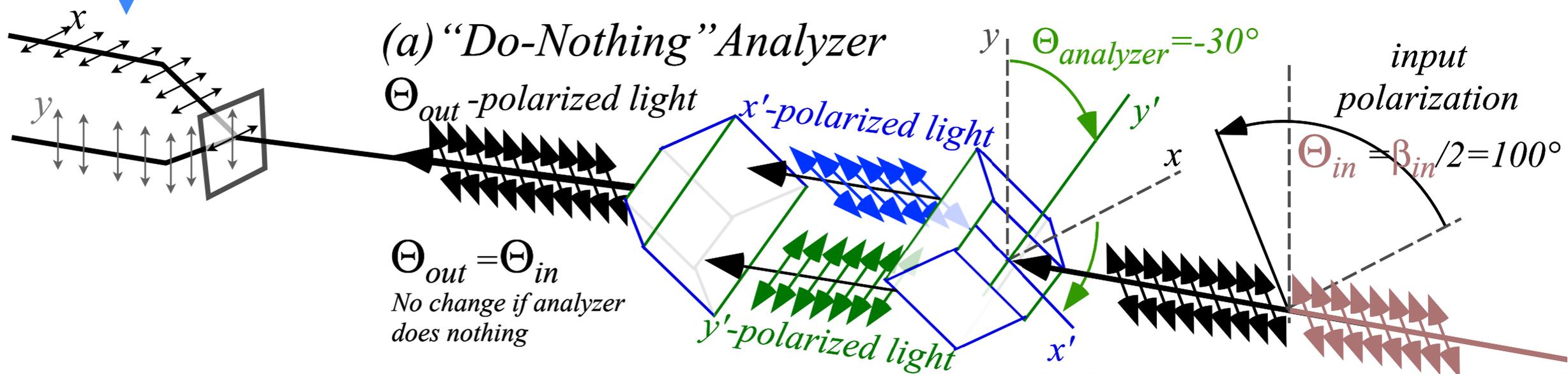
$\langle y' | \Theta_{in} \rangle = \sin(\Theta_{in} - \Theta)$

with relative angle  $\Theta_{in} - \Theta$

of  $\Theta_{in}$  to  $\Theta$ -analyzer axes- $(x', y')$



Imagine final  $xy$ -sorter analyzes output beam into  $x$  and  $y$ -components.



Amplitude in  $x$  or  $y$ -channel is sum over  $x'$  and  $y'$ -amplitudes

$$\langle x' | \Theta_{in} \rangle = \cos(\Theta_{in} - \Theta)$$

$$\langle y' | \Theta_{in} \rangle = \sin(\Theta_{in} - \Theta)$$

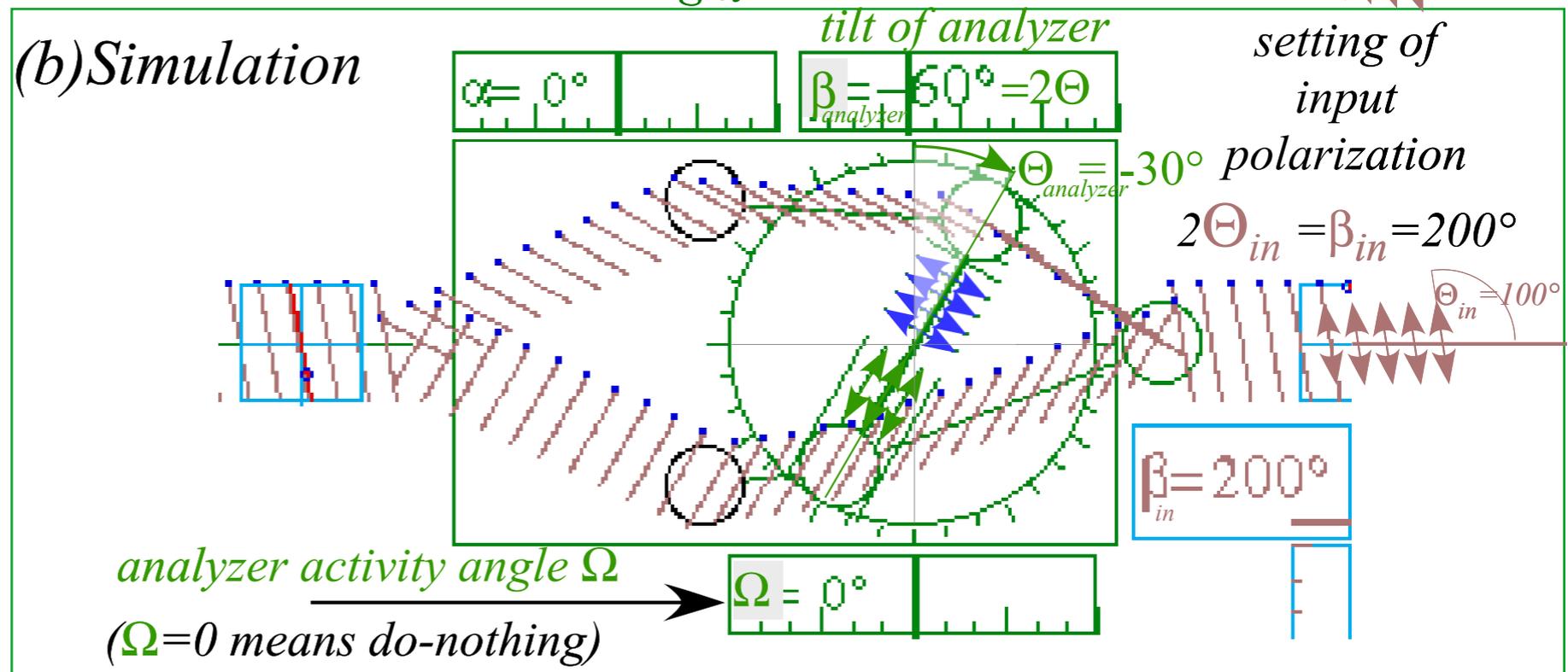
with relative angle  $\Theta_{in} - \Theta$

of  $\Theta_{in}$  to  $\Theta$ -analyzer axes- $(x', y')$

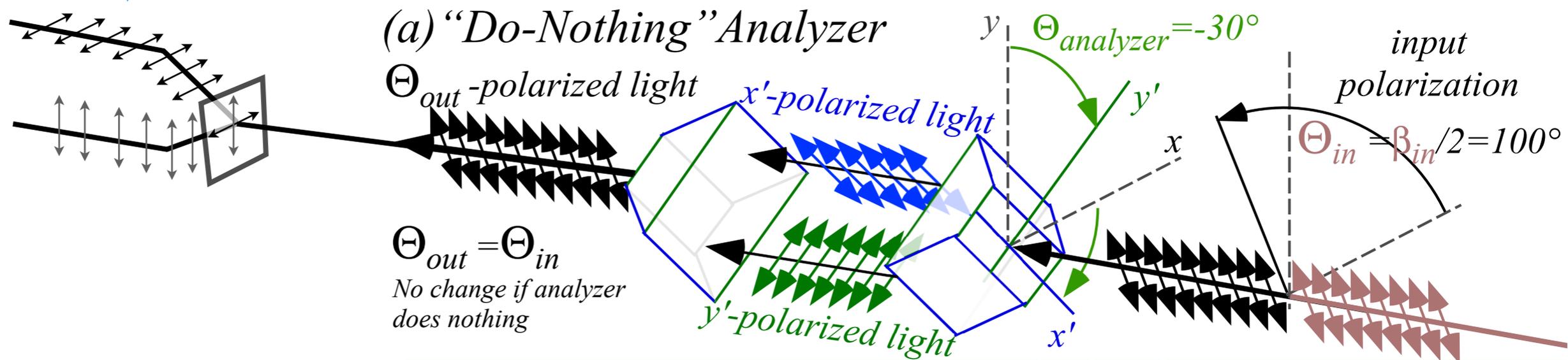
in products with final  $xy$ -sorter:

lab  $x$ -axis:  $\langle x | x' \rangle = \cos \Theta = \langle y | y' \rangle$

$y$ -axis:  $\langle y | x' \rangle = \sin \Theta = -\langle x | y' \rangle$ .



Imagine final  $xy$ -sorter analyzes output beam into  $x$  and  $y$ -components.



Amplitude in  $x$  or  $y$ -channel is sum over  $x'$  and  $y'$ -amplitudes

$$\langle x' | \Theta_{in} \rangle = \cos(\Theta_{in} - \Theta)$$

$$\langle y' | \Theta_{in} \rangle = \sin(\Theta_{in} - \Theta)$$

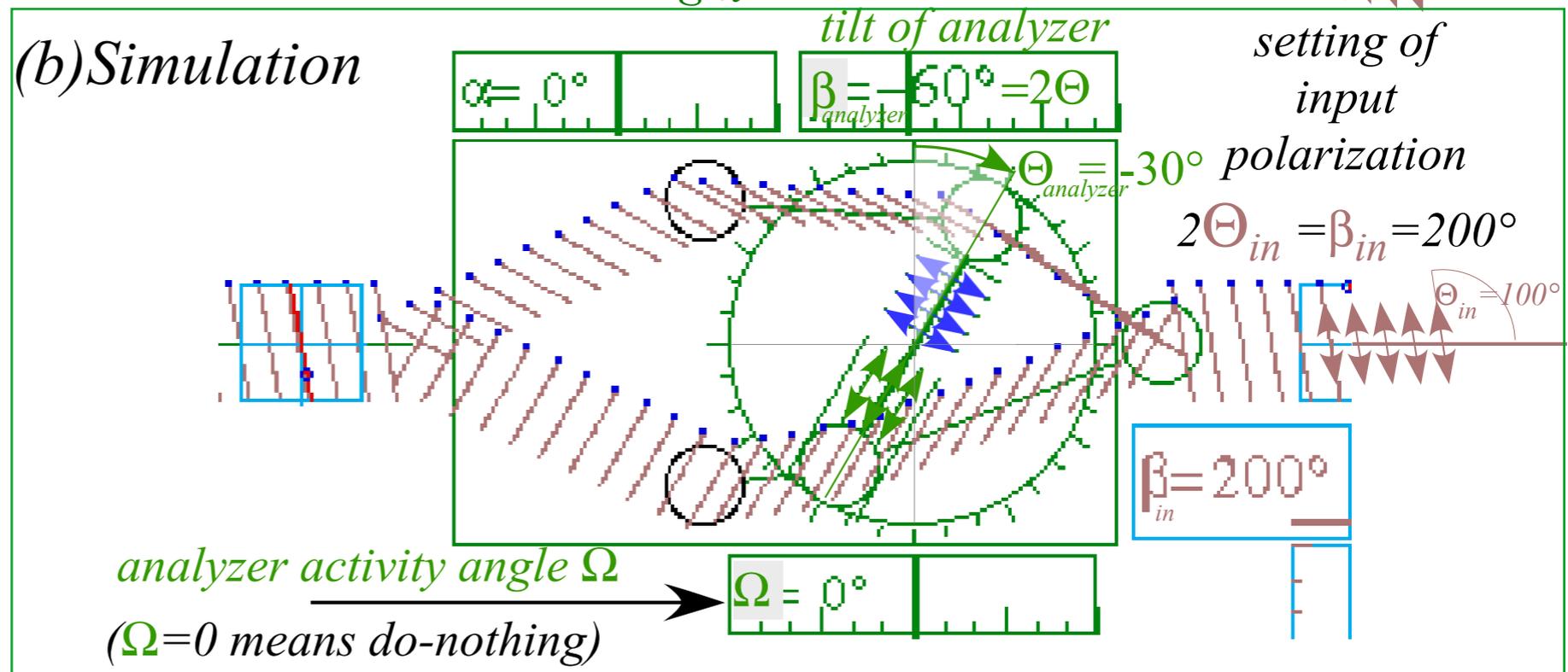
with relative angle  $\Theta_{in} - \Theta$

of  $\Theta_{in}$  to  $\Theta$ -analyzer axes- $(x', y')$

in products with final  $xy$ -sorter:

$$\text{lab } x\text{-axis: } \langle x | x' \rangle = \cos \Theta = \langle y | y' \rangle$$

$$\text{y-axis: } \langle y | x' \rangle = \sin \Theta = -\langle x | y' \rangle.$$

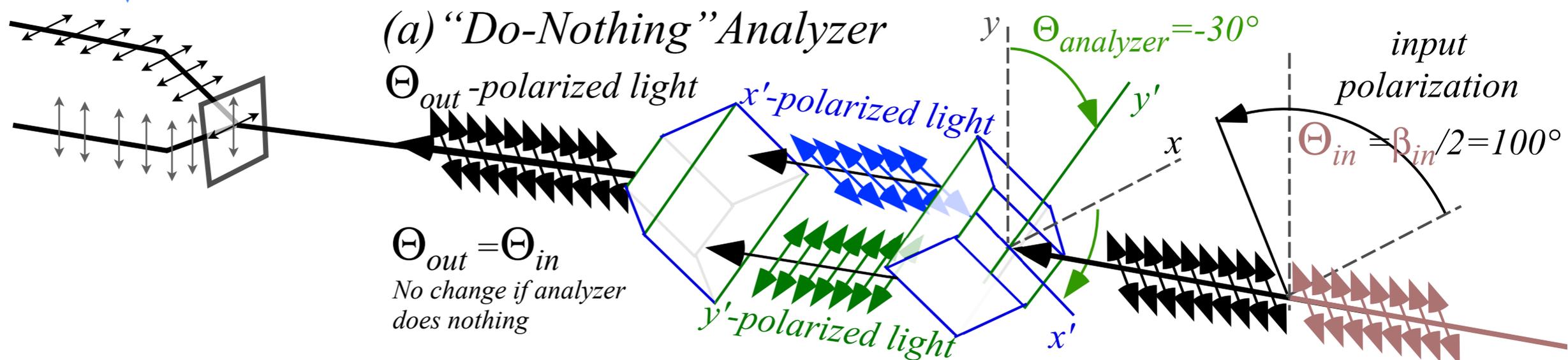


$$x\text{-Output is: } \langle x | \Theta_{out} \rangle = \langle x | x' \rangle \langle x' | \Theta_{in} \rangle + \langle x | y' \rangle \langle y' | \Theta_{in} \rangle = \cos \Theta \cos(\Theta_{in} - \Theta) - \sin \Theta \sin(\Theta_{in} - \Theta) = \cos \Theta_{in}$$

$$y\text{-Output is: } \langle y | \Theta_{out} \rangle = \langle y | x' \rangle \langle x' | \Theta_{in} \rangle + \langle y | y' \rangle \langle y' | \Theta_{in} \rangle = \sin \Theta \cos(\Theta_{in} - \Theta) - \cos \Theta \sin(\Theta_{in} - \Theta) = \sin \Theta_{in}.$$

(Recall  $\cos(a+b) = \cos a \cos b - \sin a \sin b$  and  $\sin(a+b) = \sin a \cos b + \cos a \sin b$ )

Imagine final  $xy$ -sorter analyzes output beam into  $x$  and  $y$ -components.



Amplitude in  $x$  or  $y$ -channel is sum over  $x'$  and  $y'$ -amplitudes

$$\langle x' | \Theta_{in} \rangle = \cos(\Theta_{in} - \Theta)$$

$$\langle y' | \Theta_{in} \rangle = \sin(\Theta_{in} - \Theta)$$

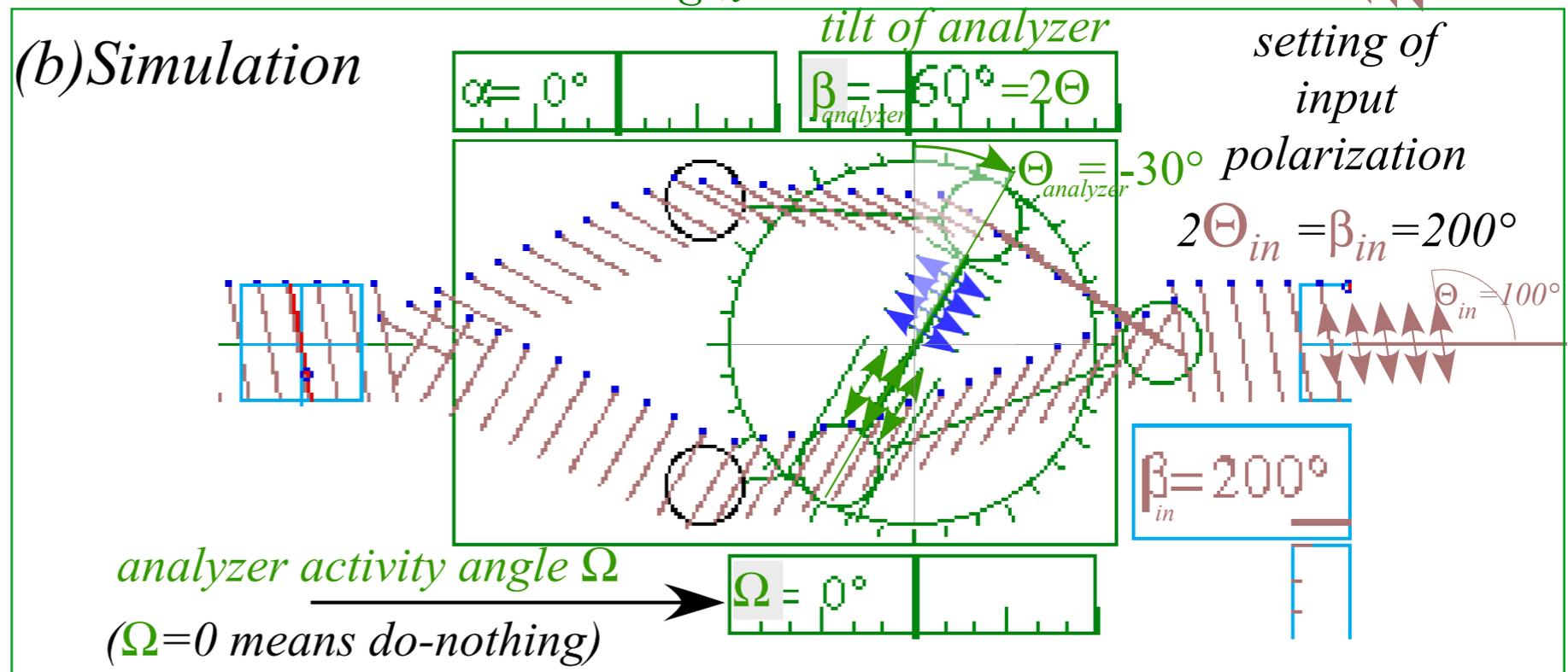
with relative angle  $\Theta_{in} - \Theta$

of  $\Theta_{in}$  to  $\Theta$ -analyzer axes- $(x', y')$

in products with final  $xy$ -sorter:

$$\text{lab } x\text{-axis: } \langle x | x' \rangle = \cos \Theta = \langle y | y' \rangle$$

$$y\text{-axis: } \langle y | x' \rangle = \sin \Theta = -\langle x | y' \rangle.$$



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$$y\text{-Output is: } \langle y | \Theta_{out} \rangle = \langle y | x' \rangle \langle x' | \Theta_{in} \rangle + \langle y | y' \rangle \langle y' | \Theta_{in} \rangle = \sin \Theta \cos(\Theta_{in} - \Theta) - \cos \Theta \sin(\Theta_{in} - \Theta) = \sin \Theta_{in}.$$

(Recall  $\cos(a+b) = \cos a \cos b - \sin a \sin b$  and  $\sin(a+b) = \sin a \cos b + \cos a \sin b$ )

Conclusion:

$$\langle x | \Theta_{out} \rangle = \cos \Theta_{out} = \cos \Theta_{in} \text{ or: } \Theta_{out} = \Theta_{in} \text{ so "Do-Nothing" Analyzer in fact does nothing.}$$

*Review: “Abstraction” of bra and ket vectors from a Transformation Matrix*

*Introducing scalar and matrix products*

*Planck's energy and N-quanta (Cavity/Beam wave mode)*

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*Sample calculations*

*Feynman amplitude axioms 2-3*

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*The “Do-Nothing” analyzer*

*Feynman amplitude axiom 4*

*Some “Do-Something” analyzers*

 *Sorter-counter, Filter, 1/2-wave plate, 1/4-wave plate*

# (1) Optical analyzer in sorter-counter configuration

Analyzer reduced to a simple sorter-counter by blocking output of  $x$ -high-road and  $y$ -low-road with counters

$$x\text{-counts} \sim |\langle x|x' \rangle|^2$$

$$= \cos^2 \theta = 0.75$$

$$y\text{-counts} \sim |\langle y|x' \rangle|^2$$

$$= \sin^2 \theta = 0.25$$

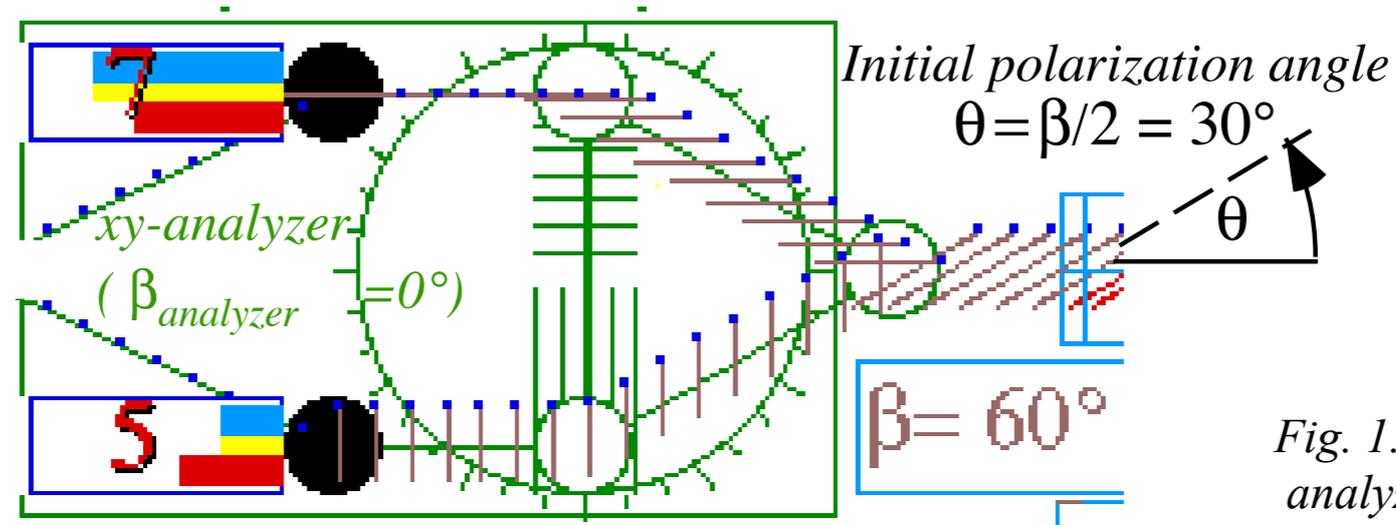


Fig. 1.3.3 Simulated polarization analyzer set up as a sorter-counter

## (1) Optical analyzer in sorter-counter configuration

Analyzer reduced to a simple sorter-counter by blocking output of  $x$ -high-road and  $y$ -low-road with counters

$$x\text{-counts} \sim |\langle x|x' \rangle|^2 = \cos^2 \theta = 0.75$$

$$y\text{-counts} \sim |\langle y|x' \rangle|^2 = \sin^2 \theta = 0.25$$

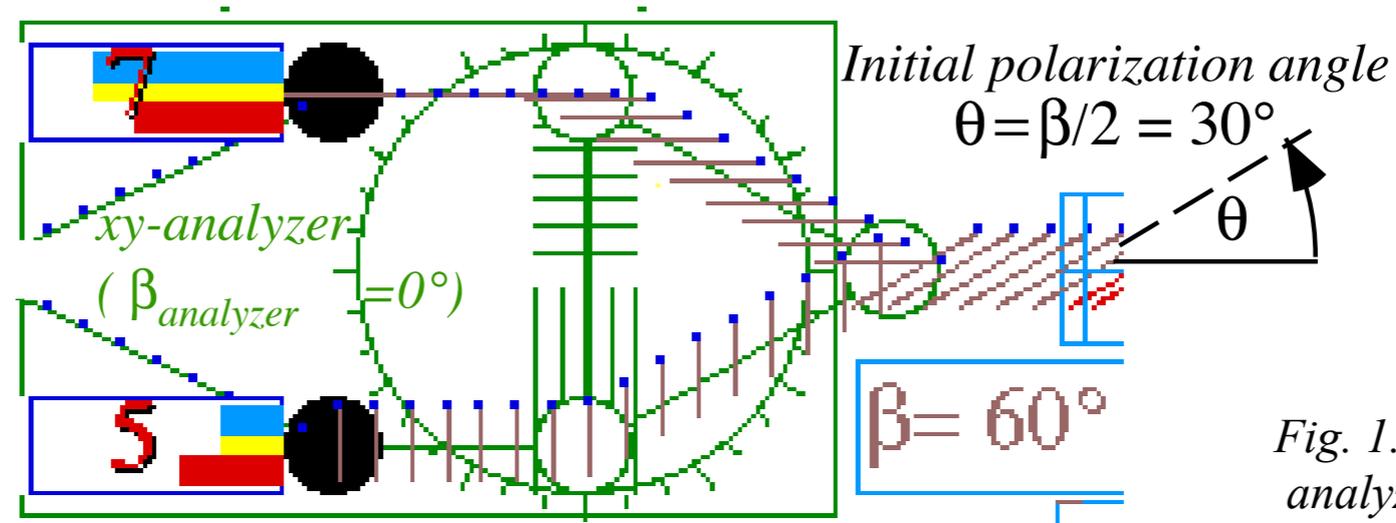


Fig. 1.3.3 Simulated polarization analyzer set up as a sorter-counter

## (2) Optical analyzer in a filter configuration (Polaroid© sunglasses)

Analyzer blocks one path which may have photon counter without affecting function.

$$x\text{-counts} \sim |\langle y|x' \rangle|^2 = 0.75$$

(Blocked and filtered out)

$$y\text{-output} \sim |\langle y|x' \rangle|^2 = \sin^2 \theta = 0.25$$

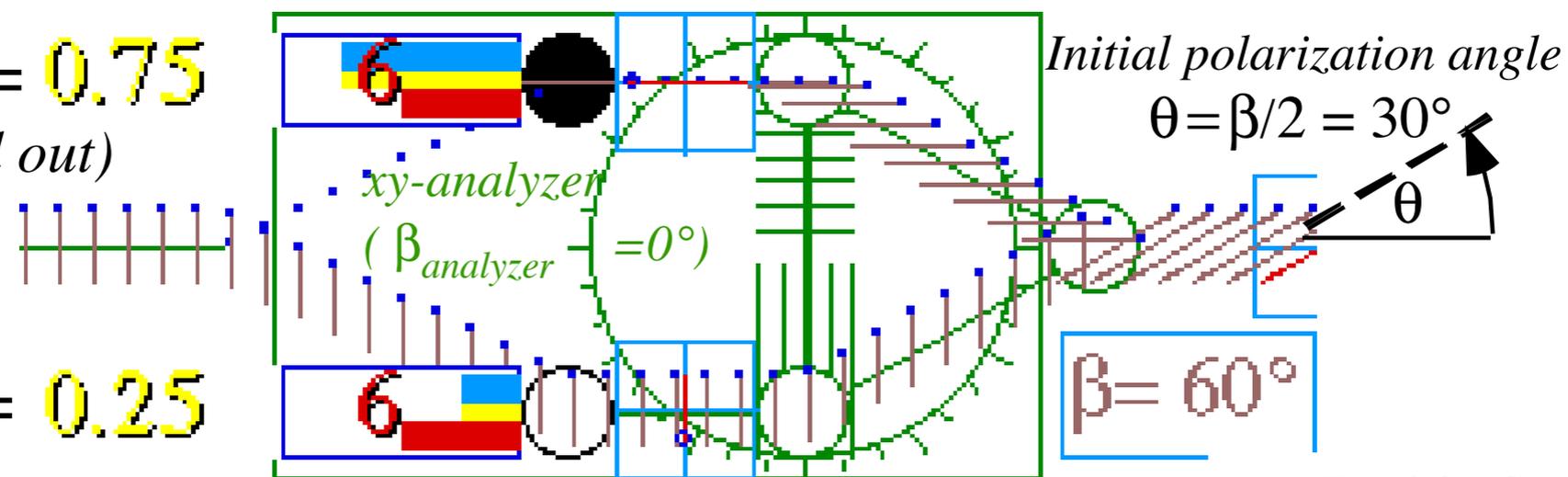


Fig. 1.3.4 Simulated polarization analyzer set up to filter out the  $x$ -polarized photons

*Review: “Abstraction” of bra and ket vectors from a Transformation Matrix*

*Introducing scalar and matrix products*

*Planck's energy and N-quanta (Cavity/Beam wave mode)*

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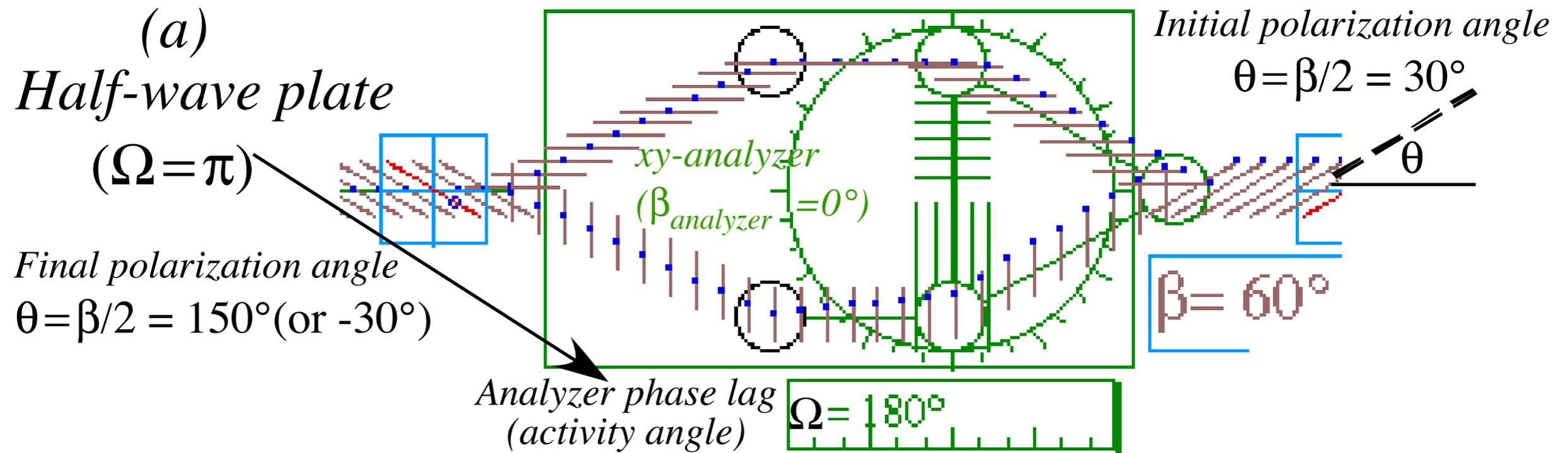
*Feynman amplitude axiom 4*

*Some “Do-Something” analyzers*

*Sorter-counter, Filter, 1/2-wave plate, 1/4-wave plate*



**(3) Optical analyzers in the "control" configuration: Half or Quarter wave plates**



**(3) Optical analyzers in the "control" configuration: Half or Quarter wave plates**

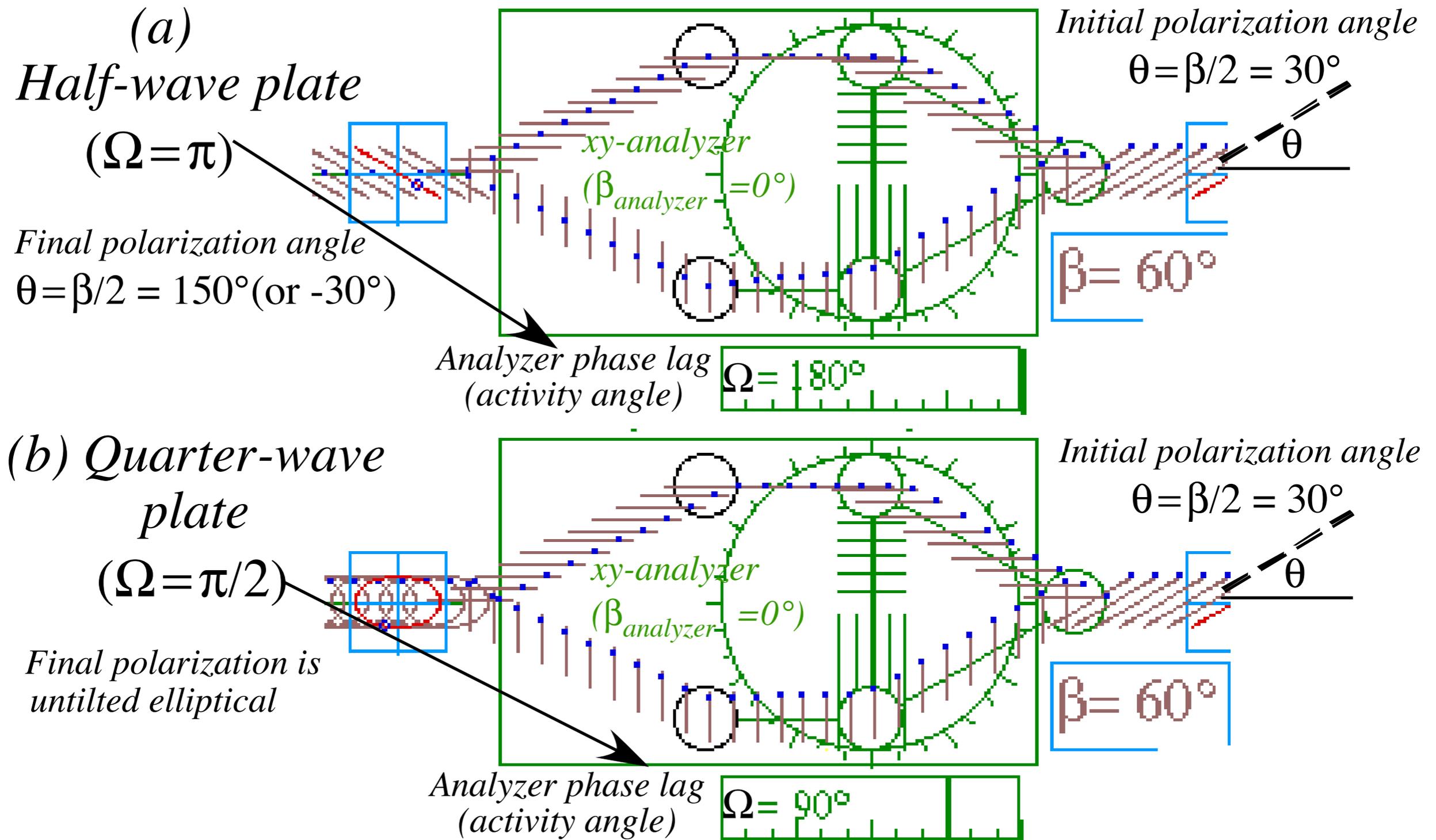
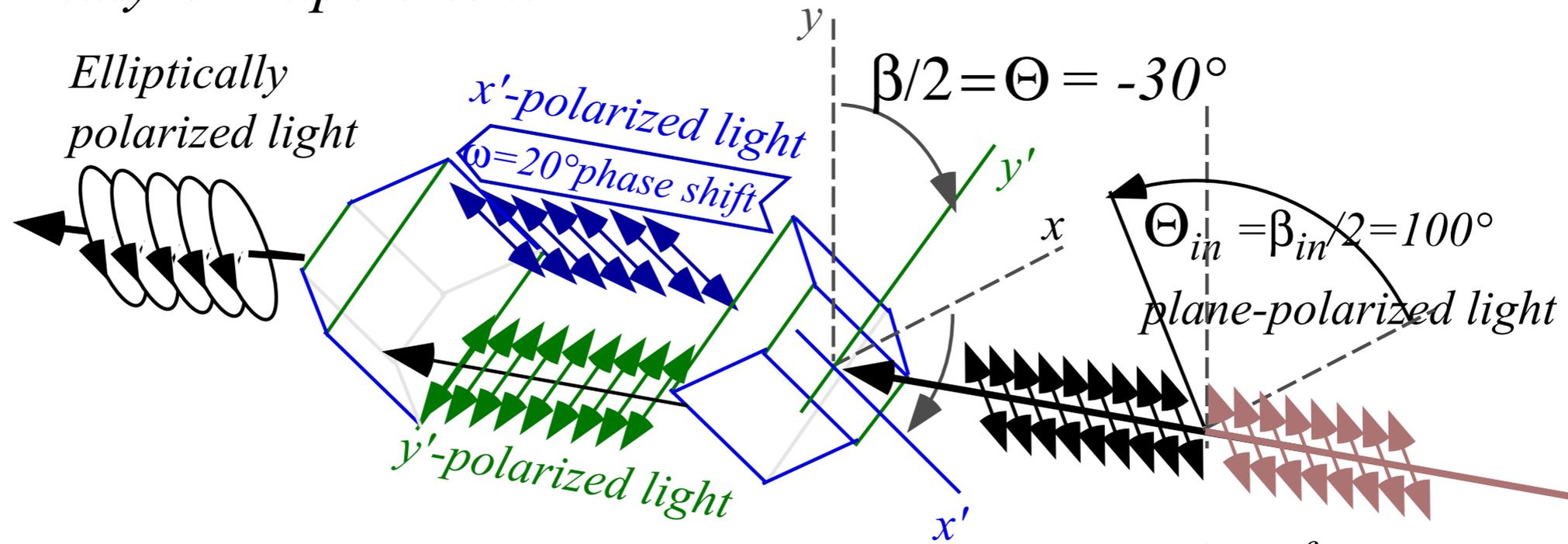
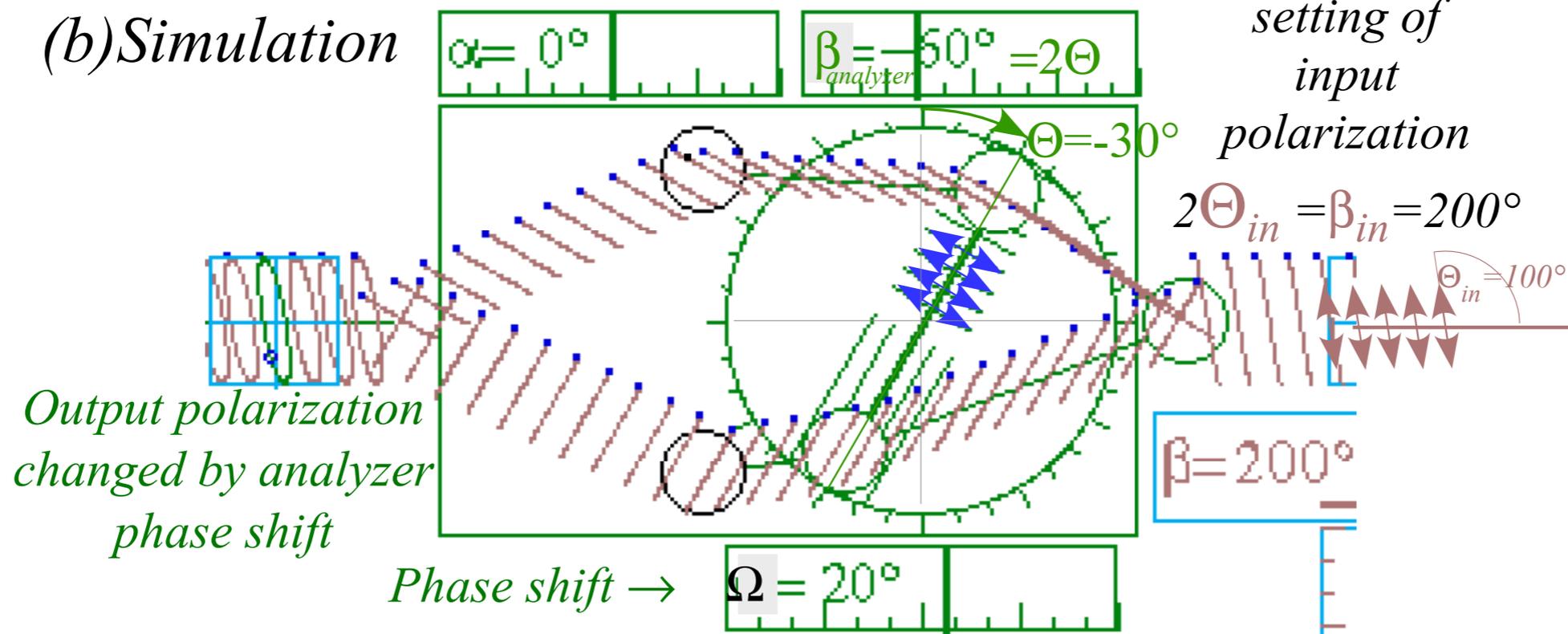


Fig. 1.3.5 Polarization control set to shift phase by (a) Half-wave ( $\Omega = \pi$ ) , (b) Quarter wave ( $\Omega = \pi/2$ )

### (a) Analyzer Experiment



### (b) Simulation



Similar to "do-nothing" analyzer but has extra phase factor  $e^{-i\Omega} = 0.94 - i 0.34$  on the  $x'$ -path .

$$x\text{-output: } \langle x | \Psi_{out} \rangle = \langle x | x' \rangle e^{-i\Omega} \langle x' | \Psi_{in} \rangle + \langle x | y' \rangle \langle y' | \Psi_{in} \rangle = e^{-i\Omega} \cos \Theta \cos(\Theta_{in} - \Theta) - \sin \Theta \sin(\Theta_{in} - \Theta)$$

$$y\text{-output: } \langle y | \Psi_{out} \rangle = \langle y | x' \rangle e^{-i\Omega} \langle x' | \Psi_{in} \rangle + \langle y | y' \rangle \langle y' | \Psi_{in} \rangle = e^{-i\Omega} \sin \Theta \cos(\Theta_{in} - \Theta) + \cos \Theta \sin(\Theta_{in} - \Theta)$$

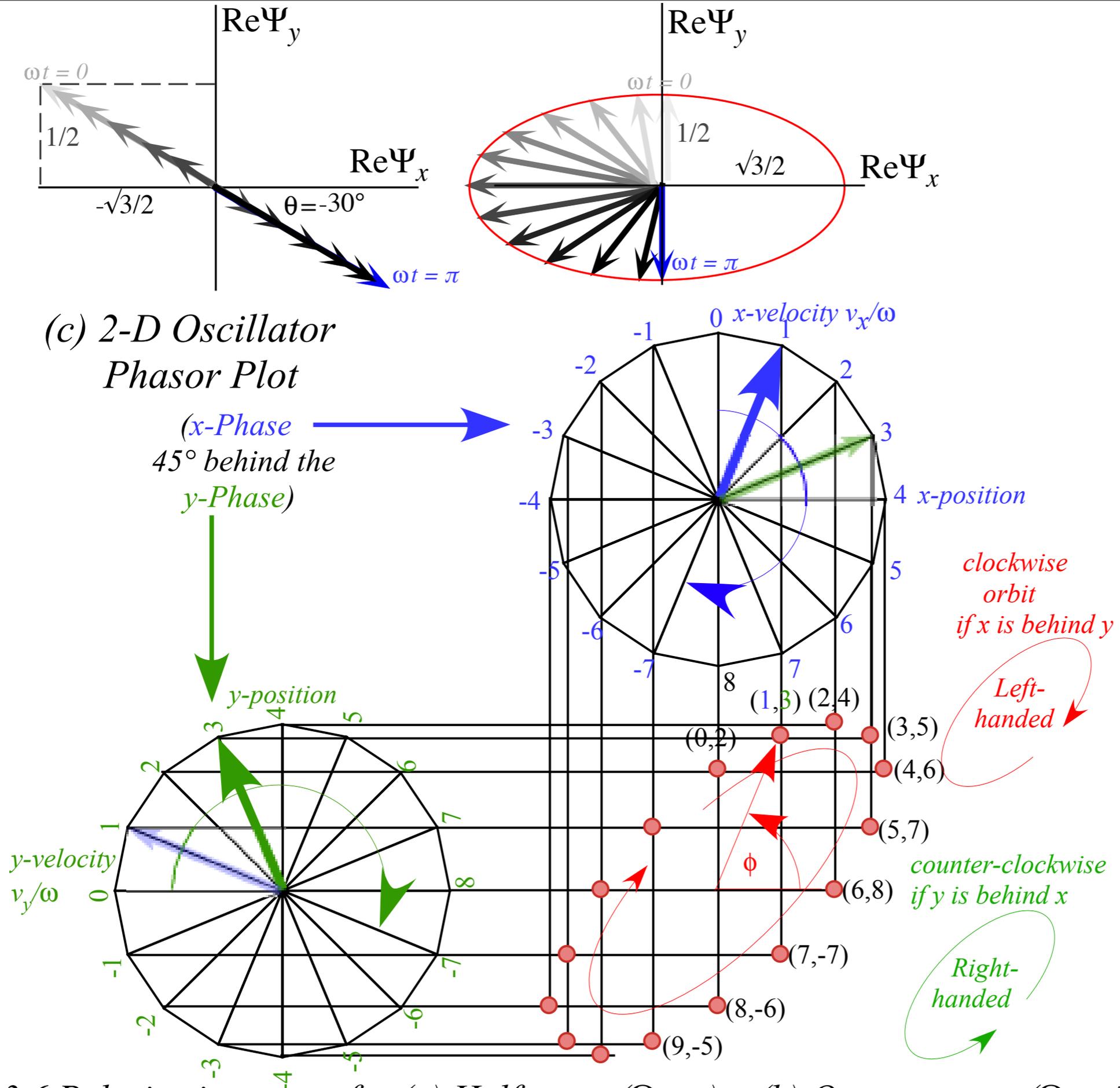
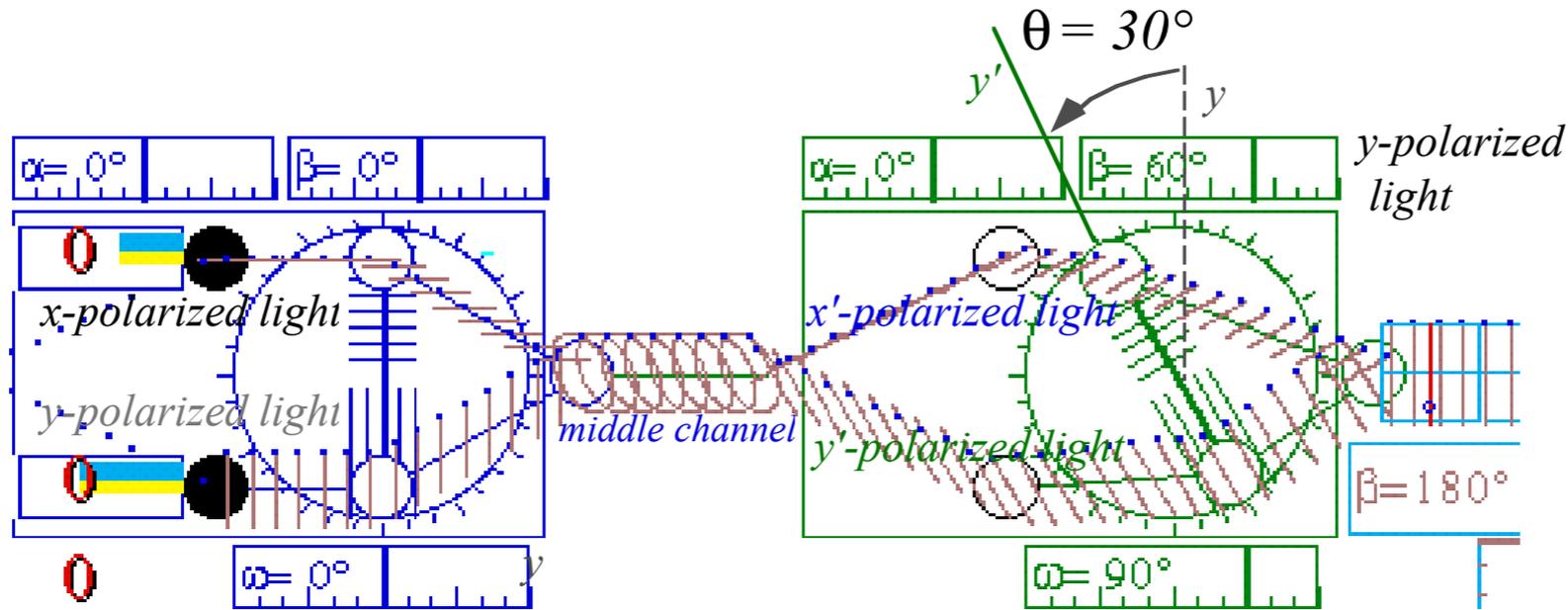


Fig. 1.3.6 Polarization states for (a) Half-wave ( $\Omega=\pi$ ) , (b) Quarter wave ( $\Omega=\pi/2$ ) (c) ( $\Omega=-\pi/4$ )

Ellipsis in Middle

1.3.1 A y-polarized light beam of unit amplitude (1 photon/sec.) enters an active analyzer that is tilted by 30° as shown below. The active analyzer puts a  $\omega = 90^\circ$  phase factor  $e^{-i\omega}$  in the x' beam.

Fill in the blanks with numbers or symbols that tell as much as possible about what is present at each channel or branch.



$$\begin{vmatrix} \langle x|x' \rangle & \langle x|y' \rangle \\ \langle y|x' \rangle & \langle y|y' \rangle \end{vmatrix} = \begin{vmatrix} \text{---} & \text{---} \\ \text{---} & \text{---} \end{vmatrix} \quad \begin{vmatrix} \langle x'|x \rangle & \langle x'|y \rangle \\ \langle y'|x \rangle & \langle y'|y \rangle \end{vmatrix} = \begin{vmatrix} \text{---} & \text{---} \\ \text{---} & \text{---} \end{vmatrix}$$

State of output x channel

$|\text{---}\rangle$   
 Amplitude=  
 Probability=  
 \_\_\_\_\_

State of x' channel

$|\text{---}\rangle$   
 Amplitude=  
 Probability=  
 \_\_\_\_\_

State of input channel

$|y\rangle$  1.0  
 Amplitude=  
1.0  
 Probability=  
1.0

State of middle channel

Probability=  
 \_\_\_\_\_

State of output y channel

$|\text{---}\rangle$   
 Amplitude=  
 Probability=  
 \_\_\_\_\_

State of y' channel

$|\text{---}\rangle$   
 Amplitude=  
 Probability=  
 \_\_\_\_\_

Note: The x' and y' quantities will not be graded. Use for work.