

Group Theory in Quantum Mechanics

Lecture 22 (4.14.15)

Harmonic oscillator symmetry $U(1) \subset U(2) \subset U(3) \dots$

(Int.J.Mol.Sci, 14, 714(2013) p.755-774 , QTCA Unit 7 Ch. 20-22)

(PSDS - Ch. 8)

1-D $\mathbf{a}^\dagger \mathbf{a}$ algebra of $U(1)$ representations

Creation-Destruction $\mathbf{a}^\dagger \mathbf{a}$ algebra

Eigenstate creationism (and destruction)

Vacuum state

1st excited state

Normal ordering for matrix calculation

Commutator derivative identities

Binomial expansion identities

Matrix $\langle \mathbf{a}^n \mathbf{a}^{\dagger n} \rangle$ calculations

Number operator and Hamiltonian operator

Expectation values of position, momentum, and uncertainty for eigenstate $|n\rangle$

Harmonic oscillator beat dynamics of mixed states

Oscillator coherent states (“Shoved” and “kicked” states)

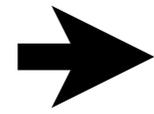
Translation operators vs. boost operators

Applying boost-translation combinations

Time evolution of coherent state

Properties of coherent state and “squeezed” states

2-D $\mathbf{a}^\dagger \mathbf{a}$ algebra of $U(2)$ representations and $R(3)$ angular momentum operators



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Q: How to convert *classical* HO Hamiltonian to *quantum* HO Hamiltonian?

$$E = H(x, p) = \frac{1}{2M} p^2 + \frac{1}{2} M \omega^2 x^2$$

1-D $\mathfrak{a}^\dagger \mathfrak{a}$ algebra of $U(1)$ representations

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$$\mathbf{H}(\mathbf{x}, \mathbf{p}) = \mathbf{p}^2/2M + V(\mathbf{x}) = \mathbf{p}^2/2M + M\omega^2 \mathbf{x}^2/2 \quad (\text{with: } \mathbf{p} = \hbar \mathbf{k})$$

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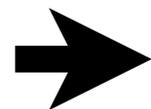
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Define

Destruction operator

and

$$\mathbf{a}^\dagger = \frac{(\mathbf{X} - i\mathbf{P})}{\sqrt{\hbar\omega}} = \frac{(\sqrt{M\omega} \mathbf{x} - i\mathbf{p} / \sqrt{M\omega})}{\sqrt{2\hbar}}$$

Creation Operator

Commutation relations between $\mathbf{a} = (\mathbf{X} + i\mathbf{P})/2$ and $\mathbf{a}^\dagger = (\mathbf{X} - i\mathbf{P})/2$ with $\mathbf{X} \equiv \sqrt{M\omega} \mathbf{x} / \sqrt{2}$ and $\mathbf{P} \equiv \mathbf{p} / \sqrt{2M}$:

$$[\mathbf{a}, \mathbf{a}^\dagger] \equiv \mathbf{a}\mathbf{a}^\dagger - \mathbf{a}^\dagger\mathbf{a} = \frac{1}{2\hbar} (\sqrt{M\omega} \mathbf{x} + i\mathbf{p} / \sqrt{M\omega}) (\sqrt{M\omega} \mathbf{x} - i\mathbf{p} / \sqrt{M\omega}) - \frac{1}{2\hbar} (\sqrt{M\omega} \mathbf{x} - i\mathbf{p} / \sqrt{M\omega}) (\sqrt{M\omega} \mathbf{x} + i\mathbf{p} / \sqrt{M\omega})$$

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$$[\mathbf{a}, \mathbf{a}^\dagger] = \frac{2i}{2\hbar} (\mathbf{p}\mathbf{x} - \mathbf{x}\mathbf{p}) = \frac{-i}{\hbar} [\mathbf{x}, \mathbf{p}]$$

Creation-Destruction $\mathbf{a}^\dagger \mathbf{a}$ algebra

$$\mathbf{a} = \frac{(\mathbf{X} + i\mathbf{P})}{\sqrt{\hbar\omega}} = \frac{(\sqrt{M\omega} \mathbf{x} + i\mathbf{p} / \sqrt{M\omega})}{\sqrt{2\hbar}}$$

Define

Destruction operator

and

$$\mathbf{a}^\dagger = \frac{(\mathbf{X} - i\mathbf{P})}{\sqrt{\hbar\omega}} = \frac{(\sqrt{M\omega} \mathbf{x} - i\mathbf{p} / \sqrt{M\omega})}{\sqrt{2\hbar}}$$

Creation Operator

Commutation relations between $\mathbf{a} = (\mathbf{X} + i\mathbf{P})/2$ and $\mathbf{a}^\dagger = (\mathbf{X} - i\mathbf{P})/2$ with $\mathbf{X} \equiv \sqrt{M\omega} \mathbf{x} / \sqrt{2}$ and $\mathbf{P} \equiv \mathbf{p} / \sqrt{2M}$:

$$[\mathbf{a}, \mathbf{a}^\dagger] \equiv \mathbf{a}\mathbf{a}^\dagger - \mathbf{a}^\dagger\mathbf{a} = \frac{1}{2\hbar} \left(\sqrt{M\omega} \mathbf{x} + i\mathbf{p} / \sqrt{M\omega} \right) \left(\sqrt{M\omega} \mathbf{x} - i\mathbf{p} / \sqrt{M\omega} \right) - \frac{1}{2\hbar} \left(\sqrt{M\omega} \mathbf{x} - i\mathbf{p} / \sqrt{M\omega} \right) \left(\sqrt{M\omega} \mathbf{x} + i\mathbf{p} / \sqrt{M\omega} \right)$$

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Recall *commutator* $[\mathbf{x}, \mathbf{p}]$ relation: $[\mathbf{x}, \mathbf{p}] \equiv \mathbf{x}\mathbf{p} - \mathbf{p}\mathbf{x} = \hbar i \mathbf{1}$

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1D-HO Hamiltonian in terms of $\mathbf{a}^\dagger \mathbf{a}$ operator

Recall: $\mathbf{H}(\mathbf{x}, \mathbf{p}) = \hbar\omega (\mathbf{a}^\dagger \mathbf{a} + \mathbf{a}\mathbf{a}^\dagger)/2$

Recall *commutator* $[\mathbf{x}, \mathbf{p}]$ relation: $[\mathbf{x}, \mathbf{p}] \equiv \mathbf{x}\mathbf{p} - \mathbf{p}\mathbf{x} = \hbar i \mathbf{1}$

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1D-HO Hamiltonian in terms of $\mathbf{a}^\dagger \mathbf{a}$ operator

$$\mathbf{H}(\mathbf{x}, \mathbf{p}) = \hbar\omega (\mathbf{a}^\dagger \mathbf{a} + \mathbf{a}\mathbf{a}^\dagger)/2 = \hbar\omega (\mathbf{a}^\dagger \mathbf{a} + \mathbf{a}^\dagger \mathbf{a} + \mathbf{1})/2$$

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1D-HO Hamiltonian in terms of $\mathbf{a}^\dagger \mathbf{a}$ operator

$$\mathbf{H}(\mathbf{x}, \mathbf{p}) = \hbar\omega (\mathbf{a}^\dagger \mathbf{a} + \mathbf{a}\mathbf{a}^\dagger)/2 = \hbar\omega (\mathbf{a}^\dagger \mathbf{a} + \mathbf{a}^\dagger \mathbf{a} + \mathbf{1})/2 = \hbar\omega \mathbf{a}^\dagger \mathbf{a} + \mathbf{1}\hbar\omega/2$$

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1-D $\mathbf{a}^\dagger\mathbf{a}$ algebra of $U(1)$ representations

Creation-Destruction $\mathbf{a}^\dagger\mathbf{a}$ algebra

Eigenstate creationism (and destruction)

Vacuum state

1st excited state

Normal ordering for matrix calculation

Commutator derivative identities

Binomial expansion identities

Matrix $\langle \mathbf{a}^n \mathbf{a}^{\dagger n} \rangle$ calculations

Number operator and Hamiltonian operator

Expectation values of position, momentum, and uncertainty for eigenstate $|n\rangle$

Harmonic oscillator beat dynamics of mixed states

Oscillator coherent states (“Shoved” and “kicked” states)

Translation operators vs. boost operators

Applying boost-translation combinations

Time evolution of coherent state

Properties of coherent state and “squeezed” states

2-D $\mathbf{a}^\dagger\mathbf{a}$ algebra of $U(2)$ representations and $R(3)$ angular momentum operators



Eigenstate creationism (and destruction)

Given 1D-HO Hamiltonian: $\mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar\omega \mathbf{a}^\dagger \mathbf{a} + \mathbf{1}\hbar\omega/2$ and commutation: $[\mathbf{a}, \mathbf{a}^\dagger] = \mathbf{1}$ or $\mathbf{a}\mathbf{a}^\dagger = \mathbf{a}^\dagger \mathbf{a} + \mathbf{1}$

Define *ground state* $|0\rangle$ as the eigenstate of $\mathbf{H}(\mathbf{x},\mathbf{p})$ with the *zero point eigenvalue* $E_0 = \hbar\omega/2$.

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QED:

$$\mathbf{H}(\mathbf{x},\mathbf{p}) |1\rangle = (\hbar\omega + \hbar\omega/2) |1\rangle = E_1 |1\rangle \text{ where: } E_1 = \hbar\omega + E_0$$

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One-quantum or *1st excited eigenket* $|1\rangle = \mathbf{a}^\dagger |0\rangle$

Eigenstate creationism (and destruction)

Given 1D-HO Hamiltonian: $\mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar\omega \mathbf{a}^\dagger \mathbf{a} + \mathbf{1}\hbar\omega/2$ and commutation: $[\mathbf{a}, \mathbf{a}^\dagger] = \mathbf{1}$ or $\mathbf{a}\mathbf{a}^\dagger = \mathbf{a}^\dagger \mathbf{a} + \mathbf{1}$

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One-quantum or *1st excited eigenket* $|1\rangle = \mathbf{a}^\dagger |0\rangle$

For kets, \mathbf{a}^\dagger is *creation operator* while \mathbf{a} is *destruction operator*.

$$\mathbf{a} |1\rangle = \mathbf{a} \mathbf{a}^\dagger |0\rangle = (\mathbf{a}^\dagger \mathbf{a} + \mathbf{1}) |0\rangle = |0\rangle$$

Eigenstate creationism (and destruction)

Given 1D-HO Hamiltonian: $\mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar\omega \mathbf{a}^\dagger \mathbf{a} + \mathbf{1}\hbar\omega/2$ and commutation: $[\mathbf{a}, \mathbf{a}^\dagger] = \mathbf{1}$ or $\mathbf{a}\mathbf{a}^\dagger = \mathbf{a}^\dagger \mathbf{a} + \mathbf{1}$

Define *ground state* $|0\rangle$ as the eigenstate of $\mathbf{H}(\mathbf{x},\mathbf{p})$ with the *zero point eigenvalue* $E_0 = \hbar\omega/2$.

$$\mathbf{H}(\mathbf{x},\mathbf{p}) |0\rangle = \hbar\omega/2 |0\rangle \quad \langle 0| \mathbf{H}(\mathbf{x},\mathbf{p}) = \hbar\omega/2 \langle 0|$$

Action by \mathbf{a} on ground ket $|0\rangle$ (or \mathbf{a}^\dagger on ground bra $\langle 0|$) gives *nothing* (zero vectors $\mathbf{0}$).

$$\mathbf{a} |0\rangle = \mathbf{0} \quad \langle 0| \mathbf{a}^\dagger = \mathbf{0}$$

But, \mathbf{a}^\dagger acts on ground ket to give $|1\rangle = \mathbf{a}^\dagger |0\rangle$ with \mathbf{H} eigenvalue $E_1 = \hbar\omega + E_0$. ($|1\rangle = \mathbf{a}^\dagger |0\rangle$, $\langle 0| \mathbf{a} = \langle 1|$.)

Proof:

$$\mathbf{H}(\mathbf{x},\mathbf{p}) \mathbf{a}^\dagger |0\rangle = \hbar\omega \mathbf{a}^\dagger \mathbf{a} \mathbf{a}^\dagger |0\rangle + \hbar\omega/2 \mathbf{a}^\dagger |0\rangle$$

$$\begin{aligned} \mathbf{H}(\mathbf{x},\mathbf{p}) \mathbf{a}^\dagger |0\rangle &= \hbar\omega \mathbf{a}^\dagger (\mathbf{a}^\dagger \mathbf{a} + \mathbf{1}) |0\rangle + \hbar\omega/2 \mathbf{a}^\dagger |0\rangle \\ &= \hbar\omega \mathbf{a}^\dagger |0\rangle + \mathbf{0} + \hbar\omega/2 \mathbf{a}^\dagger |0\rangle \end{aligned}$$

QED:

$$\mathbf{H}(\mathbf{x},\mathbf{p}) |1\rangle = (\hbar\omega + \hbar\omega/2) |1\rangle = E_1 |1\rangle \text{ where: } E_1 = \hbar\omega + E_0$$

One-quantum or *1st excited eigenket* $|1\rangle = \mathbf{a}^\dagger |0\rangle$

For kets, \mathbf{a}^\dagger is *creation operator* while \mathbf{a} is *destruction operator*.

$$\mathbf{a} |1\rangle = \mathbf{a} \mathbf{a}^\dagger |0\rangle = (\mathbf{a}^\dagger \mathbf{a} + \mathbf{1}) |0\rangle = |0\rangle$$

For bras, \mathbf{a}^\dagger is *destruction operator* while \mathbf{a} is *creation operator*.

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Vacuum state

1st excited state



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Wavefunction creationism (Vacuum state)

Coordinate representation of the “nothing” equation $\langle x|\mathbf{a}|0\rangle = 0$

with: $\mathbf{p} = \hbar\mathbf{k} = \frac{\hbar}{i} \frac{\partial}{\partial x}$

$$\langle x|\mathbf{a}|0\rangle = \frac{1}{\sqrt{2\hbar}} \left(\sqrt{M\omega} \langle x|\mathbf{x}|0\rangle + i \langle x|\mathbf{p}|0\rangle / \sqrt{M\omega} \right) = 0$$

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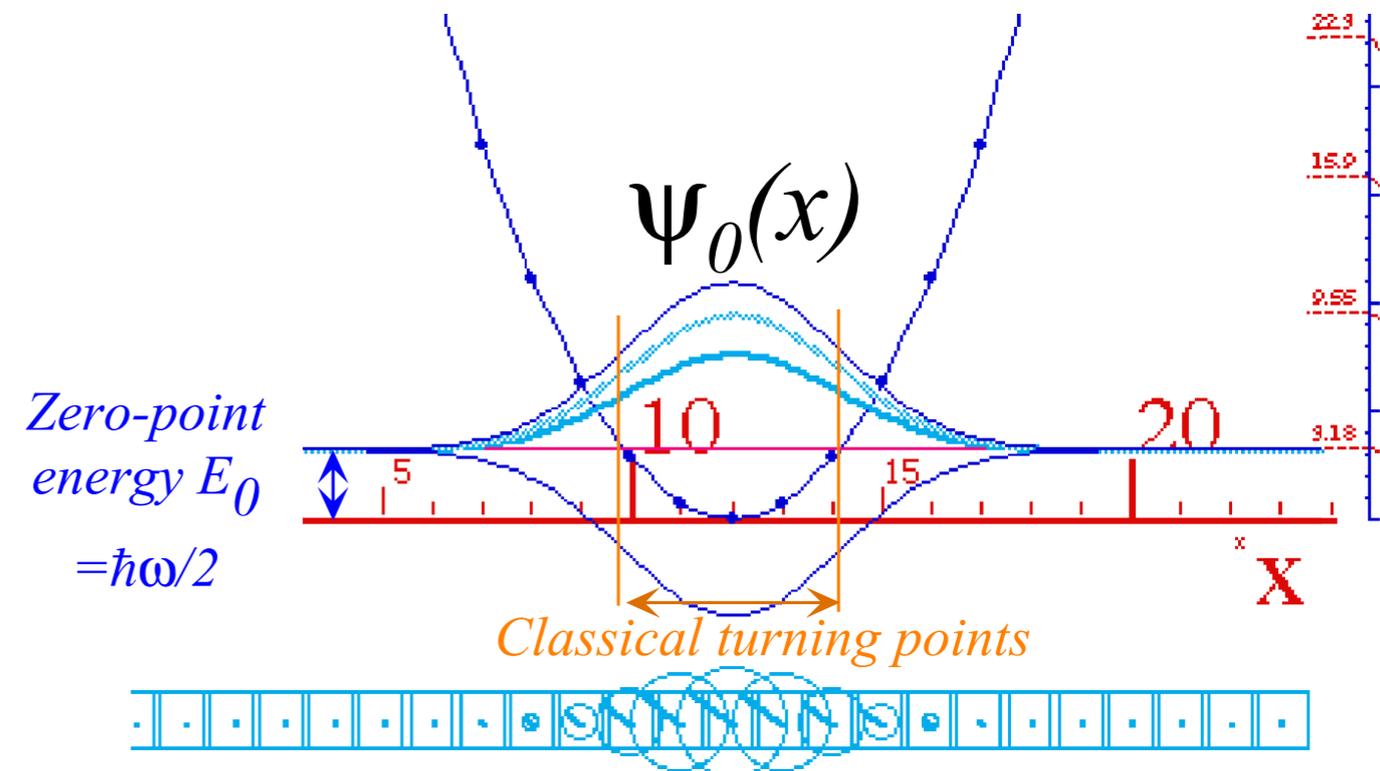
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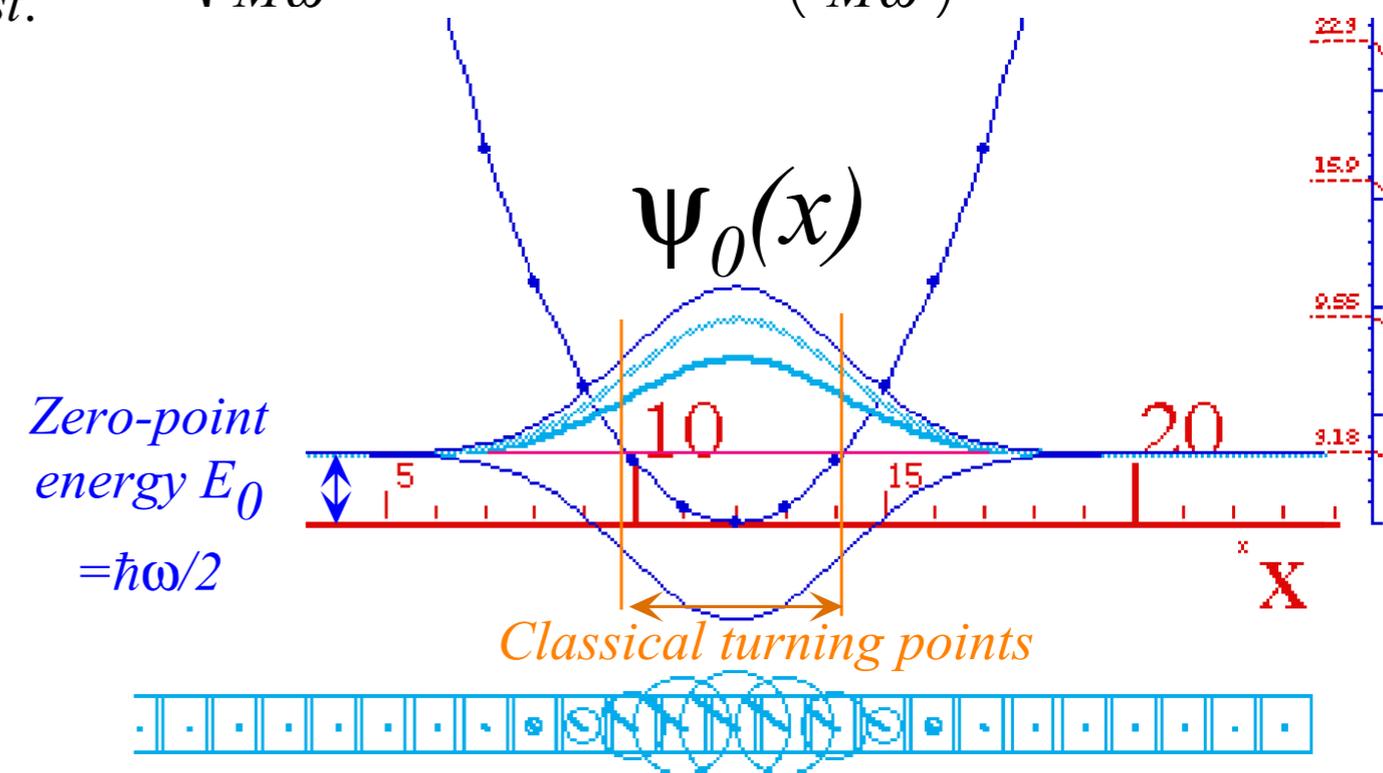
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The normalization *const.* is evaluated using a standard Gaussian integral: $\int_{-\infty}^{\infty} dx e^{-\alpha x^2} = \sqrt{\frac{\pi}{\alpha}}$

$$\langle \psi_0|\psi_0\rangle = 1 = \int_{-\infty}^{\infty} dx \frac{e^{-M\omega x^2/2\hbar}}{const.^2} = \sqrt{\frac{\pi\hbar}{M\omega}} / const.^2 \Rightarrow const. = \left(\frac{\pi\hbar}{M\omega}\right)^{1/4}$$



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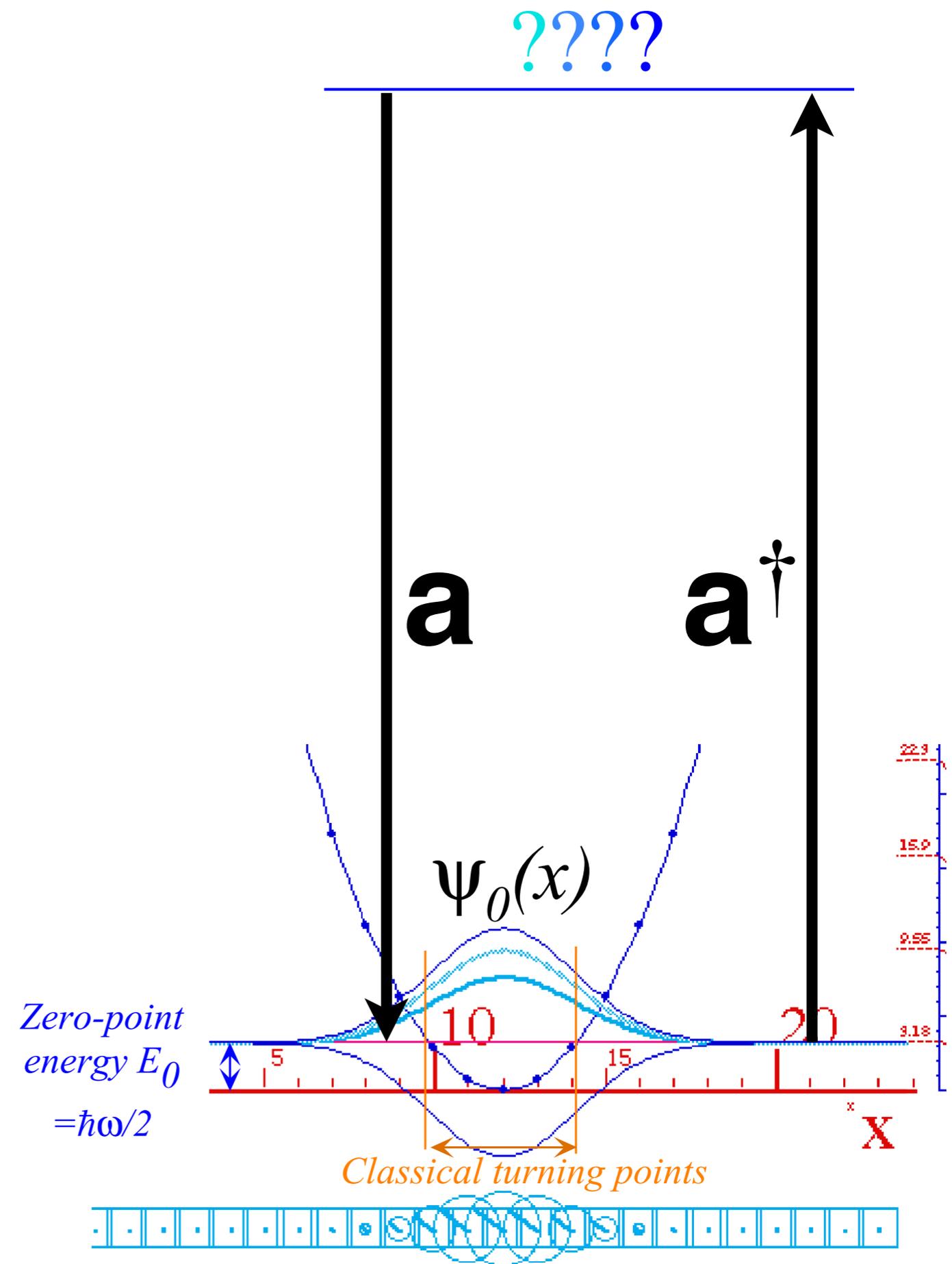
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1st excited state wavefunction $\psi_1(x) = \langle x | 1 \rangle$
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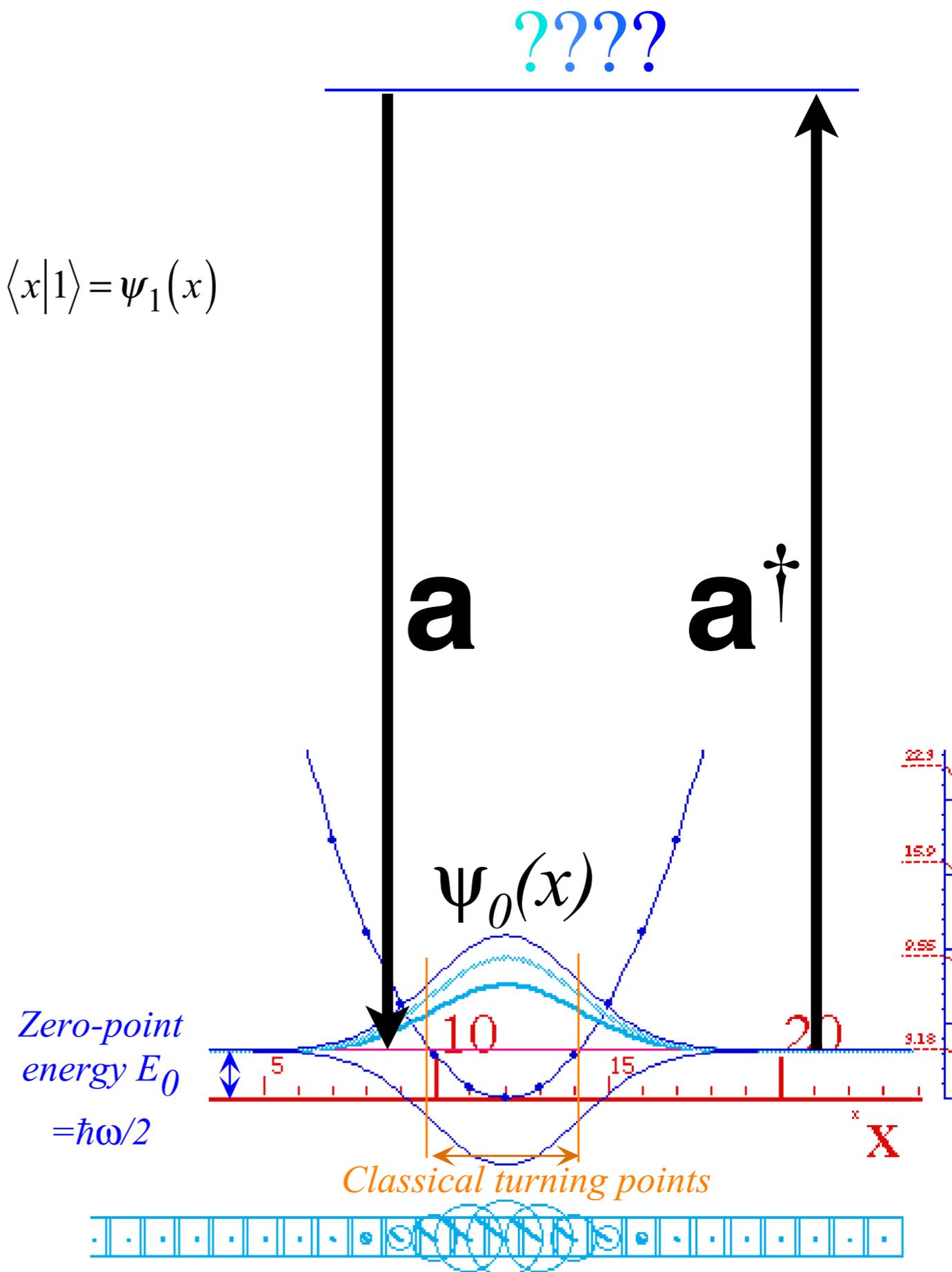


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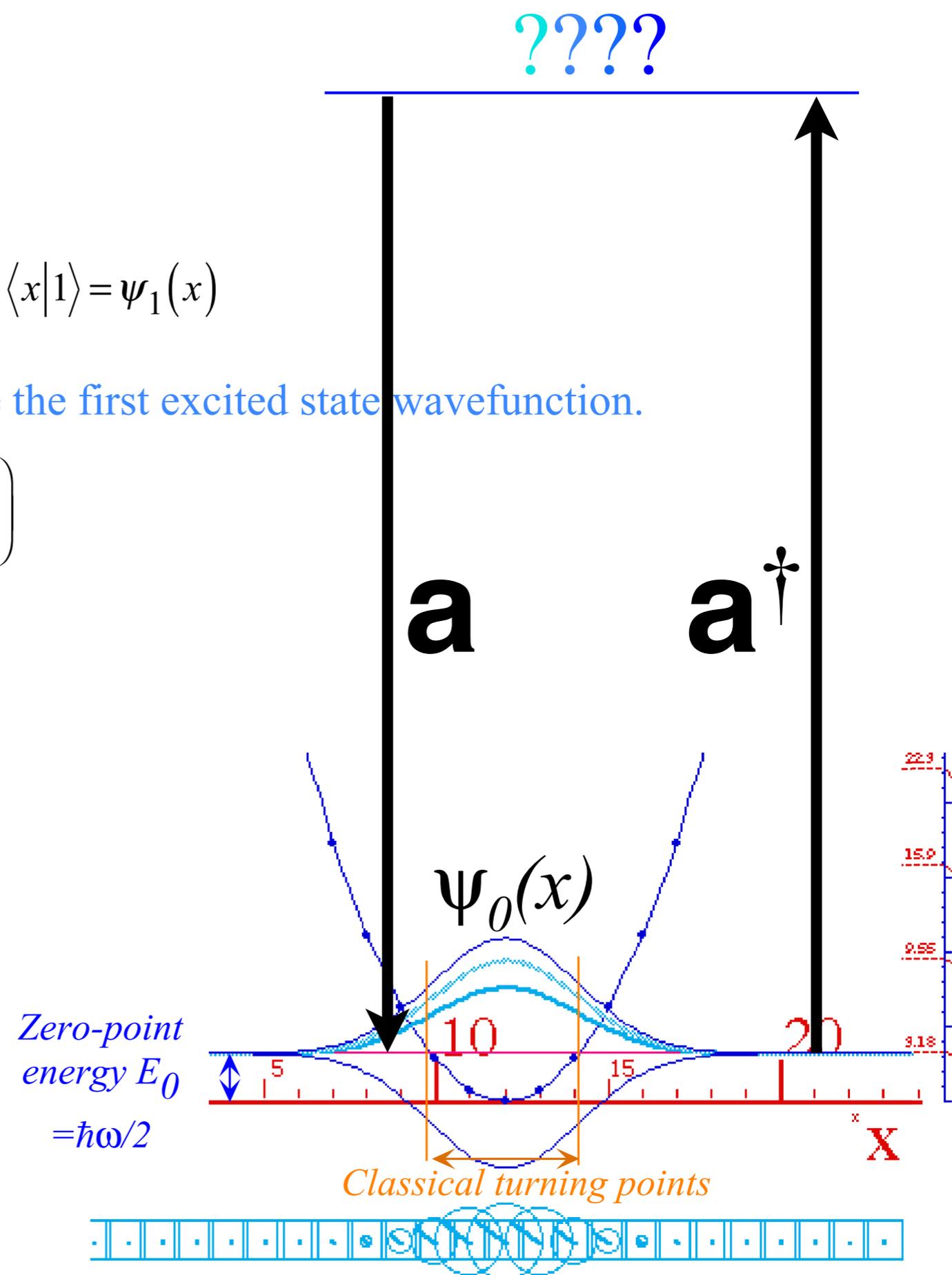
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The operator coordinate representations generate the first excited state wavefunction.

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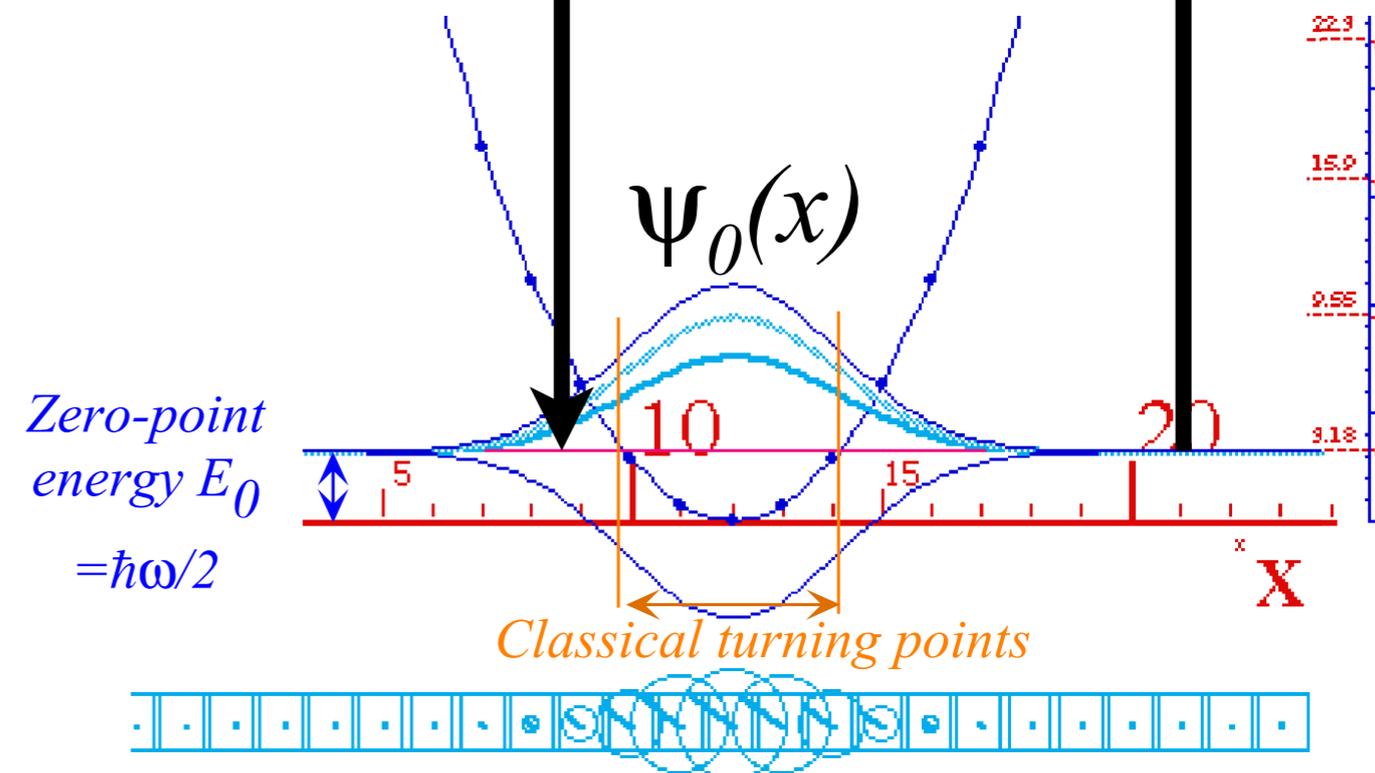
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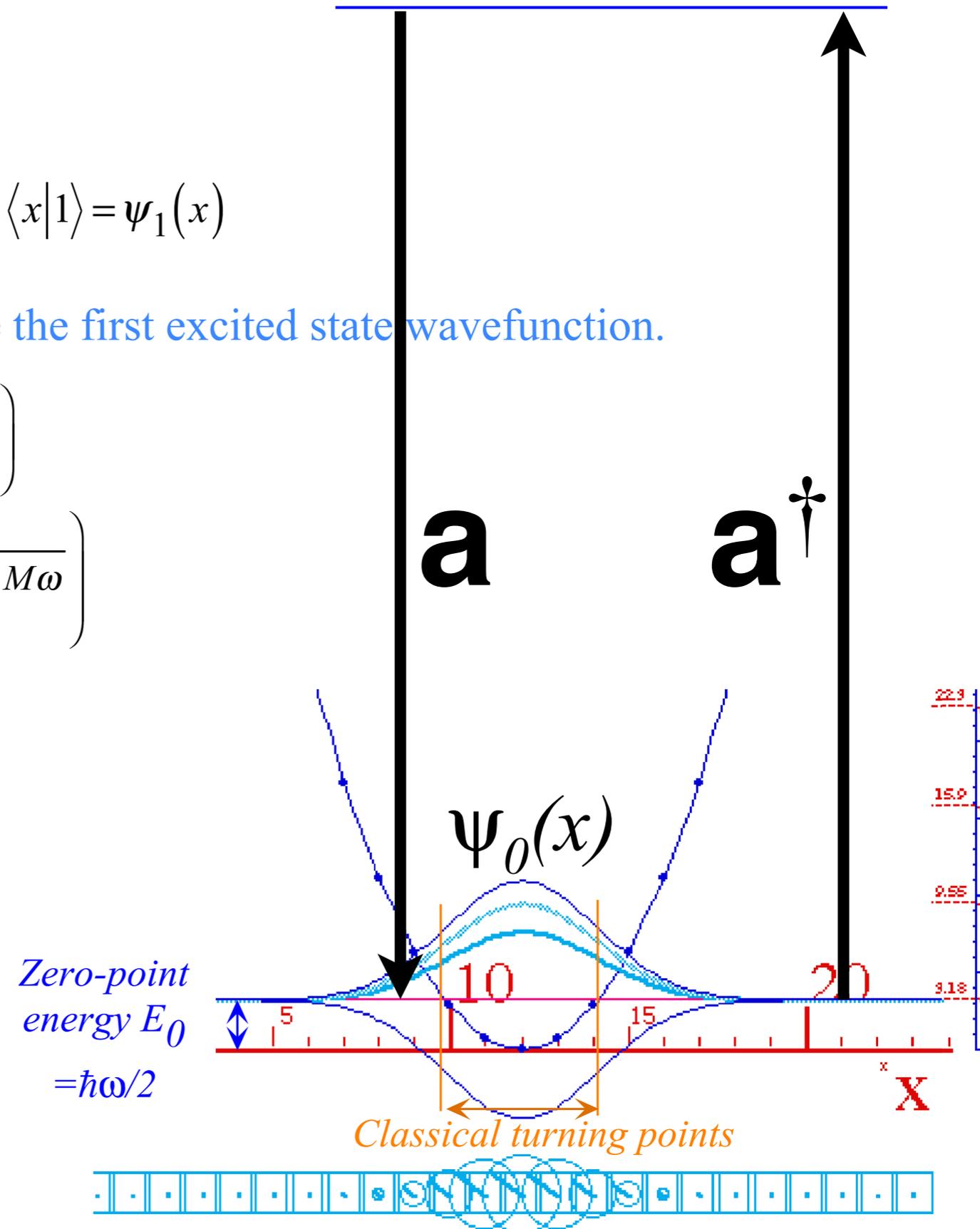
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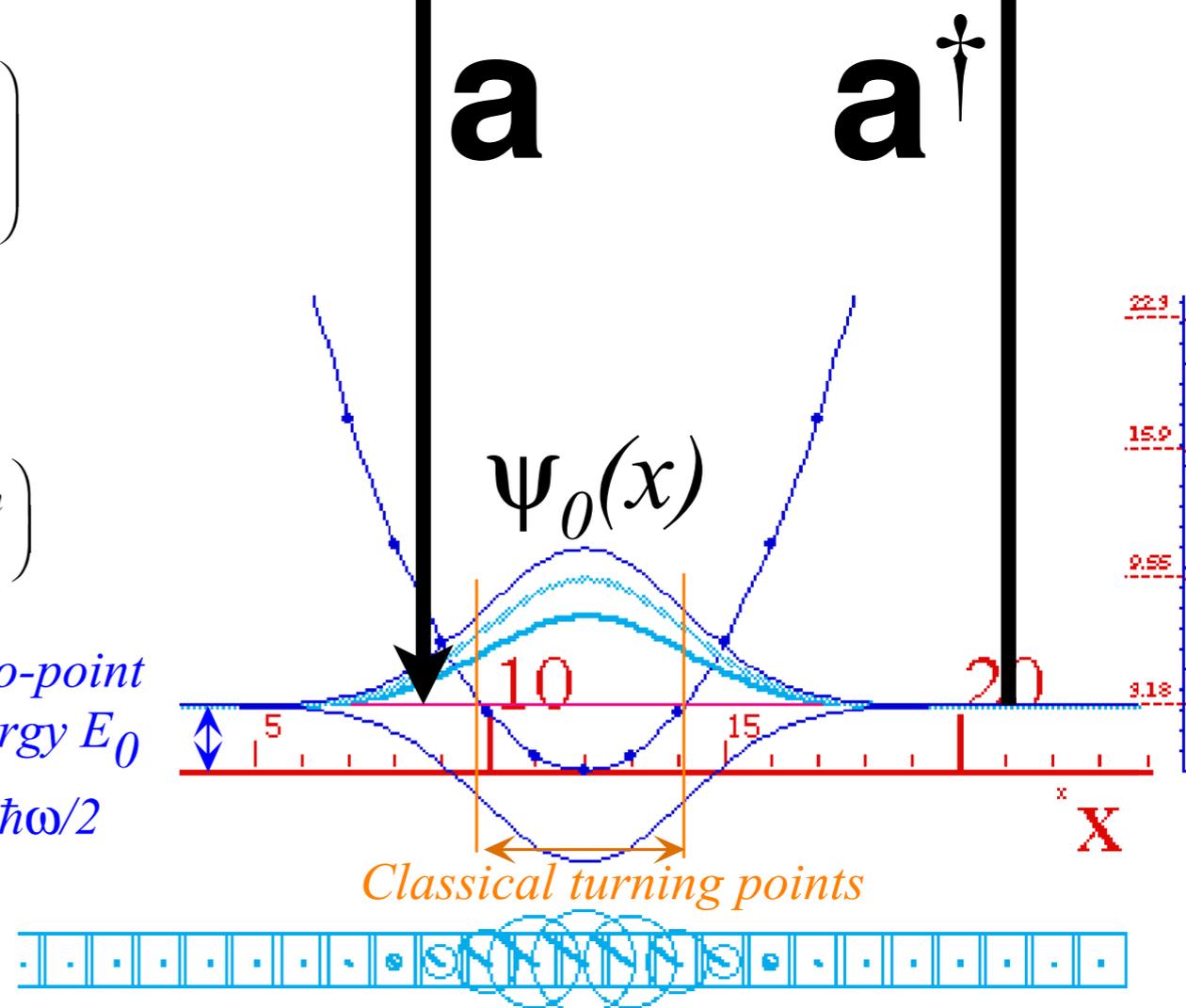
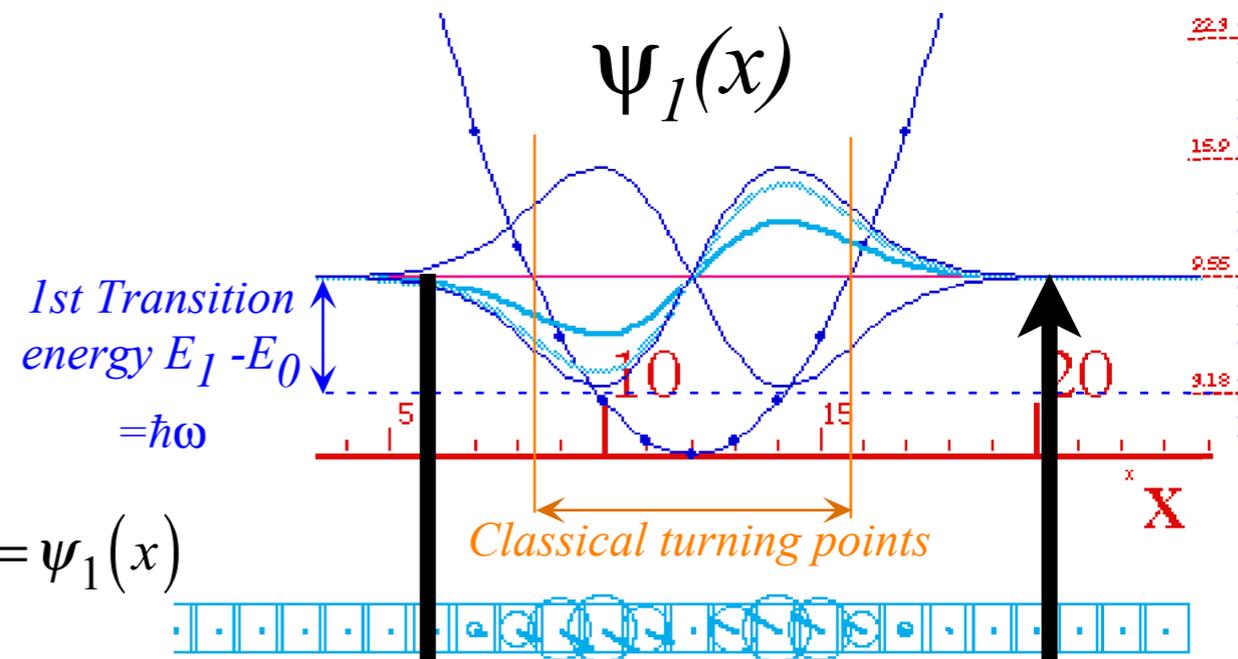
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→ *Normal ordering for matrix calculation* **←**

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Normal ordering: move destructive **a** operators to the right of creation **a**[†] to zero out on vacuum $|0\rangle$.

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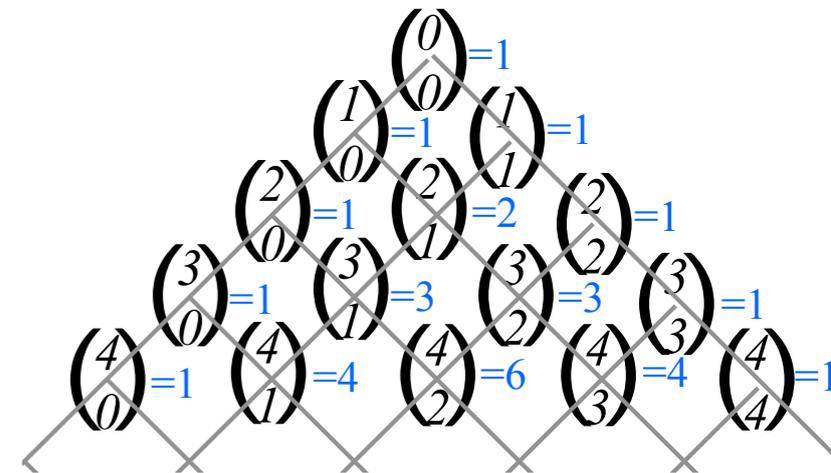
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1-D $\mathbf{a}^\dagger\mathbf{a}$ algebra of $U(1)$ representations

Creation-Destruction $\mathbf{a}^\dagger\mathbf{a}$ algebra

Eigenstate creationism (and destruction)

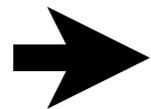
Vacuum state

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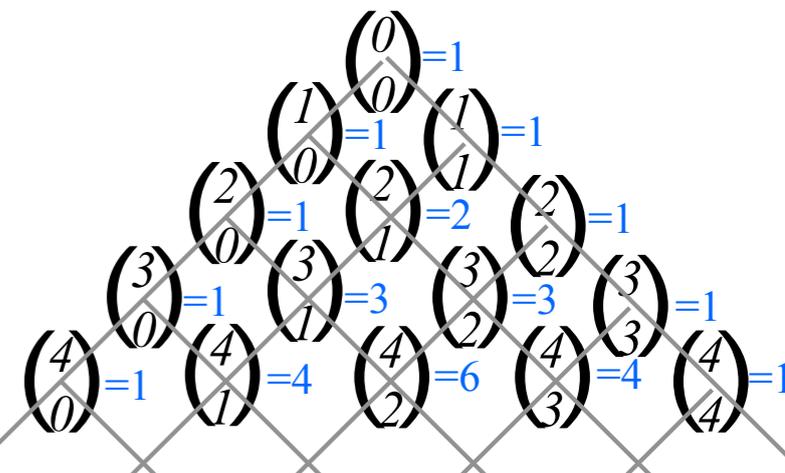
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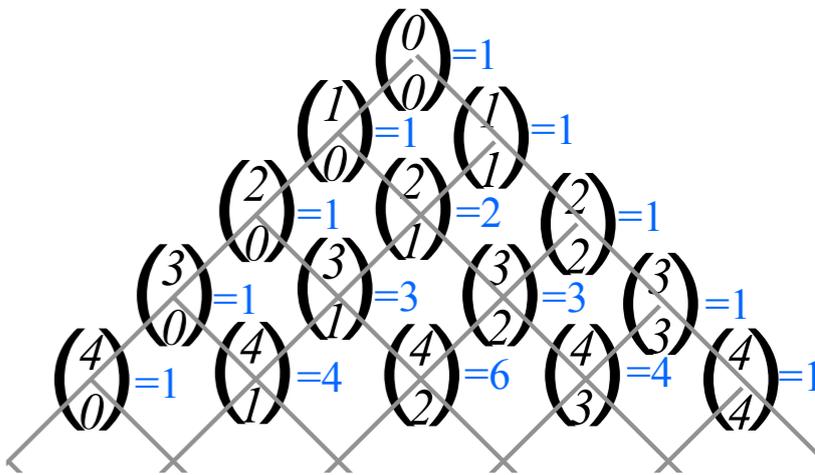
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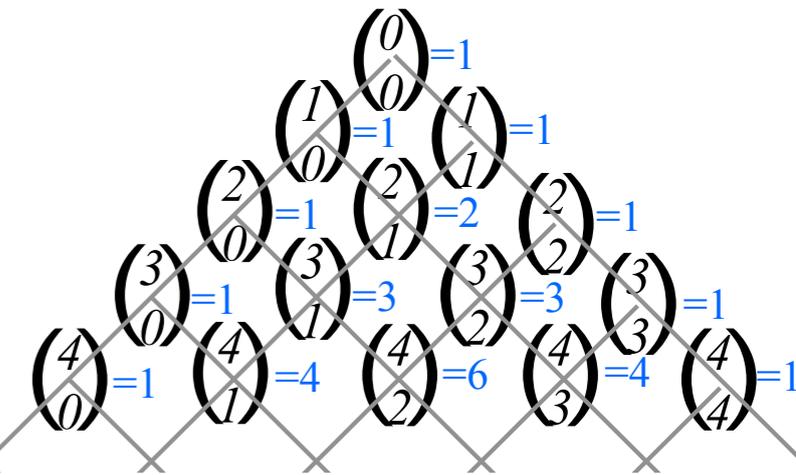
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So: $(\text{const.})^2 = n!$
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(Welcome to ∞ -dimensional... quantum space!)

1-D $\mathbf{a}^\dagger\mathbf{a}$ algebra of $U(1)$ representations

Creation-Destruction $\mathbf{a}^\dagger\mathbf{a}$ algebra

Eigenstate creationism (and destruction)

Vacuum state

1st excited state

Normal ordering for matrix calculation

Commutator derivative identities

Binomial expansion identities

Matrix $\langle \mathbf{a}^n \mathbf{a}^{\dagger n} \rangle$ calculations



Number operator and Hamiltonian operator



Expectation values of position, momentum, and uncertainty for eigenstate $|n\rangle$

Harmonic oscillator beat dynamics of mixed states

Oscillator coherent states (“Shoved” and “kicked” states)

Translation operators vs. boost operators

Applying boost-translation combinations

Time evolution of coherent state

Properties of coherent state and “squeezed” states

2-D $\mathbf{a}^\dagger\mathbf{a}$ algebra of $U(2)$ representations and $R(3)$ angular momentum operators

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Number operator $\mathbf{N} = \mathbf{a}^\dagger \mathbf{a}$ counts quanta.

$$\mathbf{a}^\dagger \mathbf{a} |n\rangle = \frac{\mathbf{a}^\dagger \mathbf{a} \mathbf{a}^{\dagger n} |0\rangle}{\sqrt{n!}} = n \frac{\mathbf{a}^\dagger \mathbf{a}^{\dagger n-1} |0\rangle}{\sqrt{n!}}$$

Matrix $\langle \mathbf{a}^n \mathbf{a}^{\dagger n} \rangle$ calculation

Derive normalization for n^{th} state obtained by $(\mathbf{a}^{\dagger})^n$ operator: Use: $\mathbf{a}^n \mathbf{a}^{\dagger n} = n! \left(\mathbf{1} + n \mathbf{a}^{\dagger} \mathbf{a} + \frac{n(n-1)}{2! \cdot 2!} \mathbf{a}^{\dagger 2} \mathbf{a}^2 + \dots \right)$

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Use: $\mathbf{a} \mathbf{a}^{\dagger n} = n \mathbf{a}^{\dagger n-1} + \mathbf{a}^{\dagger n} \mathbf{a}$

Apply creation \mathbf{a}^{\dagger} :

$$\mathbf{a}^{\dagger} |n\rangle = \frac{\mathbf{a}^{\dagger n+1} |0\rangle}{\sqrt{n!}} = \sqrt{n+1} \frac{\mathbf{a}^{\dagger n+1} |0\rangle}{\sqrt{(n+1)!}}$$

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$$\mathbf{a} |n\rangle = \frac{\mathbf{a} \mathbf{a}^{\dagger n} |0\rangle}{\sqrt{n!}} = \frac{(n \mathbf{a}^{\dagger n-1} + \mathbf{a}^{\dagger n} \mathbf{a}) |0\rangle}{\sqrt{n!}} = \sqrt{n} \frac{\mathbf{a}^{\dagger n-1} |0\rangle}{\sqrt{(n-1)!}}$$

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Feynman's mnemonic rule: Larger of two quanta goes in radical factor

$$\langle \mathbf{a}^{\dagger} \rangle = \begin{pmatrix} \cdot & & & & \\ 1 & \cdot & & & \\ & \sqrt{2} & \cdot & & \\ & & \sqrt{3} & \cdot & \\ & & & \sqrt{4} & \cdot \\ & & & & \ddots & \ddots \end{pmatrix}$$

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Hamiltonian operator

$$\mathbf{H} |n\rangle = \hbar\omega \mathbf{a}^{\dagger} \mathbf{a} |n\rangle + \hbar\omega/2 \mathbf{1} |n\rangle = \hbar\omega(n+1/2) |n\rangle$$

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Hamiltonian operator

$$\mathbf{H} |n\rangle = \hbar\omega \mathbf{a}^{\dagger} \mathbf{a} |n\rangle + \hbar\omega/2 \mathbf{1} |n\rangle = \hbar\omega(n+1/2) |n\rangle$$

$$\langle \mathbf{H} \rangle = \hbar\omega \langle \mathbf{a}^{\dagger} \mathbf{a} + \frac{1}{2} \mathbf{1} \rangle = \hbar\omega \begin{pmatrix} 0 & & & & \\ & 1 & & & \\ & & 2 & & \\ & & & 3 & \\ & & & & \ddots \end{pmatrix} + \hbar\omega \begin{pmatrix} 1/2 & & & & \\ & 1/2 & & & \\ & & 1/2 & & \\ & & & 1/2 & \\ & & & & \ddots \end{pmatrix}$$

Hamiltonian operator is $\hbar\omega \mathbf{N}$ plus zero-point energy $\mathbf{1} \hbar\omega/2$.

1-D $\mathbf{a}^\dagger\mathbf{a}$ algebra of $U(1)$ representations

Creation-Destruction $\mathbf{a}^\dagger\mathbf{a}$ algebra

Eigenstate creationism (and destruction)

Vacuum state

1st excited state

Normal ordering for matrix calculation

Commutator derivative identities

Binomial expansion identities

Matrix $\langle \mathbf{a}^n \mathbf{a}^{\dagger n} \rangle$ calculations

Number operator and Hamiltonian operator

 *Expectation values of position, momentum, and uncertainty for eigenstate $|n\rangle$* 

Harmonic oscillator beat dynamics of mixed states

Oscillator coherent states (“Shoved” and “kicked” states)

Translation operators vs. boost operators

Applying boost-translation combinations

Time evolution of coherent state

Properties of coherent state and “squeezed” states

2-D $\mathbf{a}^\dagger\mathbf{a}$ algebra of $U(2)$ representations and $R(3)$ angular momentum operators

Expectation values of position, momentum, and uncertainty for eigenstate $|n\rangle$

Operator for position \mathbf{x} : $\sqrt{\frac{M\omega}{2\hbar}}\mathbf{x} = \frac{\mathbf{a} + \mathbf{a}^\dagger}{2}$

Operator for momentum \mathbf{p} : $\sqrt{\frac{1}{2\hbar M\omega}}\mathbf{p} = \frac{\mathbf{a} - \mathbf{a}^\dagger}{2i}$

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Operator for position \mathbf{x} : $\sqrt{\frac{M\omega}{2\hbar}}\mathbf{x} = \frac{\mathbf{a} + \mathbf{a}^\dagger}{2}$

expectation for position $\langle \mathbf{x} \rangle$:

$$\bar{\mathbf{x}}|_n = \langle n | \mathbf{x} | n \rangle = \sqrt{\frac{\hbar}{2M\omega}} \langle n | (\mathbf{a} + \mathbf{a}^\dagger) | n \rangle = 0$$

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$$\bar{\mathbf{p}}|_n = \langle n | \mathbf{p} | n \rangle = i\sqrt{\frac{\hbar M\omega}{2}} \langle n | (\mathbf{a}^\dagger - \mathbf{a}) | n \rangle = 0$$

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expectation for (position)² $\langle \mathbf{x}^2 \rangle$:

$$\overline{\mathbf{x}^2}|_n = \langle n|\mathbf{x}^2|n\rangle = \frac{\hbar}{2M\omega} \langle n|(\mathbf{a} + \mathbf{a}^\dagger)^2|n\rangle$$

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expectation for (momentum)² $\langle \mathbf{p}^2 \rangle$:

$$\overline{\mathbf{p}^2}|_n = \langle n|\mathbf{p}^2|n\rangle = i^2 \frac{\hbar M\omega}{2} \langle n|(\mathbf{a}^\dagger - \mathbf{a})^2|n\rangle$$

Expectation values of position, momentum, and uncertainty for eigenstate $|n\rangle$

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Use:
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Uncertainty or standard deviation Δq of a statistical quantity q is its root mean-square difference.

$$(\Delta q)^2 = \overline{(q - \bar{q})^2} \quad \text{or:} \quad \Delta q = \sqrt{\overline{(q - \bar{q})^2}}$$

Expectation values of position, momentum, and uncertainty for eigenstate $|n\rangle$

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Heisenberg uncertainty product for the n -quantum eigenstate $|n\rangle$

$$(\Delta x \cdot \Delta p)|_n = \sqrt{\overline{\mathbf{x}^2}} \sqrt{\overline{\mathbf{p}^2}} = \sqrt{\frac{\hbar(2n+1)}{2M\omega}} \sqrt{\frac{\hbar M\omega(2n+1)}{2}}$$

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expectation for (position)² $\langle \mathbf{x}^2 \rangle$:

$$\overline{\mathbf{x}^2}|_n = \langle n | \mathbf{x}^2 | n \rangle = \frac{\hbar}{2M\omega} \langle n | (\mathbf{a} + \mathbf{a}^\dagger)^2 | n \rangle$$

$$= \frac{\hbar}{2M\omega} \langle n | (\mathbf{a}^2 + \mathbf{a}^\dagger \mathbf{a} + \mathbf{a} \mathbf{a}^\dagger + \mathbf{a}^{\dagger 2}) | n \rangle$$

$$= \frac{\hbar}{2M\omega} (2n+1)$$

Use:
 $\mathbf{a} \mathbf{a}^\dagger = \mathbf{1} + \mathbf{a}^\dagger \mathbf{a}$

Operator for momentum \mathbf{p} : $\sqrt{\frac{1}{2\hbar M\omega}}\mathbf{p} = \frac{\mathbf{a} - \mathbf{a}^\dagger}{2i}$

expectation for momentum $\langle \mathbf{p} \rangle$:

$$\bar{\mathbf{p}}|_n = \langle n | \mathbf{p} | n \rangle = i \sqrt{\frac{\hbar M\omega}{2}} \langle n | (\mathbf{a}^\dagger - \mathbf{a}) | n \rangle = 0$$

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$$\overline{\mathbf{p}^2}|_n = \langle n | \mathbf{p}^2 | n \rangle = i^2 \frac{\hbar M\omega}{2} \langle n | (\mathbf{a}^\dagger - \mathbf{a})^2 | n \rangle$$

$$= -\frac{\hbar M\omega}{2} \langle n | (\mathbf{a}^{\dagger 2} - \mathbf{a}^\dagger \mathbf{a} - \mathbf{a} \mathbf{a}^\dagger + \mathbf{a}^2) | n \rangle$$

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Uncertainty or standard deviation Δq of a statistical quantity q is its root mean-square difference.

$$\Delta x|_n = \sqrt{\overline{\mathbf{x}^2}} = \sqrt{\frac{\hbar(2n+1)}{2M\omega}} \quad (\Delta q)^2 = \overline{(q - \bar{q})^2} \quad \text{or:} \quad \Delta q = \sqrt{\overline{(q - \bar{q})^2}} \quad \Delta p|_n = \sqrt{\overline{\mathbf{p}^2}} = \sqrt{\frac{\hbar M\omega(2n+1)}{2}}$$

Heisenberg uncertainty product for the n -quantum eigenstate $|n\rangle$

$$(\Delta x \cdot \Delta p)|_n = \sqrt{\overline{\mathbf{x}^2}} \sqrt{\overline{\mathbf{p}^2}} = \sqrt{\frac{\hbar(2n+1)}{2M\omega}} \sqrt{\frac{\hbar M\omega(2n+1)}{2}}$$

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Expectation values of position, momentum, and uncertainty for eigenstate $|n\rangle$

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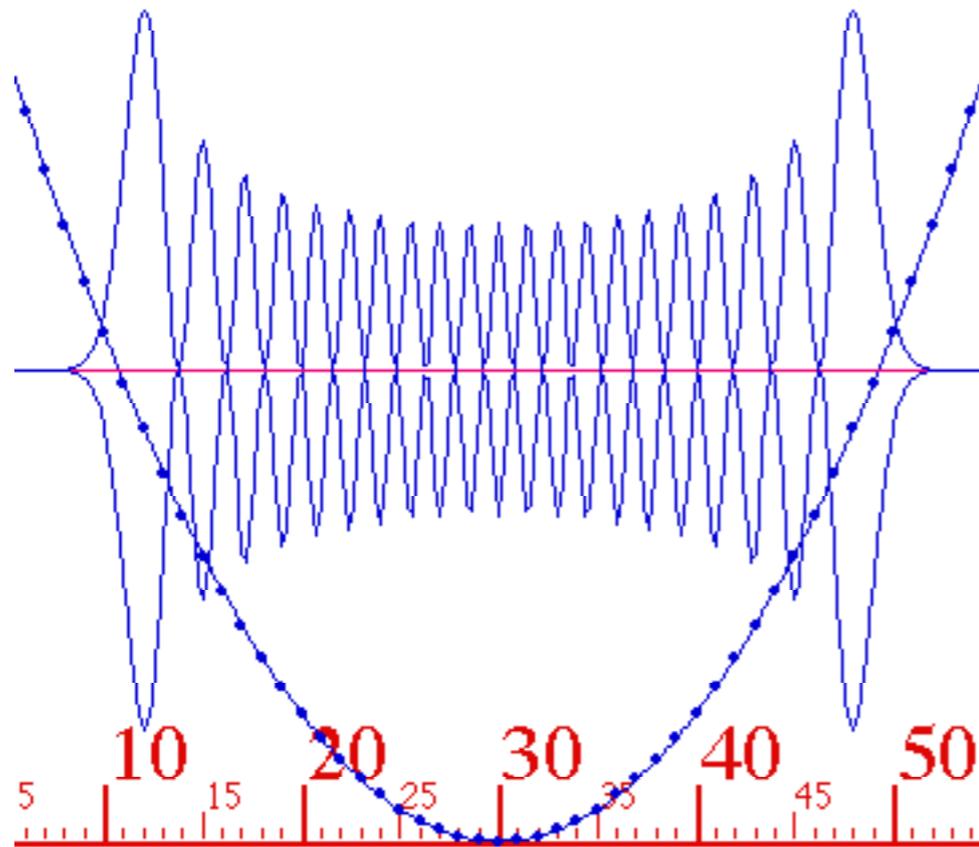
$$(\Delta x \cdot \Delta p)|_n = \sqrt{\overline{\mathbf{x}^2}} \sqrt{\overline{\mathbf{p}^2}} = \sqrt{\frac{\hbar(2n+1)}{2M\omega}} \sqrt{\frac{\hbar M\omega(2n+1)}{2}}$$

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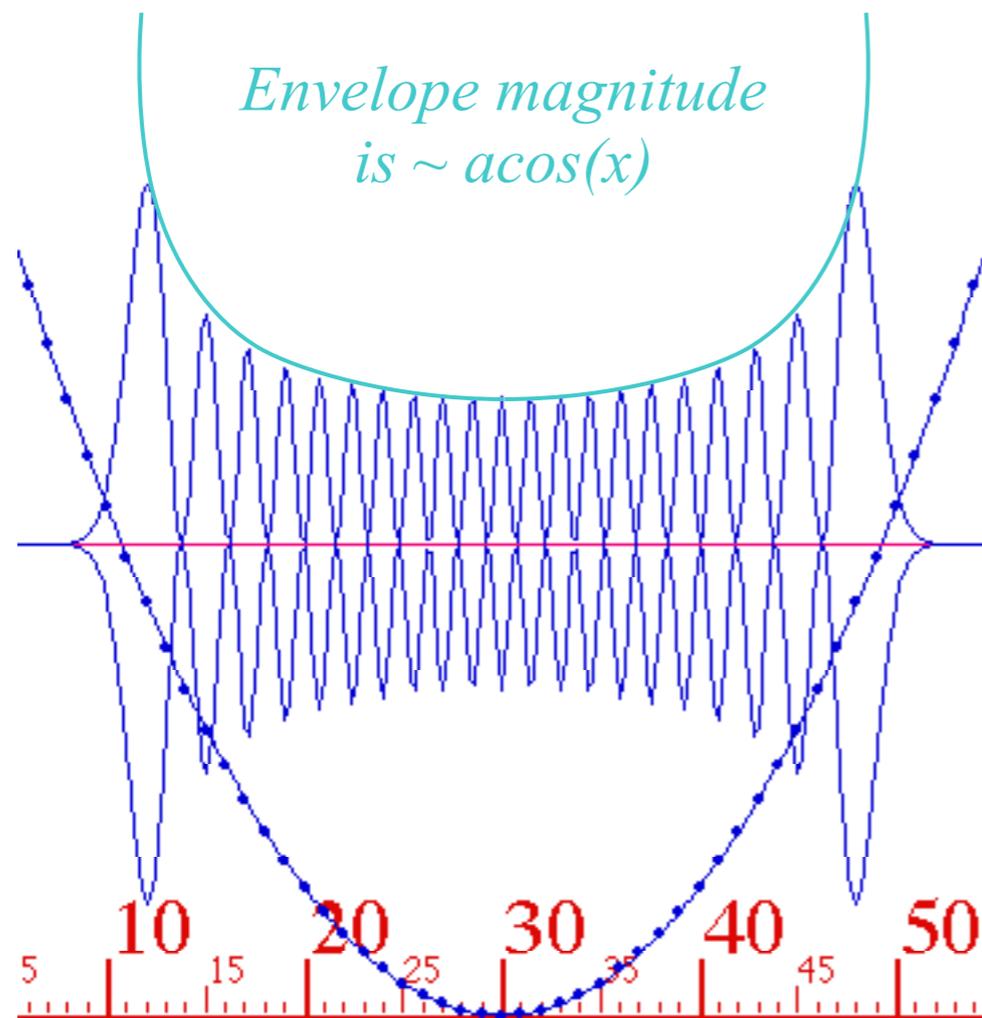
Heisenberg minimum uncertainty product occurs for the 0-quantum (ground) eigenstate.

$$(\Delta x \cdot \Delta p)|_0 = \frac{\hbar}{2}$$

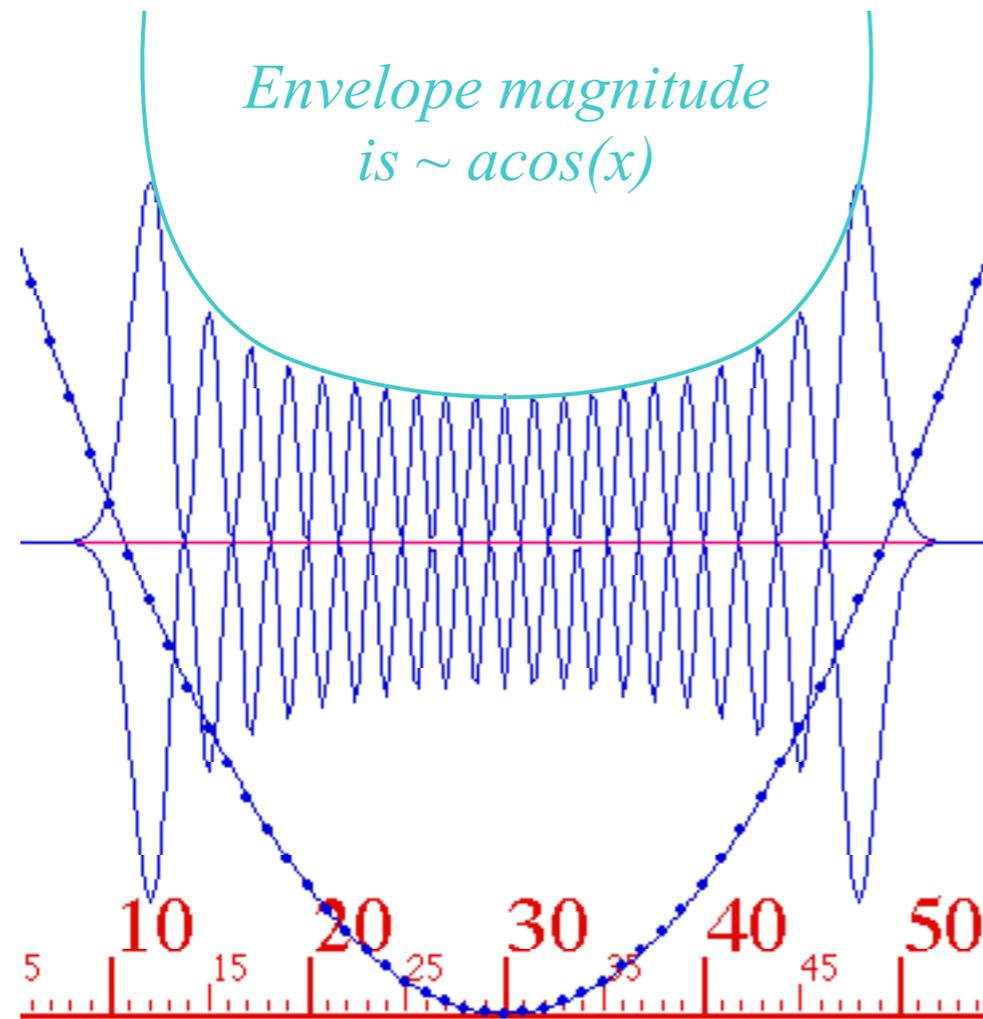
We pause for sobering considerations of the quantum world *vs.* the classical one.
Consider a “high”-quantum ($n=20$) eigenstate wavefunction:



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Consider a “high”-quantum ($n=20$) eigenstate wavefunction:



$n=20$ wave is still a long way from a classical energy value of *1 Joule*.
For a *1 Hz* oscillator, *1 Joule* would take a quantum number of roughly
 $n = 100,000,000,000,000,000,000,000,000,000,000,000,000 = 10^{35}$

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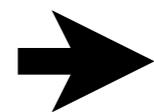
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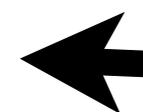
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Harmonic oscillator beat dynamics of mixed states



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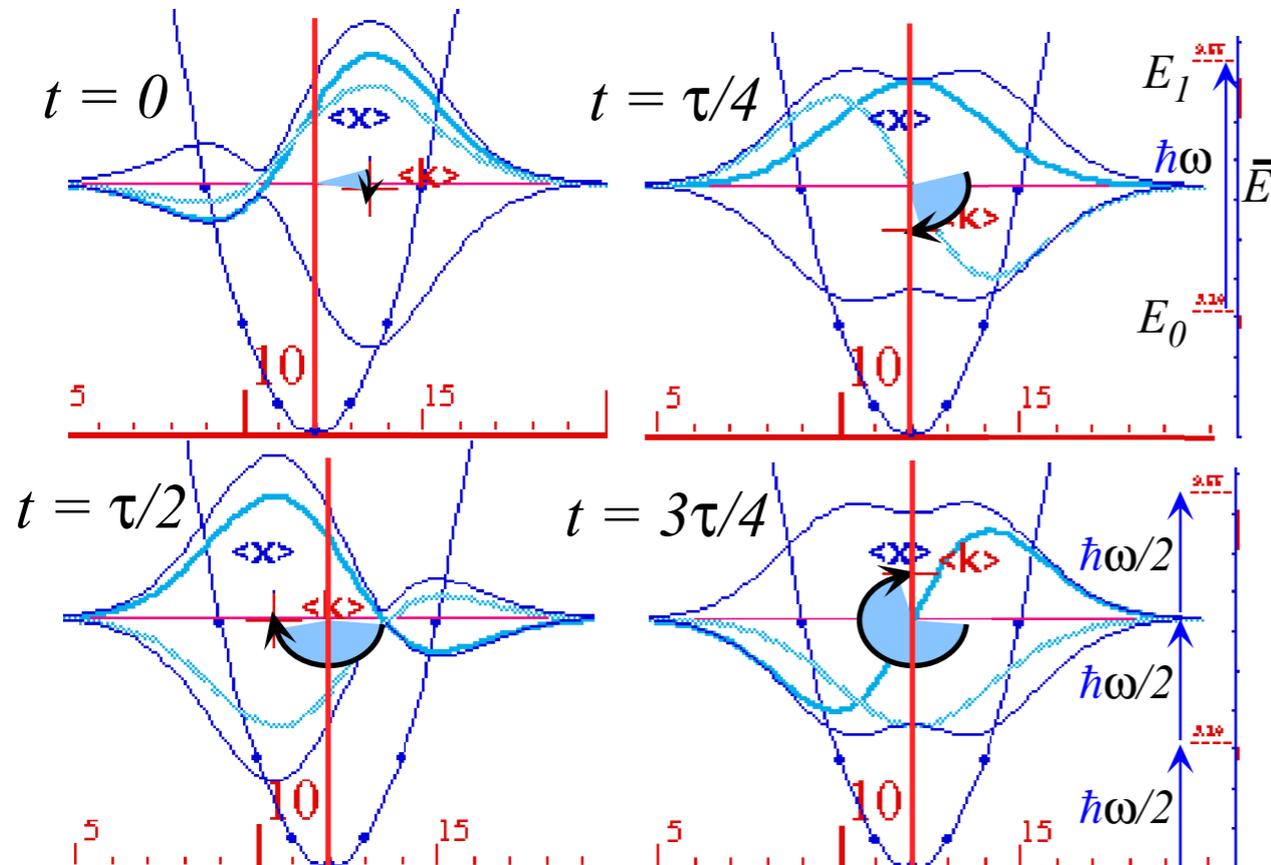
$$|\Psi\rangle = |0\rangle\langle 0|\Psi\rangle + |1\rangle\langle 1|\Psi\rangle = |0\rangle\Psi_0 + |1\rangle\Psi_1$$

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The time dependence $\Psi(x,t)$ of the mixed wave is then

$$\Psi(x,t) = \psi_0(x) e^{-i\omega_0 t} \Psi_0 + \psi_1(x) e^{-i\omega_1 t} \Psi_1 = (\psi_0(x) e^{-i\omega_0 t} + \psi_1(x) e^{-i\omega_1 t})/\sqrt{2}$$

$$\begin{aligned} |\Psi(x,t)| &= \sqrt{\Psi^* \Psi} = \sqrt{\left(e^{-i\omega_0 t} \psi_0(x) + e^{-i\omega_1 t} \psi_1(x) \right)^* \left(e^{-i\omega_0 t} \psi_0(x) + e^{-i\omega_1 t} \psi_1(x) \right) / 2} \\ &= \sqrt{\left(|\psi_0(x)|^2 + |\psi_1(x)|^2 + \psi_0(x)\psi_1(x) \left(e^{i(\omega_1-\omega_0)t} + e^{-i(\omega_1-\omega_0)t} \right) \right) / 2} \\ &= \sqrt{\left(|\psi_0(x)|^2 + |\psi_1(x)|^2 + 2\psi_0(x)\psi_1(x)\cos(\omega_1-\omega_0)t \right) / 2} \end{aligned}$$



Harmonic oscillator beat dynamics of mixed states

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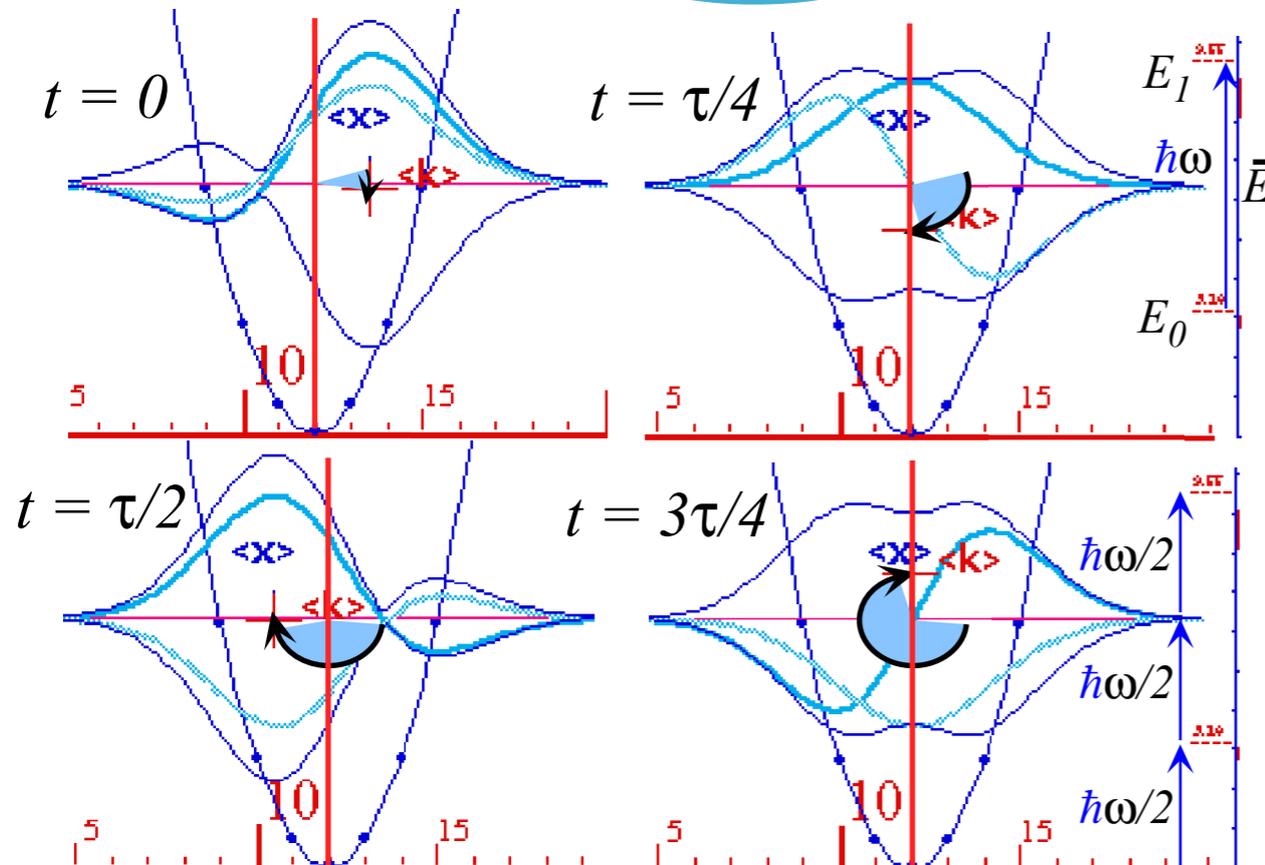
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Need some *overlap* somewhere to get some *wiggle*



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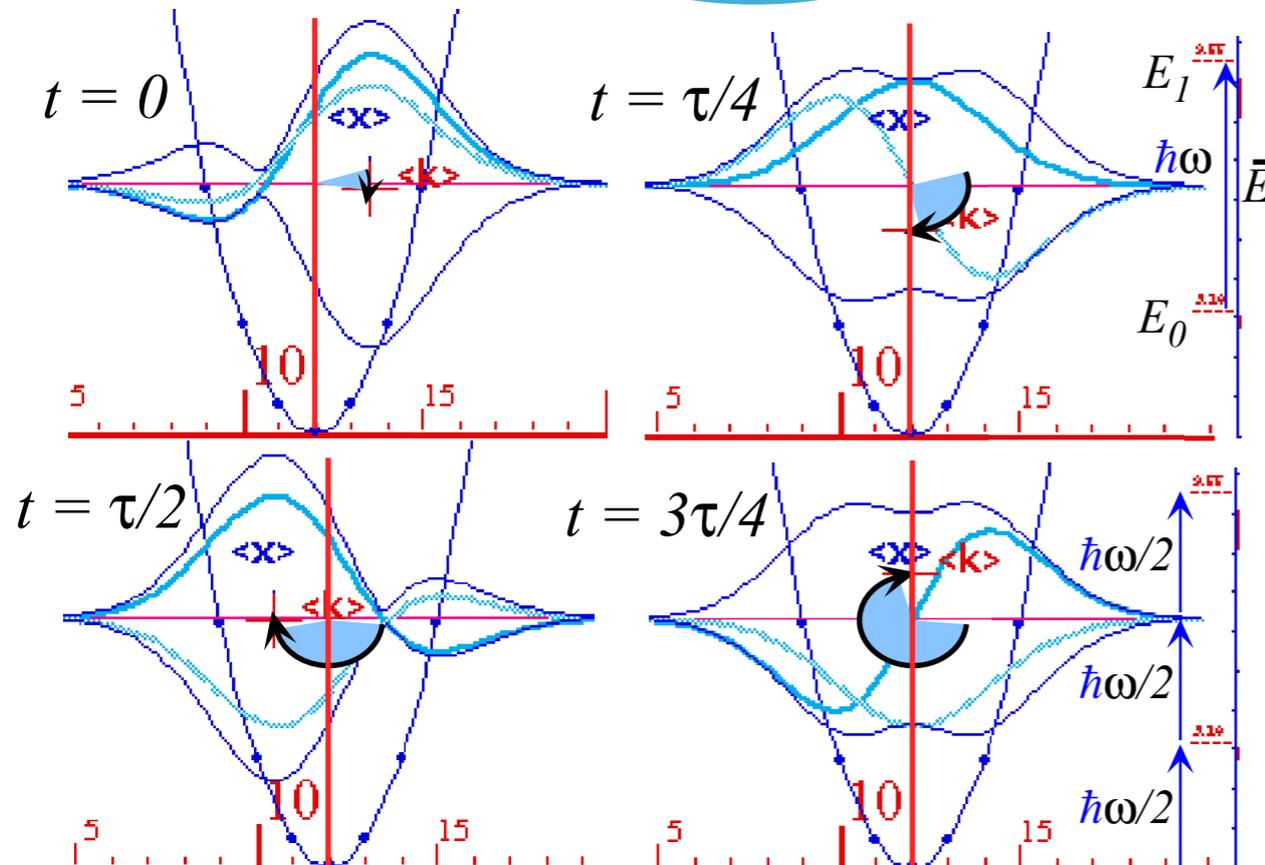
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$$\omega_{beat} = \omega_1 - \omega_0 = \omega$$

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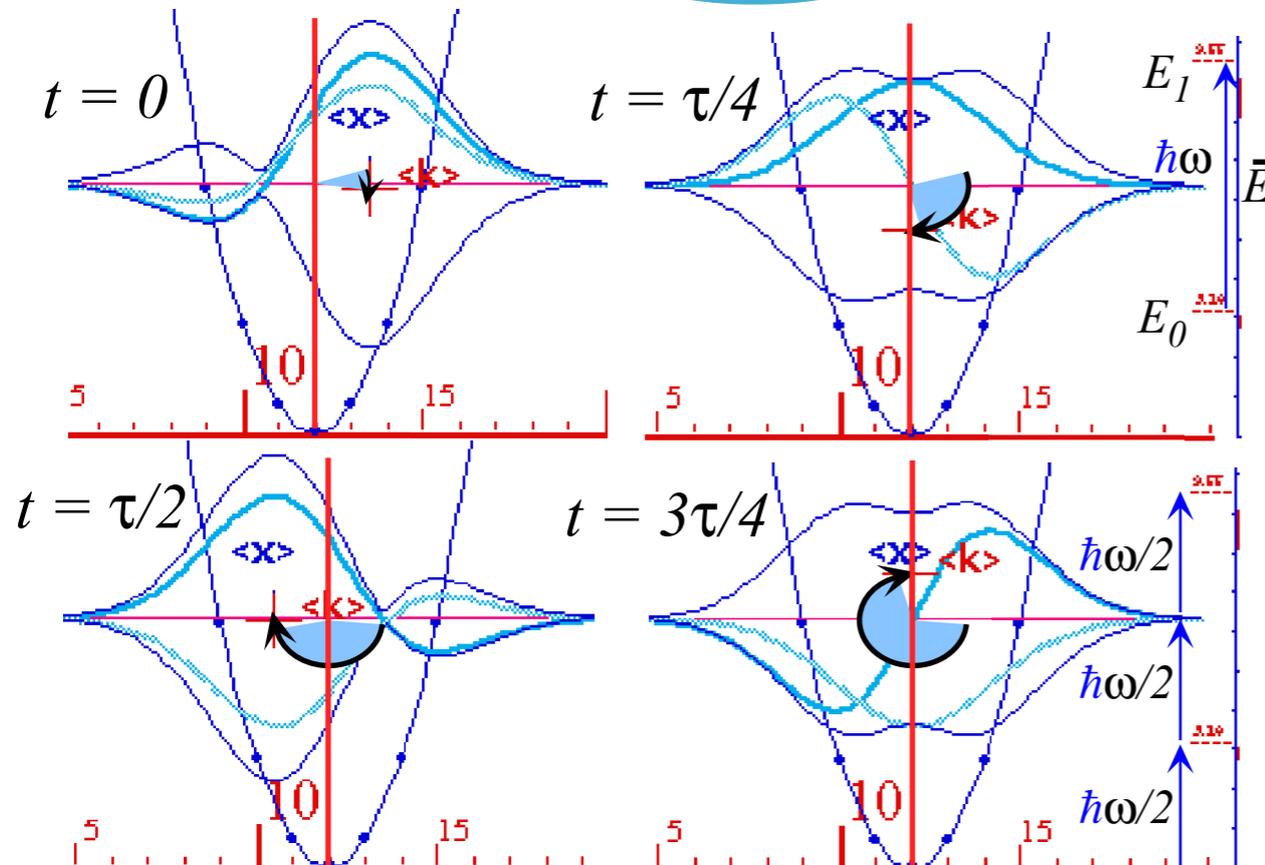
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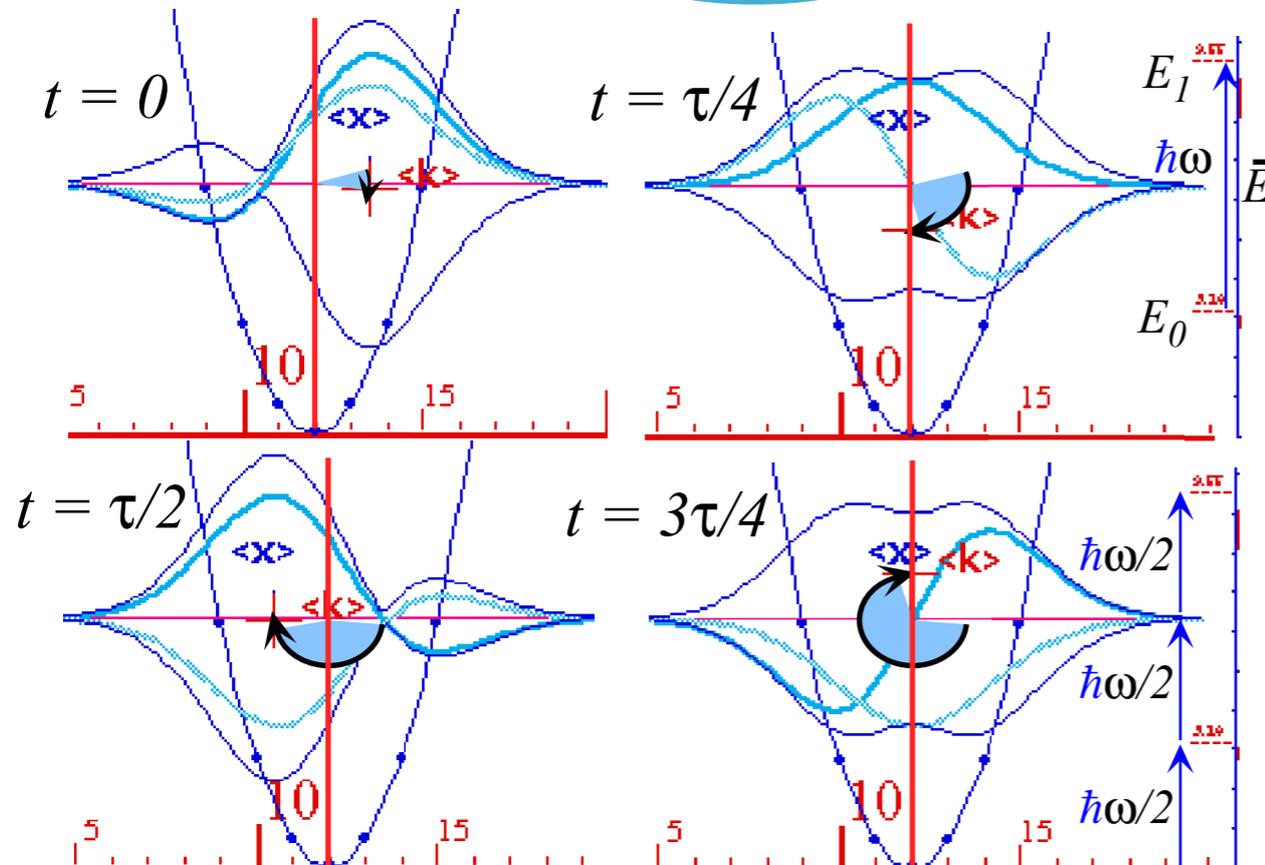
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Beat frequency $\omega =$ Transition frequency ω

Transition frequency is transition energy/ \hbar

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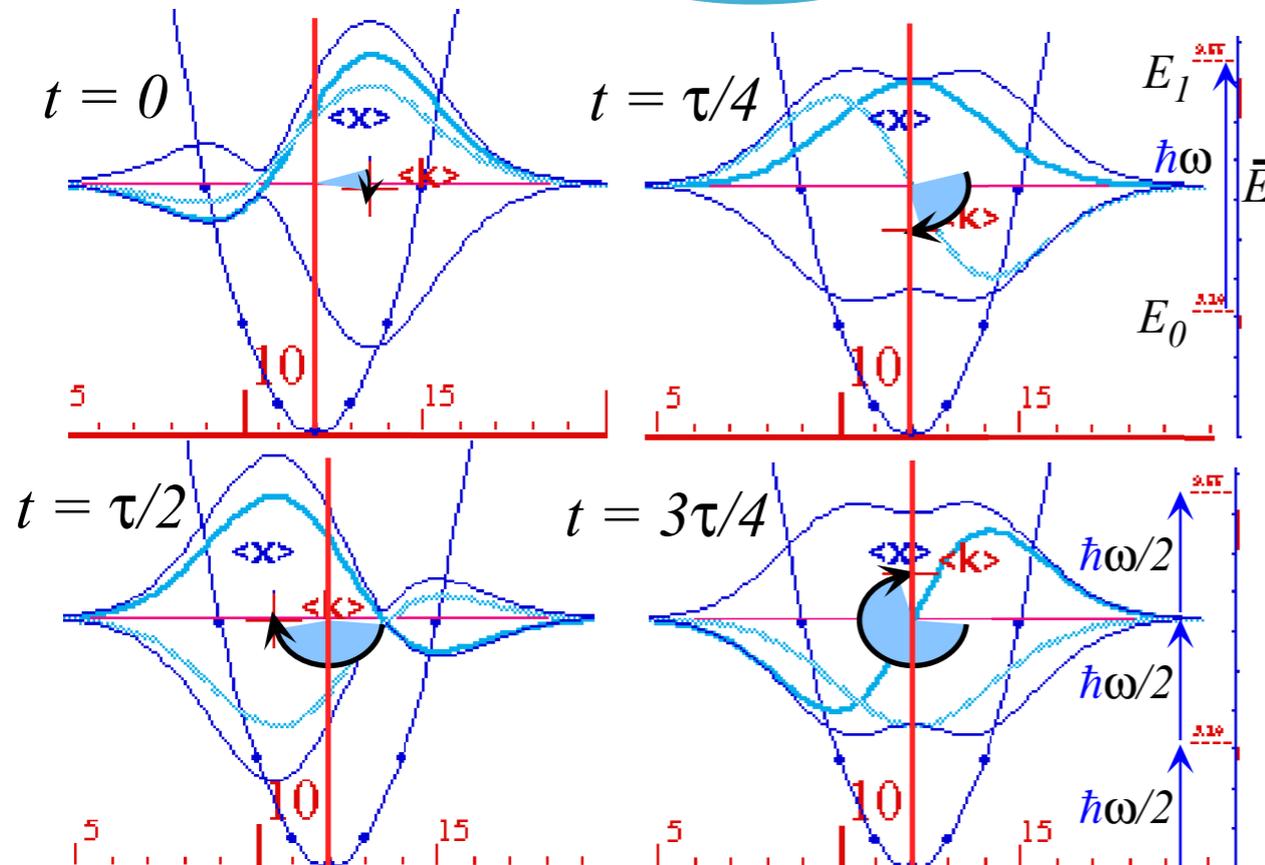
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ω is frequency of radiating antenna of a transmitter or of a receiver, i.e., of an emitter or an absorber (Usually of a dipole symmetry)

1-D $\mathbf{a}^\dagger\mathbf{a}$ algebra of $U(1)$ representations

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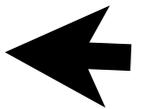
 *Oscillator coherent states (“Shoved” and “kicked” states)*

Translation operators vs. boost operators

Applying boost-translation combinations

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Translation operators and generators: (A “shove”)

Translation operator $\mathbf{T}(a)$ shoves x -wavefunctions

$$\mathbf{T}(a) \cdot \psi(x) = \psi(x-a) = \langle x | \mathbf{T}(a) | \psi \rangle = \langle x-a | \psi \rangle$$

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Increases momentum of ket-state by b units

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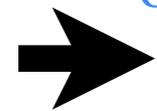
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Tiny translation $a \rightarrow da$ is identity $\mathbf{1}$ plus $\mathbf{G} \cdot da$

$$\mathbf{T}(da) = \mathbf{1} + \mathbf{G} \cdot da \quad \text{where:} \quad \mathbf{G} = \left. \frac{\partial \mathbf{T}}{\partial a} \right|_{a=0}$$

is *generator of translations*

Boost operators and generators: (A “kick”)

Boost operator $\mathbf{B}(b)$ boosts p -wavefunctions

$$\mathbf{B}(b) \cdot \psi(p) = \psi(p-b) = \langle p | \mathbf{B}(b) | \psi \rangle = \langle p-b | \psi \rangle$$

Increases momentum of ket-state by b units

$$\langle p | \mathbf{B}(b) = \langle p-b | \quad , \quad \text{or:} \quad \mathbf{B}^\dagger(b) | p \rangle = | p-b \rangle$$

Tiny boost $b \rightarrow db$ is identity $\mathbf{1}$ plus $\mathbf{K} \cdot db$

$$\mathbf{B}(db) = \mathbf{1} + \mathbf{K} \cdot db \quad \text{where:} \quad \mathbf{K} = \left. \frac{\partial \mathbf{B}}{\partial b} \right|_{b=0}$$

is *generator of boosts*

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$$\mathbf{T}(a) = \left(\mathbf{T}\left(\frac{a}{N}\right) \right)^N = \lim_{N \rightarrow \infty} \left(1 + \frac{a}{N} \mathbf{G} \right)^N = e^{a\mathbf{G}}$$

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is *generator G of translations*

is *generator K of boosts*

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$$\mathbf{G} \text{ relates to momentum } \mathbf{p} \rightarrow \frac{\hbar}{i} \frac{\partial}{\partial x} = -i\hbar \frac{\partial}{\partial x}$$

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Oscillator coherent states (“Shoved” and “kicked” states)

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$$\mathbf{T}(a) = e^{-a \frac{i}{\hbar} \mathbf{p}} = e^{a(\mathbf{a}^\dagger - \mathbf{a})\sqrt{M\omega/2\hbar}}$$

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$$\mathbf{B}(b) = e^{b \frac{i}{\hbar} \mathbf{x}} = e^{ib(\mathbf{a}^\dagger + \mathbf{a})/\sqrt{2\hbar M\omega}}$$

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Check $\mathbf{T}(a)$ on plane-wave with $p=\hbar k$ *Bottom Line*

$$\mathbf{T}(a)e^{ikx} = e^{-ia\mathbf{p}/\hbar} e^{ikx} = e^{-iak} e^{ikx} = e^{ik(x-a)}$$

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Check $\mathbf{B}(b)$ on plane-wave with $p=\hbar k$

$$\mathbf{B}(b)e^{ikx} = e^{ib\mathbf{x}/\hbar} e^{ikx} = e^{ibx/\hbar} e^{ikx} = e^{i(k+b/\hbar)x}$$

1-D $\mathbf{a}^\dagger\mathbf{a}$ algebra of $U(1)$ representations

Creation-Destruction $\mathbf{a}^\dagger\mathbf{a}$ algebra

Eigenstate creationism (and destruction)

Vacuum state

1st excited state

Normal ordering for matrix calculation

Commutator derivative identities

Binomial expansion identities

Matrix $\langle \mathbf{a}^n \mathbf{a}^{\dagger n} \rangle$ calculations

Number operator and Hamiltonian operator

Expectation values of position, momentum, and uncertainty for eigenstate $|n\rangle$

Harmonic oscillator beat dynamics of mixed states

Oscillator coherent states (“Shoved” and “kicked” states)

Translation operators vs. boost operators

 *Applying boost-translation combinations* 

Time evolution of coherent state

Properties of coherent state and “squeezed” states

2-D $\mathbf{a}^\dagger\mathbf{a}$ algebra of $U(2)$ representations and $R(3)$ angular momentum operators

Applying boost-translation combinations

T(*a*) and **B**(*b*) operations do not commute. Q. Which should come first?

??

Applying boost-translation combinations

T(*a*) and **B**(*b*) operations do not commute. Q. Which should come first? **T**(*a*) = $e^{-i a \mathbf{p} / \hbar}$ or **B**(*b*) = $e^{i b \mathbf{x} / \hbar}$??

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A. Neither and Both.

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(More like Darboux rotation $e^{-i \Theta \cdot \mathbf{J} / \hbar}$ than Euler rotation with three factors $e^{-i \mathbf{J}_z \alpha / \hbar} e^{-i \mathbf{J}_y \beta / \hbar} e^{-i \mathbf{J}_z \gamma / \hbar}$)

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May evaluate with *Baker-Campbell-Hausdorff identity* since $[\mathbf{x},\mathbf{p}] = i\hbar\mathbf{1}$ and $[[\mathbf{x},\mathbf{p}],\mathbf{x}] = [[\mathbf{x},\mathbf{p}],\mathbf{p}] = \mathbf{0}$.

$$e^{\mathbf{A}+\mathbf{B}} = e^{\mathbf{A}}e^{\mathbf{B}}e^{-[\mathbf{A},\mathbf{B}]/2} = e^{\mathbf{B}}e^{\mathbf{A}}e^{[\mathbf{A},\mathbf{B}]/2}, \text{ where: } [\mathbf{A},[\mathbf{A},\mathbf{B}]] = \mathbf{0} = [\mathbf{B},[\mathbf{A},\mathbf{B}]] \quad (\text{left as an exercise})$$

$$\begin{aligned}\mathbf{C}(a,b) &= e^{i(\mathbf{b}\mathbf{x}-\mathbf{a}\mathbf{p})/\hbar} = e^{i\mathbf{b}\mathbf{x}/\hbar}e^{-i\mathbf{a}\mathbf{p}/\hbar}e^{-ab[\mathbf{x},\mathbf{p}]/2\hbar^2} = e^{i\mathbf{b}\mathbf{x}/\hbar}e^{-i\mathbf{a}\mathbf{p}/\hbar}e^{-iab/2\hbar} \\ &= \mathbf{B}(b)\mathbf{T}(a)e^{-iab/2\hbar} = \mathbf{T}(a)\mathbf{B}(b)e^{iab/2\hbar}\end{aligned}$$

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Reordering only affects the overall phase.

$$\begin{aligned} \mathbf{C}(a,b) &= e^{i(\mathbf{b}\mathbf{x}-\mathbf{a}\mathbf{p})/\hbar} = e^{ib(\mathbf{a}^\dagger+\mathbf{a})/\sqrt{2\hbar M\omega} + a(\mathbf{a}^\dagger-\mathbf{a})\sqrt{M\omega/2\hbar}} \\ &= e^{\alpha\mathbf{a}^\dagger - \alpha^*\mathbf{a}} = e^{-|\alpha|^2/2}e^{\alpha\mathbf{a}^\dagger}e^{-\alpha^*\mathbf{a}} = e^{|\alpha|^2/2}e^{-\alpha^*\mathbf{a}}e^{\alpha\mathbf{a}^\dagger} \end{aligned}$$

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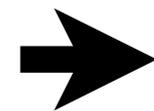
Expectation values of position, momentum, and uncertainty for eigenstate $|n\rangle$

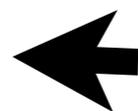
Harmonic oscillator beat dynamics of mixed states

Oscillator coherent states (“Shoved” and “kicked” states)

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 *Time evolution of coherent state*

Properties of coherent state and “squeezed” states 

2-D $\mathbf{a}^\dagger\mathbf{a}$ algebra of $U(2)$ representations and $R(3)$ angular momentum operators

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Time evolution operator for constant \mathbf{H} has general form : $\mathbf{U}(t, 0) = e^{-i\mathbf{H}t/\hbar}$

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(x_t, p_t) mimics classical oscillator

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Real and imaginary parts (x_t and $p_t/M\omega$) of α_t go clockwise on phasor circle

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Time evolution of coherent state

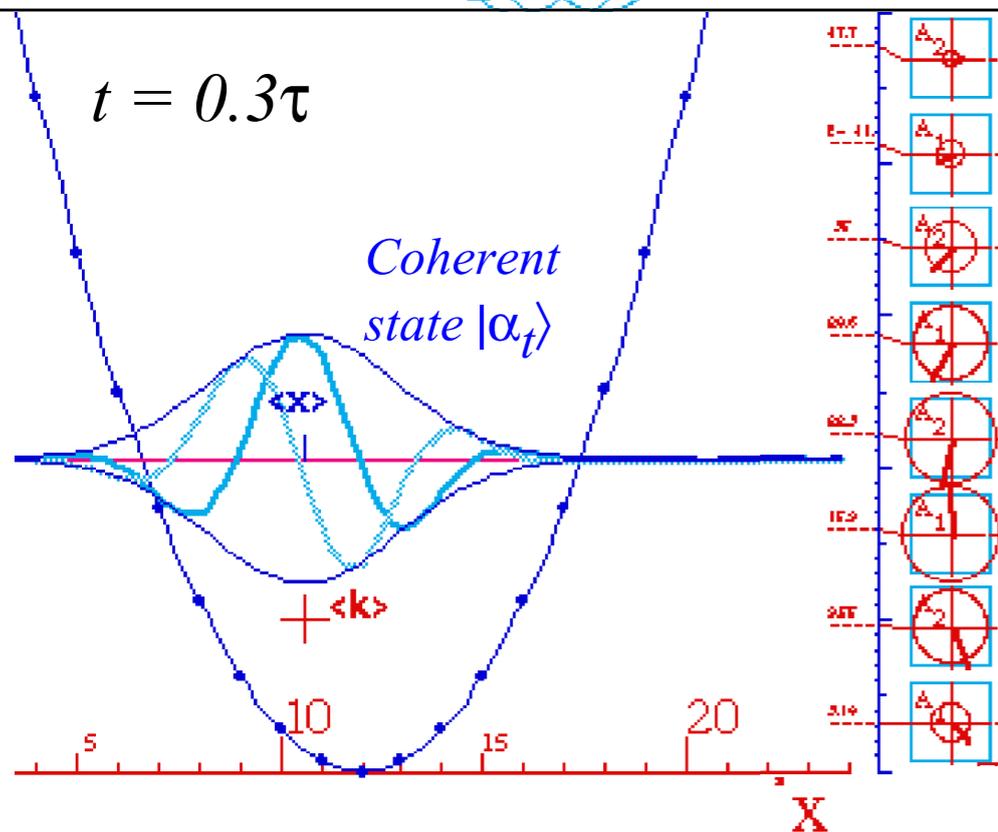
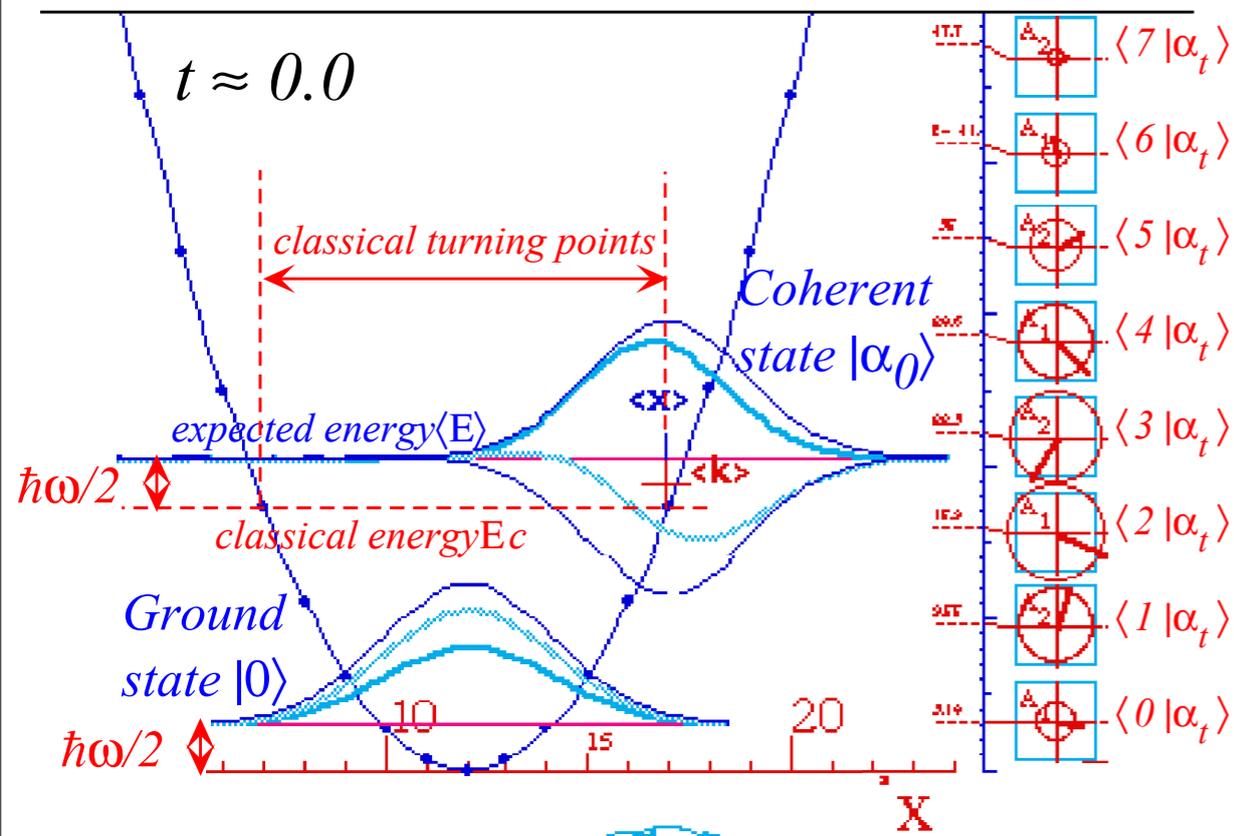
➔ *Properties of coherent state and “squeezed” states* **←**

2-D $\mathbf{a}^\dagger\mathbf{a}$ algebra of $U(2)$ representations and $R(3)$ angular momentum operators

Properties of coherent state

Coherent ket $|\alpha(x_0, p_0)\rangle$ is eigenvector of destruct-op. **a.**

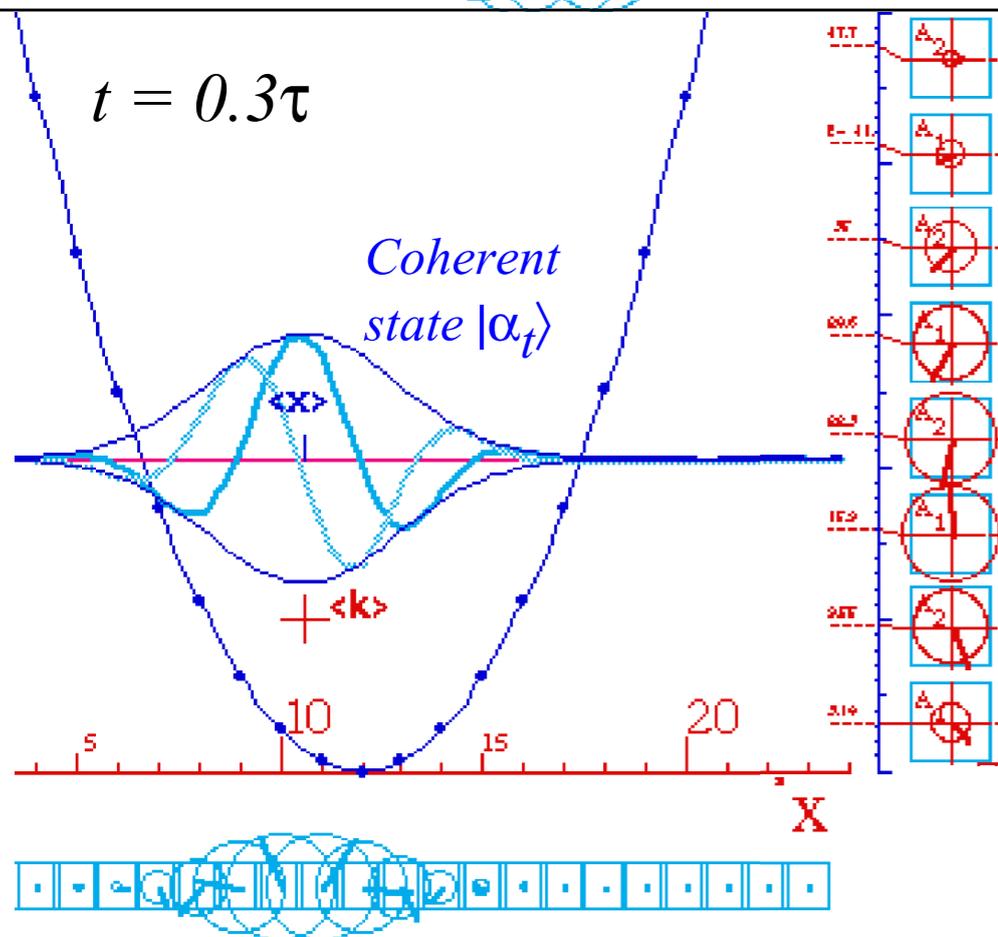
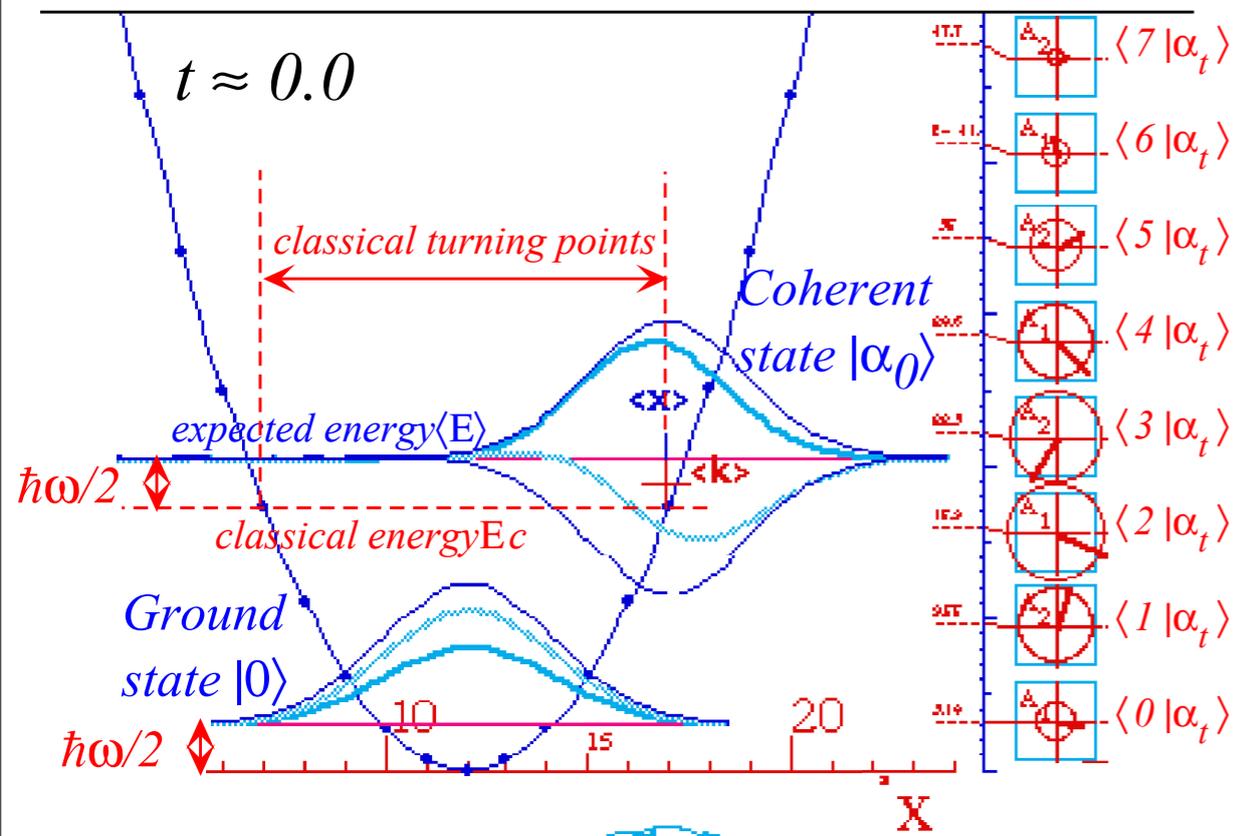
$$\mathbf{a}|\alpha_0(x_0, p_0)\rangle = e^{-|\alpha_0|^2/2} \sum_{n=0}^{\infty} \frac{(\alpha_0)^n}{\sqrt{n!}} \mathbf{a}|n\rangle$$



Properties of coherent state

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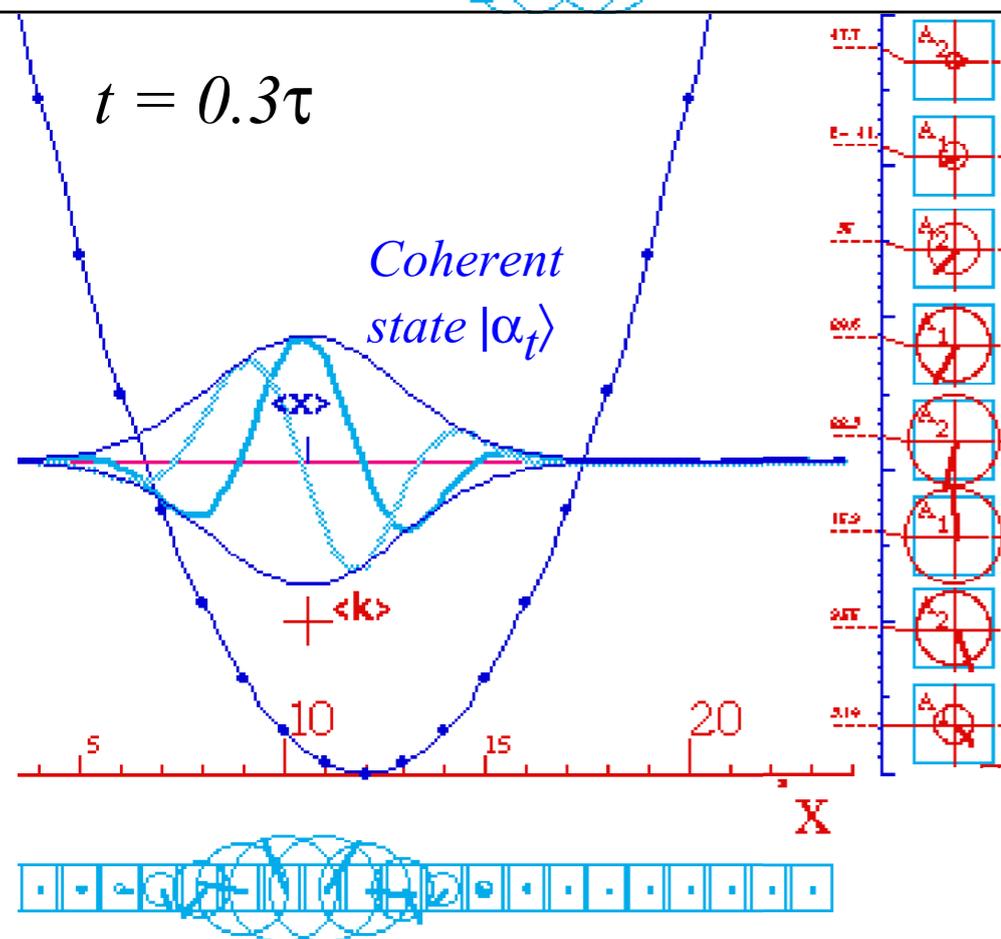
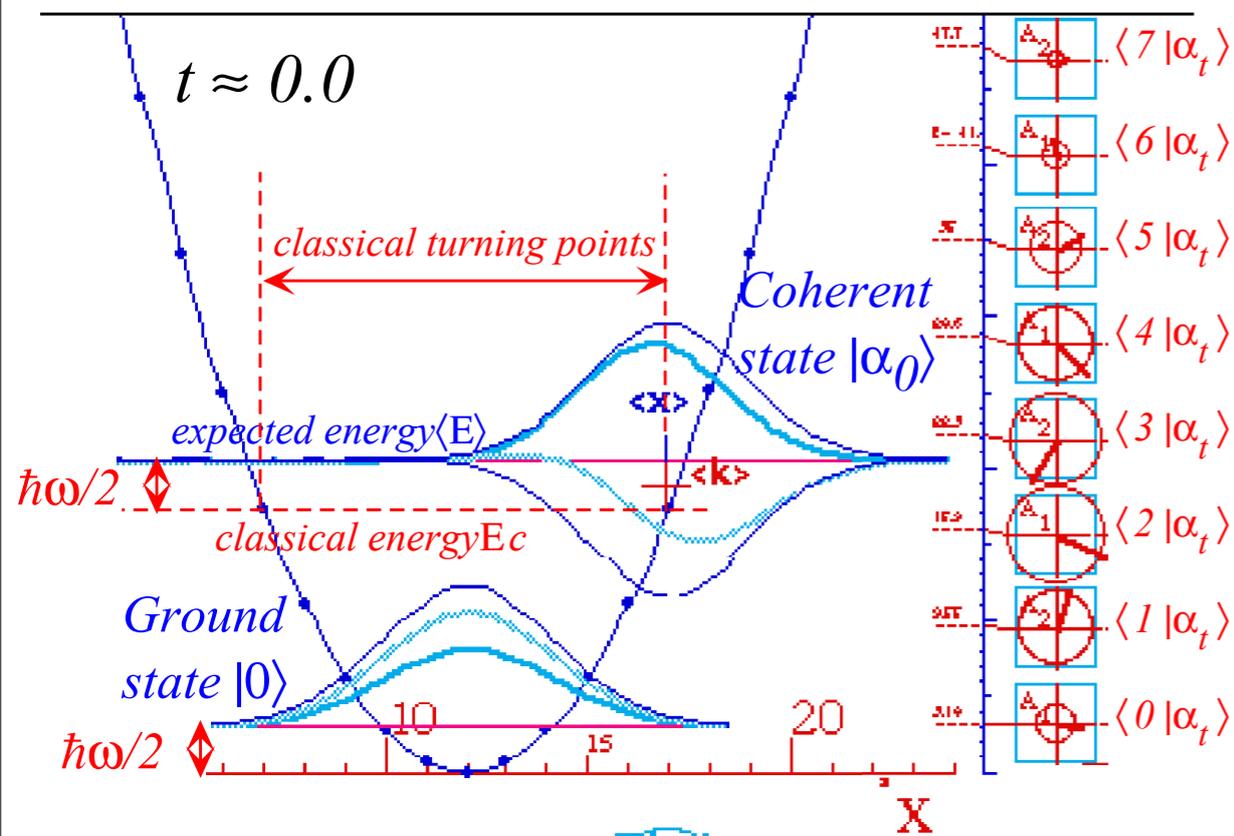
$$\begin{aligned} \mathbf{a}|\alpha_0(x_0, p_0)\rangle &= e^{-|\alpha_0|^2/2} \sum_{n=0}^{\infty} \frac{(\alpha_0)^n}{\sqrt{n!}} \mathbf{a}|n\rangle \\ &= e^{-|\alpha_0|^2/2} \sum_{n=0}^{\infty} \frac{(\alpha_0)^n}{\sqrt{n!}} \sqrt{n} |n-1\rangle \\ &= \alpha_0 |\alpha_0(x_0, p_0)\rangle \end{aligned}$$



Properties of coherent state

Coherent ket $|\alpha(x_0, p_0)\rangle$ is eigenvector of destruct-op. **a.**

$$\begin{aligned} \mathbf{a}|\alpha_0(x_0, p_0)\rangle &= e^{-|\alpha_0|^2/2} \sum_{n=0}^{\infty} \frac{(\alpha_0)^n}{\sqrt{n!}} \mathbf{a}|n\rangle \\ &= e^{-|\alpha_0|^2/2} \sum_{n=0}^{\infty} \frac{(\alpha_0)^n}{\sqrt{n!}} \sqrt{n} |n-1\rangle \\ &= \alpha_0 |\alpha_0(x_0, p_0)\rangle \quad \text{with eigenvalue } \alpha_0 \end{aligned}$$



Properties of coherent state

Coherent ket $|\alpha(x_0, p_0)\rangle$ is eigenvector of destruct-op. **a**.

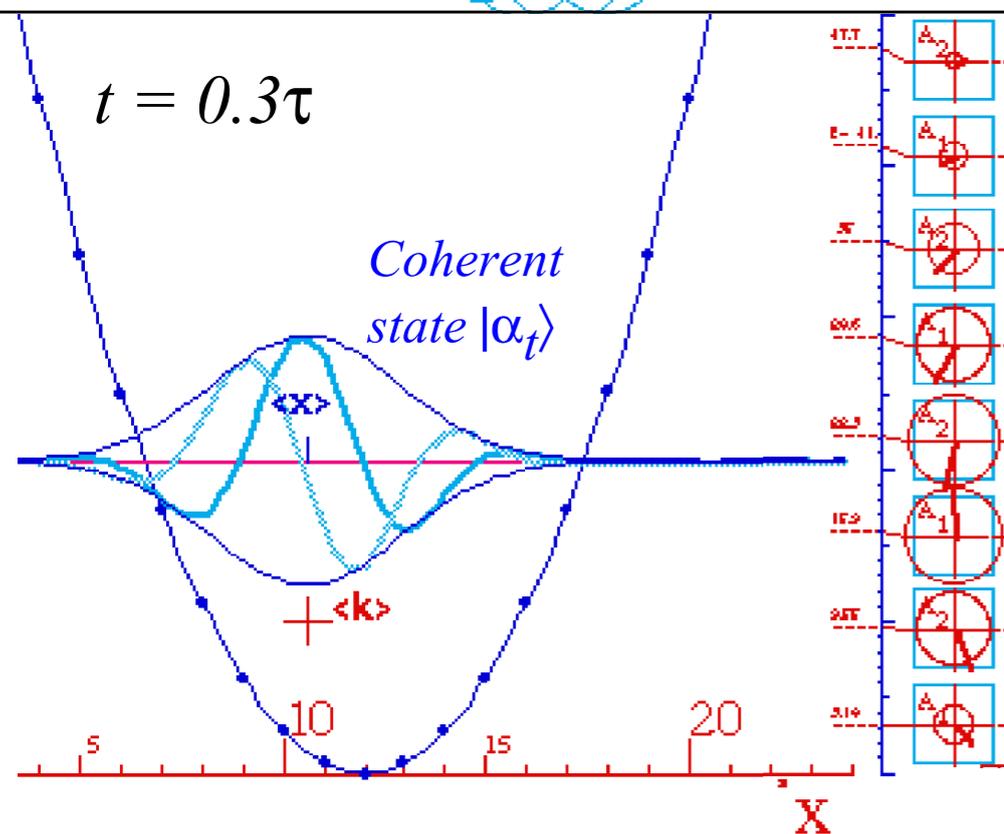
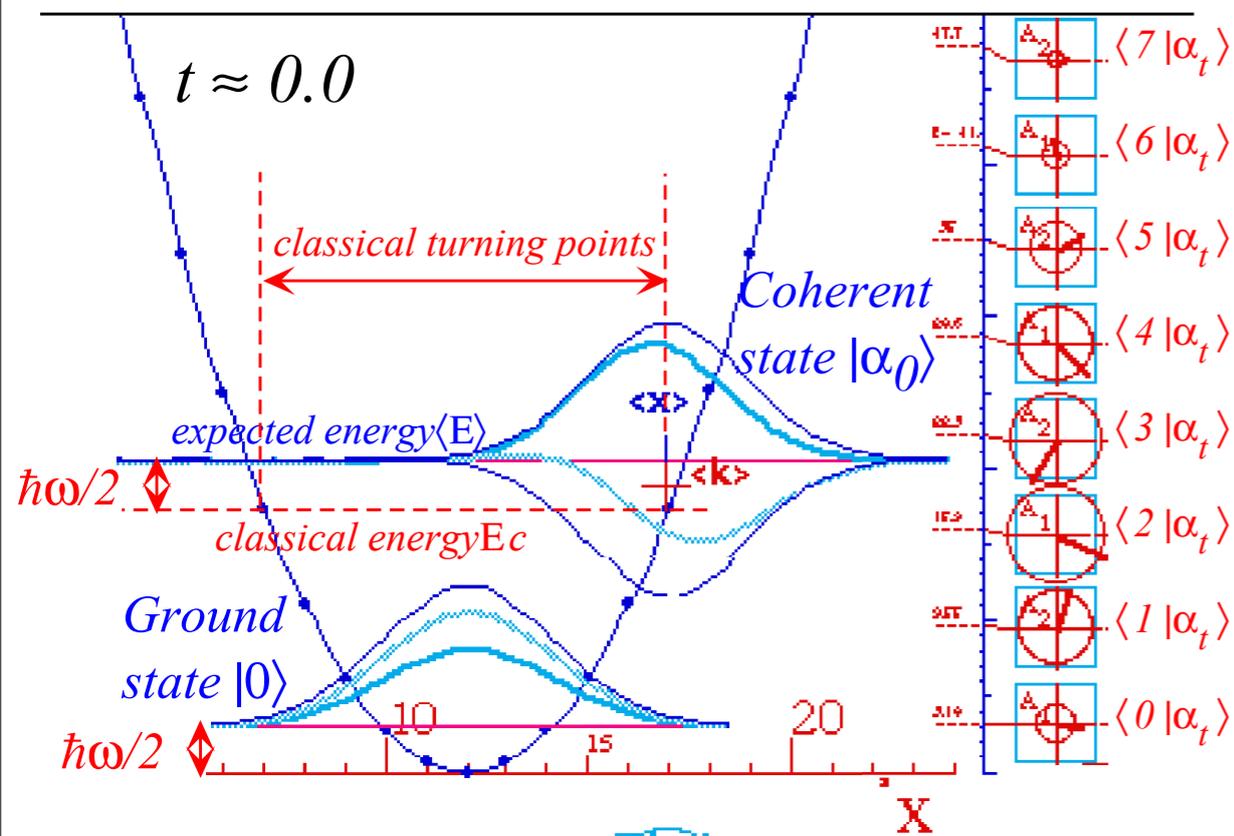
$$\begin{aligned} \mathbf{a}|\alpha_0(x_0, p_0)\rangle &= e^{-|\alpha_0|^2/2} \sum_{n=0}^{\infty} \frac{(\alpha_0)^n}{\sqrt{n!}} \mathbf{a}|n\rangle \\ &= e^{-|\alpha_0|^2/2} \sum_{n=0}^{\infty} \frac{(\alpha_0)^n}{\sqrt{n!}} \sqrt{n}|n-1\rangle \\ &= \alpha_0 |\alpha_0(x_0, p_0)\rangle \quad \text{with eigenvalue } \alpha_0 \end{aligned}$$

Coherent bra $\langle\alpha(x_0, p_0)|$ is eigenvector of create-op. **a**[†].

$$\langle\alpha_0(x_0, p_0)| \mathbf{a}^\dagger = \langle\alpha_0(x_0, p_0)| \alpha_0^*$$

Expected quantum energy has simple time independent form.

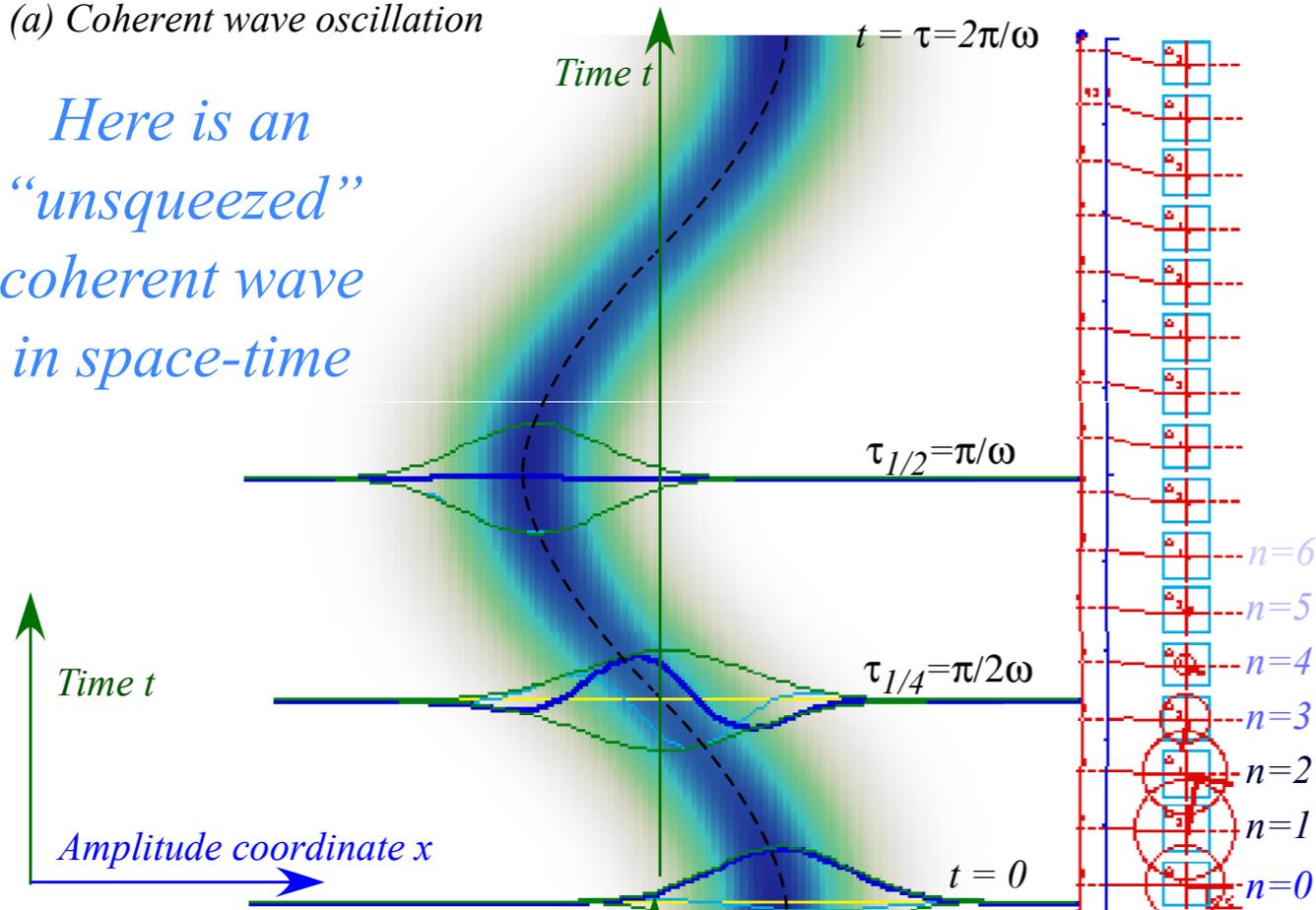
$$\begin{aligned} \langle E \rangle_{\alpha_0} &= \langle\alpha_0(x_0, p_0)| \mathbf{H} |\alpha_0(x_0, p_0)\rangle \\ &= \langle\alpha_0(x_0, p_0)| \left(\hbar\omega \mathbf{a}^\dagger \mathbf{a} + \frac{\hbar\omega}{2} \mathbf{1} \right) |\alpha_0(x_0, p_0)\rangle \\ &= \hbar\omega \alpha_0^* \alpha_0 + \frac{\hbar\omega}{2} \end{aligned}$$



Properties of “squeezed” coherent states

(a) Coherent wave oscillation

Here is an
“unsqueezed”
coherent wave
in space-time

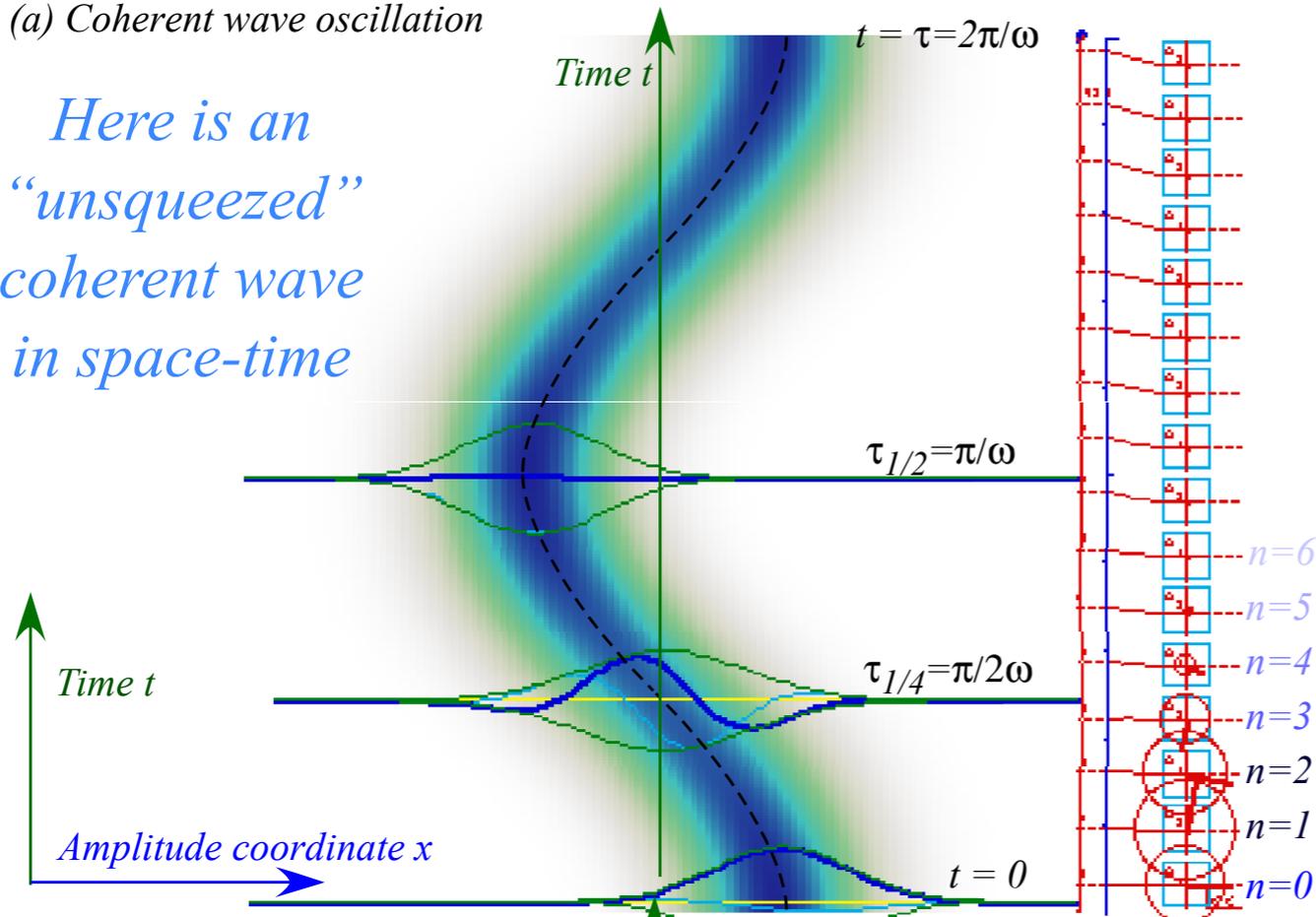


Yeah! Cosine trajectory!

Properties of “squeezed” coherent states

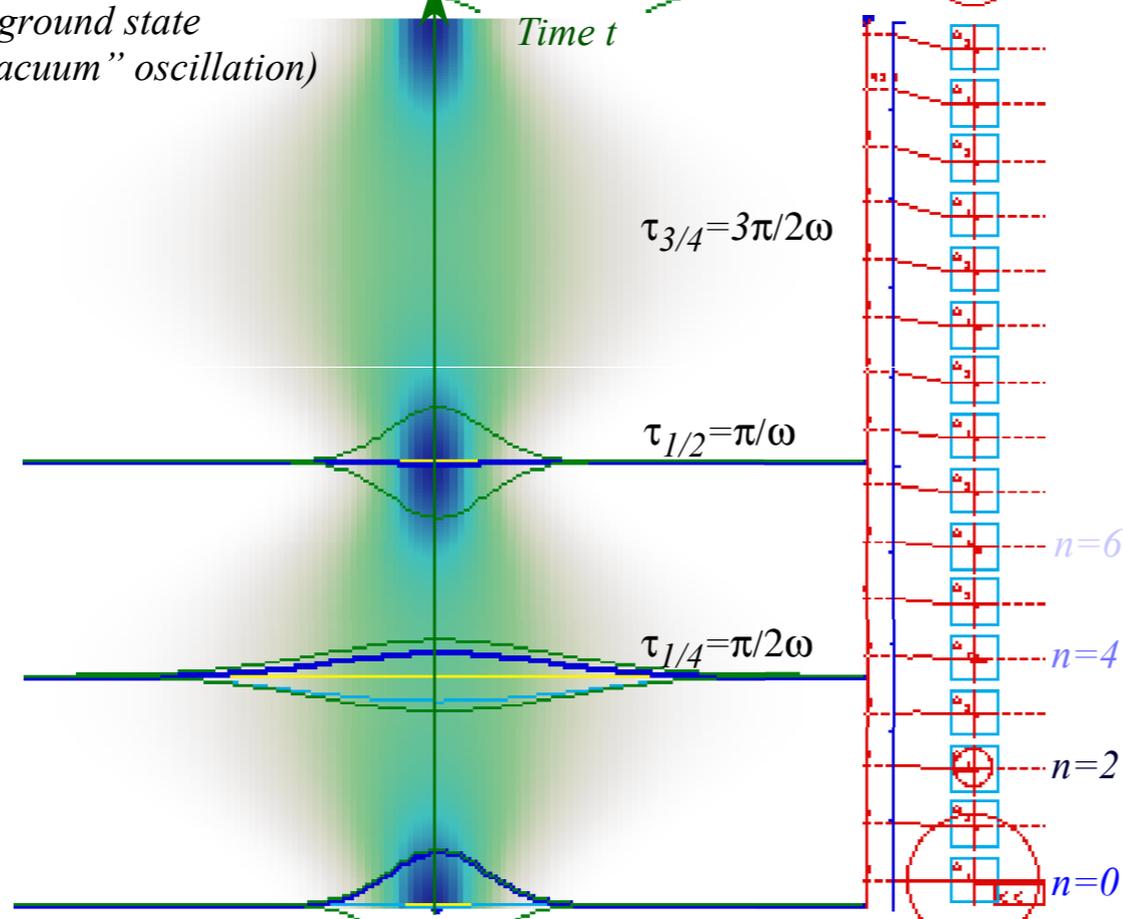
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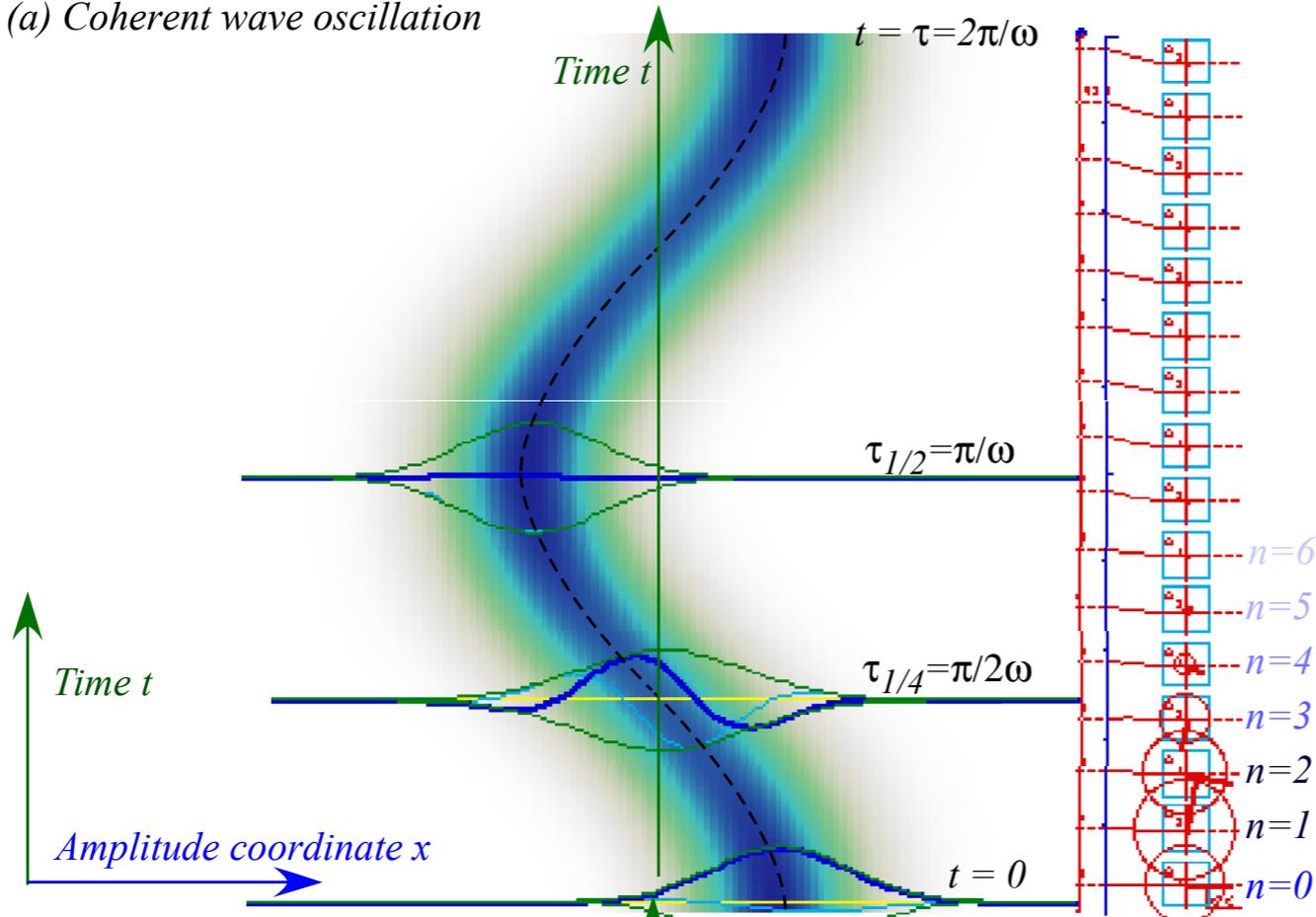
(b) Squeezed ground state
 (“Squeezed vacuum” oscillation)



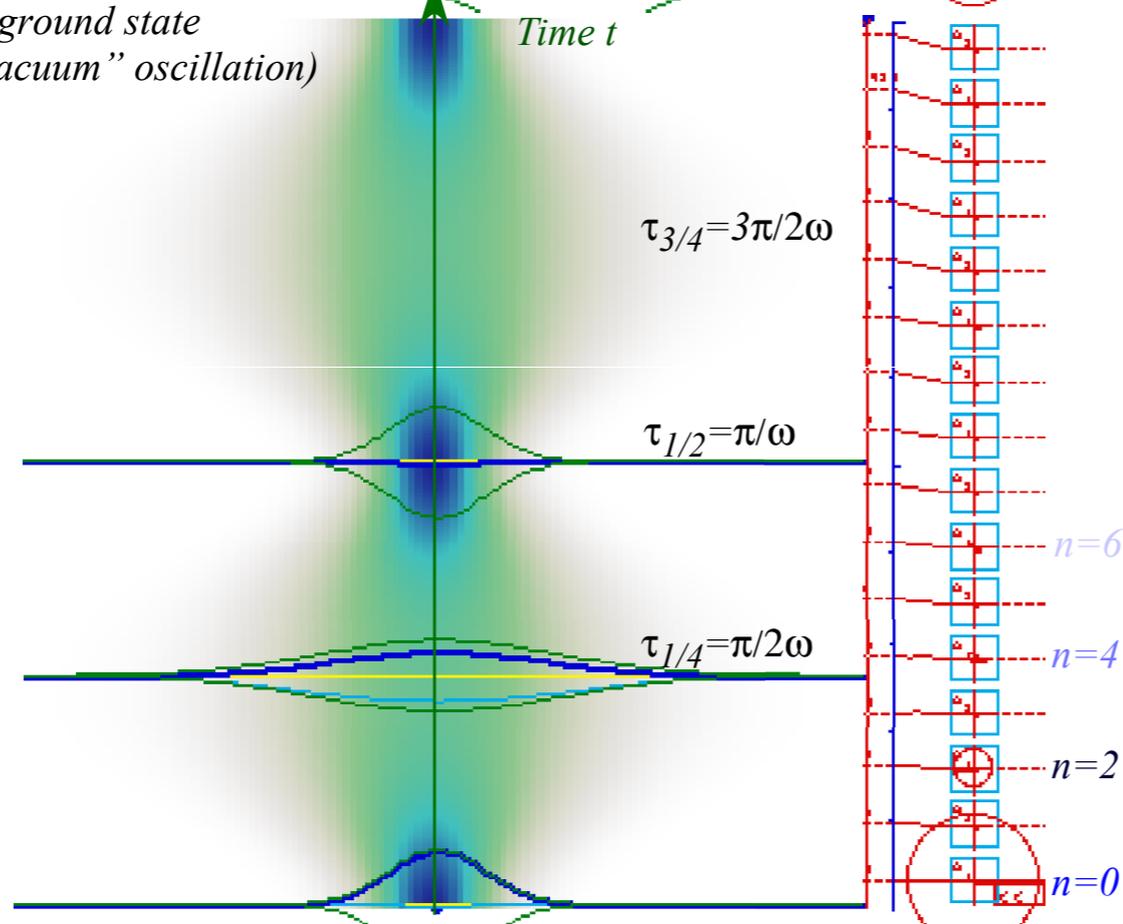
what happens if you apply
operators with non-linear “tensor”
exponents $\exp(s\mathbf{x}^2)$, $\exp(f\mathbf{p}^2)$, etc.

Properties of "squeezed" coherent states

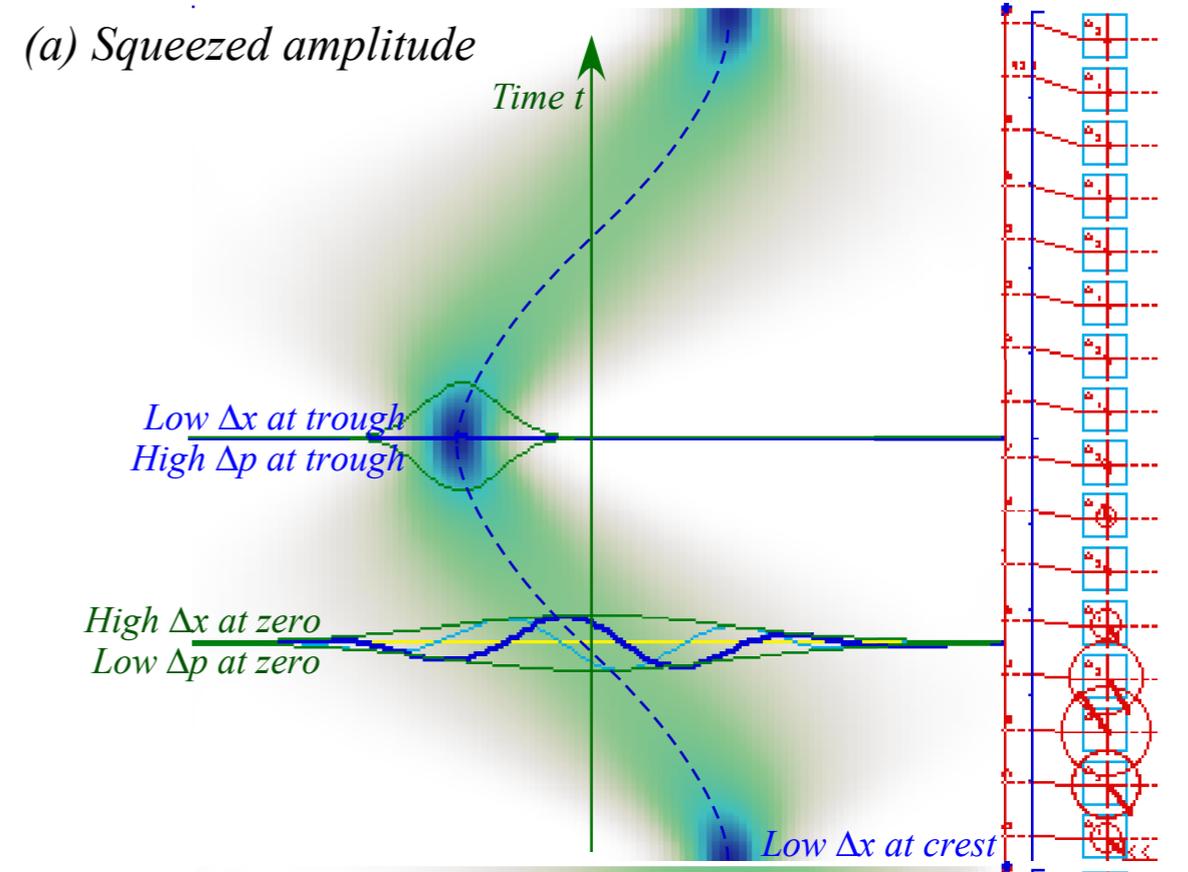
(a) Coherent wave oscillation



(b) Squeezed ground state ("Squeezed vacuum" oscillation)



(a) Squeezed amplitude



(b) Squeezed phase

