

# Group Theory in Quantum Mechanics

## Lecture 8 (2.7.13)

### Spinor and vector representations of $U(2)$ and $R(3)$ Operators

(Quantum Theory for Computer Age - Ch. 10A-B of Unit 3)

(Principles of Symmetry, Dynamics, and Spectroscopy - Sec. 1-3 of Ch. 5)

Review: How “Crazy-Thing”-Theorem makes spinor and vector representation matrices

Half-angle  $\Theta/2 = \varphi$  replacement and Darboux crank axis operators

Operator-on-Operator transformations

Product algebra for Pauli's  $\sigma_\mu$  and Hamilton's  $q_\mu = -i\sigma_\mu$

Group product algebra

Jordan-Pauli identity and  $U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - formula

Transformation  $\mathbf{R}[\Theta]\sigma_\mu\mathbf{R}[\Theta]^\dagger$  of spinor  $\sigma_\mu$ -operators

Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  of group-operators

Operator-on-Operator transformations

Geometry of groups: Hamilton's turns and It's all done with mirrors!

Group product geometry

$U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - geometry

Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  geometry

Euler  $\mathbf{R}(\alpha\beta\gamma)$  versus Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  related to Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  rotation  $\Theta = 0 - 4\pi$ -sequence  $[\varphi\vartheta]$  fixed

$R(3)$ - $U(2)$  slide rule for converting  $\mathbf{R}(\alpha\beta\gamma) \leftrightarrow \mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  Sundial

➔ Review: How “Crazy-Thing”-Theorem makes spinor and vector representation matrices  
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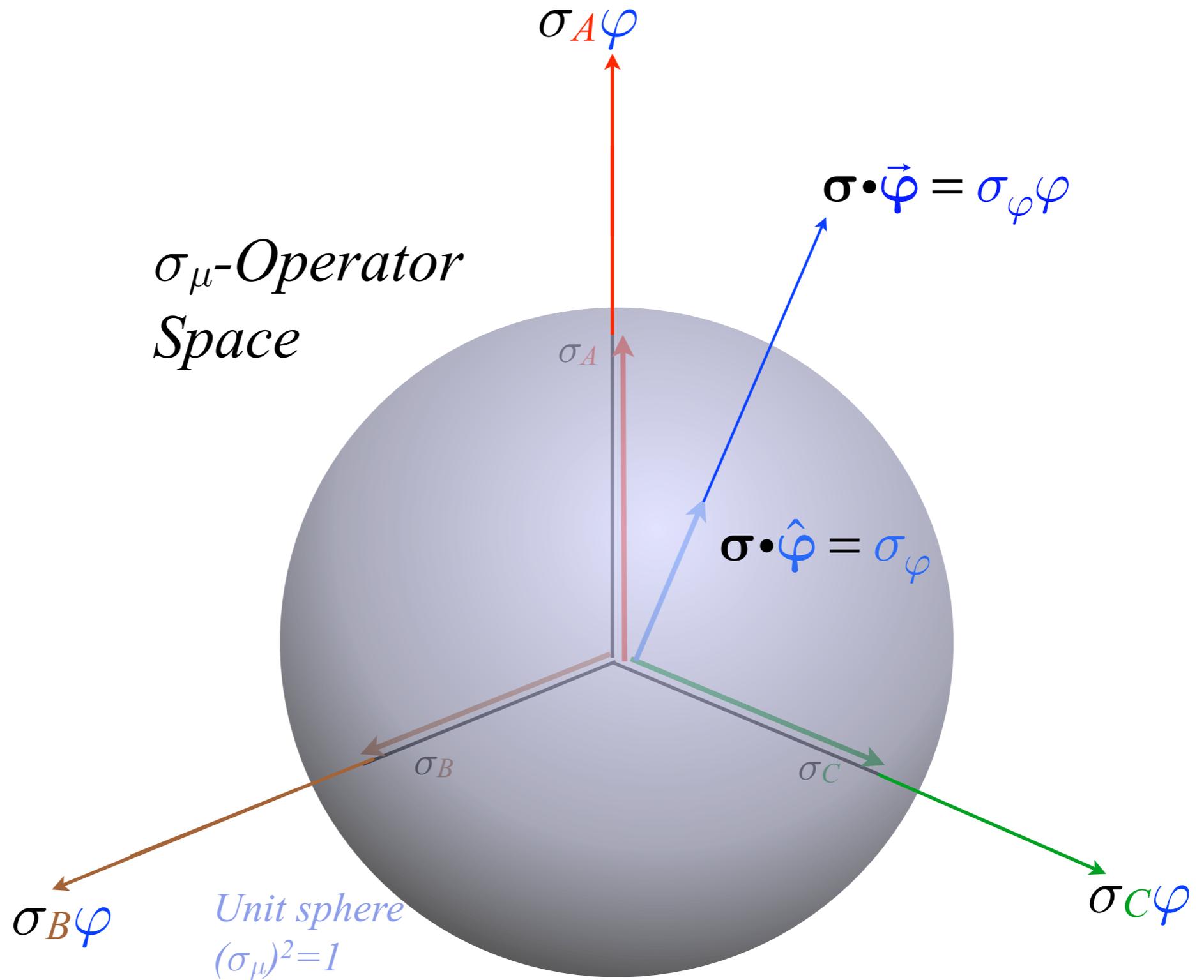
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Not-so-Crazy Thing:  $\sigma_\varphi = \sigma_A \hat{\varphi}_A + \sigma_B \hat{\varphi}_B + \sigma_C \hat{\varphi}_C = \sigma \cdot \hat{\varphi} = \sigma \cdot \vec{\varphi} / \varphi$  where:  $(\sigma_\varphi)^2 = 1$



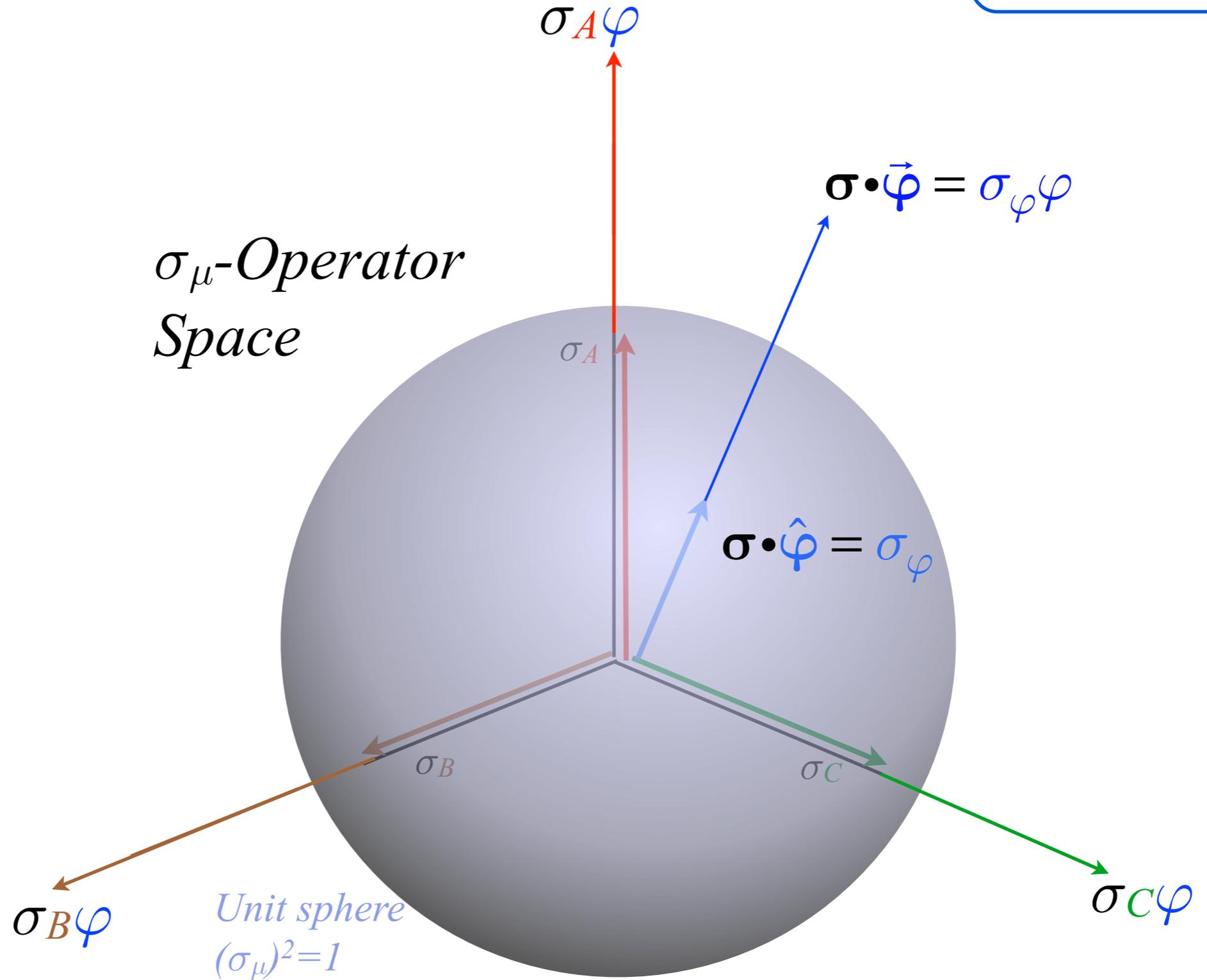


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Crazy Thing:  $i\sigma_\varphi = -i\sigma_A \hat{\varphi}_A - i\sigma_B \hat{\varphi}_B - i\sigma_C \hat{\varphi}_C = -i\boldsymbol{\sigma} \cdot \hat{\boldsymbol{\varphi}} = -i\boldsymbol{\sigma} \cdot \bar{\boldsymbol{\varphi}} / \varphi$

satisfies crazy requirement:  $(i\sigma_\varphi)^2 = (-i\sigma_\varphi)^2 = -\mathbf{1}$

The Crazy Thing Theorem:  
 If  $(i\sigma_\varphi)^2 = -\mathbf{1}$   
 Then:  
 $e(i\sigma_\varphi)\theta = \mathbf{1}\cos\theta + (i\sigma_\varphi)\sin\theta$



Not-so-Crazy Thing:  $\sigma_\varphi = \sigma_A \hat{\varphi}_A + \sigma_B \hat{\varphi}_B + \sigma_C \hat{\varphi}_C = \boldsymbol{\sigma} \cdot \hat{\boldsymbol{\varphi}} = \boldsymbol{\sigma} \cdot \vec{\boldsymbol{\varphi}} / \varphi$  where:  $(\sigma_\varphi)^2 = \mathbf{1}$

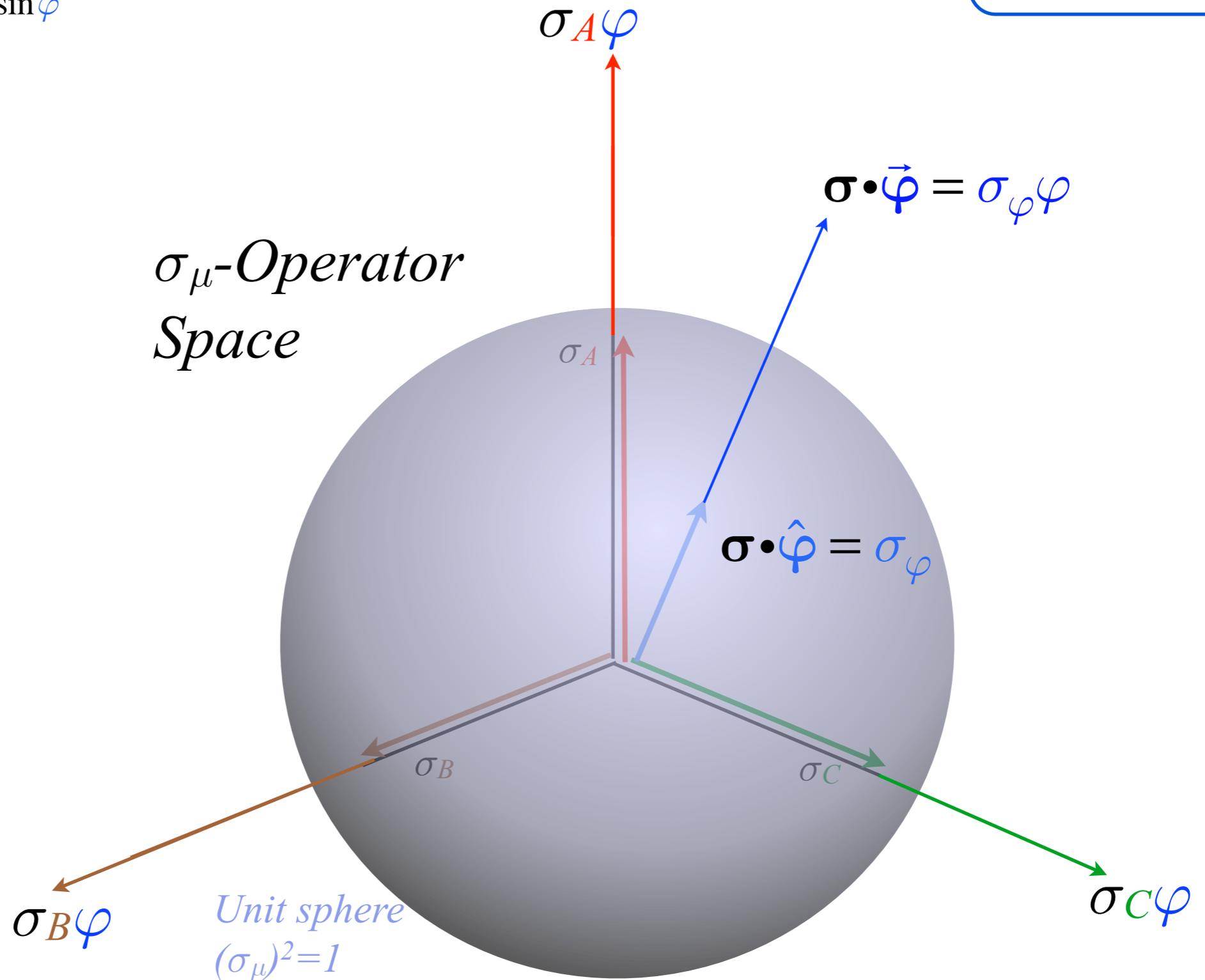
Crazy Thing:  $\text{☹} = -i\sigma_\varphi = -i\sigma_A \hat{\varphi}_A - i\sigma_B \hat{\varphi}_B - i\sigma_C \hat{\varphi}_C = -i\boldsymbol{\sigma} \cdot \hat{\boldsymbol{\varphi}} = -i\boldsymbol{\sigma} \cdot \vec{\boldsymbol{\varphi}} / \varphi$

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So:

$$e^{-i\sigma_\varphi\varphi} = \mathbf{1} \cos\varphi - i\sigma_\varphi \sin\varphi$$

The ☹  
Crazy Thing  
Theorem:  
If  $(\text{☹})^2 = -\mathbf{1}$   
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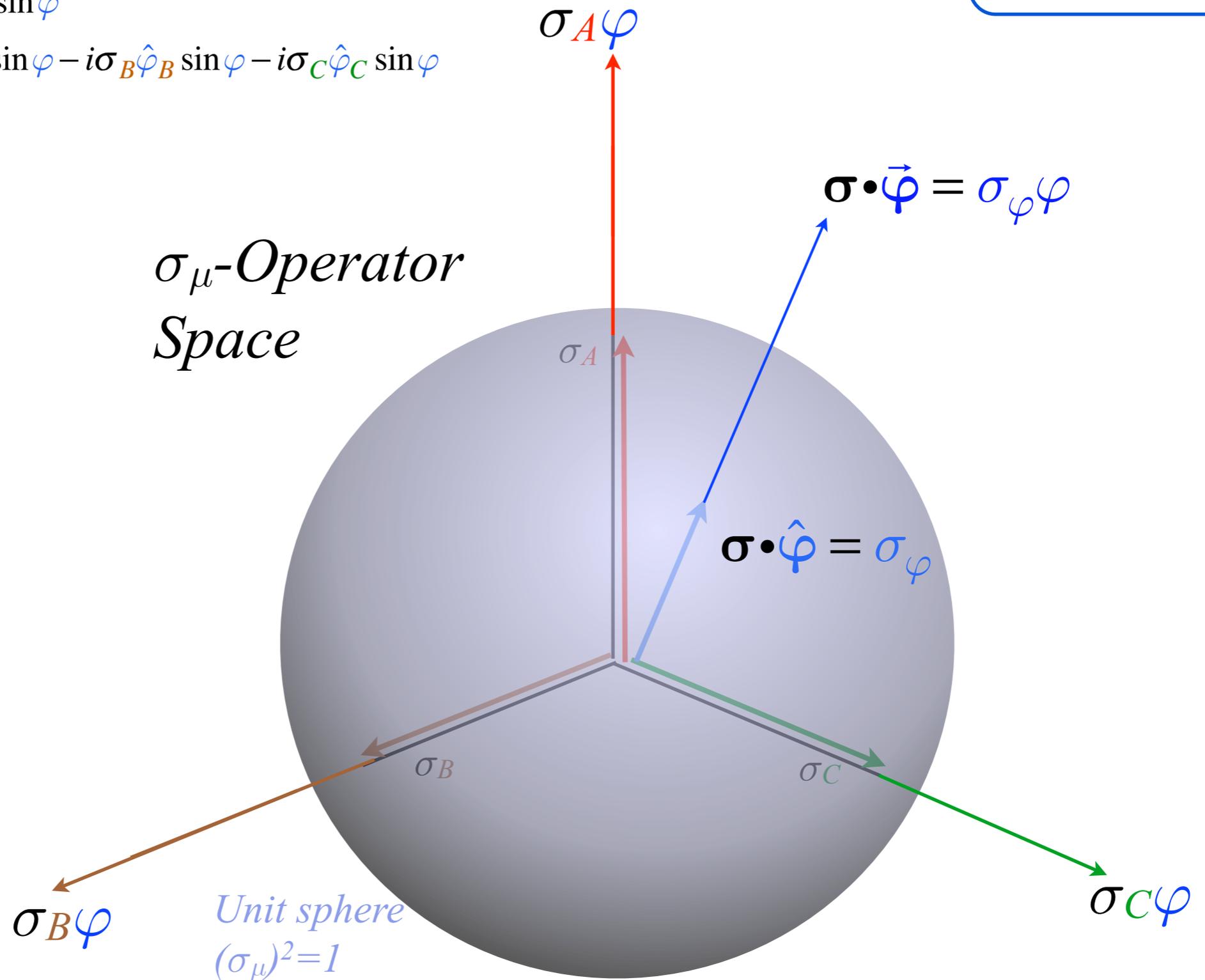
satisfies crazy requirement:  $(\text{☹️})^2 = (-i\sigma_\varphi)^2 = -1$

So:

$$e^{-i\sigma_\varphi\varphi} = 1 \cos \varphi - i\sigma_\varphi \sin \varphi$$

$$= 1 \cos \varphi - i\sigma_A \hat{\varphi}_A \sin \varphi - i\sigma_B \hat{\varphi}_B \sin \varphi - i\sigma_C \hat{\varphi}_C \sin \varphi$$

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If  $(\text{☹️})^2 = -1$   
Then:  
 $e^{(\text{☹️})\theta} = 1 \cos \theta + (\text{☹️}) \sin \theta$



Not-so-Crazy Thing:  $\sigma_\varphi = \sigma_A \hat{\varphi}_A + \sigma_B \hat{\varphi}_B + \sigma_C \hat{\varphi}_C = \sigma \cdot \hat{\varphi} = \sigma \cdot \bar{\varphi} / \varphi$  where:  $(\sigma_\varphi)^2 = \mathbf{1}$

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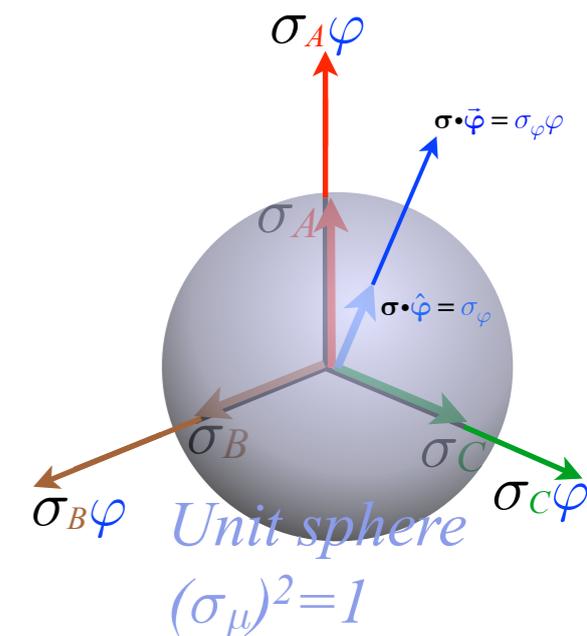
$$= \mathbf{1} \cos \varphi - i\sigma_A \hat{\varphi}_A \sin \varphi - i\sigma_B \hat{\varphi}_B \sin \varphi - i\sigma_C \hat{\varphi}_C \sin \varphi$$

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$$= \begin{pmatrix} \cos \varphi - i\hat{\varphi}_A \sin \varphi & (-i\hat{\varphi}_B - \hat{\varphi}_C) \sin \varphi \\ (-i\hat{\varphi}_B + \hat{\varphi}_C) \sin \varphi & \cos \varphi + i\hat{\varphi}_A \sin \varphi \end{pmatrix}$$

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$\sigma_\mu$ -Operator Space



Not-so-Crazy Thing:  $\sigma_\varphi = \sigma_A \hat{\varphi}_A + \sigma_B \hat{\varphi}_B + \sigma_C \hat{\varphi}_C = \sigma \cdot \hat{\varphi} = \sigma \cdot \bar{\varphi} / \varphi$  where:  $(\sigma_\varphi)^2 = \mathbf{1}$

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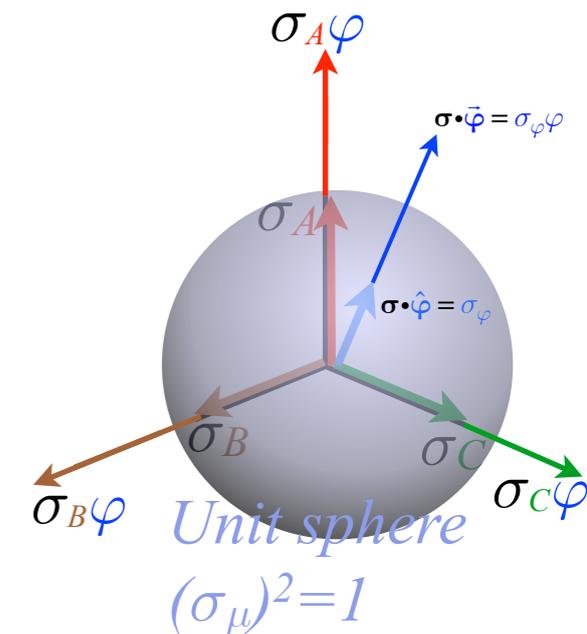
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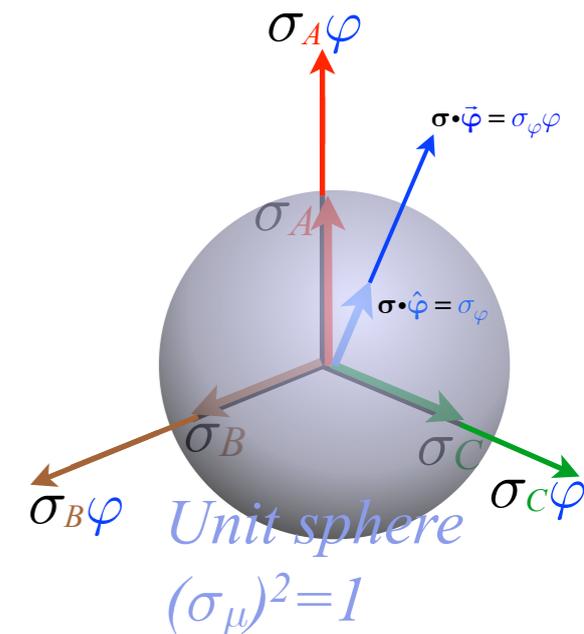
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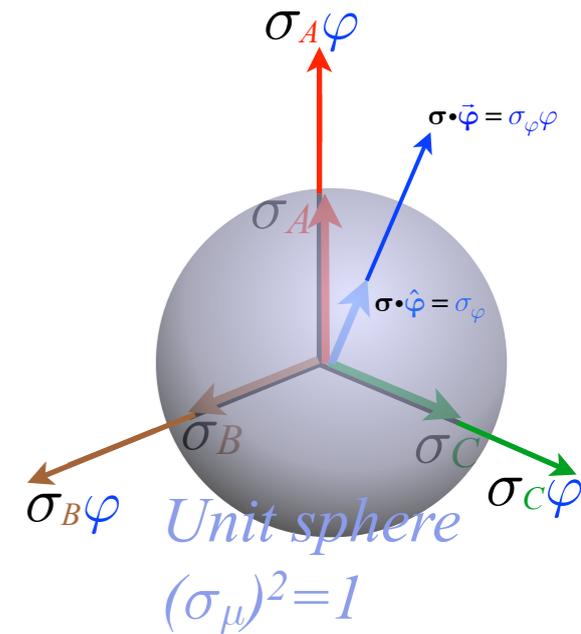
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So:

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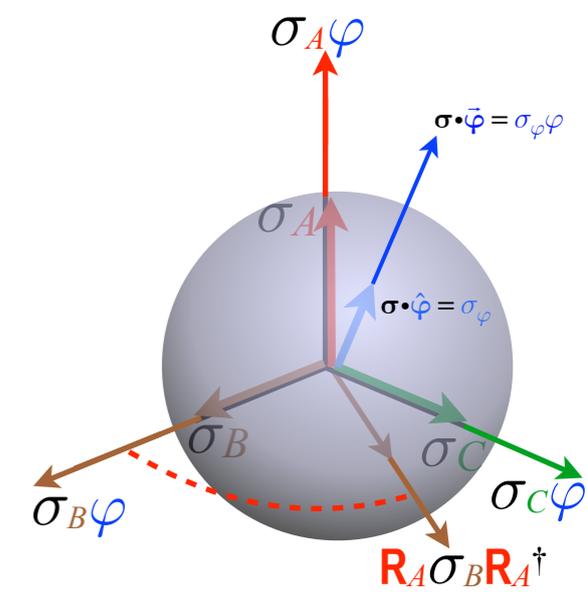
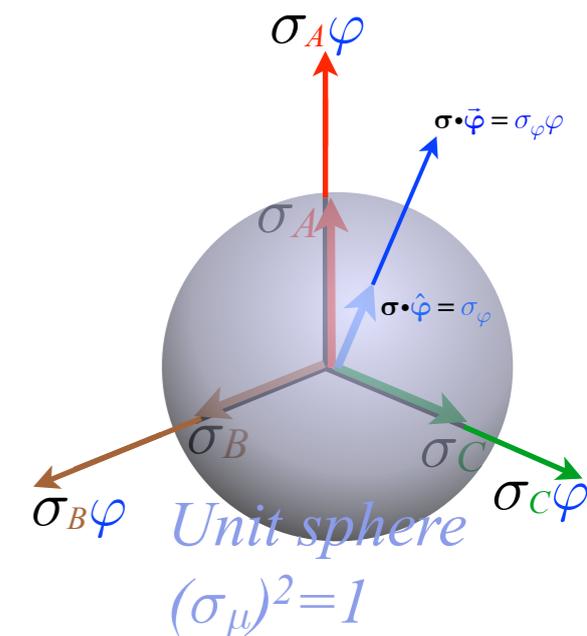
$\mathbf{R}_A \sigma_B \mathbf{R}_A^\dagger$  example:  $\sigma_B$  rotated by  $\mathbf{R}_A = e^{-i\sigma_A \varphi}$  shows 3D-space double angle  $\Theta = 2\varphi$

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Then:

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$\sigma_\mu$ -Operator Space



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$$e^{-i\sigma_\varphi\varphi} = \mathbf{1} \cos \varphi - i\sigma_\varphi \sin \varphi$$

$$= \mathbf{1} \cos \varphi - i\sigma_A \hat{\varphi}_A \sin \varphi - i\sigma_B \hat{\varphi}_B \sin \varphi - i\sigma_C \hat{\varphi}_C \sin \varphi$$

$$= \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \cos \varphi - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \hat{\varphi}_A \sin \varphi - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \hat{\varphi}_B \sin \varphi - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \hat{\varphi}_C \sin \varphi$$

$$= \begin{pmatrix} \cos \varphi - i\hat{\varphi}_A \sin \varphi & (-i\hat{\varphi}_B - \hat{\varphi}_C) \sin \varphi \\ (-i\hat{\varphi}_B + \hat{\varphi}_C) \sin \varphi & \cos \varphi + i\hat{\varphi}_A \sin \varphi \end{pmatrix}$$

$$= \begin{pmatrix} \cos \varphi - i \sin \varphi & 0 \\ 0 & \cos \varphi + i \sin \varphi \end{pmatrix} = \begin{pmatrix} e^{-i\varphi} & 0 \\ 0 & e^{+i\varphi} \end{pmatrix} \text{ for: } \hat{\varphi} = \begin{pmatrix} \hat{\varphi}_A \\ \hat{\varphi}_B \\ \hat{\varphi}_C \end{pmatrix} = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix} \text{ or: } \sigma \cdot \hat{\varphi} = \sigma_A$$

$$= \mathbf{R}_A = e^{-i\sigma_A \varphi}$$

$$= \mathbf{R}_B = e^{-i\sigma_B \varphi} = \begin{pmatrix} \cos \varphi & -i \sin \varphi \\ -i \sin \varphi & \cos \varphi \end{pmatrix} \text{ for: } \hat{\varphi} = \begin{pmatrix} \hat{\varphi}_A \\ \hat{\varphi}_B \\ \hat{\varphi}_C \end{pmatrix} = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix} \text{ or: } \sigma \cdot \hat{\varphi} = \sigma_B$$

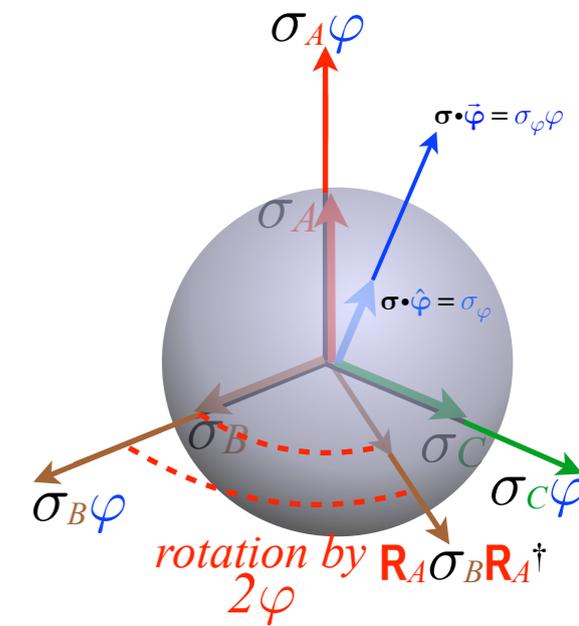
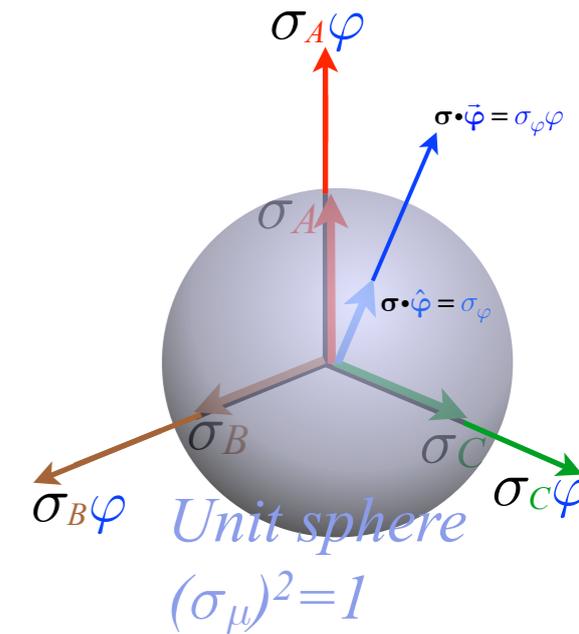
$$= \mathbf{R}_C = e^{-i\sigma_C \varphi} = \begin{pmatrix} \cos \varphi & -\sin \varphi \\ \sin \varphi & \cos \varphi \end{pmatrix} \text{ for: } \hat{\varphi} = \begin{pmatrix} \hat{\varphi}_A \\ \hat{\varphi}_B \\ \hat{\varphi}_C \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \text{ or: } \sigma \cdot \hat{\varphi} = \sigma_C$$

$\mathbf{R}_A \sigma_B \mathbf{R}_A^\dagger$  example:  $\sigma_B$  rotated by  $\mathbf{R}_A = e^{-i\sigma_A \varphi}$  shows 3D-space double angle  $\Theta = 2\varphi$

$$e^{-i\sigma_A \varphi} \sigma_B e^{+i\sigma_A \varphi} = \begin{pmatrix} e^{-i\varphi} & 0 \\ 0 & e^{+i\varphi} \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} e^{+i\varphi} & 0 \\ 0 & e^{-i\varphi} \end{pmatrix} = \begin{pmatrix} 0 & e^{-i\varphi} \\ e^{+i\varphi} & 0 \end{pmatrix} \begin{pmatrix} e^{+i\varphi} & 0 \\ 0 & e^{-i\varphi} \end{pmatrix} = \begin{pmatrix} 0 & e^{-i2\varphi} \\ e^{+i2\varphi} & 0 \end{pmatrix} = \sigma_B \cos 2\varphi + \sigma_C \sin 2\varphi$$

The Crazy Thing Theorem:  
If  $(i\sigma_\varphi)^2 = -\mathbf{1}$   
Then:  
 $e^{(i\sigma_\varphi)\theta} = \mathbf{1} \cos \theta + (i\sigma_\varphi) \sin \theta$

$\sigma_\mu$ -Operator Space





Review: How “Crazy-Thing”-Theorem makes spinor and vector representation matrices

Half-angle  $\Theta/2 = \varphi$  replacement and Darboux crank axis operators

Operator-on-Operator transformations

Product algebra for Pauli's  $\sigma_\mu$  and Hamilton's  $q_\mu = -i\sigma_\mu$

Group product algebra

Jordan-Pauli identity and  $U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - formula

Transformation  $\mathbf{R}[\Theta]\sigma_\mu\mathbf{R}[\Theta]^\dagger$  of spinor  $\sigma_\mu$ -operators

Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  of group-operators

Operator-on-Operator transformations

Geometry of groups: Hamilton's turns and It's all done with mirrors!

Group product geometry

$U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - geometry

Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  geometry

Euler  $\mathbf{R}(\alpha\beta\gamma)$  versus Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  related to Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  rotation  $\Theta = 0 - 4\pi$ -sequence  $[\varphi\vartheta]$  fixed

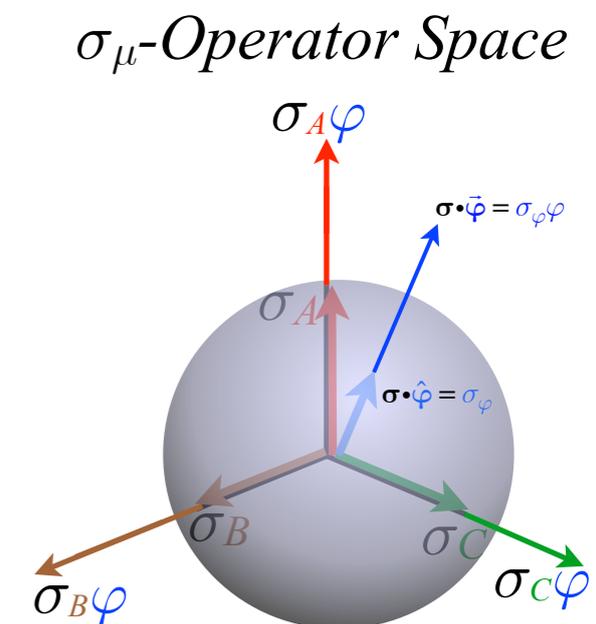
$R(3)$ - $U(2)$  slide rule for converting  $\mathbf{R}(\alpha\beta\gamma) \leftrightarrow \mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  Sundial

Half-angle  $\Theta/2 = \varphi$  replacement in rotation  $\mathbf{R}_\varphi = e^{-i\sigma_\varphi\varphi} = e^{-i\sigma \cdot \hat{\varphi}} \varphi$  where:  $\sigma_\varphi = \sigma \cdot \vec{\varphi} = \sigma_A \varphi_A + \sigma_B \varphi_B + \sigma_C \varphi_C = \sigma \cdot \hat{\varphi} \varphi$

Replace spinor angle  $\varphi$  in:  $e^{-i\sigma \cdot \vec{\varphi}} = \mathbf{R}_\varphi = \mathbf{1} \cos \varphi - i\sigma \cdot \hat{\varphi} \sin \varphi$

with 3D  $\frac{1}{2}$ -angle  $\frac{\Theta}{2} = \varphi$

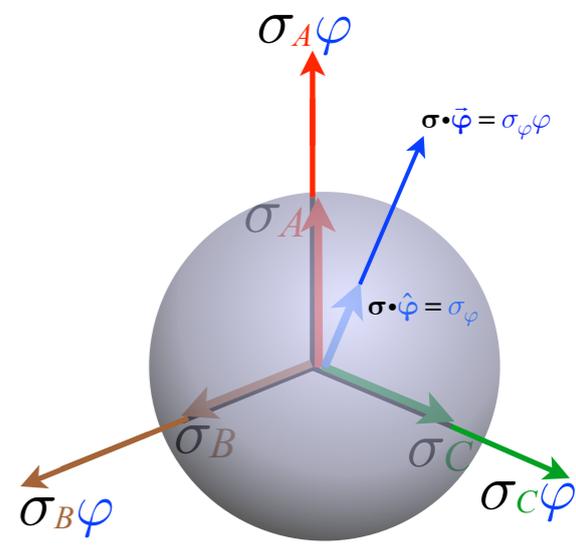


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$\sigma_\mu$ -Operator Space

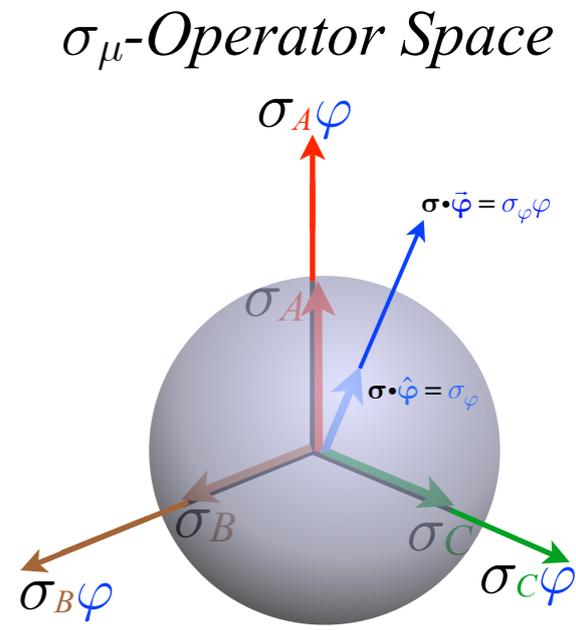


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Replace spinor angle  $\varphi$  in:  $e^{-i\sigma \cdot \vec{\varphi}} = \mathbf{R}_\varphi = \mathbf{1} \cos \varphi - i\sigma \cdot \hat{\varphi} \sin \varphi$

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Unit rotation axis vector  $\hat{\Theta} = (\hat{\Theta}_X, \hat{\Theta}_Y, \hat{\Theta}_Z) = (\cos \varphi \sin \vartheta, \sin \varphi \sin \vartheta, \cos \vartheta)$  is the same as the unit  $\hat{\varphi}$ .

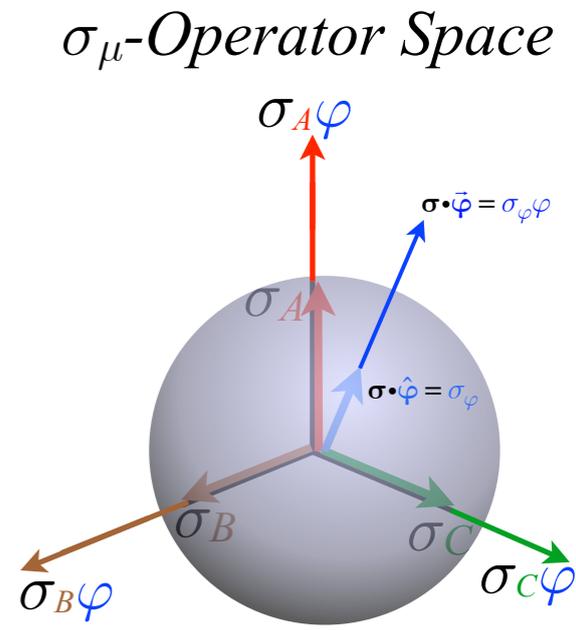


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with 3D  $\frac{1}{2}$ -angle  $\frac{\Theta}{2} = \varphi$  in:  $e^{-i\sigma \cdot \frac{\vec{\Theta}}{2}} = \mathbf{R}[\vec{\Theta}] = \mathbf{1} \cos \frac{\Theta}{2} - i\sigma \cdot \hat{\Theta} \sin \frac{\Theta}{2} = e^{-i\frac{\sigma}{2} \cdot \vec{\Theta}} = e^{-i\mathbf{S} \cdot \vec{\Theta}}$

Unit rotation axis vector  $\hat{\Theta} = (\hat{\Theta}_X, \hat{\Theta}_Y, \hat{\Theta}_Z) = (\cos\varphi \sin\vartheta, \sin\varphi \sin\vartheta, \cos\vartheta)$  is the same as the unit  $\hat{\varphi}$ .

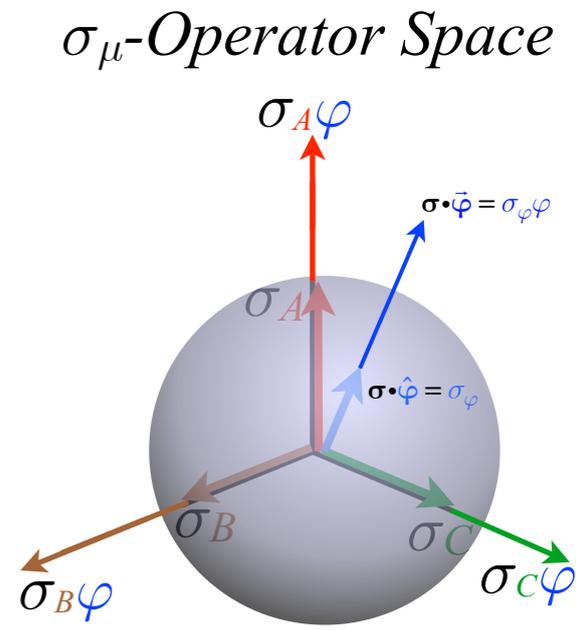


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Replace spinor angle  $\varphi$  in:  $e^{-i\sigma \cdot \vec{\varphi}} = \mathbf{R}_\varphi = \mathbf{1} \cos \varphi - i\sigma \cdot \hat{\varphi} \sin \varphi$

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Unit rotation axis vector  $\hat{\Theta} = (\hat{\Theta}_X, \hat{\Theta}_Y, \hat{\Theta}_Z) = (\cos\varphi \sin\vartheta, \sin\varphi \sin\vartheta, \cos\vartheta)$  is the same as the unit  $\hat{\varphi}$ .



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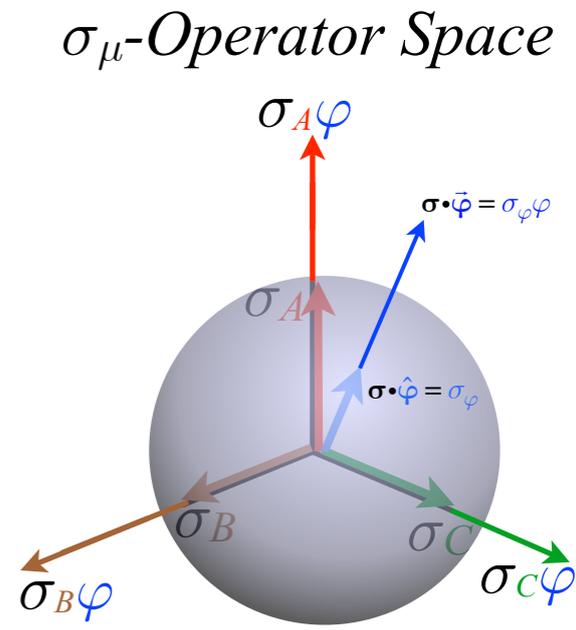
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with 3D  $\frac{1}{2}$ -angle  $\frac{\Theta}{2} = \varphi$  in:  $e^{-i\sigma \cdot \frac{\vec{\Theta}}{2}} = \mathbf{R}[\vec{\Theta}] = \mathbf{1} \cos \frac{\Theta}{2} - i\sigma \cdot \hat{\Theta} \sin \frac{\Theta}{2} = e^{-i\frac{\sigma}{2} \cdot \vec{\Theta}} = e^{-i\mathbf{S} \cdot \vec{\Theta}} = e^{-i\sigma_\varphi\varphi} = e^{-i\sigma \cdot \hat{\varphi}\varphi}$

Unit rotation axis vector  $\hat{\Theta} = (\hat{\Theta}_X, \hat{\Theta}_Y, \hat{\Theta}_Z) = (\cos\varphi \sin\vartheta, \sin\varphi \sin\vartheta, \cos\vartheta)$  is the same as the unit  $\hat{\varphi}$ .

$$\mathbf{R}[\vec{\Theta}] = \cos \frac{\Theta}{2} \mathbf{1} - i \sigma_X \hat{\Theta}_X \sin \frac{\Theta}{2} - i \sigma_Y \hat{\Theta}_Y \sin \frac{\Theta}{2} - i \sigma_Z \hat{\Theta}_Z \sin \frac{\Theta}{2}$$

$$= \cos \frac{\Theta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \hat{\Theta}_X \sin \frac{\Theta}{2} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \hat{\Theta}_Y \sin \frac{\Theta}{2} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \hat{\Theta}_Z \sin \frac{\Theta}{2}$$



Half-angle  $\Theta/2 = \varphi$  replacement in rotation  $\mathbf{R}_\varphi = e^{-i\sigma_\varphi\varphi} = e^{-i\sigma \cdot \vec{\varphi}}$  where:  $\sigma_\varphi = \sigma \cdot \vec{\varphi} = \sigma_A\varphi_A + \sigma_B\varphi_B + \sigma_C\varphi_C = \sigma \cdot \hat{\varphi}\varphi$

Replace spinor angle  $\varphi$  in:  $e^{-i\sigma \cdot \vec{\varphi}} = \mathbf{R}_\varphi = \mathbf{1} \cos \varphi - i\sigma \cdot \hat{\varphi} \sin \varphi$

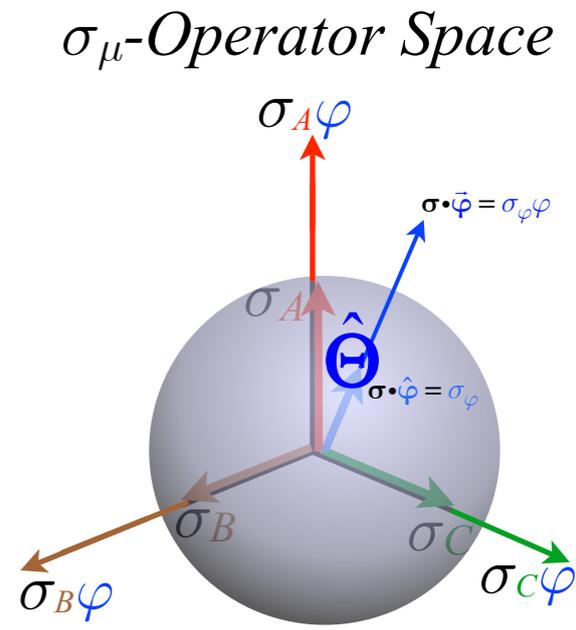
with 3D  $\frac{1}{2}$ -angle  $\frac{\Theta}{2} = \varphi$  in:  $e^{-i\sigma \cdot \frac{\vec{\Theta}}{2}} = \mathbf{R}[\vec{\Theta}] = \mathbf{1} \cos \frac{\Theta}{2} - i\sigma \cdot \hat{\Theta} \sin \frac{\Theta}{2} = e^{-i\frac{\sigma}{2} \cdot \vec{\Theta}} = e^{-i\mathbf{S} \cdot \vec{\Theta}} = e^{-i\sigma_\varphi\varphi} = e^{-i\sigma \cdot \vec{\varphi}}$

Unit rotation axis vector  $\hat{\Theta} = (\hat{\Theta}_X, \hat{\Theta}_Y, \hat{\Theta}_Z) = (\cos\varphi \sin\vartheta, \sin\varphi \sin\vartheta, \cos\vartheta)$  is the same as the unit  $\hat{\varphi}$ .

$$\mathbf{R}[\vec{\Theta}] = \cos \frac{\Theta}{2} \mathbf{1} - i \left( \sigma_X \hat{\Theta}_X \sin \frac{\Theta}{2} + \sigma_Y \hat{\Theta}_Y \sin \frac{\Theta}{2} + \sigma_Z \hat{\Theta}_Z \sin \frac{\Theta}{2} \right)$$

$$= \cos \frac{\Theta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \hat{\Theta}_X \sin \frac{\Theta}{2} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \hat{\Theta}_Y \sin \frac{\Theta}{2} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \hat{\Theta}_Z \sin \frac{\Theta}{2}$$

$$\begin{pmatrix} \langle 1 | \mathbf{R}[\vec{\Theta}] | 1 \rangle & \langle 1 | \mathbf{R}[\vec{\Theta}] | 2 \rangle \\ \langle 2 | \mathbf{R}[\vec{\Theta}] | 1 \rangle & \langle 2 | \mathbf{R}[\vec{\Theta}] | 2 \rangle \end{pmatrix} = \begin{pmatrix} \cos \frac{\Theta}{2} - i \hat{\Theta}_Z \sin \frac{\Theta}{2} & -i \sin \frac{\Theta}{2} (\hat{\Theta}_X - i \hat{\Theta}_Y) \\ -i \sin \frac{\Theta}{2} (\hat{\Theta}_X + i \hat{\Theta}_Y) & \cos \frac{\Theta}{2} + i \hat{\Theta}_Z \sin \frac{\Theta}{2} \end{pmatrix}$$



$$\begin{aligned} \hat{\Theta}_X &= \cos\varphi \sin\vartheta \\ \hat{\Theta}_Y &= \sin\varphi \sin\vartheta \\ \hat{\Theta}_Z &= \cos\vartheta \end{aligned}$$

Polar coordinates for unit axis vector  $\hat{\Theta}$

Half-angle  $\Theta/2 = \varphi$  replacement in rotation  $\mathbf{R}_\varphi = e^{-i\sigma_\varphi\varphi} = e^{-i\sigma \cdot \hat{\varphi}\varphi}$  where:  $\sigma_\varphi = \sigma \cdot \hat{\varphi} = \sigma_A\varphi_A + \sigma_B\varphi_B + \sigma_C\varphi_C = \sigma \cdot \hat{\varphi}$

Replace spinor angle  $\varphi$  in:  $e^{-i\sigma \cdot \hat{\varphi}\varphi} = \mathbf{R}_\varphi = \mathbf{1} \cos \varphi - i\sigma \cdot \hat{\varphi} \sin \varphi$

with 3D  $\frac{1}{2}$ -angle  $\frac{\Theta}{2} = \varphi$  in:  $e^{-i\sigma \cdot \frac{\hat{\Theta}}{2}\Theta} = \mathbf{R}[\hat{\Theta}] = \mathbf{1} \cos \frac{\Theta}{2} - i\sigma \cdot \hat{\Theta} \sin \frac{\Theta}{2} = e^{-i\frac{\sigma}{2} \cdot \hat{\Theta}\Theta} = e^{-i\mathbf{S} \cdot \hat{\Theta}\Theta} = e^{-i\sigma_\varphi\varphi} = e^{-i\sigma \cdot \hat{\varphi}\varphi}$

Unit rotation axis vector  $\hat{\Theta} = (\hat{\Theta}_X, \hat{\Theta}_Y, \hat{\Theta}_Z) = (\cos \varphi \sin \vartheta, \sin \varphi \sin \vartheta, \cos \vartheta)$  is the same as the unit  $\hat{\varphi}$ .

$$\mathbf{R}[\hat{\Theta}] = \cos \frac{\Theta}{2} \mathbf{1} - i \left( \sigma_X \hat{\Theta}_X \sin \frac{\Theta}{2} + \sigma_Y \hat{\Theta}_Y \sin \frac{\Theta}{2} + \sigma_Z \hat{\Theta}_Z \sin \frac{\Theta}{2} \right)$$

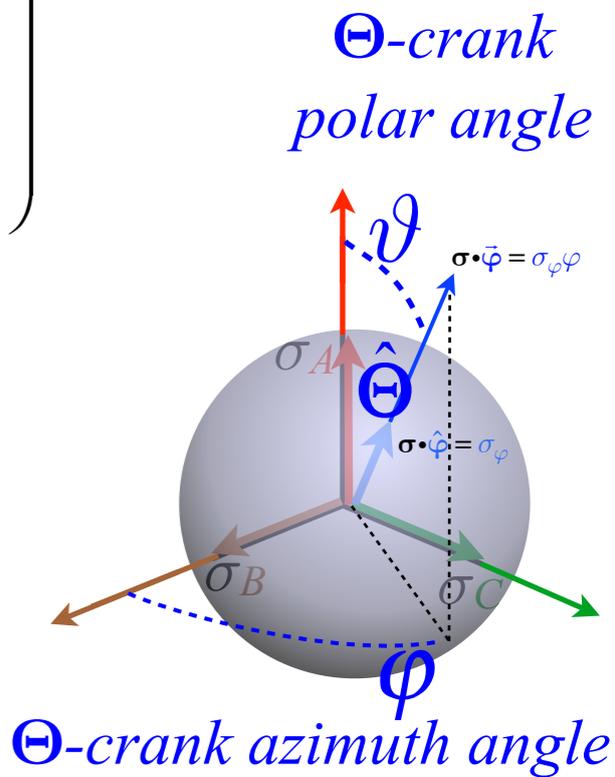
$$= \cos \frac{\Theta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \hat{\Theta}_X \sin \frac{\Theta}{2} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \hat{\Theta}_Y \sin \frac{\Theta}{2} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \hat{\Theta}_Z \sin \frac{\Theta}{2}$$

$$\begin{pmatrix} \langle 1 | \mathbf{R}[\hat{\Theta}] | 1 \rangle & \langle 1 | \mathbf{R}[\hat{\Theta}] | 2 \rangle \\ \langle 2 | \mathbf{R}[\hat{\Theta}] | 1 \rangle & \langle 2 | \mathbf{R}[\hat{\Theta}] | 2 \rangle \end{pmatrix} = \begin{pmatrix} \cos \frac{\Theta}{2} - i \hat{\Theta}_Z \sin \frac{\Theta}{2} & -i \sin \frac{\Theta}{2} (\hat{\Theta}_X - i \hat{\Theta}_Y) \\ -i \sin \frac{\Theta}{2} (\hat{\Theta}_X + i \hat{\Theta}_Y) & \cos \frac{\Theta}{2} + i \hat{\Theta}_Z \sin \frac{\Theta}{2} \end{pmatrix}$$

$$= \begin{pmatrix} \cos \frac{\Theta}{2} - i \cos \vartheta \sin \frac{\Theta}{2} & -i \sin \frac{\Theta}{2} (\cos \varphi \sin \vartheta - i \sin \varphi \sin \vartheta) \\ -i \sin \frac{\Theta}{2} (\cos \varphi \sin \vartheta + i \sin \varphi \sin \vartheta) & \cos \frac{\Theta}{2} + i \cos \vartheta \sin \frac{\Theta}{2} \end{pmatrix}$$

$$= \mathbf{R}[\varphi\vartheta\hat{\Theta}] = e^{-i\hat{\Theta} \cdot \mathbf{S}\Theta} = e^{-i\mathbf{H}t}$$

$$\begin{aligned} \hat{\Theta}_X &= \cos \varphi \sin \vartheta \\ \hat{\Theta}_Y &= \sin \varphi \sin \vartheta \\ \hat{\Theta}_Z &= \cos \vartheta \end{aligned}$$



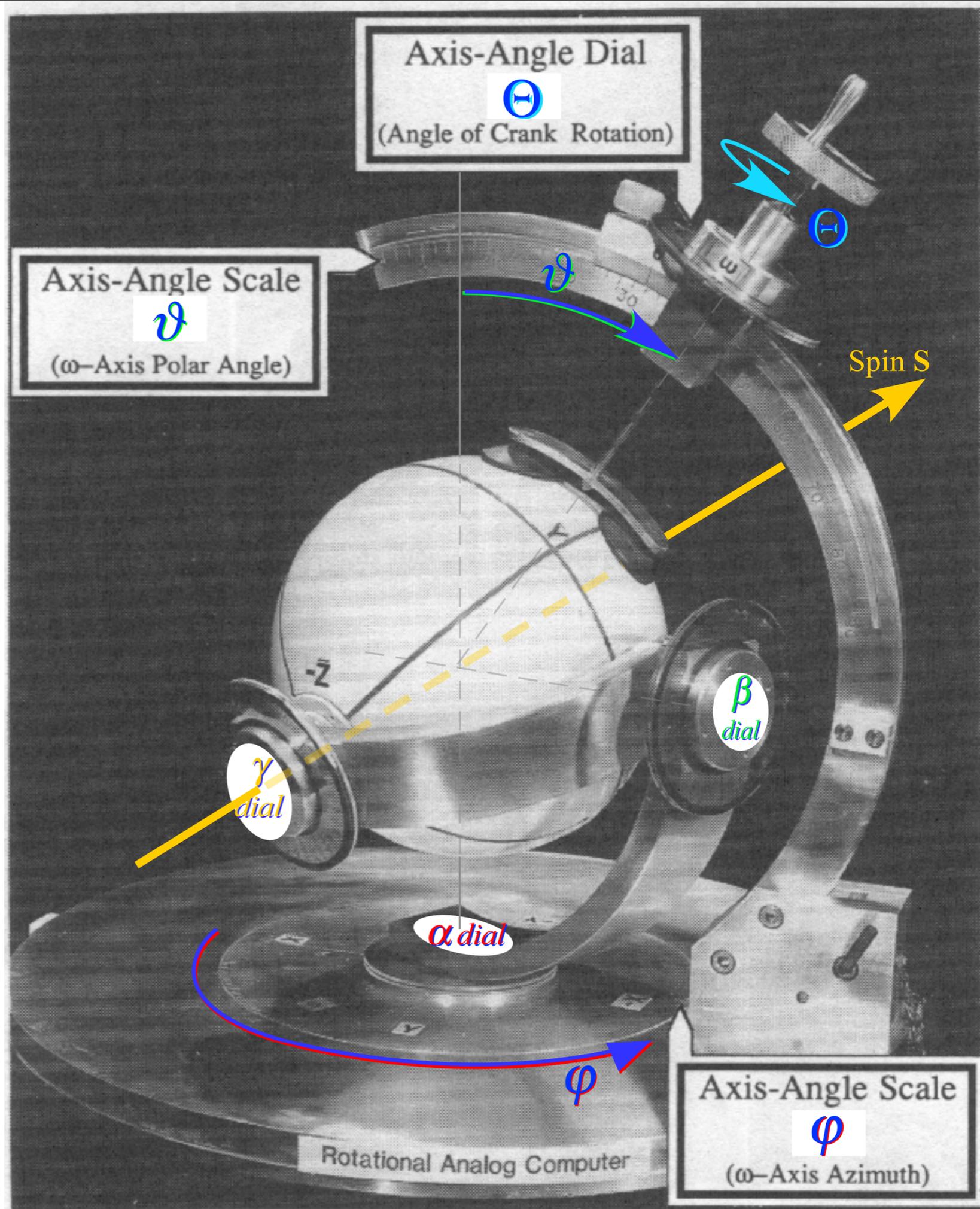
Polar coordinates for unit axis vector  $\hat{\Theta}$

*Polar coordinates  
for unit axis vector  $\hat{\Theta}$*

$$\hat{\Theta}_X = \cos\varphi \sin\vartheta$$

$$\hat{\Theta}_Y = \sin\varphi \sin\vartheta$$

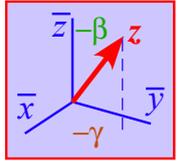
$$\hat{\Theta}_Z = \cos\vartheta$$



Here spin-rotor **S**-polar coordinates are Euler angles

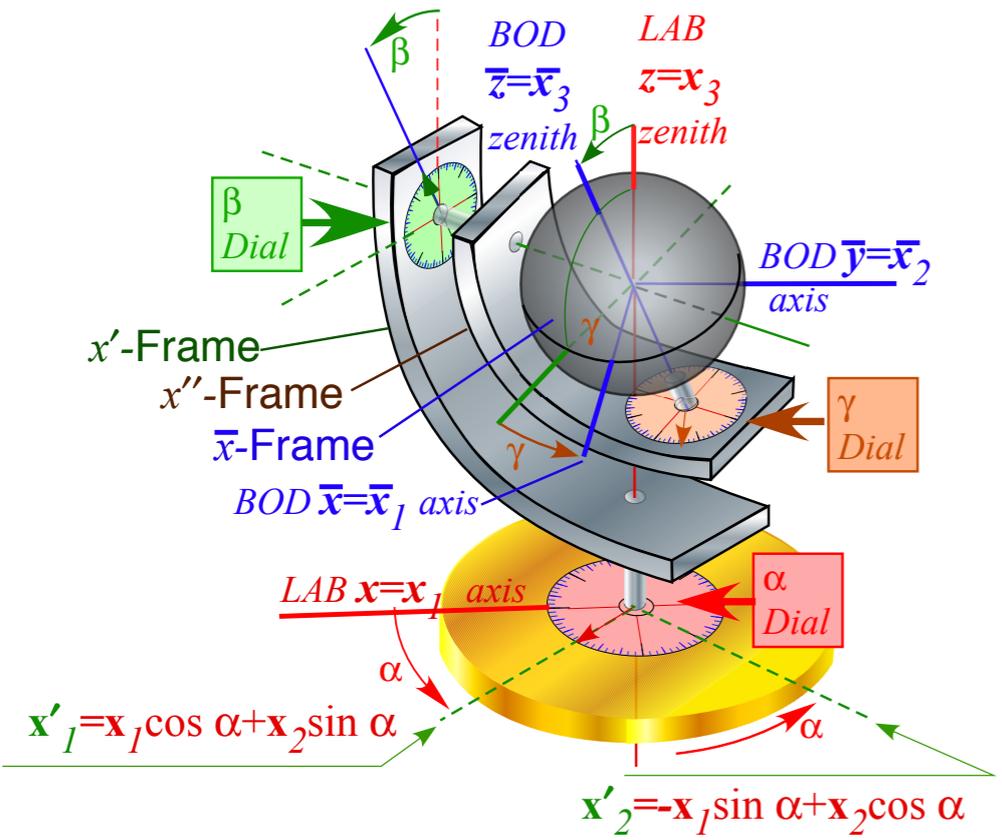
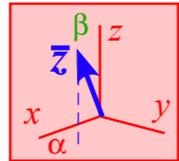
BOD frame view

Polar angles of LAB zenith  $\bar{z}=\bar{x}_3$  are  
(azimuth angle= $-\gamma$ ,  
polar angle= $-\beta$ )



LAB frame view

Polar angles of BOD zenith  $\bar{z}=\bar{x}_3$  are  
(azimuth angle= $\alpha$ ,  
polar angle= $\beta$ )

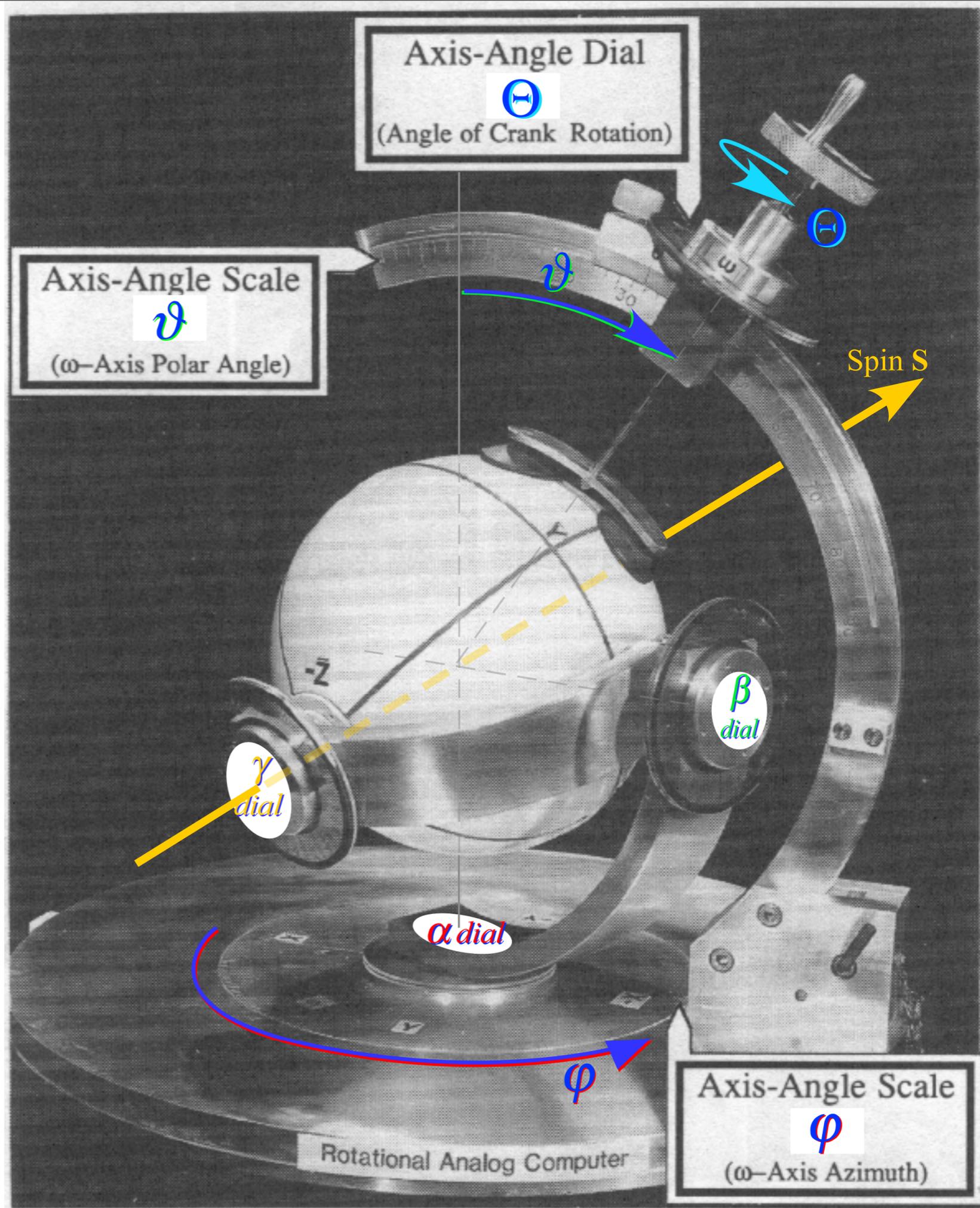


Polar coordinates for unit axis vector  $\hat{\Theta}$

$$\hat{\Theta}_X = \cos \varphi \sin \vartheta$$

$$\hat{\Theta}_Y = \sin \varphi \sin \vartheta$$

$$\hat{\Theta}_Z = \cos \vartheta$$



Review: How “Crazy-Thing”-Theorem makes spinor and vector representation matrices  
Half-angle  $\Theta/2 = \varphi$  replacement and Darboux crank axis operators

Operator-on-Operator transformations

➔ Product algebra for Pauli's  $\sigma_\mu$  and Hamilton's  $q_\mu = -i\sigma_\mu$   
Group product algebra

Jordan-Pauli identity and  $U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - formula

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Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  of group-operators

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Group product geometry

$U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - geometry

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Euler  $\mathbf{R}(\alpha\beta\gamma)$  Sundial

# Operator-on-Operator transformations

*Product algebra* Multiplication rules for Pauli's " $\sigma_\mu$ -quaternions" and Hamilton's  $\mathbf{q}_\mu = -i\sigma_\mu$ .

•	1	$\mathbf{q}_X$	$\mathbf{q}_Y$	$\mathbf{q}_Z$	•	1	$\sigma_X$	$\sigma_Y$	$\sigma_Z$
1	1	$\mathbf{q}_X$	$\mathbf{q}_Y$	$\mathbf{q}_Z$	1	1	$\sigma_X$	$\sigma_Y$	$\sigma_Z$
$\mathbf{q}_X$	$\mathbf{q}_X$	-1	$\mathbf{q}_Z$	$-\mathbf{q}_Y$	$\sigma_X$	$\sigma_X$	1	$i\sigma_Z$	$-i\sigma_Y$
$\mathbf{q}_Y$	$\mathbf{q}_Y$	$-\mathbf{q}_Z$	-1	$\mathbf{q}_X$	$\sigma_Y$	$\sigma_Y$	$-i\sigma_Z$	1	$i\sigma_X$
$\mathbf{q}_Z$	$\mathbf{q}_Z$	$\mathbf{q}_Y$	$-\mathbf{q}_X$	-1	$\sigma_Z$	$\sigma_Z$	$i\sigma_Y$	$-i\sigma_X$	1

# Operator-on-Operator transformations

Product algebra Multiplication rules for Pauli's " $\sigma_\mu$ -quaternions" and Hamilton's  $\mathbf{q}_\mu = -i\sigma_\mu$ .

•	1	$\mathbf{q}_X$	$\mathbf{q}_Y$	$\mathbf{q}_Z$
1	1	$\mathbf{q}_X$	$\mathbf{q}_Y$	$\mathbf{q}_Z$
$\mathbf{q}_X$	$\mathbf{q}_X$	-1	$\mathbf{q}_Z$	$-\mathbf{q}_Y$
$\mathbf{q}_Y$	$\mathbf{q}_Y$	$-\mathbf{q}_Z$	-1	$\mathbf{q}_X$
$\mathbf{q}_Z$	$\mathbf{q}_Z$	$\mathbf{q}_Y$	$-\mathbf{q}_X$	-1

,

•	1	$\sigma_X$	$\sigma_Y$	$\sigma_Z$
1	1	$\sigma_X$	$\sigma_Y$	$\sigma_Z$
$\sigma_X$	$\sigma_X$	1	$i\sigma_Z$	$-i\sigma_Y$
$\sigma_Y$	$\sigma_Y$	$-i\sigma_Z$	1	$i\sigma_X$
$\sigma_Z$	$\sigma_Z$	$i\sigma_Y$	$-i\sigma_X$	1

$$\mathbf{q}_\mu \mathbf{q}_\nu = -\delta_{\mu\nu} \mathbf{1} + \epsilon_{\mu\nu\lambda} \mathbf{q}_\lambda$$

$$\sigma_\mu \sigma_\nu = \delta_{\mu\nu} \mathbf{1} + i \epsilon_{\mu\nu\lambda} \sigma_\lambda$$

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1	1	$\mathbf{q}_X$	$\mathbf{q}_Y$	$\mathbf{q}_Z$
$\mathbf{q}_X$	$\mathbf{q}_X$	-1	$\mathbf{q}_Z$	$-\mathbf{q}_Y$
$\mathbf{q}_Y$	$\mathbf{q}_Y$	$-\mathbf{q}_Z$	-1	$\mathbf{q}_X$
$\mathbf{q}_Z$	$\mathbf{q}_Z$	$\mathbf{q}_Y$	$-\mathbf{q}_X$	-1

,

•	1	$\sigma_X$	$\sigma_Y$	$\sigma_Z$
1	1	$\sigma_X$	$\sigma_Y$	$\sigma_Z$
$\sigma_X$	$\sigma_X$	1	$i\sigma_Z$	$-i\sigma_Y$
$\sigma_Y$	$\sigma_Y$	$-i\sigma_Z$	1	$i\sigma_X$
$\sigma_Z$	$\sigma_Z$	$i\sigma_Y$	$-i\sigma_X$	1

$$\mathbf{q}_\mu \mathbf{q}_\nu = -\delta_{\mu\nu} \mathbf{1} + \epsilon_{\mu\nu\lambda} \mathbf{q}_\lambda$$

$$\sigma_\mu \sigma_\nu = \delta_{\mu\nu} \mathbf{1} + i \epsilon_{\mu\nu\lambda} \sigma_\lambda$$

*Commutation rules for Pauli ops:*  $\sigma_\mu$   
 $\sigma_\mu \sigma_\nu - \sigma_\nu \sigma_\mu = [\sigma_\mu, \sigma_\nu] = 2i \epsilon_{\mu\nu\lambda} \sigma_\lambda$

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1	1	$\mathbf{q}_X$	$\mathbf{q}_Y$	$\mathbf{q}_Z$
$\mathbf{q}_X$	$\mathbf{q}_X$	-1	$\mathbf{q}_Z$	$-\mathbf{q}_Y$
$\mathbf{q}_Y$	$\mathbf{q}_Y$	$-\mathbf{q}_Z$	-1	$\mathbf{q}_X$
$\mathbf{q}_Z$	$\mathbf{q}_Z$	$\mathbf{q}_Y$	$-\mathbf{q}_X$	-1

•	1	$\sigma_X$	$\sigma_Y$	$\sigma_Z$
1	1	$\sigma_X$	$\sigma_Y$	$\sigma_Z$
$\sigma_X$	$\sigma_X$	1	$i\sigma_Z$	$-i\sigma_Y$
$\sigma_Y$	$\sigma_Y$	$-i\sigma_Z$	1	$i\sigma_X$
$\sigma_Z$	$\sigma_Z$	$i\sigma_Y$	$-i\sigma_X$	1

$$\mathbf{q}_\mu \mathbf{q}_\nu = -\delta_{\mu\nu} \mathbf{1} + \epsilon_{\mu\nu\lambda} \mathbf{q}_\lambda$$

$$\sigma_\mu \sigma_\nu = \delta_{\mu\nu} \mathbf{1} + i \epsilon_{\mu\nu\lambda} \sigma_\lambda$$

Commutation rules for Pauli ops:  $\sigma_\mu$

$$\sigma_\mu \sigma_\nu - \sigma_\nu \sigma_\mu = [\sigma_\mu, \sigma_\nu] = 2i \epsilon_{\mu\nu\lambda} \sigma_\lambda$$

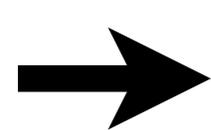
Jordan's spin-ops:  $\mathbf{J}_\mu = \mathbf{S}_\mu = \sigma_\mu/2$ .

$$\mathbf{S}_\mu \mathbf{S}_\nu - \mathbf{S}_\nu \mathbf{S}_\mu = [\mathbf{S}_\mu, \mathbf{S}_\nu] = i \epsilon_{\mu\nu\lambda} \mathbf{S}_\lambda$$

Review: How “Crazy-Thing”-Theorem makes spinor and vector representation matrices  
Half-angle  $\Theta/2 = \varphi$  replacement and Darboux crank axis operators

Operator-on-Operator transformations

Product algebra for Pauli's  $\sigma_\mu$  and Hamilton's  $q_\mu = -i\sigma_\mu$



Group product algebra

Jordan-Pauli identity and  $U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - formula

Transformation  $\mathbf{R}[\Theta]\sigma_\mu\mathbf{R}[\Theta]^\dagger$  of spinor  $\sigma_\mu$ -operators

Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  of group-operators

Operator-on-Operator transformations

Geometry of groups: Hamilton's turns and It's all done with mirrors!

Group product geometry

$U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - geometry

Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  geometry

Euler  $\mathbf{R}(\alpha\beta\gamma)$  versus Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  related to Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  rotation  $\Theta = 0 - 4\pi$ -sequence  $[\varphi\vartheta]$  fixed

$R(3)$ - $U(2)$  slide rule for converting  $\mathbf{R}(\alpha\beta\gamma) \leftrightarrow \mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  Sundial

# Operator-on-Operator transformations

Product algebra Multiplication rules for Pauli's " $\sigma_\mu$ -quaternions" and Hamilton's  $\mathbf{q}_\mu = -i\sigma_\mu$ .

•	1	$\mathbf{q}_X$	$\mathbf{q}_Y$	$\mathbf{q}_Z$	•	1	$\sigma_X$	$\sigma_Y$	$\sigma_Z$
1	1	$\mathbf{q}_X$	$\mathbf{q}_Y$	$\mathbf{q}_Z$	1	1	$\sigma_X$	$\sigma_Y$	$\sigma_Z$
$\mathbf{q}_X$	$\mathbf{q}_X$	-1	$\mathbf{q}_Z$	$-\mathbf{q}_Y$	$\sigma_X$	$\sigma_X$	1	$i\sigma_Z$	$-i\sigma_Y$
$\mathbf{q}_Y$	$\mathbf{q}_Y$	$-\mathbf{q}_Z$	-1	$\mathbf{q}_X$	$\sigma_Y$	$\sigma_Y$	$-i\sigma_Z$	1	$i\sigma_X$
$\mathbf{q}_Z$	$\mathbf{q}_Z$	$\mathbf{q}_Y$	$-\mathbf{q}_X$	-1	$\sigma_Z$	$\sigma_Z$	$i\sigma_Y$	$-i\sigma_X$	1

$$\mathbf{q}_\mu \mathbf{q}_\nu = -\delta_{\mu\nu} \mathbf{1} + \epsilon_{\mu\nu\lambda} \mathbf{q}_\lambda$$

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 $\sigma_\mu \sigma_\nu - \sigma_\nu \sigma_\mu = [\sigma_\mu, \sigma_\nu] = 2i \epsilon_{\mu\nu\lambda} \sigma_\lambda$

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 $\mathbf{S}_\mu \mathbf{S}_\nu - \mathbf{S}_\nu \mathbf{S}_\mu = [\mathbf{S}_\mu, \mathbf{S}_\nu] = i \epsilon_{\mu\nu\lambda} \mathbf{S}_\lambda$

Group products

$$\mathbf{R}[\vec{\Theta}'] \mathbf{R}[\vec{\Theta}] = \left( \cos \frac{\Theta'}{2} \mathbf{1} - i \sin \frac{\Theta'}{2} \hat{\Theta}' \cdot \boldsymbol{\sigma} \right) \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta} \cdot \boldsymbol{\sigma} \right)$$

# Operator-on-Operator transformations

Product algebra Multiplication rules for Pauli's " $\sigma_\mu$ -quaternions" and Hamilton's  $\mathbf{q}_\mu = -i\sigma_\mu$ .

•	1	$\mathbf{q}_X$	$\mathbf{q}_Y$	$\mathbf{q}_Z$	,	•	1	$\sigma_X$	$\sigma_Y$	$\sigma_Z$
1	1	$\mathbf{q}_X$	$\mathbf{q}_Y$	$\mathbf{q}_Z$		1	1	$\sigma_X$	$\sigma_Y$	$\sigma_Z$
$\mathbf{q}_X$	$\mathbf{q}_X$	-1	$\mathbf{q}_Z$	$-\mathbf{q}_Y$		$\sigma_X$	$\sigma_X$	1	$i\sigma_Z$	$-i\sigma_Y$
$\mathbf{q}_Y$	$\mathbf{q}_Y$	$-\mathbf{q}_Z$	-1	$\mathbf{q}_X$		$\sigma_Y$	$\sigma_Y$	$-i\sigma_Z$	1	$i\sigma_X$
$\mathbf{q}_Z$	$\mathbf{q}_Z$	$\mathbf{q}_Y$	$-\mathbf{q}_X$	-1		$\sigma_Z$	$\sigma_Z$	$i\sigma_Y$	$-i\sigma_X$	1

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## Group products

$$\mathbf{R}[\vec{\Theta}'] \mathbf{R}[\vec{\Theta}] = \left( \cos \frac{\Theta'}{2} \mathbf{1} - i \sin \frac{\Theta'}{2} \hat{\Theta}' \cdot \boldsymbol{\sigma} \right) \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta} \cdot \boldsymbol{\sigma} \right)$$

$$= \cos \frac{\Theta'}{2} \cos \frac{\Theta}{2} \mathbf{1} - i \left[ \cos \frac{\Theta'}{2} \sin \frac{\Theta}{2} \hat{\Theta} + \cos \frac{\Theta}{2} \sin \frac{\Theta'}{2} \hat{\Theta}' \right] \cdot \boldsymbol{\sigma} - \sin \frac{\Theta'}{2} \sin \frac{\Theta}{2} (\hat{\Theta}' \cdot \boldsymbol{\sigma})(\hat{\Theta} \cdot \boldsymbol{\sigma})$$

# Operator-on-Operator transformations

Product algebra Multiplication rules for Pauli's " $\sigma_\mu$ -quaternions" and Hamilton's  $\mathbf{q}_\mu = -i\sigma_\mu$ .

•	1	$\mathbf{q}_X$	$\mathbf{q}_Y$	$\mathbf{q}_Z$		•	1	$\sigma_X$	$\sigma_Y$	$\sigma_Z$
1	1	$\mathbf{q}_X$	$\mathbf{q}_Y$	$\mathbf{q}_Z$	,	1	1	$\sigma_X$	$\sigma_Y$	$\sigma_Z$
$\mathbf{q}_X$	$\mathbf{q}_X$	-1	$\mathbf{q}_Z$	$-\mathbf{q}_Y$		$\sigma_X$	$\sigma_X$	1	$i\sigma_Z$	$-i\sigma_Y$
$\mathbf{q}_Y$	$\mathbf{q}_Y$	$-\mathbf{q}_Z$	-1	$\mathbf{q}_X$		$\sigma_Y$	$\sigma_Y$	$-i\sigma_Z$	1	$i\sigma_X$
$\mathbf{q}_Z$	$\mathbf{q}_Z$	$\mathbf{q}_Y$	$-\mathbf{q}_X$	-1		$\sigma_Z$	$\sigma_Z$	$i\sigma_Y$	$-i\sigma_X$	1

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$$\mathbf{R}[\vec{\Theta}'] \mathbf{R}[\vec{\Theta}] = \left( \cos \frac{\Theta'}{2} \mathbf{1} - i \sin \frac{\Theta'}{2} \hat{\Theta}' \cdot \boldsymbol{\sigma} \right) \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta} \cdot \boldsymbol{\sigma} \right)$$

$$= \cos \frac{\Theta'}{2} \cos \frac{\Theta}{2} \mathbf{1} - i \left[ \cos \frac{\Theta'}{2} \sin \frac{\Theta}{2} \hat{\Theta} + \cos \frac{\Theta}{2} \sin \frac{\Theta'}{2} \hat{\Theta}' \right] \cdot \boldsymbol{\sigma} - \sin \frac{\Theta'}{2} \sin \frac{\Theta}{2} (\hat{\Theta}' \cdot \boldsymbol{\sigma})(\hat{\Theta} \cdot \boldsymbol{\sigma})$$

Jordan-Pauli identity is used to reduce  $(\hat{\Theta}' \cdot \boldsymbol{\sigma})(\hat{\Theta} \cdot \boldsymbol{\sigma})$  to  $(\hat{\Theta}' \cdot \hat{\Theta}) \mathbf{1} + (\hat{\Theta}' \times \hat{\Theta}) \cdot \boldsymbol{\sigma}$

$$\mathbf{R}[\vec{\Theta}'] \mathbf{R}[\vec{\Theta}] = \left( \cos \frac{\Theta''}{2} \right) \mathbf{1} - i \left[ \sin \frac{\Theta''}{2} \hat{\Theta}'' \right] \cdot \boldsymbol{\sigma} = \mathbf{R}[\vec{\Theta}'']$$

$$= \left( \cos \frac{\Theta'}{2} \cos \frac{\Theta}{2} - \sin \frac{\Theta'}{2} \sin \frac{\Theta}{2} \right) \mathbf{1} - i \left[ \left[ \cos \frac{\Theta'}{2} \sin \frac{\Theta}{2} \hat{\Theta} + \cos \frac{\Theta}{2} \sin \frac{\Theta'}{2} \hat{\Theta}' \right] + \sin \frac{\Theta'}{2} \sin \frac{\Theta}{2} \hat{\Theta}' \times \hat{\Theta} \right] \cdot \boldsymbol{\sigma}$$

Review: How “Crazy-Thing”-Theorem makes spinor and vector representation matrices  
Half-angle  $\Theta/2 = \varphi$  replacement and Darboux crank axis operators

Operator-on-Operator transformations

Product algebra for Pauli's  $\sigma_\mu$  and Hamilton's  $q_\mu = -i\sigma_\mu$

Group product algebra

➔ Jordan-Pauli identity and  $U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - formula  
Transformation  $\mathbf{R}[\Theta]\sigma_\mu\mathbf{R}[\Theta]^\dagger$  of spinor  $\sigma_\mu$ -operators  
Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  of group-operators  
Operator-on-Operator transformations

Geometry of groups: Hamilton's turns and It's all done with mirrors!

Group product geometry

$U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - geometry  
Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  geometry

Euler  $\mathbf{R}(\alpha\beta\gamma)$  versus Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  related to Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  rotation  $\Theta = 0 - 4\pi$ -sequence  $[\varphi\vartheta]$  fixed

$R(3)$ - $U(2)$  slide rule for converting  $\mathbf{R}(\alpha\beta\gamma) \leftrightarrow \mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  Sundial

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1	1	$\mathbf{q}_X$	$\mathbf{q}_Y$	$\mathbf{q}_Z$	,	1	1	$\sigma_X$	$\sigma_Y$	$\sigma_Z$
$\mathbf{q}_X$	$\mathbf{q}_X$	-1	$\mathbf{q}_Z$	$-\mathbf{q}_Y$		$\sigma_X$	$\sigma_X$	1	$i\sigma_Z$	$-i\sigma_Y$
$\mathbf{q}_Y$	$\mathbf{q}_Y$	$-\mathbf{q}_Z$	-1	$\mathbf{q}_X$		$\sigma_Y$	$\sigma_Y$	$-i\sigma_Z$	1	$i\sigma_X$
$\mathbf{q}_Z$	$\mathbf{q}_Z$	$\mathbf{q}_Y$	$-\mathbf{q}_X$	-1		$\sigma_Z$	$\sigma_Z$	$i\sigma_Y$	$-i\sigma_X$	1

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 $\sigma_\mu \sigma_\nu - \sigma_\nu \sigma_\mu = [\sigma_\mu, \sigma_\nu] = 2i \epsilon_{\mu\nu\lambda} \sigma_\lambda$

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 $\mathbf{S}_\mu \mathbf{S}_\nu - \mathbf{S}_\nu \mathbf{S}_\mu = [\mathbf{S}_\mu, \mathbf{S}_\nu] = i \epsilon_{\mu\nu\lambda} \mathbf{S}_\lambda$

## Group products

$$\mathbf{R}[\vec{\Theta}'] \mathbf{R}[\vec{\Theta}] = \left( \cos \frac{\Theta'}{2} \mathbf{1} - i \sin \frac{\Theta'}{2} \hat{\Theta}' \cdot \boldsymbol{\sigma} \right) \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta} \cdot \boldsymbol{\sigma} \right)$$

$$= \cos \frac{\Theta'}{2} \cos \frac{\Theta}{2} \mathbf{1} - i \left[ \cos \frac{\Theta'}{2} \sin \frac{\Theta}{2} \hat{\Theta} + \cos \frac{\Theta}{2} \sin \frac{\Theta'}{2} \hat{\Theta}' \right] \cdot \boldsymbol{\sigma} - \sin \frac{\Theta'}{2} \sin \frac{\Theta}{2} (\hat{\Theta}' \cdot \boldsymbol{\sigma})(\hat{\Theta} \cdot \boldsymbol{\sigma})$$

Jordan-Pauli identity is used to reduce  $(\hat{\Theta}' \cdot \boldsymbol{\sigma})(\hat{\Theta} \cdot \boldsymbol{\sigma})$  to  $(\hat{\Theta}' \cdot \hat{\Theta}) \mathbf{1} + (\hat{\Theta}' \times \hat{\Theta}) \cdot \boldsymbol{\sigma}$

$$\mathbf{R}[\vec{\Theta}'] \mathbf{R}[\vec{\Theta}] = \left( \cos \frac{\Theta''}{2} \right) \mathbf{1} - i \left[ \sin \frac{\Theta''}{2} \hat{\Theta}'' \right] \cdot \boldsymbol{\sigma} = \mathbf{R}[\vec{\Theta}'']$$

$$= \left( \cos \frac{\Theta'}{2} \cos \frac{\Theta}{2} - \sin \frac{\Theta'}{2} \sin \frac{\Theta}{2} \hat{\Theta}' \cdot \hat{\Theta} \right) \mathbf{1} - i \left[ \left[ \cos \frac{\Theta'}{2} \sin \frac{\Theta}{2} \hat{\Theta} + \cos \frac{\Theta}{2} \sin \frac{\Theta'}{2} \hat{\Theta}' \right] + \sin \frac{\Theta'}{2} \sin \frac{\Theta}{2} \hat{\Theta}' \times \hat{\Theta} \right] \cdot \boldsymbol{\sigma}$$

Now easy to solve for the new *product angle*  $\Theta''$ .

$$\left( \cos \frac{\Theta''}{2} \right) = \left( \cos \frac{\Theta'}{2} \cos \frac{\Theta}{2} - \sin \frac{\Theta'}{2} \sin \frac{\Theta}{2} \hat{\Theta}' \cdot \hat{\Theta} \right)$$

*U(2) and R(3) Group Product Formulae*

# Operator-on-Operator transformations

*Product algebra* Multiplication rules for Pauli's " $\sigma_\mu$ -quaternions" and Hamilton's  $\mathbf{q}_\mu = -i\sigma_\mu$ .

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1	1	$\mathbf{q}_X$	$\mathbf{q}_Y$	$\mathbf{q}_Z$	,	1	1	$\sigma_X$	$\sigma_Y$	$\sigma_Z$
$\mathbf{q}_X$	$\mathbf{q}_X$	-1	$\mathbf{q}_Z$	$-\mathbf{q}_Y$		$\sigma_X$	$\sigma_X$	1	$i\sigma_Z$	$-i\sigma_Y$
$\mathbf{q}_Y$	$\mathbf{q}_Y$	$-\mathbf{q}_Z$	-1	$\mathbf{q}_X$		$\sigma_Y$	$\sigma_Y$	$-i\sigma_Z$	1	$i\sigma_X$
$\mathbf{q}_Z$	$\mathbf{q}_Z$	$\mathbf{q}_Y$	$-\mathbf{q}_X$	-1		$\sigma_Z$	$\sigma_Z$	$i\sigma_Y$	$-i\sigma_X$	1

$$\mathbf{q}_\mu \mathbf{q}_\nu = -\delta_{\mu\nu} \mathbf{1} + \epsilon_{\mu\nu\lambda} \mathbf{q}_\lambda$$

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 $\mathbf{S}_\mu \mathbf{S}_\nu - \mathbf{S}_\nu \mathbf{S}_\mu = [\mathbf{S}_\mu, \mathbf{S}_\nu] = i \epsilon_{\mu\nu\lambda} \mathbf{S}_\lambda$

## Group products

$$\mathbf{R}[\vec{\Theta}'] \mathbf{R}[\vec{\Theta}] = \left( \cos \frac{\Theta'}{2} \mathbf{1} - i \sin \frac{\Theta'}{2} \hat{\Theta}' \cdot \boldsymbol{\sigma} \right) \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta} \cdot \boldsymbol{\sigma} \right)$$

$$= \cos \frac{\Theta'}{2} \cos \frac{\Theta}{2} \mathbf{1} - i \left[ \cos \frac{\Theta'}{2} \sin \frac{\Theta}{2} \hat{\Theta} + \cos \frac{\Theta}{2} \sin \frac{\Theta'}{2} \hat{\Theta}' \right] \cdot \boldsymbol{\sigma} - \sin \frac{\Theta'}{2} \sin \frac{\Theta}{2} (\hat{\Theta}' \cdot \boldsymbol{\sigma})(\hat{\Theta} \cdot \boldsymbol{\sigma})$$

Jordan-Pauli identity is used to reduce  $(\hat{\Theta}' \cdot \boldsymbol{\sigma})(\hat{\Theta} \cdot \boldsymbol{\sigma})$  to  $(\hat{\Theta}' \cdot \hat{\Theta}) \mathbf{1} + (\hat{\Theta}' \times \hat{\Theta}) \cdot \boldsymbol{\sigma}$

$$\mathbf{R}[\vec{\Theta}'] \mathbf{R}[\vec{\Theta}] = \left( \cos \frac{\Theta''}{2} \right) \mathbf{1} - i \left[ \sin \frac{\Theta''}{2} \hat{\Theta}'' \right] \cdot \boldsymbol{\sigma} = \mathbf{R}[\vec{\Theta}'']$$

$$= \left( \cos \frac{\Theta'}{2} \cos \frac{\Theta}{2} - \sin \frac{\Theta'}{2} \sin \frac{\Theta}{2} \hat{\Theta}' \cdot \hat{\Theta} \right) \mathbf{1} - i \left[ \left[ \cos \frac{\Theta'}{2} \sin \frac{\Theta}{2} \hat{\Theta} + \cos \frac{\Theta}{2} \sin \frac{\Theta'}{2} \hat{\Theta}' \right] + \sin \frac{\Theta'}{2} \sin \frac{\Theta}{2} \hat{\Theta}' \times \hat{\Theta} \right] \cdot \boldsymbol{\sigma}$$

Now easy to solve for the new *product angle*  $\Theta''$  and new *crank unit vector*  $\hat{\Theta}''$ .

$$\left( \cos \frac{\Theta''}{2} \right) = \left( \cos \frac{\Theta'}{2} \cos \frac{\Theta}{2} - \sin \frac{\Theta'}{2} \sin \frac{\Theta}{2} \hat{\Theta}' \cdot \hat{\Theta} \right) \quad U(2) \text{ and } R(3) \text{ Group Product Formulae}$$

$$\left[ \sin \frac{\Theta''}{2} \hat{\Theta}'' \right] = \left[ \cos \frac{\Theta'}{2} \sin \frac{\Theta}{2} \hat{\Theta} + \cos \frac{\Theta}{2} \sin \frac{\Theta'}{2} \hat{\Theta}' + \sin \frac{\Theta'}{2} \sin \frac{\Theta}{2} \hat{\Theta}' \times \hat{\Theta} \right]$$

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Product algebra for Pauli's  $\sigma_\mu$  and Hamilton's  $q_\mu = -i\sigma_\mu$

Group product algebra

→ Jordan-Pauli identity and  $U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - formula  
Transformation  $\mathbf{R}[\Theta]\sigma_\mu\mathbf{R}[\Theta]^\dagger$  of spinor  $\sigma_\mu$ -operators  
Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  of group-operators  
Operator-on-Operator transformations

Geometry of groups: Hamilton's turns and It's all done with mirrors!

Group product geometry

$U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - geometry  
Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  geometry

Euler  $\mathbf{R}(\alpha\beta\gamma)$  versus Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  related to Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

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$R(3)$ - $U(2)$  slide rule for converting  $\mathbf{R}(\alpha\beta\gamma) \leftrightarrow \mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  Sundial

*Transformation of spinor  $\sigma_\mu$ -operators*

$$\mathbf{R}[\vec{\Theta}] \sigma_L \mathbf{R}[\vec{\Theta}]^\dagger = \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta} \cdot \boldsymbol{\sigma} \right) \sigma_L \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta} \cdot \boldsymbol{\sigma} \right)^\dagger$$

*Transformation of spinor  $\sigma_\mu$ -operators*

$$\begin{aligned} \mathbf{R}[\vec{\Theta}] \sigma_L \mathbf{R}[\vec{\Theta}]^\dagger &= \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta} \cdot \boldsymbol{\sigma} \right) \sigma_L \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta} \cdot \boldsymbol{\sigma} \right)^\dagger \\ &= \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta}_K \sigma_K \right) \sigma_L \left( \cos \frac{\Theta}{2} \mathbf{1} + i \sin \frac{\Theta}{2} \hat{\Theta}_M \sigma_M \right) \end{aligned}$$

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 &= \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta}_K \sigma_K \right) \sigma_L \left( \cos \frac{\Theta}{2} \mathbf{1} + i \sin \frac{\Theta}{2} \hat{\Theta}_M \sigma_M \right) \\
 &= \sigma'_L = \sigma_L \cos \Theta - \varepsilon_{LKN} \hat{\Theta}_K \sigma_N \sin \Theta + (1 - \cos \Theta) \hat{\Theta}_L (\hat{\Theta}_M \sigma_M)
 \end{aligned}$$

*(Left as an exercise)*

Review: How “Crazy-Thing”-Theorem makes spinor and vector representation matrices  
Half-angle  $\Theta/2 = \varphi$  replacement and Darboux crank axis operators

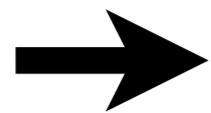
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 &= \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta}_K \sigma_K \right) \sigma_L \left( \cos \frac{\Theta}{2} \mathbf{1} + i \sin \frac{\Theta}{2} \hat{\Theta}_M \sigma_M \right) \\
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 \end{aligned}$$

*General transformation of rotational  $\mathbf{R}[\vec{\Theta}']$ -operators*

$$\begin{aligned}
 \mathbf{R}[\vec{\Theta}] \mathbf{R}[\vec{\Theta}'] \mathbf{R}[\vec{\Theta}]^\dagger &= \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta} \cdot \boldsymbol{\sigma} \right) \mathbf{R}[\vec{\Theta}'] \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta} \cdot \boldsymbol{\sigma} \right)^\dagger \\
 \mathbf{R}[\mathbf{R}[\vec{\Theta}] \cdot \vec{\Theta}'] &= \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta}_K \sigma_K \right) \mathbf{R}[\vec{\Theta}'] \left( \cos \frac{\Theta}{2} \mathbf{1} + i \sin \frac{\Theta}{2} \hat{\Theta}_M \sigma_M \right)
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*Transformation of spinor  $\sigma_\mu$ -operators*

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 \mathbf{R}[\vec{\Theta}] \sigma_L \mathbf{R}[\vec{\Theta}]^\dagger &= \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta} \cdot \boldsymbol{\sigma} \right) \sigma_L \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta} \cdot \boldsymbol{\sigma} \right)^\dagger \\
 &= \left( \cos \frac{\Theta}{2} \mathbf{1} - i \sin \frac{\Theta}{2} \hat{\Theta}_K \sigma_K \right) \sigma_L \left( \cos \frac{\Theta}{2} \mathbf{1} + i \sin \frac{\Theta}{2} \hat{\Theta}_M \sigma_M \right) \\
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 \end{aligned}$$

*This one is better seen geometrically. Algebra not so quick.*

Review: How “Crazy-Thing”-Theorem makes spinor and vector representation matrices  
Half-angle  $\Theta/2 = \varphi$  replacement and Darboux crank axis operators

Operator-on-Operator transformations

Product algebra for Pauli's  $\sigma_\mu$  and Hamilton's  $q_\mu = -i\sigma_\mu$

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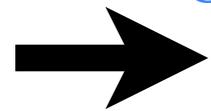
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Geometry of groups: Hamilton's turns and It's all done with mirrors!



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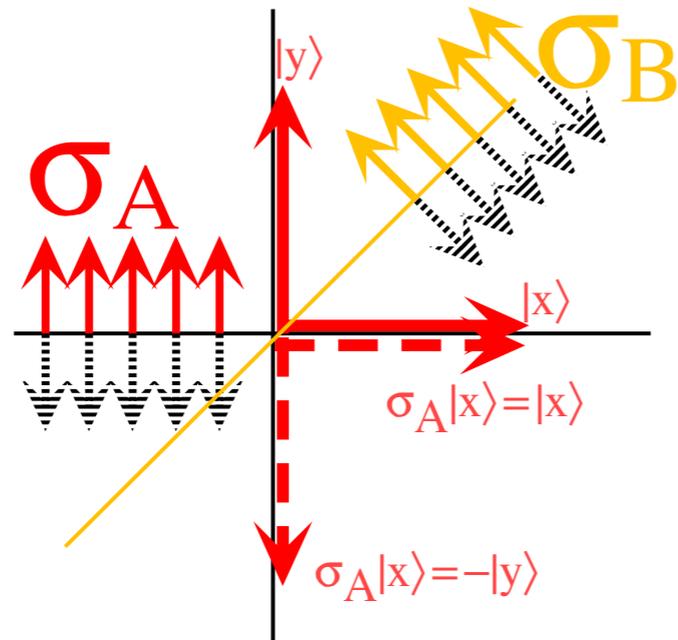
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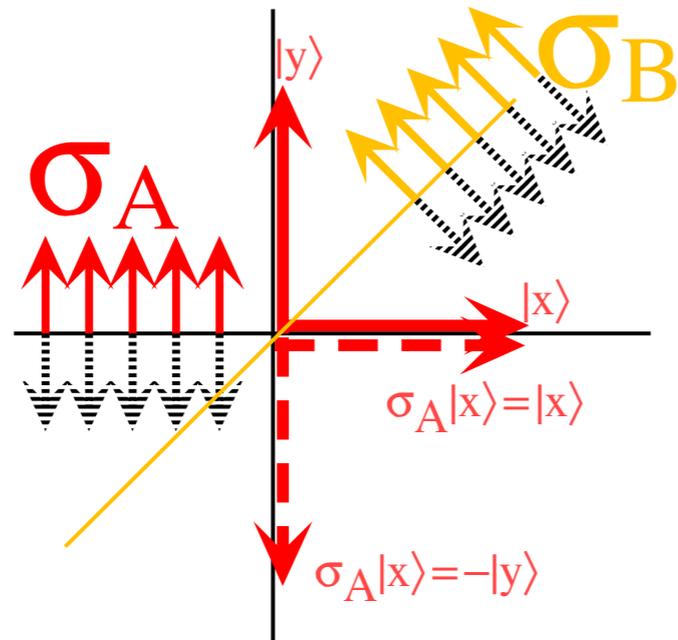
# Geometry of $U(2)$ transformations. It's all done with mirrors!

$$\sigma_A = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad \sigma_B = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$



# Geometry of $U(2)$ transformations. It's all done with mirrors!

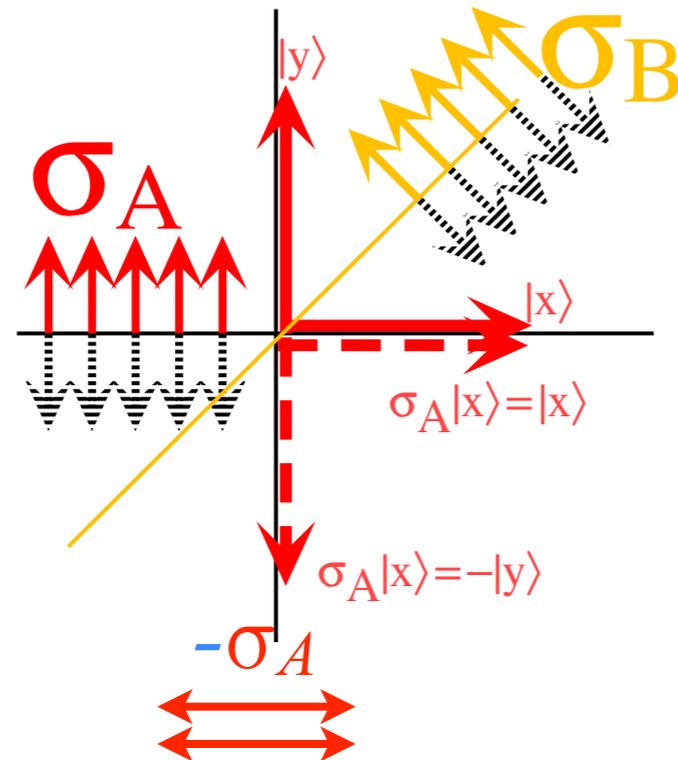
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Note that  $-\sigma_A$  is a  $y$ -plane mirror

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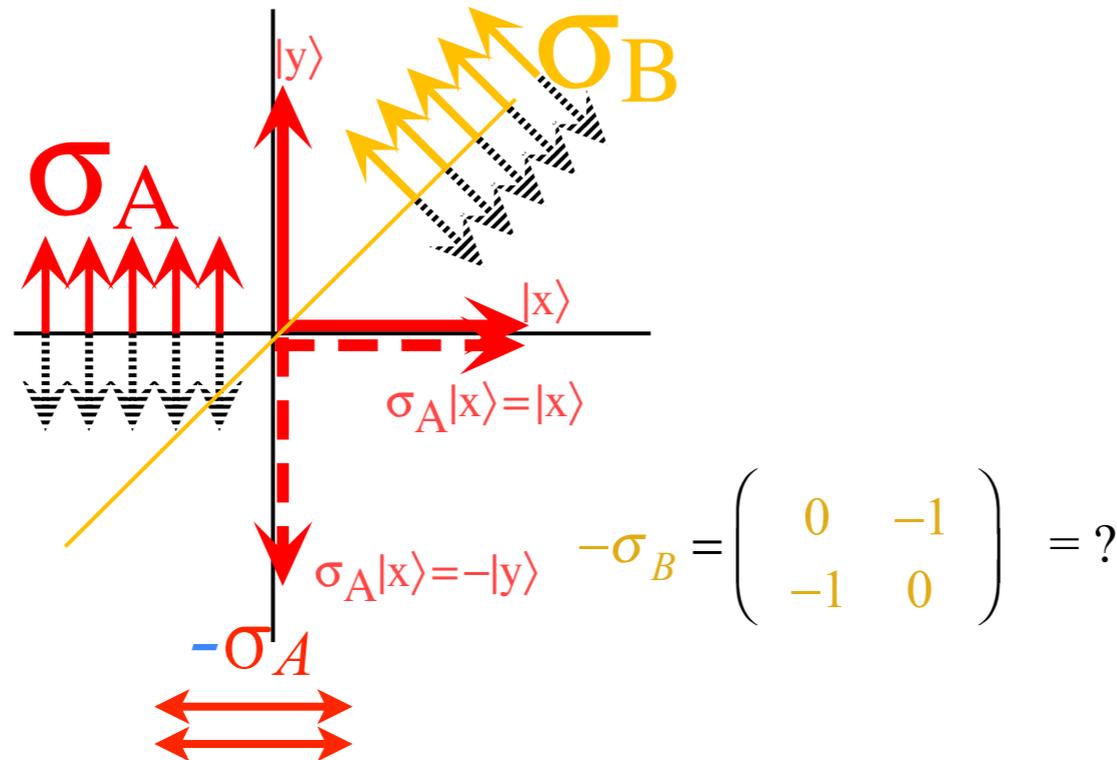


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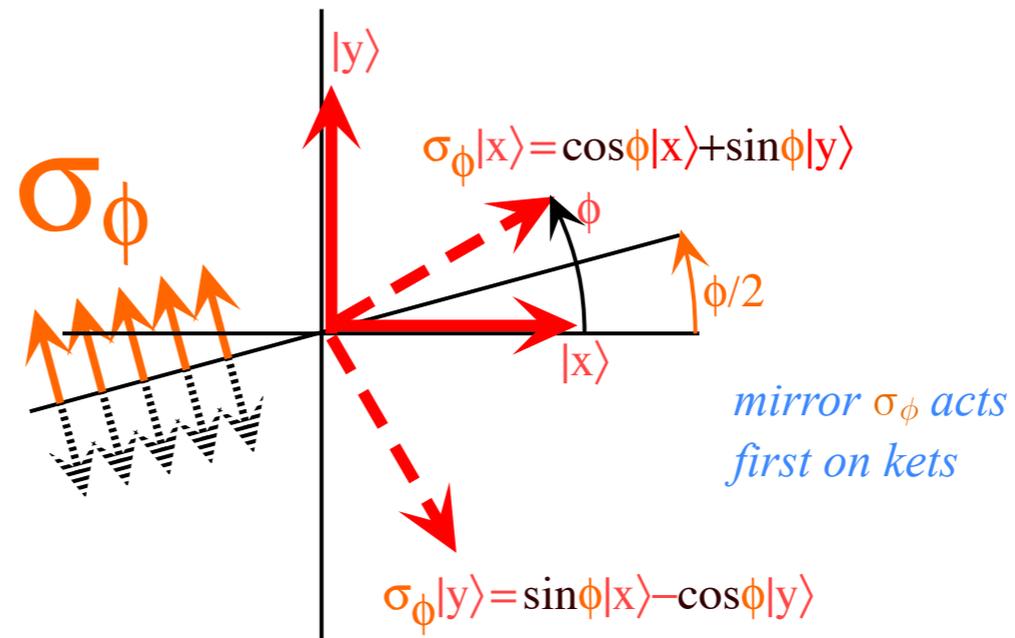
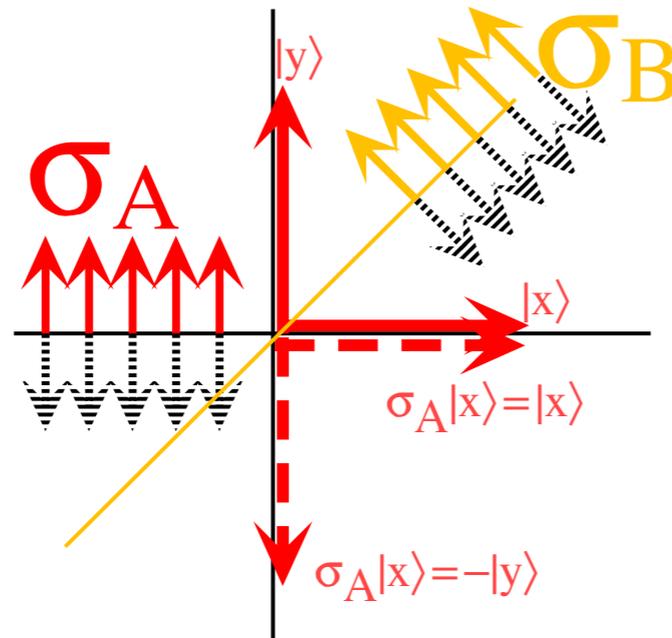


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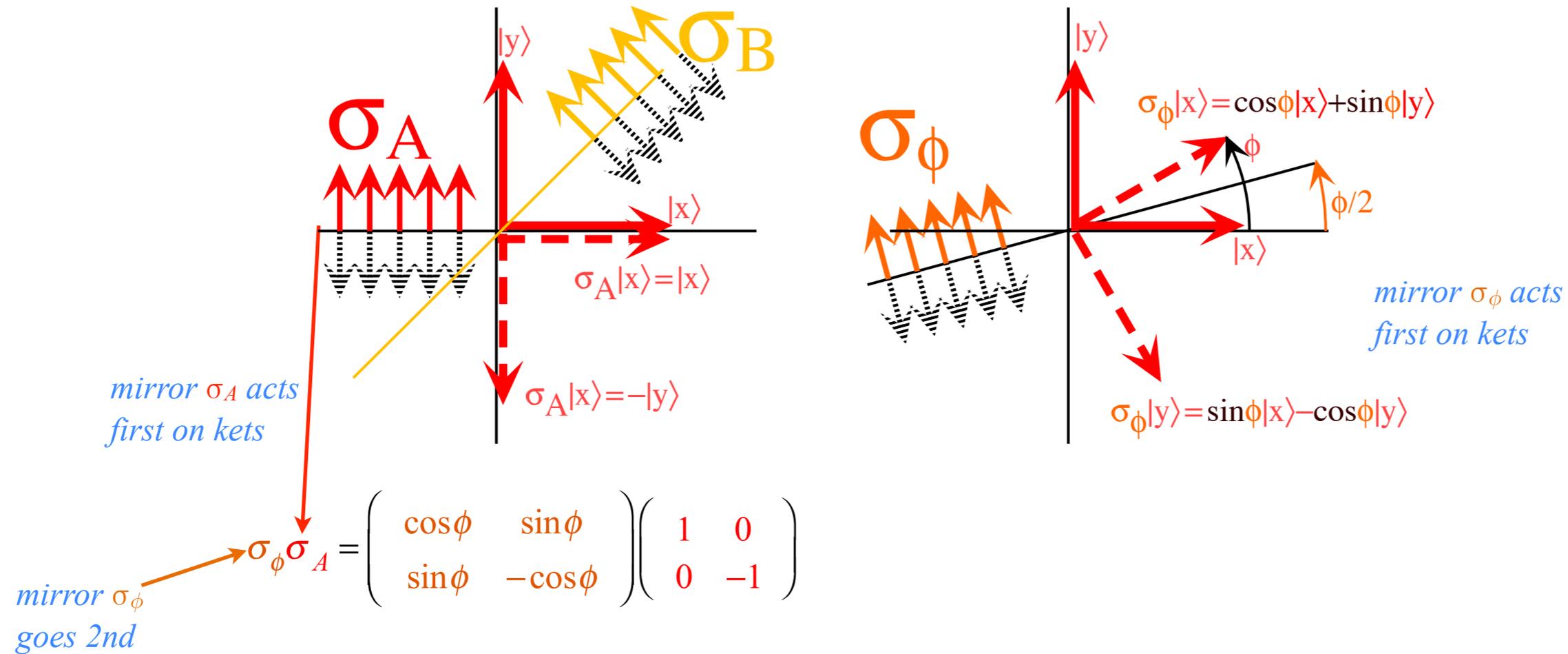
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$$\sigma_A = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}, \quad \sigma_B = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_\phi = \begin{pmatrix} \cos\phi & \sin\phi \\ \sin\phi & -\cos\phi \end{pmatrix} = \sigma_A \cos\phi + \sigma_B \sin\phi$$



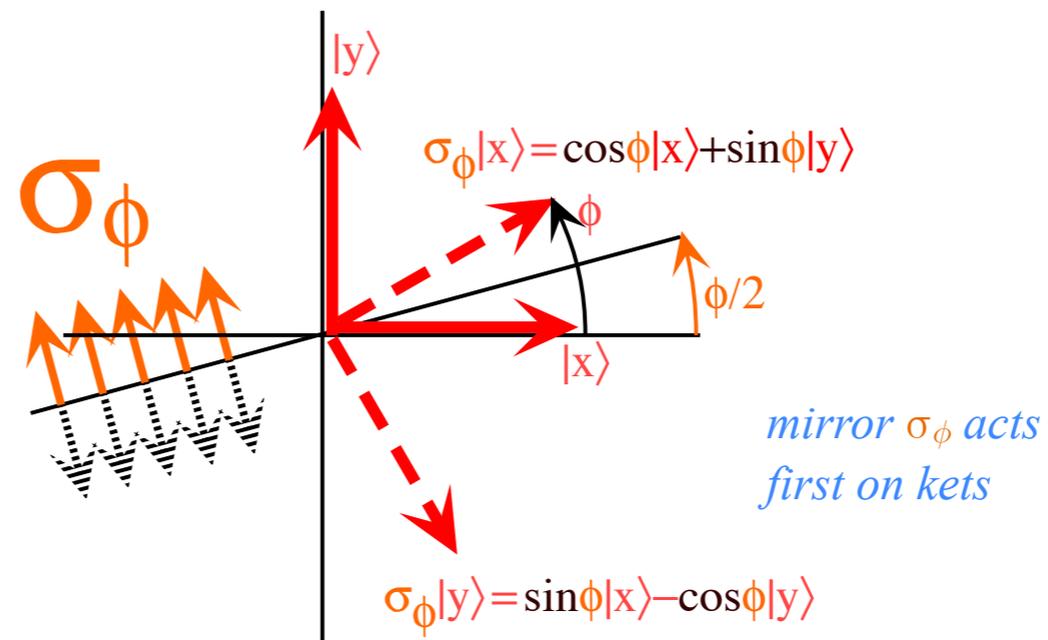
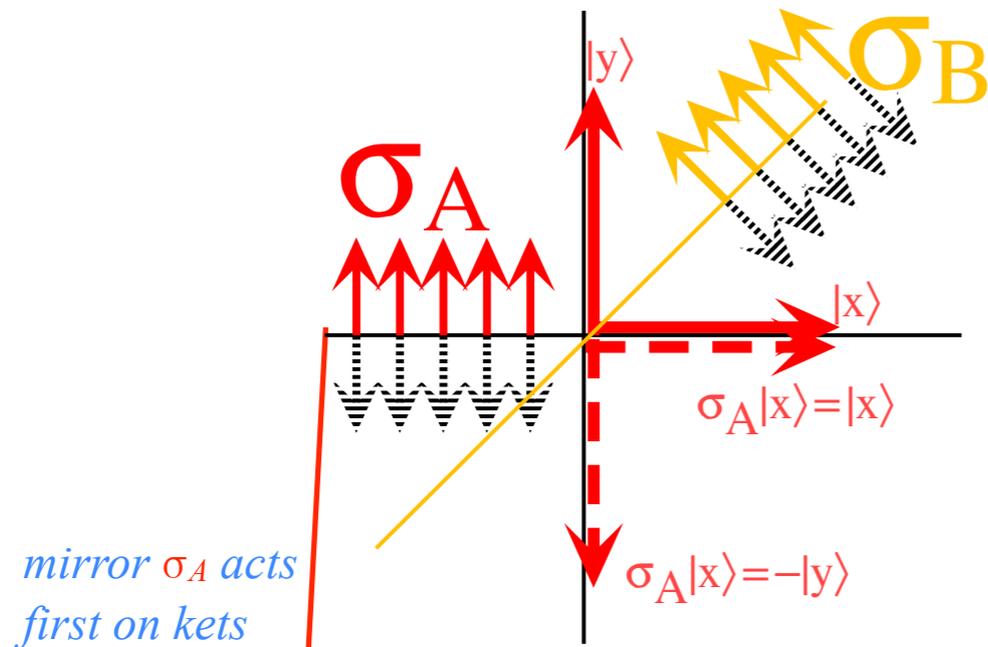
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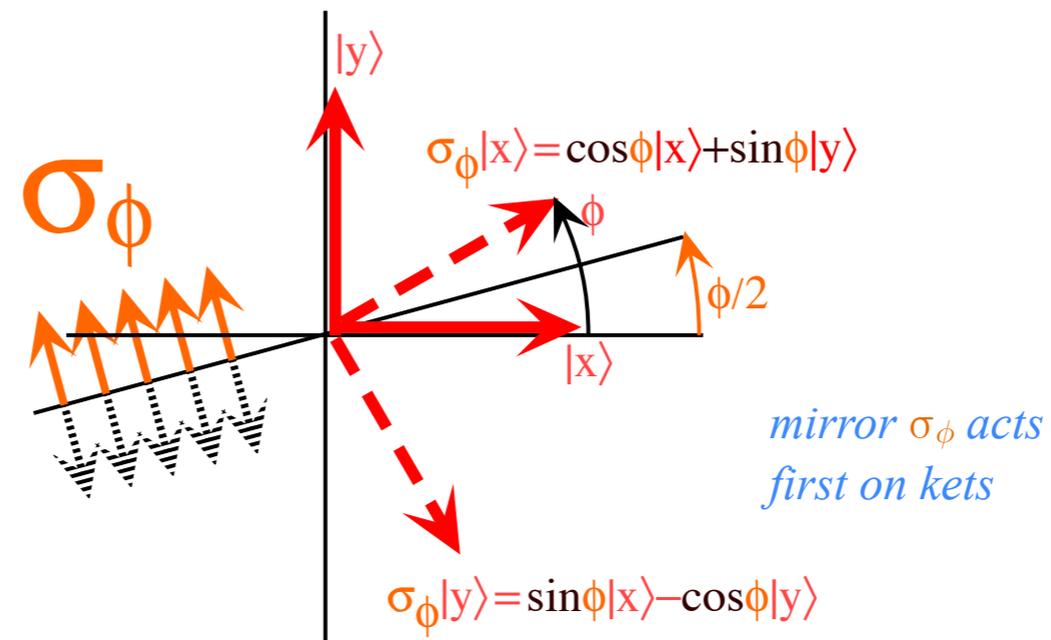
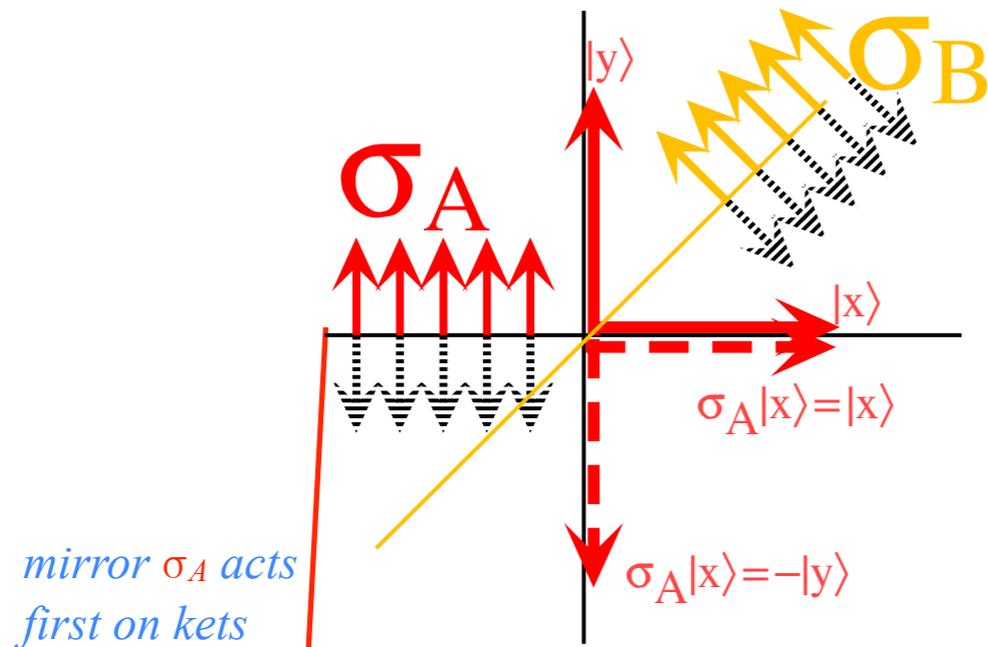


*mirror  $\sigma_\phi$  goes 2nd*

$$\sigma_\phi \sigma_A = \begin{pmatrix} \cos\phi & \sin\phi \\ \sin\phi & -\cos\phi \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} \cos\phi & -\sin\phi \\ \sin\phi & \cos\phi \end{pmatrix} = R[\phi],$$

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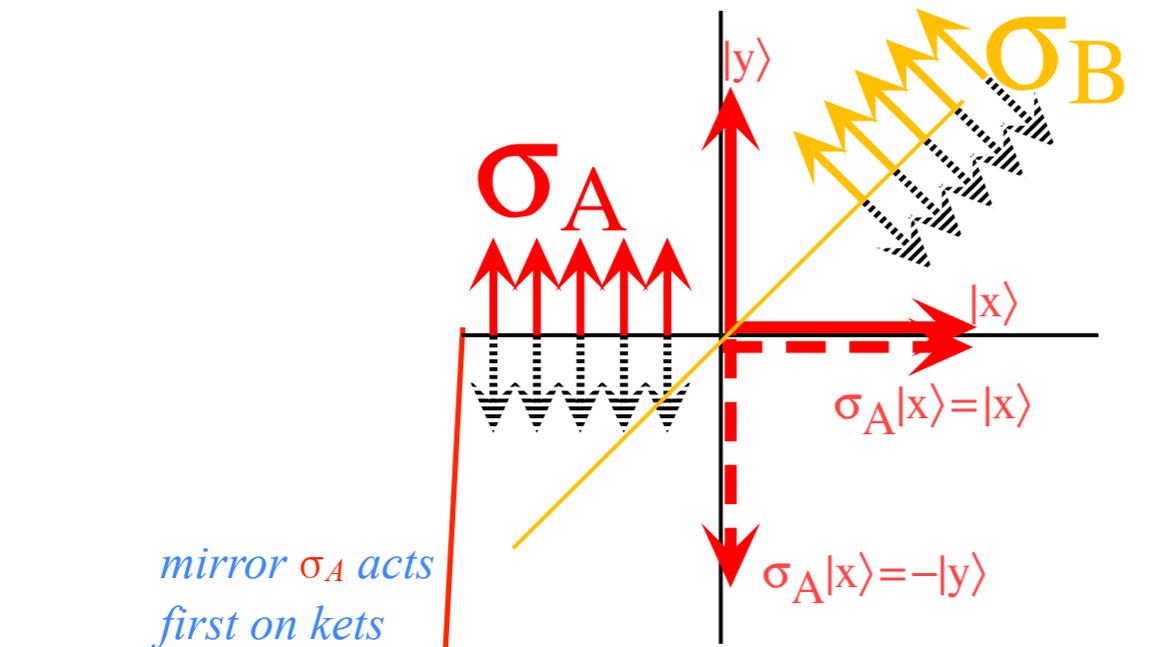
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*Rotation angle  $\phi$  is TWICE the angle  $\phi/2$  between mirror  $\sigma_A$  and mirror  $\sigma_\phi$*

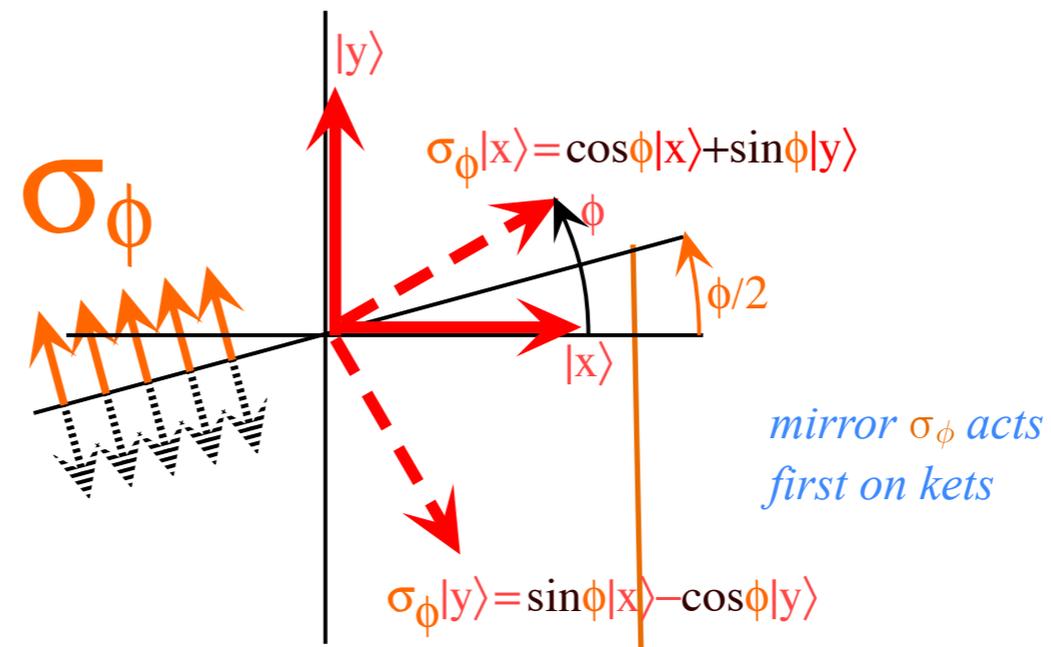
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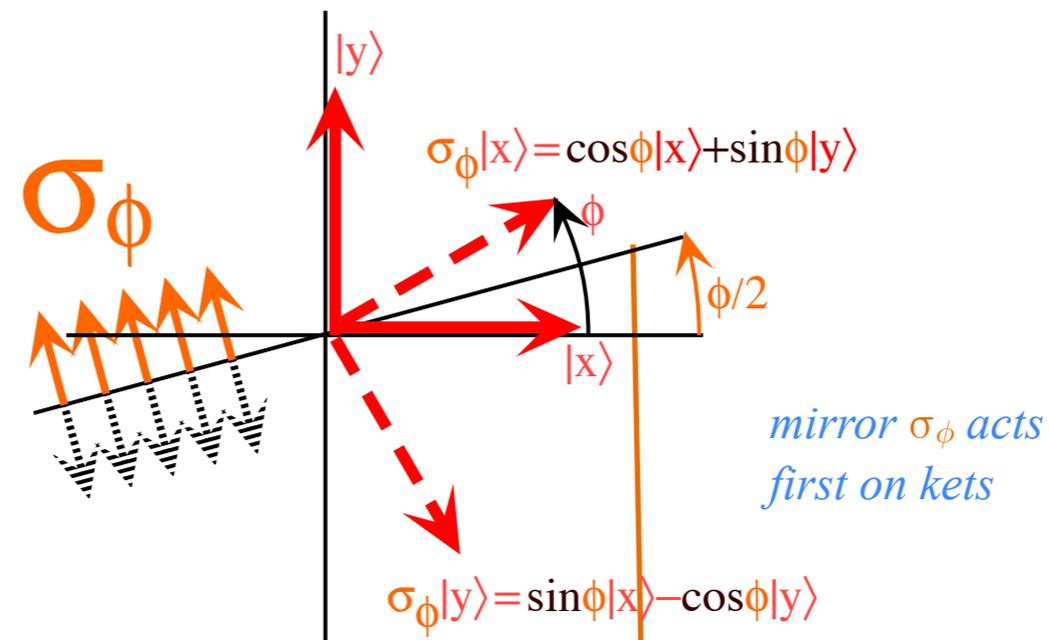
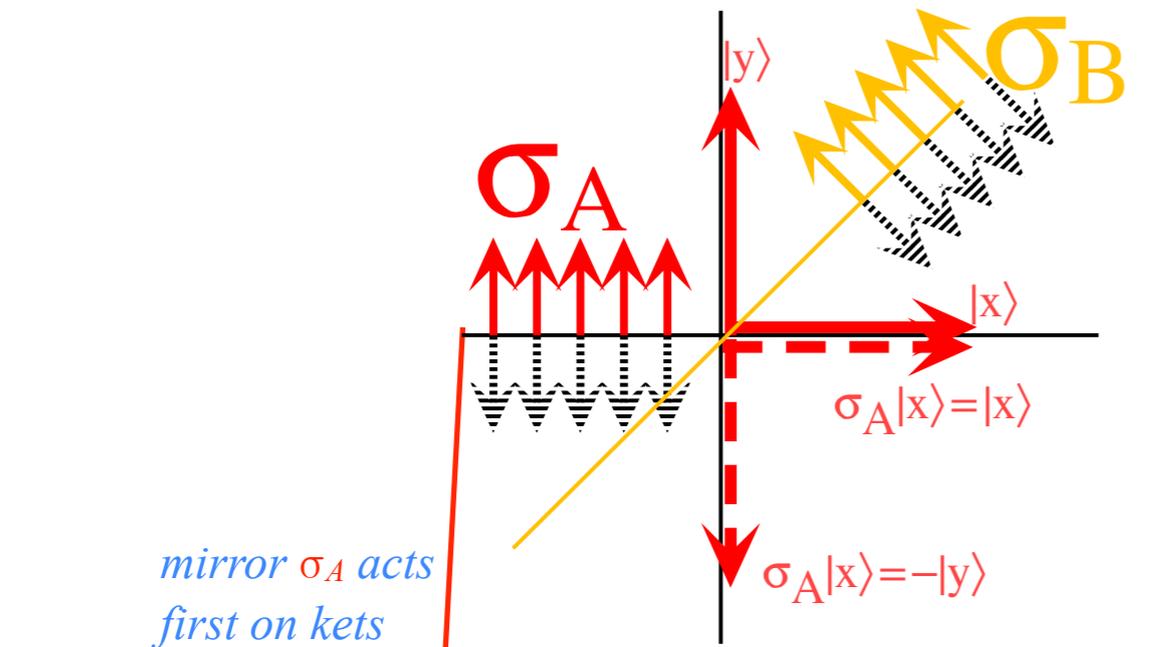
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$$\sigma_A \sigma_B = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} = \begin{pmatrix} 0 & +1 \\ -1 & 0 \end{pmatrix} = i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} = ?$$

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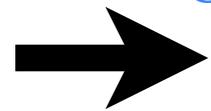
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Geometry of groups: Hamilton's turns and It's all done with mirrors!



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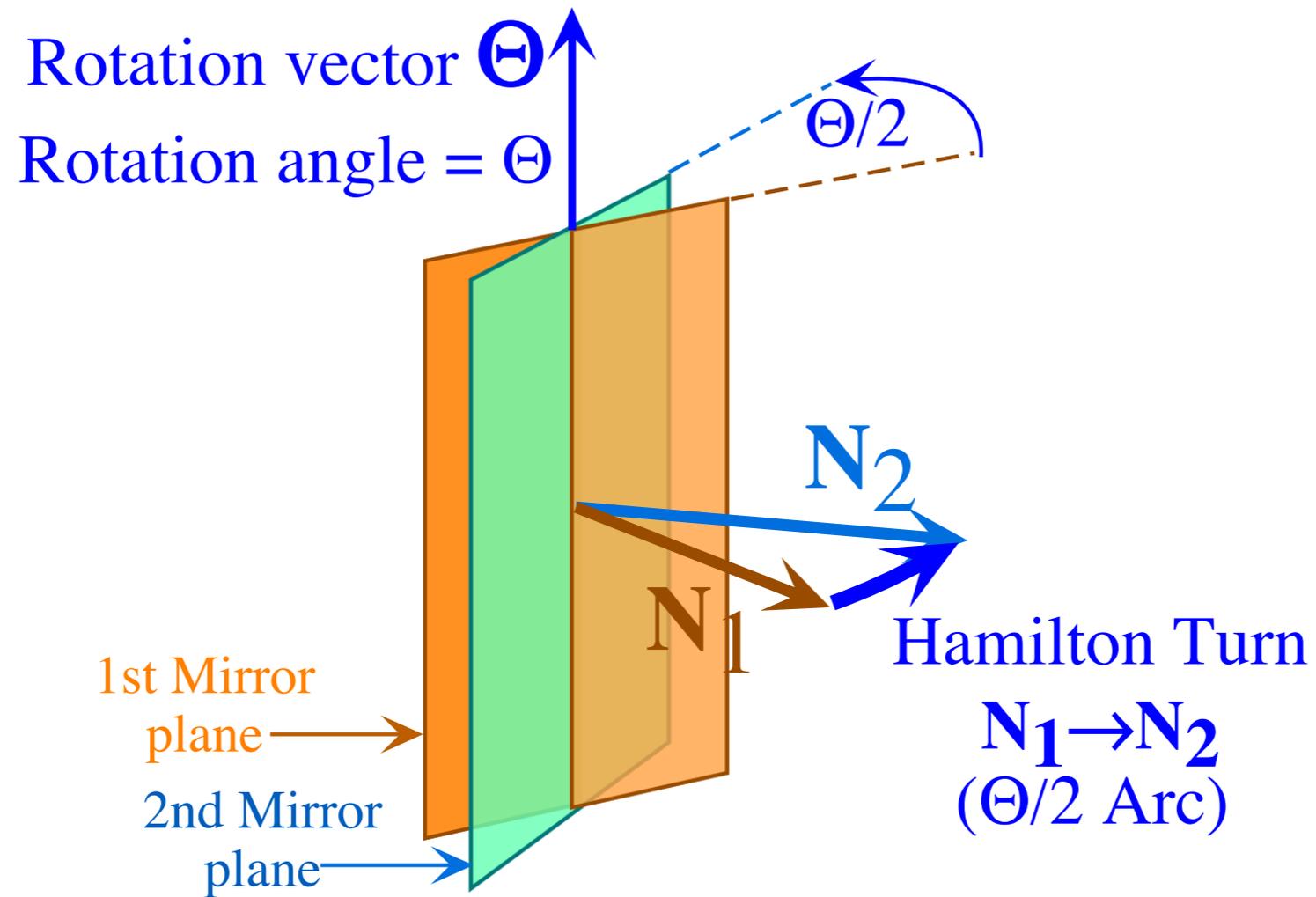
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*Fig. 10.A.7 Mirror reflection planes, normals, and Hamilton-turn arc vector.*

Geometry of  $U(2)$  group products: Hamilton's Turns

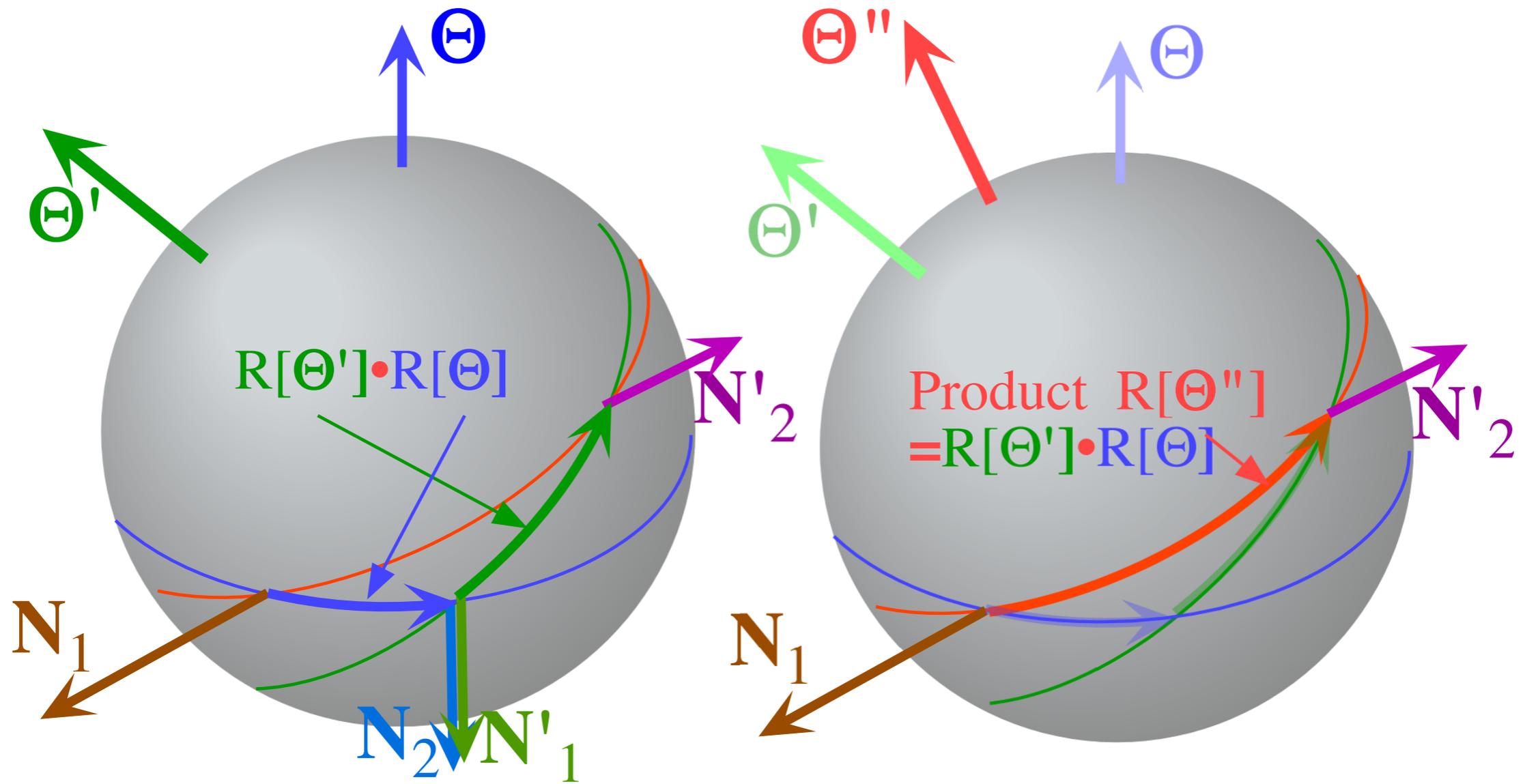


Fig. 10.A.8 Adding Hamilton-turn arcs to compute a  $U(2)$  product  $\mathbf{R}[\Theta''] = \mathbf{R}[\Theta']\mathbf{R}[\Theta]$ .

Each arc  $\Theta/2$ ,  $\Theta'/2$ , or  $\Theta''/2$  is  $1/2$  actual angle  $\Theta$ ,  $\Theta'$ , or  $\Theta''$  of rotation  $\mathbf{R}[\Theta]$ ,  $\mathbf{R}[\Theta']$ , or  $\mathbf{R}[\Theta'']$ .

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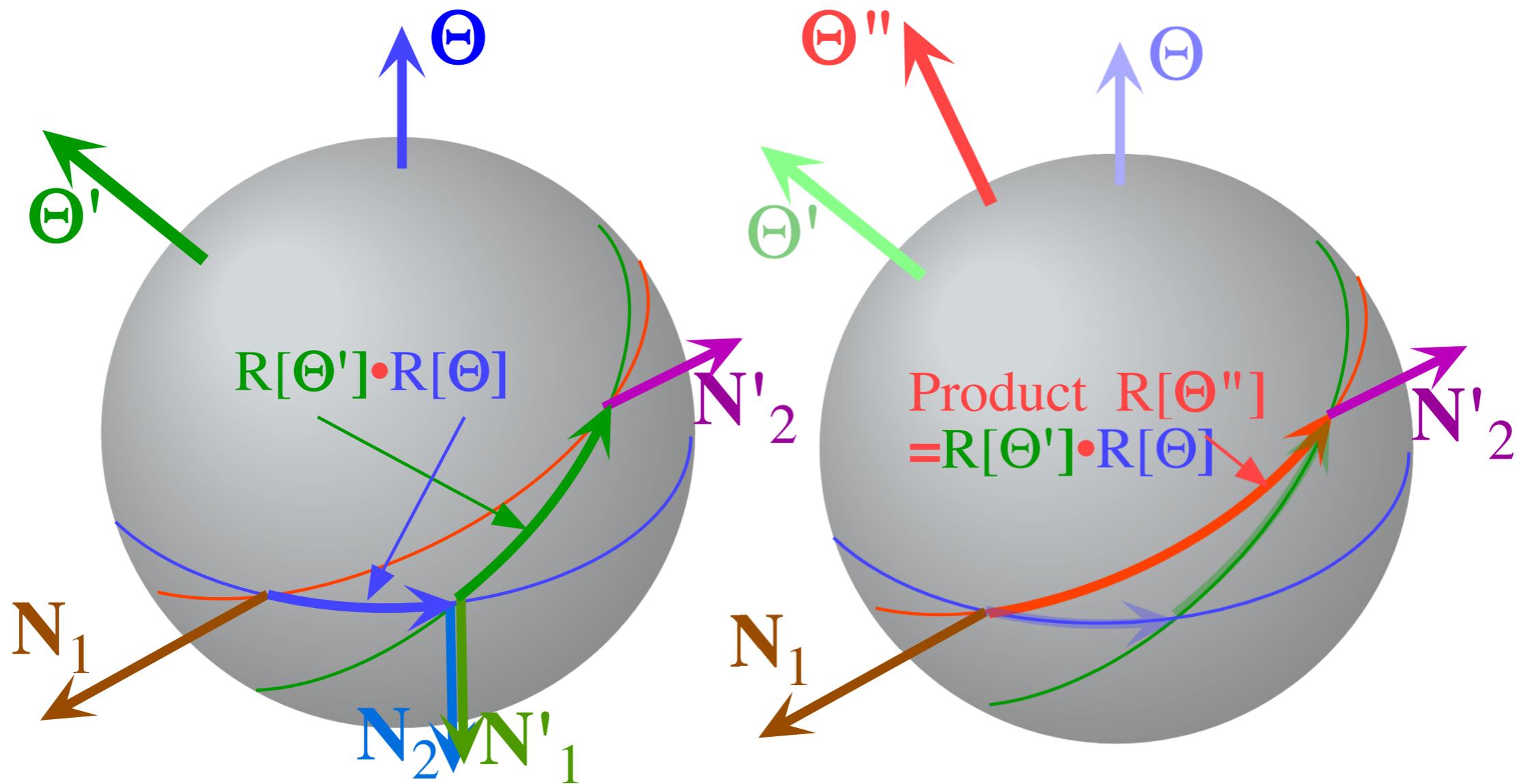


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 Arc  $\Theta/2$  between  $\mathbf{N}_1$  and  $\mathbf{N}_2$  and its supplement  $(\Theta \pm 2\pi)/2 = \Theta/2 \pm \pi$  between  $\mathbf{N}_1$  and  $-\mathbf{N}_2$  represent the same classical rotation by  $\Theta$ .

Geometry of  $U(2)$  group products: Hamilton's Turns

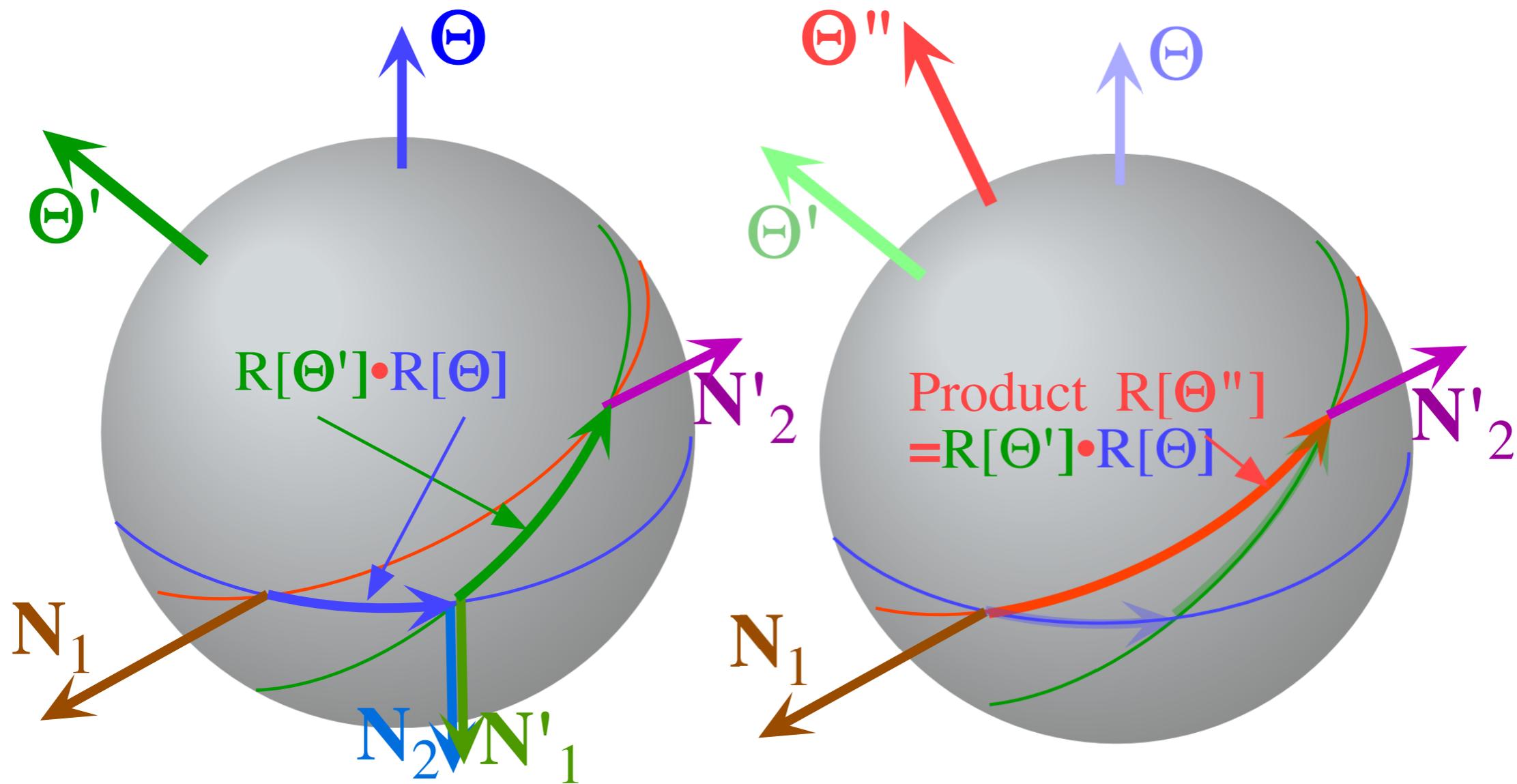


Fig. 10.A.8 Adding Hamilton-turn arcs to compute a  $U(2)$  product  $\mathbf{R}[\Theta''] = \mathbf{R}[\Theta'] \mathbf{R}[\Theta]$ .

Each arc  $\Theta/2$ ,  $\Theta'/2$ , or  $\Theta''/2$  is  $1/2$  actual angle  $\Theta$ ,  $\Theta'$ , or  $\Theta''$  of rotation  $\mathbf{R}[\Theta]$ ,  $\mathbf{R}[\Theta']$ , or  $\mathbf{R}[\Theta'']$ .  
 Arc  $\Theta/2$  between  $\mathbf{N}_1$  and  $\mathbf{N}_2$  and its supplement  $(\Theta \pm 2\pi)/2 = \Theta/2 \pm \pi$  between  $\mathbf{N}_1$  and  $-\mathbf{N}_2$  represent the same classical rotation by  $\Theta$ .

For quantum spin- $1/2$  object, the arc pointing from  $\mathbf{N}_1$  to the antipodal normal  $-\mathbf{N}_2$  represents a  $\Theta$ -rotation with an extra  $\pi$ -phase factor  $e^{\pm i\pi} = -1$ , that is,  $-\mathbf{R}[\Theta]$ .

Review: How “Crazy-Thing”-Theorem makes spinor and vector representation matrices  
Half-angle  $\Theta/2 = \varphi$  replacement and Darboux crank axis operators

Operator-on-Operator transformations

Product algebra for Pauli's  $\sigma_\mu$  and Hamilton's  $q_\mu = -i\sigma_\mu$

Group product algebra

Jordan-Pauli identity and  $U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - formula

Transformation  $\mathbf{R}[\Theta]\sigma_\mu\mathbf{R}[\Theta]^\dagger$  of spinor  $\sigma_\mu$ -operators

Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  of group-operators

Operator-on-Operator transformations

Geometry of groups: Hamilton's turns and It's all done with mirrors!

Group product geometry

  $U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - geometry  
Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  geometry

Euler  $\mathbf{R}(\alpha\beta\gamma)$  versus Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  related to Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  rotation  $\Theta = 0 - 4\pi$ -sequence  $[\varphi\vartheta]$  fixed

$R(3)$ - $U(2)$  slide rule for converting  $\mathbf{R}(\alpha\beta\gamma) \leftrightarrow \mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  Sundial

Geometry of transformation of rotational  $\mathbf{R}[\Theta]$ -operators

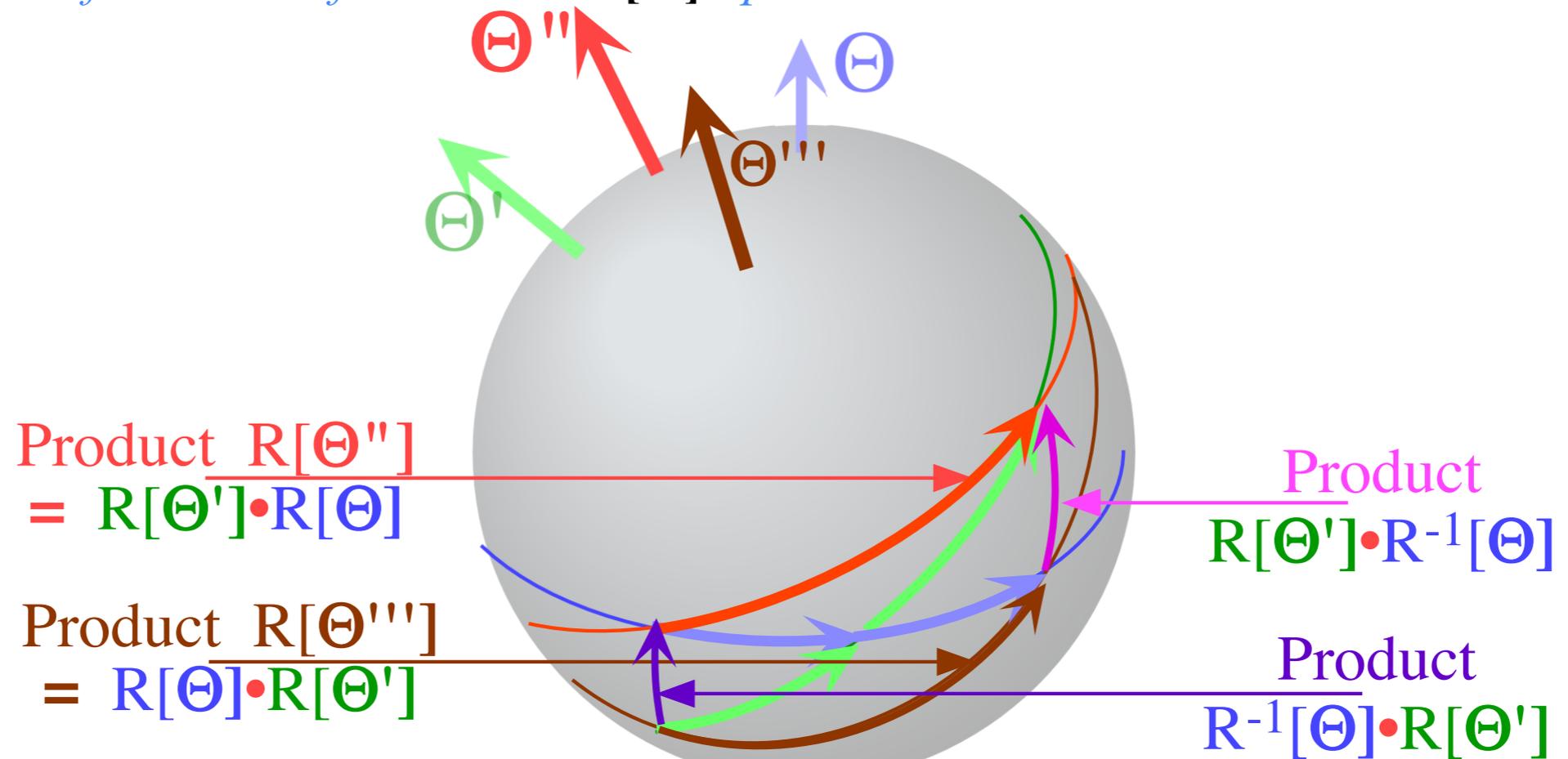


Fig. 10.A.9 Hamilton-turn arc parallelogram with  $\mathbf{R}[\Theta''] = \mathbf{R}[\Theta']\mathbf{R}[\Theta]$  and  $\mathbf{R}[\Theta'''] = \mathbf{R}[\Theta]\mathbf{R}[\Theta']$

Geometry of transformation of rotational  $\mathbf{R}[\Theta]$ -operators

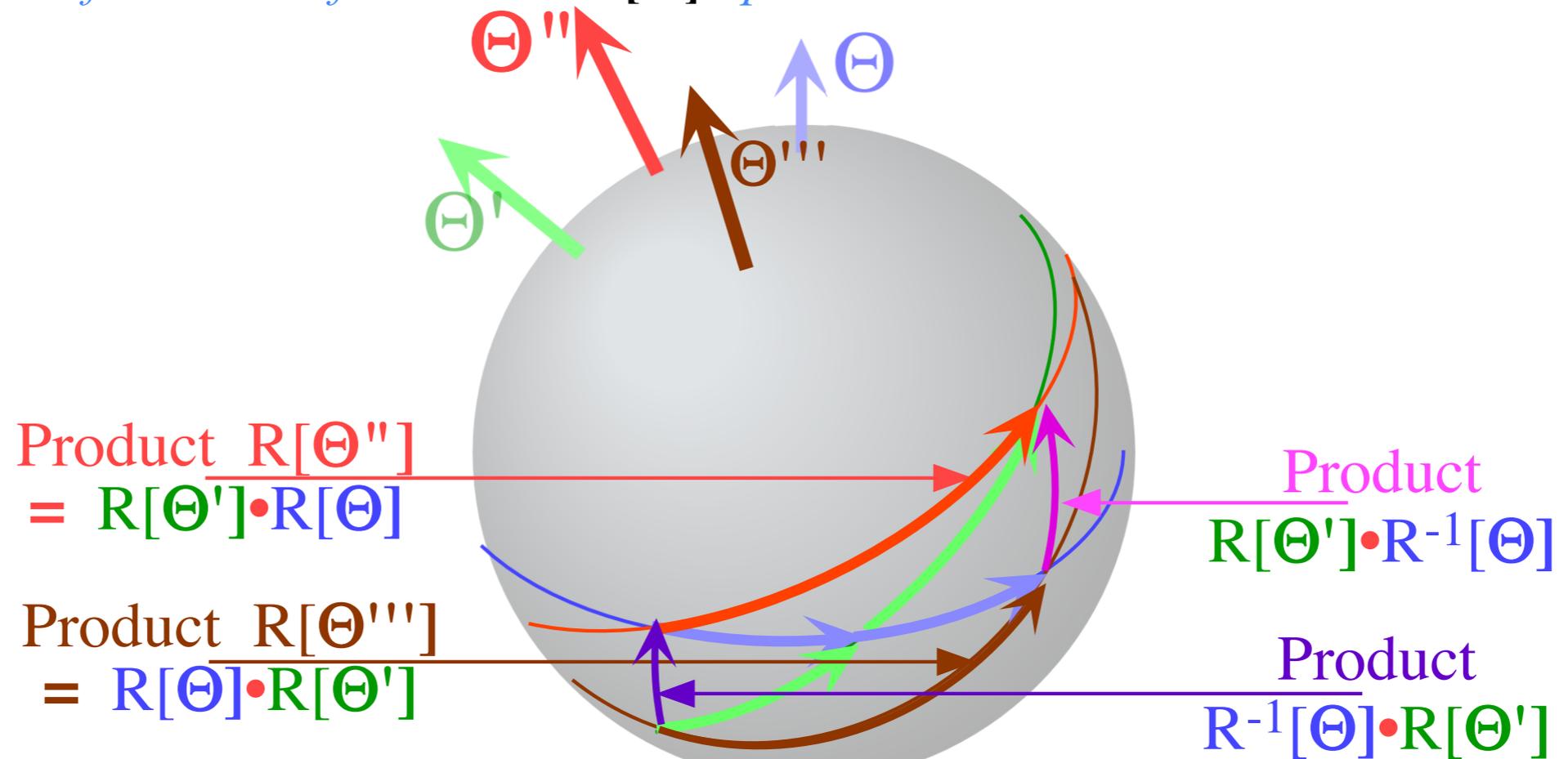


Fig. 10.A.9 Hamilton-turn arc parallelogram with  $\mathbf{R}[\Theta''] = \mathbf{R}[\Theta']\mathbf{R}[\Theta]$  and  $\mathbf{R}[\Theta'''] = \mathbf{R}[\Theta]\mathbf{R}[\Theta']$

Vectors added in the reverse order give  $\mathbf{R}[\Theta'''] = \mathbf{R}[\Theta]\mathbf{R}[\Theta']$  instead of  $\mathbf{R}[\Theta''] = \mathbf{R}[\Theta']\mathbf{R}[\Theta]$ .

Geometry of transformation of rotational  $\mathbf{R}[\Theta]$ -operators

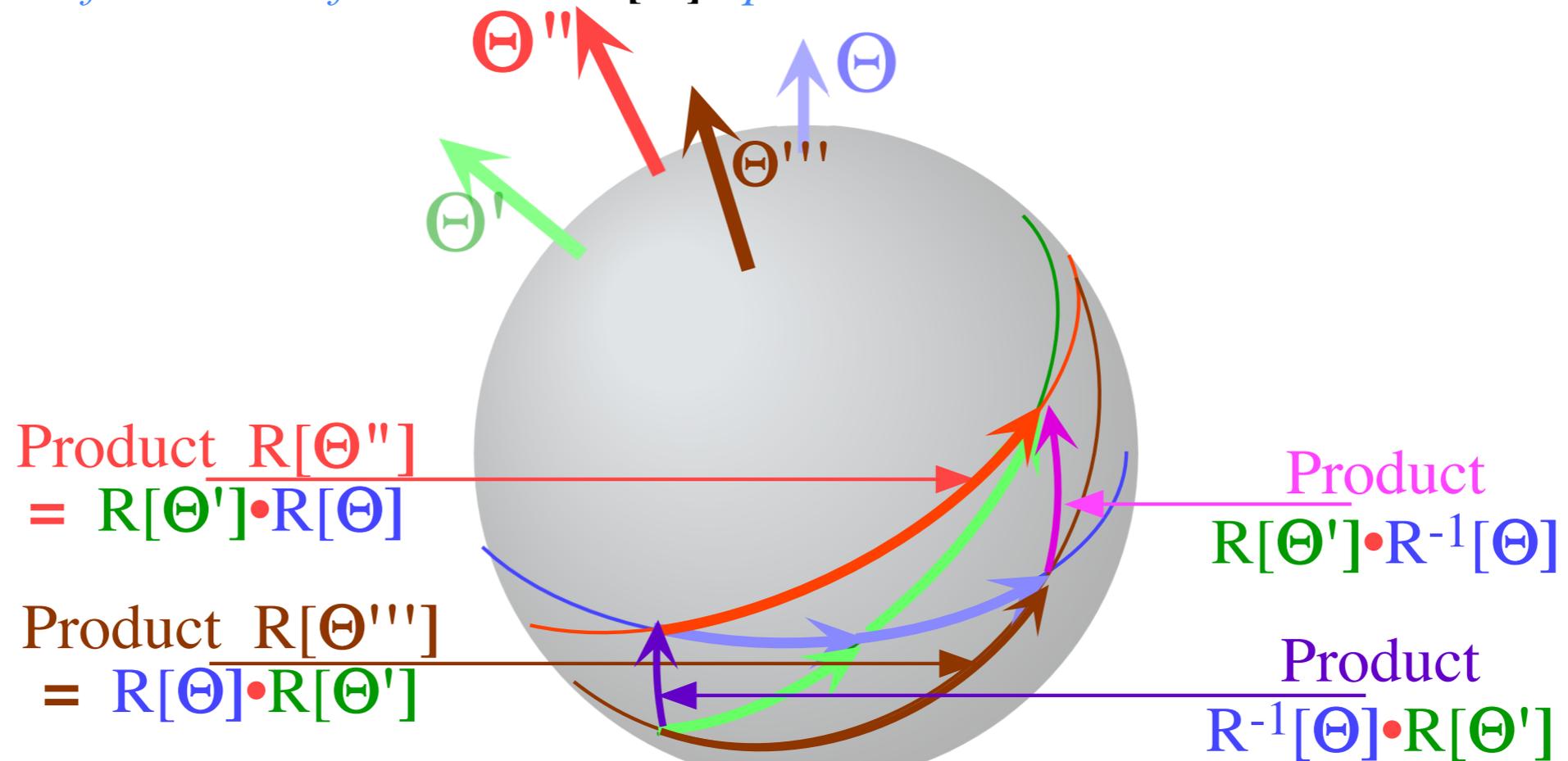


Fig. 10.A.9 Hamilton-turn arc parallelogram with  $\mathbf{R}[\Theta''] = \mathbf{R}[\Theta']\mathbf{R}[\Theta]$  and  $\mathbf{R}[\Theta'''] = \mathbf{R}[\Theta]\mathbf{R}[\Theta']$

Vectors added in the reverse order give  $\mathbf{R}[\Theta'''] = \mathbf{R}[\Theta]\mathbf{R}[\Theta']$  instead of  $\mathbf{R}[\Theta''] = \mathbf{R}[\Theta']\mathbf{R}[\Theta]$ .

A similarity transformation of rotation  $\mathbf{R}[\Theta'']$  by rotation  $\mathbf{R}[\Theta]$  gives rotation  $\mathbf{R}[\Theta''']$

$$\mathbf{R}[\Theta] \mathbf{R}[\Theta''] \mathbf{R}[-\Theta] = \mathbf{R}[\Theta''']$$

Geometry of transformation of rotational  $\mathbf{R}[\Theta]$ -operators

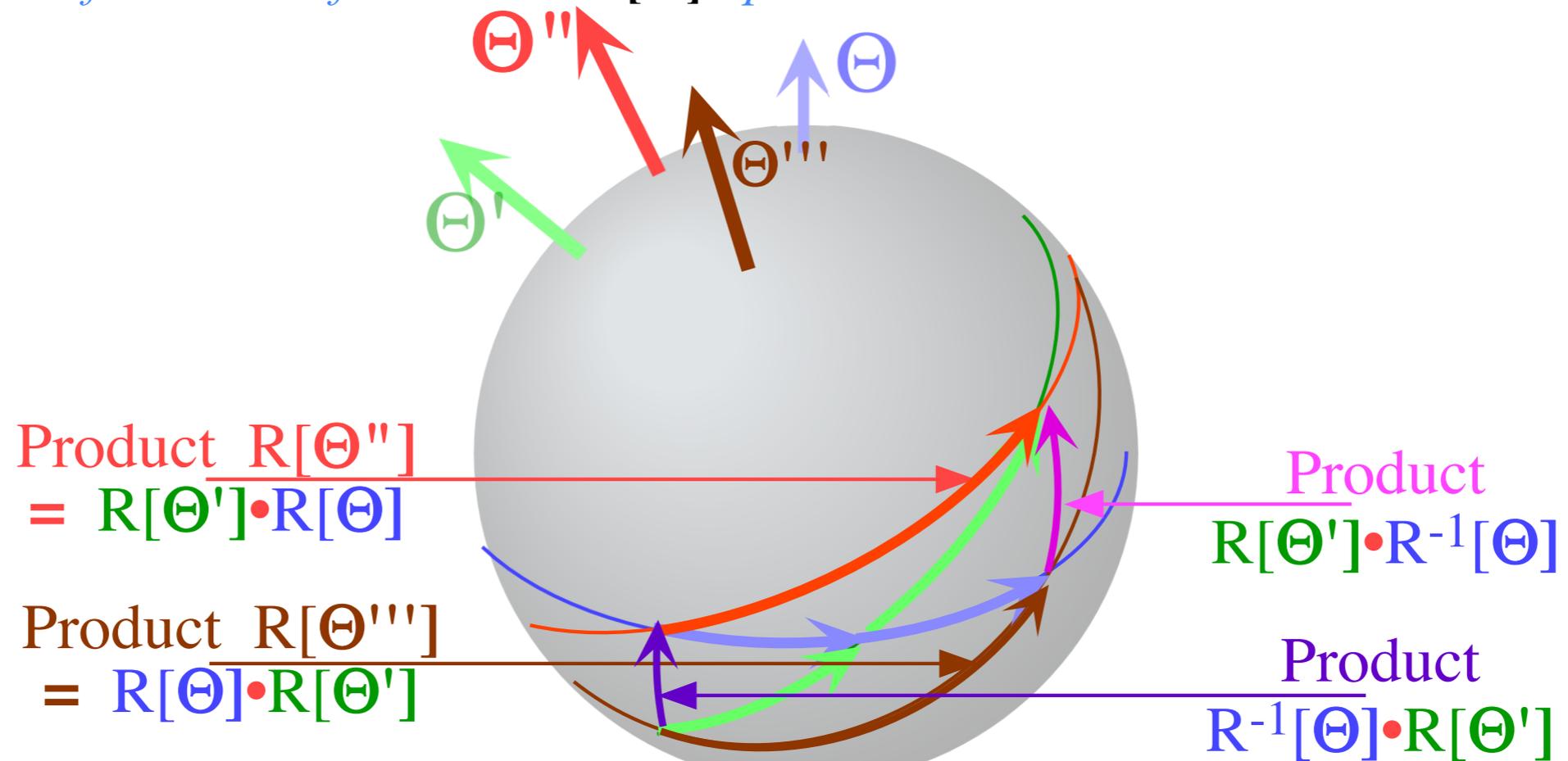


Fig. 10.A.9 Hamilton-turn arc parallelogram with  $\mathbf{R}[\Theta''] = \mathbf{R}[\Theta']\mathbf{R}[\Theta]$  and  $\mathbf{R}[\Theta'''] = \mathbf{R}[\Theta]\mathbf{R}[\Theta']$

Vectors added in the reverse order give  $\mathbf{R}[\Theta'''] = \mathbf{R}[\Theta]\mathbf{R}[\Theta']$  instead of  $\mathbf{R}[\Theta''] = \mathbf{R}[\Theta']\mathbf{R}[\Theta]$ .

A similarity transformation of rotation  $\mathbf{R}[\Theta'']$  by rotation  $\mathbf{R}[\Theta]$  gives rotation  $\mathbf{R}[\Theta''']$  and *vice-versa*:

$$\mathbf{R}[\Theta] \mathbf{R}[\Theta''] \mathbf{R}[-\Theta] = \mathbf{R}[\Theta''']$$

$$\mathbf{R}[-\Theta] \mathbf{R}[\Theta'''] \mathbf{R}[\Theta] = \mathbf{R}[\Theta'']$$

Geometry of transformation of rotational  $\mathbf{R}[\Theta]$ -operators

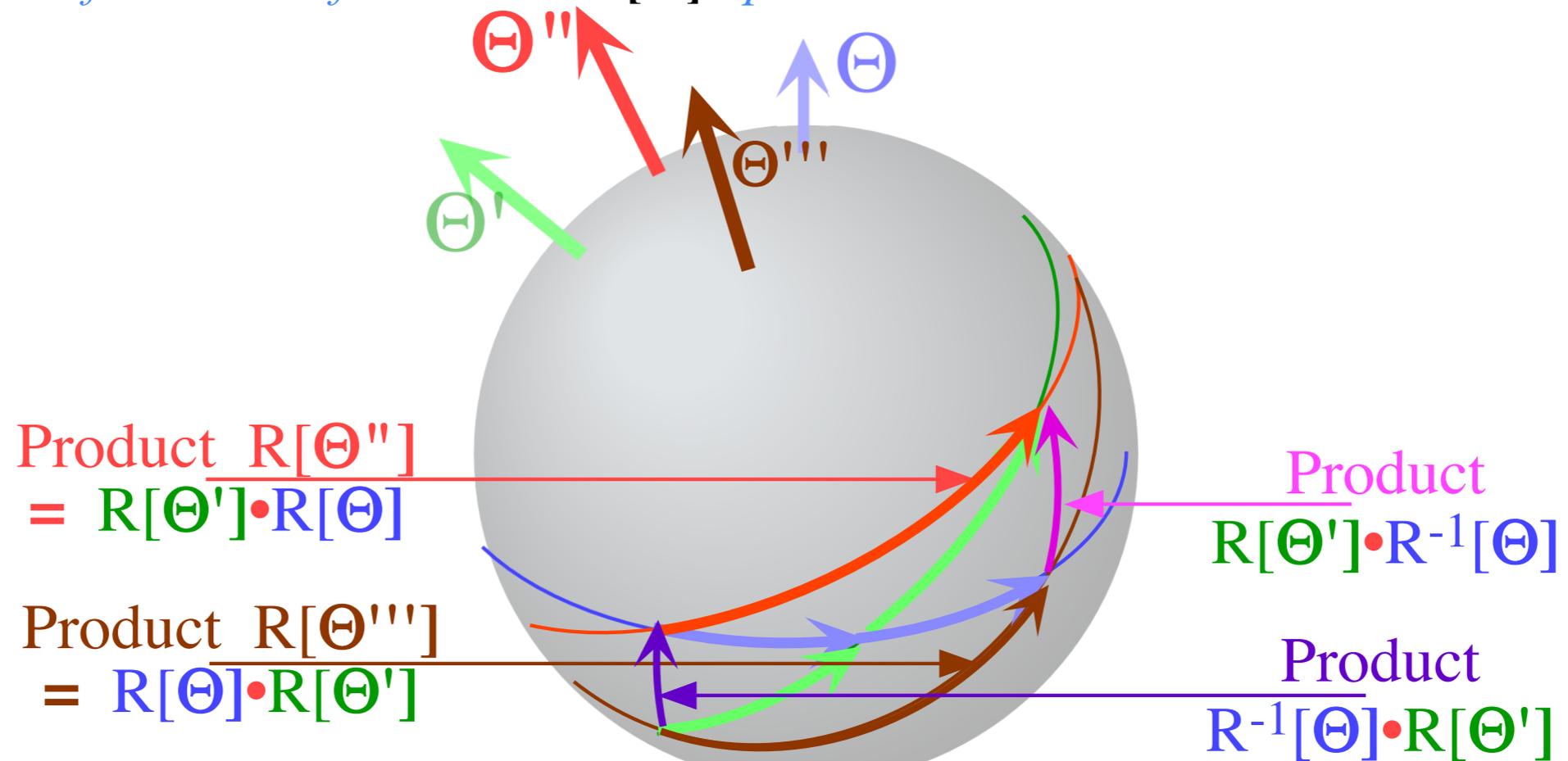


Fig. 10.A.9 Hamilton-turn arc parallelogram with  $\mathbf{R}[\Theta''] = \mathbf{R}[\Theta']\mathbf{R}[\Theta]$  and  $\mathbf{R}[\Theta'''] = \mathbf{R}[\Theta]\mathbf{R}[\Theta']$

Vectors added in the reverse order give  $\mathbf{R}[\Theta'''] = \mathbf{R}[\Theta]\mathbf{R}[\Theta']$  instead of  $\mathbf{R}[\Theta''] = \mathbf{R}[\Theta']\mathbf{R}[\Theta]$ .

A similarity transformation of rotation  $\mathbf{R}[\Theta'']$  by rotation  $\mathbf{R}[\Theta]$  gives rotation  $\mathbf{R}[\Theta''']$  and *vice-versa*:

$$\mathbf{R}[\Theta] \mathbf{R}[\Theta''] \mathbf{R}[-\Theta] = \mathbf{R}[\Theta''']$$

$$\mathbf{R}[-\Theta] \mathbf{R}[\Theta'''] \mathbf{R}[\Theta] = \mathbf{R}[\Theta'']$$

Everything associated with rotation  $\mathbf{R}[\Theta'']$  is rotated by full angle  $\Theta$  around axis  $\Theta$ .

Geometry of transformation of rotational  $\mathbf{R}[\Theta]$ -operators

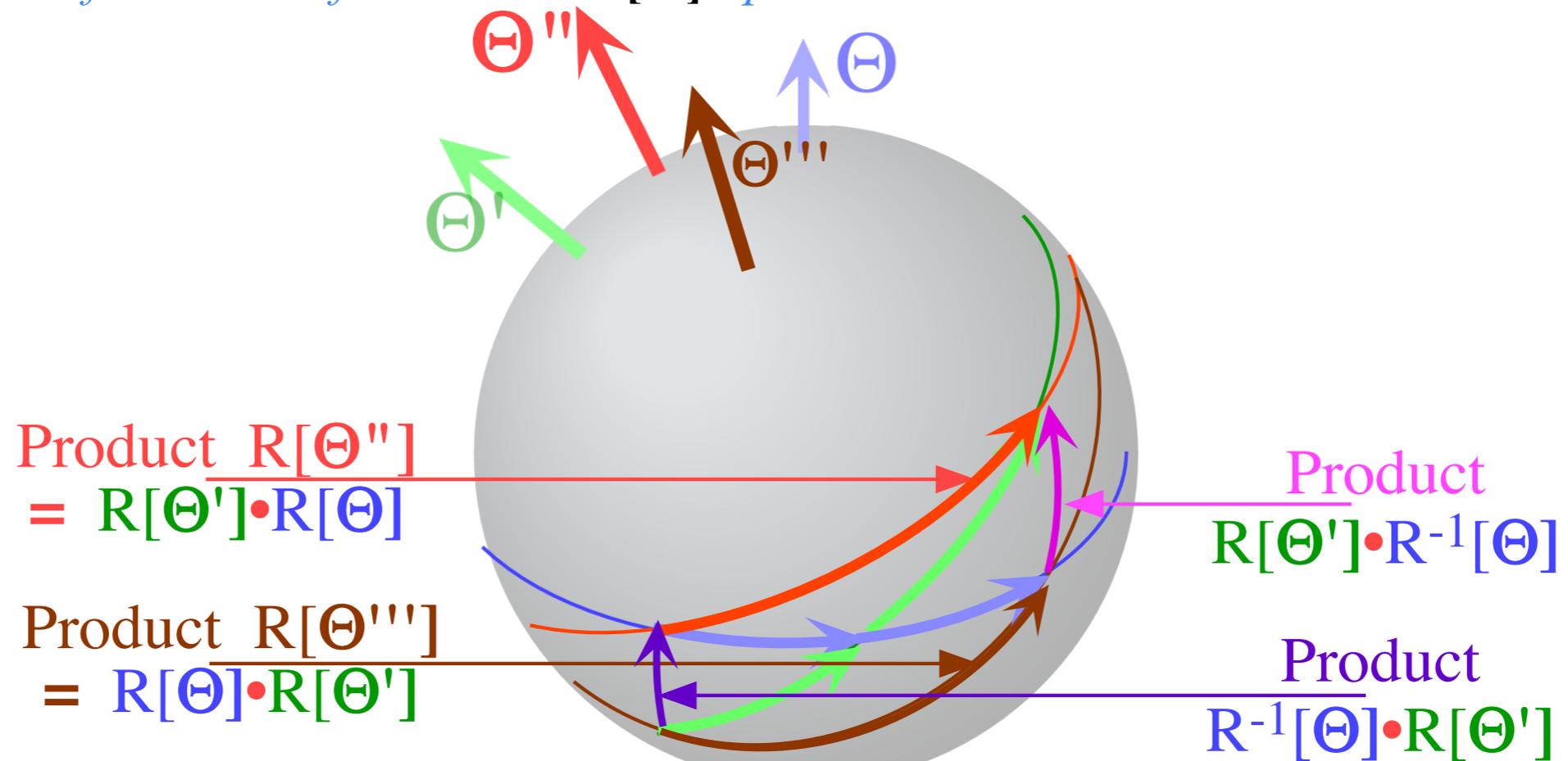


Fig. 10.A.9 Hamilton-turn arc parallelogram with  $\mathbf{R}[\Theta''] = \mathbf{R}[\Theta']\mathbf{R}[\Theta]$  and  $\mathbf{R}[\Theta'''] = \mathbf{R}[\Theta]\mathbf{R}[\Theta']$

Vectors added in the reverse order give  $\mathbf{R}[\Theta'''] = \mathbf{R}[\Theta]\mathbf{R}[\Theta']$  instead of  $\mathbf{R}[\Theta''] = \mathbf{R}[\Theta']\mathbf{R}[\Theta]$ .

A similarity transformation of rotation  $\mathbf{R}[\Theta'']$  by rotation  $\mathbf{R}[\Theta]$  gives rotation  $\mathbf{R}[\Theta''']$  and *vice-versa*:

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Everything associated with rotation  $\mathbf{R}[\Theta'']$  is rotated by full angle  $\Theta$  around axis  $\Theta$ .

Crank vector  $\Theta$  and its turn arc moved by two  $\mathbf{R}[\Theta]$  turn arcs into turn arc of  $\mathbf{R}[\Theta''']$  below it.

Geometry of transformation of rotational  $\mathbf{R}[\Theta]$ -operators

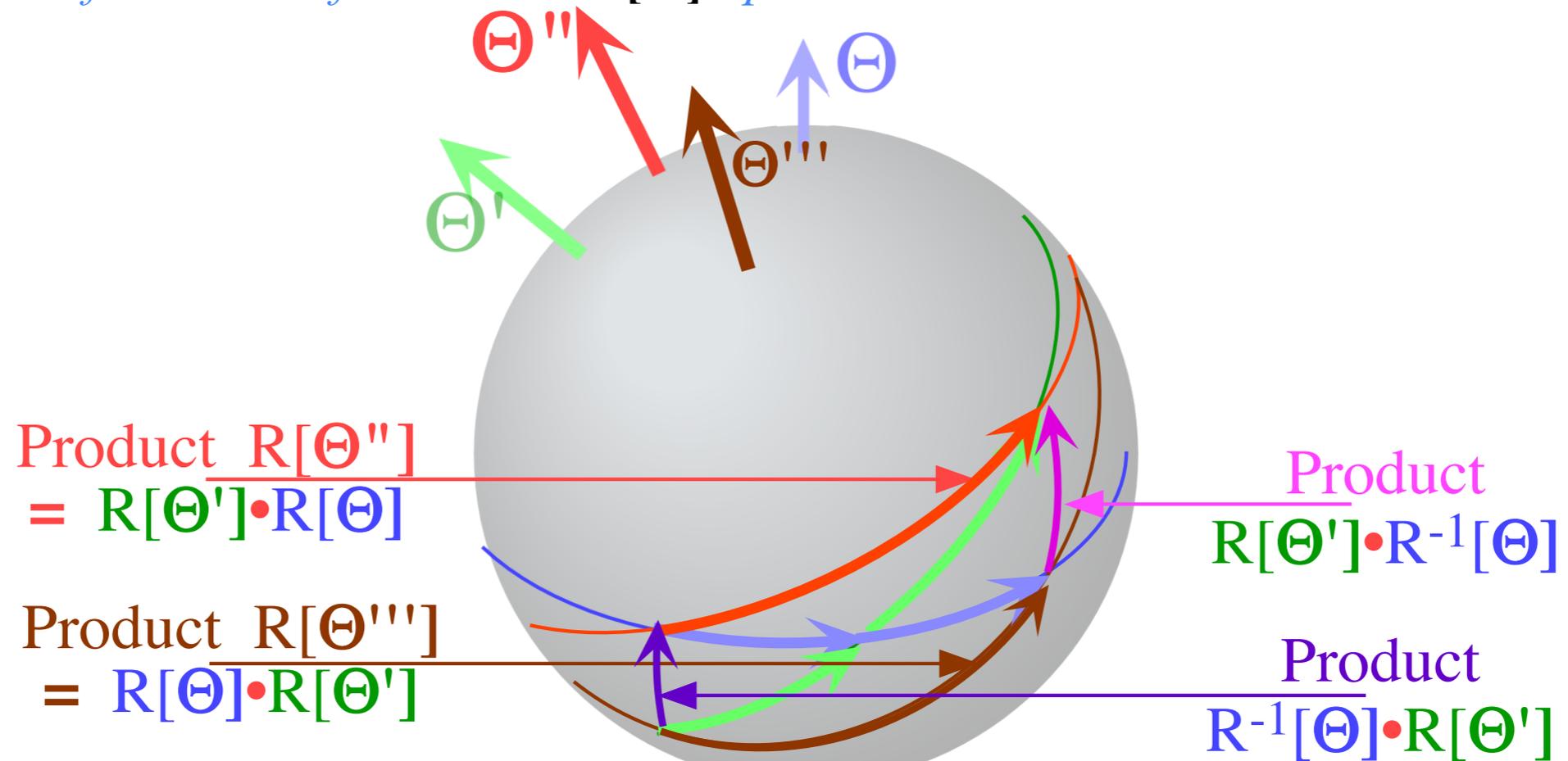


Fig. 10.A.9 Hamilton-turn arc parallelogram with  $\mathbf{R}[\Theta''] = \mathbf{R}[\Theta']\mathbf{R}[\Theta]$  and  $\mathbf{R}[\Theta'''] = \mathbf{R}[\Theta]\mathbf{R}[\Theta']$

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A similarity transformation of rotation  $\mathbf{R}[\Theta'']$  by rotation  $\mathbf{R}[\Theta]$  gives rotation  $\mathbf{R}[\Theta''']$  and *vice-versa*:

$$\mathbf{R}[\Theta] \mathbf{R}[\Theta''] \mathbf{R}[-\Theta] = \mathbf{R}[\Theta'''] \qquad \mathbf{R}[-\Theta] \mathbf{R}[\Theta'''] \mathbf{R}[\Theta] = \mathbf{R}[\Theta'']$$

Everything associated with rotation  $\mathbf{R}[\Theta'']$  is rotated by full angle  $\Theta$  around axis  $\Theta$ .

Crank vector  $\Theta$  and its turn arc moved by two  $\mathbf{R}[\Theta]$  turn arcs into turn arc of  $\mathbf{R}[\Theta''']$  below it.

Another similarity transformation of rotation  $\mathbf{R}[\Theta''']$  by rotation  $\mathbf{R}[\Theta']$  to  $\mathbf{R}[\Theta'']$

$$\mathbf{R}[\Theta'] \mathbf{R}[\Theta'''] \mathbf{R}[-\Theta'] = \mathbf{R}[\Theta'']$$

Geometry of transformation of rotational  $\mathbf{R}[\Theta]$ -operators

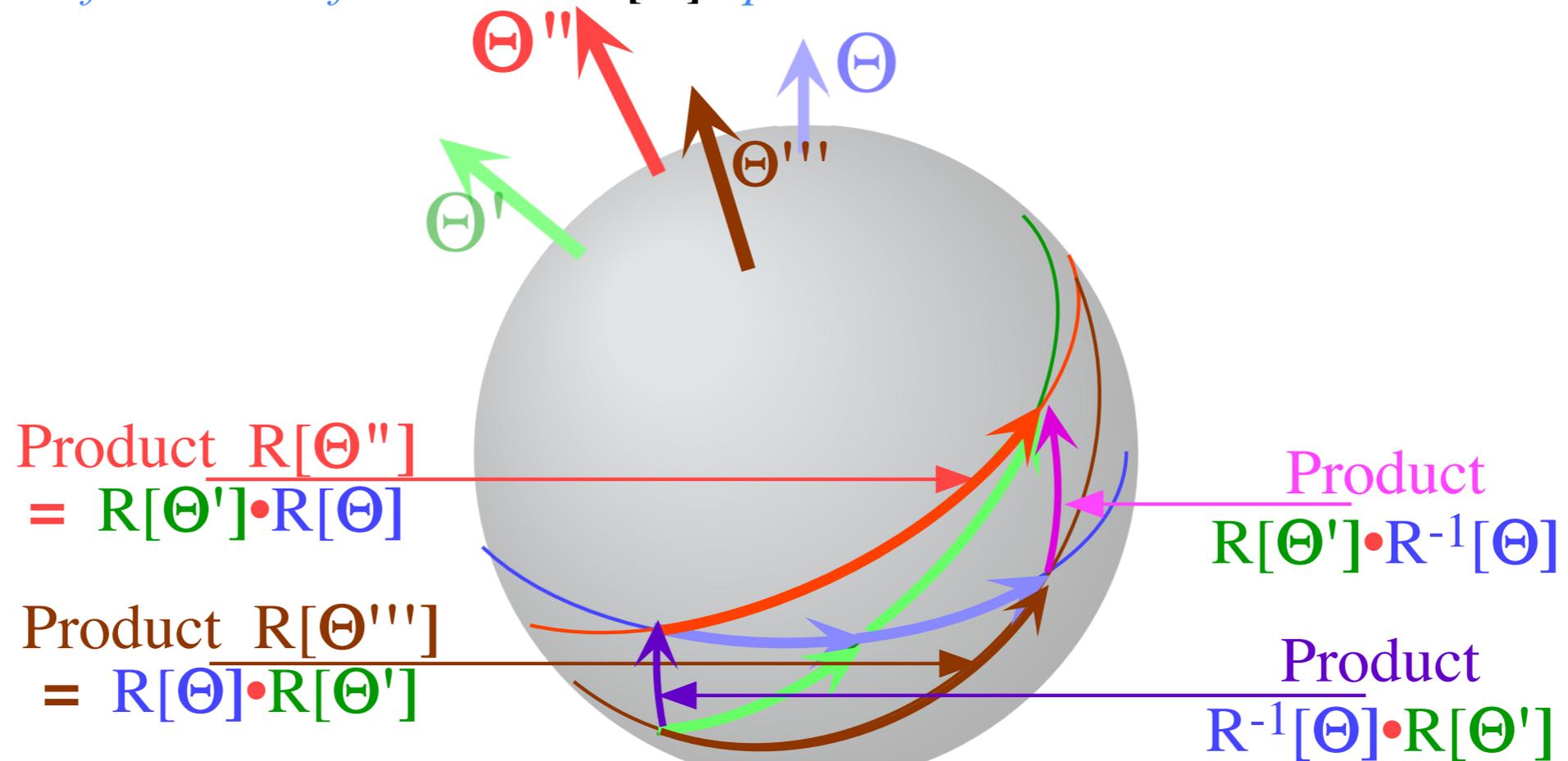


Fig. 10.A.9 Hamilton-turn arc parallelogram with  $\mathbf{R}[\Theta''] = \mathbf{R}[\Theta']\mathbf{R}[\Theta]$  and  $\mathbf{R}[\Theta'''] = \mathbf{R}[\Theta]\mathbf{R}[\Theta']$

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A similarity transformation of rotation  $\mathbf{R}[\Theta'']$  by rotation  $\mathbf{R}[\Theta]$  gives rotation  $\mathbf{R}[\Theta''']$  and *vice-versa*:

$$\mathbf{R}[\Theta] \mathbf{R}[\Theta''] \mathbf{R}[-\Theta] = \mathbf{R}[\Theta'''] \quad \mathbf{R}[-\Theta] \mathbf{R}[\Theta'''] \mathbf{R}[\Theta] = \mathbf{R}[\Theta'']$$

Everything associated with rotation  $\mathbf{R}[\Theta'']$  is rotated by full angle  $\Theta$  around axis  $\Theta$ .

Crank vector  $\Theta$  and its turn arc moved by two  $\mathbf{R}[\Theta]$  turn arcs into turn arc of  $\mathbf{R}[\Theta''']$  below it.

Another similarity transformation of rotation  $\mathbf{R}[\Theta''']$  by rotation  $\mathbf{R}[\Theta']$  to  $\mathbf{R}[\Theta'']$  and *vice-versa*:

$$\mathbf{R}[\Theta'] \mathbf{R}[\Theta'''] \mathbf{R}[-\Theta'] = \mathbf{R}[\Theta''] \quad \mathbf{R}[-\Theta'] \mathbf{R}[\Theta''] \mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$$

Geometry of transformation of rotational  $\mathbf{R}[\Theta]$ -operators

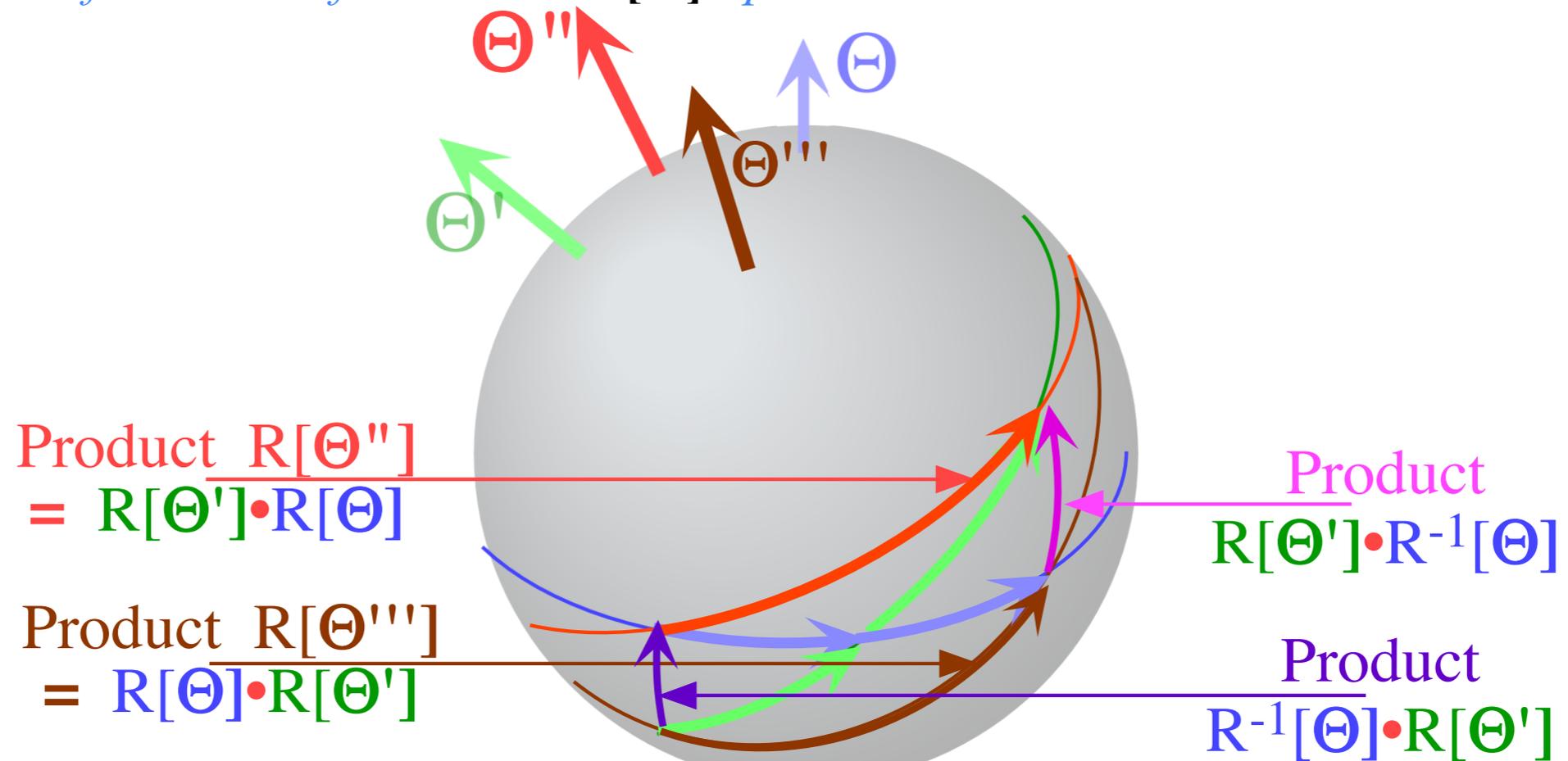


Fig. 10.A.9 Hamilton-turn arc parallelogram with  $\mathbf{R}[\Theta''] = \mathbf{R}[\Theta']\mathbf{R}[\Theta]$  and  $\mathbf{R}[\Theta'''] = \mathbf{R}[\Theta]\mathbf{R}[\Theta']$

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A similarity transformation of rotation  $\mathbf{R}[\Theta'']$  by rotation  $\mathbf{R}[\Theta]$  gives rotation  $\mathbf{R}[\Theta''']$  and *vice-versa*:

$$\mathbf{R}[\Theta] \mathbf{R}[\Theta''] \mathbf{R}[-\Theta] = \mathbf{R}[\Theta'''] \quad \mathbf{R}[-\Theta] \mathbf{R}[\Theta'''] \mathbf{R}[\Theta] = \mathbf{R}[\Theta'']$$

Everything associated with rotation  $\mathbf{R}[\Theta'']$  is rotated by full angle  $\Theta$  around axis  $\Theta$ .

Crank vector  $\Theta$  and its turn arc moved by two  $\mathbf{R}[\Theta]$  turn arcs into turn arc of  $\mathbf{R}[\Theta''']$  below it.

Another similarity transformation of rotation  $\mathbf{R}[\Theta''']$  by rotation  $\mathbf{R}[\Theta']$  to  $\mathbf{R}[\Theta'']$  and *vice-versa*:

$$\mathbf{R}[\Theta'] \mathbf{R}[\Theta'''] \mathbf{R}[-\Theta'] = \mathbf{R}[\Theta''] \quad \mathbf{R}[-\Theta'] \mathbf{R}[\Theta''] \mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$$

Many ( $\infty$ ) rotations transform  $\mathbf{R}[\Theta'']$  into  $\mathbf{R}[\Theta''']$ . Of these, there is one with the least angle  $\Theta_{min}$ .

Review: How “Crazy-Thing”-Theorem makes spinor and vector representation matrices  
Half-angle  $\Theta/2 = \varphi$  replacement and Darboux crank axis operators

Operator-on-Operator transformations

Product algebra for Pauli's  $\sigma_\mu$  and Hamilton's  $q_\mu = -i\sigma_\mu$

Group product algebra

Jordan-Pauli identity and  $U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - formula

Transformation  $\mathbf{R}[\Theta]\sigma_\mu\mathbf{R}[\Theta]^\dagger$  of spinor  $\sigma_\mu$ -operators

Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  of group-operators

Operator-on-Operator transformations

Geometry of groups: Hamilton's turns and It's all done with mirrors!

Group product geometry

$U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - geometry

Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  geometry

→ Euler  $\mathbf{R}(\alpha\beta\gamma)$  versus Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

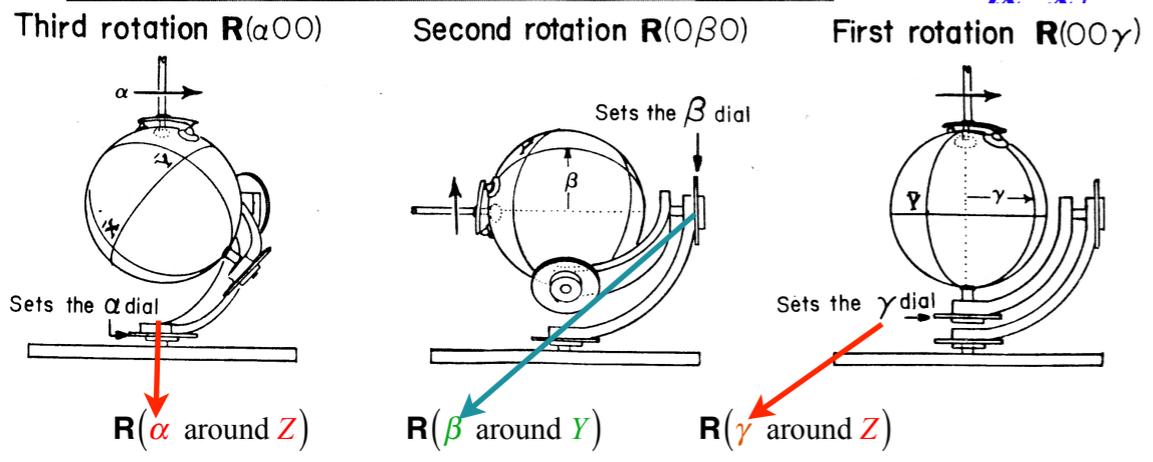
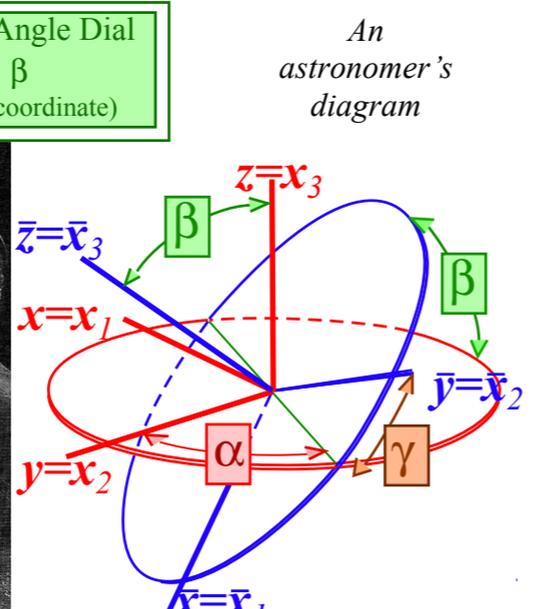
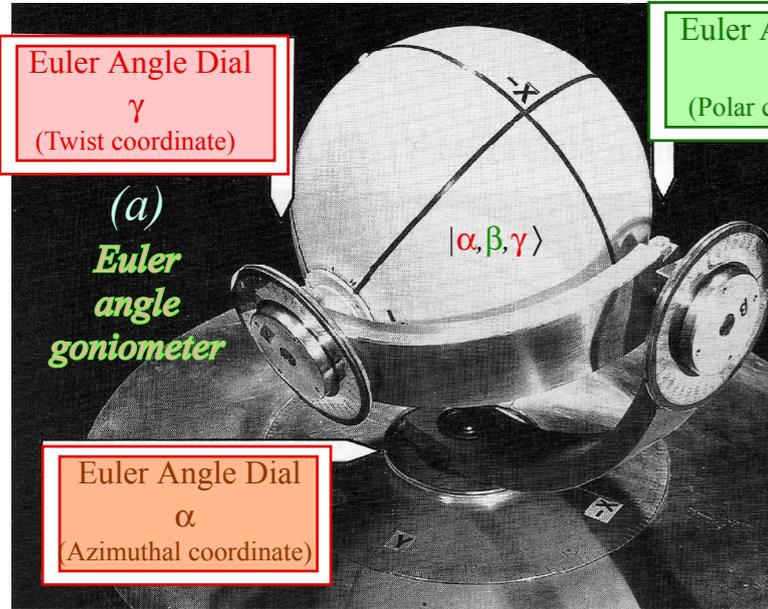
Euler  $\mathbf{R}(\alpha\beta\gamma)$  related to Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  rotation  $\Theta = 0 - 4\pi$ -sequence  $[\varphi\vartheta]$  fixed

$R(3)$ - $U(2)$  slide rule for converting  $\mathbf{R}(\alpha\beta\gamma) \leftrightarrow \mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  Sundial

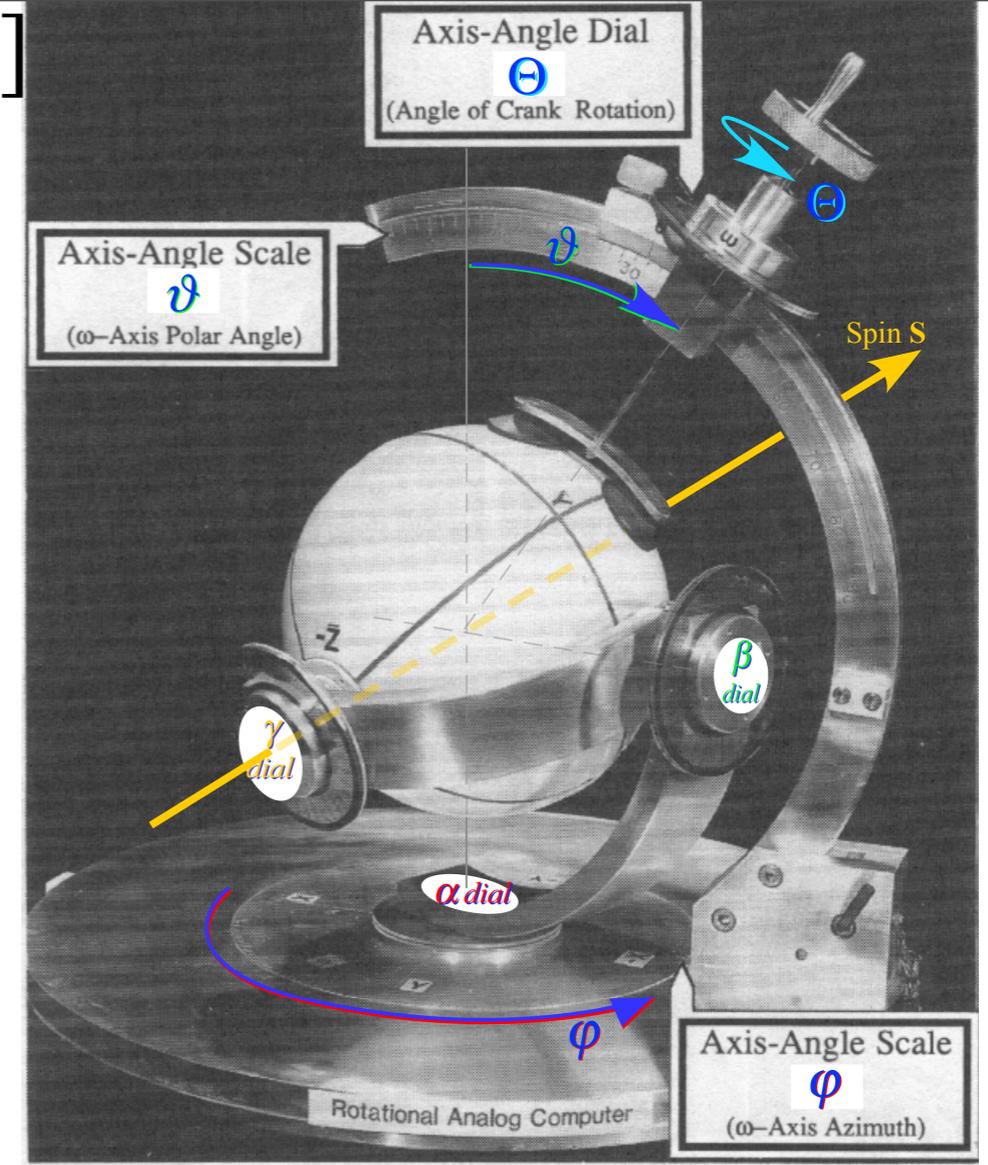
# Euler $R(\alpha\beta\gamma)$ versus Darboux $R[\varphi\vartheta\Theta]$



$$R(\alpha\beta\gamma) = \begin{pmatrix} e^{-i\frac{\alpha}{2}} & 0 \\ 0 & e^{i\frac{\alpha}{2}} \end{pmatrix} \begin{pmatrix} \cos\frac{\beta}{2} & -\sin\frac{\beta}{2} \\ \sin\frac{\beta}{2} & \cos\frac{\beta}{2} \end{pmatrix} \begin{pmatrix} e^{-i\frac{\gamma}{2}} & 0 \\ 0 & e^{i\frac{\gamma}{2}} \end{pmatrix} = \begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} \end{pmatrix}$$

$$= \cos\frac{\alpha+\gamma}{2} \cos\frac{\beta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \sin\frac{\gamma-\alpha}{2} \sin\frac{\beta}{2} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \cos\frac{\gamma-\alpha}{2} \sin\frac{\beta}{2} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \sin\frac{\alpha+\gamma}{2} \cos\frac{\beta}{2}$$

Euler  $R(\alpha\beta\gamma)$  is simpler to form than  $\Theta$ -axis Darboux  $R[\varphi\vartheta\Theta]$ .

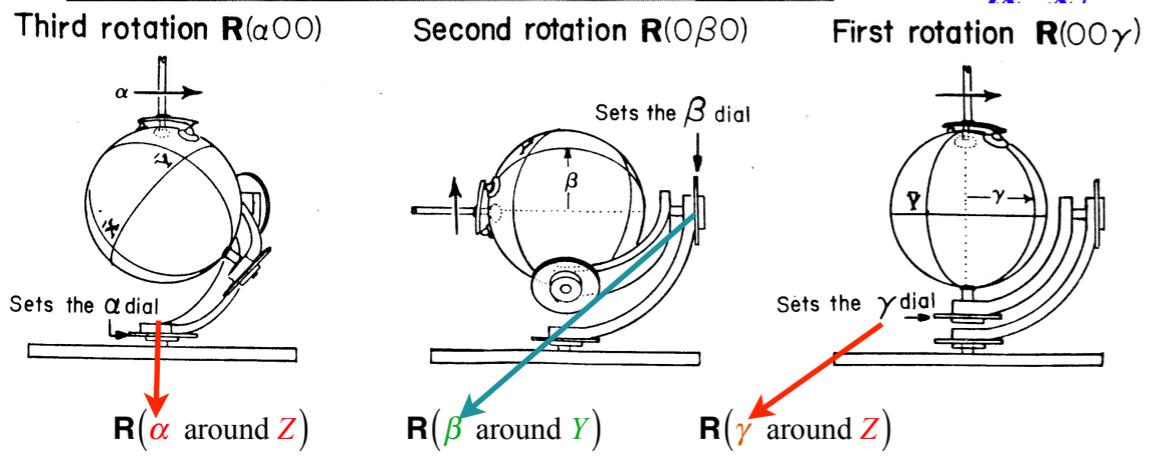
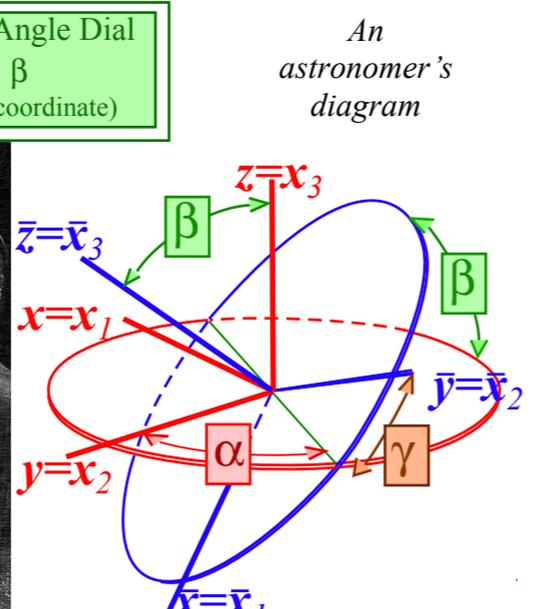
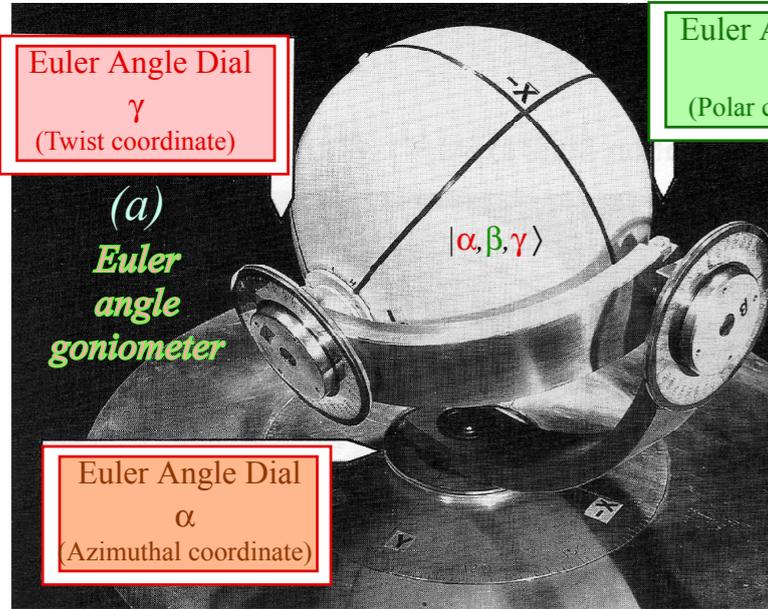


$$R[\vec{\Theta}] = \begin{pmatrix} \cos\frac{\Theta}{2} - i\hat{\Theta}_Z \sin\frac{\Theta}{2} & -i\sin\frac{\Theta}{2}(\hat{\Theta}_X - i\hat{\Theta}_Y) \\ -i\sin\frac{\Theta}{2}(\hat{\Theta}_X + i\hat{\Theta}_Y) & \cos\frac{\Theta}{2} + i\hat{\Theta}_Z \sin\frac{\Theta}{2} \end{pmatrix} = R[\varphi\vartheta\Theta] = e^{-i\mathbf{H}t}$$

$$= \cos\frac{\Theta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \underbrace{\hat{\Theta}_X \sin\frac{\Theta}{2}}_{\cos\varphi \sin\vartheta} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \underbrace{\hat{\Theta}_Y \sin\frac{\Theta}{2}}_{\sin\varphi \sin\vartheta} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \underbrace{\hat{\Theta}_Z \sin\frac{\Theta}{2}}_{\cos\vartheta}$$

$$= \begin{pmatrix} \cos\frac{\Theta}{2} - i\cos\vartheta \sin\frac{\Theta}{2} & -\sin\frac{\Theta}{2}(\sin\varphi \sin\vartheta + i\cos\varphi \sin\vartheta) \\ \sin\frac{\Theta}{2}(\sin\varphi \sin\vartheta - i\cos\varphi \sin\vartheta) & \cos\frac{\Theta}{2} + i\cos\vartheta \sin\frac{\Theta}{2} \end{pmatrix}$$

# Euler $R(\alpha\beta\gamma)$ versus Darboux $R[\varphi\vartheta\Theta]$

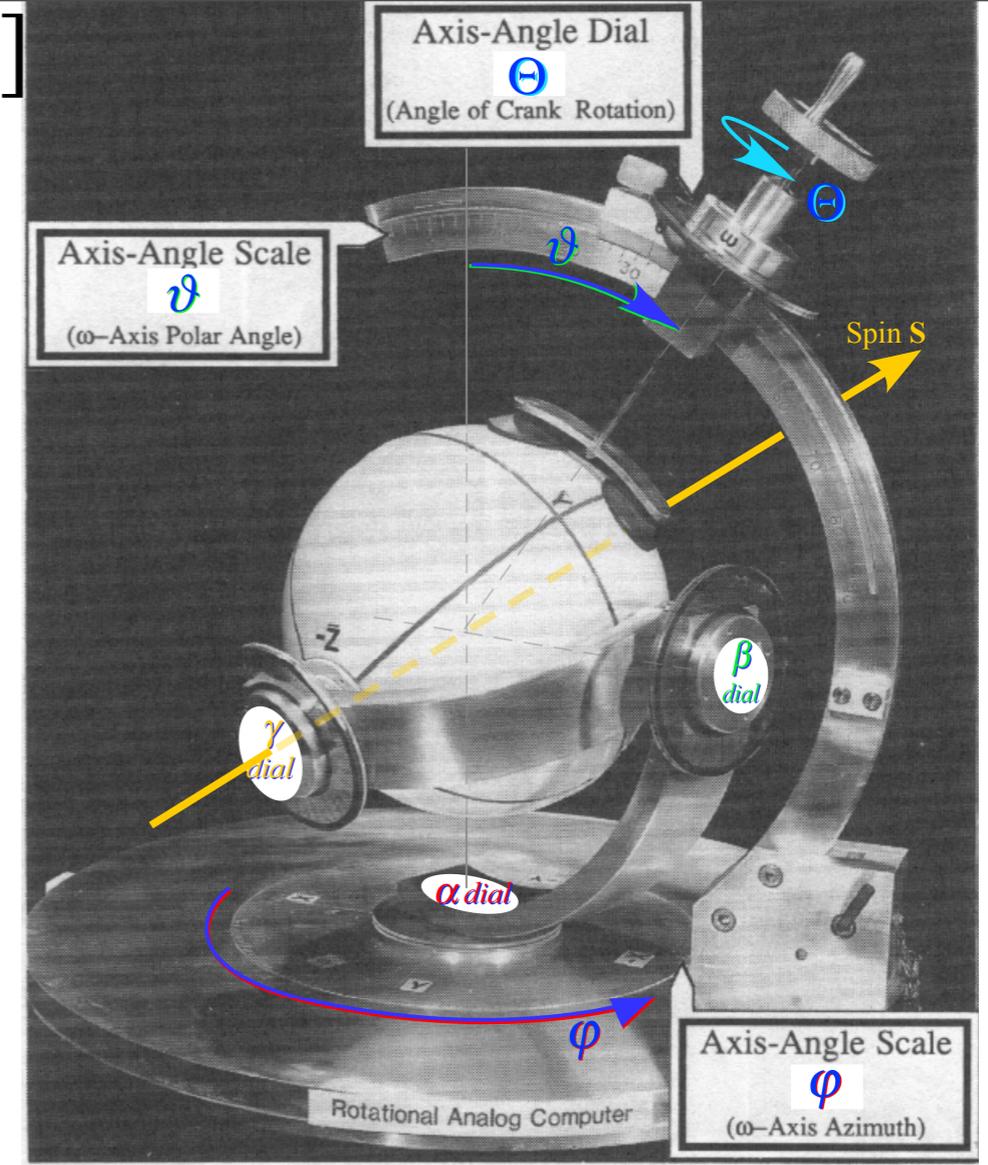


$$R(\alpha\beta\gamma) = \begin{pmatrix} e^{-i\frac{\alpha}{2}} & 0 \\ 0 & e^{i\frac{\alpha}{2}} \end{pmatrix} \begin{pmatrix} \cos\frac{\beta}{2} & -\sin\frac{\beta}{2} \\ \sin\frac{\beta}{2} & \cos\frac{\beta}{2} \end{pmatrix} \begin{pmatrix} e^{-i\frac{\gamma}{2}} & 0 \\ 0 & e^{i\frac{\gamma}{2}} \end{pmatrix} = \begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} \end{pmatrix}$$

$$= \cos\frac{\alpha+\gamma}{2} \cos\frac{\beta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \sin\frac{\gamma-\alpha}{2} \sin\frac{\beta}{2} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \cos\frac{\gamma-\alpha}{2} \sin\frac{\beta}{2} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \sin\frac{\alpha+\gamma}{2} \cos\frac{\beta}{2}$$

Euler  $R(\alpha\beta\gamma)$  is simpler to form than  $\Theta$ -axis Darboux  $R[\varphi\vartheta\Theta]$ .

Euler *state definition*:  
 $|\alpha\beta\gamma\rangle = R(\alpha\beta\gamma)|000\rangle$  ( $\alpha\beta\gamma$  make better coordinates)

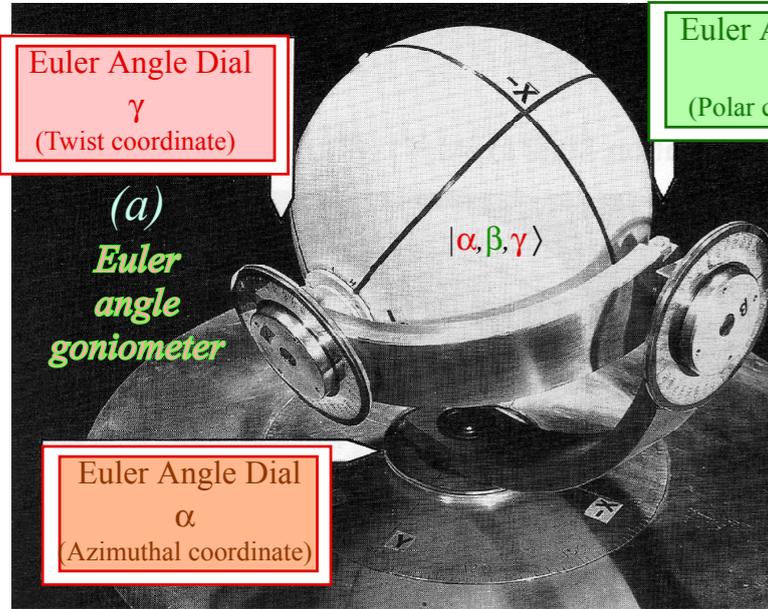


$$R[\bar{\Theta}] = \begin{pmatrix} \cos\frac{\Theta}{2} - i\hat{\Theta}_Z \sin\frac{\Theta}{2} & -i\sin\frac{\Theta}{2}(\hat{\Theta}_X - i\hat{\Theta}_Y) \\ -i\sin\frac{\Theta}{2}(\hat{\Theta}_X + i\hat{\Theta}_Y) & \cos\frac{\Theta}{2} + i\hat{\Theta}_Z \sin\frac{\Theta}{2} \end{pmatrix} = R[\varphi\vartheta\Theta] = e^{-iHt}$$

$$= \cos\frac{\Theta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \underbrace{\hat{\Theta}_X \sin\frac{\Theta}{2}}_{\cos\varphi \sin\vartheta} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \underbrace{\hat{\Theta}_Y \sin\frac{\Theta}{2}}_{\sin\varphi \sin\vartheta} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \underbrace{\hat{\Theta}_Z \sin\frac{\Theta}{2}}_{\cos\vartheta}$$

$$= \begin{pmatrix} \cos\frac{\Theta}{2} - i\cos\vartheta \sin\frac{\Theta}{2} & -\sin\frac{\Theta}{2}(\sin\varphi \sin\vartheta + i\cos\varphi \sin\vartheta) \\ \sin\frac{\Theta}{2}(\sin\varphi \sin\vartheta - i\cos\varphi \sin\vartheta) & \cos\frac{\Theta}{2} + i\cos\vartheta \sin\frac{\Theta}{2} \end{pmatrix}$$

# Euler $R(\alpha\beta\gamma)$ versus Darboux $R[\varphi\vartheta\Theta]$

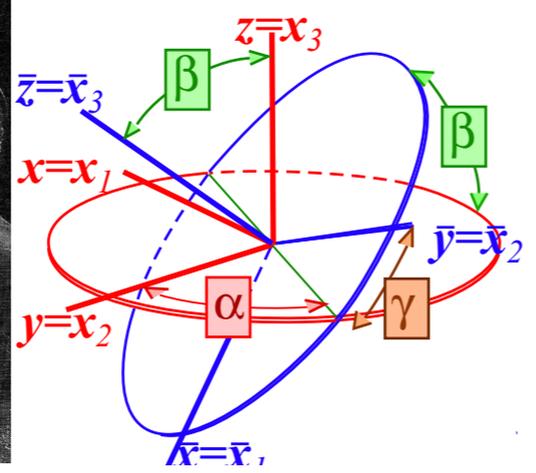


Euler Angle Dial  $\beta$   
(Polar coordinate)

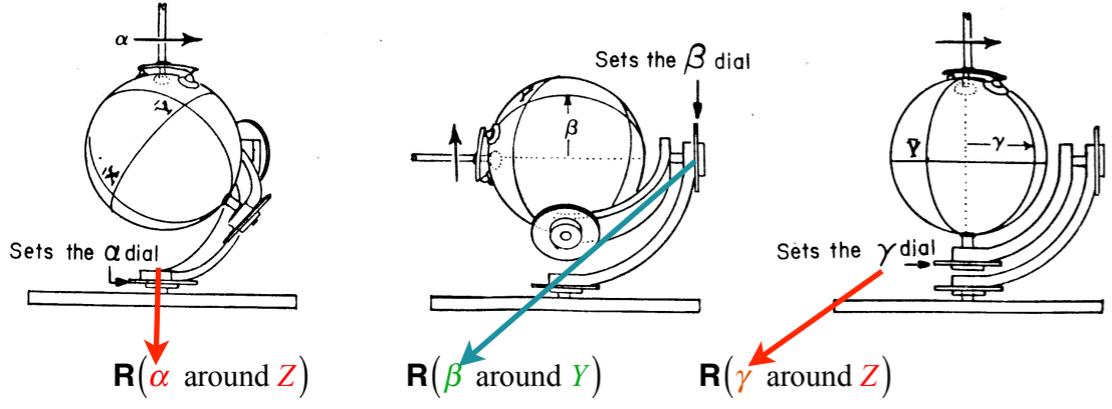
Euler Angle Dial  $\gamma$   
(Twist coordinate)

Euler Angle Dial  $\alpha$   
(Azimuthal coordinate)

An astronomer's diagram



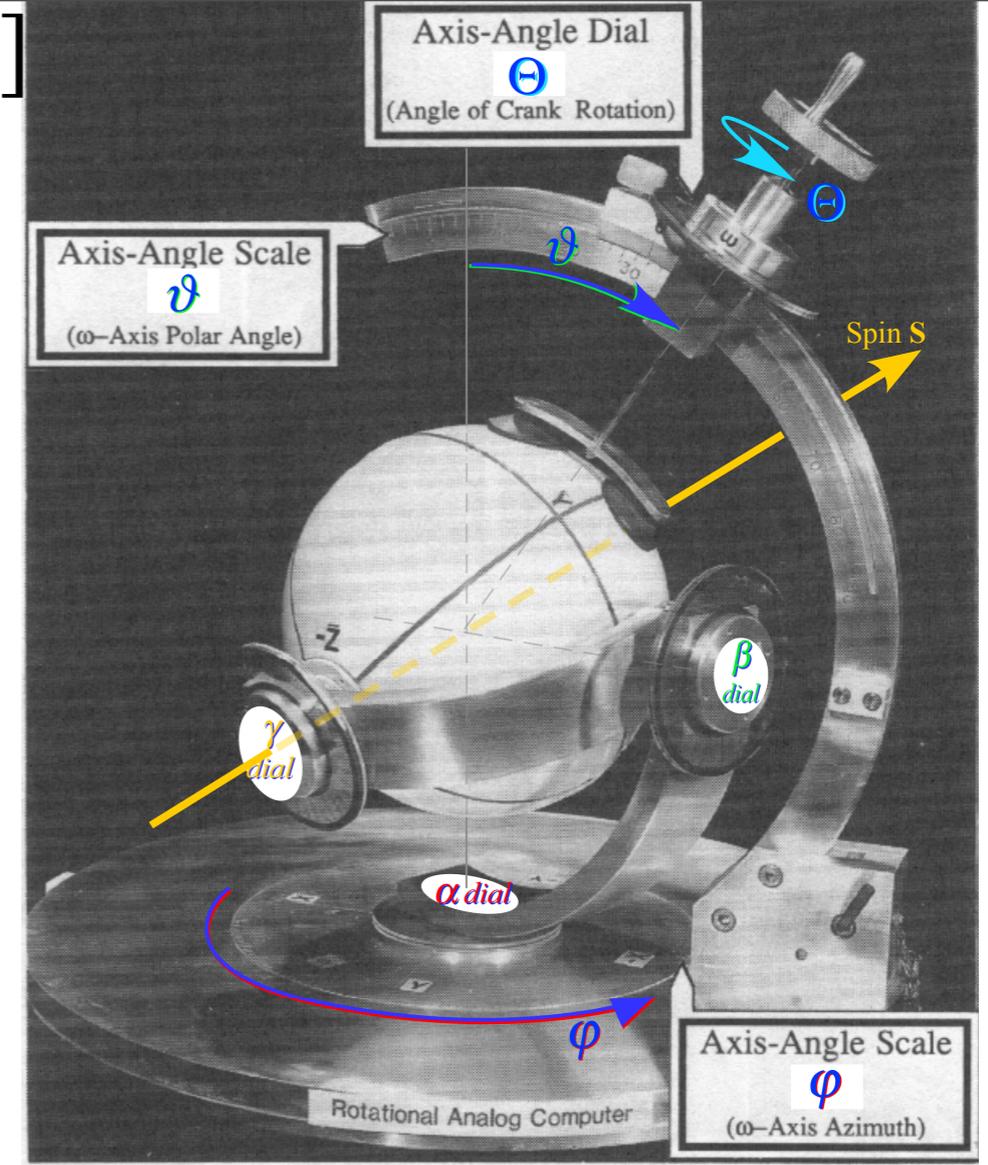
Third rotation  $R(\alpha 0 0)$     Second rotation  $R(0 \beta 0)$     First rotation  $R(0 0 \gamma)$



$$R(\alpha\beta\gamma) = \begin{pmatrix} e^{-i\frac{\alpha}{2}} & 0 \\ 0 & e^{i\frac{\alpha}{2}} \end{pmatrix} \begin{pmatrix} \cos\frac{\beta}{2} & -\sin\frac{\beta}{2} \\ \sin\frac{\beta}{2} & \cos\frac{\beta}{2} \end{pmatrix} \begin{pmatrix} e^{-i\frac{\gamma}{2}} & 0 \\ 0 & e^{i\frac{\gamma}{2}} \end{pmatrix} = \begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} \end{pmatrix}$$

$$= \cos\frac{\alpha+\gamma}{2} \cos\frac{\beta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \sin\frac{\gamma-\alpha}{2} \sin\frac{\beta}{2} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \cos\frac{\gamma-\alpha}{2} \sin\frac{\beta}{2} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \sin\frac{\alpha+\gamma}{2} \cos\frac{\beta}{2}$$

Euler  $R(\alpha\beta\gamma)$  is simpler to form than  $\Theta$ -axis Darboux  $R[\varphi\vartheta\Theta]$ .  
Euler *state definition* lets us relate  $R(\alpha\beta\gamma)$  to  $R[\varphi\vartheta\Theta]$  ...  
 $|\alpha\beta\gamma\rangle = R(\alpha\beta\gamma)|000\rangle$  ( $\alpha\beta\gamma$  make better coordinates)

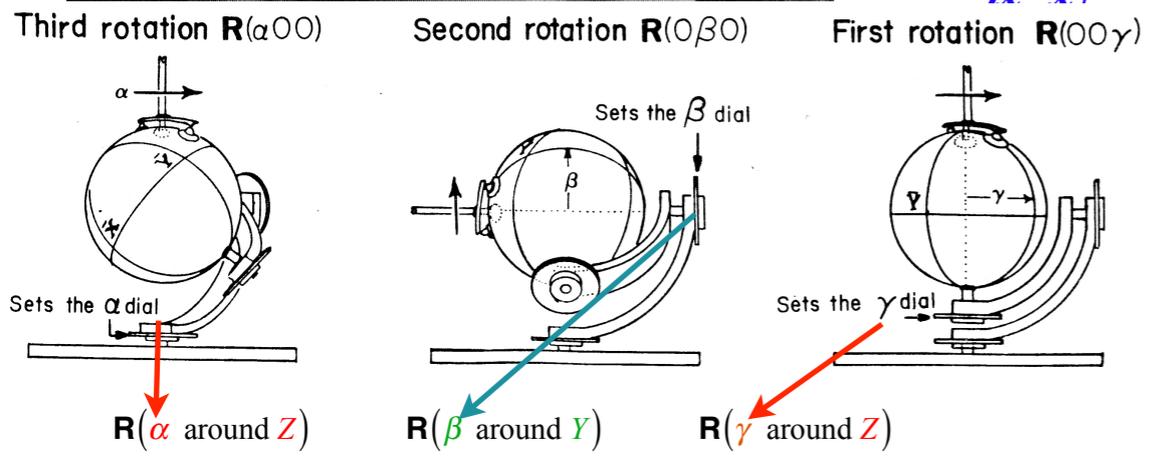
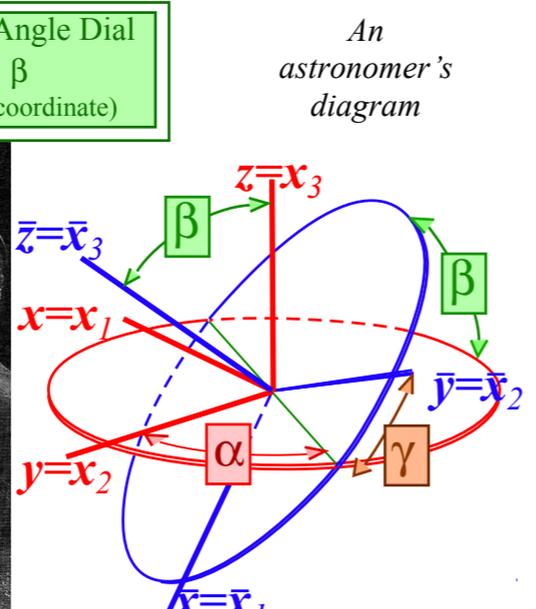
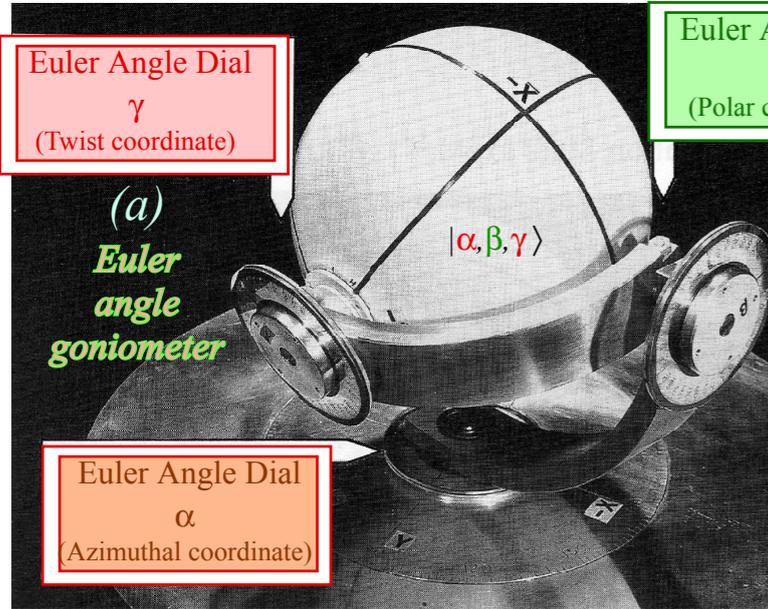


$$R[\bar{\Theta}] = \begin{pmatrix} \cos\frac{\Theta}{2} - i\hat{\Theta}_Z \sin\frac{\Theta}{2} & -i\sin\frac{\Theta}{2}(\hat{\Theta}_X - i\hat{\Theta}_Y) \\ -i\sin\frac{\Theta}{2}(\hat{\Theta}_X + i\hat{\Theta}_Y) & \cos\frac{\Theta}{2} + i\hat{\Theta}_Z \sin\frac{\Theta}{2} \end{pmatrix} = R[\varphi\vartheta\Theta] = e^{-iHt}$$

$$= \cos\frac{\Theta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \underbrace{\hat{\Theta}_X \sin\frac{\Theta}{2}}_{\cos\varphi \sin\vartheta} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \underbrace{\hat{\Theta}_Y \sin\frac{\Theta}{2}}_{\sin\varphi \sin\vartheta} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \underbrace{\hat{\Theta}_Z \sin\frac{\Theta}{2}}_{\cos\vartheta}$$

$$= \begin{pmatrix} \cos\frac{\Theta}{2} - i\cos\vartheta \sin\frac{\Theta}{2} & -\sin\frac{\Theta}{2}(\sin\varphi \sin\vartheta + i\cos\varphi \sin\vartheta) \\ \sin\frac{\Theta}{2}(\sin\varphi \sin\vartheta - i\cos\varphi \sin\vartheta) & \cos\frac{\Theta}{2} + i\cos\vartheta \sin\frac{\Theta}{2} \end{pmatrix}$$

# Euler $R(\alpha\beta\gamma)$ versus Darboux $R[\varphi\vartheta\Theta]$

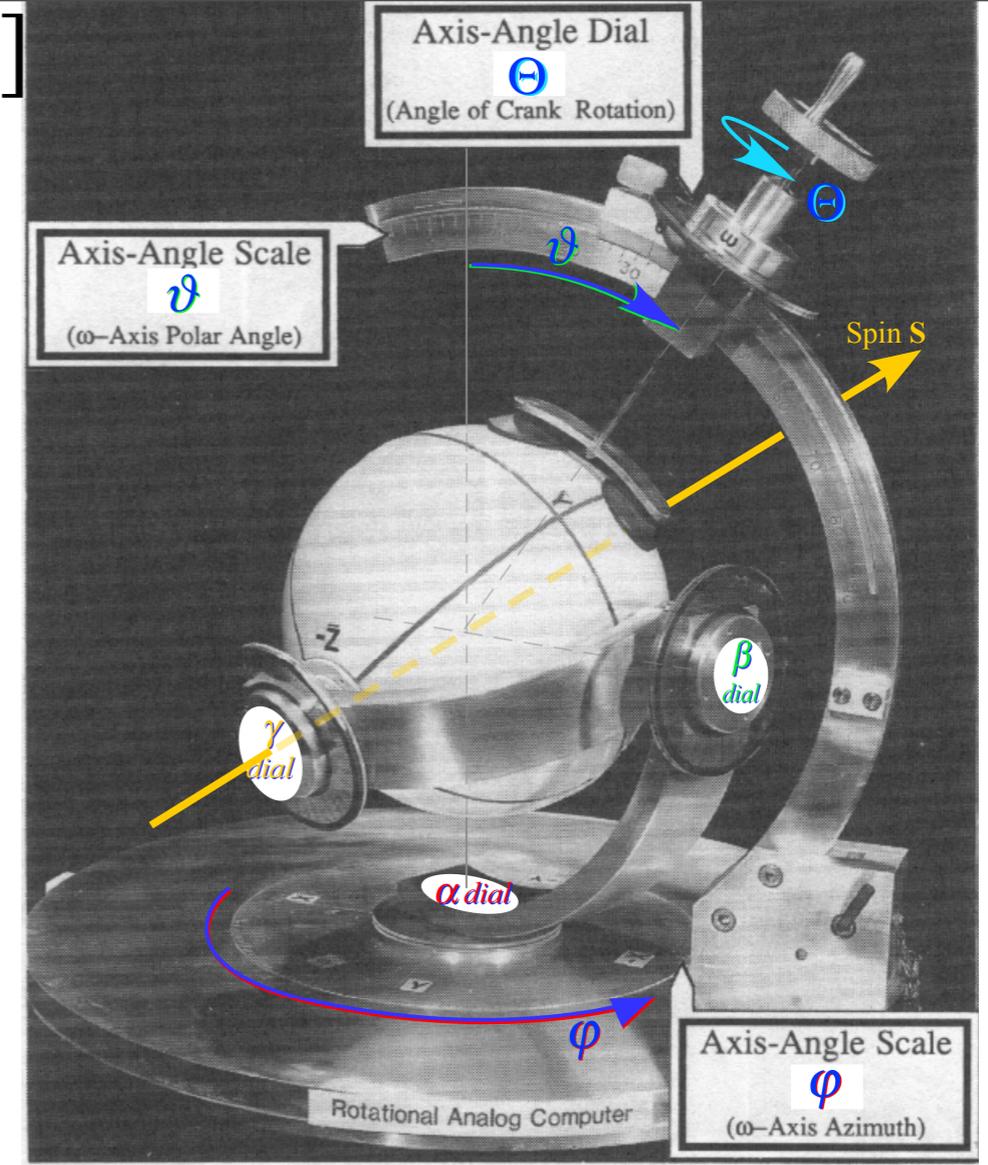


$$R(\alpha\beta\gamma) = \begin{pmatrix} e^{-i\frac{\alpha}{2}} & 0 \\ 0 & e^{i\frac{\alpha}{2}} \end{pmatrix} \begin{pmatrix} \cos\frac{\beta}{2} & -\sin\frac{\beta}{2} \\ \sin\frac{\beta}{2} & \cos\frac{\beta}{2} \end{pmatrix} \begin{pmatrix} e^{-i\frac{\gamma}{2}} & 0 \\ 0 & e^{i\frac{\gamma}{2}} \end{pmatrix} = \begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} \end{pmatrix}$$

$$= \cos\frac{\alpha+\gamma}{2} \cos\frac{\beta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \sin\frac{\gamma-\alpha}{2} \sin\frac{\beta}{2} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \cos\frac{\gamma-\alpha}{2} \sin\frac{\beta}{2} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \sin\frac{\alpha+\gamma}{2} \cos\frac{\beta}{2}$$

Euler  $R(\alpha\beta\gamma)$  is simpler to form than  $\Theta$ -axis Darboux  $R[\varphi\vartheta\Theta]$ .  
 Euler *state definition* lets us relate  $R(\alpha\beta\gamma)$  to  $R[\varphi\vartheta\Theta]$  ...  
 $|\alpha\beta\gamma\rangle = R(\alpha\beta\gamma)|000\rangle$  ( $\alpha\beta\gamma$  make better coordinates)

$$\begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} x_1 + ip_1 \\ x_2 + ip_2 \end{pmatrix}$$

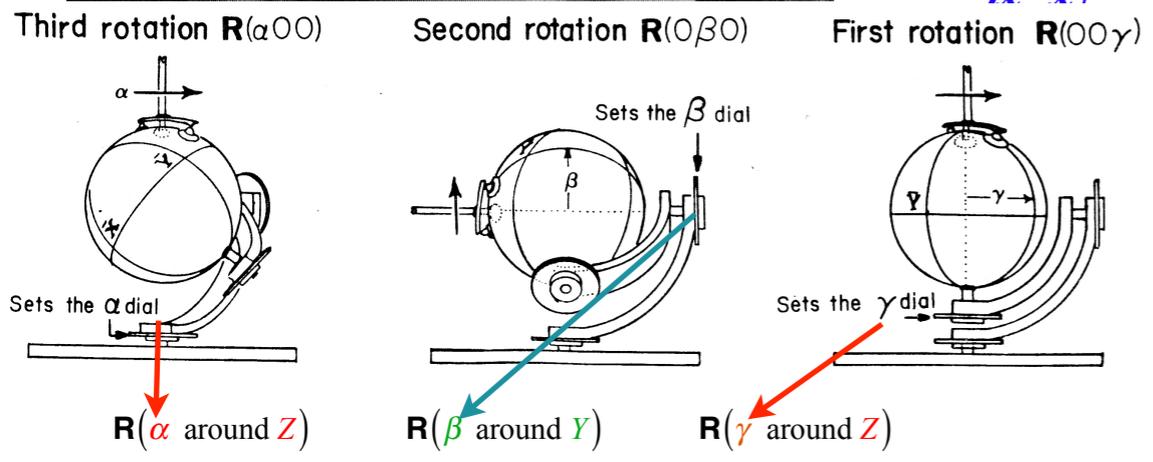
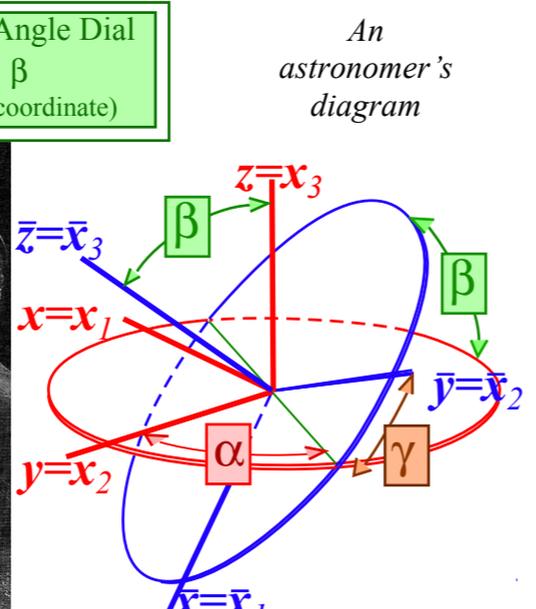
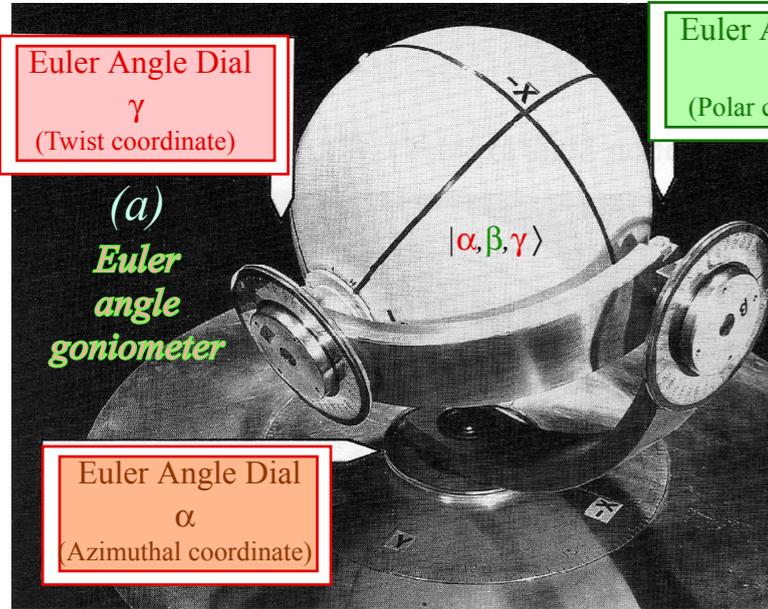


$$R[\bar{\Theta}] = \begin{pmatrix} \cos\frac{\Theta}{2} - i\hat{\Theta}_Z \sin\frac{\Theta}{2} & -i\sin\frac{\Theta}{2}(\hat{\Theta}_X - i\hat{\Theta}_Y) \\ -i\sin\frac{\Theta}{2}(\hat{\Theta}_X + i\hat{\Theta}_Y) & \cos\frac{\Theta}{2} + i\hat{\Theta}_Z \sin\frac{\Theta}{2} \end{pmatrix} = R[\varphi\vartheta\Theta] = e^{-iHt}$$

$$= \cos\frac{\Theta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \left( \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \hat{\Theta}_X \sin\frac{\Theta}{2} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \hat{\Theta}_Y \sin\frac{\Theta}{2} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \hat{\Theta}_Z \sin\frac{\Theta}{2} \right)$$

$$= \begin{pmatrix} \cos\frac{\Theta}{2} - i\cos\vartheta \sin\frac{\Theta}{2} & -\sin\frac{\Theta}{2}(\sin\varphi \sin\vartheta + i\cos\varphi \sin\vartheta) \\ \sin\frac{\Theta}{2}(\sin\varphi \sin\vartheta - i\cos\varphi \sin\vartheta) & \cos\frac{\Theta}{2} + i\cos\vartheta \sin\frac{\Theta}{2} \end{pmatrix}$$

# Euler $R(\alpha\beta\gamma)$ versus Darboux $R[\varphi\vartheta\Theta]$



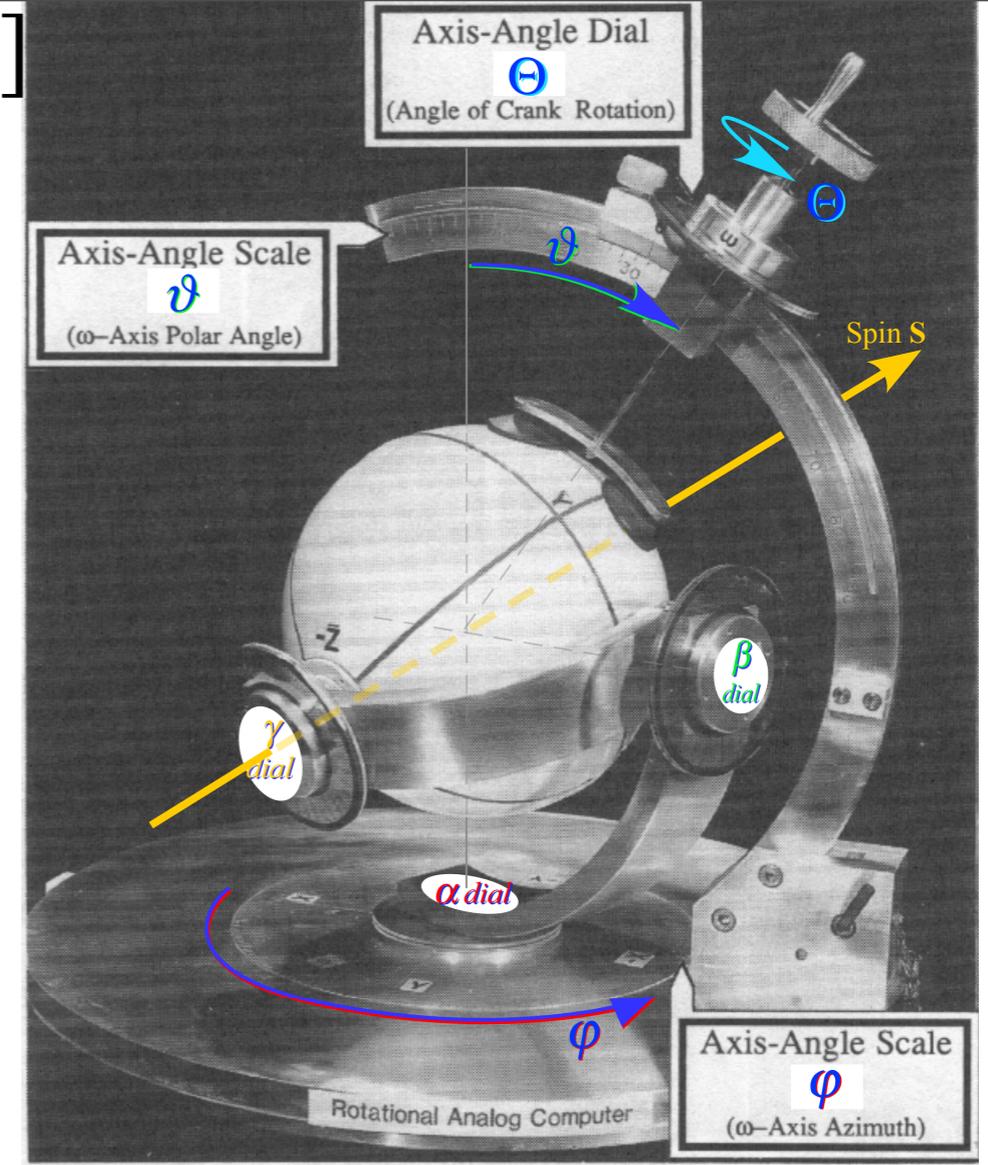
$$R(\alpha\beta\gamma) = \begin{pmatrix} e^{-i\frac{\alpha}{2}} & 0 \\ 0 & e^{i\frac{\alpha}{2}} \end{pmatrix} \begin{pmatrix} \cos\frac{\beta}{2} & -\sin\frac{\beta}{2} \\ \sin\frac{\beta}{2} & \cos\frac{\beta}{2} \end{pmatrix} \begin{pmatrix} e^{-i\frac{\gamma}{2}} & 0 \\ 0 & e^{i\frac{\gamma}{2}} \end{pmatrix} = \begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} \end{pmatrix}$$

$$= \cos\frac{\alpha+\gamma}{2} \cos\frac{\beta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \sin\frac{\gamma-\alpha}{2} \sin\frac{\beta}{2} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \cos\frac{\gamma-\alpha}{2} \sin\frac{\beta}{2} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \sin\frac{\alpha+\gamma}{2} \cos\frac{\beta}{2}$$

Euler  $R(\alpha\beta\gamma)$  is simpler to form than  $\Theta$ -axis Darboux  $R[\varphi\vartheta\Theta]$ .  
Euler *state definition* lets us relate  $R(\alpha\beta\gamma)$  to  $R[\varphi\vartheta\Theta]$  ...  
 $|\alpha\beta\gamma\rangle = R(\alpha\beta\gamma)|000\rangle$  ( $\alpha\beta\gamma$  make better coordinates)

$$\begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} x_1+ip_1 \\ x_2+ip_2 \end{pmatrix}$$

$x_1 = \cos[(\gamma+\alpha)/2] \cos\beta/2 = \cos\Theta/2$

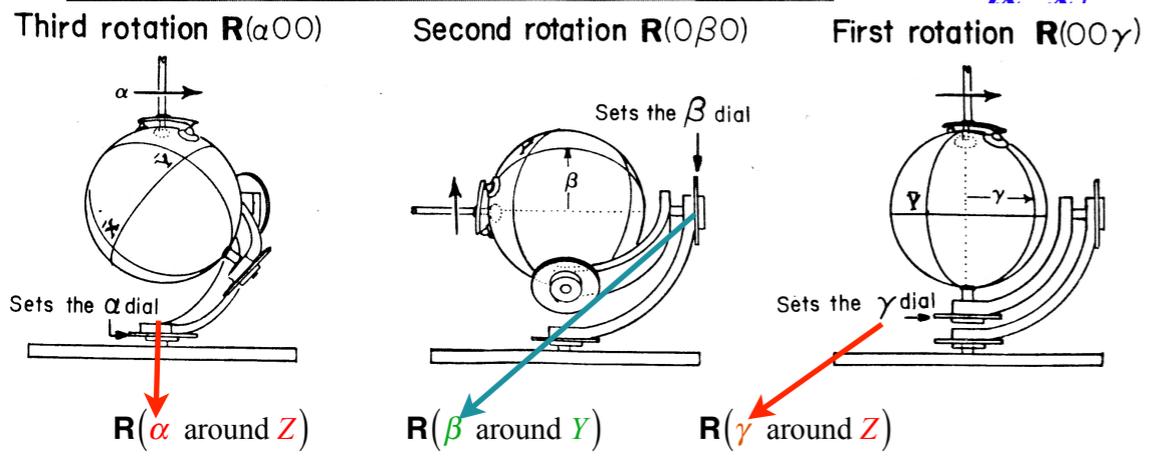
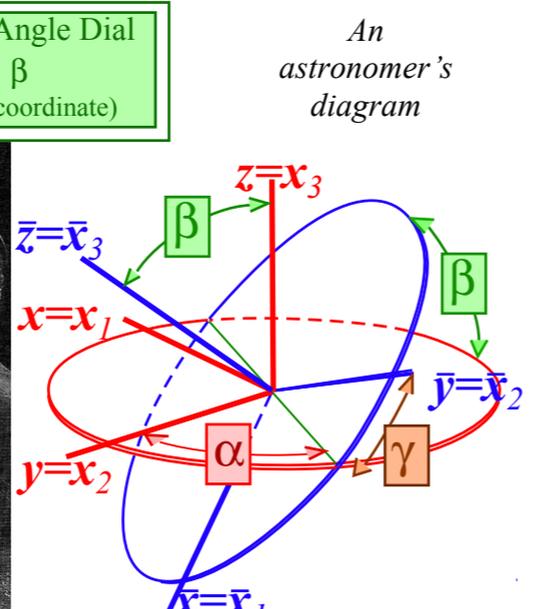
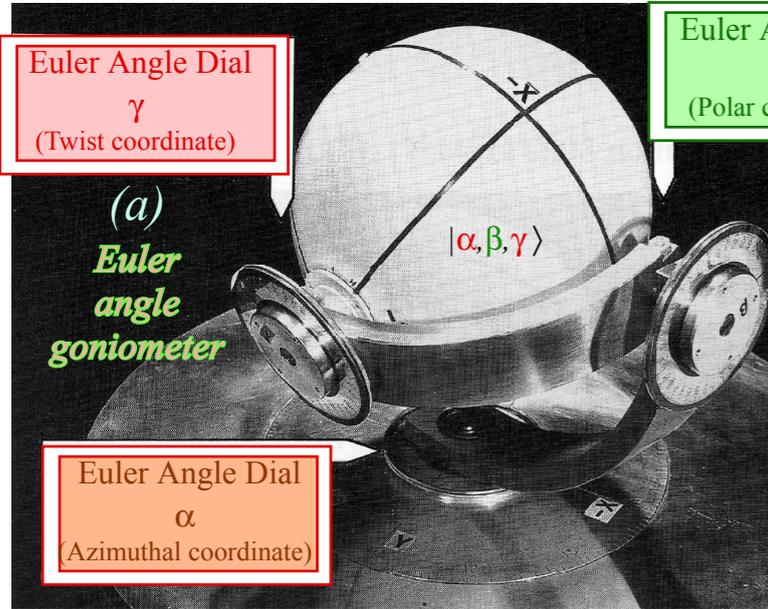


$$R[\vec{\Theta}] = \begin{pmatrix} \cos\frac{\Theta}{2} - i\hat{\Theta}_Z \sin\frac{\Theta}{2} & -i\sin\frac{\Theta}{2}(\hat{\Theta}_X - i\hat{\Theta}_Y) \\ -i\sin\frac{\Theta}{2}(\hat{\Theta}_X + i\hat{\Theta}_Y) & \cos\frac{\Theta}{2} + i\hat{\Theta}_Z \sin\frac{\Theta}{2} \end{pmatrix} = R[\varphi\vartheta\Theta] = e^{-iHt}$$

$$= \cos\frac{\Theta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \hat{\Theta}_X \sin\frac{\Theta}{2} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \hat{\Theta}_Y \sin\frac{\Theta}{2} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \hat{\Theta}_Z \sin\frac{\Theta}{2}$$

$$= \begin{pmatrix} \cos\frac{\Theta}{2} - i\cos\vartheta \sin\frac{\Theta}{2} & -\sin\frac{\Theta}{2}(\sin\varphi \sin\vartheta + i\cos\varphi \sin\vartheta) \\ \sin\frac{\Theta}{2}(\sin\varphi \sin\vartheta - i\cos\varphi \sin\vartheta) & \cos\frac{\Theta}{2} + i\cos\vartheta \sin\frac{\Theta}{2} \end{pmatrix}$$

# Euler $R(\alpha\beta\gamma)$ versus Darboux $R[\varphi\vartheta\Theta]$



$$R(\alpha\beta\gamma) = \begin{pmatrix} e^{-i\frac{\alpha}{2}} & 0 \\ 0 & e^{i\frac{\alpha}{2}} \end{pmatrix} \begin{pmatrix} \cos\frac{\beta}{2} & -\sin\frac{\beta}{2} \\ \sin\frac{\beta}{2} & \cos\frac{\beta}{2} \end{pmatrix} \begin{pmatrix} e^{-i\frac{\gamma}{2}} & 0 \\ 0 & e^{i\frac{\gamma}{2}} \end{pmatrix} = \begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} \end{pmatrix}$$

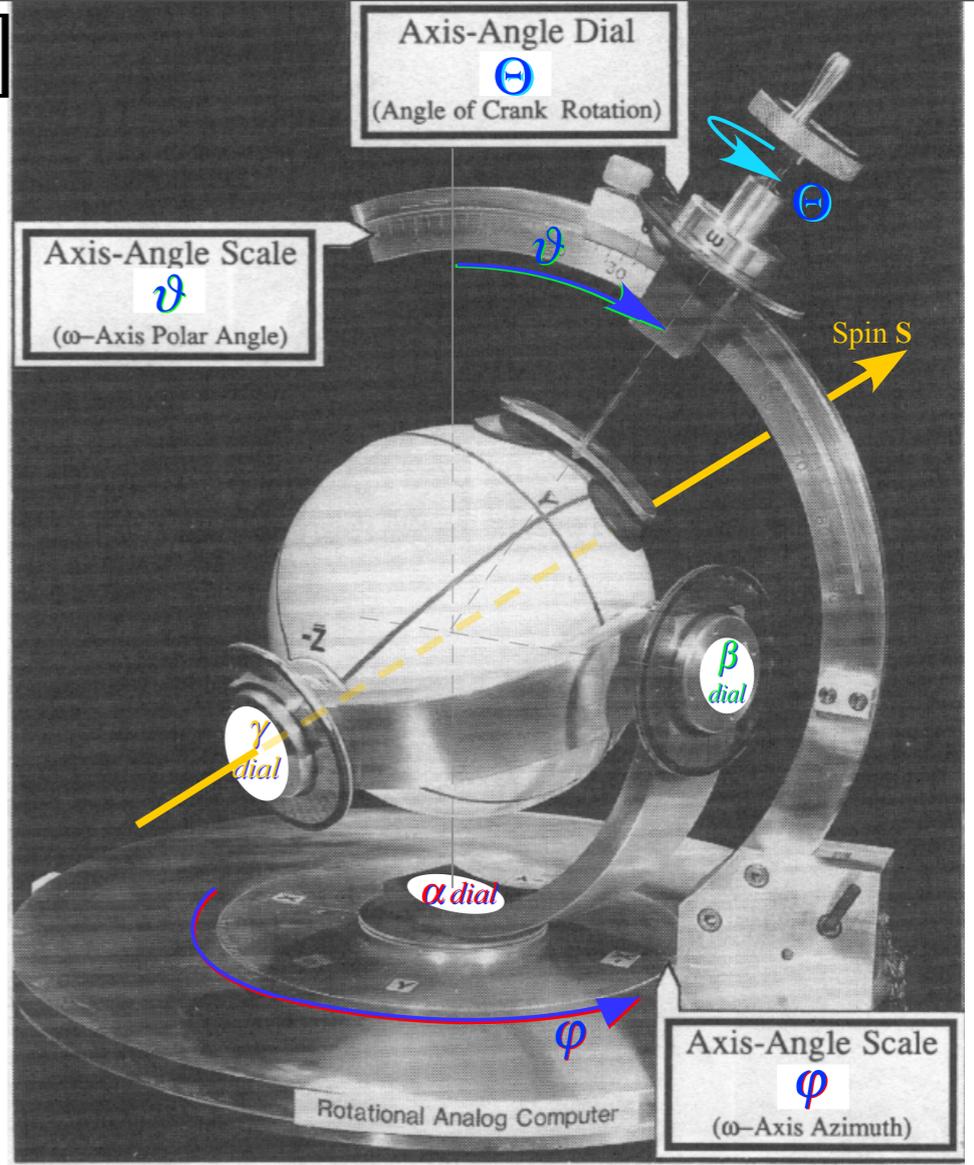
$$= \cos\frac{\alpha+\gamma}{2} \cos\frac{\beta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \sin\frac{\gamma-\alpha}{2} \sin\frac{\beta}{2} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \cos\frac{\gamma-\alpha}{2} \sin\frac{\beta}{2} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \sin\frac{\alpha+\gamma}{2} \cos\frac{\beta}{2}$$

Euler  $R(\alpha\beta\gamma)$  is simpler to form than  $\Theta$ -axis Darboux  $R[\varphi\vartheta\Theta]$ .  
 Euler *state definition* lets us relate  $R(\alpha\beta\gamma)$  to  $R[\varphi\vartheta\Theta]$  ...  
 $|\alpha\beta\gamma\rangle = R(\alpha\beta\gamma)|000\rangle$  ( $\alpha\beta\gamma$  make better coordinates)

$$\begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} x_1+ip_1 \\ x_2+ip_2 \end{pmatrix}$$

$$x_1 = \cos[(\gamma+\alpha)/2] \cos\beta/2 = \cos\Theta/2$$

$$-p_2 = \sin[(\gamma-\alpha)/2] \sin\beta/2 = \hat{\Theta}_X \sin\Theta/2 = \cos\varphi \sin\vartheta \sin\Theta/2$$



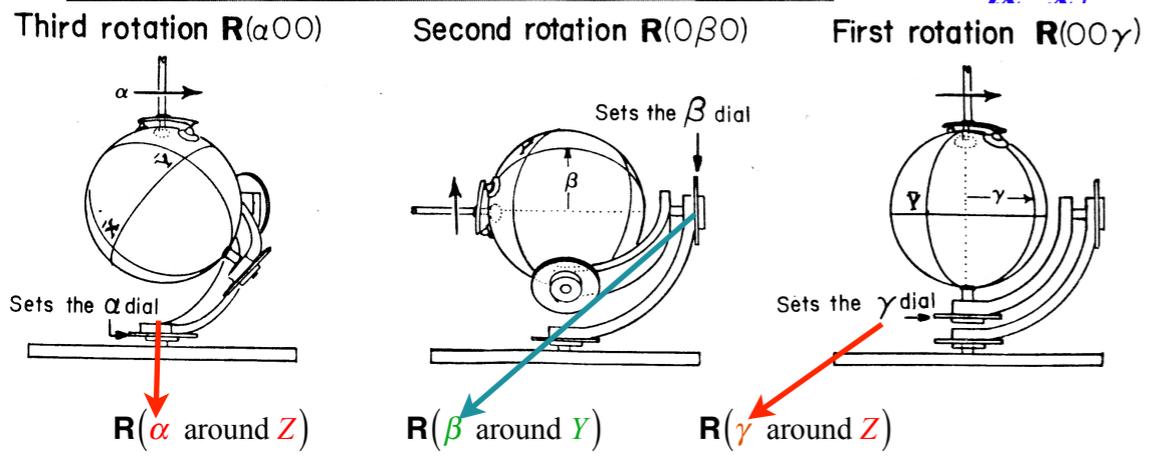
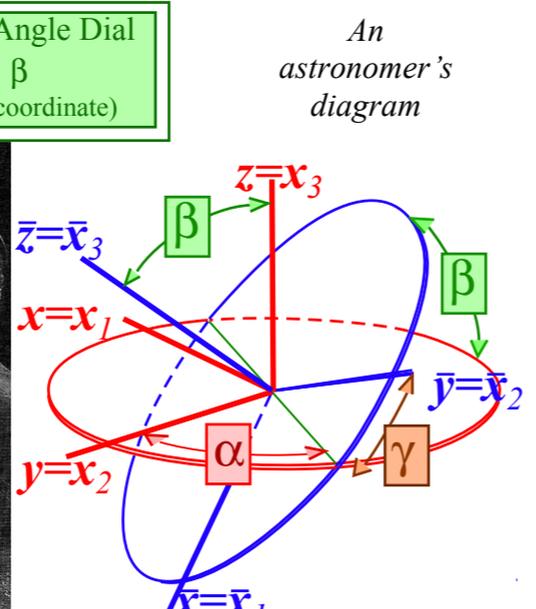
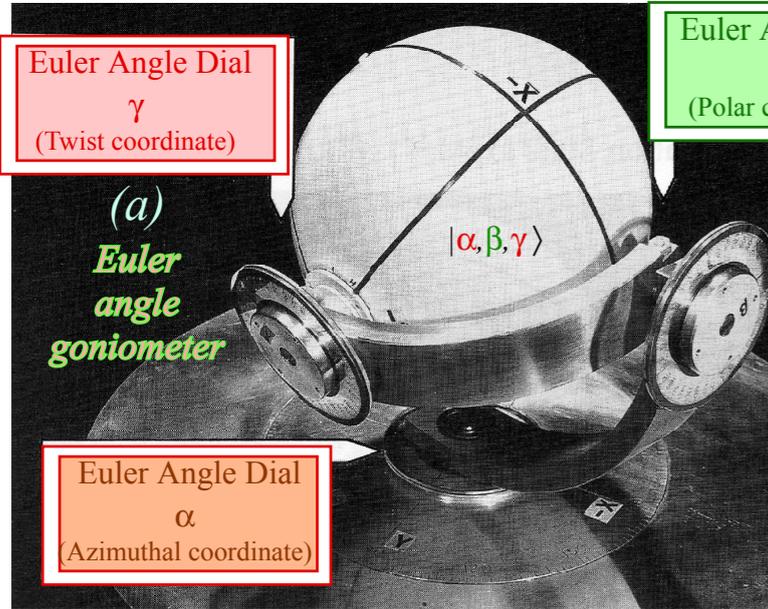
$$R[\vec{\Theta}] = \begin{pmatrix} \cos\frac{\Theta}{2} - i\hat{\Theta}_Z \sin\frac{\Theta}{2} & -i\sin\frac{\Theta}{2}(\hat{\Theta}_X - i\hat{\Theta}_Y) \\ -i\sin\frac{\Theta}{2}(\hat{\Theta}_X + i\hat{\Theta}_Y) & \cos\frac{\Theta}{2} + i\hat{\Theta}_Z \sin\frac{\Theta}{2} \end{pmatrix} = R[\varphi\vartheta\Theta] = e^{-iHt}$$

$$= \cos\frac{\Theta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \hat{\Theta}_X \sin\frac{\Theta}{2} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \hat{\Theta}_Y \sin\frac{\Theta}{2} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \hat{\Theta}_Z \sin\frac{\Theta}{2}$$

$$= \begin{pmatrix} \cos\frac{\Theta}{2} - i\cos\vartheta \sin\frac{\Theta}{2} & -\sin\frac{\Theta}{2}(\sin\varphi \sin\vartheta + i\cos\varphi \sin\vartheta) \\ \sin\frac{\Theta}{2}(\sin\varphi \sin\vartheta - i\cos\varphi \sin\vartheta) & \cos\frac{\Theta}{2} + i\cos\vartheta \sin\frac{\Theta}{2} \end{pmatrix}$$

$$\cos\varphi \sin\vartheta \sin\Theta/2$$

# Euler $R(\alpha\beta\gamma)$ versus Darboux $R[\varphi\vartheta\Theta]$



$$R(\alpha\beta\gamma) = \begin{pmatrix} e^{-i\frac{\alpha}{2}} & 0 \\ 0 & e^{i\frac{\alpha}{2}} \end{pmatrix} \begin{pmatrix} \cos\frac{\beta}{2} & -\sin\frac{\beta}{2} \\ \sin\frac{\beta}{2} & \cos\frac{\beta}{2} \end{pmatrix} \begin{pmatrix} e^{-i\frac{\gamma}{2}} & 0 \\ 0 & e^{i\frac{\gamma}{2}} \end{pmatrix} = \begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} \end{pmatrix}$$

$$= \cos\frac{\alpha+\gamma}{2} \cos\frac{\beta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \sin\frac{\gamma-\alpha}{2} \sin\frac{\beta}{2} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \cos\frac{\gamma-\alpha}{2} \sin\frac{\beta}{2} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \sin\frac{\alpha+\gamma}{2} \cos\frac{\beta}{2}$$

Euler  $R(\alpha\beta\gamma)$  is simpler to form than  $\Theta$ -axis Darboux  $R[\varphi\vartheta\Theta]$ .

Euler *state definition* lets us relate  $R(\alpha\beta\gamma)$  to  $R[\varphi\vartheta\Theta]$  ...

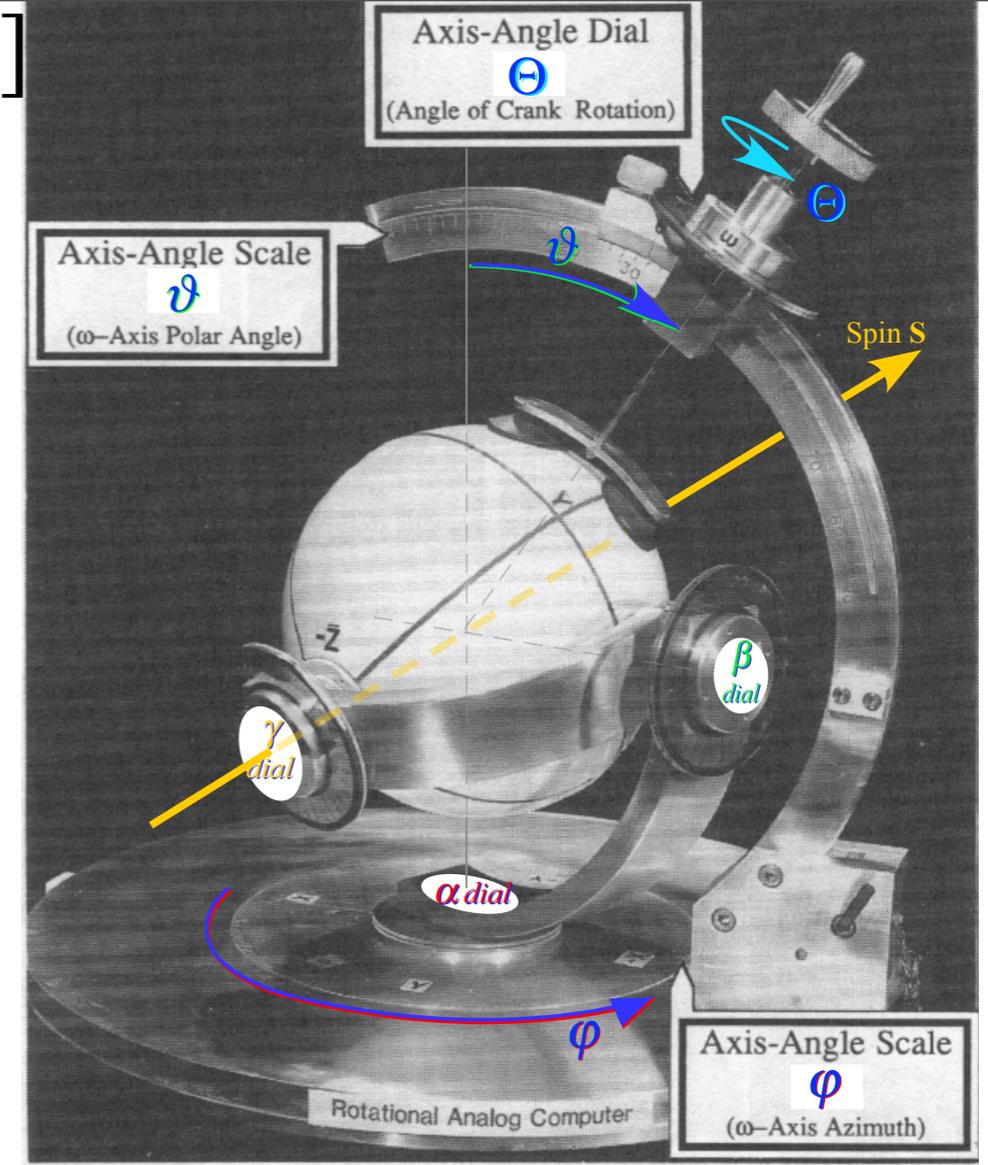
$|\alpha\beta\gamma\rangle = R(\alpha\beta\gamma)|000\rangle$  ( $\alpha\beta\gamma$  make better coordinates)

$$\begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} x_1+ip_1 \\ x_2+ip_2 \end{pmatrix}$$

$x_1 = \cos[(\gamma+\alpha)/2] \cos\beta/2 = \cos\Theta/2$

$-p_2 = \sin[(\gamma-\alpha)/2] \sin\beta/2 = \hat{\Theta}_X \sin\Theta/2 = \cos\varphi \sin\vartheta \sin\Theta/2$

$x_2 = \cos[(\gamma-\alpha)/2] \sin\beta/2 = \hat{\Theta}_Y \sin\Theta/2 = \sin\varphi \sin\vartheta \sin\Theta/2$

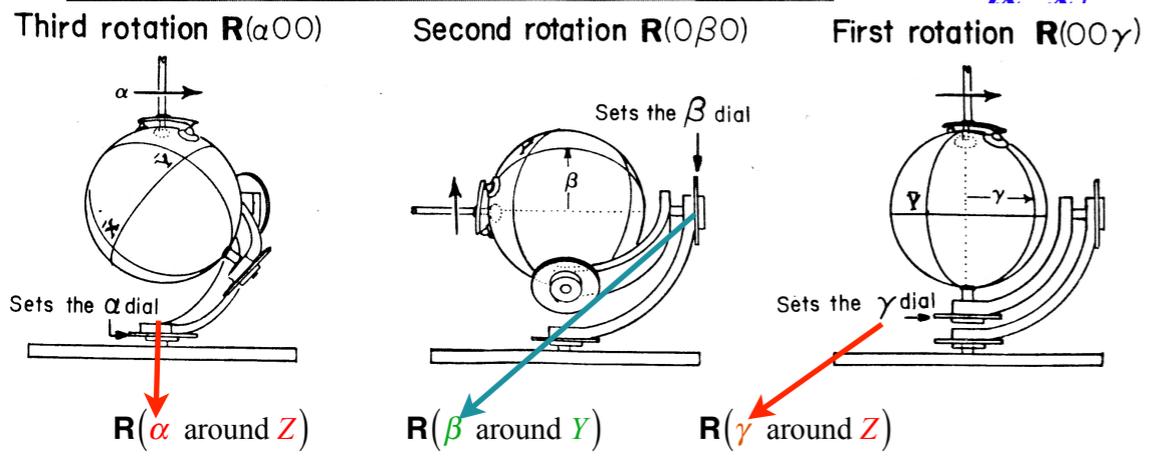
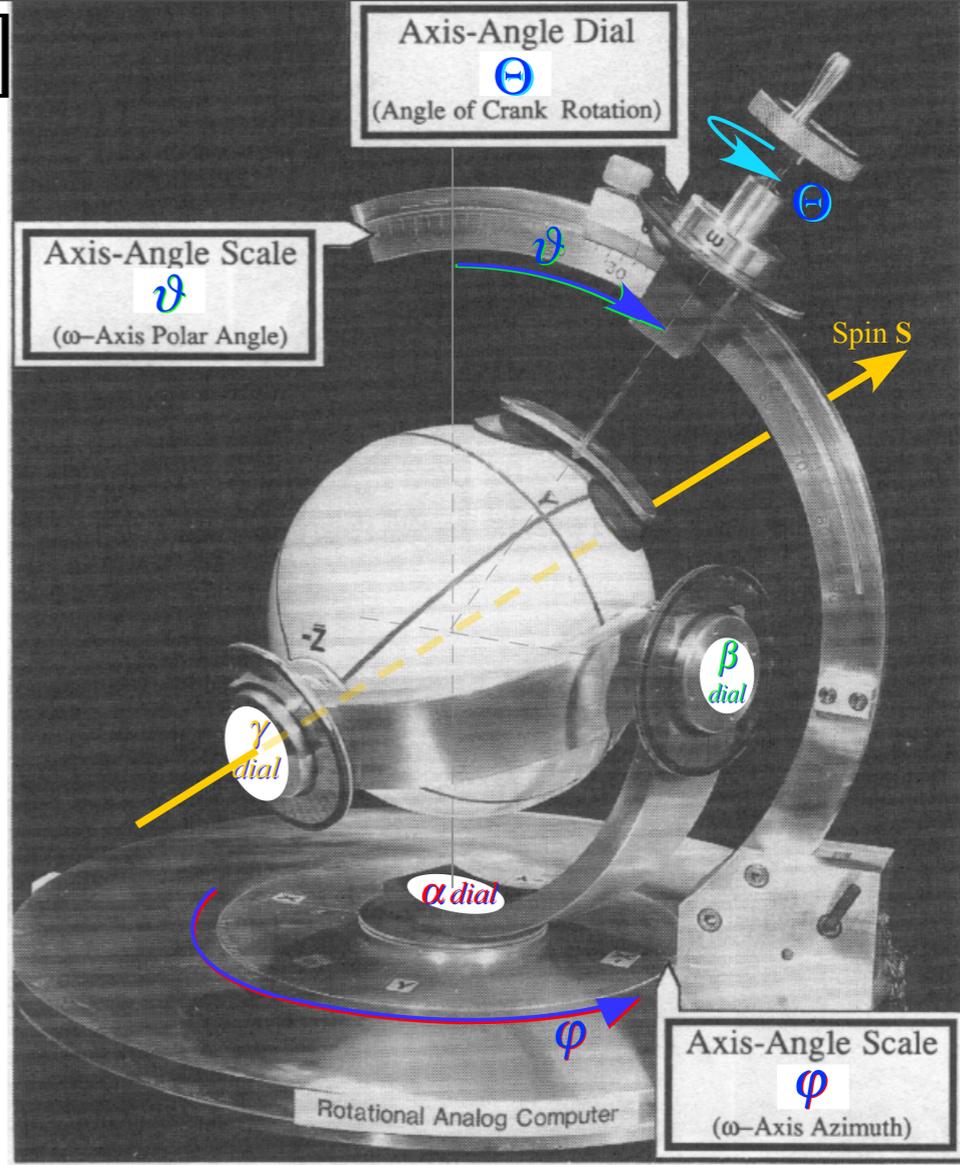
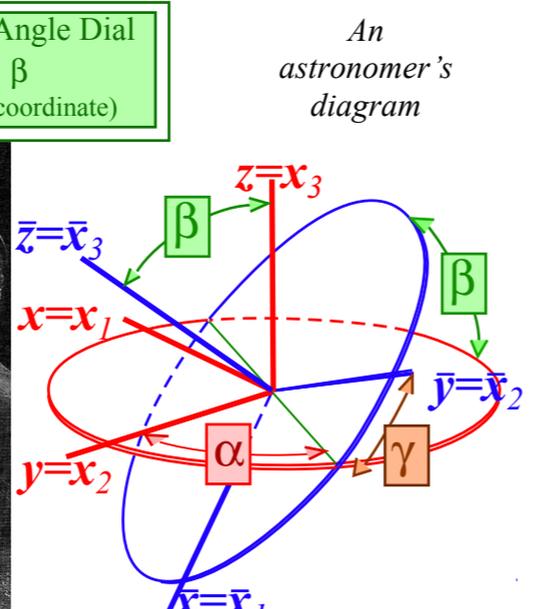
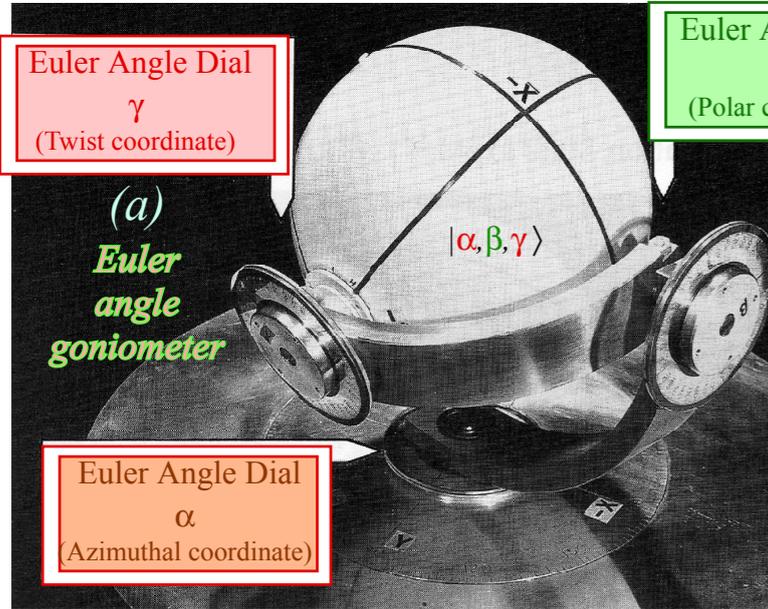


$$R[\vec{\Theta}] = \begin{pmatrix} \cos\frac{\Theta}{2} - i\hat{\Theta}_Z \sin\frac{\Theta}{2} & -i\sin\frac{\Theta}{2}(\hat{\Theta}_X - i\hat{\Theta}_Y) \\ -i\sin\frac{\Theta}{2}(\hat{\Theta}_X + i\hat{\Theta}_Y) & \cos\frac{\Theta}{2} + i\hat{\Theta}_Z \sin\frac{\Theta}{2} \end{pmatrix} = R[\varphi\vartheta\Theta] = e^{-i\mathbf{H}t}$$

$$= \cos\frac{\Theta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \hat{\Theta}_X \sin\frac{\Theta}{2} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \hat{\Theta}_Y \sin\frac{\Theta}{2} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \hat{\Theta}_Z \sin\frac{\Theta}{2}$$

$$= \begin{pmatrix} \cos\frac{\Theta}{2} - i\cos\vartheta \sin\frac{\Theta}{2} & -\sin\frac{\Theta}{2}(\sin\varphi \sin\vartheta + i\cos\varphi \sin\vartheta) \\ \sin\frac{\Theta}{2}(\sin\varphi \sin\vartheta - i\cos\varphi \sin\vartheta) & \cos\frac{\Theta}{2} + i\cos\vartheta \sin\frac{\Theta}{2} \end{pmatrix}$$

# Euler $R(\alpha\beta\gamma)$ versus Darboux $R[\varphi\vartheta\Theta]$



$$R(\alpha\beta\gamma) = \begin{pmatrix} e^{-i\frac{\alpha}{2}} & 0 \\ 0 & e^{i\frac{\alpha}{2}} \end{pmatrix} \begin{pmatrix} \cos\frac{\beta}{2} & -\sin\frac{\beta}{2} \\ \sin\frac{\beta}{2} & \cos\frac{\beta}{2} \end{pmatrix} \begin{pmatrix} e^{-i\frac{\gamma}{2}} & 0 \\ 0 & e^{i\frac{\gamma}{2}} \end{pmatrix} = \begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} \end{pmatrix}$$

$$R[\bar{\Theta}] = \begin{pmatrix} \cos\frac{\Theta}{2} - i\hat{\Theta}_Z \sin\frac{\Theta}{2} & -i\sin\frac{\Theta}{2}(\hat{\Theta}_X - i\hat{\Theta}_Y) \\ -i\sin\frac{\Theta}{2}(\hat{\Theta}_X + i\hat{\Theta}_Y) & \cos\frac{\Theta}{2} + i\hat{\Theta}_Z \sin\frac{\Theta}{2} \end{pmatrix} = R[\varphi\vartheta\Theta] = e^{-iHt}$$

$$= \cos\frac{\Theta}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - i \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \hat{\Theta}_X \sin\frac{\Theta}{2} - i \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} \hat{\Theta}_Y \sin\frac{\Theta}{2} - i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \hat{\Theta}_Z \sin\frac{\Theta}{2}$$

$$= \begin{pmatrix} \cos\frac{\Theta}{2} - i\cos\vartheta \sin\frac{\Theta}{2} & -\sin\frac{\Theta}{2}(\sin\varphi \sin\vartheta + i\cos\varphi \sin\vartheta) \\ \sin\frac{\Theta}{2}(\sin\varphi \sin\vartheta - i\cos\varphi \sin\vartheta) & \cos\frac{\Theta}{2} + i\cos\vartheta \sin\frac{\Theta}{2} \end{pmatrix}$$

$$\begin{aligned} x_1 &= \cos[(\gamma+\alpha)/2] \cos\beta/2 = \cos\Theta/2 \\ -p_2 &= \sin[(\gamma-\alpha)/2] \sin\beta/2 = \hat{\Theta}_X \sin\Theta/2 = \cos\varphi \sin\vartheta \sin\Theta/2 \\ x_2 &= \cos[(\gamma-\alpha)/2] \sin\beta/2 = \hat{\Theta}_Y \sin\Theta/2 = \sin\varphi \sin\vartheta \sin\Theta/2 \\ -p_1 &= \sin[(\gamma+\alpha)/2] \cos\beta/2 = \hat{\Theta}_Z \sin\Theta/2 = \cos\vartheta \sin\Theta/2 \end{aligned}$$

Euler  $R(\alpha\beta\gamma)$  is simpler to form than  $\Theta$ -axis Darboux  $R[\varphi\vartheta\Theta]$ .  
 Euler *state definition* lets us relate  $R(\alpha\beta\gamma)$  to  $R[\varphi\vartheta\Theta]$  ...  
 $|\alpha\beta\gamma\rangle = R(\alpha\beta\gamma)|000\rangle$  ( $\alpha\beta\gamma$  make better coordinates)

$$\begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} x_1 + ip_1 \\ x_2 + ip_2 \end{pmatrix}$$

Review: How “Crazy-Thing”-Theorem makes spinor and vector representation matrices  
Half-angle  $\Theta/2 = \varphi$  replacement and Darboux crank axis operators

Operator-on-Operator transformations

Product algebra for Pauli's  $\sigma_\mu$  and Hamilton's  $q_\mu = -i\sigma_\mu$

Group product algebra

Jordan-Pauli identity and  $U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - formula

Transformation  $\mathbf{R}[\Theta]\sigma_\mu\mathbf{R}[\Theta]^\dagger$  of spinor  $\sigma_\mu$ -operators

Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  of group-operators

Operator-on-Operator transformations

Geometry of groups: Hamilton's turns and It's all done with mirrors!

Group product geometry

$U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - geometry

Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  geometry

Euler  $\mathbf{R}(\alpha\beta\gamma)$  versus Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$



Euler  $\mathbf{R}(\alpha\beta\gamma)$  related to Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  rotation  $\Theta=0-4\pi$ -sequence  $[\varphi\vartheta]$  fixed

$R(3)$ - $U(2)$  slide rule for converting  $\mathbf{R}(\alpha\beta\gamma) \leftrightarrow \mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  Sundial

# Euler $\mathbf{R}(\alpha\beta\gamma)$ related to Darboux $\mathbf{R}[\varphi\vartheta\Theta]$

Euler *state definition* lets us relate  $\mathbf{R}(\alpha\beta\gamma)$  to  $\mathbf{R}[\varphi\vartheta\Theta]$  ...

$$|\alpha\beta\gamma\rangle = \mathbf{R}(\alpha\beta\gamma)|000\rangle \quad \alpha\beta\gamma \text{ make better coordinates but: } \mathbf{R}(\alpha\beta\gamma)|000\rangle = \mathbf{R}(\alpha\beta\gamma)|\mathbf{1}\rangle = \mathbf{R}[\varphi\vartheta\Theta]|\mathbf{1}\rangle$$

$$\begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} x_1 + ip_1 \\ x_2 + ip_2 \end{pmatrix} \quad \begin{aligned} x_1 &= \cos[(\gamma+\alpha)/2] \cos\beta/2 = \cos\Theta/2 \\ -p_2 &= \sin[(\gamma-\alpha)/2] \sin\beta/2 = \hat{\Theta}_X \sin\Theta/2 = \cos\varphi \sin\vartheta \sin\Theta/2 \\ x_2 &= \cos[(\gamma-\alpha)/2] \sin\beta/2 = \hat{\Theta}_Y \sin\Theta/2 = \sin\varphi \sin\vartheta \sin\Theta/2 \\ -p_1 &= \sin[(\gamma+\alpha)/2] \cos\beta/2 = \hat{\Theta}_Z \sin\Theta/2 = \cos\vartheta \sin\Theta/2 \end{aligned}$$

Solving these relations yields *Euler angles*  $(\alpha\beta\gamma)$  in terms of *Darboux axis angles*  $[\varphi\vartheta\Theta]$

$$\alpha = \varphi - \pi/2 + T, \quad \text{where: } T = \tan^{-1}(\tan(\Theta/2) \cos\vartheta)$$

$$\beta = 2\sin^{-1}(\sin\Theta/2 \sin\vartheta)$$

$$\gamma = \pi/2 - \varphi + T$$

# Euler $\mathbf{R}(\alpha\beta\gamma)$ related to Darboux $\mathbf{R}[\varphi\vartheta\Theta]$

Euler *state definition* lets us relate  $\mathbf{R}(\alpha\beta\gamma)$  to  $\mathbf{R}[\varphi\vartheta\Theta]$  ...

$$|\alpha\beta\gamma\rangle = \mathbf{R}(\alpha\beta\gamma)|000\rangle \quad \alpha\beta\gamma \text{ make better coordinates but: } \mathbf{R}(\alpha\beta\gamma)|000\rangle = \mathbf{R}(\alpha\beta\gamma)|1\rangle = \mathbf{R}[\varphi\vartheta\Theta]|1\rangle$$

$$\begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} x_1 + ip_1 \\ x_2 + ip_2 \end{pmatrix} \quad \begin{aligned} x_1 &= \cos[(\gamma+\alpha)/2] \cos\beta/2 = \cos\Theta/2 \\ -p_2 &= \sin[(\gamma-\alpha)/2] \sin\beta/2 = \hat{\Theta}_x \sin\Theta/2 = \cos\varphi \sin\vartheta \sin\Theta/2 \\ x_2 &= \cos[(\gamma-\alpha)/2] \sin\beta/2 = \hat{\Theta}_y \sin\Theta/2 = \sin\varphi \sin\vartheta \sin\Theta/2 \\ -p_1 &= \sin[(\gamma+\alpha)/2] \cos\beta/2 = \hat{\Theta}_z \sin\Theta/2 = \cos\vartheta \sin\Theta/2 \end{aligned}$$

Solving these relations yields *Euler angles*  $(\alpha\beta\gamma)$  in terms of *Darboux axis angles*  $[\varphi\vartheta\Theta]$

$$\alpha = \varphi - \pi/2 + T, \quad \text{where: } T = \tan^{-1}(\tan(\Theta/2) \cos\vartheta)$$

$$\beta = 2\sin^{-1}(\sin\Theta/2 \sin\vartheta)$$

$$\gamma = \pi/2 - \varphi + T$$

Inverse relations have *Darboux axis angles*  $[\varphi\vartheta\Theta]$  in terms of *Euler angles*  $(\alpha\beta\gamma)$

$$\varphi = (\alpha - \gamma + \pi)/2$$

$$\vartheta = \tan^{-1}[\tan\beta/2 / \sin(\alpha+\gamma)/2]$$

$$\Theta = 2 \cos^{-1}[\cos\beta/2 \cos(\alpha+\gamma)/2]$$

# Euler $\mathbf{R}(\alpha\beta\gamma)$ related to Darboux $\mathbf{R}[\varphi\vartheta\Theta]$

Euler *state definition* lets us relate  $\mathbf{R}(\alpha\beta\gamma)$  to  $\mathbf{R}[\varphi\vartheta\Theta]$  ...

$$|\alpha\beta\gamma\rangle = \mathbf{R}(\alpha\beta\gamma)|000\rangle \quad \alpha\beta\gamma \text{ make better coordinates but: } \mathbf{R}(\alpha\beta\gamma)|000\rangle = \mathbf{R}(\alpha\beta\gamma)|1\rangle = \mathbf{R}[\varphi\vartheta\Theta]|1\rangle$$

$$\begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} x_1 + ip_1 \\ x_2 + ip_2 \end{pmatrix} \quad \begin{array}{l} x_1 = \cos[(\gamma+\alpha)/2] \cos\beta/2 = \cos\Theta/2 \\ -p_2 = \sin[(\gamma-\alpha)/2] \sin\beta/2 = \hat{\Theta}_x \sin\Theta/2 = \cos\varphi \sin\vartheta \sin\Theta/2 \\ x_2 = \cos[(\gamma-\alpha)/2] \sin\beta/2 = \hat{\Theta}_y \sin\Theta/2 = \sin\varphi \sin\vartheta \sin\Theta/2 \\ -p_1 = \sin[(\gamma+\alpha)/2] \cos\beta/2 = \hat{\Theta}_z \sin\Theta/2 = \cos\vartheta \sin\Theta/2 \end{array}$$

Solving these relations yields *Euler angles*  $(\alpha\beta\gamma)$  in terms of *Darboux axis angles*  $[\varphi\vartheta\Theta]$

$$\alpha = \varphi - \pi/2 + T, \quad \text{where: } T = \tan^{-1}(\tan(\Theta/2) \cos\vartheta)$$

$$\beta = 2\sin^{-1}(\sin\Theta/2 \sin\vartheta)$$

$$\gamma = \pi/2 - \varphi + T$$

Inverse relations have *Darboux axis angles*  $[\varphi\vartheta\Theta]$  in terms of *Euler angles*  $(\alpha\beta\gamma)$

$$\varphi = (\alpha - \gamma + \pi)/2$$

$$\vartheta = \tan^{-1}[\tan\beta/2 / \sin(\alpha+\gamma)/2]$$

$$\Theta = 2 \cos^{-1}[\cos\beta/2 \cos(\alpha+\gamma)/2]$$

Example: *Euler angles*  $(\alpha=50^\circ \beta=60^\circ \gamma=70^\circ)$

$$\varphi = (50^\circ - 70^\circ + 180^\circ)/2 = 80^\circ$$

$$\vartheta = \tan^{-1}[\tan 60^\circ/2 / \sin(50^\circ + \gamma)/2] = 33.7^\circ$$

$$\Theta = 2 \cos^{-1}[\cos 60^\circ/2 \cos(50^\circ + \gamma)/2] = 128.7^\circ$$

# Euler $\mathbf{R}(\alpha\beta\gamma)$ related to Darboux $\mathbf{R}[\varphi\vartheta\Theta]$

Euler *state definition* lets us relate  $\mathbf{R}(\alpha\beta\gamma)$  to  $\mathbf{R}[\varphi\vartheta\Theta]$  ...

$$|\alpha\beta\gamma\rangle = \mathbf{R}(\alpha\beta\gamma)|000\rangle \quad \alpha\beta\gamma \text{ make better coordinates but: } \mathbf{R}(\alpha\beta\gamma)|000\rangle = \mathbf{R}(\alpha\beta\gamma)|1\rangle = \mathbf{R}[\varphi\vartheta\Theta]|1\rangle$$

$$\begin{pmatrix} e^{-i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} & -e^{-i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} \\ e^{i\frac{\alpha-\gamma}{2}} \sin\frac{\beta}{2} & e^{i\frac{\alpha+\gamma}{2}} \cos\frac{\beta}{2} \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} x_1+ip_1 \\ x_2+ip_2 \end{pmatrix} \quad \begin{aligned} x_1 &= \cos[(\gamma+\alpha)/2] \cos\beta/2 = \cos\Theta/2 \\ -p_2 &= \sin[(\gamma-\alpha)/2] \sin\beta/2 = \hat{\Theta}_x \sin\Theta/2 = \cos\varphi \sin\vartheta \sin\Theta/2 \\ x_2 &= \cos[(\gamma-\alpha)/2] \sin\beta/2 = \hat{\Theta}_y \sin\Theta/2 = \sin\varphi \sin\vartheta \sin\Theta/2 \\ -p_1 &= \sin[(\gamma+\alpha)/2] \cos\beta/2 = \hat{\Theta}_z \sin\Theta/2 = \cos\vartheta \sin\Theta/2 \end{aligned}$$

Solving these relations yields *Euler angles*  $(\alpha\beta\gamma)$  in terms of *Darboux axis angles*  $[\varphi\vartheta\Theta]$

$$\alpha = \varphi - \pi/2 + T, \quad \text{where: } T = \tan^{-1}(\tan(\Theta/2) \cos\vartheta)$$

$$\beta = 2\sin^{-1}(\sin\Theta/2 \sin\vartheta)$$

$$\gamma = \pi/2 - \varphi + T$$

Inverse relations have *Darboux axis angles*  $[\varphi\vartheta\Theta]$  in terms of *Euler angles*  $(\alpha\beta\gamma)$

$$\varphi = (\alpha - \gamma + \pi)/2$$

$$\vartheta = \tan^{-1}[\tan\beta/2 / \sin(\alpha+\gamma)/2]$$

$$\Theta = 2 \cos^{-1}[\cos\beta/2 \cos(\alpha+\gamma)/2]$$

Example: *Euler angles*  $(\alpha=50^\circ \beta=60^\circ \gamma=70^\circ)$

$$\varphi = (50^\circ - 70^\circ + 180^\circ)/2 = 80^\circ$$

$$\vartheta = \tan^{-1}[\tan 60^\circ/2 / \sin(50^\circ+\gamma)/2] = 33.7^\circ$$

$$\Theta = 2 \cos^{-1}[\cos 60^\circ/2 \cos(50^\circ+\gamma)/2] = 128.7^\circ$$

Reverse check:  $(\alpha\beta\gamma)$  in terms of  $[\varphi\vartheta\Theta]$

$$\alpha = 80^\circ - 90^\circ + \tan^{-1}(\tan(128.7^\circ/2) \cos 33.7^\circ) = 50.007^\circ$$

$$\beta = 2\sin^{-1}(\sin 128.7^\circ/2 \sin 33.7^\circ) = 60.022^\circ$$

$$\gamma = \pi/2 - 128.7^\circ + \tan^{-1}(\tan(128.7^\circ/2)) = 70.007^\circ$$

*Review: How “Crazy-Thing”-Theorem makes spinor and vector representation matrices  
Half-angle  $\Theta/2 = \varphi$  replacement and Darboux crank axis operators*

*Operator-on-Operator transformations*

*Product algebra for Pauli's  $\sigma_\mu$  and Hamilton's  $q_\mu = -i\sigma_\mu$*

*Group product algebra*

*Jordan-Pauli identity and  $U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - formula*

*Transformation  $\mathbf{R}[\Theta]\sigma_\mu\mathbf{R}[\Theta]^\dagger$  of spinor  $\sigma_\mu$ -operators*

*Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  of group-operators*

*Operator-on-Operator transformations*

*Geometry of groups: Hamilton's turns and It's all done with mirrors!*

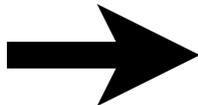
*Group product geometry*

*$U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - geometry*

*Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  geometry*

*Euler  $\mathbf{R}(\alpha\beta\gamma)$  versus Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$*

*Euler  $\mathbf{R}(\alpha\beta\gamma)$  related to Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$*

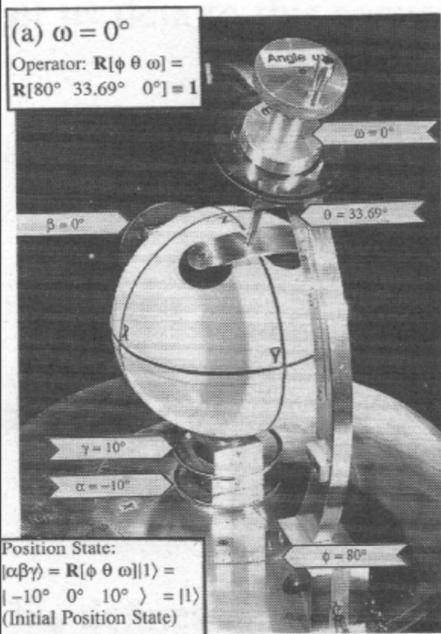
 *Euler  $\mathbf{R}(\alpha\beta\gamma)$  rotation  $\Theta=0-4\pi$ -sequence  $[\varphi\vartheta]$  fixed*

*$R(3)$ - $U(2)$  slide rule for converting  $\mathbf{R}(\alpha\beta\gamma) \leftrightarrow \mathbf{R}[\varphi\vartheta\Theta]$*

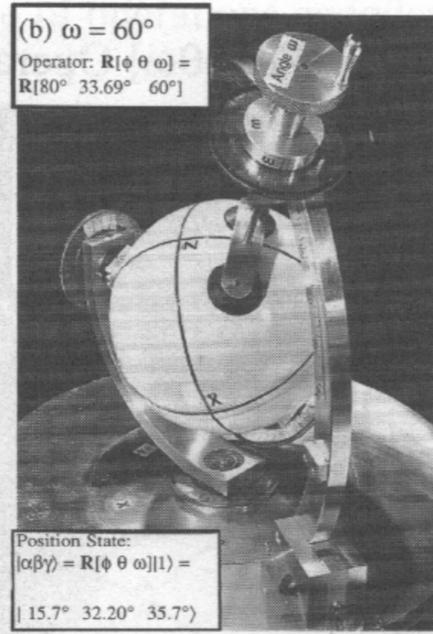
*Euler  $\mathbf{R}(\alpha\beta\gamma)$  Sundial*

# Euler $\mathbf{R}(\alpha\beta\gamma)$ rotation $\Theta=0-4\pi$ -sequence $[\varphi\vartheta]$ fixed

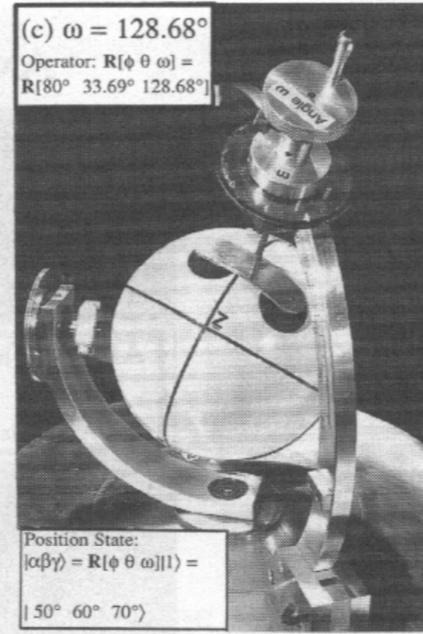
$\Theta=0^\circ$



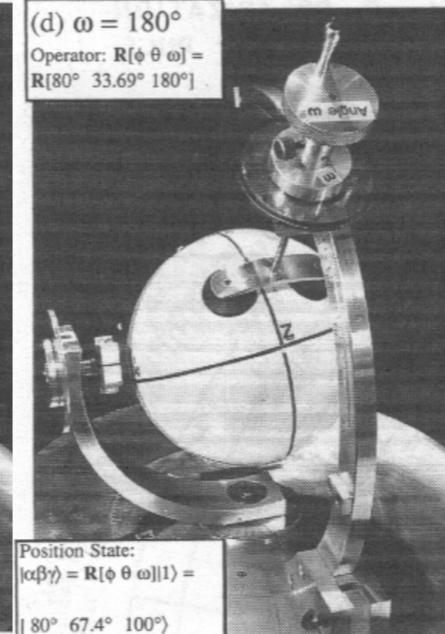
$\Theta=60^\circ$



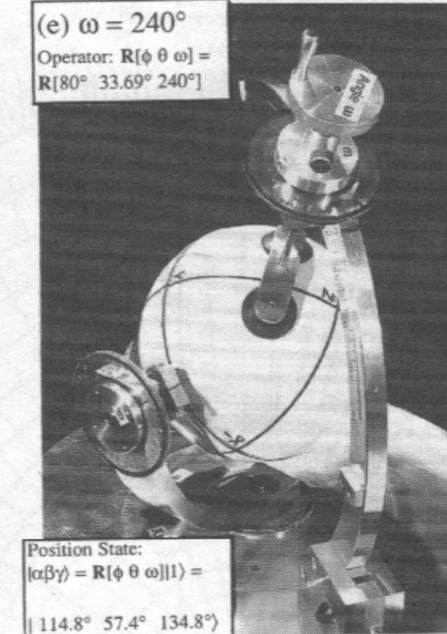
$\Theta=128.7^\circ$



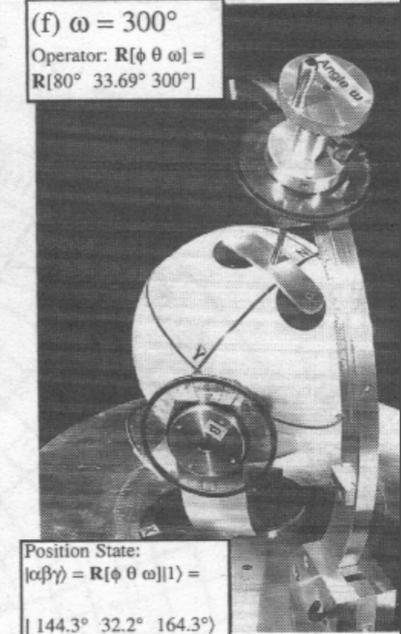
$\Theta=180^\circ$



$\Theta=240^\circ$

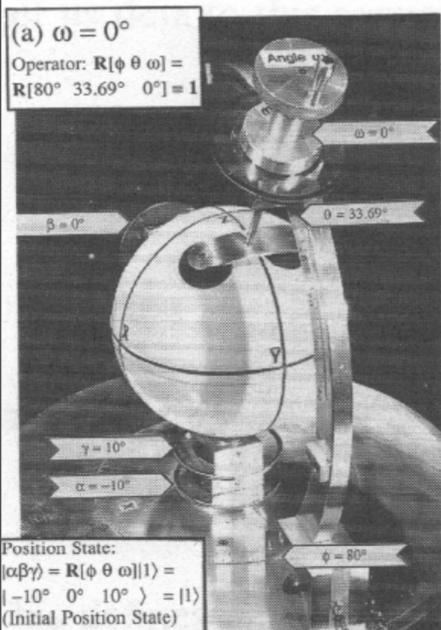


$\Theta=300^\circ$

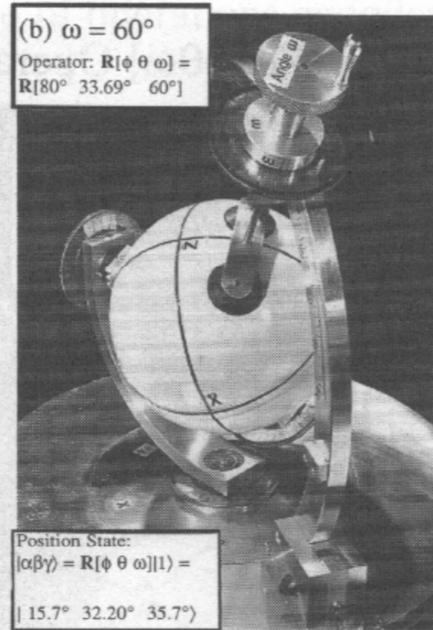


# Euler $\mathbf{R}(\alpha\beta\gamma)$ rotation $\Theta=0-4\pi$ -sequence $[\varphi\vartheta]$ fixed

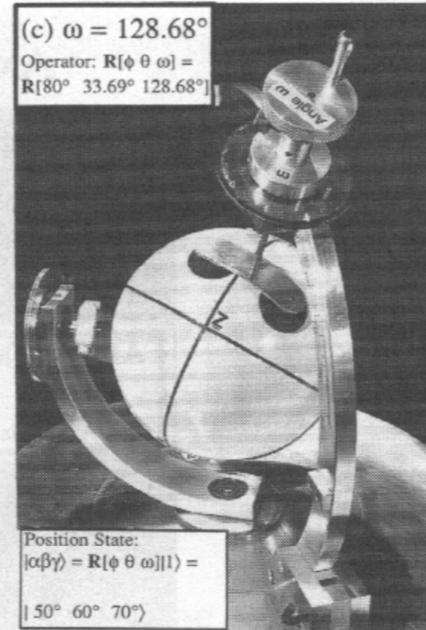
$\Theta=0^\circ$



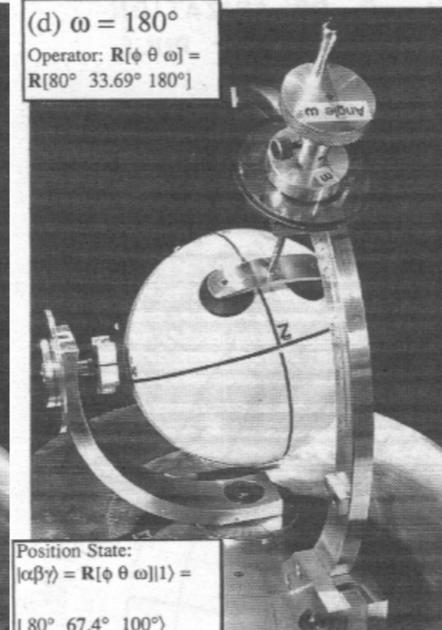
$\Theta=60^\circ$



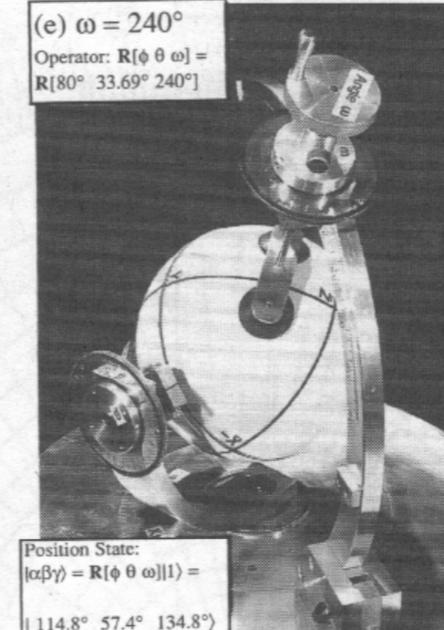
$\Theta=128.7^\circ$



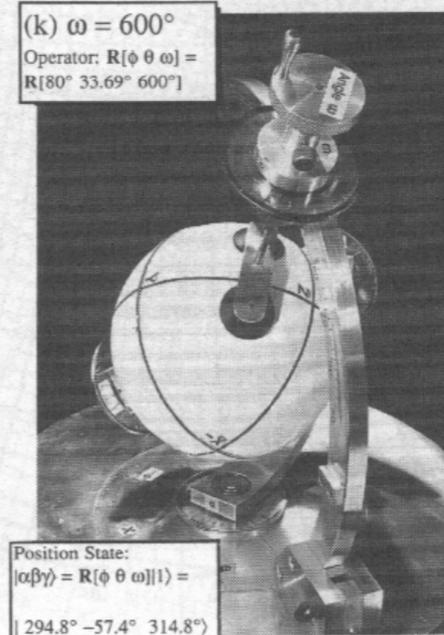
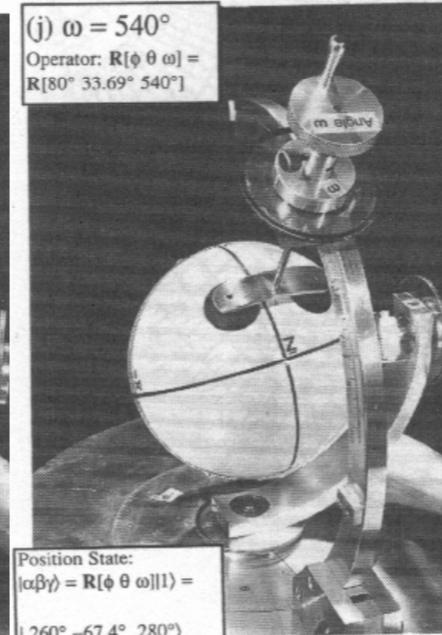
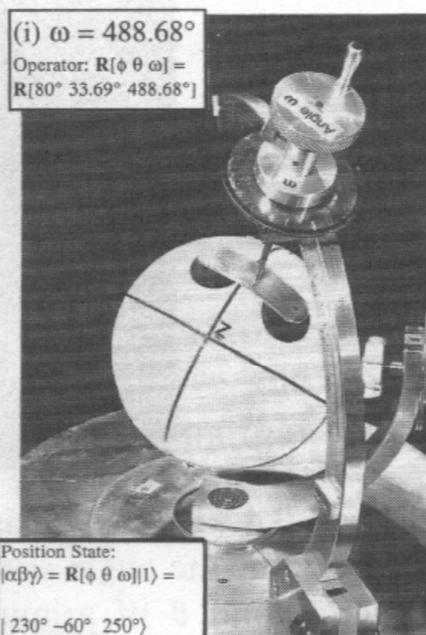
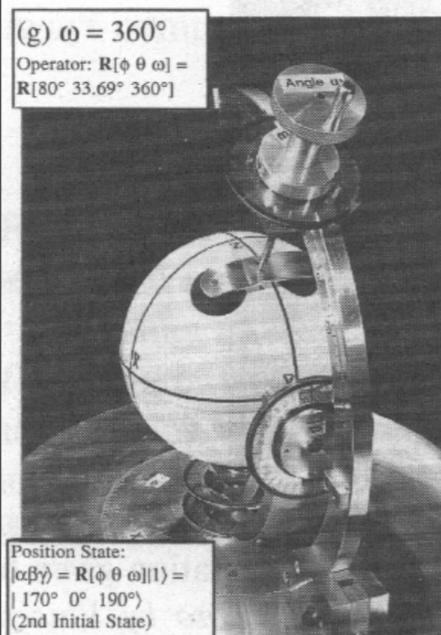
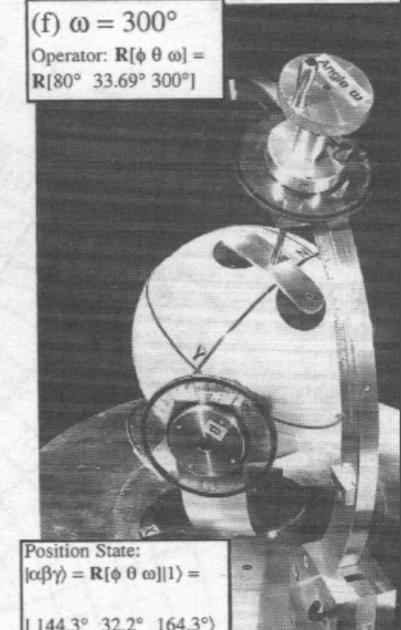
$\Theta=180^\circ$



$\Theta=240^\circ$



$\Theta=300^\circ$



$\Theta=360^\circ$

$\Theta=420^\circ$

$\Theta=488.7^\circ$

$\Theta=540^\circ$

$\Theta=600^\circ$

$\Theta=660^\circ$

Review: How “Crazy-Thing”-Theorem makes spinor and vector representation matrices  
Half-angle  $\Theta/2 = \varphi$  replacement and Darboux crank axis operators

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Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  of group-operators

Operator-on-Operator transformations

Geometry of groups: Hamilton's turns and It's all done with mirrors!

Group product geometry

$U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - geometry

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Euler  $\mathbf{R}(\alpha\beta\gamma)$  versus Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  related to Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  rotation  $\Theta=0-4\pi$ -sequence  $[\varphi\vartheta]$  fixed

→  $R(3)$ - $U(2)$  slide rule for converting  $\mathbf{R}(\alpha\beta\gamma) \leftrightarrow \mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  Sundial

# $R(3)-U(2)$ slide rule for converting $\mathbf{R}(\alpha\beta\gamma) \leftrightarrow \mathbf{R}[\phi\vartheta\Theta]$

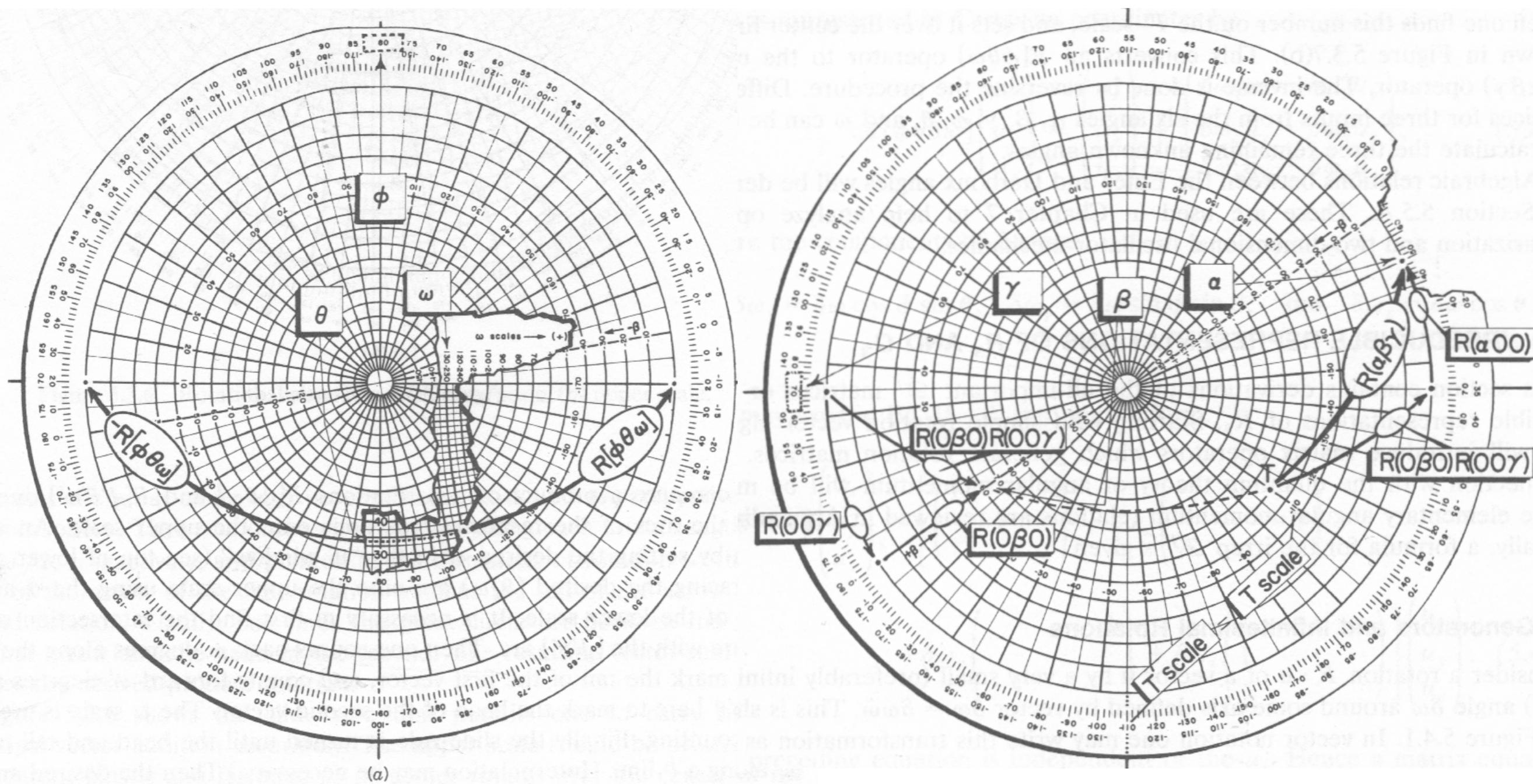


Figure 5.3.7 Setting the rotational slide rule. (a) Darboux or axis angles. (b) Euler angles.

Review: How “Crazy-Thing”-Theorem makes spinor and vector representation matrices  
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Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  of group-operators

Operator-on-Operator transformations

Geometry of groups: Hamilton's turns and It's all done with mirrors!

Group product geometry

$U(2)$  product  $\mathbf{R}[\Theta]\mathbf{R}[\Theta'] = \mathbf{R}[\Theta''']$ - geometry

Transformation  $\mathbf{R}[\Theta]\mathbf{R}[\Theta']\mathbf{R}[\Theta]^\dagger$  geometry

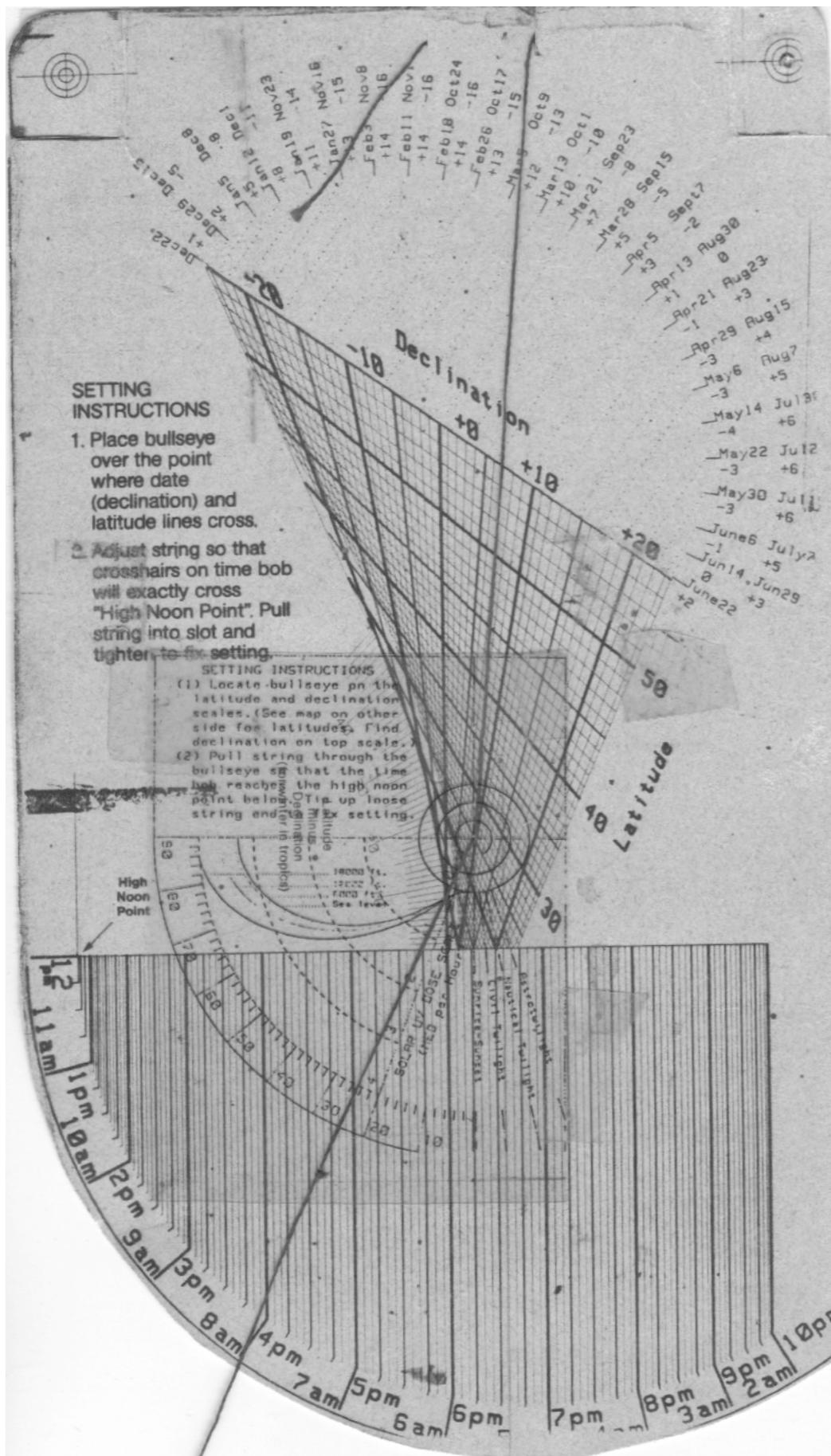
Euler  $\mathbf{R}(\alpha\beta\gamma)$  versus Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  related to Darboux  $\mathbf{R}[\varphi\vartheta\Theta]$

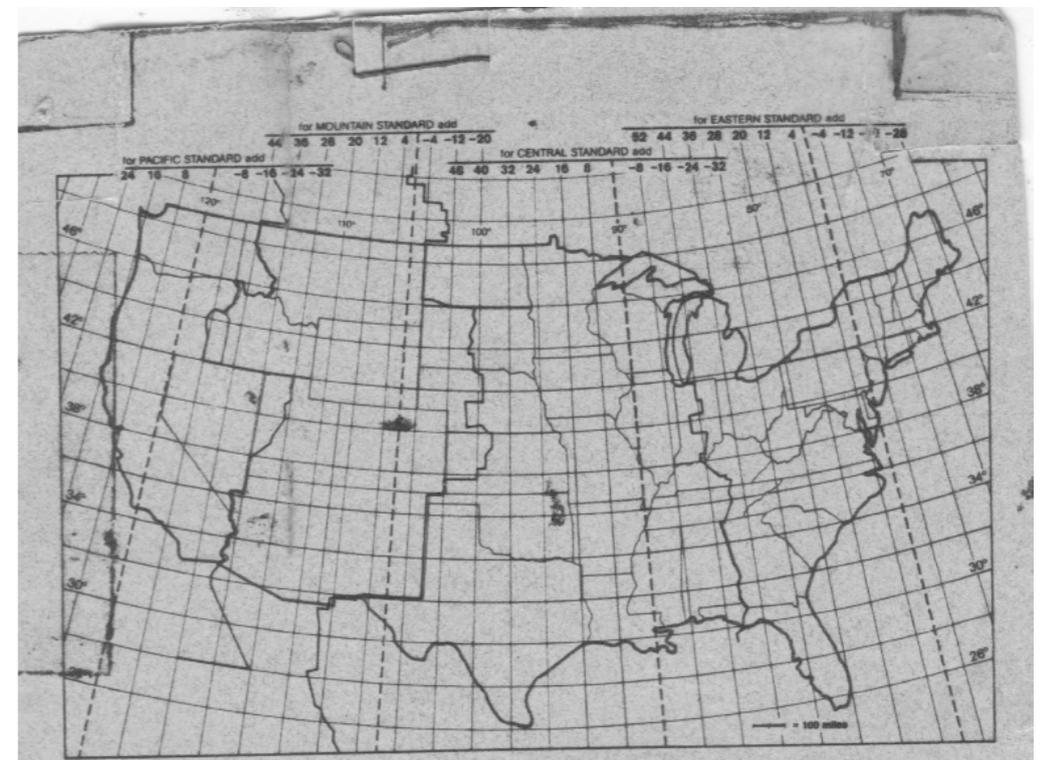
Euler  $\mathbf{R}(\alpha\beta\gamma)$  rotation  $\Theta=0-4\pi$ -sequence  $[\varphi\vartheta]$  fixed

→  $R(3)$ - $U(2)$  slide rule for converting  $\mathbf{R}(\alpha\beta\gamma) \leftrightarrow \mathbf{R}[\varphi\vartheta\Theta]$

Euler  $\mathbf{R}(\alpha\beta\gamma)$  Sundial



*Euler R(αβγ) Sundial*



FYV +16

**INSTRUCTIONS**

- Follow "Setting Instructions" on other side.
- Fold aiming tabs into place.
- Holding card vertical, tilt card so that sunlight passes through hole in tab and strikes target on opposite tab.
- Allow time bob to come to rest.
- Gently tilt card or hold time bob to keep it in position. Read SOLAR time under crosshairs.
- To convert SOLAR time to CIVIL (standard) or DAYLIGHT time, use the following formula:  
 CIVIL time = SOLAR time + date correction (see calendar) + map correction (see map)  
 DAYLIGHT time = CIVIL time + 1 hour

**SOLAR COMPUTER™**

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