Lecture 8 Mon. 9.23.2019

Quadratic form geometry and development of mechanics of Lagrange and Hamilton (Ch. 12 of Unit 1 and Ch. 4-5 of Unit 7)

Review of partial differential calculus Chain rule and order $\partial^2 \Psi / \partial x \partial y = \partial^2 \Psi / \partial y \partial x$ *symmetry*

Scaling transformation between Lagrangian and Hamiltonian views of KE Introducing 0th Lagrange and 0th Hamilton differential equations of mechanics Introducing 1st Lagrange and 1st Hamilton differential equations of mechanics

Introducing the Poincare' and Legendre contact transformations Geometry of Legendre contact transformation (Preview of Unit 8 relativistic quantum mechanics) Example from thermodynamics Legendre transform: special case of General Contact Transformation (lights, camera, ACTION!)

An elementary contact transformation from sophomore physics Algebra-calculus development of "The Volcanoes of Io" and "The Atoms of NIST" Intuitive-geometric development of """"""""""""""""""""""""""""""""""

 $Link \Rightarrow CoulIt - Simulation of the Volcanoes of Io$

 $Link \Rightarrow RelaWavity - Physical Terms H(p) \& L(u)$

This Lecture's Reference Link Listing

<u>Web Resources - front page</u> <u>UAF Physics UTube channel</u> Quantum Theory for the Computer Age

Principles of Symmetry, Dynamics, and Spectroscopy

Classical Mechanics with a Bang! Modern Physics and its Classical Foundations 2017 Group Theory for QM 2018 Adv CM 2018 AMOP 2019 Advanced Mechanics

Lecture #8

"RelaWavity" Web Simulations:

<u>2-CW laser wave, Lagrangian vs Hamiltonian,</u> <u>Physical Terms Lagrangian L(u) vs Hamiltonian H(p)</u> <u>Coullt Web Simulation of the Volcanoes of Io</u> BohrIt Multi-Panel Plot: Relativistically shifted Time-Space plots of 2 CW light waves NASA Astronomy Picture of the Day -<u>Io: The Prometheus Plume (Just Image)</u>
NASA Galileo - *Io's Alien Volcanoes*New Horizons - *Volcanic Eruption Plume on Jupiter's moon IO*NASA Galileo - A Hawaiian-Style Volcano on Io

AMOP Ch 0 Space-Time Symmetry - 2019 Seminar at Rochester Institute of Optics, Aux. slides-2018

Select, exciting, and related Research & Articles of Interest:

These *Are* hot off the presses. Out in MISC for quick reference. *Burning a hole in reality—design for a new laser may be powerful enough to pierce space-time - Sumner-Daily KOS-2019 Trampoline mirror may push laser pulse through fabric of the Universe - Lee-ArsTechnica-2019 Achieving_Extreme_Light_Intensities_using_Optically_Curved_Relativistic_Plasma_Mirrors_-_Vincenti-prl-2019 A_Soft_Matter_Computer_for_Soft_Robots_-_Garrad-sr-2019 Correlated_Insulator_Behaviour_at_Half-Filling_in_Magic-Angle_Graphene_Superlattices_-_cao-n-2018*

Sorting ultracold atoms in a three-dimensional optical lattice in a realization of Maxwell's Demon - Kumar-n-2018 Synthetic three-dimensional atomic structures assembled atom by atom - Barredo-n-2018 Older ones: Wave-particle duality of C60 molecules - Arndt-Itn-1999 Optical Vortex Knots - One Photon At A Time - Tempone-Wiltshire-Sr-2018 Baryon Deceleration by Strong Chromofields in Ultrarelativistic_, Nuclear Collisions - Mishustin-PhysRevC-2007, APS Link & Abstract

<u>Hadronic Molecules - Guo-x-2017</u> Hidden-charm pentaguark and tetraguark states - Chen-pr-2016

Running Reference Link Listing

Lectures #7 *through* #7

In reverse order

BoxIt Web Simulations:

<u>Generic/Default</u> <u>Most Basic A-Type</u> <u>Basic A-Type w/reference lines</u> <u>Basic A-Type A-Type with Potential energy</u> <u>A-Type with Potential energy and Stokes Plot</u> <u>A-Type w/3 time rates of change</u> <u>A-Type w/3 time rates of change with Stokes Plot</u> <u>B-Type (A=1.0, B=-0.05, C=0.0, D=1.0)</u>

RelaWavity Web Elliptical Motion Simulations:

Orbits with b/a=0.125 Orbits with b/a=0.5 Orbits with b/a=0.7 Exegesis with b/a=0.125 Exegesis with b/a=0.5 Exegesis with b/a=0.7 Contact Ellipsometry

Pirelli Site: Phasors animimation CMwithBang Lecture #6, page=70 (9.10.18)

Running Reference Link Listing

Lectures #6 through #1

In reverse order

<u>RelaWavity Web Simulation: Contact Ellipsometry</u> <u>BoxIt Web Simulation: Elliptical Motion (A-Type)</u> <u>CMwBang Course: Site Title Page</u> <u>Pirelli Relativity Challenge: Describing Wave Motion With Complex Phasors</u> UAF Physics UTube channel

Velocity Amplification in Collision Experiments Involving Superballs - Harter, 1971 <u>MIT OpenCourseWare: High School/Physics/Impulse and Momentum</u> <u>Hubble Site: Supernova - SN 1987A</u>

BounceItIt Web Animation - Scenarios:

49:1 y vs t, 49:1 V2 vs V1, 1:500:1 - 1D Gas Model w/ faux restorative force (Cool), 1:500:1 - 1D Gas (Warm), 1:500:1 - 1D Gas Model (Cool, Zoomed in),
Farey Sequence - Wolfram Fractions - Ford-AMM-1938
Monstermash BounceItIt Animations: 1000:1 - V2 vs V1, 1000:1 with t vs x - Minkowski Plot
Quantum Revivals of Morse Oscillators and Farey-Ford Geometry - Li-Harter-2013
Quantum Revivals of Morse Oscillators and Farey-Ford Geometry - Li-Harter-cpl-2015 Quant. Revivals of Morse Oscillators and Farey-Ford Geometry - Li-Harter-cpl-2015
Quant. Revivals of Morse Oscillators and Farey-Ford Geom. - Harter-Li-CPL-2015 (Publ.)
Velocity Amplification in Collision Experiments Involving Superballs-Harter-1971
WaveIt Web Animation - Scenarios: Quantum Carpet, Quantum Carpet_wMBars, Quantum Carpet BCar, Quantum Carpet BCar_wMBars
Wave Node Dynamics and Revival Symmetry in Quantum Rotors - Harter-JMS-2001 Wave Node Dynamics and Revival Symmetry in Ouantum Rotors - Harter-ims-2001 (Publ.)

AJP article on superball dynamics <u>AAPT Summer Reading List</u> <u>Scitation.org - AIP publications</u> HarterSoft Youtube Channel

BounceIt Web Animation - Scenarios:

Generic Scenario: <u>2-Balls dropped no Gravity (7:1) - V vs V Plot (Power=4)</u> 1-Ball dropped w/Gravity=0.5 w/Potential Plot: <u>Power=1, Power=4</u> <u>7:1 - V vs V Plot: Power=1</u> <u>3-Ball Stack (10:3:1) w/Newton plot (y vs t) - Power=4</u> <u>3-Ball Stack (10:3:1) w/Newton plot (y vs t) - Power=1 w/Gaps</u> <u>4-Ball Stack (27:9:3:1) w/Newton plot (y vs t) - Power=4</u> <u>4-Newton's Balls (1:1:1:1) w/Newtonian plot (y vs t) - Power=4</u> <u>5-Ball Totally Inelastic (1:1:1:1:1) w/Gaps: Newtonian plot (t vs x), V6 vs V5 plot</u> <u>5-Ball Totally Inelastic Pile-up w/ 5-Stationary-Balls - Minkowski plot (t vs x1) w/Gaps</u>

BounceIt Dual plots

 $m_1:m_2 = 3:1$ $v_2 v_5 v_1 and V_2 v_5 V_1, (v_1, v_2) = (1, 0.1), (v_1, v_2) = (1, 0)$ $y_2 v_5 y_1 plots: (v_1, v_2) = (1, 0.1), (v_1, v_2) = (1, 0), (v_1, v_2) = (1, -1)$ Estrangian plot $V_2 v_5 V_1$: $(v_1, v_2) = (0, 1), (v_1, v_2) = (1, -1)$

$m_1:m_2 = 4:1$

 $\frac{v2 vs vl}{v2 vs yl}$

 $m_1:m_2 = 100:1, (v_1, v_2) = (1, 0): V2 vs V1 Estrangian plot, y2 vs y1 plot$

With g=0 and 70:10 mass ratio

With non zero g, velocity dependent damping and mass ratio of 70:35

 $M_1=49, M_2=1$ with Newtonian time plot

 $M_1=49, M_2=1$ with V_2 vs V_1 plot

Example with friction

Low force constant with drag displaying a Pass-thru, Fall-Thru, Bounce-Off m1:m2=3:1 and (v1, v2) = (1, 0) Comparison with Estrangian

X2 paper: <u>Velocity Amplification in Collision Experiments Involving Superballs - Harter, et. al. 1971 (pdf)</u> Car Collision Web Simulator: <u>https://modphys.hosted.uark.edu/markup/CMMotionWeb.html</u>; with Scenarios: <u>1007</u> Superball Collision Web Simulator: <u>https://modphys.hosted.uark.edu/markup/BounceItWeb.html</u>; with Scenarios: <u>1007</u> <u>BounceIt web simulation with g=0 and 70:10 mass ratio</u> <u>With non zero g, velocity dependent damping and mass ratio of 70:35</u> Elastic Collision Dual Panel Space vs Space: <u>Space vs Time (Newton)</u>, <u>Time vs. Space(Minkowski)</u> Inelastic Collision Dual Panel Space vs Space: <u>Space vs Time (Newton)</u>, <u>Time vs. Space(Minkowski)</u> Matrix Collision Simulator: $M_1=49$, $M_2=1$ V₂ vs V₁ plot <<Under Construction>>

More Advanced QM and classical references at the end of this Lecture

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$$f(x_1, y_1) = f(x_0, y_0) + \frac{\partial f}{\partial x}(x_0, y_0) \Delta x + \frac{\partial f}{\partial y}(x_0, y_0) \Delta y + \frac{\partial f}{\partial y} \frac{\partial f}{\partial x}(x_0, y_0) \Delta x \Delta y$$
$$= f(x_0, y_0) + \frac{\partial f}{\partial y}(x_0, y_0) \Delta y + \frac{\partial f}{\partial x}(x_0, y_0) \Delta x + \frac{\partial f}{\partial x} \frac{\partial f}{\partial y}(x_0, y_0) \Delta y \Delta x$$

$$f(x_{1}, y_{1}) = f(x_{0}, y_{0}) + \frac{\partial f}{\partial x}(x_{0}, y_{0})\Delta x + \frac{\partial f}{\partial y}(x_{0}, y_{0})\Delta y + \frac{\partial f}{\partial y}\frac{\partial f}{\partial x}(x_{0}, y_{0})\Delta x\Delta y$$
$$= f(x_{0}, y_{0}) + \frac{\partial f}{\partial y}(x_{0}, y_{0})\Delta y + \frac{\partial f}{\partial x}(x_{0}, y_{0})\Delta x + \frac{\partial f}{\partial x}\frac{\partial f}{\partial y}(x_{0}, y_{0})\Delta y\Delta x$$

If f(x, y) is continuous around (x_0, y_0) and (x_1, y_1) then $\frac{\partial}{\partial y} \frac{\partial f}{\partial x}$ equals $\frac{\partial}{\partial x} \frac{\partial f}{\partial y}$

$$f(x_{1}, y_{1}) = f(x_{0}, y_{0}) + \frac{\partial f}{\partial x}(x_{0}, y_{0})\Delta x + \frac{\partial f}{\partial y}(x_{0}, y_{0})\Delta y + \frac{\partial f}{\partial y}\frac{\partial f}{\partial x}(x_{0}, y_{0})\Delta x\Delta y$$
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If f(x, y) is continuous around (x_0, y_0) and (x_1, y_1) then $\frac{\partial}{\partial y} \frac{\partial f}{\partial x}$ equals $\frac{\partial}{\partial x} \frac{\partial f}{\partial y}$

1. Chain rules

 $\begin{bmatrix} f(x_1, y_1) - f(x_0, y_0) \end{bmatrix} = df = \frac{\partial f}{\partial x} (x_0, y_0) dx + \frac{\partial f}{\partial y} (x_0, y_0) dy \dots_{(keep \ 1^{st} - order \ terms \ only!)}$ $\frac{df}{dt} = \frac{\partial f}{\partial x} (x_0, y_0) \frac{dx}{dt} + \frac{\partial f}{\partial y} (x_0, y_0) \frac{dy}{dt}$ $\dot{f} = \frac{\partial f}{\partial x} \dot{x} + \frac{\partial f}{\partial y} \dot{y} \qquad (shorthand \ notation)$

$$f(x_{1}, y_{1}) = f(x_{0}, y_{0}) + \frac{\partial f}{\partial x}(x_{0}, y_{0})\Delta x + \frac{\partial f}{\partial y}(x_{0}, y_{0})\Delta y + \frac{\partial g}{\partial y}\partial x(x_{0}, y_{0})\Delta x \Delta y$$
$$= f(x_{0}, y_{0}) + \frac{\partial f}{\partial y}(x_{0}, y_{0})\Delta y + \frac{\partial f}{\partial x}(x_{0}, y_{0})\Delta x + \frac{\partial g}{\partial x}\partial y(x_{0}, y_{0})\Delta y \Delta x$$

If f(x, y) is continuous around (x_0, y_0) and (x_1, y_1) then $\frac{\partial}{\partial y} \frac{\partial f}{\partial x}$ equals $\frac{\partial}{\partial x} \frac{\partial f}{\partial y}$

 $+ \partial_{y} f \dot{y}$

1. Chain rules

$$f(x_{1}, y_{1}) - f(x_{0}, y_{0})] = df = \frac{\partial f}{\partial x}(x_{0}, y_{0})dx + \frac{\partial f}{\partial y}(x_{0}, y_{0})dy \dots_{(keep \ 1^{st} - order \ terms \ only!)}$$
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$$\dot{f} = \frac{\partial f}{\partial x}\dot{x} + \frac{\partial f}{\partial y}\dot{y} \qquad (shorthand \ notation) = \partial_{x} f \dot{x}$$

2. Symmetry of partial deriv. ordering $\frac{\partial}{\partial y} \frac{\partial f}{\partial x} = \frac{\partial}{\partial x} \frac{\partial f}{\partial y} \quad \text{or:} \quad \frac{\partial^2 f}{\partial y \partial x} = \frac{\partial^2 f}{\partial x \partial y} \quad \text{or:} \quad \partial_y \partial_x f = \partial_x \partial_y f$

 $(pay attention to the 2^{nd} - order terms, too!)$

(shorthand notation)

$$f(x_{1}, y_{1}) = f(x_{0}, y_{0}) + \frac{\partial f}{\partial x}(x_{0}, y_{0})\Delta x + \frac{\partial f}{\partial y}(x_{0}, y_{0})\Delta y + \frac{\partial g}{\partial y}\partial x(x_{0}, y_{0})\Delta x \Delta y$$

= $f(x_{0}, y_{0}) + \frac{\partial f}{\partial y}(x_{0}, y_{0})\Delta y + \frac{\partial f}{\partial x}(x_{0}, y_{0})\Delta x + \frac{\partial g}{\partial x}\partial y(x_{0}, y_{0})\Delta y \Delta x$

If f(x, y) is continuous around (x_0, y_0) and (x_1, y_1) then $\frac{\partial}{\partial y} \frac{\partial f}{\partial x}$ equals $\frac{\partial}{\partial x} \frac{\partial f}{\partial y}$

 $+ \partial_{y} f \dot{y}$

(pay attention to the 2nd - order terms, too!)

1. Chain rules

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2. Symmetry of partial deriv. ordering $\frac{\partial}{\partial y}\frac{\partial f}{\partial x} = \frac{\partial}{\partial x}\frac{\partial f}{\partial y} \quad \text{or:} \quad \frac{\partial^2 f}{\partial y \partial x} = \frac{\partial^2 f}{\partial x \partial y} \quad \text{or:} \quad \partial_y \partial_x f = \partial_x \partial_y f$ (shorthand notation)

Let :
$$\vec{\nabla} = \begin{pmatrix} \partial_x & \partial_y \end{pmatrix}$$
 so : $\vec{\nabla}f \cdot \mathbf{dr} = \begin{pmatrix} \partial_x f & \partial_y f \end{pmatrix} \cdot \begin{pmatrix} dx \\ dy \end{pmatrix} = \partial_x f \, dx + \partial_y f \, dy = df$

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An elementary contact transformation from sophomore physics Algebra-calculus development of "The Volcanoes of Io" and "The Atoms of NIST" Intuitive-geometric development of """"""""""""""""""""""""""""""""""" Three ways to express energy: Consider kinetic energy (KE) first

1. Lagrangian is explicit function of velocity:
$$\mathbf{v} = \begin{pmatrix} v_1 \\ v_2 \end{pmatrix}$$

 $L(v_k...) = \frac{1}{2} (m_1 v_1^2 + m_2 v_2^2 + ...) = L(\mathbf{v}...) = \frac{1}{2} \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + ... = \frac{1}{2} \begin{pmatrix} v_1 & v_2 \end{pmatrix} \begin{pmatrix} m_1 & 0 \\ 0 & m_2 \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} + ...$

2. "Estrangian" is explicit function of **R**-rescaled velocity:
or: "speedinum"
$$V = \mathbf{R} \cdot \mathbf{v}$$
 or: $\begin{pmatrix} V_1 \\ V_2 \end{pmatrix} = \begin{pmatrix} \sqrt{m_1} & 0 \\ 0 & \sqrt{m_2} \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix}$
 $E(V_k \dots) = \frac{1}{2} (V_1^2 + V_2^2 + \dots) = E(\mathbf{V} \dots) = \frac{1}{2} \mathbf{V} \cdot \mathbf{1} \cdot \mathbf{V} + \dots = \frac{1}{2} \begin{pmatrix} V_1 & V_2 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} V_1 \\ V_2 \end{pmatrix} + \dots$

3. Hamiltonian is explicit function of
$$\mathbf{M} = \mathbf{R}^2$$
-rescaled velocity:
or: momentum p $\mathbf{p} = \mathbf{M} \cdot \mathbf{v}$ or: $\begin{pmatrix} p_1 \\ p_2 \end{pmatrix} = \begin{pmatrix} m_1 & 0 \\ 0 & m_2 \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} = \begin{pmatrix} m_1 v_1 \\ m_2 v_2 \end{pmatrix}$
 $H(p_k \dots) = \frac{1}{2}(\frac{p_1^2}{m_1} + \frac{p_2^2}{m_2} + \dots) = H(\mathbf{p} \dots) = \frac{1}{2}\mathbf{p} \cdot \mathbf{M}^{-1} \cdot \mathbf{p} + \dots = \frac{1}{2}\begin{pmatrix} p_1 & p_2 \end{pmatrix} \begin{pmatrix} 1/m_1 & 0 \\ 0 & 1/m_2 \end{pmatrix} \begin{pmatrix} p_1 \\ p_2 \end{pmatrix} + \dots$

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Introducing the (partial $\frac{\partial 2}{\partial 2}$) differential equations of mechanics

Starts out with simple demands for explicit-dependence, "loyalty" or "fealty to the colors"

Lagrangian and *Estrangian* have <u>no</u> explicit dependence on *momentum* **p**=**M**•v

$$\frac{\partial L}{\partial p_k} \equiv 0 \equiv \frac{\partial E}{\partial p_k}$$

Hamiltonian and *Estrangian* have <u>no</u> explicit dependence on velocity v=M⁻¹•p

$$\frac{\partial H}{\partial v_k} \equiv 0 \equiv \frac{\partial E}{\partial v_k}$$

Lagrangian and Hamiltonian have <u>no</u> explicit dependence On speedinum V=M^{1/2}•v

$$\frac{\partial L}{\partial V_k} \equiv 0 \equiv \frac{\partial H}{\partial V_k}$$

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Lagrangian and *Estrangian* have <u>no</u> explicit dependence on *momentum* **p**=**M**•v *Hamiltonian* and *Estrangian* have <u>no</u> explicit dependence on *velocity* v=M⁻¹•p Lagrangian and Hamiltonian have <u>no</u> explicit dependence ON speedinum V=M^{1/2}•V

$$\frac{\partial L}{\partial p_k} \equiv 0 \equiv \frac{\partial E}{\partial p_k} \qquad \qquad \frac{\partial H}{\partial v_k} \equiv 0 \equiv \frac{\partial E}{\partial v_k} \qquad \qquad \frac{\partial L}{\partial V_k} \equiv 0 \equiv \frac{\partial H}{\partial V_k}$$

Such non-dependencies hold in spite of "under-the-table" matrix and partial-differential connections[†]

$$\nabla_{v}L = \frac{\partial L}{\partial \mathbf{v}} = \frac{\partial}{\partial \mathbf{v}} \frac{\mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v}}{2}$$
$$= \mathbf{M} \cdot \mathbf{v} = \mathbf{p}$$
$$\begin{pmatrix} \frac{\partial L}{\partial v_{1}} \\ \frac{\partial L}{\partial v_{2}} \end{pmatrix} = \begin{pmatrix} m_{1} & 0 \\ 0 & m_{2} \end{pmatrix} \begin{pmatrix} v_{1} \\ v_{2} \end{pmatrix} = \begin{pmatrix} p_{1} \\ p_{2} \end{pmatrix}$$
$$Lagrange's \ 1^{st} \ equation(s)$$
$$\frac{\partial L}{\partial v_{k}} = p_{k} \quad \text{or:} \quad \frac{\partial L}{\partial \mathbf{v}} = \mathbf{p}$$

 $\nabla_{p}H = \mathbf{v} = \frac{\partial H}{\partial \mathbf{p}} = \frac{\partial}{\partial \mathbf{p}} \frac{\mathbf{p} \cdot \mathbf{M}^{-1} \cdot \mathbf{p}}{2}$ $= \mathbf{M}^{-1} \cdot \mathbf{p} = \mathbf{v}$ $\begin{pmatrix} \frac{\partial H}{\partial p_{1}} \\ \frac{\partial H}{\partial p_{2}} \end{pmatrix} = \begin{pmatrix} m_{1}^{-1} & 0 \\ 0 & m_{2}^{-1} \end{pmatrix} \begin{pmatrix} p_{1} \\ p_{2} \end{pmatrix} = \begin{pmatrix} v_{1} \\ v_{2} \end{pmatrix}$ $Hamilton 's \ 1^{st} \ equation(s)$ $\frac{\partial H}{\partial p_{k}} = v_{k} \quad \text{or:} \quad \frac{\partial H}{\partial \mathbf{p}} = \mathbf{v}$

Estrangian is neglected for now. (*It is related to dual ellipse geometry in Lecture 7 p. 71-79 and 80-85*)

> *†non-dependency due to stationary-value effects as shown on* <u>*p.39*</u> *to* <u>*p.48*</u>







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Given matrix relation: $\mathbf{p}=\mathbf{M}\cdot\mathbf{v}$ or its inverse: $\mathbf{v}=\mathbf{M}^{-1}\cdot\mathbf{p}$ you might be tempted to rewrite

Q-forms $L(\mathbf{v}..) = (1/2)\mathbf{v}\cdot\mathbf{M}\cdot\mathbf{v}$ or $H(\mathbf{p}..) = (1/2)\mathbf{p}\cdot\mathbf{M}^{-1}\cdot\mathbf{p}$ to be $H = (1/2)\mathbf{p}\cdot\mathbf{v}$ or equivalently $L = (1/2)\mathbf{v}\cdot\mathbf{p}$.

Given matrix relation: $\mathbf{p}=\mathbf{M}\cdot\mathbf{v}$ or its inverse: $\mathbf{v}=\mathbf{M}^{-1}\cdot\mathbf{p}$ you might be tempted to rewrite

Q-forms $L(\mathbf{v}..) = (1/2)\mathbf{v}\cdot\mathbf{M}\cdot\mathbf{v}$ or $H(\mathbf{p}..) = (1/2)\mathbf{p}\cdot\mathbf{M}^{-1}\cdot\mathbf{p}$ to be $H = (1/2)\mathbf{p}\cdot\mathbf{v}$ or equivalently $L = (1/2)\mathbf{v}\cdot\mathbf{p}$. Numerically-CORRECT, but Differentially-WRONG!

Given matrix relation: $\mathbf{p}=\mathbf{M}\cdot\mathbf{v}$ or its inverse: $\mathbf{v}=\mathbf{M}\cdot\mathbf{l}\cdot\mathbf{p}$ you might be tempted to rewrite

Q-forms $L(\mathbf{v}..) = (1/2)\mathbf{v}\cdot\mathbf{M}\cdot\mathbf{v}$ or $H(\mathbf{p}..) = (1/2)\mathbf{p}\cdot\mathbf{M}^{-1}\cdot\mathbf{p}$ to be $H = (1/2)\mathbf{p}\cdot\mathbf{v}$ or equivalently $L = (1/2)\mathbf{v}\cdot\mathbf{p}$.

Numerically-CORRECT, but Differentially-WRONG! (In classical physics **p**•**v** and **v**•**p** are identical)

Instead try: $H(\mathbf{p}..) = \mathbf{p} \cdot \mathbf{v} \cdot (1/2) \mathbf{v} \cdot \mathbf{p} = \mathbf{p} \cdot \mathbf{v} \cdot L(\mathbf{v}..)$ or else: $L(\mathbf{v}..) = \mathbf{p} \cdot \mathbf{v} \cdot H(\mathbf{p}..)$

Given matrix relation: $\mathbf{p}=\mathbf{M}\cdot\mathbf{v}$ or its inverse: $\mathbf{v}=\mathbf{M}^{-1}\cdot\mathbf{p}$ you might be tempted to rewrite

Q-forms $L(\mathbf{v}..)=(1/2)\mathbf{v}\cdot\mathbf{M}\cdot\mathbf{v}$ or $H(\mathbf{p}..)=(1/2)\mathbf{p}\cdot\mathbf{M}^{-1}\cdot\mathbf{p}$ to be $H=(1/2)\mathbf{p}\cdot\mathbf{v}$ or equivalently $L=(1/2)\mathbf{v}\cdot\mathbf{p}$. Numerically-CORRECT, but Differentially-WRONG!

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That is ... the Legendre contact transformation $L(\mathbf{v}) = \mathbf{p} \cdot \mathbf{v} - H(\mathbf{p}) \quad or: \quad H(\mathbf{p}) = \mathbf{p} \cdot \mathbf{v} - L(\mathbf{v})$
Introducing the Poincare' and Legendre contact transformations

Given matrix relation: $\mathbf{p}=\mathbf{M}\cdot\mathbf{v}$ or its inverse: $\mathbf{v}=\mathbf{M}\cdot\mathbf{P}$ you might be tempted to rewrite

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Now explicit dependency (non)-relations give the right derivatives

$$\frac{\partial L(\mathbf{v})}{\partial \mathbf{p}} = \frac{\partial}{\partial \mathbf{p}} \mathbf{p} \cdot \mathbf{v} - \frac{\partial H(\mathbf{p})}{\partial \mathbf{p}} \qquad \qquad \frac{\partial H(\mathbf{p})}{\partial \mathbf{v}} = \frac{\partial}{\partial \mathbf{v}} \mathbf{p} \cdot \mathbf{v} - \frac{\partial L(\mathbf{v})}{\partial \mathbf{v}} \\ 0 = \mathbf{v} - \frac{\partial H(\mathbf{p})}{\partial \mathbf{p}} \qquad \qquad 0 = \mathbf{p} - \frac{\partial L(\mathbf{v})}{\partial \mathbf{v}} \end{cases}$$

Introducing the Poincare' and Legendre contact transformations

Given matrix relation: $\mathbf{p}=\mathbf{M}\cdot\mathbf{v}$ or its inverse: $\mathbf{v}=\mathbf{M}\cdot\mathbf{l}\cdot\mathbf{p}$ you might be tempted to rewrite

Q-forms $L(\mathbf{v}_{..}) = (1/2)\mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v}$ or $H(\mathbf{p}_{..}) = (1/2)\mathbf{p} \cdot \mathbf{M}^{-1} \cdot \mathbf{p}$ to be $H = (1/2)\mathbf{p} \cdot \mathbf{v}$ or equivalently $L = (1/2)\mathbf{v} \cdot \mathbf{p}$. Numerically-CORRECT, but Differentially-WRONG!

Instead try: $H(\mathbf{p}_{..}) = \mathbf{p} \cdot \mathbf{v}_{-(1/2)} \mathbf{v} \cdot \mathbf{p} = \mathbf{p} \cdot \mathbf{v}_{-L}(\mathbf{v}_{..})$ or else: $L(\mathbf{v}_{..}) = \mathbf{p} \cdot \mathbf{v}_{-H}(\mathbf{p}_{..})$

That is ... the Legendre contact transformation $L(\mathbf{v}) = \mathbf{p} \cdot \mathbf{v} - H(\mathbf{p})$ or: $H(\mathbf{p}) = \mathbf{p} \cdot \mathbf{v} - L(\mathbf{v})$

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$$\frac{\partial L(\mathbf{v})}{\partial \mathbf{p}} = \frac{\partial}{\partial \mathbf{p}} \mathbf{p} \cdot \mathbf{v} - \frac{\partial H(\mathbf{p})}{\partial \mathbf{p}} \qquad \qquad \frac{\partial H(\mathbf{p})}{\partial \mathbf{v}} = \frac{\partial}{\partial \mathbf{v}} \mathbf{p} \cdot \mathbf{v} - \frac{\partial L(\mathbf{v})}{\partial \mathbf{v}} \\ 0 = \mathbf{v} - \frac{\partial H(\mathbf{p})}{\partial \mathbf{p}} \qquad \qquad 0 = \mathbf{p} - \frac{\partial L(\mathbf{v})}{\partial \mathbf{v}} \end{cases}$$

That is

Hamilton's 1st equation(s) and Lagrange's 1st equation(s)

$$\mathbf{v} = \frac{\partial H(\mathbf{p})}{\partial \mathbf{p}}$$
 $\mathbf{p} = \frac{\partial L(\mathbf{v})}{\partial \mathbf{v}}$



Scaling transformation between Lagrangian and Hamiltonian views of KE Introducing 0th Lagrange and 0th Hamilton differential equations of mechanics Introducing 1st Lagrange and 1st Hamilton differential equations of mechanics

Introducing the Poincare' and Legendre contact transformations

Geometry of Legendre contact transformation (Preview of Unit 8 relativistic quantum mechanics) Example from thermodynamics Legendre transform: special case of General Contact Transformation (lights,camera, ACTION!)

An elementary contact transformation from sophomore physics Algebra-calculus development of "The Volcanoes of Io" and "The Atoms of NIST" Intuitive-geometric development of """"""""""""""""""""""""""""""""""



Energy (E) Coordinate angle v=atan(u/c)

Preview of Unit 8: Geometry of Legendre contact transformation persists in relativistic quantum mechanics!

(In fact it is due to the wave mechanics and phase invariance principles.)



 $\underline{\text{Link}} \Rightarrow \text{RelaWavity} - \text{Physical Terms H(p) \& L(u)}$

Coordinate angle v=atan(u/c)

Preview of Unit 8: Geometry of Legendre contact transformation persists in relativistic quantum mechanics!

(In fact it is due to the wave mechanics and phase invariance principles.)

More to the point it's due to Evenson Axiom: "All colors go c



 $Link \Rightarrow RelaWavity - Physical Terms H(p) \& L(u)$

How Legendre contact transformations work...(to make $\frac{\partial H}{\partial v} = 0$ or $\frac{\partial L}{\partial p} = 0$)

Secant lines $L(\mathbf{v}) = p \cdot v - H$ of fixed slope $p = \frac{\partial L}{\partial v}$ and decreasing intercept $-H(v_{-2}) > -H(v_{-1}) > ...$ for increasing velocity $v_{-2} > v_{-1} > ... > v_0$ lead to unique tangent to $L(\mathbf{v})$ -curve at the tangent contact point $v = v_0$ that has max $H(p v_0)$ Thus $\frac{\partial H}{\partial v} = 0$



Unit 1 Fig. 12.4

How Legendre contact transformations work...(to make $\frac{\partial H}{\partial v} = 0$ or $\frac{\partial L}{\partial p} = 0$)

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Unit 1 Fig. 12.4

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Unit 1 Fig. 12.4

How L	legendre co	ntact tr	ransform	ations	S WORK(to make $\frac{\partial I}{\partial t}$	$\frac{H}{v} = 0 Or \frac{\partial L}{\partial p} = 0)$
Secant line and decrea for increa lead to un tangent co Thus $\frac{\partial H}{\partial v} =$	thes $L(\mathbf{v}) = p \cdot v - H$ the easing intercept - the easing velocity the end of the easing velocity the end of the easing velocity the easing velocit	$ \begin{array}{c} of fixed \\ -H(v_{-2}) > -H \\ v_{-2} > v_{-1} > \\ L(\mathbf{v}) - curve \\ that has \underline{n} \end{array} $	slope $p = \frac{\partial L}{\partial v}$ $(v_{-1}) > \dots$ $\dots > v_0$ at the $\max H(p v_0)$		(Similarly)	Unit 1 Fig. 12.4
(a) Secant	lines: $L(v) = p \cdot v - H$	H S	(b)	Secant	lines: $H(p) = p \cdot v \cdot L(v)$	
$L(\mathbf{v})$	for fixed slope p and varying H/		(0)	<i>H(</i> p)	for fixed slope v and varying L	stope
		Tangent i extreme v	line points to value -H(v ₀)		Tangent line points to extreme value $-L(p_0)$	dL(p)/dp = 0
	V_3	of interce dH(v,	dv = 0		of intercept -L thus: dL(p)/dp = 0	momentum
$-H(v_{-3})$	V-2 V-1 V	, v ₁	<u>velo</u> city v -L((p_{α})	<i>p</i> ₋₃ <i>p</i> ₋₂ <i>p</i> ₋₁	$p_0 p_1 \dots p$
$-H(v_{-2})$ $-H(v_{-1})/2$ $-H(v_{0})$	extre	me values.	-L -L -/	(p_{-2}) $L(p_{-1})$ $L(p_{0})$		
$\frac{\partial H}{\partial v} = 0 \text{ at each point } v = \frac{\partial H}{\partial p} \text{ of } L(v) \text{ with slope } p = \frac{\partial L}{\partial v}$						



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Example of Legendre contact transformation in thermodynamics

Internal energy U(S, V) is defined as a function of entropy S and volume V.

A new function *enthalpy* H(S,P) depends on entropy and *pressure* P.

It is a Legendre transform $H(S,P) = P \cdot V + U$ of energy U(S,V) to new variable $P = -(\frac{\partial U}{\partial V})_S$.

Example of Legendre contact transformation in thermodynamics

Lagrangian L(r,v)position rvelocity vInternal energy U(S,V) is defined as a function of entropy S and volume V.Hamiltonian H(r,p)position rmomentum pA new function enthalpy H(S,P) depends on entropy and pressure P. $p = (\frac{\partial L}{\partial v})_r$ $H(r,p) = p \cdot v - L$ Lagrangian L(r,v) $p = (\frac{\partial L}{\partial v})_r$ It is a Legendre transform $H(S,P) = P \cdot V + U$ of energy U(S,V) to new variable $P = -(\frac{\partial U}{\partial V})_S$.

Example of Legendre contact transformation in thermodynamics

Lagrangian L(r,v)position rvelocity vInternal energy U(S, V) is defined as a function of entropy S and volume V.Hamiltonian H(r,p)position rmomentum pA new function enthalpy H(S,P) depends on entropy and pressure P. $p = (\frac{\partial L}{\partial v})_r$ It is a Legendre transform $H(S,P) = P \cdot V + U$ of energy U(S,V) to new variable $P = -(\frac{\partial U}{\partial v})_s$.Except for \pm signs, it's our Hamiltonian $H(p) = p \cdot v - L(v)$ going from Lagrangian L(v)to use new variable momentum $p = (\frac{\partial L}{\partial v})_x$.

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Legendre transform: special case of General Contact Transformation Active-Contact-Transformation Generator or Action function: S(x, y:X, Y)=const. does mapping. (b)(a) YY(X) is mapped from y(x) as an envelope of contacting S=const. curves. y(x) $(x_0, y(x_0))$ $S(x_2, y_2, X, Y) = 10$ $(X_0, Y(X_0))$ Unit 1 Fig. 12.7 $S(x_0, y_0, X, Y) = 10$ $X_0 X_1 X_2$ x_0 x_{l} *x*₂ X $\boldsymbol{\chi}$











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An elementary contact transformation from sophomore physics Algebra-calculus development of "The Volcanoes of Io" and "The Atoms of NIST" Intuitive-geometric development of " " " and " " ""







 $Link \Rightarrow Coullt - Simulation of the Volcanoes of Io$

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Introducing the Poincare' and Legendre contact transformations Geometry of Legendre contact transformation (Preview of Unit 8 relativistic quantum mechanics) Example from thermodynamics Legendre transform: special case of General Contact Transformation (lights,camera, ACTION!)



Unit 1 Fig. 12.5

UP-1 formulas for trajectories in constant gravity g

 $x(t) = (v_0 \cos \alpha)t \qquad \qquad y(t) = (v_0 \sin \alpha)t - \frac{1}{2}gt^2$ $\dot{x}(0) = v_x(0) = v_0 \cos \alpha \qquad \qquad \dot{y}(0) = v_y(0) = v_0 \sin \alpha$

Substitute time $t=x/(v_0 \cos \alpha)$ into y(t)

$$y(x) = \frac{v_0 \sin \alpha}{v_0 \cos \alpha} x - \frac{g x^2}{2 v_0^2 \cos^2 \alpha}$$
$$y(x) = x \tan \alpha - \frac{g x^2}{2 v_0^2 \cos^2 \alpha}$$





Envelopes of the v_0 -trajectory region contain extremal *contact points* with each trajectory where: $\frac{\partial S(v_0, \alpha : x, y)}{\partial \alpha} = 0$



Envelopes of the v_0 -trajectory region contain extremal *contact points* with each trajectory where: $\frac{\partial S(v_0, \alpha; x, y)}{\partial \alpha} = 0$ $x \frac{\partial \tan \alpha}{\partial \alpha} - \frac{gx^2}{2v_0^2} \frac{\partial \cos^{-2} \alpha}{\partial \alpha} = 0 = \frac{x}{\cos^2 \alpha} - \frac{gx^2}{2v_0^2} \frac{2\sin \alpha}{\cos^3 \alpha}$



Envelopes of the v_0 -trajectory region contain extremal contact points with each trajectory where: $\frac{\partial S(v_0, \alpha : x, y)}{\partial \alpha} = 0$ $x \frac{\partial \tan \alpha}{\partial \alpha} - \frac{gx^2}{2v_0^2} \frac{\partial \cos^{-2} \alpha}{\partial \alpha} = 0 = \frac{x}{\cos^2 \alpha} - \frac{gx^2}{2v_0^2} \frac{2\sin \alpha}{\cos^3 \alpha}$ gives: $\tan \alpha = \frac{v_0^2}{gx}$ or: $x = \frac{v_0^2}{g \tan \alpha}$.



Envelopes of the v_0 -trajectory region contain extremal contact points with each trajectory where: $\frac{\partial S(v_0, \alpha; x, y)}{\partial \alpha} = 0$ $x \frac{\partial \tan \alpha}{\partial \alpha} - \frac{gx^2}{2v_0^2} \frac{\partial \cos^{-2} \alpha}{\partial \alpha} = 0 = \frac{x}{\cos^2 \alpha} - \frac{gx^2}{2v_0^2} \frac{2\sin \alpha}{\cos^3 \alpha} \qquad \tan \alpha = \frac{v_0^2}{gx} \text{ or: } x = \frac{v_0^2}{g \tan \alpha}.$ $y_{env}(x) = x \tan \alpha - \frac{gx^2}{2v_0^2} (1 + \tan^{2} \alpha) \Rightarrow y_{env}(x) = x \frac{v_0^2}{gx} - \frac{gx^2}{2v_0^2} (1 + \frac{v_0^4}{g^2 x^2})$



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Introducing the Poincare' and Legendre contact transformations Geometry of Legendre contact transformation (Preview of Unit 8 relativistic quantum mechanics) Example from thermodynamics Legendre transform: special case of General Contact Transformation (lights,camera, ACTION!)

An elementary contact transformation from sophomore physics Algebra-calculus development of "The Volcanoes of Io" and "The Atoms of NIST" Intuitive-geometric development of """"""""""""""""""""""""""""""" The Plumes of Prometheus NASA-Galileo Project Io fly-by on August 18, 1997



NASA Astronomy Picture of the Day - Io: The Prometheus Plume (Just Image)

NASA Galileo - Io's Alien Volcanoes

New Horizons - Volcanic Eruption Plume on Jupiter's moon IO

NASA Galileo - A Hawaiian-Style Volcano on Io
IO'S ALIEN VOLCANOES



Inform Inspire Involve science.nasa.gov Space Science News home Pretty bad sketch of plumes (LasVegas model of planetary ejecta?)

Io's ALIEN VOLCANOE's se guys need a geometry lesson?

Need to fly parabola Scientists are eager for a closer look at the solar system's strangest and nost active file geometry... volcanoes when Galileo flies by Io on October 11.

October 4, 1999: Thirty years ago, before the Voyager probes visited Jupiter, if you had described Io to a literary critic it would have been declared overwrought science fiction. Jupiter's strange moon is literally bursting with volcanoes. Dozens of active vents pepper the landscape which also includes gigantic frosty plains, towering mountains and volcanic rings the size of California. The volcanoes themselves are the hottest spots in the solar system with temperatures exceeding 1800 K (1527 C). The plumes which rise 300 km into space are so large they can be seen from Earth by the Hubble Space Telescope. Confounding common sense, these high-rising ejecta seem to be made up of, not blisteringly hot lava, but frozen sulfur



dioxide. And to top it all off, Io bears a striking resemblance to a pepperoni pizza. Simply unbelievable.

Right: Digital Radiance simulation of Pillan Patera just before the Galileo flyby. <u>click for animation</u> → .

NASA Astronomy Picture of the Day - Io: The Prometheus Plume (Just Image)

NASA Galileo - Io's Alien Volcanoes

New Horizons - Volcanic Eruption Plume on Jupiter's moon IO

NASA Galileo - A Hawaiian-Style Volcano on Io

...conventional parabolic geometry...carried to extremes...

Recall Lecture 6 p.26 and p. 48-49 for kite geometry and application



Unit 1 Fig. 9.4

























AMOP reference links (Updated list given on 2nd and 3rd pages of each class presentation)

<u>Web Resources - front page</u> UAF Physics UTube channel Quantum Theory for the Computer Age

Principles of Symmetry, Dynamics, and Spectroscopy

2014 AMOP 2017 Group Theory for QM 2018 AMOP

Classical Mechanics with a Bang!

Modern Physics and its Classical Foundations

Representaions Of Multidimensional Symmetries In Networks - harter-jmp-1973

Alternative Basis for the Theory of Complex Spectra

Alternative_Basis_for_the_Theory_of_Complex_Spectra_I - harter-pra-1973

Alternative Basis for the Theory of Complex Spectra II - harter-patterson-pra-1976

Alternative_Basis_for_the_Theory_of_Complex_Spectra_III_-_patterson-harter-pra-1977

Frame Transformation Relations And Multipole Transitions In Symmetric Polyatomic Molecules - RMP-1978

Asymptotic eigensolutions of fourth and sixth rank octahedral tensor operators - Harter-Patterson-JMP-1979

Rotational energy surfaces and high-J eigenvalue structure of polyatomic molecules - Harter - Patterson - 1984

Galloping waves and their relativistic properties - ajp-1985-Harter

Rovibrational Spectral Fine Structure Of Icosahedral Molecules - Cpl 1986 (Alt Scan)

Theory of hyperfine and superfine levels in symmetric polyatomic molecules.

- I) Trigonal and tetrahedral molecules: Elementary spin-1/2 cases in vibronic ground states PRA-1979-Harter-Patterson (Alt scan)
- II) Elementary cases in octahedral hexafluoride molecules Harter-PRA-1981 (Alt scan)

Rotation-vibration spectra of icosahedral molecules.

- I) Icosahedral symmetry analysis and fine structure harter-weeks-jcp-1989 (Alt scan)
- II) Icosahedral symmetry, vibrational eigenfrequencies, and normal modes of buckminsterfullerene weeks-harter-jcp-1989 (Alt scan)
- III) Half-integral angular momentum harter-reimer-jcp-1991

Rotation-vibration scalar coupling zeta coefficients and spectroscopic band shapes of buckminsterfullerene - Weeks-Harter-CPL-1991 (Alt scan) Nuclear spin weights and gas phase spectral structure of 12C60 and 13C60 buckminsterfullerene -Harter-Reimer-Cpl-1992 - (Alt1, Alt2 Erratum) Gas Phase Level Structure of C60 Buckyball and Derivatives Exhibiting Broken Icosahedral Symmetry - reimer-diss-1996

Fullerene symmetry reduction and rotational level fine structure/ the Buckyball isotopomer 12C 13C59 - jcp-Reimer-Harter-1997 (HiRez) Wave Node Dynamics and Revival Symmetry in Quantum Rotors - harter - jms - 2001

Molecular Symmetry and Dynamics - Ch32-Springer Handbooks of Atomic, Molecular, and Optical Physics - Harter-2006

Resonance and Revivals

- I) QUANTUM ROTOR AND INFINITE-WELL DYNAMICS ISMSLi2012 (Talk) OSU knowledge Bank
- II) <u>Comparing Half-integer Spin and Integer Spin Alva-ISMS-Ohio2013-R777 (Talks)</u>
- III) Quantum Resonant Beats and Revivals in the Morse Oscillators and Rotors (2013-Li-Diss)

Resonance and Revivals in Quantum Rotors - Comparing Half-integer Spin and Integer Spin - Alva-ISMS-Ohio2013-R777 (Talk)

Molecular Eigensolution Symmetry Analysis and Fine Structure - IJMS-harter-mitchell-2013

Quantum Revivals of Morse Oscillators and Farey-Ford Geometry - Li-Harter-cpl-2013

QTCA Unit 10 Ch 30 - 2013

AMOP Ch 0 Space-Time Symmetry - 2019

*Index/Search is disabled - a web based A.M.O.P. oriented reference page, with thumbnail/previews, greater control over the information display. <u>https://modphys.hosted.uark.edu/markup/AMOP_References.html</u> AMOP reference links (Updated list given on 2nd and 3rd pages of each class presentation)

(Int.J.Mol.Sci, 14, 714(2013) p.755-774, QTCA Unit 7 Ch. 23-26), (PSDS - Ch. 5, 7)

Int.J.Mol.Sci, 14, 714(2013), QTCA Unit 8 Ch. 23-25, QTCA Unit 9 Ch. 26, PSDS Ch. 5, PSDS Ch. 7

Intro spin ½ coupling <u>Unit 8 Ch. 24 p3</u> H atom hyperfine-B-level crossing <u>Unit 8 Ch. 24 p15</u>

Hyperf. theory <u>Ch. 24 p48.</u>

Hyperf. theory Ch. 24 p48. <u>Deeper theory ends p53</u>

Intro 2p3p coupling <u>Unit 8 Ch. 24 p17</u>. Intro LS-jj coupling <u>Unit 8 Ch. 24 p22</u>. CG coupling derived (start) <u>Unit 8 Ch. 24 p39</u>. CG coupling derived (formula) <u>Unit 8 Ch. 24 p44</u>. Lande' g-factor

<u>Unit 8 Ch. 24 p26</u>.

Irrep Tensor building <u>Unit 8 Ch. 25 p5</u>.

Irrep Tensor Tables Unit 8 Ch. 25 p12.

Wigner-Eckart tensor Theorem. <u>Unit 8 Ch. 25 p17</u>.

Tensors Applied to d,f-levels. <u>Unit 8 Ch. 25 p21</u>.

Tensors Applied to high J levels. <u>Unit 8 Ch. 25 p63</u>. *Intro 3-particle coupling.* <u>Unit 8 Ch. 25 p28</u>.

Intro 3,4-particle Young Tableaus <u>GrpThLect29 p42</u>.

Young Tableau Magic Formulae <u>GrpThLect29 p46-48</u>.

*Index/Search is disabled - a web based A.M.O.P. oriented reference page, with thumbnail/previews, greater control over the information display. <u>https://modphys.hosted.uark.edu/markup/AMOP_References.html</u> and eventually full on Apache-SOLR Index and search for nuanced, whole-site content/metadata level searching.

AMOP reference links (Updated list given on 2nd and 3rd and 4th pages of each class presentation)

Predrag Cvitanovic's: Birdtrack Notation, Calculations, and Simplification

Chaos_Classical_and_Quantum_- 2018-Cvitanovic-ChaosBook Group Theory - PUP_Lucy_Day_- Diagrammatic_notation_- Ch4 Simplification_Rules_for_Birdtrack_Operators_- Alcock-Zeilinger-Weigert-zeilinger-jmp-2017 Group Theory - Birdtracks_Lies_and_Exceptional_Groups_- Cvitanovic-2011 Simplification_rules_for_birdtrack_operators-_jmp-alcock-zeilinger-2017 Birdtracks for SU(N) - 2017-Keppeler

Frank Rioux's: <u>UMA</u> method of vibrational induction

Quantum_Mechanics_Group_Theory_and_C60 - Frank_Rioux - Department_of_Chemistry_Saint_Johns_U Symmetry_Analysis_for_H20-_H20GrpTheory-_Rioux Quantum_Mechanics-Group_Theory_and_C60 - JChemEd-Rioux-1994 Group_Theory_Problems-_Rioux-_SymmetryProblemsX Comment_on_the_Vibrational_Analysis_for_C60_and_Other_Fullerenes_Rioux-RSP

Supplemental AMOP Techniques & Experiment

Many Correlation Tables are Molien Sequences - Klee (Draft 2016)

High-resolution_spectroscopy_and_global_analysis_of_CF4_rovibrational_bands_to_model_its_atmospheric_absorption-_carlos-Boudon-jqsrt-2017 Symmetry and Chirality - Continuous_Measures_-_Avnir

Special Topics & Colloquial References

r-process_nucleosynthesis_from_matter_ejected_in_binary_neutron_star_mergers-PhysRevD-Bovard-2017