Lecture 26 Mon. 11.19.2018

Geometry and Symmetry of Coulomb Orbital Dynamics (Ch. 2-4 of Unit 5)

Rutherford scattering and hyperbolic orbit geometry Backward vs forward scattering angles and orbit construction example Parabolic "kite" and orbital envelope geometry Differential and total scattering cross-sections Eccentricity vector $\boldsymbol{\varepsilon}$ and $(\boldsymbol{\varepsilon}, \boldsymbol{\lambda})$ -geometry of orbital mechanics Projection $\boldsymbol{\varepsilon}$ •r geometry of $\boldsymbol{\varepsilon}$ -vector and orbital radius r

Review and connection to usual orbital algebra (previous lecture) Projection ε •**p** *geometry of* ε *-vector and momentum* **p**=m**v**

General geometric orbit construction using ε -vector and (γ, \mathbf{R}) -parameters Derivation of ε -construction by analytic geometry

Coulomb orbit algebra of ε -vector and Kepler dynamics of momentum $\mathbf{p}=m\mathbf{v}$ Example of complete (\mathbf{r},\mathbf{p}) -geometry of elliptical orbit

Connection formulas for (γ, \mathbf{R}) *-parameters with* (a, b) *and* (ε, λ)

Excerpts from Lect. 27

A running collection of links to course-relevant sites and articles

		Clusses
Comprehensive Harter-Soft Resource Listing	Classical Mechanics with a Bang!	<u>2014 AMOP</u>
UAF Physics YouTube channel	Quantum Theory for the Computer Age	2017 Group Theory for QM
LearnIt Physics Web Applications Pr	rinciples of Symmetry, Dynamics, and Spectroscopy	<u>2018 AMOP</u>
Neat external material to start the class: <u>AIP publications</u>	Modern Physics and its Classical Foundations	2018 Adv Mechanics
<u>AJP article on superball dynamics</u> <u>AAPT summer reading</u> These <i>are</i> hot off the presses: <u>Sorting ultracold atoms in a 3D optical lattice in a rea</u> <u>Synthetic three-dimensional atomic structures asser</u>	<u>AMOP Ch 0 Space-Time Symm</u> <u>Seminar at Rochester Institute</u> <u>Springer AMO Handbook - Ch 3</u> <u>alization of Maxwell's demon - Kumar-Nature-Letters-20</u> <u>mbled atom by atom - Berredo-Nature-Letters-2018</u>	<u>etry - 2019</u> e of Optics, Auxiliary slides, June 19, 2018 32 - Harter-Reimer-2019 18
Slightly Older ones: <u>Wave–particle duality of C60 molecules</u> <u>Optical vortex knots – One Photon at a Time</u>	"Relawavity" and quantum bas 2-CW laser wave - Bohrlt We Lagrangian vs Hamiltonian -	sis of <i>Lagrangian</i> & <i>Hamiltonian</i> mechanics: <u>eb App</u> <u>RelaWavity Web App</u>
 Older Links from Lectures 14-20 http://thearmchaircritic.blogspot.com/2011/11/punkin-chunkin.htt http://www.sussexcountyonline.com/news/photos/punkinchunkir Shooting-range-for-medieval-siege-weapons-Anybody-knows https://modphys.hosted.uark.edu/markup/TrebuchetWeb.html https://modphys.hosted.uark.edu/markup/TrebuchetWeb.html?s https://modphys.hosted.uark.edu/markup/TrebuchetWeb.html?s https://modphys.hosted.uark.edu/markup/TrebuchetWeb.html?s https://modphys.hosted.uark.edu/markup/TrebuchetWeb.html?s https://modphys.hosted.uark.edu/markup/TrebuchetWeb.html?s The trebuchet, Chevedden, Sci Am 1995 'Simple' Pendulum Sim: https://modphys.hosted.uark.edu/markup/C Google search on: "Satelite view of Patricia" (Images) Physics Girl Channel - Fun with Vortex Rings in the Pool iBall demo - Quasi-periodicity: https://youtu.be/_intDtULxDc https://modphys.hosted.uark.edu/markup/CoulltWeb.html?scenthttps://modphys.hosted.uark.edu/markup/CoulltWeb.html?scenthttps://modphys.hosted.uark.edu/markup/CoulltWeb.html?scenthttps://modphys.hosted.uark.edu/markup/CoulltWeb.html?scenthttps://modphys.hosted.uark.edu/markup/CoulltWeb.html?scenthttps://modphys.hosted.uark.edu/markup/CoulltWeb.html?scenthttps://modphys.hosted.uark.edu/markup/CoulltWeb.html?scenthttps://modphys.hosted.uark.edu/markup/CoulltWeb.html?scenthttps://modphys.hosted.uark.edu/markup/CoulltWeb.html?scenthttps://modphys.hosted.uark.edu/markup/CoulltWeb.html?scenthttps://modphys.hosted.uark.edu/markup/CoulltWeb.html?scenthttps://modphys.hosted.uark.edu/markup/CoulltWeb.html?scenthttps://modphys.hosted.uark.edu/markup/CoulltWeb.html?scenthttps://modphys.hosted.uark.edu/markup/CoulltWeb.html?scenthttps://modphys.hosted.uark.edu/markup/CoulltWeb.html?scenthttps://modphys.hosted.uark.edu/markup/CoulltWeb.html?scenthttps://modphys.hosted.uark.	Imile h.htmlOlder Links from Leace Advanced Atomic and Molect BoxIt Web Simulations Pure A-Type w/Cosine Pure B-Type w/Cosine Pure B-Type w/Cosine Pure B-Type w/Freq ratios Mixed AB-Type 2:1 Freq rat Pure A-Type A=4.9, B=0.0 Pure B-Type: A=4.0, B=0.0 Pure B-Type: A=4.0, B=0.0 Pure B-Type: A=4.0, B=0.0 Pure B-Type: A=4.0, B=0.0 Pure B-Type A_D=4.055, B Mixed AB-Type A_D=4.055, B Mixed AB-Type A_D=4.055, B Mixed AB-Type A_D=4.055, B Mixed AB-Type A_D=4.055, B Mixed AB Type A=4.0, BU2 Mixed AB Type A=4.0, BU2 Mixed AB Type A=4.0, BU2 Mixed AB Type A=4.0, BU2 Mixed AB Type A=5.086 B Mixed AB Type A=4.0, BU2 Mixed AB Type A=5.086 B Mixed AB Type A=5.086 B Mixed AB Type A=2.0, BU2 Mixed AB Type A=2.0, BU2 	ctures 21-23 cular Optical Physics 2018 Class #9, pages: $5, 61$ atio C=0, & D=4.0 2, C=0, & D=4.0 3=0, C=0.1 2=0.866, CU2=0, & D=1.0 w/Stokes & Freq rats =-0.27 C=0 D=2.024 w/Stokes plot B=0.2403 C=0.4162 D=4.277 w/Stokes plot x=0.325 B=0.375 C=0.825 D=0.05 w/Stokes plot Bang! 2018 velopment Mechanics - 2017 Lectures: <u>6</u> , <u>7</u> , <u>8</u> , mputer Age <u>Unit 3 Ch.7-10, page=90</u> W,V}, R=Reality) <u>https://www.babylonjs.com/</u>

continued γ

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Older Links from Lectures 24-25

Jerklt Web App: 2-, 2+, Amp50Omega147-, Amp50Omega296, Amp50Omega602, Gap(1) MolVibes Web App: C3vN3 Wavelt Web App: Dim = <u>3 w/Wave Components;</u> Static Char Table: <u>6</u>, <u>12</u>, <u>12(b)</u>, <u>16</u>, <u>36</u>, <u>256</u> Quantum Carpet with N=20: <u>Gaussian, Boxcar</u> Quantum Revivals of Morse Oscillators and Farey-Ford Geometry - Li-Harter-CPL-2015 QTCA Unit <u>5</u> Ch14 2013 Lester. R. Ford, Am. Math. Monthly <u>45</u>,586(<u>1938</u>) John Farey, Phil. Mag.(<u>1816</u>) <u>Wolfram</u> Harter, J. Mol. Spec. 210, <u>166-182</u> (2001) Harter, Li IMSS (2013) Li, Harter, Chem.Phys.Letters (2015) OscillatorPE Web App: IHO Scenario 2, Coulomb Scenario 3 RelaWavity Web App/Simulator/Calculator: <u>Elliptical - IHO orbits</u>

Links to supplement Lecture 26

Coullt Web App Simulations: p19, p32, p72, p73, p92, R=-0.375, R=+0.5, Rutherford

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Pick an "impact parameter" line y = b. Draw circle of radius a around center point C=(-a,b) tangent to y-axis. Draw "focus-locus" line OCF.



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Copy angle \angle BCF (equal to $\Phi/2$) to make angle \angle FCB' (also equal to $\Phi/2$) Resulting line CB' is outgoing asymptote at scattering angle Θ .



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Locate secondary focus O' by drawing circle around point C of diameter CO thru point O. Diameter O'CO is $2a\varepsilon$. Hyperbolic orbit points P now found using constant 2a=PO-PO'



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Rutherford scattering geometry





https://modphys.hosted.uark.edu/CoulItWeb.html?scenario=Rutherford

Chapter 1 Orbit Families and Action Families of particle orbits are drawn in a varying color which represents the classical action or Hamiltoon's characteristic function $SH = \int p$ Volcanoes of Io (Paths=180, No color quant.) Parabolic Fountain (Uniform) dq.(Sometimes SH is called 'reduced action'.) The color is chosen by first calculating c = SH modulo h-bar (You can change Planck's constant from Space Bomb (Coulomb) Exploding Starlet (IHO) its default value h/2n = 1.0) The chromatic value c assigns the hue by its position on the color wheel (e.g.; c=0 is red, c=0.2 is a yellow, c=0.5 is a Synchrotron Motion (Crossed E & B fields) green, etc.). Chapter 2 Rutherford Scattering A parallel beam of iso-energetic alpha particles undergo Rutherford scattering from a coulomb field of a nucleus as calculated in these demos. Rutherford scattering 2-Electron Orbits It is also the ideal pattern of paths followed by intergalactic hydrogen in perturbed by the solar wind. Chapter 3 Coulomb Field (H atom) Orbits in an attractive Coulomb field are calculated here. You may select the initial position (x(0),y(0)) by moving the mouse to a desired Atomic Orbits launch point, and then select the initial momentum (px(0), py(0)) by pressing the mouse button and dragging. **Chapter 4 Molecular Ion Orbits** Orbits around two fixed nuclei are calculated here. A set of elliptic coordinates are drawn in the background. After running a few trajectories Molecular Ion Orbits you may notice that their caustics conform to one or two of the elliptic coordinate lines. Oscillator Scattering) 2-Particle Orbits 2-Particle Collision











tangent slope=-5/2 slo

slope=1/2





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Also: Approximate model of deep-space H-atom scattering from solar wind as our Sun travels around galaxy. Lyman- α shock wave found just inside Mars orbital radius 2a~1.2Au.

Fig. 5.3.2 Family of iso-energetic Rutherford scattering orbits with varying impact parameter.

Incremental window $d\sigma = b \cdot db$ normal to beam axis at $x = -\infty$ scatters to area $dA = R^2 \sin \Theta d\Theta d\varphi = R^2 d\Omega$ onto a sphere at $R = +\infty$ where is called the *incremental solid angle* $d\Omega = \sin \Theta d\Theta d\varphi$



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Ratio $\frac{d\sigma}{d\Omega} = \frac{b \, db \, d\phi}{\sin \Theta d\Theta d\phi} = \frac{b}{\sin \Theta} \frac{db}{d\Theta}$ is called the *differential scattering cross-section (DSC)*



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Also: Approximate

H-atom scattering

as our Sun travels

Lyman- α *shock wave*

found just inside Mars

orbital radius 2a~1.2Au.

from solar wind

around galaxy.

model of deep-space

Ratio
$$\frac{d\sigma}{d\Omega} = \frac{b \, db \, d\phi}{\sin \Theta \, d\Theta \, d\phi} = \frac{b}{\sin \Theta} \frac{db}{d\Theta}$$
 is called the *differential scattering cross-section (DSC)*
Geometry: $b = a \cot \frac{\Theta}{2} = \frac{k}{2E} \cot \frac{\Theta}{2}$
with: $\frac{db}{d\Theta} = \frac{-a}{2} \csc^2 \frac{\Theta}{2} = \frac{-a}{2\sin^2 \frac{\Theta}{2}}$
(*Never forget*!: $a = \frac{-k}{2E}$)







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Isotropic field V=V(r) guarantees conservation *angular momentum vector* **L**

 $\mathbf{L} = \mathbf{r} \times \mathbf{p} = m \, \mathbf{r} \times \dot{\mathbf{r}}$

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Coulomb V = -k/r also conserves *eccentricity vector* ε

 $\mathbf{\hat{\varepsilon}} = \hat{\mathbf{r}} - \frac{\mathbf{p} \times \mathbf{L}}{km} = \frac{\mathbf{r}}{r} - \frac{\mathbf{p} \times (\mathbf{r} \times \mathbf{p})}{km}$

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Consider dot product of ε with a radial vector **r**:

 $\mathbf{\varepsilon} \bullet \mathbf{r} = \frac{\mathbf{r} \bullet \mathbf{r}}{r} - \frac{\mathbf{r} \bullet \mathbf{p} \times \mathbf{L}}{km} = r - \frac{\mathbf{r} \times \mathbf{p} \bullet \mathbf{L}}{km} = r - \frac{\mathbf{L} \bullet \mathbf{L}}{km}$

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 $\mathbf{A} = km \cdot \varepsilon$ is known as the *Laplace-Hamilton-Gibbs-Runge-Lenz vector*.

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...or of ε with momentum vector **p**:

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(...for sake of comparison...) IHO $V = (k/2)r^2$ also conserves Stokes vector S $S_A = \frac{1}{2}(x_1^2 + p_1^2 - x_2^2 - p_2^2)$ $S_B = x_1 p_1 + x_2 p_2$ $S_C = x_1 p_2 - x_2 p_1$



Isotropic field V=V(r) guarantees conservation *angular momentum vector* **L**

 $\mathbf{L} = \mathbf{r} \times \mathbf{p} = m \mathbf{r} \times \dot{\mathbf{r}}$

Coulomb V = -k/r also conserves *eccentricity vector* ε

$$\mathbf{\varepsilon} = \hat{\mathbf{r}} - \frac{\mathbf{p} \times \mathbf{L}}{km} = \frac{\mathbf{r}}{r} - \frac{\mathbf{p} \times (\mathbf{r} \times \mathbf{p})}{km}$$

(...for sake of comparison...) IHO $V = (k/2)r^2$ also conserves Stokes vector S $S_A = \frac{1}{2}(x_1^2 + p_1^2 - x_2^2 - p_2^2)$ $S_B = x_1 p_1 + x_2 p_2$ $S_C = x_1 p_2 - x_2 p_1$



Eccentricity vector $\boldsymbol{\varepsilon}$ and (ε, λ) -geometry of orbital mechanics Projection $\boldsymbol{\varepsilon} \cdot \mathbf{r}$ geometry of $\boldsymbol{\varepsilon}$ -vector and orbital radius \mathbf{r} Review and connection to usual orbital algebra (previous lecture) Projection $\boldsymbol{\varepsilon} \cdot \mathbf{p}$ geometry of $\boldsymbol{\varepsilon}$ -vector and momentum $\mathbf{p}=m\mathbf{v}$

General geometric orbit construction using ε -vector and (γ, R) -parameters Derivation of ε -construction by analytic geometry

Coulomb orbit algebra of ε -vector and Kepler dynamics of momentum $\mathbf{p}=m\mathbf{v}$ Example of complete (\mathbf{r}, \mathbf{p}) -geometry of elliptical orbit



Eccentricity vector ε and (ε,λ)-geometry of orbital mechanics
Projection ε•r geometry of ε-vector and orbital radius r
Review and connection to usual orbital algebra (previous lecture)
Projection ε•p geometry of ε-vector and momentum p=mv

General geometric orbit construction using ε -vector and (γ, R) -parameters Derivation of ε -construction by analytic geometry

Coulomb orbit algebra of ε -vector and Kepler dynamics of momentum $\mathbf{p}=m\mathbf{v}$ Example of complete (\mathbf{r}, \mathbf{p}) -geometry of elliptical orbit





NOTE: Lengths of vectors **p** and -**p** are not drawn so they correctly show that momentum $\mathbf{p}=m\mathbf{v}$ grows as radial distance $r=|\mathbf{r}|$ falls. (To be shown on p. 92-102)



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Dot product of ε with momentum vector **p**:

 $\boldsymbol{\varepsilon} \bullet \mathbf{p} = \frac{\mathbf{p} \bullet \mathbf{r}}{r} - \frac{\mathbf{p} \bullet \mathbf{p} \times \mathbf{L}}{km}$ $= \mathbf{p} \bullet \hat{\mathbf{r}} = p_r = \boldsymbol{\varepsilon} p_x$

This says: "Projection p_r of \mathbf{p} onto radial \mathbf{r} or \mathbf{r}' lines equals eccentricity $\boldsymbol{\varepsilon}$ times projection p_x of \mathbf{p} onto orbit major axis : $(\hat{\mathbf{x}} = \hat{\boldsymbol{\varepsilon}})$ "

Focal geometry demands: "Momentum \mathbf{p} must bisect angle $\measuredangle_{\mathbf{r}}^{\mathbf{r}}$, between radial \mathbf{r} or \mathbf{r}' lines."

 $p_r = \epsilon p_x$ $p_{\rm x}$ b ae a

Hyperbola has eccentricity $\varepsilon > 1$ (*Here*: $\varepsilon = 5/4 = 1.25$)

Eccentricity vector $\boldsymbol{\varepsilon}$ and (ε, λ) -geometry of orbital mechanics Projection $\boldsymbol{\varepsilon} \cdot \mathbf{r}$ geometry of $\boldsymbol{\varepsilon}$ -vector and orbital radius \mathbf{r} Review and connection to usual orbital algebra (previous lecture) Projection $\boldsymbol{\varepsilon} \cdot \mathbf{p}$ geometry of $\boldsymbol{\varepsilon}$ -vector and momentum $\mathbf{p} = m\mathbf{v}$

General geometric orbit construction using ε -vector and (γ, R) -parameters Derivation of ε -construction by analytic geometry

Coulomb orbit algebra of ε -vector and Kepler dynamics of momentum $\mathbf{p}=m\mathbf{v}$ Example of complete (\mathbf{r} , \mathbf{p})-geometry of elliptical orbit



Next several pages give step-by-step constructions of ε -vector and Coulomb orbit and trajectory physics

Fig. 5.4.2 Construction of eccentricity vector ε *and*

General geometric orbit construction using ε -vector and (γ, \mathbf{R}) -parameters

Pick launch point P (radius vector r) and elevation angle γ from radius (momentum initial p direction) inital momentum elevation angle γ r

General geometric orbit construction using ε -vector and (γ, \mathbf{R}) -parameters



General geometric orbit construction using ε -vector and (γ, \mathbf{R}) -parameters *Copy double angle 2* γ *(* \angle **FPQ***) onto* \angle **PFT** Copy F-center circle around launch point P Pick launch point P *Extend* ∠PFT *chord* PT *to make R***-ratio scale line** *Copy elevation angle* γ (\angle FPP') *onto* \angle P'PQ (radius vector **r**) and elevation angle γ from radius Extend resulting line QPQ' to make focus locus (momentum initial **p** direction) inital momentum р *wpied* elevation angle γ inital momentum D p elevation angle γ *copied* F **P**' inital momentum p elevation angle γ FOCUS Reason for focus locu Line **r** from 1st focus **F**/"reflects r off line **p** (or **P'P**) toward 2nd focus **F** somewhere so incident-angle γ equals reflected-angle γ









Eccentricity vector $\boldsymbol{\varepsilon}$ and (ε, λ) -geometry of orbital mechanics Projection $\boldsymbol{\varepsilon} \cdot \mathbf{r}$ geometry of $\boldsymbol{\varepsilon}$ -vector and orbital radius \mathbf{r} Review and connection to usual orbital algebra (previous lecture) Projection $\boldsymbol{\varepsilon} \cdot \mathbf{p}$ geometry of $\boldsymbol{\varepsilon}$ -vector and momentum $\mathbf{p} = m\mathbf{v}$

General geometric orbit construction using ε -vector and (γ, R) -parameters Derivation of ε -construction by analytic geometry

Coulomb orbit algebra of ε -vector and Kepler dynamics of momentum $\mathbf{p}=m\mathbf{v}$ Example of complete (\mathbf{r}, \mathbf{p}) -geometry of elliptical orbit














Play embedded movie with controls above.

or follow this link to try your hand at ε-construction using the CoulIt Web App

Just click and drag in main window to set new initial conditions. The Lenz vector will display as part of an overlay.

Rutherford scattering and hyperbolic orbit geometry Backward vs forward scattering angles and orbit construction example Parabolic "kite" and orbital envelope geometry Differential and total scattering cross-sections

Eccentricity vector $\boldsymbol{\varepsilon}$ and (ε, λ) -geometry of orbital mechanics Projection $\boldsymbol{\varepsilon} \cdot \mathbf{r}$ geometry of $\boldsymbol{\varepsilon}$ -vector and orbital radius \mathbf{r} Review and connection to usual orbital algebra (previous lecture) Projection $\boldsymbol{\varepsilon} \cdot \mathbf{p}$ geometry of $\boldsymbol{\varepsilon}$ -vector and momentum $\mathbf{p} = m\mathbf{v}$

General geometric orbit construction using ε -vector and (γ, R) -parameters Derivation of ε -construction by analytic geometry

Coulomb orbit algebra of ε -vector and Kepler dynamics of momentum $\mathbf{p}=m\mathbf{v}$ Example of complete (\mathbf{r},\mathbf{p}) -geometry of elliptical orbit

Connection formulas for (γ, \mathbf{R}) *-parameters with* (a, b) *and* (ε, λ)

Finding time derivatives of orbital coordinates r, ϕ , x, y, and eventually velocity **v** or momentum $\mathbf{p} = m\mathbf{v}$ Radius r: $\lambda = mr^2 \frac{d\phi}{dt} = mr^2 \dot{\phi}$

$$r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi}$$

Finding time derivatives of orbital coordinates r, ϕ , x, y, and eventually velocity **v** or momentum $\mathbf{p} = m\mathbf{v}$ Radius r: $\lambda = mr^2 \frac{d\phi}{dt} = mr^2 \dot{\phi}$

$$r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi}$$

$$\dot{\phi} = \frac{L}{mr^2} = \frac{L}{m} \frac{1}{r^2}$$

Finding time derivatives of orbital coordinates r, ϕ , x, y, and eventually velocity **v** or momentum $\mathbf{p} = m\mathbf{v}$ Radius r: Polar angle ϕ using: $L = mr^2 \frac{d\phi}{d\phi} = mr^2 \dot{\phi}$

$$r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi}$$

$$\dot{\phi} = \frac{L}{mr^2} = \frac{L}{m} \frac{1}{r^2} = \frac{L}{m} \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos \phi)^2$$

using:
$$\frac{1}{r^2} = \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos\phi)^2$$

Finding time derivatives of orbital coordinates r, ϕ , x, y, and eventually velocity **v** or momentum $\mathbf{p} = m\mathbf{v}$ Radius r: Polar angle ϕ using: $L = mr^2 \frac{d\phi}{\phi} = mr^2 \dot{\phi}$

$$r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi}$$
$$\dot{r} = \frac{dr}{dt} = \frac{L^2}{km} \frac{-\frac{d}{dt}(-\varepsilon \cos \phi)}{(1 - \varepsilon \cos \phi)^2}$$

Polar angle
$$\phi$$
 using: $L = mr^2 \frac{d\phi}{dt} = mr^2 \dot{\phi}$
 $\dot{\phi} = \frac{L}{mr^2} = \frac{L}{m} \frac{1}{r^2} = \frac{L}{m} \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos \phi)^2$

using:
$$\frac{1}{r^2} = \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos\phi)^2$$

Finding time derivatives of orbital coordinates r, ϕ , x, y, and eventually velocity **v** or momentum $\mathbf{p} = m\mathbf{v}$ Radius r: Polar angle ϕ using: $I = mr^2 \frac{d\phi}{\phi} = mr^2 \dot{\phi}$

$$r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi}$$

$$\dot{r} = \frac{dr}{dt} = \frac{L^2}{km} \frac{-\frac{d}{dt}(-\varepsilon \cos \phi)}{(1 - \varepsilon \cos \phi)^2}$$

$$Polar angle \phi using: \ L = mr^2 \frac{\pi}{dt} = mr^2 \phi$$

$$\dot{\phi} = \frac{L}{mr^2} = \frac{L}{m} \frac{1}{r^2} = \frac{L}{m} \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos \phi)^2$$

$$r\dot{\phi} = \frac{L}{mr}$$

$$using: \ \frac{1}{r^2} = \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos \phi)^2$$

Finding time derivatives of orbital coordinates r, ϕ , x, y, and eventually velocity **v** or momentum $\mathbf{p} = m\mathbf{v}$ Radius r: Polar angle ϕ using: $L = mr^2 \frac{d\phi}{d\phi} = mr^2 \dot{\phi}$

$$r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi}$$
$$\dot{r} = \frac{dr}{dt} = \frac{L^2}{km} \frac{-\frac{d}{dt}(-\varepsilon \cos \phi)}{(1 - \varepsilon \cos \phi)^2}$$

Polar angle
$$\phi$$
 using: $L = mr^2 \frac{d\phi}{dt} = mr^2 \dot{\phi}$
 $\dot{\phi} = \frac{L}{mr^2} = \frac{L}{m} \frac{1}{r^2} = \frac{L}{m} \left(\frac{km}{L^2}\right)^2 (1 - \epsilon \cos \phi)^2$
 $r\dot{\phi} = \frac{L}{mr} = \frac{L}{m} \frac{1}{r} = \frac{L}{m} \left(\frac{km}{L^2}\right) (1 - \epsilon \cos \phi) = \frac{k}{L} (1 - \epsilon \cos \phi)$
 $using: \frac{1}{r} = \left(\frac{km}{L^2}\right) (1 - \epsilon \cos \phi)$

Finding time derivatives of orbital coordinates r, ϕ , x, y, and eventually velocity **v** or momentum $\mathbf{p} = m\mathbf{v}$ Radius r: Polar angle ϕ using: $I - mr^2 \frac{d\phi}{\phi} - mr^2 \dot{\phi}$

$$r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi}$$
$$\dot{r} = \frac{dr}{dt} = \frac{L^2}{km} \frac{-\frac{d}{dt}(-\varepsilon \cos \phi)}{(1 - \varepsilon \cos \phi)^2}$$
$$\dot{r} = \frac{L^2}{km} \frac{-\varepsilon \sin \phi \dot{\phi}}{(1 - \varepsilon \cos \phi)^2}$$

Polar angle
$$\phi$$
 using: $L = mr^2 \frac{d\phi}{dt} = mr^2 \dot{\phi}$
 $\dot{\phi} = \frac{L}{mr^2} = \frac{L}{m} \frac{1}{r^2} = \frac{L}{m} \left(\frac{km}{L^2}\right)^2 (1 - \epsilon \cos \phi)^2$
 $r\dot{\phi} = \frac{L}{mr} = \frac{L}{m} \frac{1}{r} = \frac{L}{m} \left(\frac{km}{L^2}\right) (1 - \epsilon \cos \phi) = \frac{k}{L} (1 - \epsilon \cos \phi)$
 $using: \frac{1}{r^2} = \left(\frac{km}{L^2}\right)^2 (1 - \epsilon \cos \phi)^2$

Finding time derivatives of orbital coordinates r, ϕ , x, y, and eventually velocity **v** or momentum **p**=m**v**

$$\begin{aligned} \text{Radius } r: \\ r &= \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi} \\ \dot{r} &= \frac{dr}{dt} = \frac{L^2}{km} \frac{-\frac{d}{dt}(-\varepsilon \cos \phi)}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} &= \frac{dr}{dt} = \frac{L^2}{km} \frac{-\frac{d}{dt}(-\varepsilon \cos \phi)}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} &= \frac{L^2}{km} \frac{-\varepsilon \sin \phi \dot{\phi}}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} &= -\frac{L^2}{km} \left(\frac{km}{L^2}\right)^2 r^2 \dot{\phi} \varepsilon \sin \phi \end{aligned}$$

$$\begin{aligned} \text{Polar angle } \phi \text{ using: } L &= mr^2 \frac{d\phi}{dt} = mr^2 \dot{\phi} \\ \dot{\phi} &= \frac{L}{mr^2} = \frac{L}{m} \frac{1}{r^2} = \frac{L}{m} \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos \phi)^2 \\ r \dot{\phi} &= \frac{L}{mr} = \frac{L}{m} \frac{1}{r} = \frac{L}{m} \left(\frac{km}{L^2}\right) (1 - \varepsilon \cos \phi) = \frac{k}{L} (1 - \varepsilon \cos \phi)^2 \\ using: \frac{1}{r^2} &= \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos \phi)^2 \\ using: \frac{1}{(1 - \varepsilon \cos \phi)^2} &= \left(\frac{km}{L^2}\right)^2 r^2 \end{aligned}$$

Finding time derivatives of orbital coordinates r, ϕ , x, y, and eventually velocity **v** or momentum $\mathbf{p} = m\mathbf{v}$ Radius r: Polar angle ϕ using: $I - mr^2 \frac{d\phi}{\phi} - mr^2 \dot{\phi}$

Radius r:

$$r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi}$$

$$\dot{r} = \frac{dr}{dt} = \frac{L^2}{km} \frac{-\frac{d}{dt}(-\varepsilon \cos \phi)}{(1 - \varepsilon \cos \phi)^2}$$

$$\dot{r} = \frac{L^2}{km} \frac{-\frac{d}{dt}(-\varepsilon \cos \phi)}{(1 - \varepsilon \cos \phi)^2}$$

$$\dot{r} = -\frac{L^2}{km} \frac{-\varepsilon \sin \phi \dot{\phi}}{(1 - \varepsilon \cos \phi)^2}$$

$$r = -\frac{L^2}{km} \frac{km}{L^2} \frac{1}{2} r^2 \dot{\phi} \varepsilon \sin \phi$$

$$using: \frac{1}{(1 - \varepsilon \cos \phi)^2} = \left(\frac{km}{L^2}\right)^2 r^2$$

$$using: \frac{1}{(1 - \varepsilon \cos \phi)^2} = \left(\frac{km}{L^2}\right)^2 r^2$$

$$\dot{r} = -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi$$

$$again using: L = mr^2 \dot{\phi}$$

Finding time derivatives of orbital coordinates r, ϕ , x, y, and eventually velocity **v** or momentum **p**=m**v** $2 d\phi$ $2 \dot{\phi}$

$$\begin{aligned} \text{Radius } r: \\ r &= \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi} \\ \dot{r} &= \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi} \\ \dot{r} &= \frac{dr}{dt} = \frac{L^2}{km} \frac{-\frac{d}{dt} (-\varepsilon \cos \phi)}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} &= \frac{L^2}{km} \frac{-\frac{d}{dt} (-\varepsilon \cos \phi)}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} &= \frac{L^2}{km} \frac{-\varepsilon \sin \phi \dot{\phi}}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} &= -\frac{L^2}{km} \left(\frac{km}{L^2}\right)^2 r^2 \dot{\phi} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\sigma} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\sigma} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\sigma} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\sigma} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\sigma} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\sigma} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\sigma} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\sigma} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\sigma} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\sigma} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\sigma} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\sigma} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\sigma} \varepsilon \sin \phi \\ \dot{r} &= -\frac{k}{L^2} mr^2 \dot{\sigma} \varepsilon \sin \phi \\ \dot{r}$$

Finding time derivatives of orbital coordinates r, ϕ , x, y, and eventually velocity **v** or momentum **p**=m**v** Radius r:

Radius r:

$$r = \frac{\lambda}{1 - \varepsilon \cos\phi} = \frac{L^2/km}{1 - \varepsilon \cos\phi}$$

$$r = \frac{\lambda}{1 - \varepsilon \cos\phi} = \frac{L^2/km}{1 - \varepsilon \cos\phi}$$

$$\dot{r} = \frac{dr}{dt} = \frac{L^2}{km} \frac{-\frac{d}{dt}(-\varepsilon \cos\phi)}{(1 - \varepsilon \cos\phi)^2}$$

$$\dot{r} = \frac{L^2}{km} \frac{-\frac{d}{dt}(-\varepsilon \cos\phi)}{(1 - \varepsilon \cos\phi)^2}$$

$$\dot{r} = \frac{L^2}{km} \frac{-\varepsilon \sin\phi}{(1 - \varepsilon \cos\phi)^2}$$

$$\dot{r} = -\frac{L^2}{km} \left(\frac{km}{L^2}\right)^2 r^2 \dot{\phi} \varepsilon \sin\phi$$

$$using: \frac{1}{mr} = \frac{L}{mr} \frac{1}{mr} = \frac{L}{mr} \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos\phi)^2$$

$$using: \frac{1}{r^2} = \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos\phi)^2$$

$$\dot{r} = -\frac{L^2}{km} \left(\frac{km}{L^2}\right)^2 r^2 \dot{\phi} \varepsilon \sin\phi$$

$$using: \frac{1}{(1 - \varepsilon \cos\phi)^2} = \left(\frac{km}{L^2}\right)^2 r^2$$

$$\dot{r} = -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin\phi = -\frac{k}{L} \varepsilon \sin\phi$$

$$again using: L = mr^2 \dot{\phi}$$
Cartesian $x = r \cos\phi$:
$$\dot{x} = \frac{dx}{dt} = -\dot{r} \cos\phi - \sin\phi r\dot{\phi}$$

$$\dot{r} = -\frac{k}{L} \varepsilon \sin\phi \cos\phi - \sin\phi \frac{k}{L} (1 - \varepsilon \cos\phi)$$

$$\dot{r} = -\frac{k}{L} \varepsilon \sin\phi \sin\phi + \cos\phi r\dot{\phi}$$

$$\dot{r} = -\frac{k}{L} \varepsilon \sin\phi \sin\phi + \cos\phi \frac{k}{L} (1 - \varepsilon \cos\phi)$$

Finding time derivatives of orbital coordinates r, ϕ , x, y, and eventually velocity **v** or momentum $\mathbf{p} = m\mathbf{v}$ Radius r: Polar angle ϕ using: $\mathbf{I} = m\mathbf{v}^2 \frac{d\phi}{d\phi} = m\mathbf{v}^2 \dot{\phi}$

$$r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi}$$

$$r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi}$$

$$\dot{r} = \frac{dr}{dt} = \frac{L^2}{km} \frac{-\frac{d}{dt}(-\varepsilon \cos \phi)}{(1 - \varepsilon \cos \phi)^2}$$

$$\dot{r} = \frac{L}{mr} = \frac{L}{mr} \frac{1}{r^2} = \frac{L}{m} \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos \phi)^2$$

$$r = \frac{L}{mr} = \frac{L}{mr} \frac{1}{r} = \frac{L}{m} \left(\frac{km}{L^2}\right) (1 - \varepsilon \cos \phi) = \frac{k}{L} (1 - \varepsilon \cos \phi)^2$$

$$r = \frac{L^2}{km} \left(\frac{km}{L^2}\right)^2 r^2 \dot{\phi} \varepsilon \sin \phi$$

$$using: \frac{1}{(1 - \varepsilon \cos \phi)^2} = \left(\frac{km}{L^2}\right)^2 r^2$$

$$r = -\frac{L^2}{km} \left(\frac{km}{L^2}\right)^2 r^2 \dot{\phi} \varepsilon \sin \phi$$

$$using: \frac{1}{(1 - \varepsilon \cos \phi)^2} = \left(\frac{km}{L^2}\right)^2 r^2$$

$$r^2$$

$$r = -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = \left(\frac{k}{L} \varepsilon \sin \phi\right)$$

$$again using: L = mr^2 \dot{\phi}$$

$$Cartesian x = r \cos \phi:$$

$$\dot{x} = \frac{dx}{dt} = -\dot{r} \cos \phi - \sin \phi \left(\frac{k}{L}(1 - \varepsilon \cos \phi)\right)$$

$$= -\frac{k}{L} \sin \phi$$

$$r = -\frac{k}{L} \cos \phi - \sin \phi$$

$$r = -\frac{k}{L} \cos \phi$$

Finding time derivatives of orbital coordinates r, ϕ , x, y, and eventually velocity **v** or momentum **p**=m**v**

$$\begin{aligned} \text{Radius } r: \\ r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi} \\ \dot{r} = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi} \\ \dot{r} = \frac{dr}{dt} = \frac{L^2}{km} - \frac{d}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} = \frac{dr}{dt} = \frac{L^2}{km} - \frac{d}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} = -\frac{L^2}{km} - \frac{\varepsilon \sin \phi \phi}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} = -\frac{L^2}{km} - \frac{\varepsilon \sin \phi \phi}{L^2} \\ \dot{r} = -\frac{L^2}{km} - \frac{k}{L^2} - \frac{k}{2} \sin \phi \\ \dot{r} = -\frac{k}{L} - \frac{k}{mr^2} - \frac{k}{2} \sin \phi \\ \dot{r} = -\frac{k}{L} - \frac{k}{mr^2} - \frac{k}{2} \sin \phi \\ \dot{r} = -\frac{k}{L} - \frac{k}{mr^2} - \frac{k}{2} \sin \phi \\ \dot{r} = -\frac{k}{L} - \frac{k}{mr^2} - \frac{k}{2} \sin \phi \\ \dot{r} = -\frac{k}{L} - \frac{k}{mr^2} - \frac{k}{2} \sin \phi \\ \dot{r} = -\frac{k}{L} \cos \phi - \varepsilon \\ \dot{r} = -\frac{k}{L} \sin \phi \\ \dot{r} = -\frac{k}{L} \sin \phi \\ \dot{r} = -\frac{k}{L} \sin \phi \\ \dot{r} = -\frac{k}{L} \cos \phi - \varepsilon \\ \dot{r} = -\frac{k}{L} \sin \phi \\ \dot{r} = -\frac{k}{L} \cos \phi - \varepsilon \\ \dot{r} = -\frac{k}{L} \sin \phi \\ \dot{r} = -\frac{k}{L} \cos \phi - \varepsilon \\ \dot{r} = -\frac{k}{L} \sin \phi \\ \dot{r} = -\frac{k}{L} \cos \phi - \varepsilon \\ \dot{r} = -\frac{k}{L} \cos \phi - \varepsilon \\ \dot{r} = -\frac{k}{L} \sin \phi \\ \dot{r} = -\frac{k}{L} \cos \phi - \varepsilon \\ \dot{r} =$$

Rutherford scattering and hyperbolic orbit geometry Backward vs forward scattering angles and orbit construction example Parabolic "kite" and orbital envelope geometry Differential and total scattering cross-sections

Eccentricity vector $\boldsymbol{\varepsilon}$ and (ε, λ) -geometry of orbital mechanics Projection $\boldsymbol{\varepsilon} \cdot \mathbf{r}$ geometry of $\boldsymbol{\varepsilon}$ -vector and orbital radius \mathbf{r} Review and connection to usual orbital algebra (previous lecture) Projection $\boldsymbol{\varepsilon} \cdot \mathbf{p}$ geometry of $\boldsymbol{\varepsilon}$ -vector and momentum $\mathbf{p} = m\mathbf{v}$

General geometric orbit construction using ε -vector and (γ, R) -parameters Derivation of ε -construction by analytic geometry

Coulomb orbit algebra of ε-vector and Kepler dynamics of momentum p=mv
 Example of complete (r,p)-geometry of elliptical orbit

Connection formulas for (γ, \mathbf{R}) *-parameters with* (a, b) *and* (ε, λ)

















Algebra of ε -construction geometry The eccentricity parameter relates ratios $R = \frac{KE}{PE}$ and $\frac{b^2}{a^2}$

 $\varepsilon^2 = 1 + 4R(R+1)\sin^2\gamma$

and $\frac{b^2}{a^2}$ 1.Cartesian (a,b), 2.Physics (E,L), 3.Polar (ε, λ) Now we relate a 4th pair: 4.Initial (γ, R)

Three pairs of parameters for Coulomb orbits:

$$= 1 - \frac{b^2}{a^2} \quad \text{for ellipse} \quad (\varepsilon < 1) \quad \text{where:} \quad 4R(R+1)\sin^2\gamma = -\frac{b^2}{a^2} = \varepsilon^2 - 1$$
$$= 1 + \frac{b^2}{a^2} \quad \text{for hyperbola} \ (\varepsilon > 1) \quad \text{where:} \quad 4R(R+1)\sin^2\gamma = +\frac{b^2}{a^2} = \varepsilon^2 - 1$$





0 > R = KE/PE > -1 scale subtends angle 2γ with length $2r \sin\gamma$ as is derived before on p. 66-70. Note similarity of (**R**,**r**)-triangle in **r**-circle of radius r to that in **p**-circle of diameter p above. Rutherford scattering and hyperbolic orbit geometry Backward vs forward scattering angles and orbit construction example Parabolic "kite" and orbital envelope geometry Differential and total scattering cross-sections

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$$= 1 + \frac{b^{2}}{a^{2}} \quad \text{for hyperbola } (\varepsilon > 1)$$

Three pairs of parameters for Coulomb orbits: 1.Cartesian (a,b), 2.Physics (E,L), 3.Polar (ε , λ) Now we relate a 4th pair: 4.Initial (γ ,**R**) Algebra of ε -construction geometry The eccentricity parameter relates ratios $R = \frac{KE}{PE}$ and $\frac{b^2}{a^2}$ Three pairs of parameters for Coulomb orbits: 1.Cartesian (a,b), 2.Physics (E,L), 3.Polar (ε , λ) Now we relate a 4th pair: 4.Initial (γ ,**R**)

$$\varepsilon^{2} = 1 + 4R(R+1)\sin^{2}\gamma$$

$$= 1 - \frac{b^{2}}{a^{2}} \text{ for ellipse } (\varepsilon < 1) \text{ where: } 4R(R+1)\sin^{2}\gamma = -\frac{b^{2}}{a^{2}} = \varepsilon^{2} - 1 \text{ implying: } R(R+1) < 0$$

$$= 1 + \frac{b^{2}}{a^{2}} \text{ for hyperbola } (\varepsilon > 1) \text{ where: } 4R(R+1)\sin^{2}\gamma = +\frac{b^{2}}{a^{2}} = \varepsilon^{2} - 1 \text{ implying: } R(R+1) > 0$$

Three pairs of parameters for Coulomb orbits:Three pairs of parameters for Coulomb orbits:I. Cartesian (a,b), 2.Physics (E,L), 3.Polar (ε , λ)Now we relate a 4th pair: 4.Initial (γ , R) $\varepsilon^2 = 1 + 4R(R+1)\sin^2\gamma$ $= 1 - \frac{b^2}{a^2}$ for ellipse ($\varepsilon < 1$) where: $4R(R+1)\sin^2\gamma = -\frac{b^2}{a^2} = \varepsilon^2 - 1$ implying: R(R+1) < 0 (or: $-R^2 > R$) $= 1 + \frac{b^2}{a^2}$ for hyperbola ($\varepsilon > 1$) where: $4R(R+1)\sin^2\gamma = +\frac{b^2}{a^2} = \varepsilon^2 - 1$ implying: R(R+1) > 0 (or: $-R^2 < R$)
(or: -1 > R > 0)

Algebra of ε -construction geometry The eccentricty parameter relates ratios $R = \frac{KE}{PE}$ and $\frac{b^2}{a^2}$ Three pairs of parameters for Coulomb orbits: 1. Cartesian (a,b), 2. Physics (E,L), 3. Polar (ε, λ) Now we relate a 4th pair: 4. Initial (γ, R) $\varepsilon^2 = 1+4R(R+1)\sin^2\gamma$ $= 1 - \frac{b^2}{a^2}$ for ellipse $(\varepsilon < 1)$ where: $4R(R+1)\sin^2\gamma = -\frac{b^2}{a^2} = \varepsilon^2 - 1$ implying: R(R+1) < 0 (or: $-R^2 > R$) (or: 0 > R > -1) $= 1 + \frac{b^2}{a^2}$ for hyperbola $(\varepsilon > 1)$ where: $4R(R+1)\sin^2\gamma = +\frac{b^2}{a^2} = \varepsilon^2 - 1$ implying: R(R+1) > 0 (or: $-R^2 < R$) $(or: -R^2 < R)$ $(or: -R^2 < R)$
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Latus radius is similarly related:

$$\lambda = \frac{b^2}{a} = \mp 2r R \sin^2 \gamma$$





Eccentricity vector $\boldsymbol{\varepsilon}$ *and* (ε, λ) *-geometry of orbital mechanics* Analytic geometry derivation of ε -construction *Connection formulas for* (a,b) *and* (ε,λ) *with* (γ, R) Detailed ruler & compass construction of ε -vector and orbits $(R=-0.375 \ elliptic \ orbit)$ (R=+0.5 hyperbolic orbit)Properties of Coulomb trajectory families and envelopes Graphical ε -development of orbits Launch angle fixed-Varied launch energy - Launch energy fixed-Varied launch angle Launch optimization and orbit family envelopes



Excerpts from Lect. 27



Fig. 11. Sample computer trajectories. (a) Family of hyperbolic orbits with R = 1. (b) Family of hyperbolic orbits with R < 1.

Fig. 7. Coordinate grids for orbital analog computers. (a) Elliptical orbit scale (0 > R > -1). This can be used with the apparatus in Figs. 8 or 9. Radial lines marked $\pm 10^\circ$, $\pm 20^\circ$, ..., are each the focus locus for orbits with an initial velocity $\pm 10^\circ$, $\pm 20^\circ$, ..., above the horizon line. The circle marked 20° , 40° , ..., 340° can be taken as the Earth's surface, or any circle inside this one can be taken to be the surface of any celestial body. The *R* values apply correctly in either case, while the velocity values are marked for the former case only. (b) Hyperbolic orbit scale $(0 < R < \infty)$ and $(-\infty < R < -1)$. This can only be used with the apparatus shown in Fig. 9. Outer circles locate foci for orbits of particles attracted to the force center, while inner circles locate foci for orbits in a repulsive field. In either case a radial line marked $\pm 10^\circ$, $\pm 20^\circ$, ..., is the focus locus for an orbit with the initial velocity an angle $\pm 10^\circ$, $\pm 20^\circ$, ..., above the nadir line.



Coulomb envelope geometry











Fig. 5.4.5 in Unit 5 of CMwBANG!





<u>CoulIt Web Simulation Repulsive</u> <u>Coulomb Burst - Tight</u>





Eccentricity vector ε and (ε, λ) -geometry of orbital mechanics Analytic geometry derivation of ε -construction Connection formulas for (a,b) and (ε,λ) with (γ, \mathbf{R}) Detailed ruler & compass construction of ε -vector and orbits $(R=-0.375 \ elliptic \ orbit)$ $(R=+0.5 \ hyperbolic \ orbit)$ Properties of Coulomb trajectory families and envelopes \blacktriangleright Graphical ε -development of orbits Launch angle fixed-Varied launch energy Launch energy fixed-Varied launch angle \blacktriangleright Launch optimization and orbit family envelopes





Fig. 11. Sample computer trajectories. (a) Family of hyperbolic orbits with R = 1. (b) Family of hyperbolic orbits with R < 1.

(b)

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The Lenz vector and orbital analog computers*

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A single geometrical diagram based on the Lenz vector shows the qualitative and quantitative features of all three types of Coulomb orbits. A simple analog computer can be made for an overhead projector by using this theory, and a number of interesting effects can be efficiently demonstrated.

(1)

I. INTRODUCTION: THE ECCENTRICITY VECTOR

Occasionally, a geometrical construction and accompanying picture is worth a great deal more to the physicist or the physics student than pages of equations and solutions, especially now that computer graphics are so available. Since Newton's time the geometrical approach has come to be regarded as more clumsy than other methods of thought, and some very pretty pictures and proofs of physical phenomena have undoubtedly been lost. An example of such a construction involving Rutherford scattering was discussed in a recent article¹ by my students and myself, and the following is an improvement of this which describes general Coulombic orbit mechanics.

The generalization we shall describe below is based partly on a more recently discovered quantity called the Lenz-Runge vector^{2,3} or the "eccentricity" vector ϵ defined by Eq. (1). There **r** is the position vector of the orbiting particle, **L** is its angular momentum, **p** is its linear momentum, *m* is its mass, and *k* is the gravitational (or electrostatic) coefficient:

$$\epsilon = \mathbf{r}/r - \mathbf{L} \times \mathbf{p}/km.$$

Lately this quantity has received a flurry of attention in group theoretical studies of the hydrogen atom⁴; however, we shall use only its geometrical and classical properties.

In particular, the main property of ϵ is that it is a constant vector for any particle moving according to a Coulomb field. Vector ϵ points along the major axis of ellipse, parabola, or hyperbola, whichever is the appropriate orbit of the particle. Furthermore, the magnitude ϵ of this vector is the eccentricity of the orbit.

To show that this is consistent with the usual formulation, we take the dot product of this vector ϵ with the position **r** as in Eq. (2). This then reduces to the following equation (3) of a conic section in polar coordinates, which is the general orbit equation⁵:

$$\epsilon r \cos \theta = \epsilon \cdot \mathbf{r} = r - \mathbf{L} \times \mathbf{p} \cdot \mathbf{r} / km$$

$$= r + \mathbf{L} \cdot \mathbf{L} / km,$$

$$r = -(L^2 / km)(1 - \epsilon \cos \theta)^{-1}.$$
(3)

In Sec. II a simple geometric construction using these properties is shown to describe qualitatively and quantitatively the Coulomb orbits for all three cases: namely, the attractive case (k < 0) with positive energy, with negative energy, and the repulsive case (k > 0).

S(star)

Fig. 1. Initial position and momentum must be given before construction of the resulting orbit is possible.

Finally, it is shown how this construction leads to an analog computer of orbits that can be made for a few dollars to fit onto an overhead projector. This device can be appreciated by elementary classes (even *large* elementary classes if you use the right projector) when they know only a little about conic sections, since the more tedious mathematics is built into the device.

II. COMPUTING ORBITS BY RULER AND COMPASS

We start by simply listing three steps of an orbit construction while demonstrating their application to a particular case of a satellite orbiting a star. Then a general proof of the steps will be provided along with further discussion and applications.

Suppose you are given the initial position and velocity of a satellite relative to some very massive star. If these quantities are given in a pictorial form which shows the angle γ between momentum vector vector $\mathbf{p} = m\mathbf{v}$ and the radius line PS in Fig. 1, and if the magnitude of \mathbf{v} is given by the ratio R = T/V of the kinetic energy $(T = mv^2/2)$ to the potential energy (V = k/r), then the construction below proceeds immediately. Otherwise, these quantities must be calculated before proceeding. (In the potential energy of the star's gravity we have k = -GMm, where M is the star mass and G is the universal constant of gravitation.) Note that R is minus the squared ratio of initial velocity to the escape velocity in



Fig. 2. Doubling the angle between the momentum and the position vectors gives a line QP which must contain the orbit focus. MARKING PEN TIP WITH GROOVE

Fig. 8. Orbital computer design: cheaper model that computes elliptical orbits only using the scale of Fig. 7(a). A Plexiglas sheet that is about $\frac{1}{3}$ in. thick has a string hole at the orbit's center. A transparency (Xerox, 3-M, etc.) of Fig. 7(a) is taped in position on the underside. (*Caution:* One should avoid marking pens that permanently mark plastic.)

8 go with Fig. 7(a) and can be assembled in a few minutes by using odds and ends. A more sophisticated apparatus is shown in Fig. 9. The simple apparatus of Fig. 8 produces the well-known elliptical orbits and trajectories of planets and satellites, but not the hyperbolic trajectories characteristic of higher-than-escape velocity meteors or of the repulsive Coulomb force problems. The second apparatus (Fig. 9) is designed to handle all cases, provided that the appropriate focal point scale is inserted.

The operation of either plotter begins with the positioning of second focal point S' according to the scale on the plotting board. Then the marking pen is poked into a small identation at P and held while the strings to S and S' are tightened. Finally, you slide the pen out along the board in such a way that the strings stay tight and the desired trajectory is drawn.

The apparatus in Fig. 8 will thus make an ellipse since the sum of distances SP and S'P is constant. The apparatus in Fig. 9 does the same when the spool brakes are



Fig. 9. Orbital computer design: a more elaborate model that computes general Coulomb orbits. The ellipse drawing mode is obtained when, first, the clamp is opened to allow the string to slide and then the spool brakes are tightened after the initial adjustment has been made with the use of the scale in Fig. 7(a). The hyperbola drawing mode is obtained when the string clamp and spool clutch are tightened but the brakes are released. The spools must turn together after the initial adjustment has been made with the use of the scale in Fig. 7(b). One hand can maintain the string tension while drawing the orbit, and the other hand can control the paying out of string. (Alternately, springs on the spool axis accomplish the same thing.)





Fig. 10. Sample computer trajectory problem. One finds the minimumenergy trajectory for a given range ρ . Initial velocity v and θ follow easily from the geometry of the computer scale in Fig. 7(a).

tightened and the string clamp on the marking pen is loosened.

The apparatus in Fig. 9 will produce a hyperbola if the *difference* between distances SP and S'P is constant. This is accomplished by tightening the string clamp and the clutch so that two string spools reel equal amounts of string in or out and the constant difference in length is maintained.

IV. SOME USES FOR COMPUTERS

For the computer to be set up to draw elliptical orbits, there is one important question that can be answered immediately: To throw with *minimum* initial velocity a free falling spacecraft between two fixed points near the earth and a distance ρ apart, what initial angle θ and speed vare needed? (We imagine a fixed coordinate system here, not a rotating one.)

Measuring this range by a great circle angle ρ , we see that the focus of the orbit must be on a line through the Earth's center, making an angle $\rho/2$ with the launch point P. The smallest R circle [recall Fig. 7(a)] intersecting this line is the one tangent to it and represents the solution to the problem. Indeed, the algebraic solution to this problem follows from the diagram in Fig. 10 and is given there. Note that the angle θ approaches 45° as the range becomes small compared to the radius of the Earth.

Note that one may change the radius of the starting point P by simply reinterpreting the scale of the computer. For example, if the starting point is located at a height of, say, four times the Earth's radius, then the velocities marked on this scale are all divided by the square root of this factor, in this case, by 2.

With the computer set up to draw hyperbolic orbits and the appropriate scales available, there are a number of interesting problems to examine. For example, the attractive-field positive-energy scale allows one to exhibit the paths of meteorites. Given the impact direction and speed, one can extrapolate to find its origin.

The hyperbolic computer setup can be used to demonstrate Rutherford scattering for either the repulsive field (see Ref. 1) or the attractive field. At the same time