

Lecture 21

Fri. 11.04.2016

Introduction to coupled oscillation and eigenmodes

(Ch. 2-4 of Unit 4 11.12.15)

2D harmonic oscillator equations

Lagrangian and matrix forms and Reciprocity symmetry

2D harmonic oscillator equation eigensolutions

Geometric method

Matrix-algebraic eigensolutions with example $M = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix}$

Secular equation

Hamilton-Cayley equation and projectors

Idempotent projectors (how eigenvalues \Rightarrow eigenvectors)

Operator orthonormality and Completeness (Idempotent means: $\mathbf{P} \cdot \mathbf{P} = \mathbf{P}$)

Spectral Decompositions

Functional spectral decomposition

Orthonormality vs. Completeness vis-a-vis Operator vs. State

Lagrange functional interpolation formula

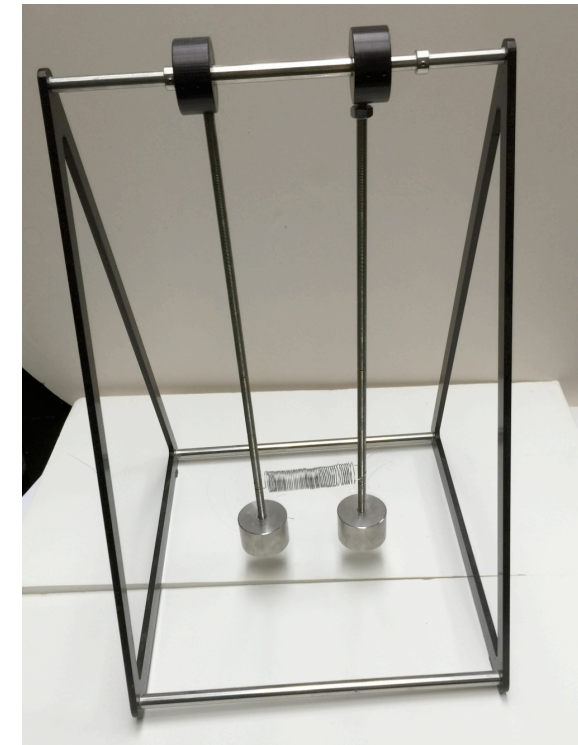
Diagonalizing Transformations (D-Tran) from projectors

2D-HO eigensolution example with bilateral (B-Type) symmetry

Mixed mode beat dynamics and fixed $\pi/2$ phase

2D-HO eigensolution example with asymmetric (A-Type) symmetry

Initial state projection, mixed mode beat dynamics with variable phase



➔ *2D harmonic oscillator equations*
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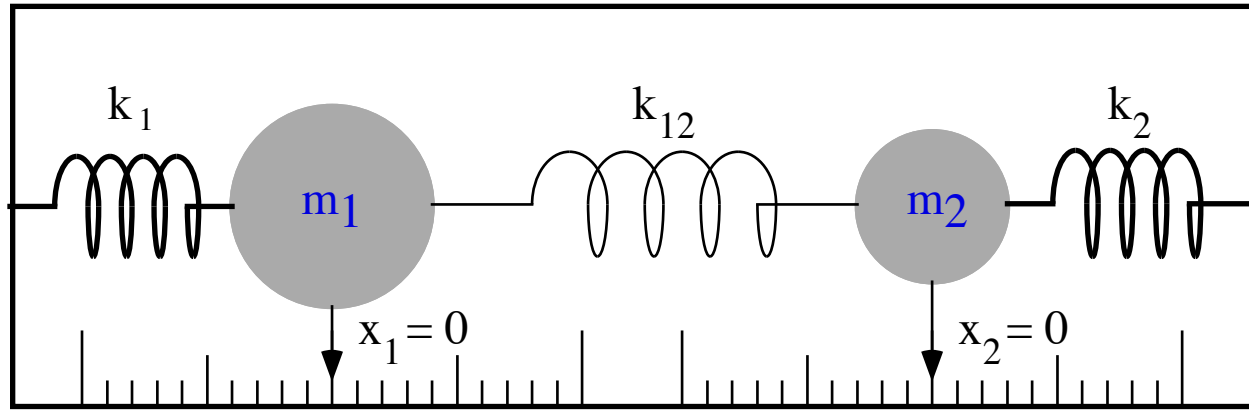


Fig. 3.3.1 Two 1-dimensional coupled oscillators

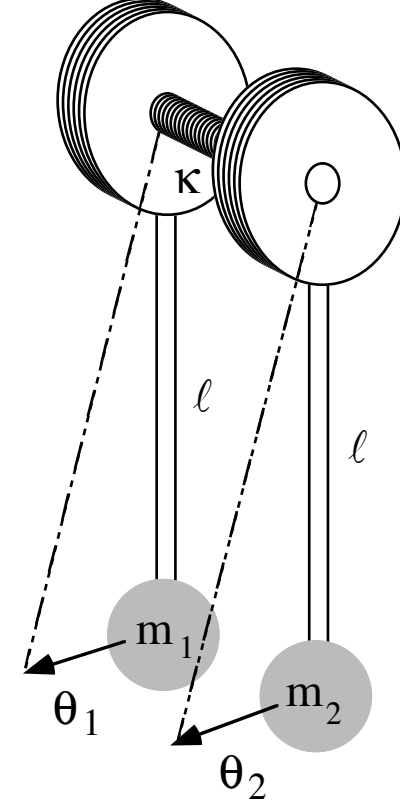
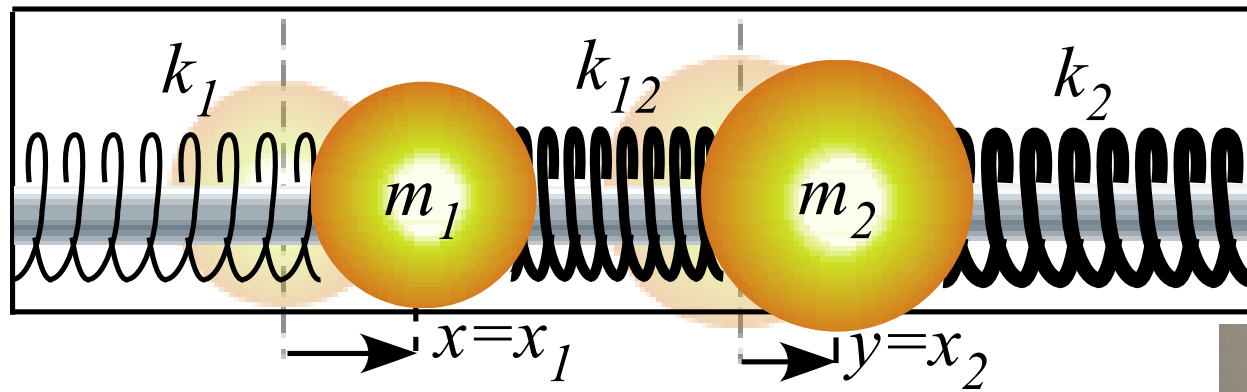


Fig. 3.3.2 Coupled pendulums

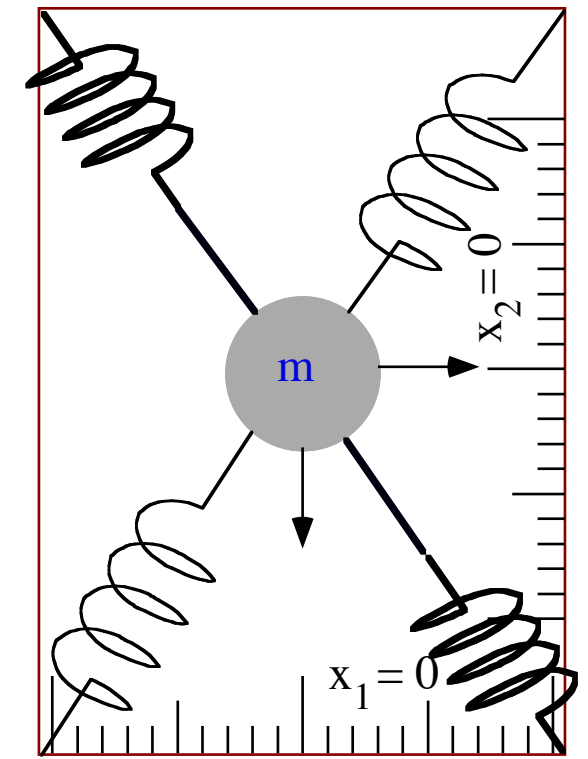


Fig. 3.3.3 One 2-dimensional coupled oscillator



2D harmonic oscillator energy

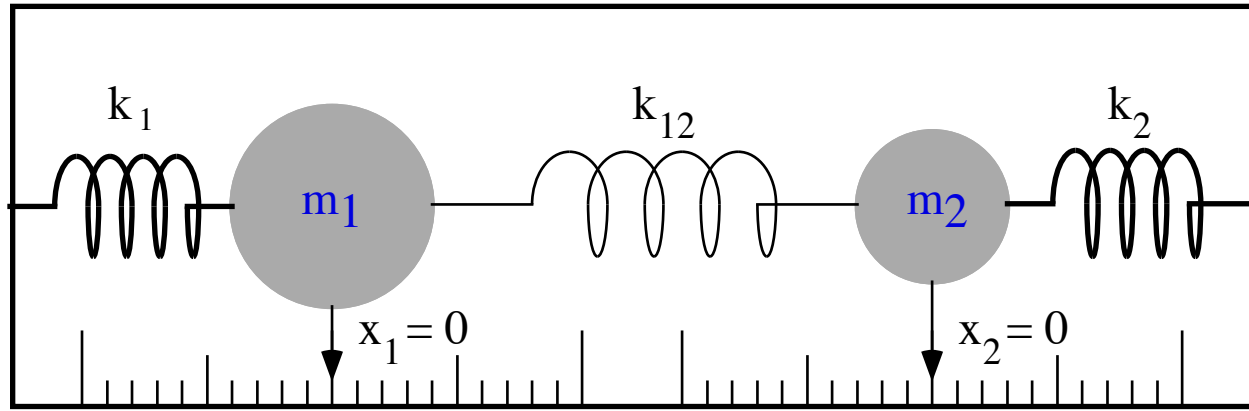


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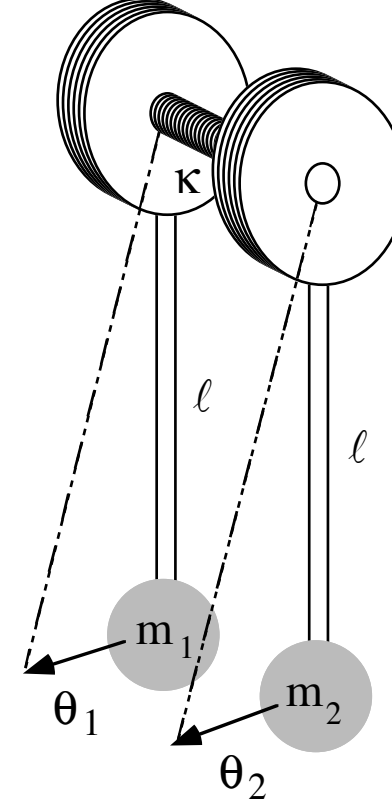
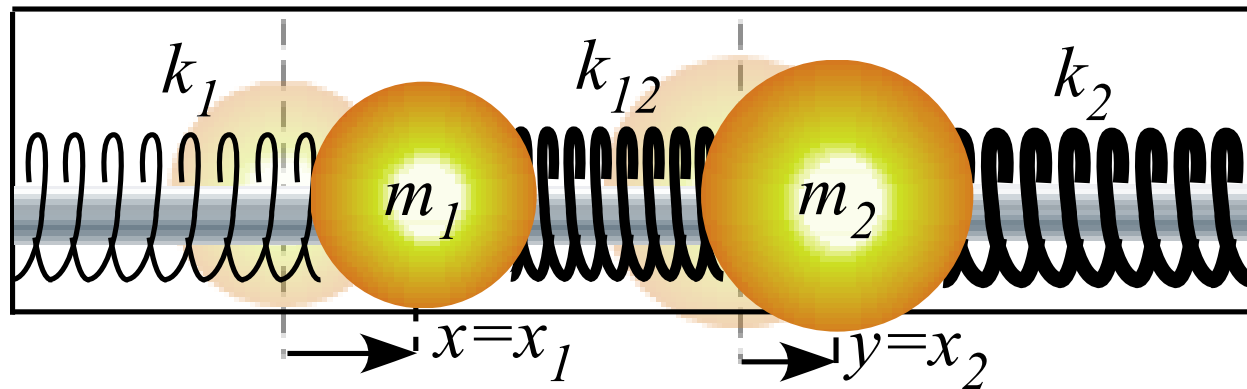


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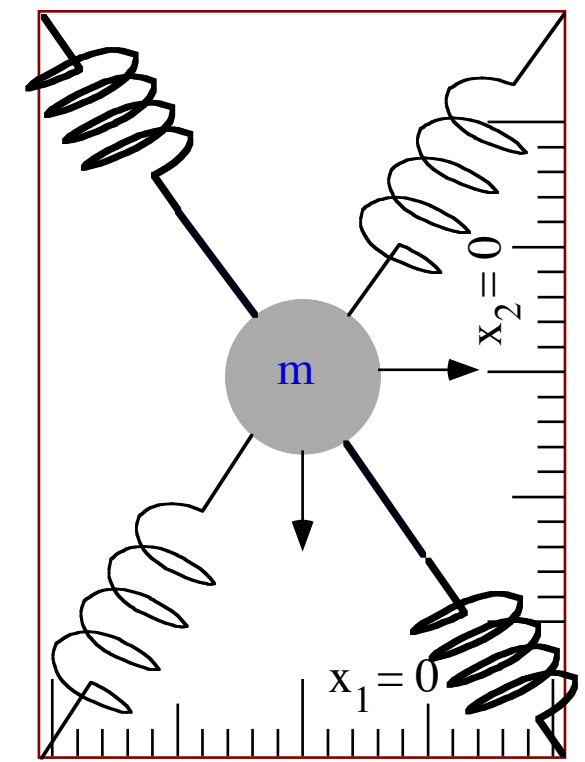


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2D HO kinetic energy $T(v_1, v_2)$

$$T = \frac{1}{2} m_1 \dot{x}_1^2 + \frac{1}{2} m_2 \dot{x}_2^2$$

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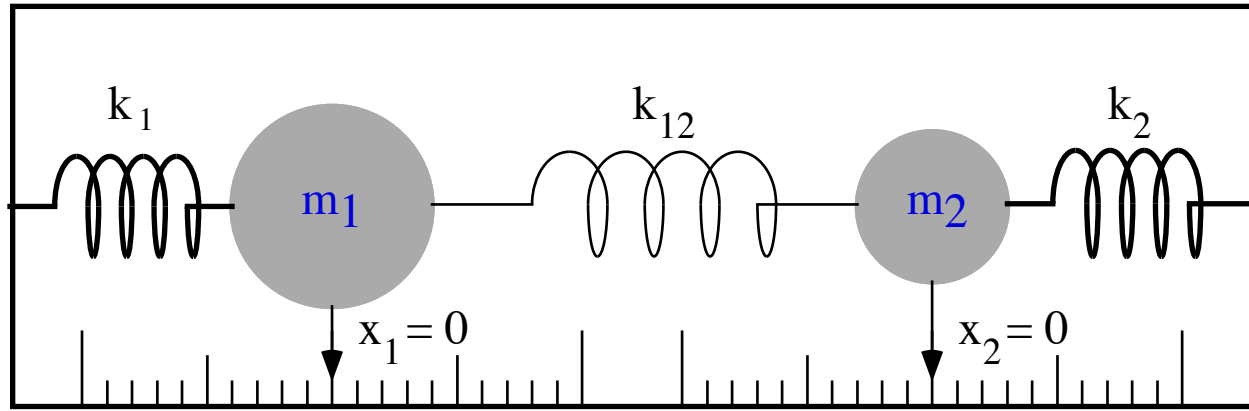


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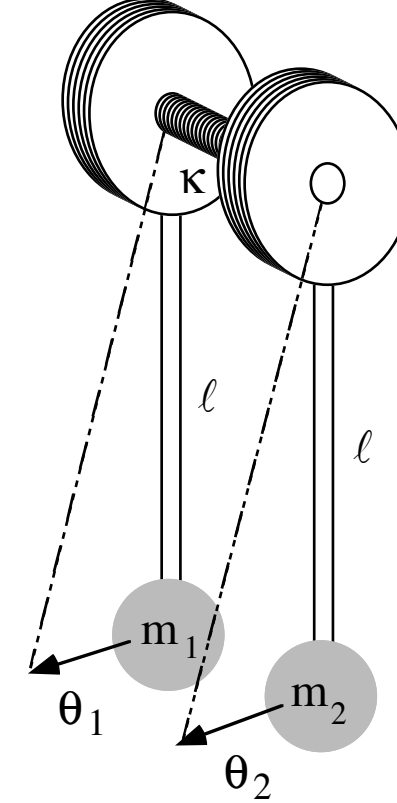
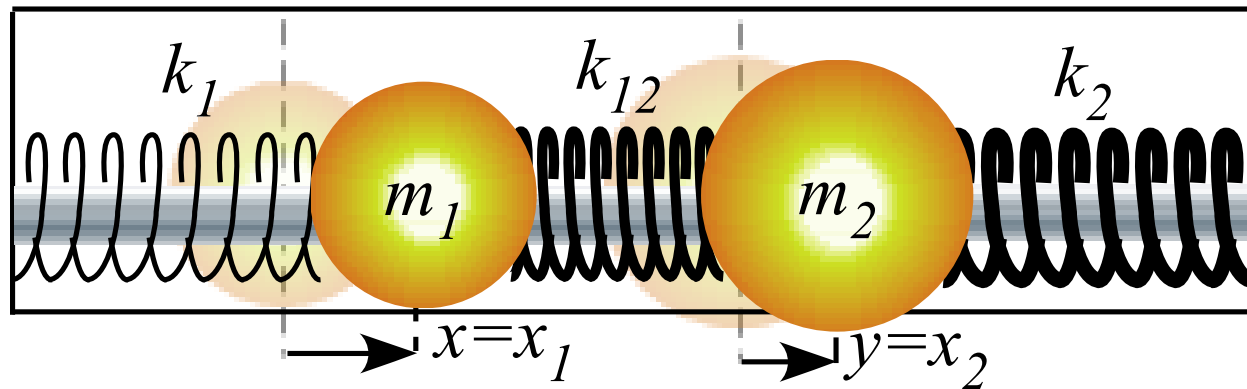


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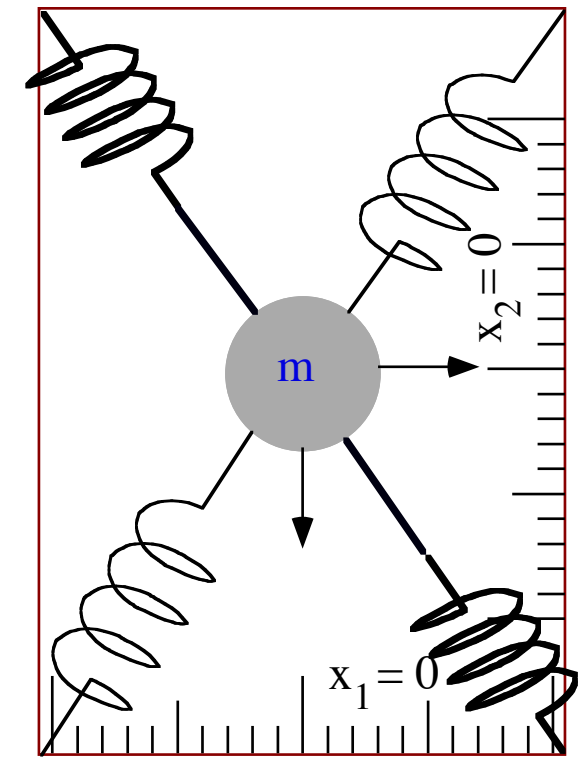


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$$\begin{aligned} V &= \frac{1}{2} k_1 x_1^2 + \frac{1}{2} k_2 x_2^2 + \frac{1}{2} k_{12} (x_1 - x_2)^2 \\ &= \frac{1}{2} (k_1 + k_{12}) x_1^2 - k_{12} x_1 x_2 + \frac{1}{2} (k_2 + k_{12}) x_2^2 \end{aligned}$$

Lagrangian $L=T-V$

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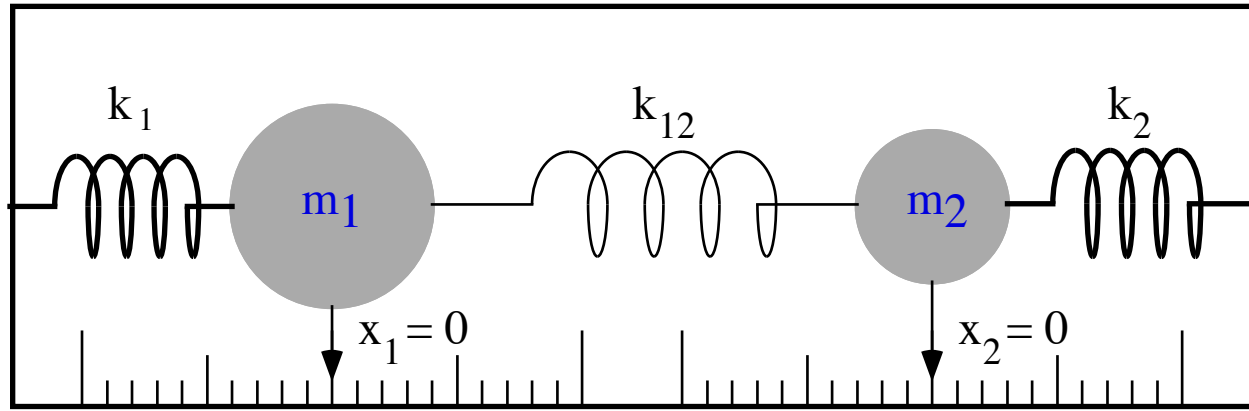


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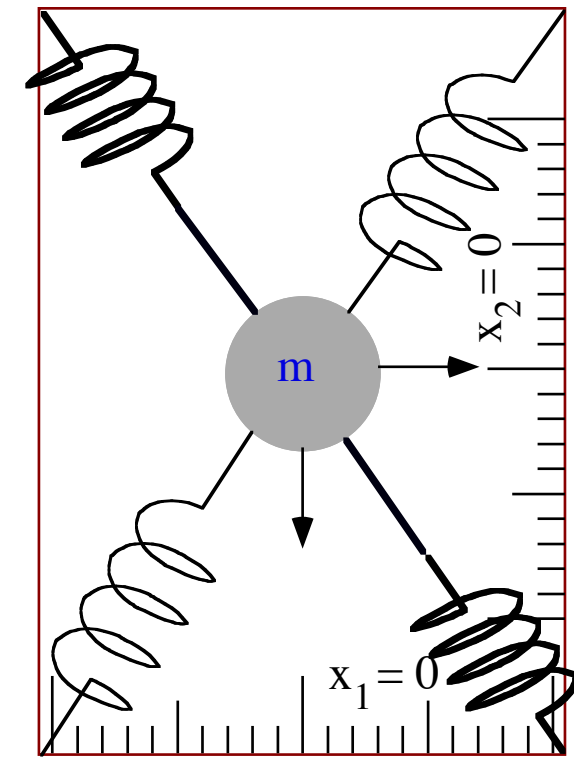
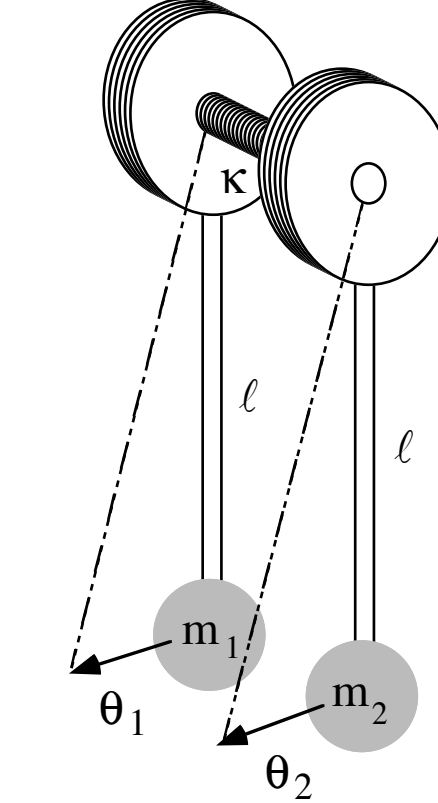
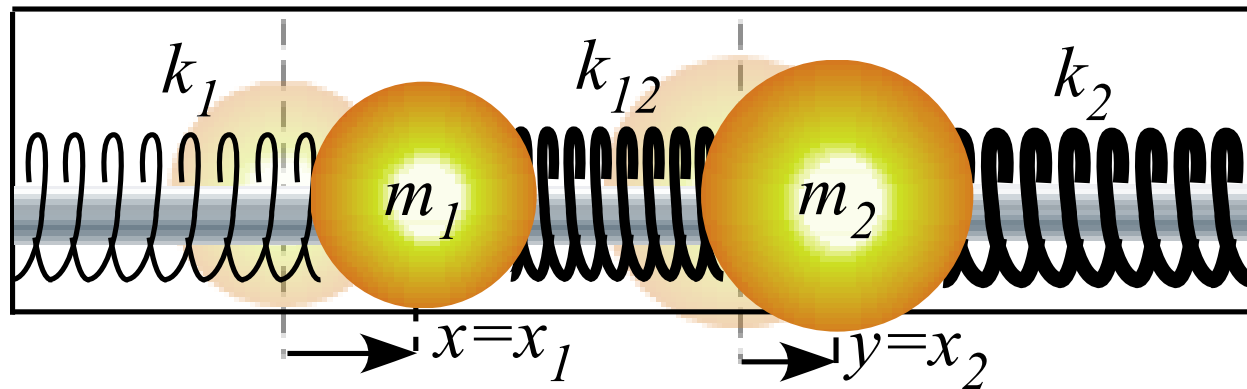


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2D HO Lagrange equations

$$\frac{d}{dt} \left(\frac{\partial T}{\partial \dot{x}_1} \right) = m_1 \ddot{x}_1 = F_1 = - \frac{\partial V}{\partial x_1} = - (k_1 + k_{12}) x_1 + k_{12} x_2$$

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2D harmonic oscillator equations

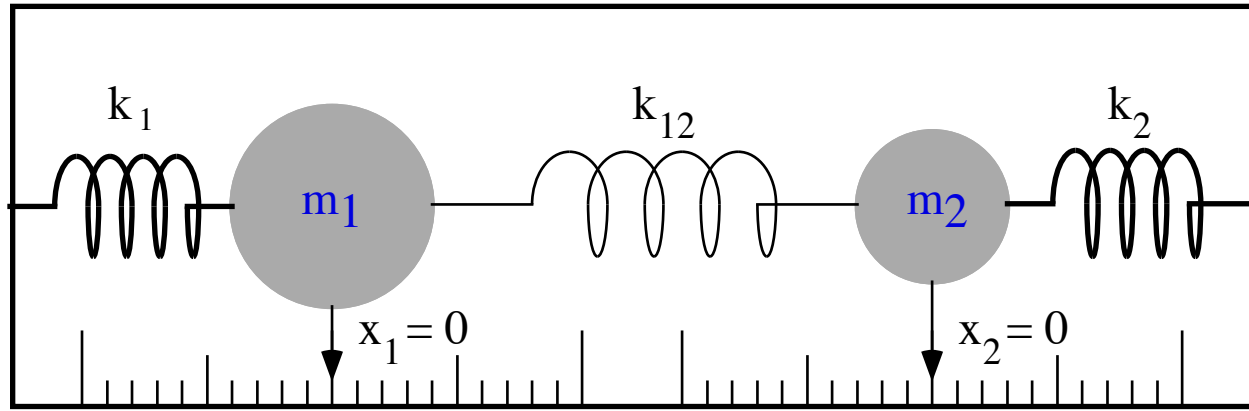


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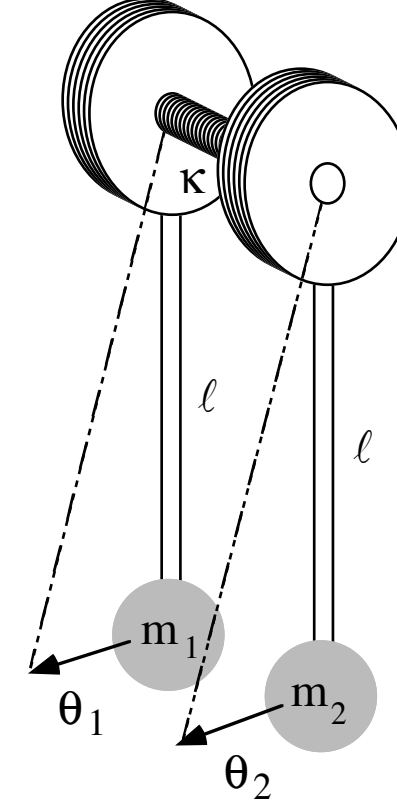
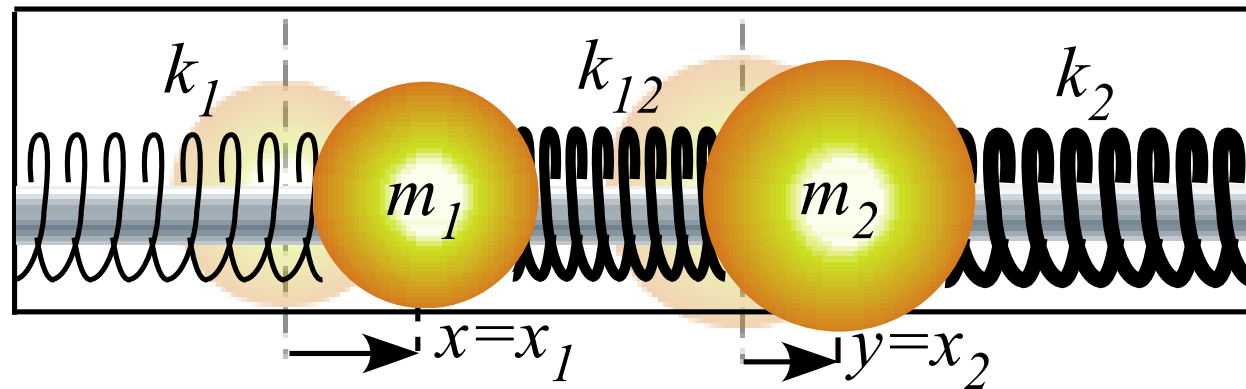


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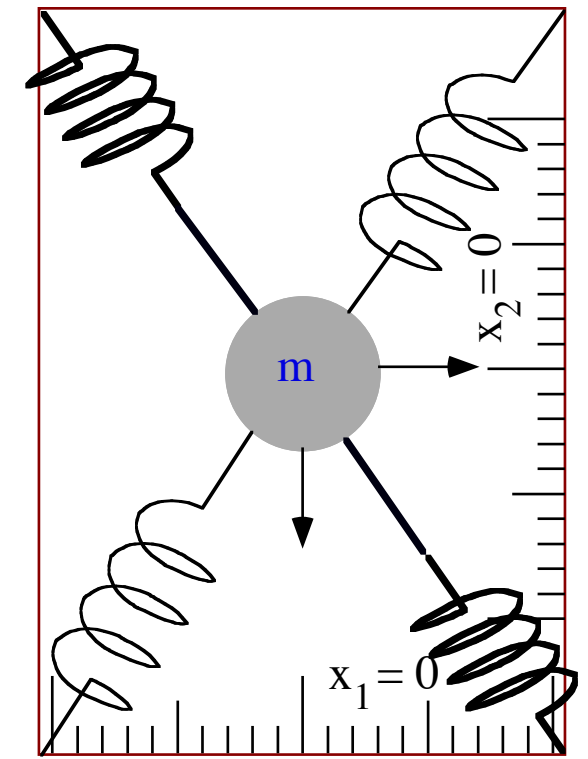


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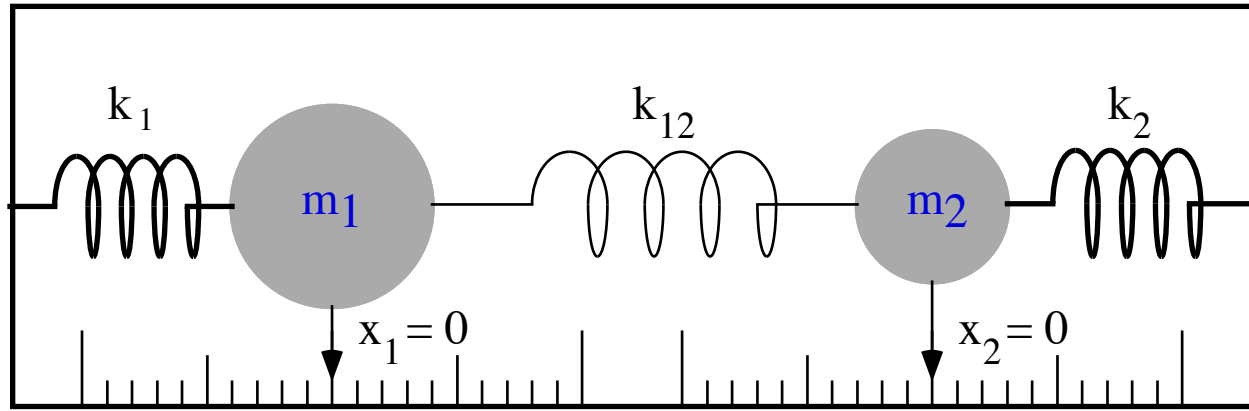


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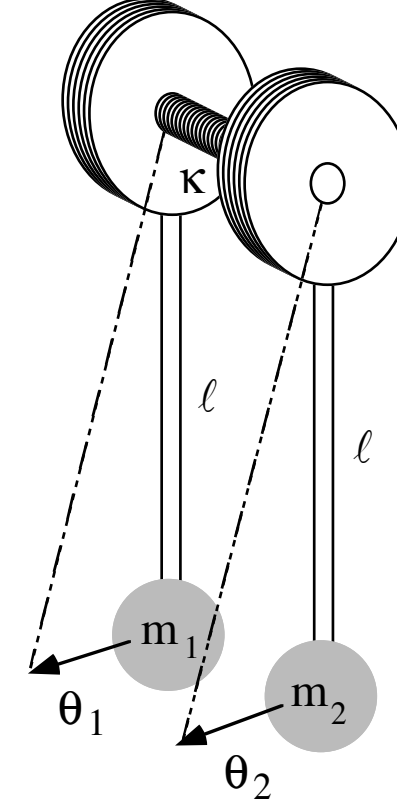
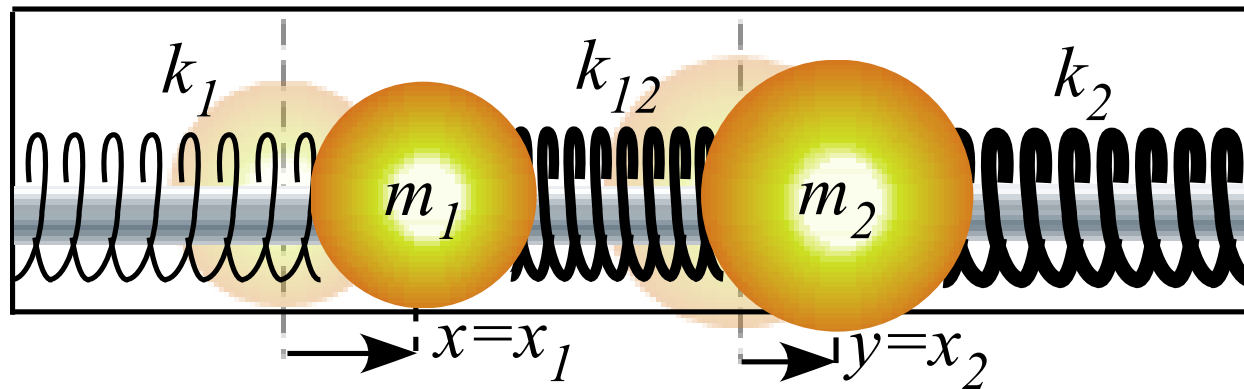


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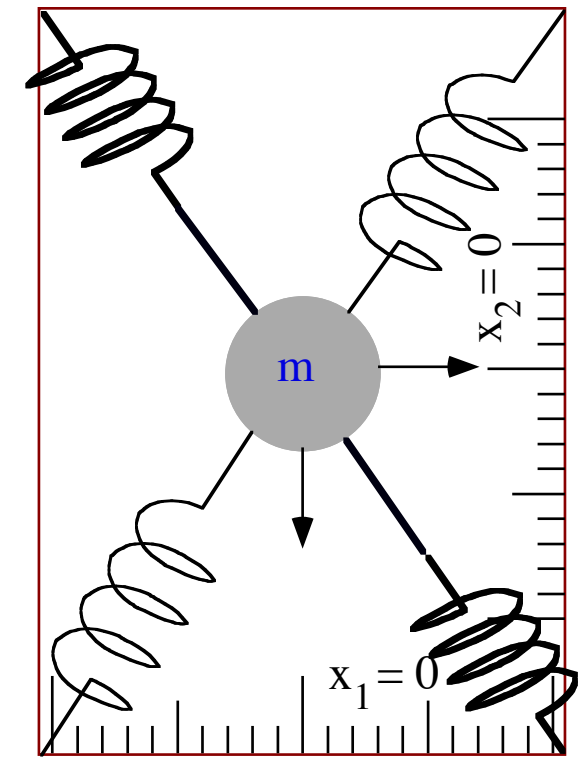


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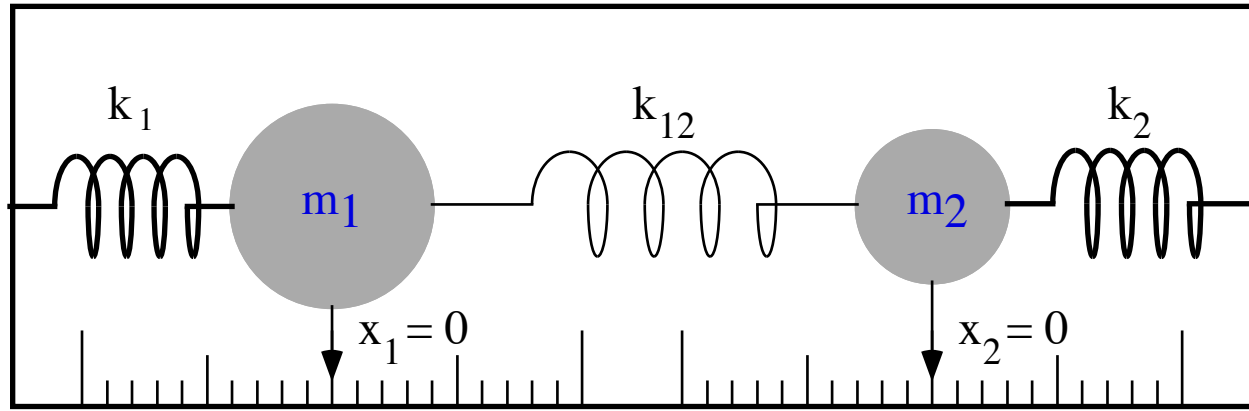


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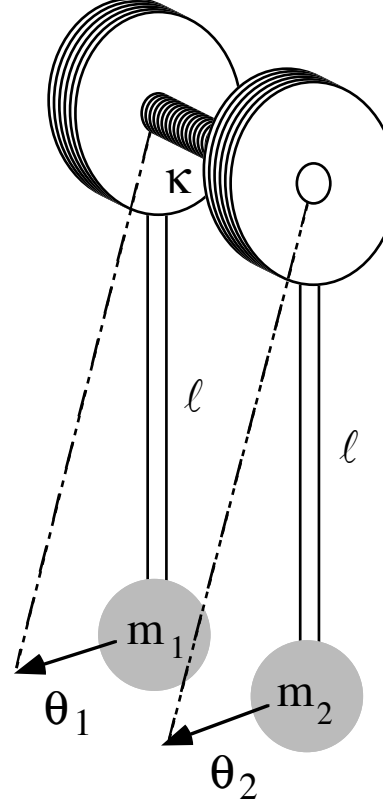
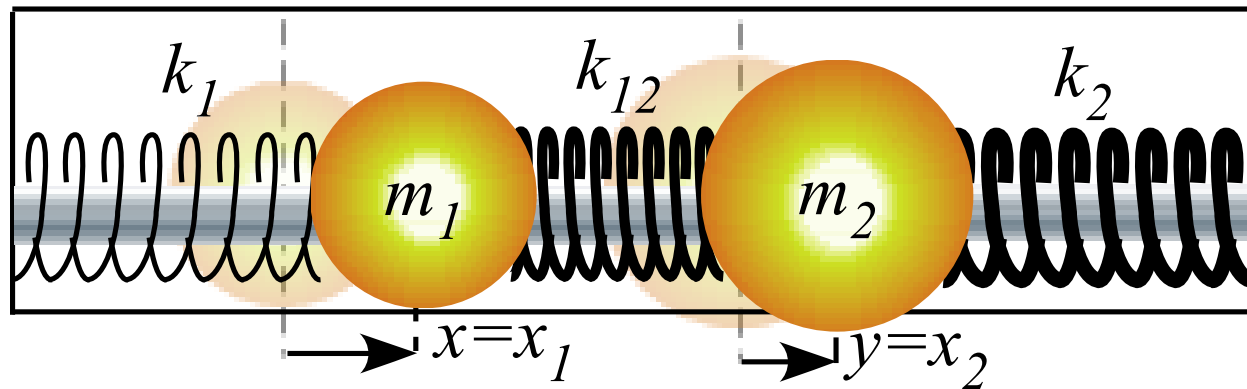


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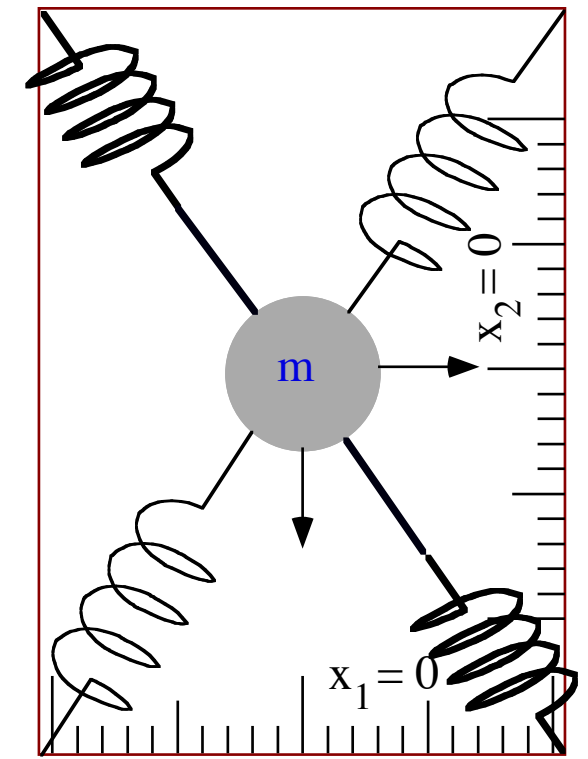


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where: $\mathbf{K} = \begin{pmatrix} k_1 + k_{12} & -k_{12} \\ -k_{12} & k_2 + k_{12} \end{pmatrix}$

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2D harmonic oscillator equations

Lagrangian and matrix forms \rightarrow *Reciprocity symmetry* \leftarrow

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Matrix equations and reciprocity symmetry

General form of 2D-HO equation of motion has force matrix components: $\kappa_{11} = k_1 + k_{12}$, $\kappa_{22} = k_2 + k_{12}$

$$\begin{pmatrix} F_1 \\ F_2 \end{pmatrix} = \begin{pmatrix} m_1 \ddot{x}_1 \\ m_2 \ddot{x}_2 \end{pmatrix} = - \begin{pmatrix} \kappa_{11} & \kappa_{12} \\ \kappa_{21} & \kappa_{22} \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}$$

Off-diagonal force constants satisfy *Reciprocity Relations*: $-\kappa_{12} = k_{12} = \frac{\partial F_1}{\partial x_2} = -\frac{\partial^2 V}{\partial x_2 \partial x_1} = -\frac{\partial^2 V}{\partial x_1 \partial x_2} = \frac{\partial F_2}{\partial x_1} = k_{21} = -\kappa_{21}$

Rescaling and symmetrization

Each coordinate (x_1, x_2) is rescaled $(q_1 = s_1 x_1, q_2 = s_2 x_2)$ to symmetrize mass factors on \ddot{q}_j -terms.

$$\begin{aligned} -\frac{m_1}{s_1} \ddot{q}_1 &= \kappa_{11} \frac{q_1}{s_1} + \kappa_{12} \frac{q_2}{s_2} & -\ddot{q}_1 &= \frac{\kappa_{11}}{m_1} q_1 + \frac{\kappa_{12} s_1}{m_1 s_2} q_2 \equiv \mathbf{K}_{11} q_1 + \mathbf{K}_{12} q_2 \\ -\frac{m_2}{s_2} \ddot{q}_2 &= \kappa_{12} \frac{q_1}{s_1} + \kappa_{22} \frac{q_2}{s_2} & -\ddot{q}_2 &= \frac{\kappa_{12} s_2}{m_2 s_1} q_1 + \frac{\kappa_{22}}{m_2} q_2 \equiv \mathbf{K}_{21} q_1 + \mathbf{K}_{22} q_2 \end{aligned}$$

New constants K_{ij} have pseudo-reciprocity symmetry for a special scale factor ratio: $\frac{s_2}{s_1} = \sqrt{\frac{m_2}{m_1}}$

$$\mathbf{K}_{21} = \frac{\kappa_{12} s_2}{m_2 s_1} = \mathbf{K}_{12} = \frac{\kappa_{12} s_1}{m_1 s_2} = \frac{-k_{12}}{\sqrt{m_1 m_2}}$$

Diagonal constants K_{jj} are not affected by scaling. To be equal requires: $\frac{\mathbf{K}_{11}}{m_1} = \frac{\mathbf{K}_{22}}{m_2}$ or: $\frac{\mathbf{K}_{11}}{\mathbf{K}_{22}} = \frac{m_1}{m_2}$

$$\mathbf{K}_{11} = \frac{\kappa_{11}}{m_1} = \frac{k_1 + k_{12}}{m_1} \quad \mathbf{K}_{22} = \frac{\kappa_{22}}{m_2} = \frac{k_2 + k_{12}}{m_2}$$

Caution is advised since such forced symmetry may give modes with imaginary frequency.

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2D harmonic oscillator equation solutions

1. May rewrite equation $\mathbf{M} \cdot |\ddot{\mathbf{x}}\rangle = -\mathbf{K} \cdot |\mathbf{x}\rangle$ in acceleration matrix form: $|\ddot{\mathbf{x}}\rangle = -\mathbf{A}|\mathbf{x}\rangle$ where: $\mathbf{A} = \mathbf{M}^{-1} \cdot \mathbf{K}$

$$\begin{pmatrix} \ddot{x}_1 \\ \ddot{x}_2 \end{pmatrix} = - \begin{pmatrix} m_1 & 0 \\ 0 & m_2 \end{pmatrix}^{-1} \begin{pmatrix} k_1 + k_{12} & -k_{12} \\ -k_{12} & k_2 + k_{12} \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = - \begin{pmatrix} \frac{k_1 + k_{12}}{m_1} & \frac{-k_{12}}{m_1} \\ \frac{-k_{12}}{m_2} & \frac{k_2 + k_{12}}{m_2} \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}$$

2. Need to find *eigenvectors* $|\mathbf{e}_1\rangle, |\mathbf{e}_2\rangle, \dots$ of acceleration matrix such that: $\mathbf{A}|\mathbf{e}_n\rangle = \varepsilon_n|\mathbf{e}_n\rangle = \omega_n^2|\mathbf{e}_n\rangle$

Then equations decouple to: $|\ddot{\mathbf{e}}_n\rangle = -\mathbf{A}|\mathbf{e}_n\rangle = -\varepsilon_n|\mathbf{e}_n\rangle = -\omega_n^2|\mathbf{e}_n\rangle$ where ε_n is an *eigenvalue*

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To introduce eigensolutions we take a simple case of unit masses ($m_1=1=m_2$)

So equation of motion is simply: $|\ddot{\mathbf{x}}\rangle = -\mathbf{K}|\mathbf{x}\rangle$

Eigenvectors $|\mathbf{x}\rangle = |\mathbf{e}_n\rangle$ are in special directions where $|\ddot{\mathbf{x}}\rangle = -\mathbf{K}|\mathbf{x}\rangle$ is in same direction as $|\mathbf{x}\rangle$

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➔ *Geometric method*



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2D HO potential energy $V(x_1, x_2)$ quadratic form defines layers of elliptical V -contours

$$V = \frac{1}{2}(k_1 + k_{12})x_1^2 - k_{12}x_1x_2 + \frac{1}{2}(k_2 + k_{12})x_2^2 = \frac{1}{2}\langle \mathbf{x} | \mathbf{K} | \mathbf{x} \rangle = \frac{1}{2}\mathbf{x} \cdot \mathbf{K} \cdot \mathbf{x} = \begin{pmatrix} x_1 & x_2 \end{pmatrix} \begin{pmatrix} k_1 + k_{12} & -k_{12} \\ -k_{12} & k_2 + k_{12} \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}$$

(a) PE Contours

What direction $|\mathbf{x}\rangle = |\mathbf{e}_n\rangle$
is the *same* as $\mathbf{K}|\mathbf{x}\rangle$??
Not most directions!

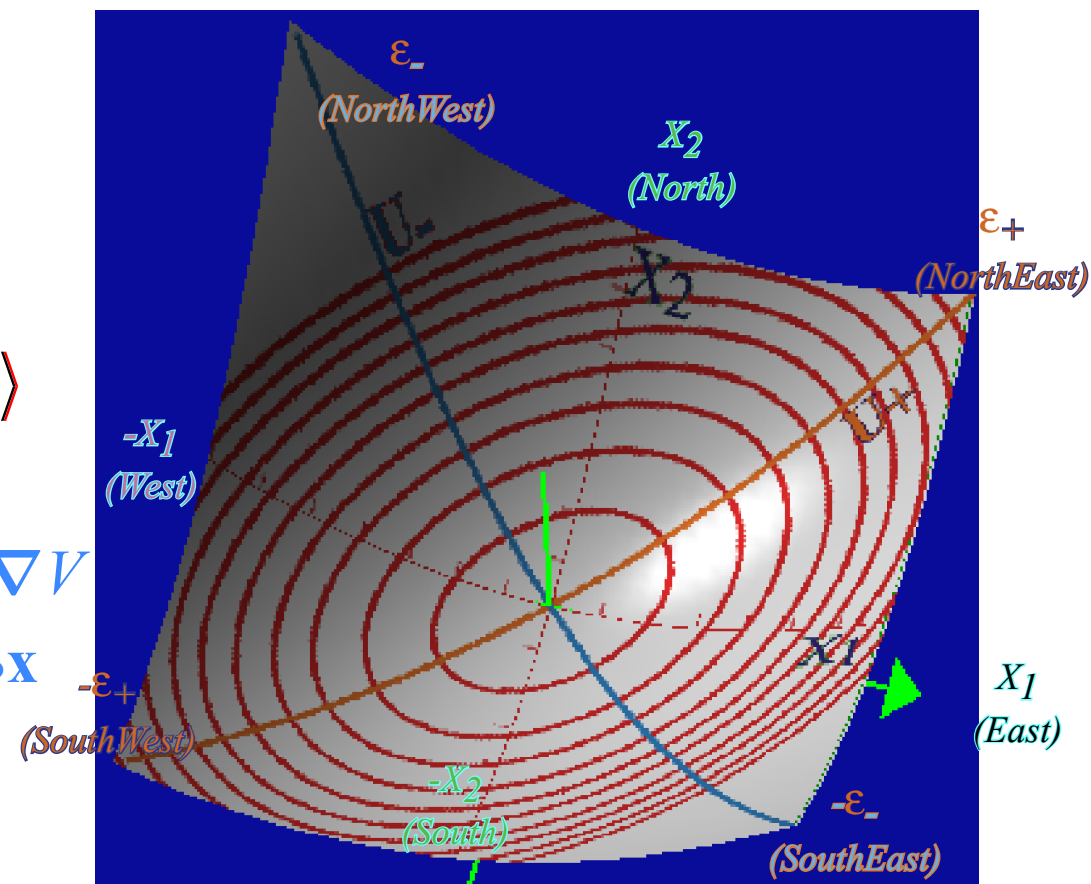
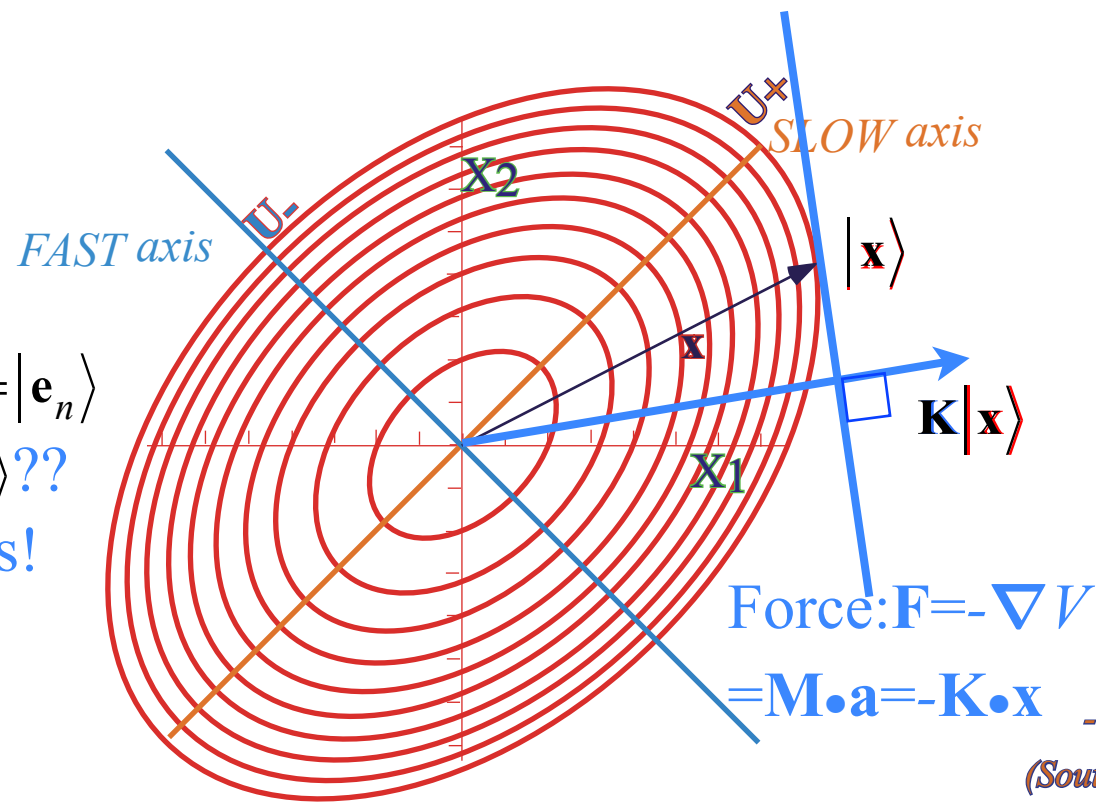
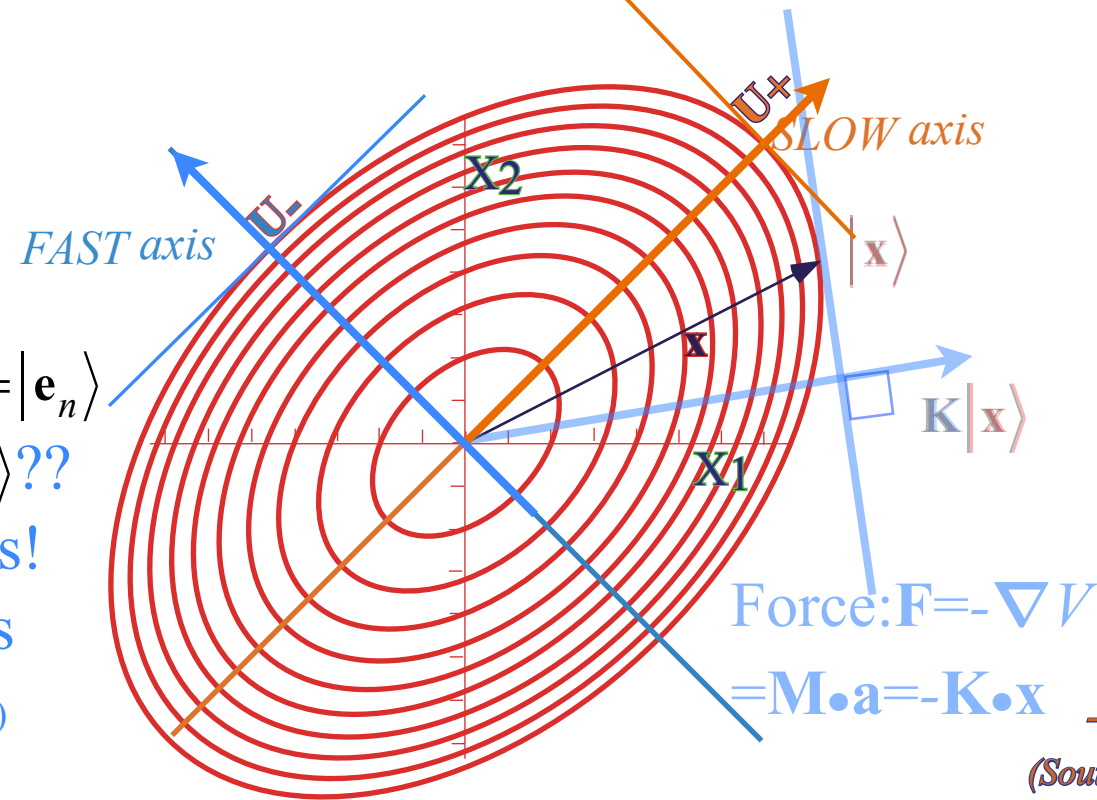


Fig. 3.3.4 Plot of potential function $V(x_1, x_2)$ showing elliptical $V(x_1, x_2) = \text{const.}$ level curves.

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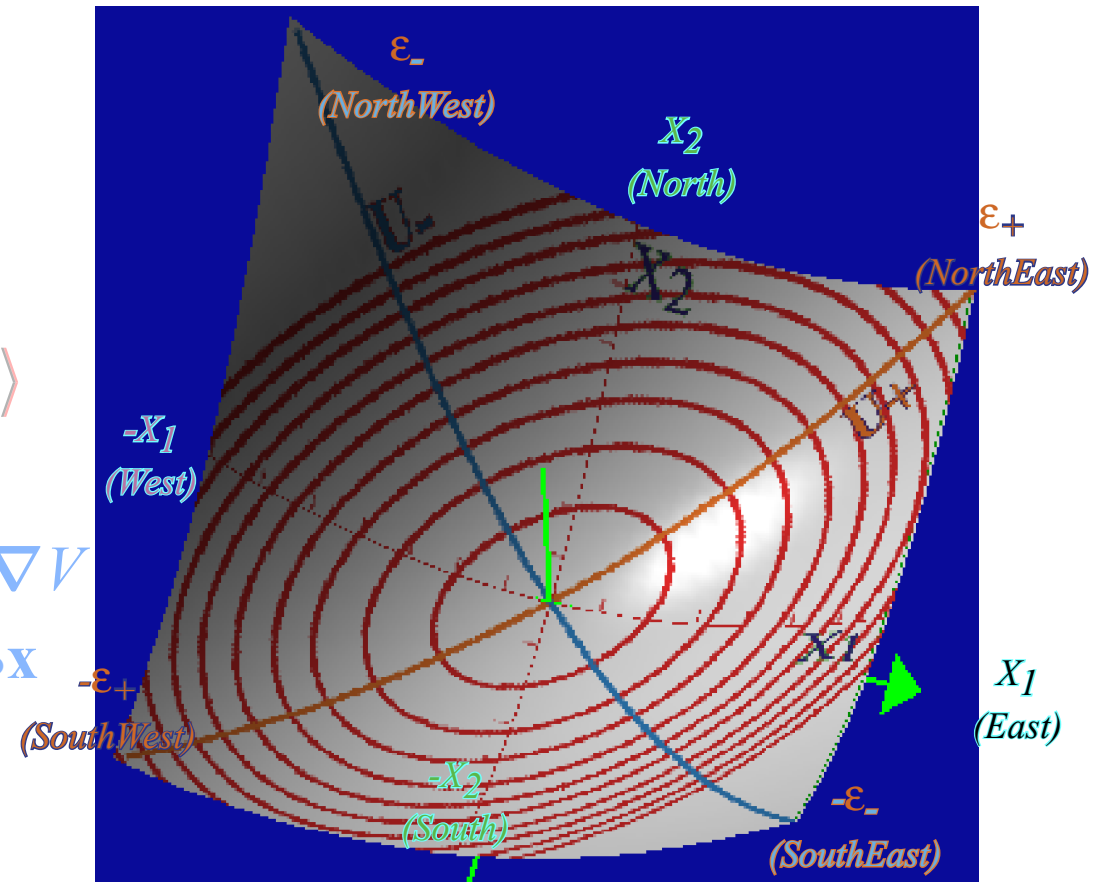
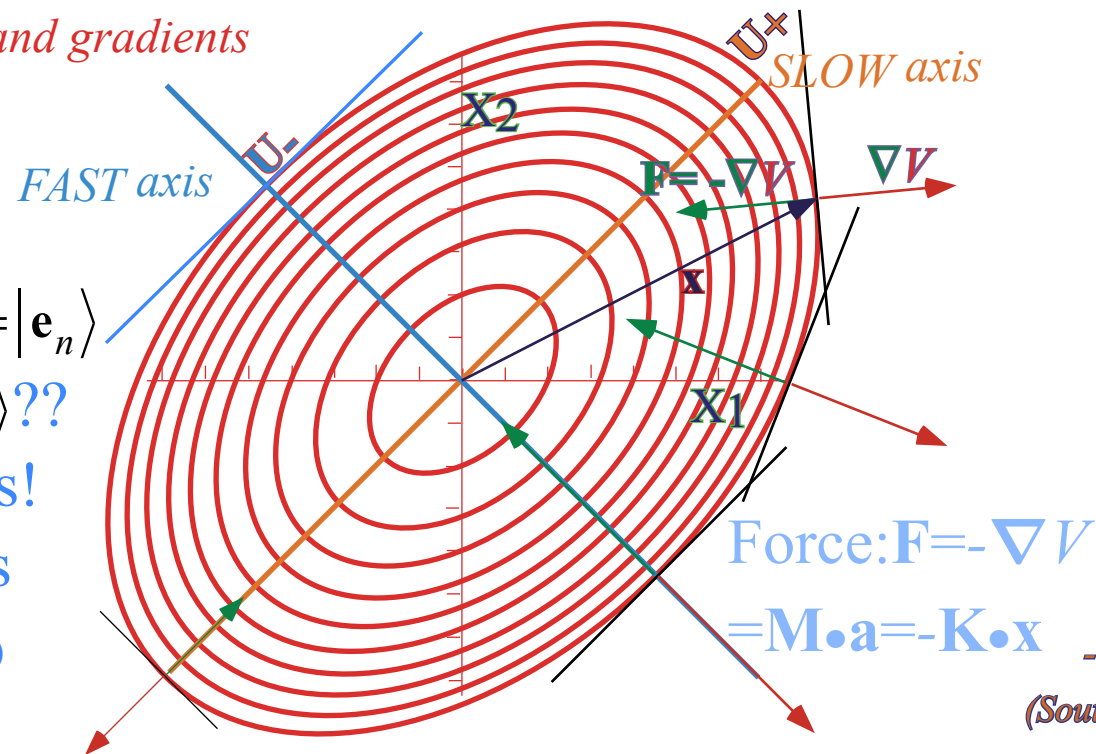


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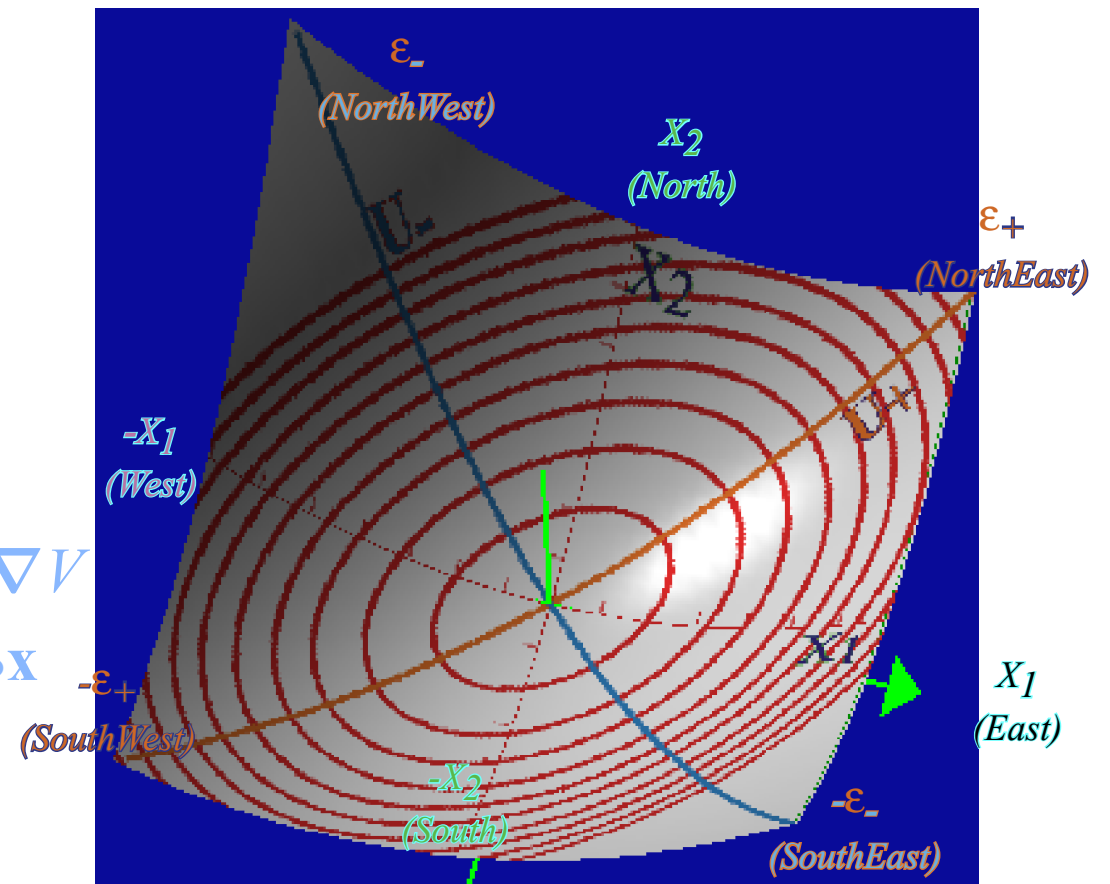
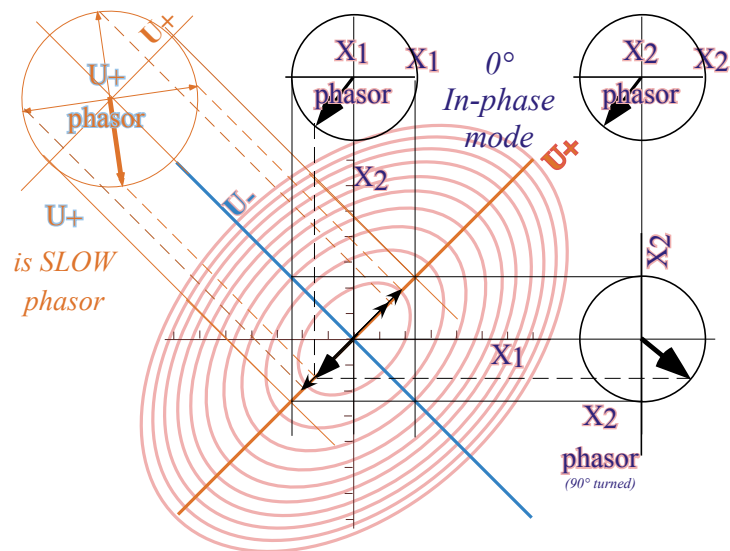
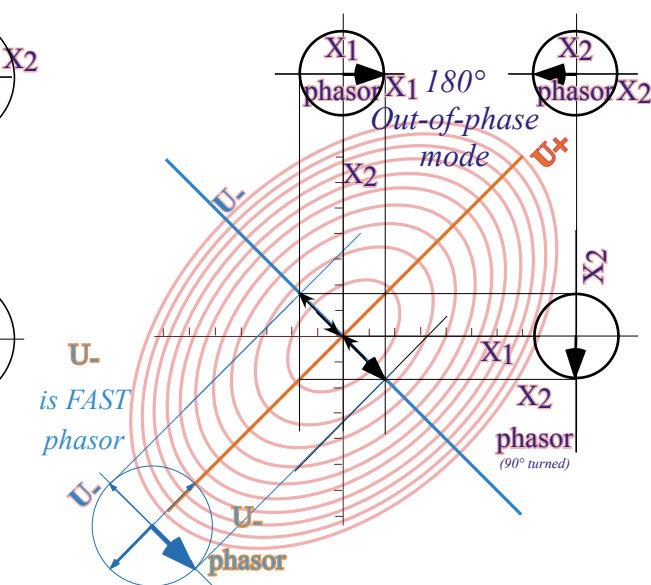


Fig. 3.3.4 Plot of potential function $V(x_1, x_2)$ showing elliptical $V(x_1, x_2) = \text{const.}$ level curves.

(b) Symmetric $U+$ Coordinate SLOW Mode



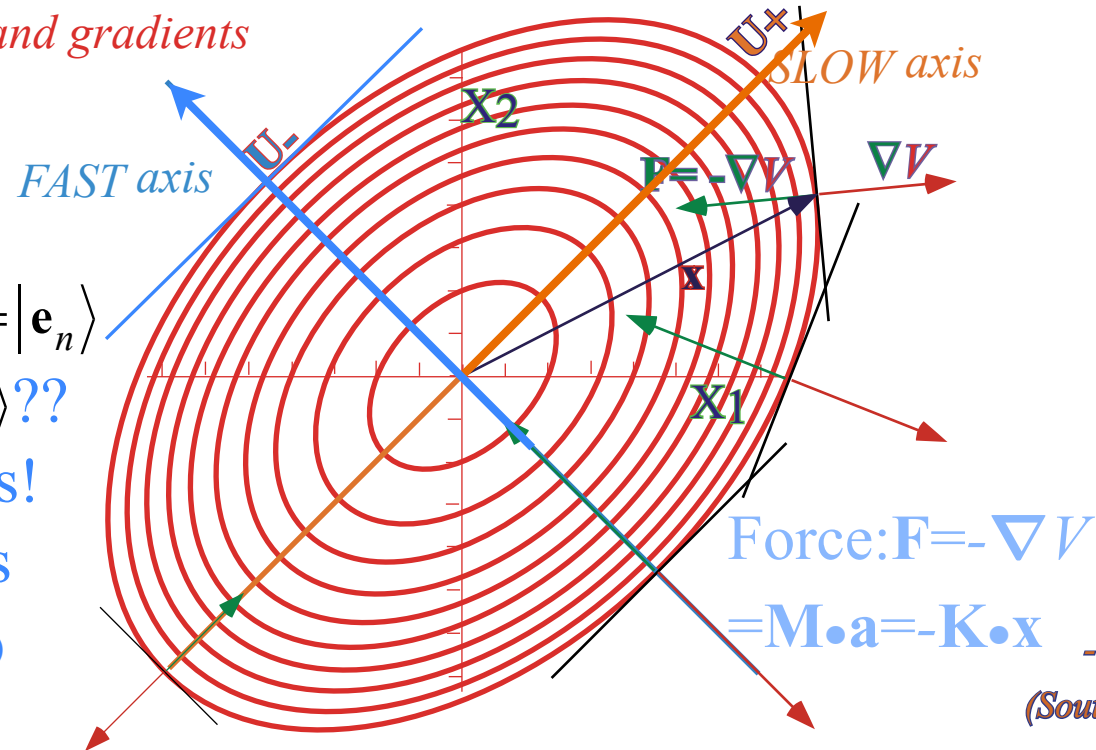
(c) Anti-symmetric $U-$ Coordinate FAST Mode



2D HO potential energy $V(x_1, x_2)$ quadratic form defines layers of elliptical V -contours (Here: $k_1 = k = k_2$)

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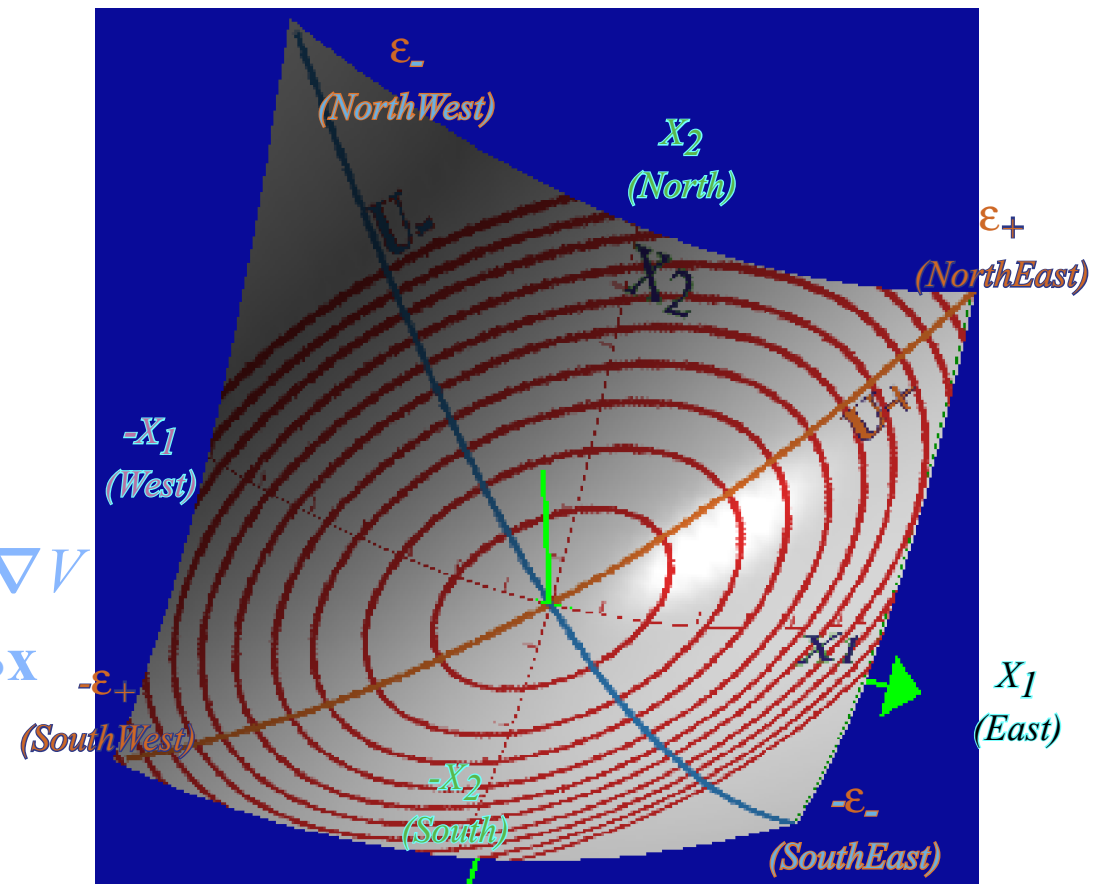


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(b) Symmetric $U+$ Coordinate SLOW Mode

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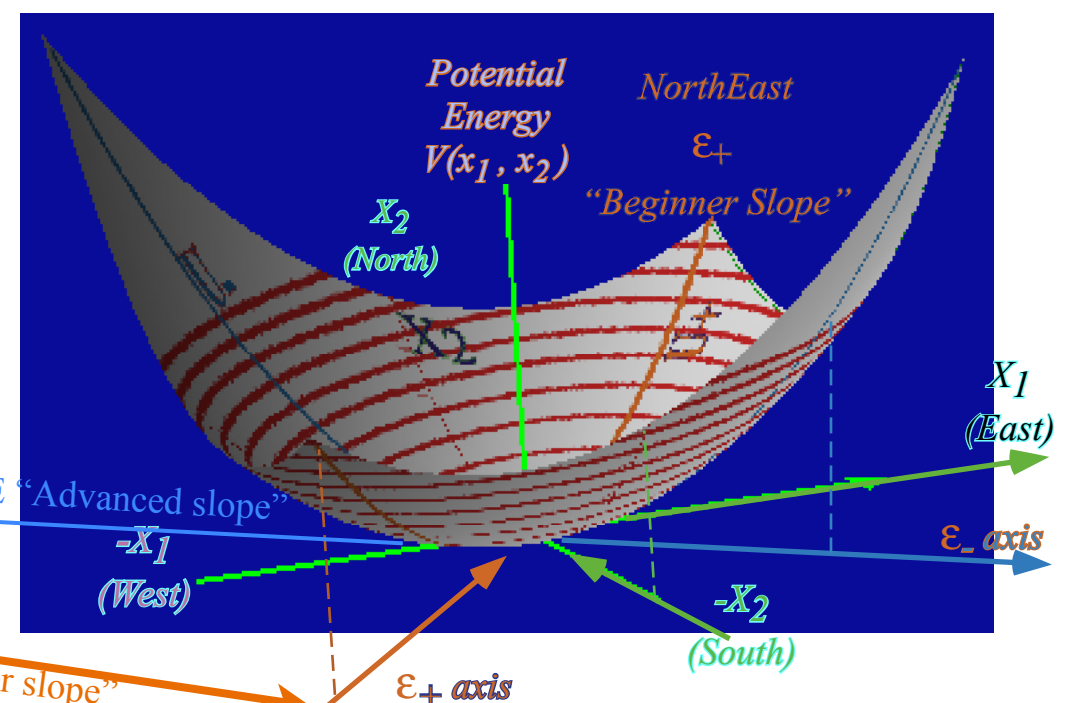
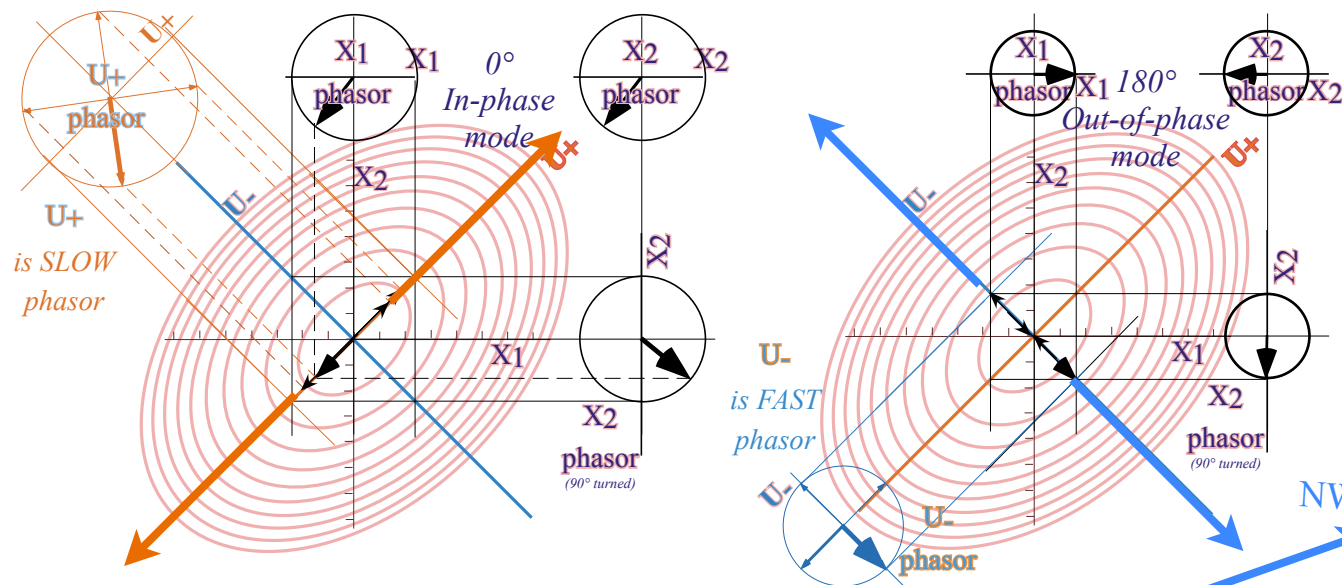


Fig. 3.3.5 Topography lines of potential function $V(x_1, x_2)$ and orthogonal ϵ_+ and ϵ_- normal mode slopes

With Bilateral symmetry ($k_1 = k = k_2$) the extremal axes lie at $\pm 45^\circ$

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➔ *Matrix-algebraic eigensolutions with example $M = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix}$* ←

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Matrix-algebraic method for finding eigenvector and eigenvalues *With example matrix* $\mathbf{M} = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix}$

An *eigenvector* $|\epsilon_k\rangle$ of \mathbf{M} is in a direction that is left unchanged by \mathbf{M} .

$$\mathbf{M}|\epsilon_k\rangle = \epsilon_k|\epsilon_k\rangle, \text{ or: } (\mathbf{M} - \epsilon_k\mathbf{1})|\epsilon_k\rangle = \mathbf{0}$$

$$\mathbf{M}|\epsilon\rangle = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \epsilon \begin{pmatrix} x \\ y \end{pmatrix} \text{ or: } \begin{pmatrix} 4-\epsilon & 1 \\ 3 & 2-\epsilon \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

ϵ_k is *eigenvalue* associated with eigenvector $|\epsilon_k\rangle$ direction.

A change of basis to $\{|\epsilon_1\rangle, |\epsilon_2\rangle, \dots, |\epsilon_n\rangle\}$ called *diagonalization* gives

$$\begin{pmatrix} \langle \epsilon_1 | \mathbf{M} | \epsilon_1 \rangle & \langle \epsilon_1 | \mathbf{M} | \epsilon_2 \rangle & \cdots & \langle \epsilon_1 | \mathbf{M} | \epsilon_n \rangle \\ \langle \epsilon_2 | \mathbf{M} | \epsilon_1 \rangle & \langle \epsilon_2 | \mathbf{M} | \epsilon_2 \rangle & \cdots & \langle \epsilon_2 | \mathbf{M} | \epsilon_n \rangle \\ \vdots & \vdots & \ddots & \vdots \\ \langle \epsilon_n | \mathbf{M} | \epsilon_1 \rangle & \langle \epsilon_n | \mathbf{M} | \epsilon_2 \rangle & \cdots & \langle \epsilon_n | \mathbf{M} | \epsilon_n \rangle \end{pmatrix} = \begin{pmatrix} \epsilon_1 & 0 & \cdots & 0 \\ 0 & \epsilon_2 & \cdots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \cdots & \epsilon_n \end{pmatrix}$$

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$$\det|\mathbf{M} - \epsilon \mathbf{1}| = 0 = (-1)^n (\epsilon^n + a_1 \epsilon^{n-1} + a_2 \epsilon^{n-2} + \dots + a_{n-1} \epsilon + a_n)$$

where:

$$a_1 = -\text{Trace} \mathbf{M}, \dots, a_k = (-1)^k \sum \text{diagonal k-by-k minors of } \mathbf{M}, \dots, a_n = (-1)^n \det|\mathbf{M}|$$

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Only possible non-zero $\{x, y\}$ if denominator is zero, too!

$$0 = \det|\mathbf{M} - \epsilon \cdot \mathbf{1}| = \det \left[\begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix} - \epsilon \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \right] = \det \begin{pmatrix} 4-\epsilon & 1 \\ 3 & 2-\epsilon \end{pmatrix}$$

$$0 = (4-\epsilon)(2-\epsilon) - 1 \cdot 3 = 8 - 6\epsilon + \epsilon^2 - 3 = \epsilon^2 - 6\epsilon + 5$$

$$0 = \epsilon^2 - \text{Trace}(\mathbf{M})\epsilon + \det(\mathbf{M})$$

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Only possible non-zero $\{x, y\}$ if denominator is zero, too!

$$0 = \det|\mathbf{M} - \epsilon \cdot \mathbf{1}| = \det \left[\begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix} - \epsilon \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \right] = \det \begin{pmatrix} 4 - \epsilon & 1 \\ 3 & 2 - \epsilon \end{pmatrix}$$

$$0 = (4 - \epsilon)(2 - \epsilon) - 1 \cdot 3 = 8 - 6\epsilon + \epsilon^2 - 3 = \epsilon^2 - 6\epsilon + 5$$

$$0 = \epsilon^2 - \text{Trace}(\mathbf{M})\epsilon + \det(\mathbf{M}) = \epsilon^2 - 6\epsilon + 5$$

$$0 = (\epsilon - 1)(\epsilon - 5) \text{ so let: } \epsilon_1 = 1 \text{ and: } \epsilon_2 = 5$$

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Secular equation has n -factors, one for each eigenvalue.

$$\det|\mathbf{M} - \epsilon \mathbf{1}| = 0 = (-1)^n (\epsilon - \epsilon_1)(\epsilon - \epsilon_2) \dots (\epsilon - \epsilon_n)$$

Each ϵ replaced by \mathbf{M} and each ϵ_k by $\epsilon_k \mathbf{1}$ gives *Hamilton-Cayley* matrix equation.

$$\mathbf{0} = (\mathbf{M} - \epsilon_1 \mathbf{1})(\mathbf{M} - \epsilon_2 \mathbf{1}) \dots (\mathbf{M} - \epsilon_n \mathbf{1})$$

Obviously true if \mathbf{M} has diagonal form. (But, that's circular logic. Faith needed!)

$$\mathbf{M}|\epsilon\rangle = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \epsilon \begin{pmatrix} x \\ y \end{pmatrix} \text{ or: } \begin{pmatrix} 4 - \epsilon & 1 \\ 3 & 2 - \epsilon \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

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$$0 = (\epsilon - 1)(\epsilon - 5) \text{ so let: } \epsilon_1 = 1 \text{ and: } \epsilon_2 = 5$$

$$0 = \mathbf{M}^2 - 6\mathbf{M} + 5\mathbf{1} = (\mathbf{M} - 1 \cdot \mathbf{1})(\mathbf{M} - 5 \cdot \mathbf{1})$$

$$\begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix}^2 - 6 \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix} + 5 \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

2D harmonic oscillator equations

Lagrangian and matrix forms and Reciprocity symmetry

2D harmonic oscillator equation eigensolutions

Geometric method

Matrix-algebraic eigensolutions with example $M = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix}$

Secular equation

➔ *Hamilton-Cayley equation and projectors* **←**

Idempotent projectors (how eigenvalues \Rightarrow eigenvectors)

Operator orthonormality and Completeness (Idempotent means: $\mathbf{P} \cdot \mathbf{P} = \mathbf{P}$)

Spectral Decompositions

Functional spectral decomposition

Orthonormality vs. Completeness vis-a`-vis Operator vs. State

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Mixed mode beat dynamics and fixed $\pi/2$ phase

2D-HO eigensolution example with asymmetric (A-Type) symmetry

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Matrix-algebraic method for finding eigenvector and eigenvalues With example matrix $\mathbf{M} = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix}$

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A change of basis to $\{|\epsilon_1\rangle, |\epsilon_2\rangle, \dots, |\epsilon_n\rangle\}$ called *diagonalization* gives

$$\begin{pmatrix} \langle \epsilon_1 | \mathbf{M} | \epsilon_1 \rangle & \langle \epsilon_1 | \mathbf{M} | \epsilon_2 \rangle & \dots & \langle \epsilon_1 | \mathbf{M} | \epsilon_n \rangle \\ \langle \epsilon_2 | \mathbf{M} | \epsilon_1 \rangle & \langle \epsilon_2 | \mathbf{M} | \epsilon_2 \rangle & \dots & \langle \epsilon_2 | \mathbf{M} | \epsilon_n \rangle \\ \vdots & \vdots & \ddots & \vdots \\ \langle \epsilon_n | \mathbf{M} | \epsilon_1 \rangle & \langle \epsilon_n | \mathbf{M} | \epsilon_2 \rangle & \dots & \langle \epsilon_n | \mathbf{M} | \epsilon_n \rangle \end{pmatrix} = \begin{pmatrix} \epsilon_1 & 0 & \dots & 0 \\ 0 & \epsilon_2 & \dots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ 0 & 0 & \dots & \epsilon_n \end{pmatrix}$$

First step in finding eigenvalues: Solve *secular equation*

$$\det|\mathbf{M} - \epsilon \mathbf{1}| = 0 = (-1)^n (\epsilon^n + a_1 \epsilon^{n-1} + a_2 \epsilon^{n-2} + \dots + a_{n-1} \epsilon + a_n)$$

where:

$$a_1 = -\text{Trace} \mathbf{M}, \dots, a_k = (-1)^k \sum \text{diagonal k-by-k minors of } \mathbf{M}, \dots, a_n = (-1)^n \det|\mathbf{M}|$$

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$$\det|\mathbf{M} - \epsilon \mathbf{1}| = 0 = (-1)^n (\epsilon - \epsilon_1)(\epsilon - \epsilon_2) \dots (\epsilon - \epsilon_n)$$

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Obviously true if \mathbf{M} has diagonal form. (But, that's circular logic. Faith needed!)

Replace j^{th} HC-factor by $(\mathbf{1})$ to make *projection operators*

$$\begin{aligned} \mathbf{p}_1 &= (\mathbf{1})(\mathbf{M} - \epsilon_2 \mathbf{1}) \dots (\mathbf{M} - \epsilon_n \mathbf{1}) \\ \mathbf{p}_2 &= (\mathbf{M} - \epsilon_1 \mathbf{1})(\mathbf{1}) \dots (\mathbf{M} - \epsilon_n \mathbf{1}) \\ &\vdots \\ \mathbf{p}_n &= (\mathbf{M} - \epsilon_1 \mathbf{1})(\mathbf{M} - \epsilon_2 \mathbf{1}) \dots (\mathbf{1}) \end{aligned}$$

$$\mathbf{M}|\epsilon\rangle = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \epsilon \begin{pmatrix} x \\ y \end{pmatrix} \text{ or: } \begin{pmatrix} 4 - \epsilon & 1 \\ 3 & 2 - \epsilon \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

Trying to solve by Kramer's inversion:

$$x = \frac{\det \begin{pmatrix} 0 & 1 \\ 0 & 2 - \epsilon \end{pmatrix}}{\det \begin{pmatrix} 4 - \epsilon & 1 \\ 3 & 2 - \epsilon \end{pmatrix}} \quad \text{and} \quad y = \frac{\det \begin{pmatrix} 4 - \epsilon & 0 \\ 3 & 0 \end{pmatrix}}{\det \begin{pmatrix} 4 - \epsilon & 1 \\ 3 & 2 - \epsilon \end{pmatrix}}$$

Only possible non-zero $\{x, y\}$ if denominator is zero, too!

$$0 = \det|\mathbf{M} - \epsilon \cdot \mathbf{1}| = \det \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix} - \epsilon \det \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \det \begin{pmatrix} 4 - \epsilon & 1 \\ 3 & 2 - \epsilon \end{pmatrix}$$

$$0 = (4 - \epsilon)(2 - \epsilon) - 1 \cdot 3 = 8 - 6\epsilon + \epsilon^2 - 3 = \epsilon^2 - 6\epsilon + 5$$

$$0 = \epsilon^2 - \text{Trace}(\mathbf{M})\epsilon + \det(\mathbf{M}) = \epsilon^2 - 6\epsilon + 5$$

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Matrix-algebraic method for finding eigenvector and eigenvalues With example matrix $\mathbf{M} = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix}$

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First step in finding eigenvalues: Solve *secular equation*

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Each ϵ replaced by \mathbf{M} and each ϵ_k by $\epsilon_k\mathbf{1}$ gives *Hamilton-Cayley* matrix equation.

$$\mathbf{0} = (\mathbf{M} - \epsilon_1\mathbf{1})(\mathbf{M} - \epsilon_2\mathbf{1}) \dots (\mathbf{M} - \epsilon_n\mathbf{1})$$

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$$\mathbf{p}_2 = (\mathbf{M} - \epsilon_1\mathbf{1})(\mathbf{1}) \dots (\mathbf{M} - \epsilon_n\mathbf{1})$$

\vdots

$$\mathbf{p}_n = (\mathbf{M} - \epsilon_1\mathbf{1})(\mathbf{M} - \epsilon_2\mathbf{1}) \dots (\mathbf{1})$$

(Assume distinct e-values here: *Non-degeneracy clause*)
 $\epsilon_j \neq \epsilon_k \neq \dots$

Each \mathbf{p}_k contains *eigen-bra-kets* since: $(\mathbf{M} - \epsilon_k\mathbf{1})\mathbf{p}_k = \mathbf{0}$ or: $\mathbf{M}\mathbf{p}_k = \epsilon_k\mathbf{p}_k = \mathbf{p}_k\mathbf{M}$.

$$\mathbf{M}|\epsilon\rangle = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \epsilon \begin{pmatrix} x \\ y \end{pmatrix} \text{ or: } \begin{pmatrix} 4 - \epsilon & 1 \\ 3 & 2 - \epsilon \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

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$$\mathbf{M}\mathbf{p}_1 = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix} \cdot \begin{pmatrix} -1 & 1 \\ 3 & -3 \end{pmatrix} = 1 \cdot \begin{pmatrix} -1 & 1 \\ 3 & -3 \end{pmatrix} = 1 \cdot \mathbf{p}_1$$

$$\mathbf{M}\mathbf{p}_2 = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix} \cdot \begin{pmatrix} 3 & 1 \\ 3 & 1 \end{pmatrix} = 5 \cdot \begin{pmatrix} 3 & 1 \\ 3 & 1 \end{pmatrix} = 5 \cdot \mathbf{p}_2$$

Matrix-algebraic method for finding eigenvector and eigenvalues With example matrix $\mathbf{M} = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix}$

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$$\begin{aligned} \mathbf{p}_1 &= (\mathbf{1})(\mathbf{M} - \epsilon_2\mathbf{1}) \dots (\mathbf{M} - \epsilon_n\mathbf{1}) \\ \mathbf{p}_2 &= (\mathbf{M} - \epsilon_1\mathbf{1})(\mathbf{1}) \dots (\mathbf{M} - \epsilon_n\mathbf{1}) \\ &\vdots \\ \mathbf{p}_n &= (\mathbf{M} - \epsilon_1\mathbf{1})(\mathbf{M} - \epsilon_2\mathbf{1}) \dots (\mathbf{1}) \end{aligned} \quad \begin{array}{l} \text{(Assume distinct e-values here: } \\ \text{Non-degeneracy clause)} \\ \epsilon_j \neq \epsilon_k \neq \dots \end{array}$$

Each \mathbf{p}_k contains *eigen-bra-kets* since: $(\mathbf{M} - \epsilon_k\mathbf{1})\mathbf{p}_k = \mathbf{0}$ or: $\mathbf{M}\mathbf{p}_k = \epsilon_k\mathbf{p}_k = \mathbf{p}_k\mathbf{M}$.

Notice \mathbf{p}_k commutes with \mathbf{M} ...

since $\mathbf{M}^1, \mathbf{M}^2, \dots$ commute with \mathbf{M} .

$$\mathbf{M}|\epsilon\rangle = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \epsilon \begin{pmatrix} x \\ y \end{pmatrix} \text{ or: } \begin{pmatrix} 4-\epsilon & 1 \\ 3 & 2-\epsilon \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

Trying to solve by Kramer's inversion:

$$x = \frac{\det \begin{pmatrix} 0 & 1 \\ 0 & 2-\epsilon \end{pmatrix}}{\det \begin{pmatrix} 4-\epsilon & 1 \\ 3 & 2-\epsilon \end{pmatrix}} \quad \text{and} \quad y = \frac{\det \begin{pmatrix} 4-\epsilon & 0 \\ 3 & 0 \end{pmatrix}}{\det \begin{pmatrix} 4-\epsilon & 1 \\ 3 & 2-\epsilon \end{pmatrix}}$$

Only possible non-zero $\{x, y\}$ if denominator is zero, too!

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$$\mathbf{p}_2 = (\mathbf{M} - 1\mathbf{1})(\mathbf{1}) = \begin{pmatrix} 4-1 & 1 \\ 3 & 2-1 \end{pmatrix} = \begin{pmatrix} 3 & 1 \\ 3 & 1 \end{pmatrix}$$

$$\mathbf{M}\mathbf{p}_1 = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix} \cdot \begin{pmatrix} -1 & 1 \\ 3 & -3 \end{pmatrix} = 1 \cdot \begin{pmatrix} -1 & 1 \\ 3 & -3 \end{pmatrix} = 1 \cdot \mathbf{p}_1$$

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➔ *Idempotent projectors* (↔ *low eigenvalues ⇒ eigenvectors*)

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Initial state projection, mixed mode beat dynamics with variable phase

Matrix-algebraic method for finding eigenvector and eigenvalues

$$\mathbf{p}_j \mathbf{p}_k = \mathbf{p}_j \prod_{m \neq k} (\mathbf{M} - \varepsilon_m \mathbf{1}) = \prod_{m \neq k} (\mathbf{p}_j \mathbf{M} - \varepsilon_m \mathbf{p}_j \mathbf{1}) \quad \mathbf{M} \mathbf{p}_k = \varepsilon_k \mathbf{p}_k = \mathbf{p}_k \mathbf{M}$$

Multiplication properties of \mathbf{p}_j :

$$\mathbf{p}_j \mathbf{p}_k = \prod_{m \neq k} (\varepsilon_j \mathbf{p}_j - \varepsilon_m \mathbf{p}_j) = \mathbf{p}_j \prod_{m \neq k} (\varepsilon_j - \varepsilon_m) = \begin{cases} \mathbf{0} & \text{if } j \neq k \\ \mathbf{p}_k \prod_{m \neq k} (\varepsilon_k - \varepsilon_m) & \text{if } j = k \end{cases}$$

With example matrix

$$\mathbf{M} = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix}$$

$$\mathbf{p}_1 = (\mathbf{M} - 5 \cdot \mathbf{1}) = \begin{pmatrix} -1 & 1 \\ 3 & -3 \end{pmatrix}$$

$$\mathbf{p}_2 = (\mathbf{M} - 1 \cdot \mathbf{1}) = \begin{pmatrix} 3 & 1 \\ 3 & 1 \end{pmatrix}$$

$$\mathbf{p}_1 \mathbf{p}_2 = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

Matrix-algebraic method for finding eigenvector and eigenvalues

With example matrix

$$\mathbf{M} = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix}$$

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$$\mathbf{p}_2 = (\mathbf{M} - 1 \cdot \mathbf{1}) = \begin{pmatrix} 3 & 1 \\ 3 & 1 \end{pmatrix}$$

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Last step:

make *Idempotent Projectors*: $\mathbf{P}_k = \frac{\mathbf{p}_k}{\prod_{m \neq k} (\varepsilon_k - \varepsilon_m)}$
 (Idempotent means: $\mathbf{P} \cdot \mathbf{P} = \mathbf{P}$)

$$\mathbf{P}_k = \frac{\mathbf{p}_k}{\prod_{m \neq k} (\varepsilon_k - \varepsilon_m)} = \frac{\prod_{m \neq k} (\mathbf{M} - \varepsilon_m \mathbf{1})}{\prod_{m \neq k} (\varepsilon_k - \varepsilon_m)}$$

$$\mathbf{P}_1 = \frac{(\mathbf{M} - 5 \cdot \mathbf{1})}{(1 - 5)} = \frac{1}{4} \begin{pmatrix} 1 & -1 \\ -3 & 3 \end{pmatrix}$$

$$\mathbf{P}_2 = \frac{(\mathbf{M} - 1 \cdot \mathbf{1})}{(5 - 1)} = \frac{1}{4} \begin{pmatrix} 3 & 1 \\ 3 & 1 \end{pmatrix}$$

Matrix-algebraic method for finding eigenvector and eigenvalues

With example matrix

$$\mathbf{M} = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix}$$

$$\mathbf{p}_j \mathbf{p}_k = \mathbf{p}_j \prod_{m \neq k} (\mathbf{M} - \varepsilon_m \mathbf{1}) = \prod_{m \neq k} (\mathbf{p}_j \mathbf{M} - \varepsilon_m \mathbf{p}_j \mathbf{1}) \quad \mathbf{M} \mathbf{p}_k = \varepsilon_k \mathbf{p}_k = \mathbf{p}_k \mathbf{M}$$

$$\mathbf{p}_1 = (\mathbf{M} - 5 \cdot \mathbf{1}) = \begin{pmatrix} -1 & 1 \\ 3 & -3 \end{pmatrix}$$

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➔ *Idempotent projectors (how eigenvalues \Rightarrow eigenvectors)* ←

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$$\mathbf{p}_j \mathbf{p}_k = \mathbf{p}_j \prod_{m \neq k} (\mathbf{M} - \varepsilon_m \mathbf{1}) = \prod_{m \neq k} (\mathbf{p}_j \mathbf{M} - \varepsilon_m \mathbf{p}_j \mathbf{1}) \quad \mathbf{M} \mathbf{p}_k = \varepsilon_k \mathbf{p}_k = \mathbf{p}_k \mathbf{M}$$

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With example matrix $\mathbf{M} = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix}$

$$\mathbf{p}_1 = (\mathbf{M} - 5 \cdot \mathbf{1}) = \begin{pmatrix} -1 & 1 \\ 3 & -3 \end{pmatrix} \quad \mathbf{p}_1 \mathbf{p}_2 = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

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Matrix-algebraic method for finding eigenvector and eigenvalues

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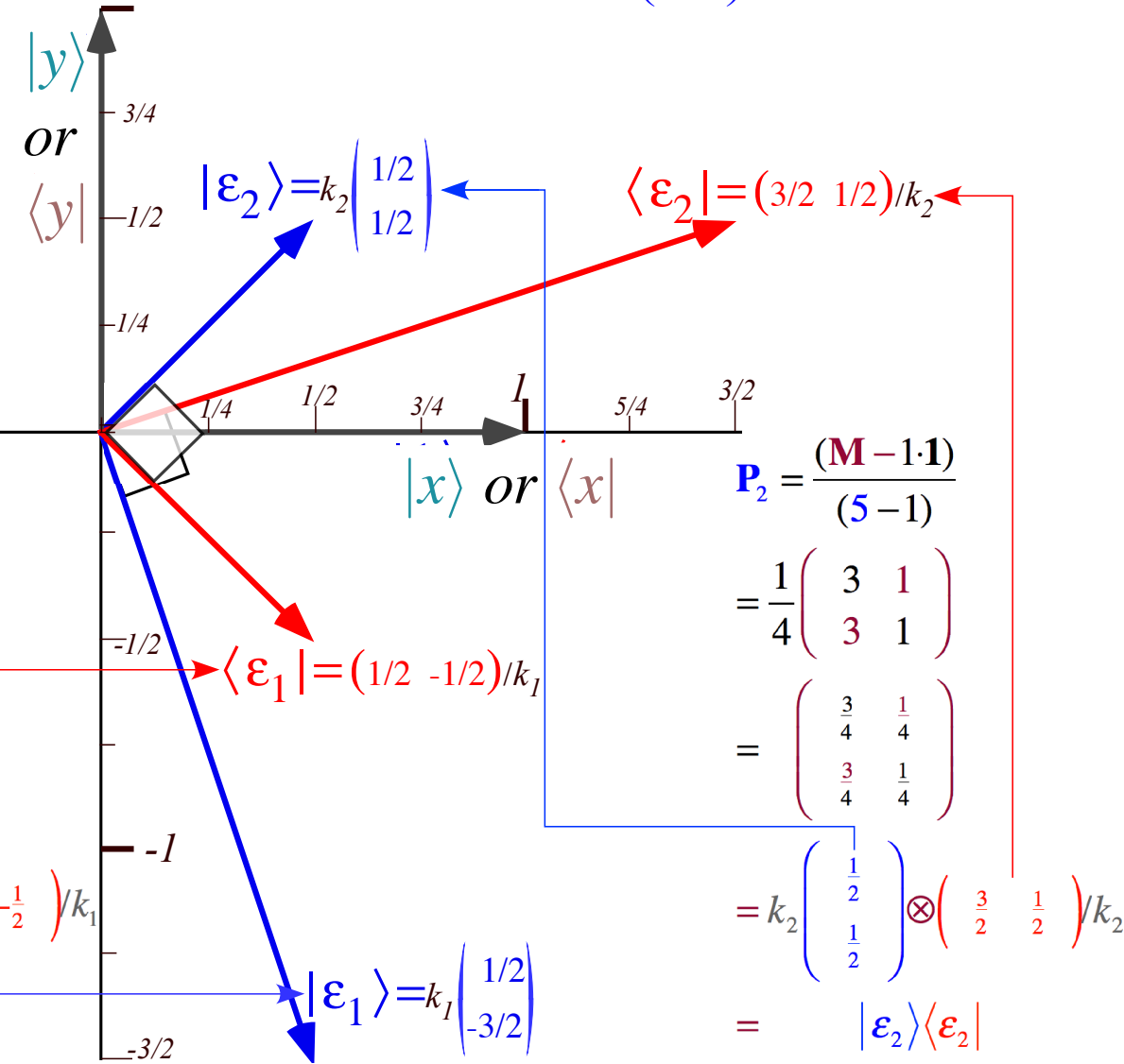
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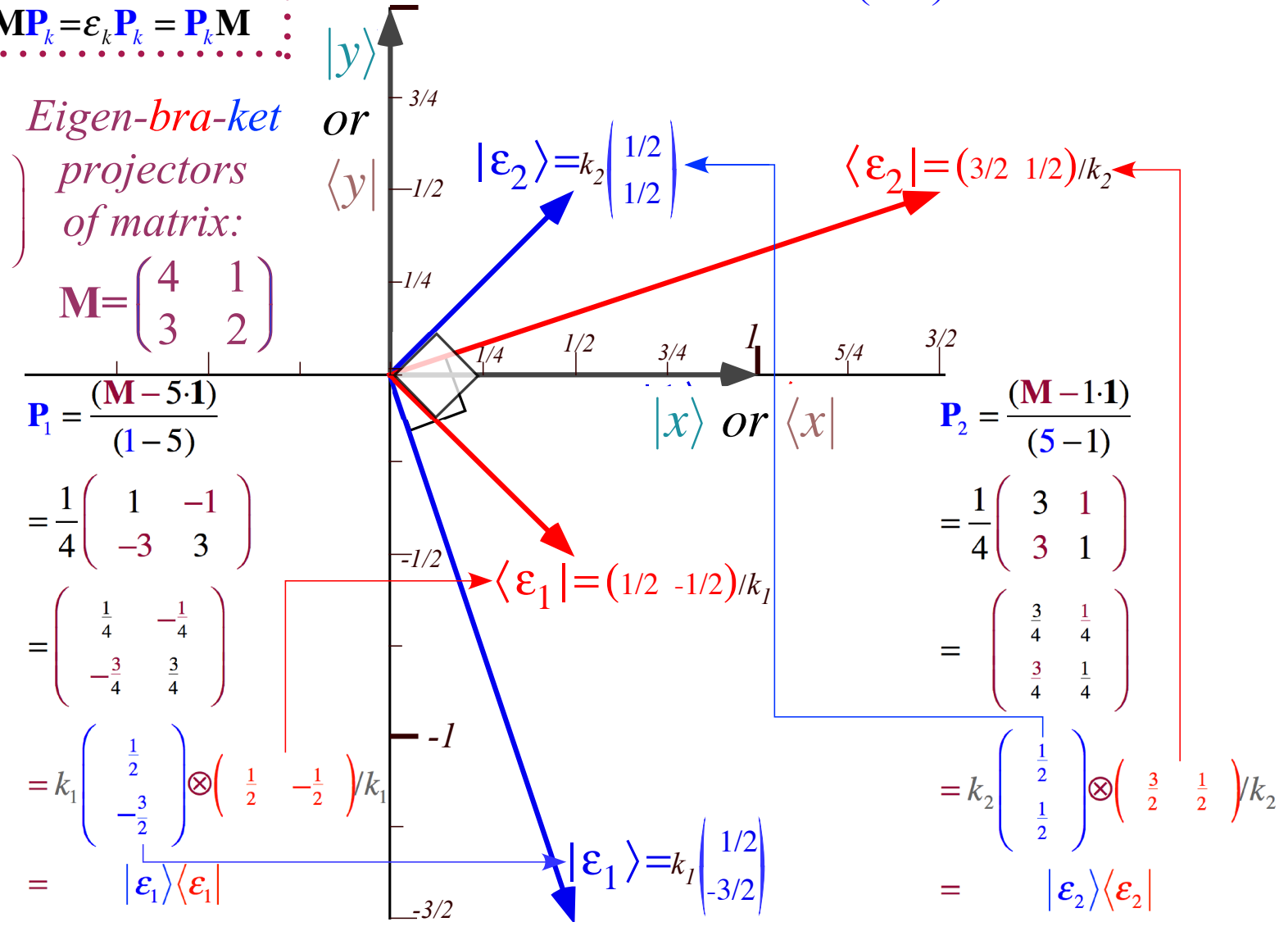
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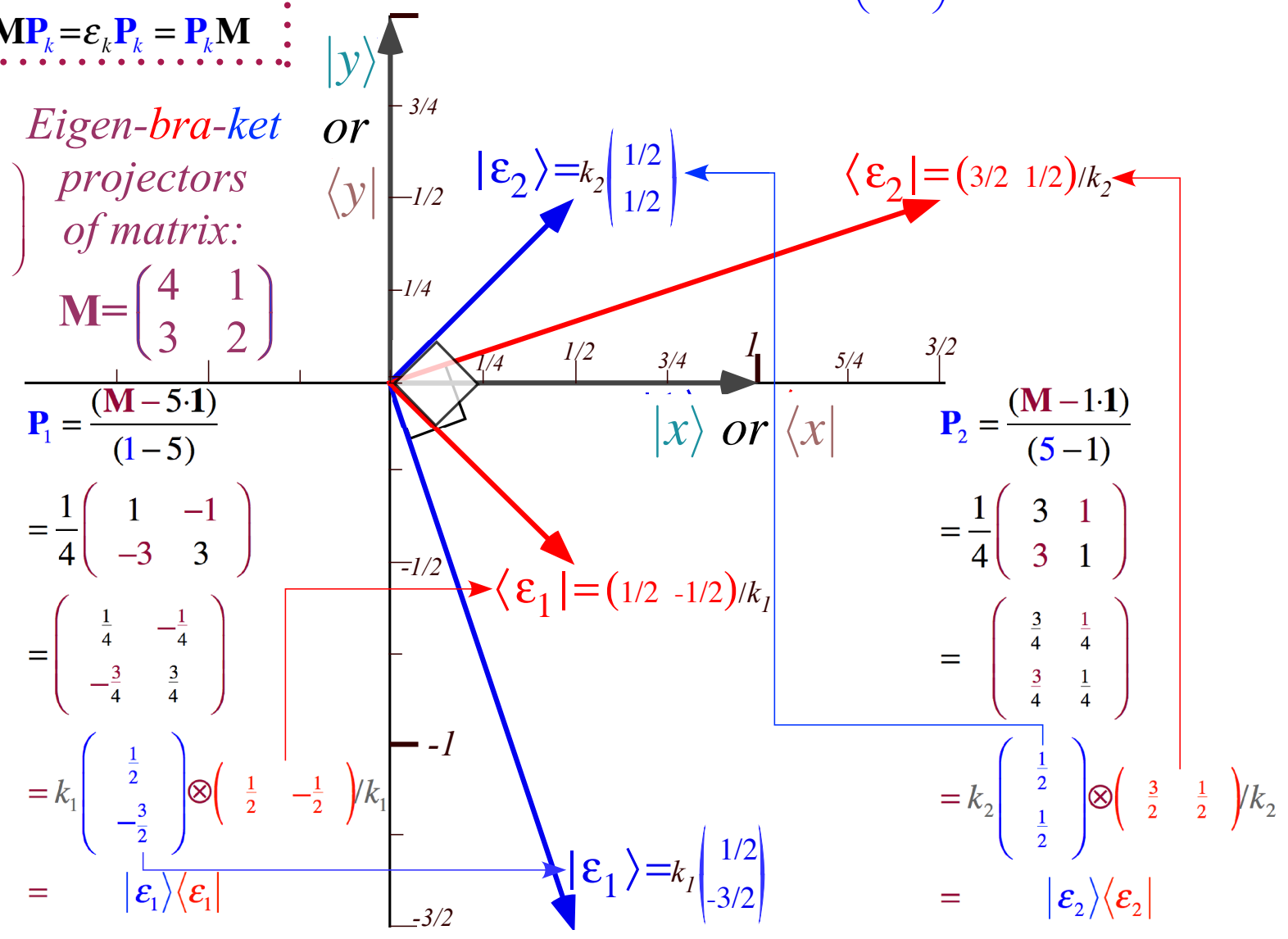
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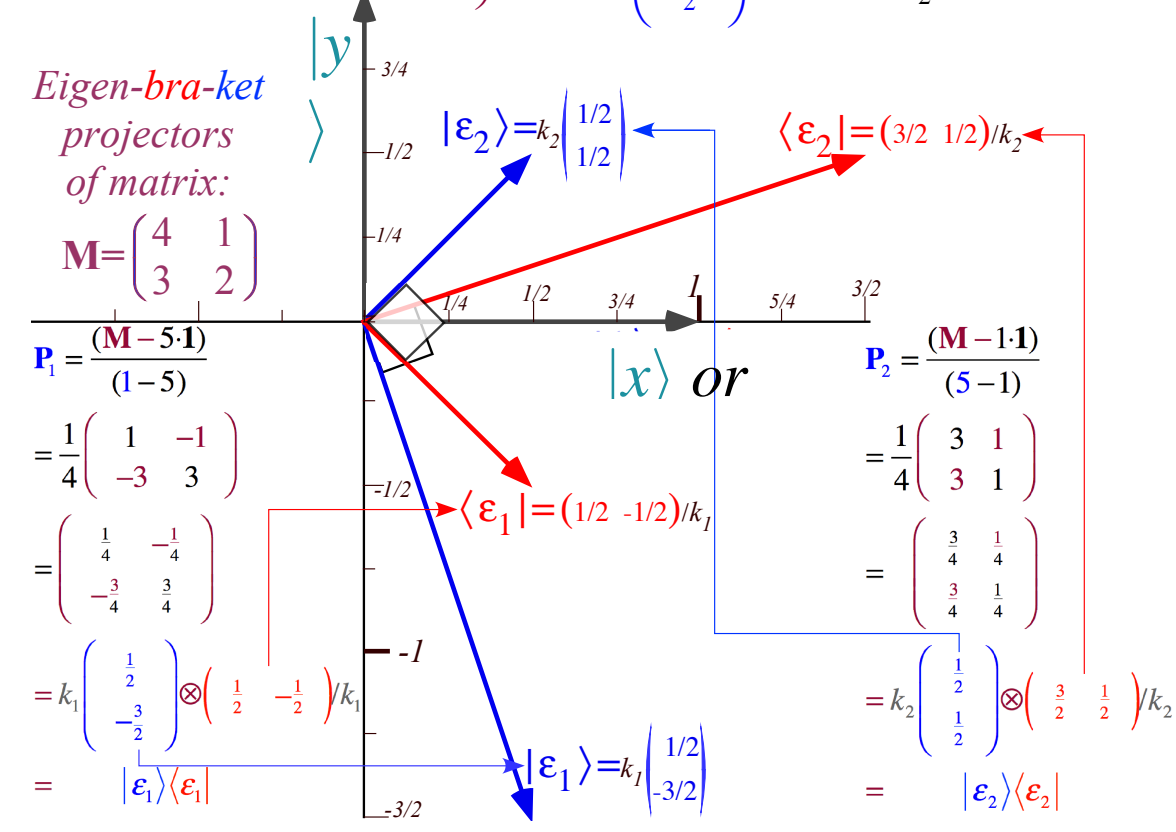
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Eigen-operators $M P_k = \epsilon_k P_k$ then give *Spectral Decomposition* of operator M

$$M = M P_1 + M P_2 + \dots + M P_n = \epsilon_1 P_1 + \epsilon_2 P_2 + \dots + \epsilon_n P_n$$



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$$\mathbf{P}_1 + \mathbf{P}_2 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = |\epsilon_1 \rangle \langle \epsilon_1| + |\epsilon_2 \rangle \langle \epsilon_2|$$

$$\mathbf{M} = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix} = 1\mathbf{P}_1 + 5\mathbf{P}_2 = 1|\epsilon_1 \rangle \langle \epsilon_1| + 5|\epsilon_2 \rangle \langle \epsilon_2| = 1 \begin{pmatrix} \frac{1}{4} & -\frac{1}{4} \\ -\frac{3}{4} & \frac{3}{4} \end{pmatrix} + 5 \begin{pmatrix} \frac{3}{4} & \frac{1}{4} \\ \frac{3}{4} & \frac{1}{4} \end{pmatrix}$$

With example matrix $\mathbf{M} = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix}$

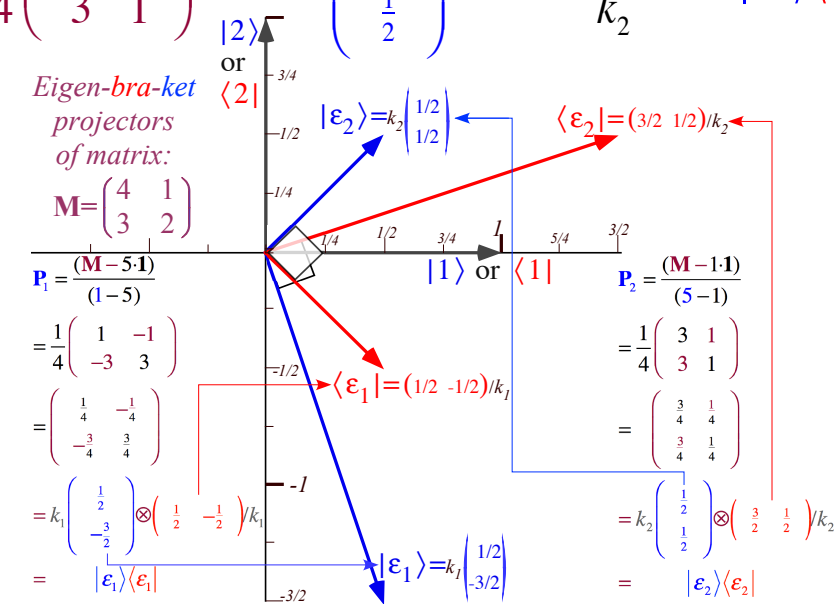
$$\mathbf{p}_1 = (\mathbf{M} - 5 \cdot \mathbf{1}) = \begin{pmatrix} -1 & 1 \\ 3 & -3 \end{pmatrix} \quad \mathbf{p}_1 \mathbf{p}_2 = \begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix}$$

$$\mathbf{p}_2 = (\mathbf{M} - 1 \cdot \mathbf{1}) = \begin{pmatrix} 3 & 1 \\ 3 & 1 \end{pmatrix}$$

Factoring bra-kets into "Ket-Bras:

$$\mathbf{P}_1 = \frac{(\mathbf{M} - 5 \cdot \mathbf{1})}{(1 - 5)} = \frac{1}{4} \begin{pmatrix} 1 & -1 \\ -3 & 3 \end{pmatrix} = k_1 \begin{pmatrix} \frac{1}{2} \\ -\frac{3}{2} \end{pmatrix} \otimes \frac{\begin{pmatrix} \frac{1}{2} & -\frac{1}{2} \end{pmatrix}}{k_1} = |\epsilon_1 \rangle \langle \epsilon_1|$$

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Eigen-operators $\mathbf{M} \mathbf{P}_k = \epsilon_k \mathbf{P}_k$ then give *Spectral Decomposition* of operator \mathbf{M}

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➔ *Functional spectral decomposition* **←**

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$$\mathbf{p}_j \mathbf{p}_k = \mathbf{p}_j \prod_{m \neq k} (\mathbf{M} - \varepsilon_m \mathbf{1}) = \prod_{m \neq k} (\mathbf{p}_j \mathbf{M} - \varepsilon_m \mathbf{p}_j \mathbf{1}) \quad \mathbf{M} \mathbf{p}_k = \varepsilon_k \mathbf{p}_k = \mathbf{p}_k \mathbf{M}$$

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The \mathbf{P}_j are *Mutually Ortho-Normal* as are bra-ket $\langle \varepsilon_j |$ and $|\varepsilon_j \rangle$ inside \mathbf{P}_j 's

$$\begin{pmatrix} \langle \varepsilon_1 | \varepsilon_1 \rangle & \langle \varepsilon_1 | \varepsilon_2 \rangle \\ \langle \varepsilon_2 | \varepsilon_1 \rangle & \langle \varepsilon_2 | \varepsilon_2 \rangle \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

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$$\mathbf{1} = \mathbf{P}_1 + \mathbf{P}_2 + \dots + \mathbf{P}_n = |\varepsilon_1 \rangle \langle \varepsilon_1 | + |\varepsilon_2 \rangle \langle \varepsilon_2 | + \dots + |\varepsilon_n \rangle \langle \varepsilon_n |$$

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Matrix and operator Spectral Decompositions

$$\mathbf{p}_j \mathbf{p}_k = \mathbf{p}_j \prod_{m \neq k} (\mathbf{M} - \varepsilon_m \mathbf{1}) = \prod_{m \neq k} (\mathbf{p}_j \mathbf{M} - \varepsilon_m \mathbf{p}_j \mathbf{1}) \quad \mathbf{M} \mathbf{p}_k = \varepsilon_k \mathbf{p}_k = \mathbf{p}_k \mathbf{M}$$

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Example:

$$\mathbf{M}^{50} = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix} = 1^{50} \begin{pmatrix} \frac{1}{4} & -\frac{1}{4} \\ -\frac{3}{4} & \frac{3}{4} \end{pmatrix} + 5^{50} \begin{pmatrix} \frac{3}{4} & \frac{1}{4} \\ \frac{3}{4} & \frac{1}{4} \end{pmatrix} = \frac{1}{4} \begin{pmatrix} 1+3 \cdot 5^{50} & 5^{50}-1 \\ 3 \cdot 5^{50}-3 & 5^{50}+3 \end{pmatrix}$$

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implies:
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Eigen-bra-ket projectors of matrix:

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$\{|x \rangle, |y \rangle\}$ -orthonormality with $\{|\varepsilon_1 \rangle, |\varepsilon_2 \rangle\}$ -completeness

$$\langle x | y \rangle = \delta_{x,y} = \langle x | \mathbf{1} | y \rangle = \langle x | \varepsilon_1 \rangle \langle \varepsilon_1 | y \rangle + \langle x | \varepsilon_2 \rangle \langle \varepsilon_2 | y \rangle.$$

$\{|\varepsilon_1 \rangle, |\varepsilon_2 \rangle\}$ -orthonormality with $\{|x \rangle, |y \rangle\}$ -completeness

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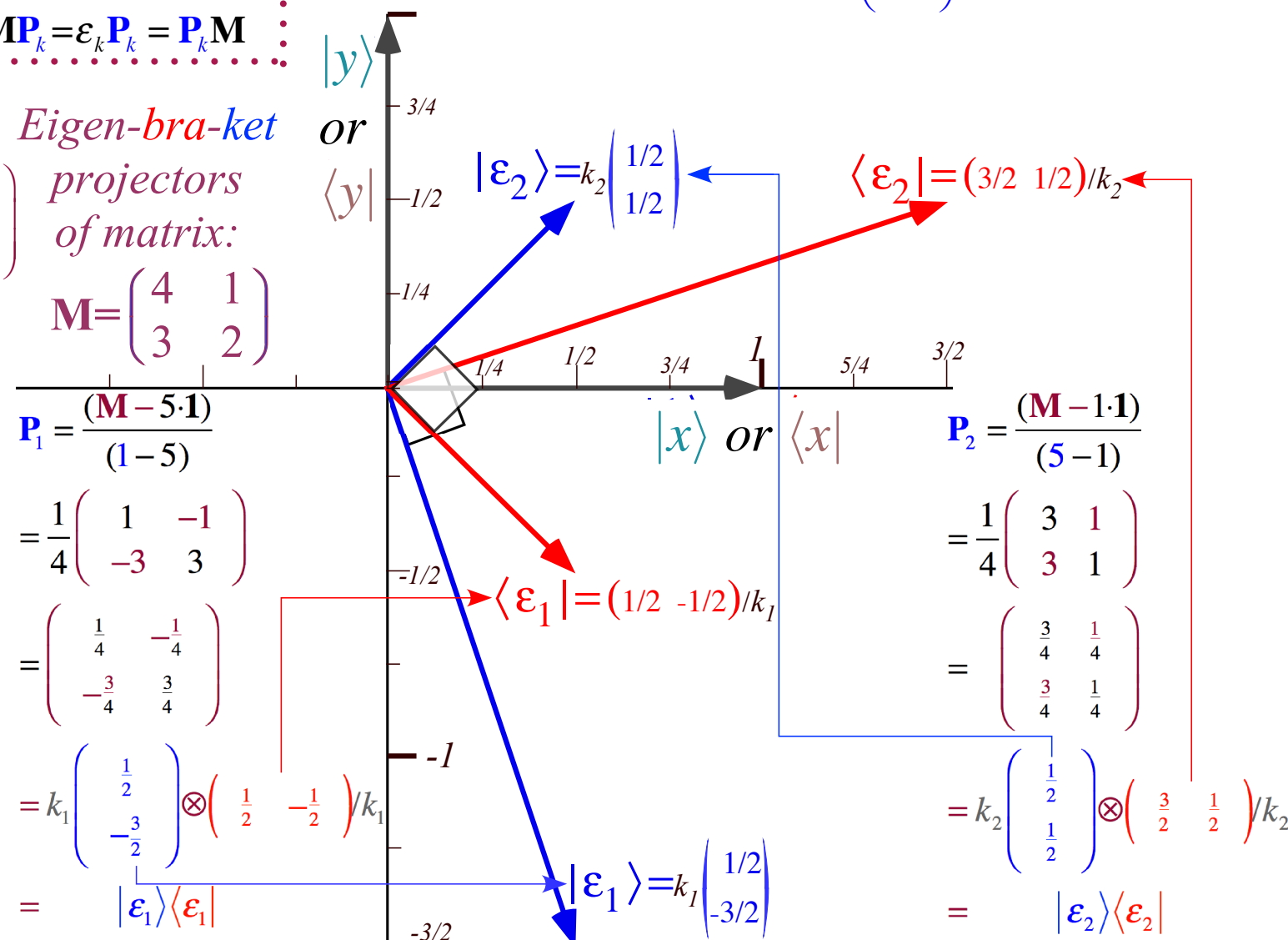
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Factoring bra-kets into "Ket-Bras:

$$\mathbf{P}_1 = \frac{(\mathbf{M} - 5 \cdot \mathbf{1})}{(1 - 5)} = \frac{1}{4} \begin{pmatrix} 1 & -1 \\ -3 & 3 \end{pmatrix} = k_1 \begin{pmatrix} \frac{1}{2} \\ -\frac{3}{2} \end{pmatrix} \otimes \frac{\begin{pmatrix} \frac{1}{2} & -\frac{1}{2} \end{pmatrix}}{k_1} = |\varepsilon_1 \rangle \langle \varepsilon_1 |$$

"Gauge" scale factors that only affect plots

$$\mathbf{P}_2 = \frac{(\mathbf{M} - 1 \cdot \mathbf{1})}{(5 - 1)} = \frac{1}{4} \begin{pmatrix} 3 & 1 \\ 3 & 1 \end{pmatrix} = k_2 \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \end{pmatrix} \otimes \frac{\begin{pmatrix} \frac{3}{2} & \frac{1}{2} \end{pmatrix}}{k_2} = |\varepsilon_2 \rangle \langle \varepsilon_2 |$$



Orthonormality vs. Completeness vis-a-vis Operator vs. State

Operator expressions for orthonormality appear quite different from expressions for completeness.

$$\mathbf{P}_j \mathbf{P}_k = \delta_{jk} \mathbf{P}_k = \begin{cases} \mathbf{0} & \text{if } j \neq k \\ \mathbf{P}_k & \text{if } j = k \end{cases}$$

$$\mathbf{1} = \mathbf{P}_1 + \mathbf{P}_2 + \dots + \mathbf{P}_n$$

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State vector representations of orthonormality are quite **similar** to representations of completeness.

Like 2-sides of the same coin.

$$\begin{aligned} & \{|x\rangle, |y\rangle\}\text{-orthonormality with } \{|\epsilon_1\rangle, |\epsilon_2\rangle\}\text{-completeness} \\ & \langle x|y\rangle = \delta_{x,y} = \langle x|\mathbf{1}|y\rangle = \langle x|\epsilon_1\rangle\langle\epsilon_1|y\rangle + \langle x|\epsilon_2\rangle\langle\epsilon_2|y\rangle. \end{aligned}$$

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$$\langle x | y \rangle = \delta(x,y) = \psi_1(x) \psi_1^*(y) + \psi_2(x) \psi_2^*(y) + \dots$$

Dirac δ -function

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However Schrodinger wavefunction notation $\psi(x) = \langle x | \psi \rangle$ shows quite a difference...

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$$\langle \epsilon_i | \epsilon_j \rangle = \delta_{i,j} = \dots + \psi_i^*(x) \psi_j(x) + \psi_2(y) \psi_2^*(y) + \dots \rightarrow \int dx \psi_i^*(x) \psi_j(x)$$

However Schrodinger wavefunction notation $\psi(x) = \langle x | \psi \rangle$ shows quite a difference...

...particularly in the orthonormality integral.

2D harmonic oscillator equations

Lagrangian and matrix forms and Reciprocity symmetry

2D harmonic oscillator equation eigensolutions

Geometric method

Matrix-algebraic eigensolutions with example $M = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix}$

Secular equation

Hamilton-Cayley equation and projectors

Idempotent projectors (how eigenvalues \Rightarrow eigenvectors)

Operator orthonormality and Completeness (Idempotent means: $\mathbf{P} \cdot \mathbf{P} = \mathbf{P}$)

Spectral Decompositions

Functional spectral decomposition

Orthonormality vs. Completeness vis-a`-vis Operator vs. State

➔ *Lagrange functional interpolation formula* **←**

Diagonalizing Transformations (D-Ttran) from projectors

2D-HO eigensolution example with bilateral (B-Type) symmetry

Mixed mode beat dynamics and fixed $\pi/2$ phase

2D-HO eigensolution example with asymmetric (A-Type) symmetry

Initial state projection, mixed mode beat dynamics with variable phase

A Proof of Projector Completeness (Truer-than-true by Lagrange interpolation)

Compare matrix completeness relation and functional spectral decompositions

$$\mathbf{1} = \mathbf{P}_1 + \mathbf{P}_2 + \dots + \mathbf{P}_n = \sum_{\epsilon_k} \mathbf{P}_k = \sum_{\epsilon_k} \frac{\prod_{m \neq k} (\mathbf{M} - \epsilon_m \mathbf{1})}{\prod_{m \neq k} (\epsilon_k - \epsilon_m)} \quad f(\mathbf{M}) = f(\epsilon_1)\mathbf{P}_1 + f(\epsilon_2)\mathbf{P}_2 + \dots + f(\epsilon_n)\mathbf{P}_n = \sum_{\epsilon_k} f(\epsilon_k)\mathbf{P}_k = \sum_{\epsilon_k} f(\epsilon_k) \frac{\prod_{m \neq k} (\mathbf{M} - \epsilon_m \mathbf{1})}{\prod_{m \neq k} (\epsilon_k - \epsilon_m)}$$

with *Lagrange interpolation formula* of function $f(x)$ approximated by its value at N points x_1, x_2, \dots, x_N .

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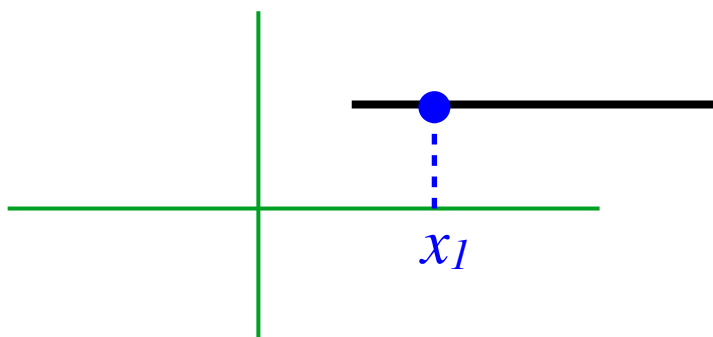
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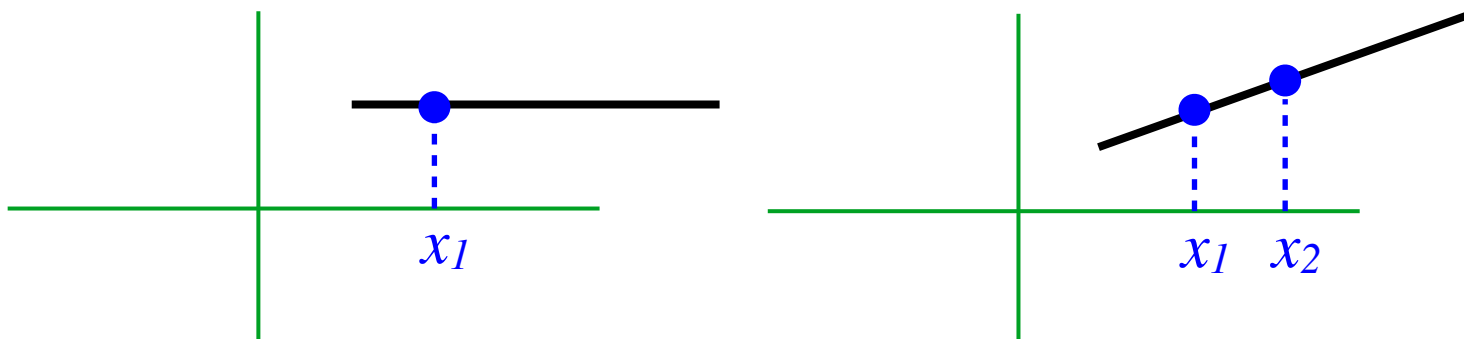
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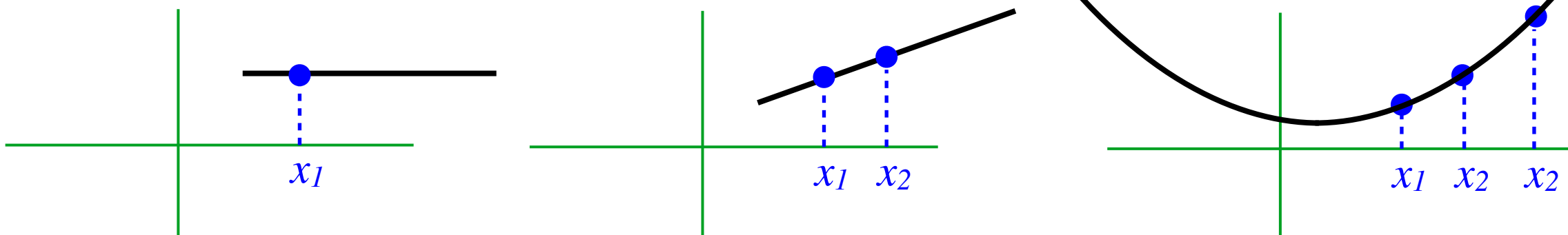
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All distinct values $\epsilon_1 \neq \epsilon_2 \neq \dots \neq \epsilon_N$ satisfy $\sum \mathbf{P}_k = \mathbf{1}$.

A Proof of Projector Completeness (Truer-than-true)

Compare matrix *completeness relation* and *functional spectral decompositions*

$$\mathbf{1} = \mathbf{P}_1 + \mathbf{P}_2 + \dots + \mathbf{P}_n = \sum_{\epsilon_k} \mathbf{P}_k = \sum_{\epsilon_k} \frac{\prod_{m \neq k} (\mathbf{M} - \epsilon_m \mathbf{1})}{\prod_{m \neq k} (\epsilon_k - \epsilon_m)} \quad f(\mathbf{M}) = f(\epsilon_1)\mathbf{P}_1 + f(\epsilon_2)\mathbf{P}_2 + \dots + f(\epsilon_n)\mathbf{P}_n = \sum_{\epsilon_k} f(\epsilon_k)\mathbf{P}_k = \sum_{\epsilon_k} f(\epsilon_k) \frac{\prod_{m \neq k} (\mathbf{M} - \epsilon_m \mathbf{1})}{\prod_{m \neq k} (\epsilon_k - \epsilon_m)}$$

with *Lagrange interpolation formula* of function $f(x)$ approximated by its value at N points x_1, x_2, \dots, x_N .

$$L(f(x)) = \sum_{k=1}^N f(x_k) P_k(x) \quad \text{where: } P_k(x) = \frac{\prod_{j \neq k} (x - x_j)}{\prod_{j \neq k} (x_k - x_j)}$$

Each polynomial term $P_m(x)$ has zeros at each point $x=x_j$ *except* where $x=x_m$. Then $P_m(x_m)=1$. So at each of these points this L-approximation becomes exact: $L(f(x_j))=f(x_j)$.

If $f(x)$ happens to be a polynomial of degree $N-1$ or less, then $L(f(x))=f(x)$ may be exact everywhere.

$$1 = \sum_{m=1}^N P_m(x) \quad x = \sum_{m=1}^N x_m P_m(x) \quad x^2 = \sum_{m=1}^N x_m^2 P_m(x)$$

One point determines a constant level line, two separate points uniquely determine a sloping line, three separate points uniquely determine a parabola, etc.

Lagrange interpolation formula \rightarrow *Completeness formula* as $x \rightarrow \mathbf{M}$ and as $x_k \rightarrow \epsilon_k$ and as $P_k(x_k) \rightarrow \mathbf{P}_k$

All distinct values $\epsilon_1 \neq \epsilon_2 \neq \dots \neq \epsilon_N$ satisfy $\sum \mathbf{P}_k = \mathbf{1}$. Completeness is *truer than true* as is seen for $N=2$.

$$\mathbf{P}_1 + \mathbf{P}_2 = \frac{\prod_{j \neq 1} (\mathbf{M} - \epsilon_j \mathbf{1})}{\prod_{j \neq 1} (\epsilon_1 - \epsilon_j)} + \frac{\prod_{j \neq 2} (\mathbf{M} - \epsilon_j \mathbf{1})}{\prod_{j \neq 2} (\epsilon_2 - \epsilon_j)} = \frac{(\mathbf{M} - \epsilon_2 \mathbf{1})}{(\epsilon_1 - \epsilon_2)} + \frac{(\mathbf{M} - \epsilon_1 \mathbf{1})}{(\epsilon_2 - \epsilon_1)} = \frac{(\mathbf{M} - \epsilon_2 \mathbf{1}) - (\mathbf{M} - \epsilon_1 \mathbf{1})}{(\epsilon_1 - \epsilon_2)} = \frac{-\epsilon_2 \mathbf{1} + \epsilon_1 \mathbf{1}}{(\epsilon_1 - \epsilon_2)} = \mathbf{1} \text{ (for all } \epsilon_j \text{)}$$

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However, only *select* values ϵ_k work for eigen-forms $\mathbf{M}\mathbf{P}_k = \epsilon_k \mathbf{P}_k$ or orthonormality $\mathbf{P}_j \mathbf{P}_k = \delta_{jk} \mathbf{P}_k$.

2D harmonic oscillator equations

Lagrangian and matrix forms and Reciprocity symmetry

2D harmonic oscillator equation eigensolutions

Geometric method

Matrix-algebraic eigensolutions with example $M = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix}$

Secular equation

Hamilton-Cayley equation and projectors

Idempotent projectors (how eigenvalues \Rightarrow eigenvectors)

Operator orthonormality and Completeness (Idempotent means: $\mathbf{P} \cdot \mathbf{P} = \mathbf{P}$)

Spectral Decompositions

Functional spectral decomposition

Orthonormality vs. Completeness vis-a`-vis Operator vs. State

Lagrange functional interpolation formula

➔ *Diagonalizing Transformations (D-Tran) from projectors* **←**

2D-HO eigensolution example with bilateral (B-Type) symmetry

Mixed mode beat dynamics and fixed $\pi/2$ phase

2D-HO eigensolution example with asymmetric (A-Type) symmetry

Initial state projection, mixed mode beat dynamics with variable phase

Diagonalizing Transformations (D-Ttran) from projectors

Given our eigenvectors and their Projectors.

$$\mathbf{P}_1 = \frac{(\mathbf{M} - 5 \cdot \mathbf{1})}{(1 - 5)} = \frac{1}{4} \begin{pmatrix} 1 & -1 \\ -3 & 3 \end{pmatrix} = k_1 \begin{pmatrix} \frac{1}{2} \\ -\frac{3}{2} \end{pmatrix} \otimes \frac{\begin{pmatrix} \frac{1}{2} & -\frac{1}{2} \end{pmatrix}}{k_1} = |\boldsymbol{\varepsilon}_1\rangle\langle\boldsymbol{\varepsilon}_1|$$
$$\mathbf{P}_2 = \frac{(\mathbf{M} - 1 \cdot \mathbf{1})}{(5 - 1)} = \frac{1}{4} \begin{pmatrix} 3 & 1 \\ 3 & 1 \end{pmatrix} = k_2 \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \end{pmatrix} \otimes \frac{\begin{pmatrix} \frac{3}{2} & \frac{1}{2} \end{pmatrix}}{k_2} = |\boldsymbol{\varepsilon}_2\rangle\langle\boldsymbol{\varepsilon}_2|$$

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Load distinct bras $\langle\varepsilon_1|$ and $\langle\varepsilon_2|$ into d-tran **rows**, kets $|\varepsilon_1\rangle$ and $|\varepsilon_2\rangle$ into inverse d-tran **columns**.

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Load distinct bras $\langle\epsilon_1|$ and $\langle\epsilon_2|$ into d-tran **rows**, kets $|\epsilon_1\rangle$ and $|\epsilon_2\rangle$ into inverse d-tran **columns**.

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$(\epsilon_1, \epsilon_2) \leftarrow (1, 2)$ *d-Tran matrix*

$(1, 2) \leftarrow (\epsilon_1, \epsilon_2)$ *INVERSE d-Tran matrix*

$$\begin{pmatrix} \langle\epsilon_1|x\rangle & \langle\epsilon_1|y\rangle \\ \langle\epsilon_2|x\rangle & \langle\epsilon_2|y\rangle \end{pmatrix} = \begin{pmatrix} \frac{1}{2} & -\frac{1}{2} \\ \frac{3}{2} & \frac{1}{2} \end{pmatrix}, \quad \begin{pmatrix} \langle x|\epsilon_1\rangle & \langle x|\epsilon_2\rangle \\ \langle y|\epsilon_1\rangle & \langle y|\epsilon_2\rangle \end{pmatrix} = \begin{pmatrix} \frac{1}{2} & \frac{1}{2} \\ -\frac{3}{2} & \frac{1}{2} \end{pmatrix}$$

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$$\mathbf{P}_2 = \frac{(\mathbf{M} - 1\mathbf{1})}{(5-1)} = \frac{1}{4} \begin{pmatrix} 3 & 1 \\ 3 & 1 \end{pmatrix} = k_2 \begin{pmatrix} \frac{1}{2} \\ \frac{1}{2} \end{pmatrix} \otimes \frac{\begin{pmatrix} \frac{3}{2} & \frac{1}{2} \end{pmatrix}}{k_2} = |\varepsilon_2\rangle\langle\varepsilon_2|$$

Load distinct bras $\langle\varepsilon_1|$ and $\langle\varepsilon_2|$ into d-tran **rows**, kets $|\varepsilon_1\rangle$ and $|\varepsilon_2\rangle$ into inverse d-tran **columns**.

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$(1,2) \leftarrow (\varepsilon_1, \varepsilon_2)$ *INVERSE d-Tran matrix*

$$\begin{pmatrix} \langle\varepsilon_1|x\rangle & \langle\varepsilon_1|y\rangle \\ \langle\varepsilon_2|x\rangle & \langle\varepsilon_2|y\rangle \end{pmatrix} = \begin{pmatrix} \frac{1}{2} & -\frac{1}{2} \\ \frac{3}{2} & \frac{1}{2} \end{pmatrix}, \quad \begin{pmatrix} \langle x|\varepsilon_1\rangle & \langle x|\varepsilon_2\rangle \\ \langle y|\varepsilon_1\rangle & \langle y|\varepsilon_2\rangle \end{pmatrix} = \begin{pmatrix} \frac{1}{2} & \frac{1}{2} \\ -\frac{3}{2} & \frac{1}{2} \end{pmatrix}$$

Use *Dirac labeling for all components* so transformation is OK

$$\begin{pmatrix} \langle\varepsilon_1|x\rangle & \langle\varepsilon_1|y\rangle \\ \langle\varepsilon_2|x\rangle & \langle\varepsilon_2|y\rangle \end{pmatrix} \cdot \begin{pmatrix} \langle x|\mathbf{K}|x\rangle & \langle x|\mathbf{K}|y\rangle \\ \langle y|\mathbf{K}|x\rangle & \langle y|\mathbf{K}|y\rangle \end{pmatrix} \cdot \begin{pmatrix} \langle x|\varepsilon_1\rangle & \langle x|\varepsilon_2\rangle \\ \langle y|\varepsilon_1\rangle & \langle y|\varepsilon_2\rangle \end{pmatrix} = \begin{pmatrix} \langle\varepsilon_1|\mathbf{K}|\varepsilon_1\rangle & \langle\varepsilon_1|\mathbf{K}|\varepsilon_2\rangle \\ \langle\varepsilon_2|\mathbf{K}|\varepsilon_1\rangle & \langle\varepsilon_2|\mathbf{K}|\varepsilon_2\rangle \end{pmatrix}$$

$$\begin{pmatrix} \frac{1}{2} & -\frac{1}{2} \\ \frac{3}{2} & \frac{1}{2} \end{pmatrix} \cdot \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix} \cdot \begin{pmatrix} \frac{1}{2} & \frac{1}{2} \\ -\frac{3}{2} & \frac{1}{2} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 5 \end{pmatrix}$$

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Given our eigenvectors and their Projectors.

$$\mathbf{P}_1 = \frac{(\mathbf{M} - 5\mathbf{1})}{(1-5)} = \frac{1}{4} \begin{pmatrix} 1 & -1 \\ -3 & 3 \end{pmatrix} = k_1 \begin{pmatrix} \frac{1}{2} \\ -\frac{3}{2} \end{pmatrix} \otimes \frac{\begin{pmatrix} \frac{1}{2} & -\frac{1}{2} \end{pmatrix}}{k_1} = |\epsilon_1\rangle\langle\epsilon_1|$$

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Load distinct bras $\langle\epsilon_1|$ and $\langle\epsilon_2|$ into d-tran **rows**, kets $|\epsilon_1\rangle$ and $|\epsilon_2\rangle$ into inverse d-tran **columns**.

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Use *Dirac labeling for all components* so transformation is OK

$$\begin{pmatrix} \langle\epsilon_1|x\rangle & \langle\epsilon_1|y\rangle \\ \langle\epsilon_2|x\rangle & \langle\epsilon_2|y\rangle \end{pmatrix} \cdot \begin{pmatrix} \langle x|\mathbf{K}|x\rangle & \langle x|\mathbf{K}|y\rangle \\ \langle y|\mathbf{K}|x\rangle & \langle y|\mathbf{K}|y\rangle \end{pmatrix} \cdot \begin{pmatrix} \langle x|\epsilon_1\rangle & \langle x|\epsilon_2\rangle \\ \langle y|\epsilon_1\rangle & \langle y|\epsilon_2\rangle \end{pmatrix} = \begin{pmatrix} \langle\epsilon_1|\mathbf{K}|\epsilon_1\rangle & \langle\epsilon_1|\mathbf{K}|\epsilon_2\rangle \\ \langle\epsilon_2|\mathbf{K}|\epsilon_1\rangle & \langle\epsilon_2|\mathbf{K}|\epsilon_2\rangle \end{pmatrix}$$

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Check inverse-d-tran is really inverse of *your* d-tran.

$$\begin{pmatrix} \langle\epsilon_1|1\rangle & \langle\epsilon_1|2\rangle \\ \langle\epsilon_2|1\rangle & \langle\epsilon_2|2\rangle \end{pmatrix} \cdot \begin{pmatrix} \langle 1|\epsilon_1\rangle & \langle 1|\epsilon_2\rangle \\ \langle 2|\epsilon_1\rangle & \langle 2|\epsilon_2\rangle \end{pmatrix} = \begin{pmatrix} \langle\epsilon_1|\mathbf{1}|\epsilon_1\rangle & \langle\epsilon_1|\mathbf{1}|\epsilon_2\rangle \\ \langle\epsilon_2|\mathbf{1}|\epsilon_1\rangle & \langle\epsilon_2|\mathbf{1}|\epsilon_2\rangle \end{pmatrix}$$

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Diagonalizing Transformations (D-Tran) from projectors

Given our eigenvectors and their Projectors.

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$(\epsilon_1, \epsilon_2) \leftarrow (1, 2)$ *d-Tran matrix*

$(1, 2) \leftarrow (\epsilon_1, \epsilon_2)$ **INVERSE** *d-Tran matrix*

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Check inverse-d-tran is really inverse of *your* d-tran. In standard quantum matrices inverses are “easy”

$$\begin{pmatrix} \langle\epsilon_1|x\rangle & \langle\epsilon_1|y\rangle \\ \langle\epsilon_2|x\rangle & \langle\epsilon_2|y\rangle \end{pmatrix} \cdot \begin{pmatrix} \langle x|\epsilon_1\rangle & \langle x|\epsilon_2\rangle \\ \langle y|\epsilon_1\rangle & \langle y|\epsilon_2\rangle \end{pmatrix} = \begin{pmatrix} \langle\epsilon_1|\mathbf{1}|\epsilon_1\rangle & \langle\epsilon_1|\mathbf{1}|\epsilon_2\rangle \\ \langle\epsilon_2|\mathbf{1}|\epsilon_1\rangle & \langle\epsilon_2|\mathbf{1}|\epsilon_2\rangle \end{pmatrix} \begin{pmatrix} \langle\epsilon_1|x\rangle & \langle\epsilon_1|y\rangle \\ \langle\epsilon_2|x\rangle & \langle\epsilon_2|y\rangle \end{pmatrix} = \begin{pmatrix} \langle x|\epsilon_1\rangle & \langle x|\epsilon_2\rangle \\ \langle y|\epsilon_1\rangle & \langle y|\epsilon_2\rangle \end{pmatrix}^\dagger = \begin{pmatrix} \langle x|\epsilon_1\rangle^* & \langle y|\epsilon_1\rangle^* \\ \langle x|\epsilon_2\rangle^* & \langle y|\epsilon_2\rangle^* \end{pmatrix} = \begin{pmatrix} \langle x|\epsilon_1\rangle & \langle x|\epsilon_2\rangle \\ \langle y|\epsilon_1\rangle & \langle y|\epsilon_2\rangle \end{pmatrix}^{-1}$$

$$\begin{pmatrix} \frac{1}{2} & -\frac{1}{2} \\ \frac{3}{2} & \frac{1}{2} \end{pmatrix} \cdot \begin{pmatrix} \frac{1}{2} & \frac{1}{2} \\ -\frac{3}{2} & \frac{1}{2} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

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Idempotent projectors (how eigenvalues \Rightarrow eigenvectors)

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Orthonormality vs. Completeness vis-a`-vis Operator vs. State

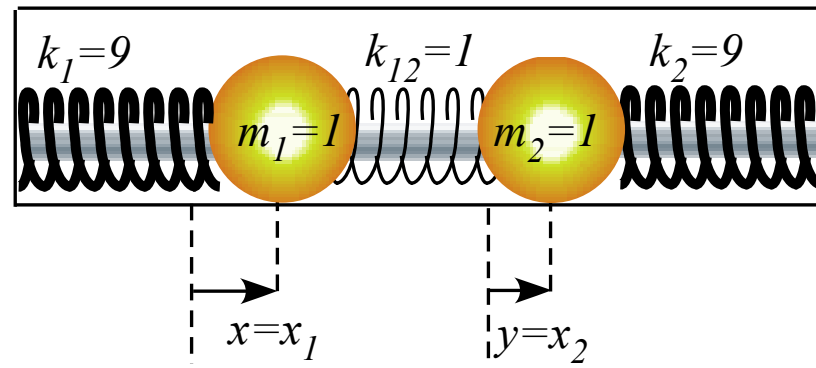
Lagrange functional interpolation formula

Diagonalizing Transformations (D-Tran) from projectors

➔ *2D-HO eigensolution example with bilateral (B-Type) symmetry* **←**
Mixed mode beat dynamics and fixed $\pi/2$ phase

2D-HO eigensolution example with asymmetric (A-Type) symmetry
Initial state projection, mixed mode beat dynamics with variable phase

Analyzing 2D-HO beats and mixed mode eigen-solutions



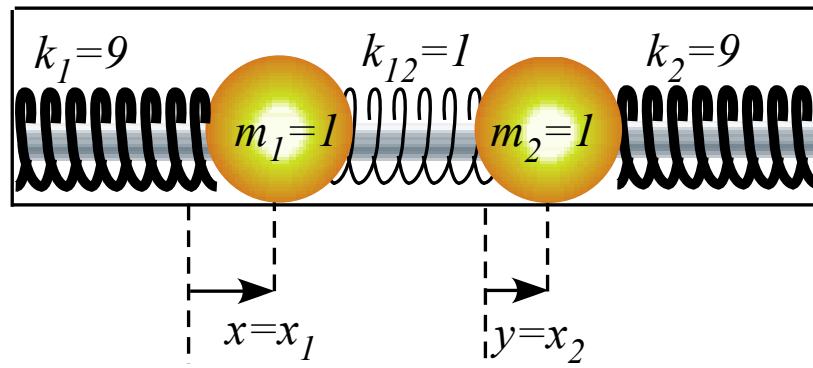
$$\mathbf{K} = \begin{pmatrix} K_{11} & K_{12} \\ K_{12} & K_{22} \end{pmatrix} = \begin{pmatrix} k_1 + k_{12} & -k_{12} \\ -k_{12} & k_2 + k_{12} \end{pmatrix} = \begin{pmatrix} 10 & -1 \\ -1 & 10 \end{pmatrix}$$

$Det(\mathbf{K}) = 10 \cdot 10 - 1 = 99$
 $Trace(\mathbf{K}) = 10 + 10 = 20$

The \mathbf{K} secular equation $K^2 - Trace(\mathbf{K})K + Det(\mathbf{K}) = K^2 - 20K + 99 = 0 = (K - 9)(K - 11) = (K - K_1)(K - K_2)$

Eigenvalues K_k and squared eigenfrequencies $\omega_0(\epsilon_k)^2$ $K_1 = \omega_0^2(\epsilon_1) = 9, \quad K_2 = \omega_0^2(\epsilon_2) = 11,$

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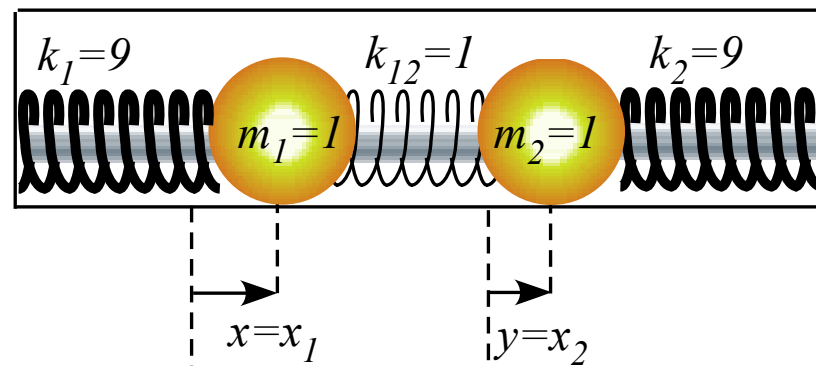
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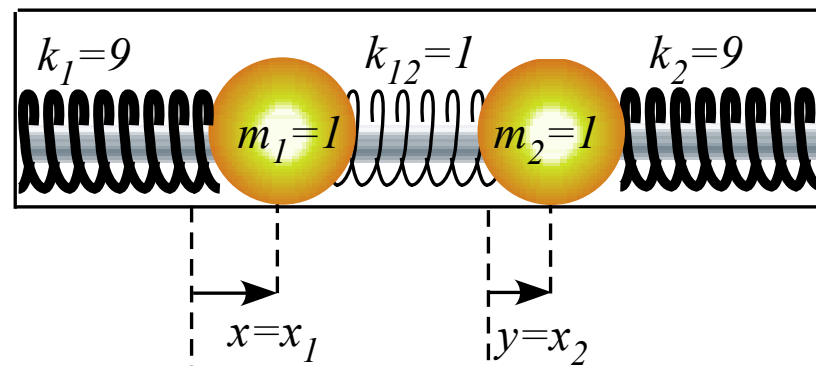
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Eigenbra vectors: $\langle\epsilon_1| = \left(1/\sqrt{2} \quad +1/\sqrt{2} \right)$, $\langle\epsilon_2| = \left(1/\sqrt{2} \quad -1/\sqrt{2} \right)$

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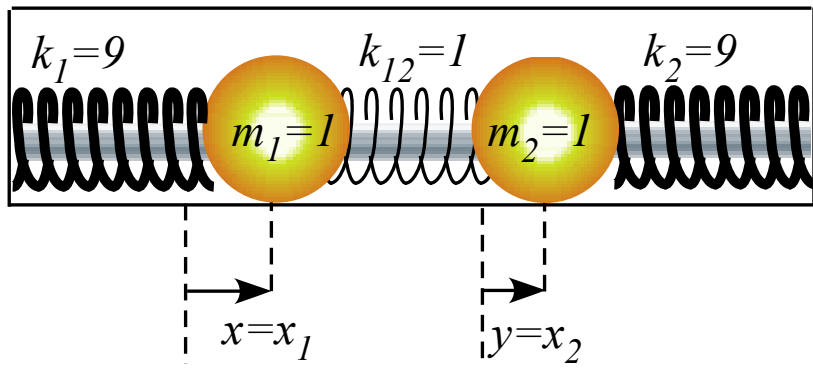
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$$|x(t)\rangle = |\epsilon_1\rangle \langle\epsilon_1|x(0)\rangle e^{-i\omega_1 t} + |\epsilon_2\rangle \langle\epsilon_2|x(0)\rangle e^{-i\omega_2 t}$$

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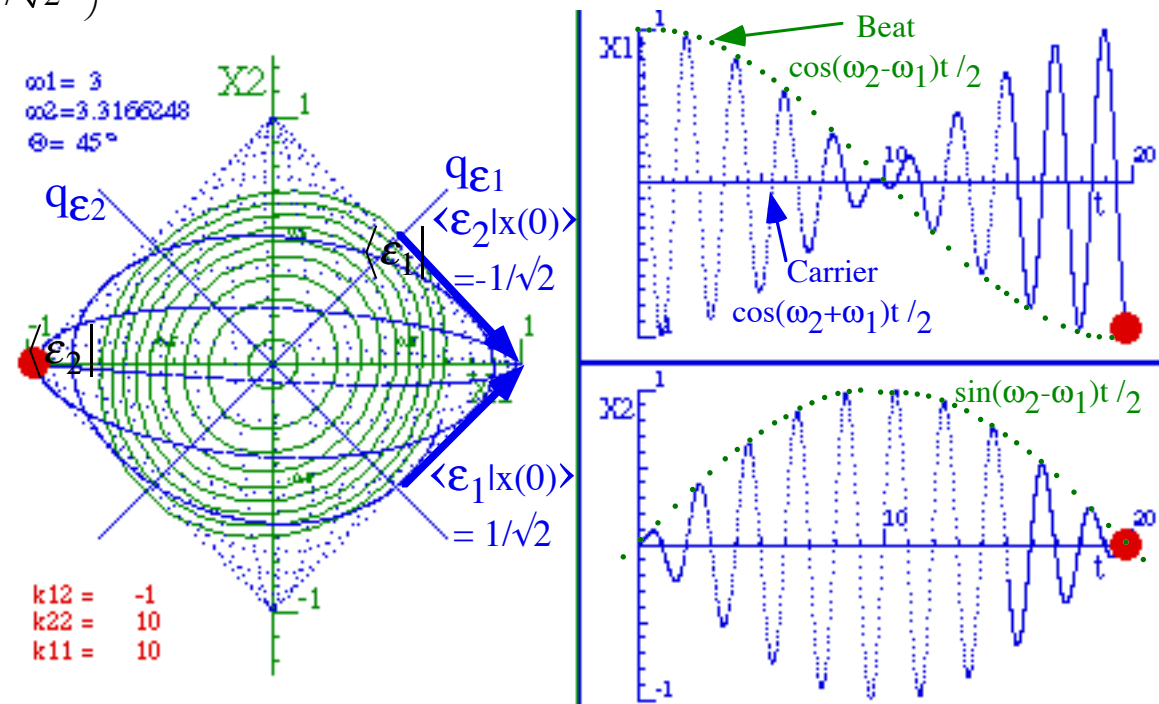
Mixed mode dynamics

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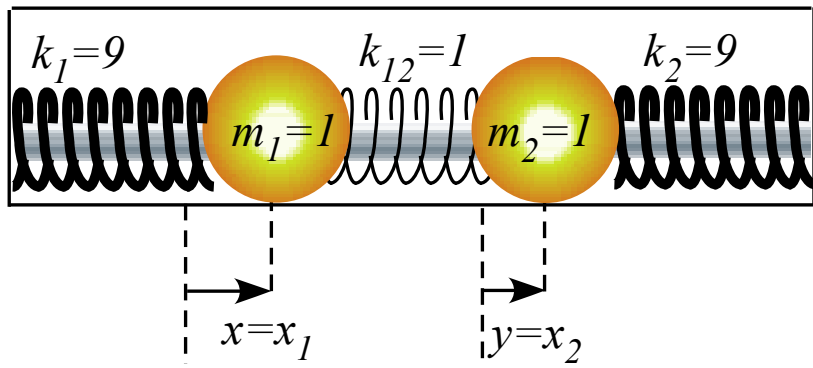
100% modulation (SWR=0) $\frac{e^{ia} + e^{ib}}{2} = e^{\frac{i(a+b)}{2}} \frac{e^{\frac{i(a-b)}{2}} + e^{-\frac{i(a-b)}{2}}}{2}$

$$\begin{pmatrix} x_1(t) \\ x_2(t) \end{pmatrix} = \begin{pmatrix} \frac{e^{-i\omega_1 t} + e^{-i\omega_2 t}}{2} \\ \frac{e^{-i\omega_1 t} - e^{-i\omega_2 t}}{2} \end{pmatrix} = \frac{e^{-\frac{i(\omega_1 + \omega_2)t}}{2}} \begin{pmatrix} e^{-\frac{i(\omega_1 - \omega_2)t}{2}} + e^{\frac{i(\omega_1 - \omega_2)t}{2}} \\ e^{-\frac{i(\omega_1 - \omega_2)t}{2}} - e^{\frac{i(\omega_1 - \omega_2)t}{2}} \end{pmatrix}$$



BoxIt (Beating) Simulation

Fig. 3.3.9 Beats in weakly coupled symmetric oscillators with equal mode magnitudes.



Analyzing 2D-HO beats and mixed mode eigen-solutions

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$$= \begin{pmatrix} 1/\sqrt{2} \\ -1/\sqrt{2} \end{pmatrix} \begin{pmatrix} 1/\sqrt{2} & -1/\sqrt{2} \end{pmatrix} = |\epsilon_2\rangle\langle\epsilon_2|$$

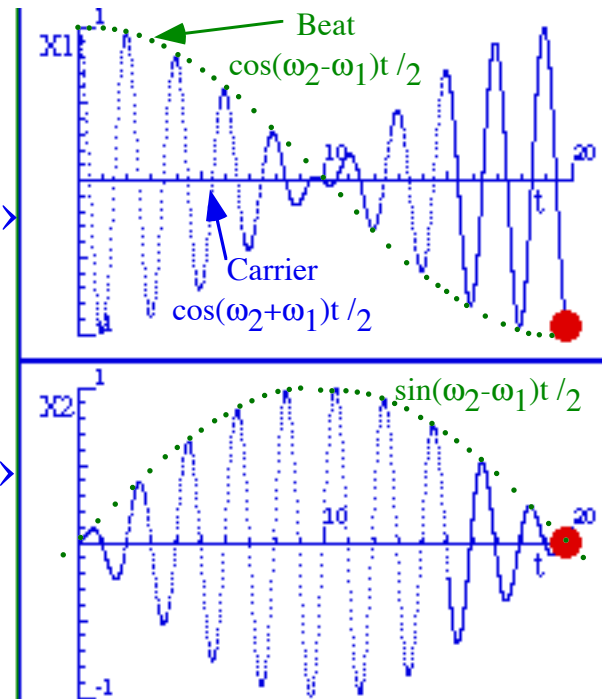
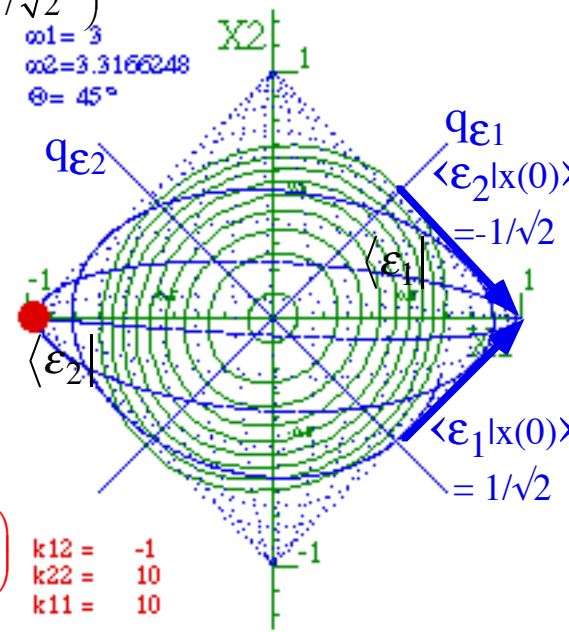
Eigenbra vectors: $\langle\epsilon_1| = \left(1/\sqrt{2} \quad +1/\sqrt{2} \right)$, $\langle\epsilon_2| = \left(1/\sqrt{2} \quad -1/\sqrt{2} \right)$

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100% modulation (SWR=0) $\frac{e^{ia} + e^{ib}}{2} = e^{i\frac{a+b}{2}} \frac{e^{i\frac{a-b}{2}} + e^{-i\frac{a-b}{2}}}{2} = e^{i\frac{a+b}{2}} \cos\left(\frac{a-b}{2}\right)$



BoxIt (Beating) Simulation

Note the i phase

Fig. 3.3.9 Beats in weakly coupled symmetric oscillators with equal mode magnitudes.

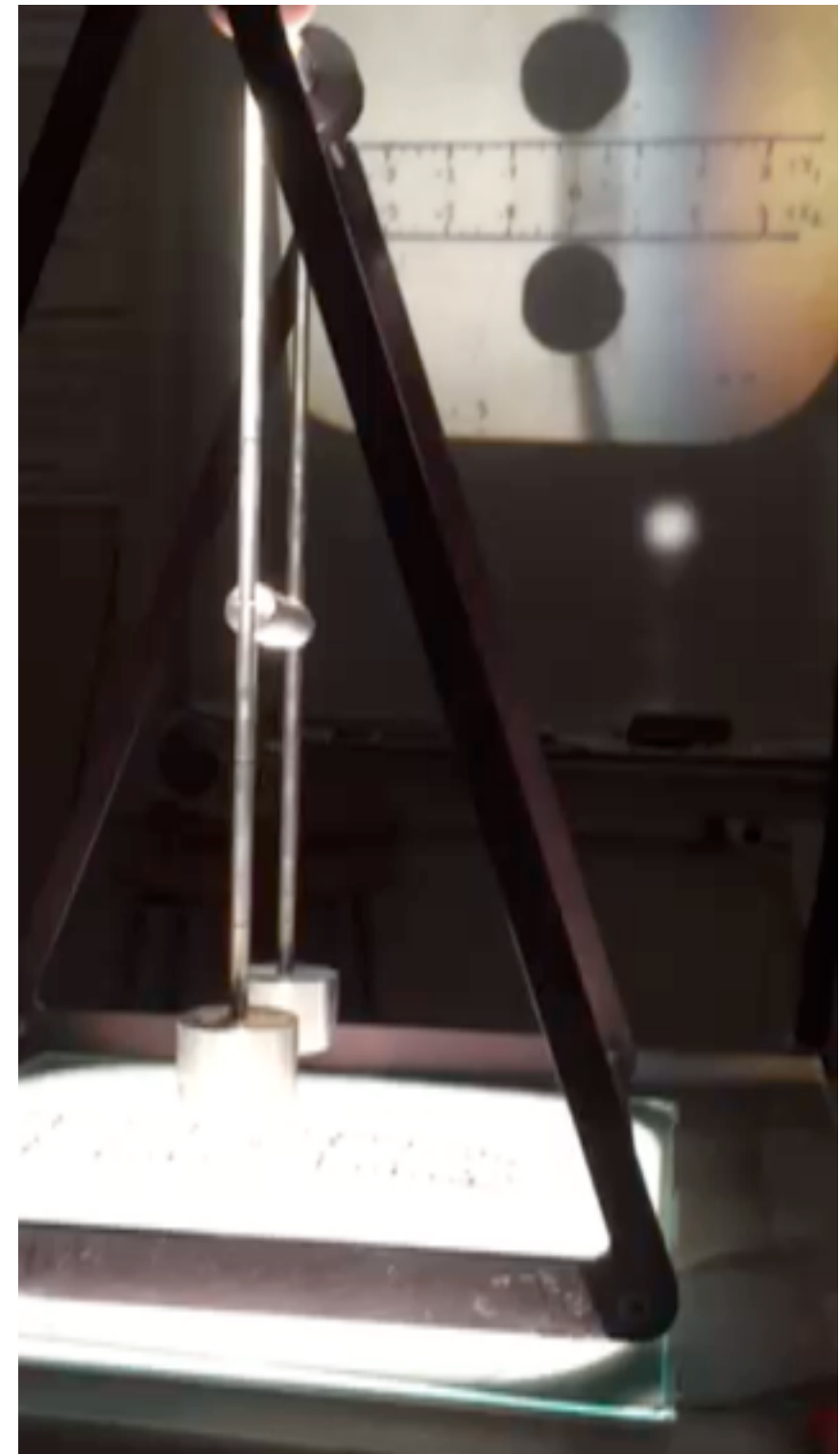
Videos of Coupled Pendula aided by Overhead Projector



[View on YouTube](#) 

*Launch embedded videos
using your browser/App
or*

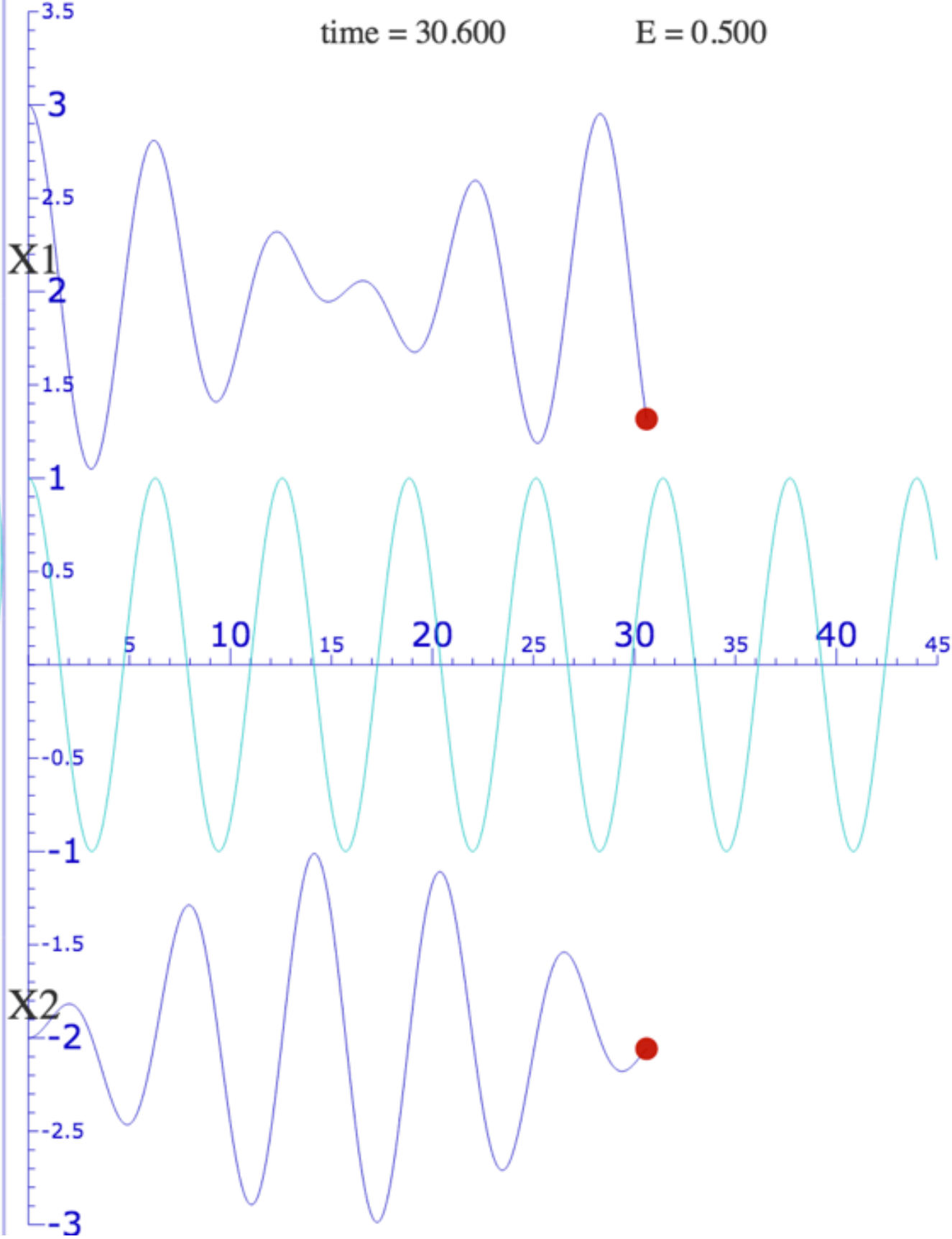
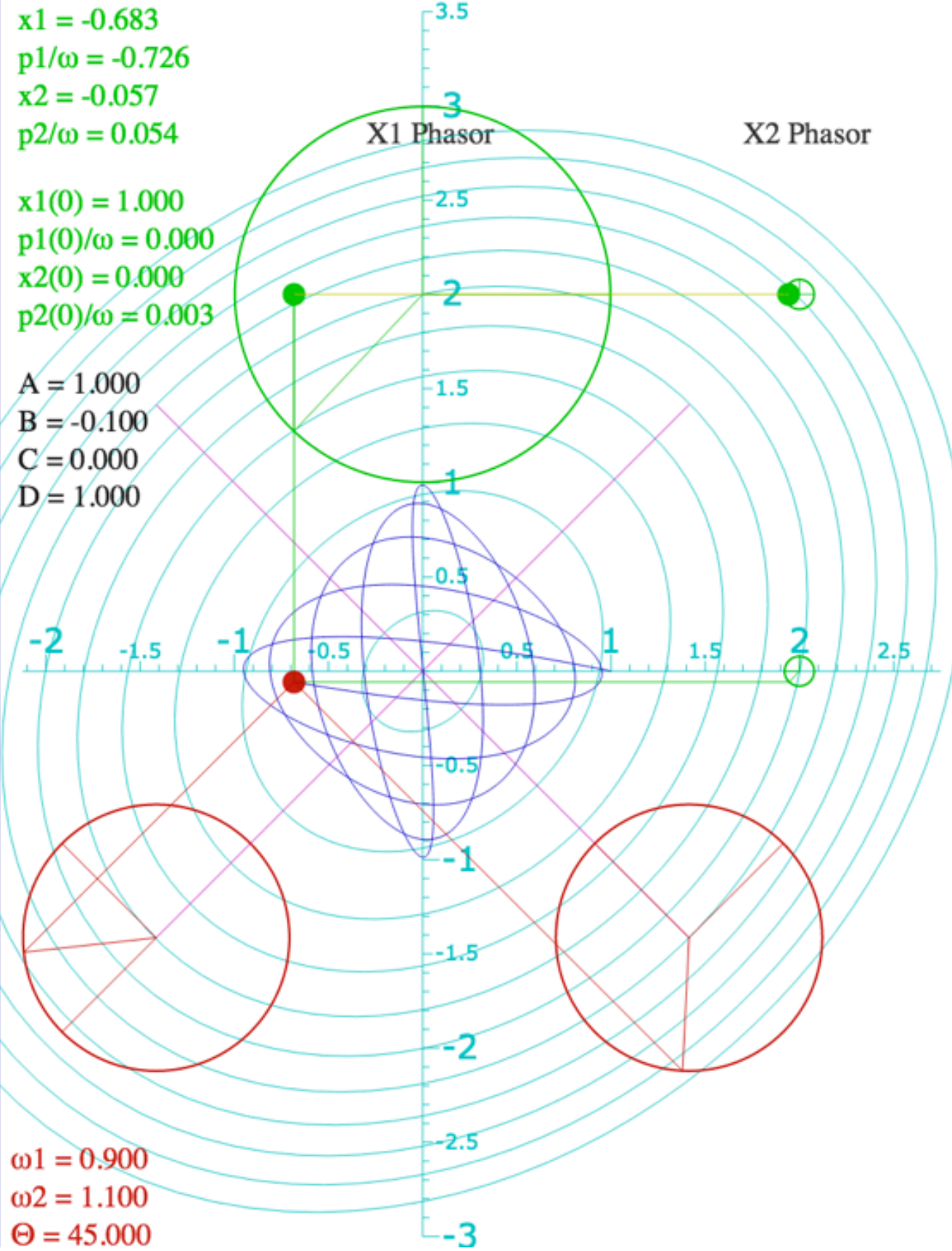
⇐ view on YouTube ⇒



[View on YouTube](#) 

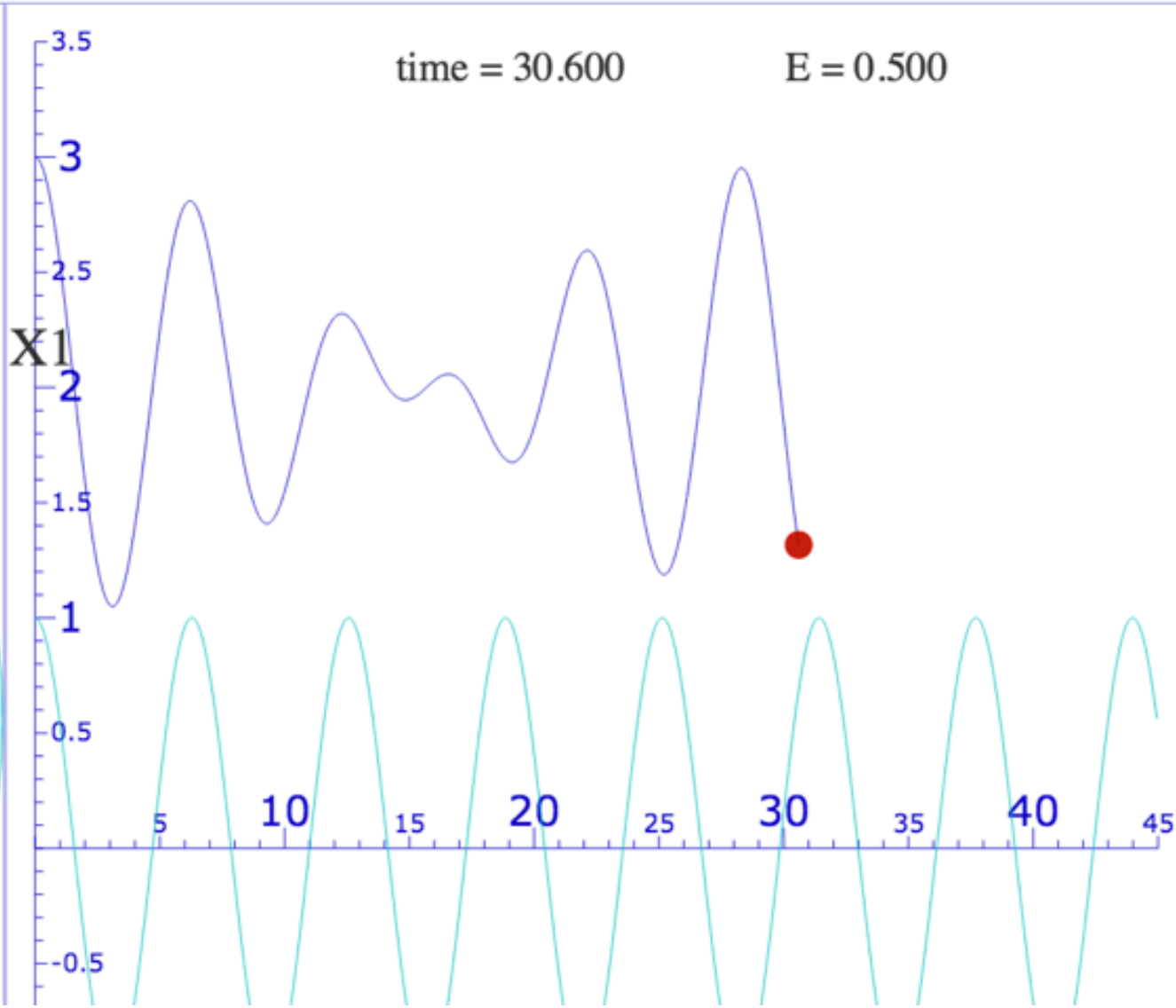
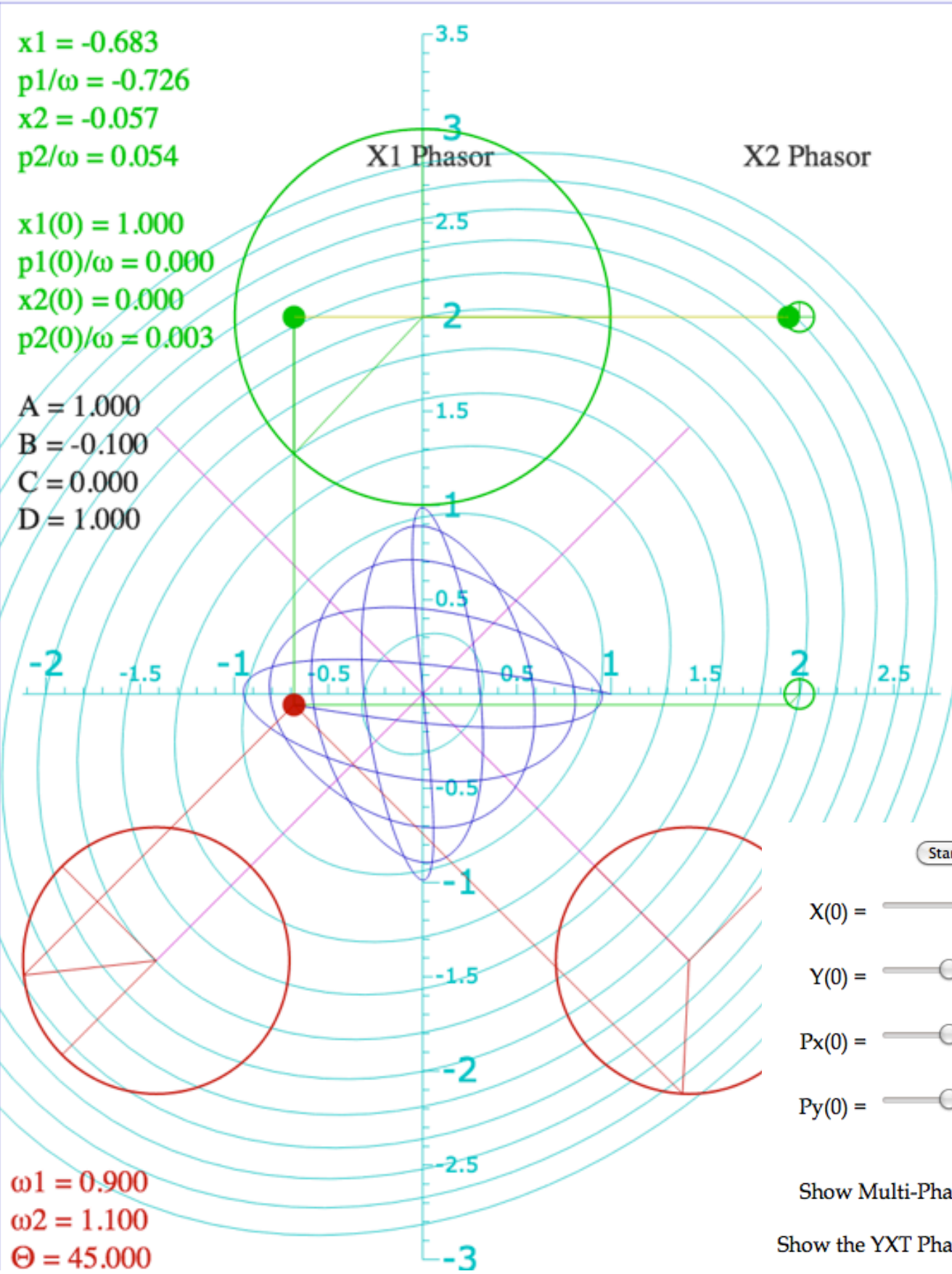
Stronger coupling on the right, illustrated indirectly by a darker looking spring on screen

$x1 = -0.683$
 $p1/\omega = -0.726$
 $x2 = -0.057$
 $p2/\omega = 0.054$
 $x1(0) = 1.000$
 $p1(0)/\omega = 0.000$
 $x2(0) = 0.000$
 $p2(0)/\omega = 0.003$
 $A = 1.000$
 $B = -0.100$
 $C = 0.000$
 $D = 1.000$



[BoxIt \(Beating\) Web Simulation](#)
(A=1, B=-0.1, C=0, D=1)

$x1 = -0.683$
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 $C = 0.000$
 $D = 1.000$



$X(0) =$ $A =$ $\text{Number of Derivatives} =$
 $Y(0) =$ $B =$
 $Px(0) =$ $C =$
 $Py(0) =$ $D =$

- Show Multi-Phasor View
- Show the YXT Phasor View
- Draw Main Phasors
- Draw Vector Heads
- wantVectorHeads, wantTimeRateTangents
- Draw PE Levels
- Draw Box Lines
- Draw Modal Phasors
- Draw Time Rate Tangents
- Left Phasor Rides on Right Phasor
- Left Phasor Rides on Right Phasor
- Normalize Phasors
- Print $\omega1:\omega2$ fractions

BoxIt (Beating) Web Simulation
(A=1, B=-0.1, C=0, D=1)

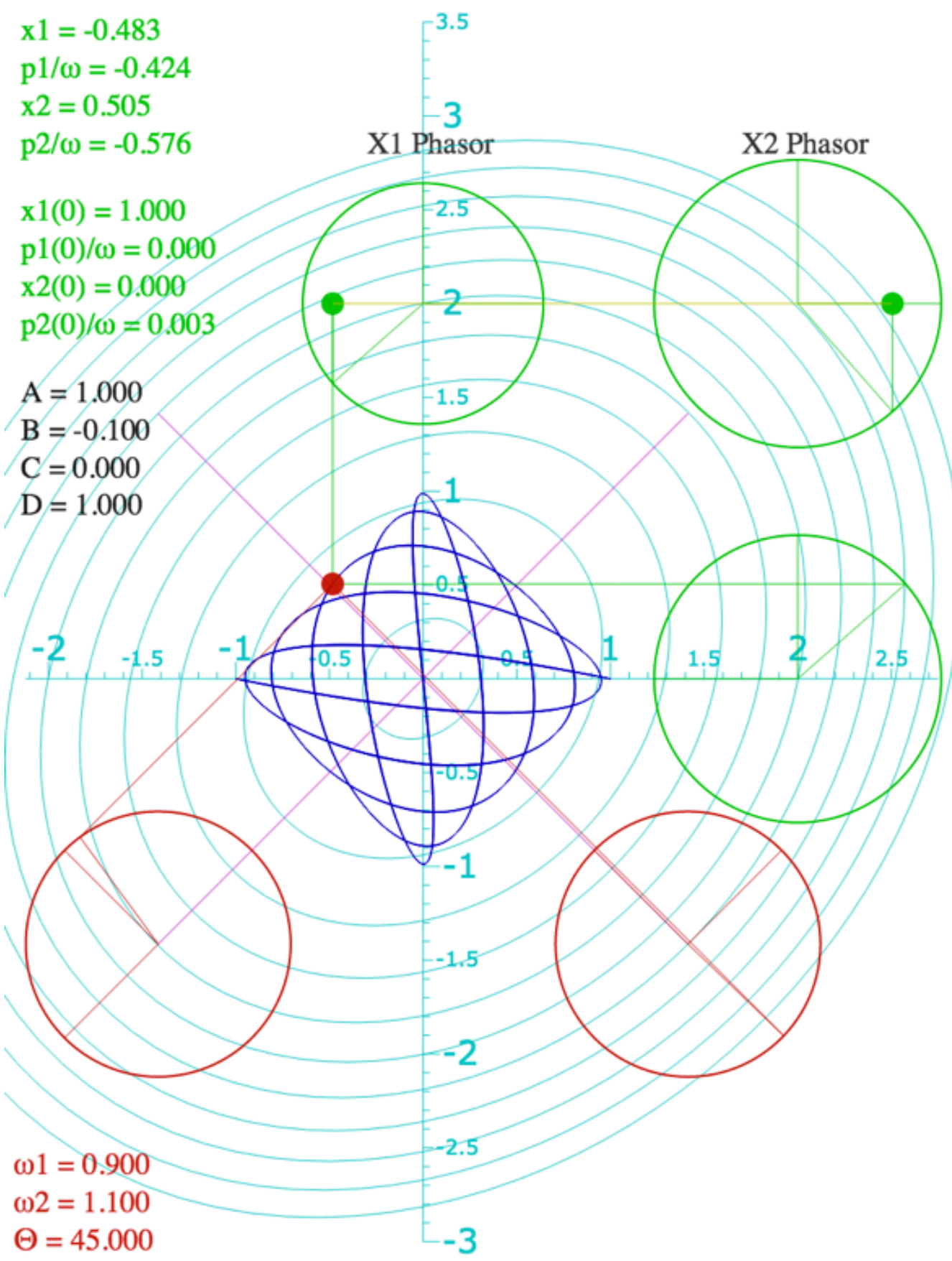
$\omega1 = 0.900$
 $\omega2 = 1.100$
 $\Theta = 45.000$

$x_1 = -0.483$
 $p_1/\omega = -0.424$
 $x_2 = 0.505$
 $p_2/\omega = -0.576$

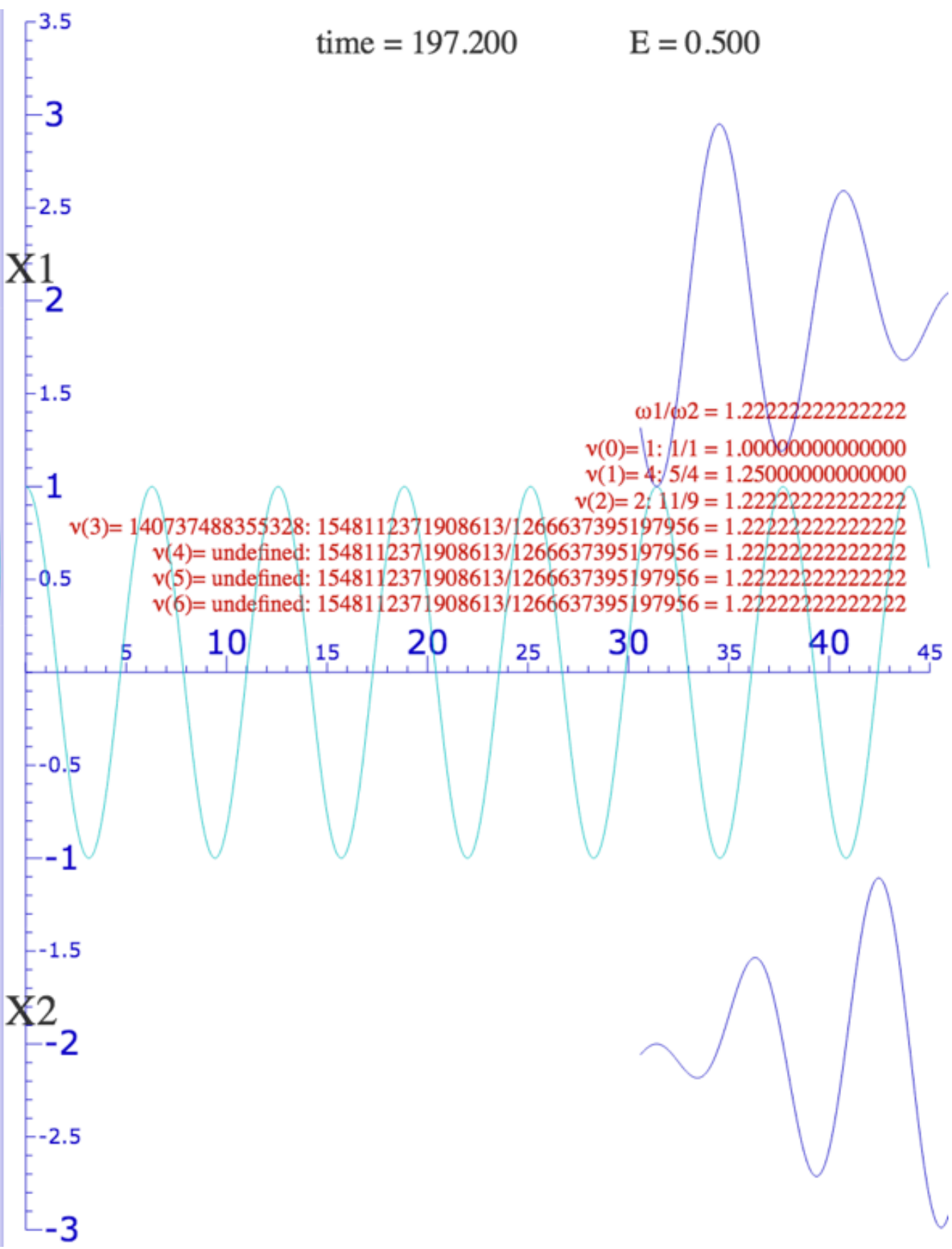
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 $x_2(0) = 0.000$
 $p_2(0)/\omega = 0.003$

$A = 1.000$
 $B = -0.100$
 $C = 0.000$
 $D = 1.000$

$\omega_1 = 0.900$
 $\omega_2 = 1.100$
 $\Theta = 45.000$



time = 197.200 E = 0.500



$\omega_1/\omega_2 = 1.22222222222222$
 $v(0) = 1: 1/1 = 1.00000000000000$
 $v(1) = 4: 5/4 = 1.25000000000000$
 $v(2) = 2: 11/9 = 1.22222222222222$
 $v(3) = 140737488355328: 1548112371908613/1266637395197956 = 1.22222222222222$
 $v(4) = \text{undefined}: 1548112371908613/1266637395197956 = 1.22222222222222$
 $v(5) = \text{undefined}: 1548112371908613/1266637395197956 = 1.22222222222222$
 $v(6) = \text{undefined}: 1548112371908613/1266637395197956 = 1.22222222222222$

BoxIt (Beating) Web Simulation ($A=1, B=-0.1, C=0, D=1$) with frequency ratios

2D harmonic oscillator equations

Lagrangian and matrix forms and Reciprocity symmetry

2D harmonic oscillator equation eigensolutions

Geometric method

Matrix-algebraic eigensolutions with example $M = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix}$

Secular equation

Hamilton-Cayley equation and projectors

Idempotent projectors (how eigenvalues \Rightarrow eigenvectors)

Operator orthonormality and Completeness (Idempotent means: $\mathbf{P} \cdot \mathbf{P} = \mathbf{P}$)

Spectral Decompositions

Functional spectral decomposition

Orthonormality vs. Completeness vis-a`-vis Operator vs. State

Lagrange functional interpolation formula

Diagonalizing Transformations (D-Ttran) from projectors

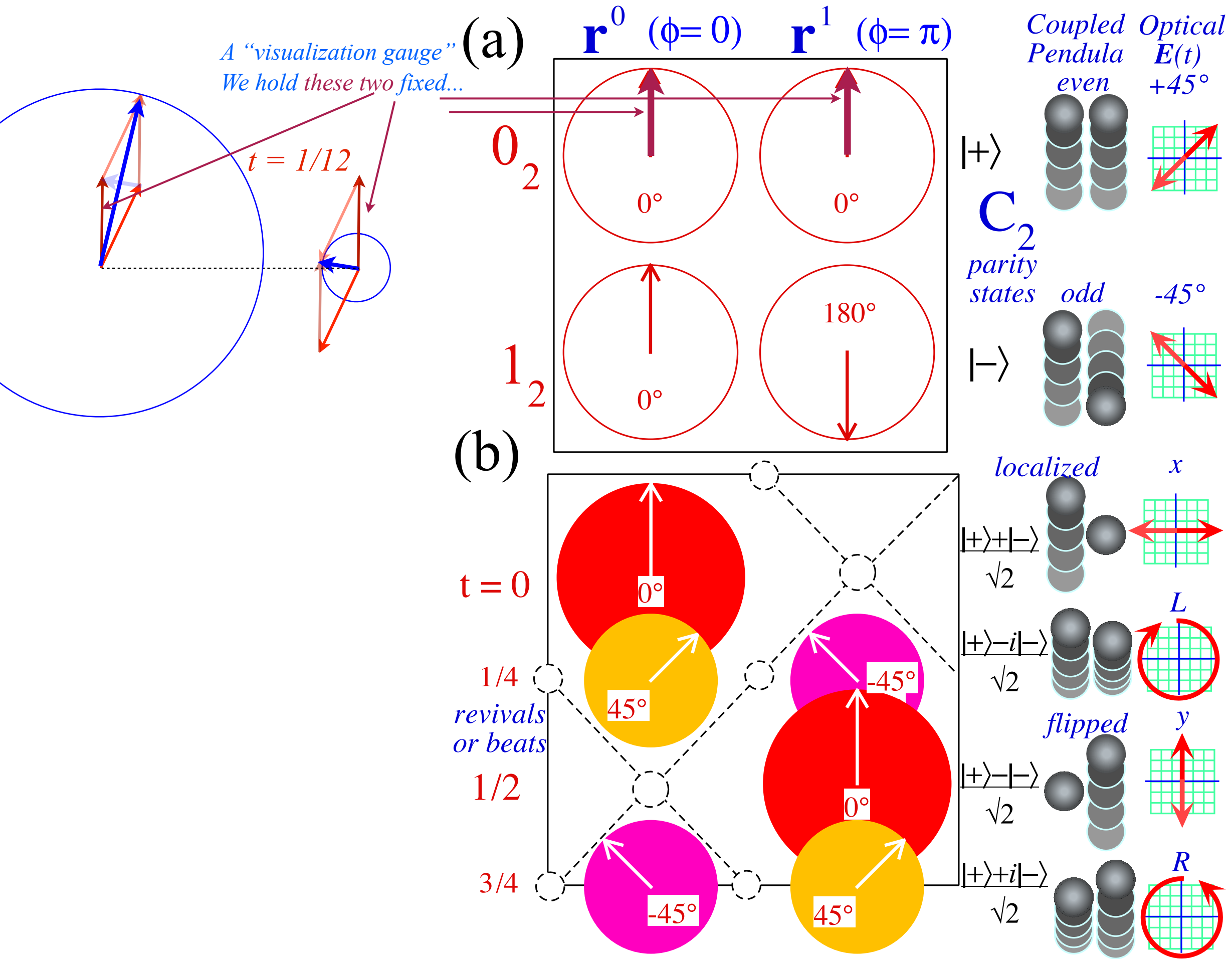
2D-HO eigensolution example with bilateral (B-Type) symmetry

➔ Mixed mode beat dynamics and fixed $\pi/2$ phase ←

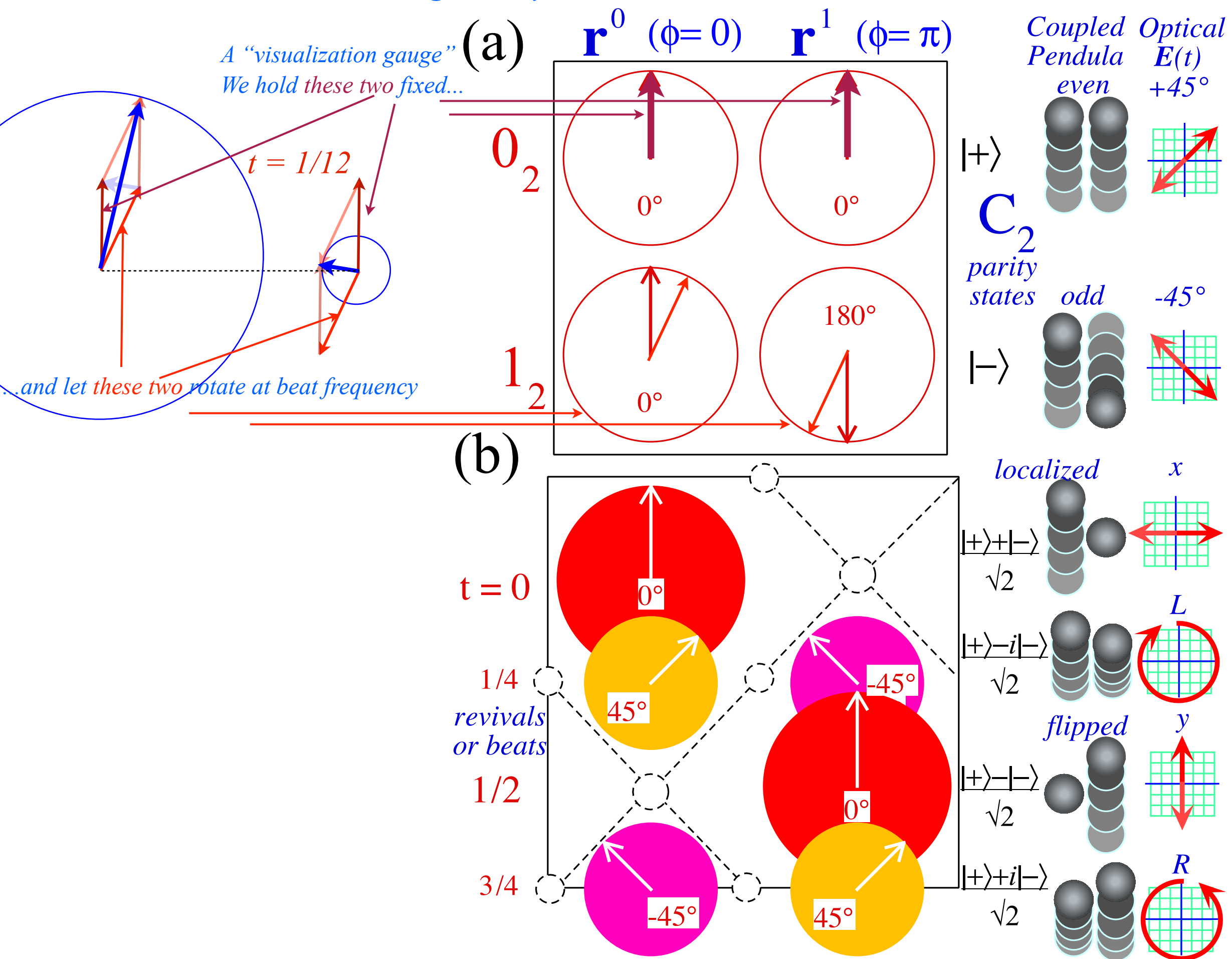
2D-HO eigensolution example with asymmetric (A-Type) symmetry

Initial state projection, mixed mode beat dynamics with variable phase

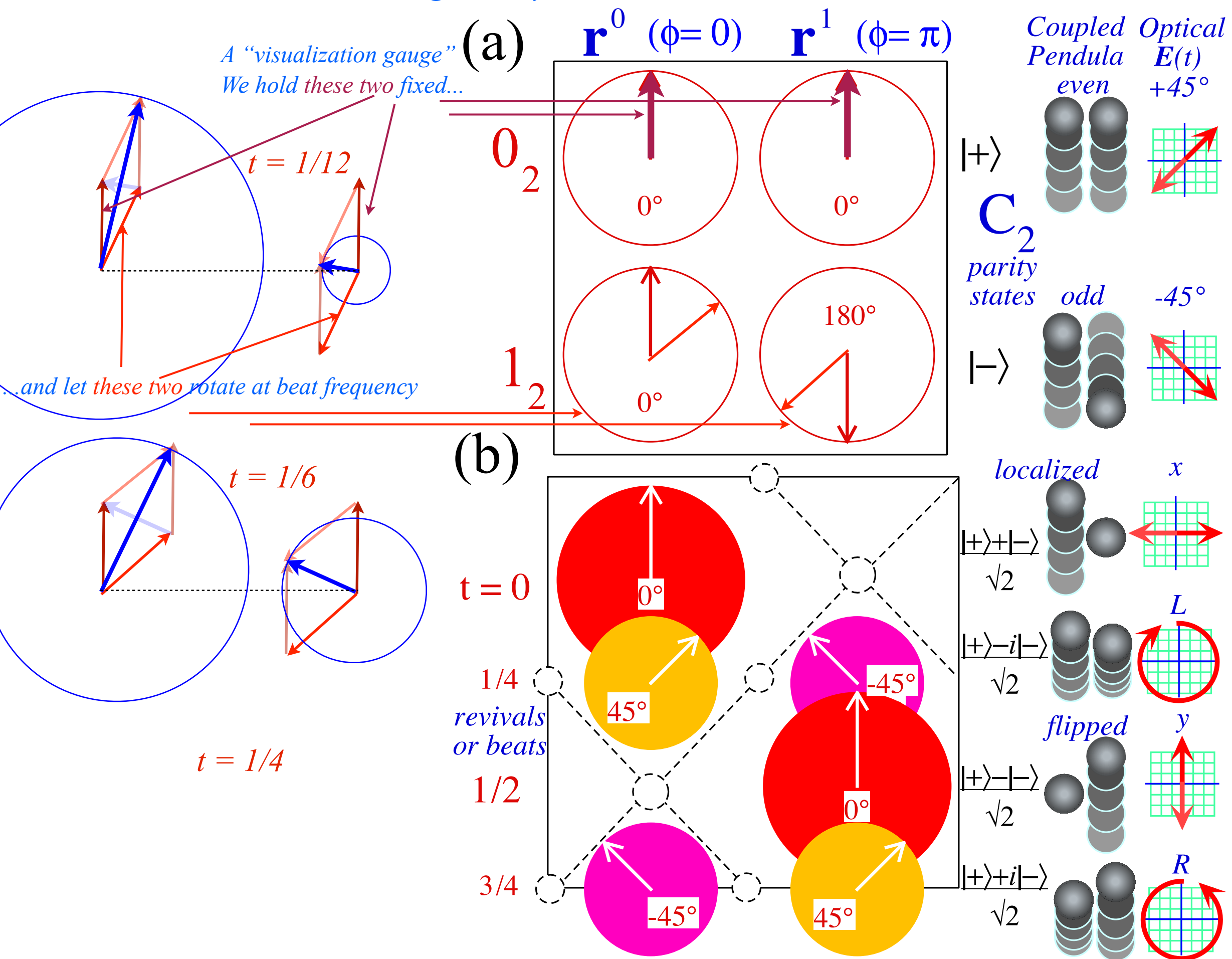
2D-HO beats and mixed mode geometry



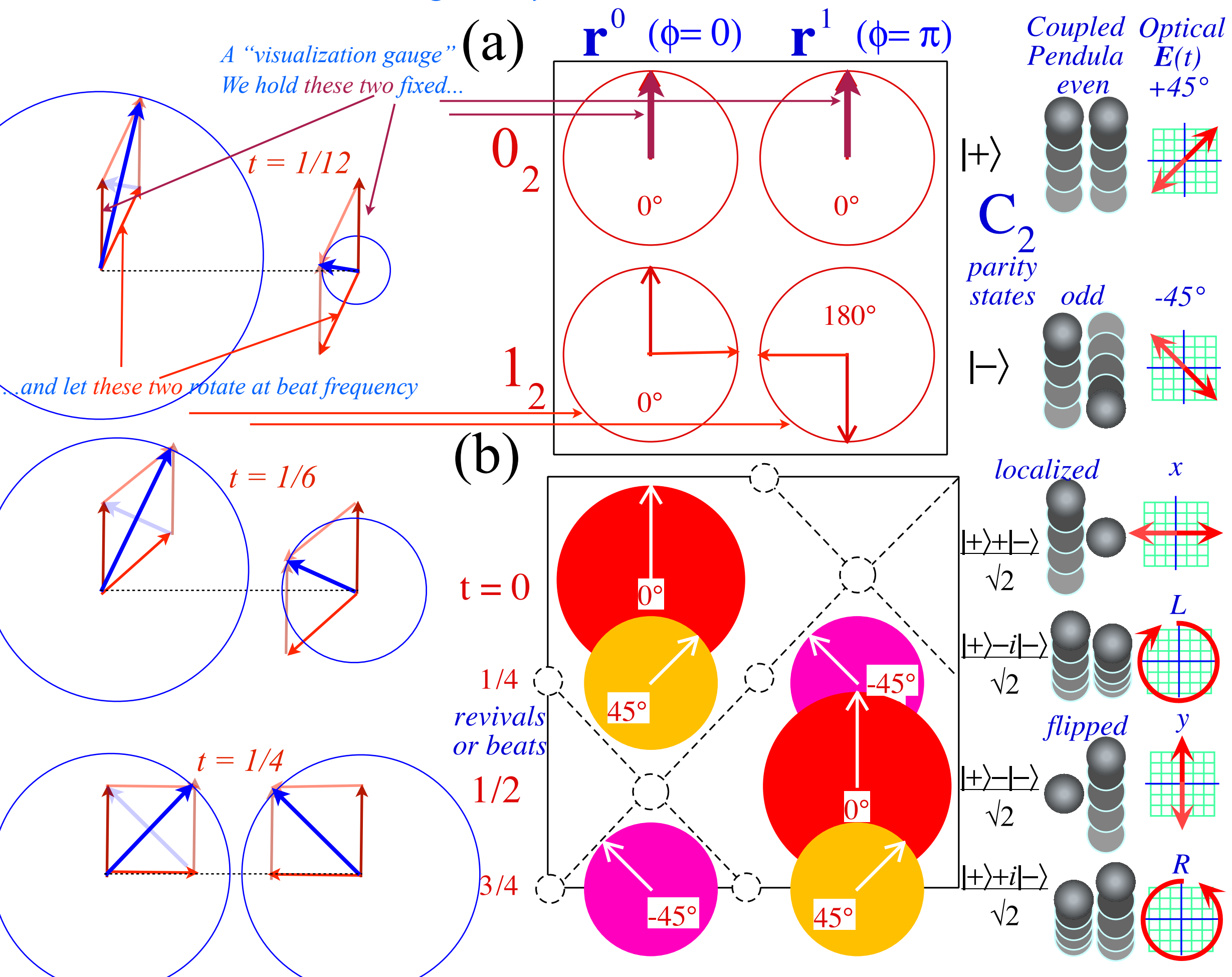
2D-HO beats and mixed mode geometry



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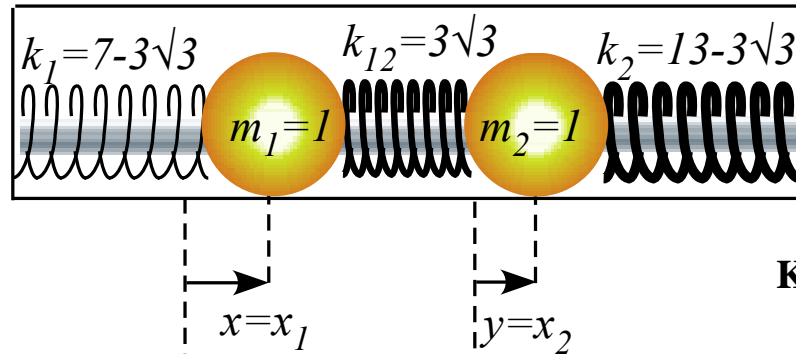
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Mixed mode beat dynamics and fixed $\pi/2$ phase

➔ *2D-HO eigensolution example with asymmetric (A-Type) symmetry* **←**
Initial state projection, mixed mode beat dynamics with variable phase

Spectral decomposition of 2D-HO mode dynamics for lower symmetry

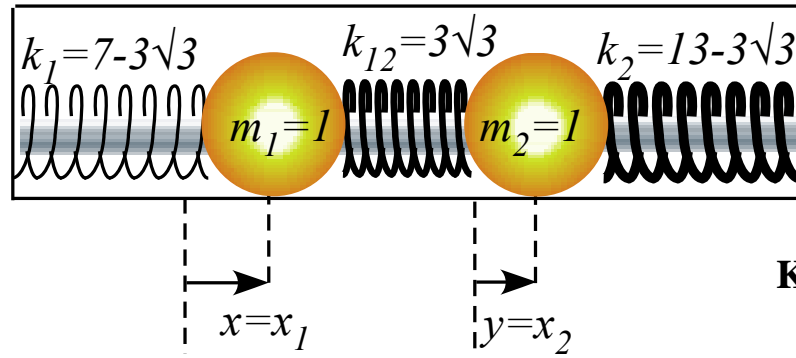


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$$\text{Det}(\mathbf{K}) = 7 \cdot 13 - 27 = 91 - 27 = 64$$

$$\text{Trace}(\mathbf{K}) = 7 + 13 = 20$$

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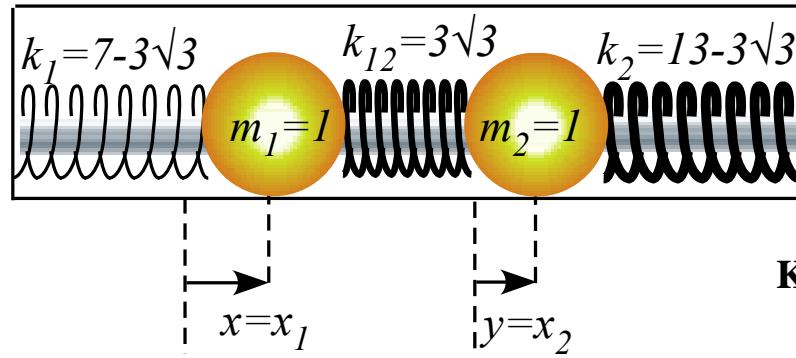
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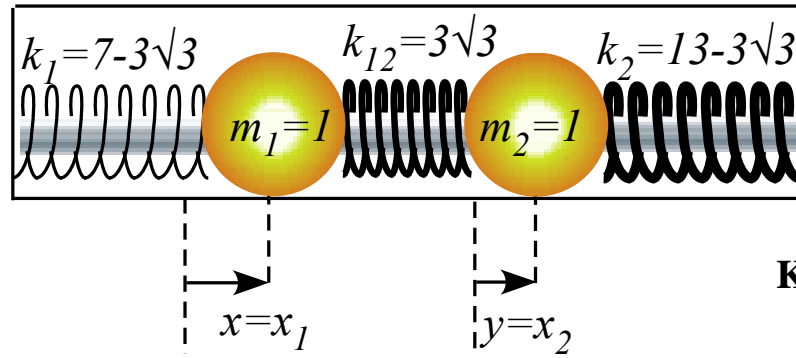


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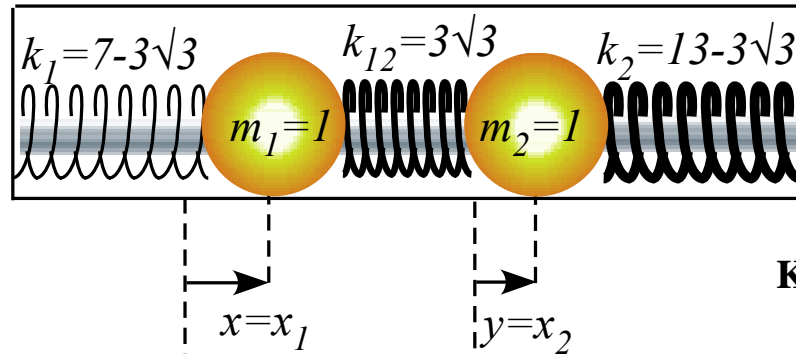


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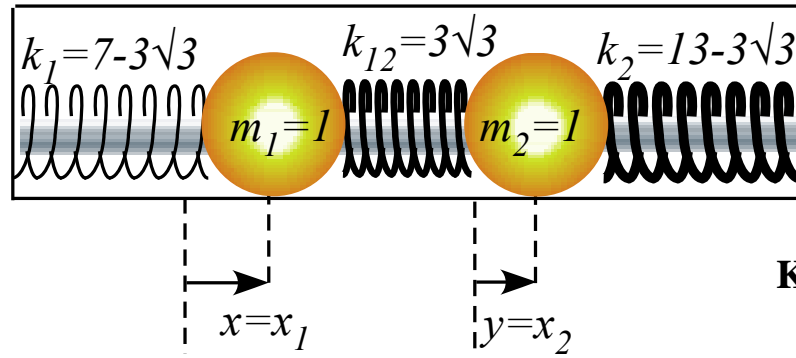
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$$\begin{aligned} \mathbf{P}_1 &= \frac{\begin{pmatrix} K_{11} - K_2 & K_{12} \\ K_{12} & K_{22} - K_2 \end{pmatrix}}{K_1 - K_2} = \frac{\begin{pmatrix} 7 - 16 & -3\sqrt{3} \\ -3\sqrt{3} & 13 - 16 \end{pmatrix}}{4 - 16} = \frac{\begin{pmatrix} 9 & +3\sqrt{3} \\ +3\sqrt{3} & 3 \end{pmatrix}}{12} \\ &= \frac{\begin{pmatrix} 3 & \sqrt{3} \\ \sqrt{3} & 1 \end{pmatrix}}{4} = \begin{pmatrix} \sqrt{3}/2 \\ 1/2 \end{pmatrix} \begin{pmatrix} \sqrt{3}/2 & 1/2 \end{pmatrix} = |\epsilon_1\rangle\langle\epsilon_1| \end{aligned}$$

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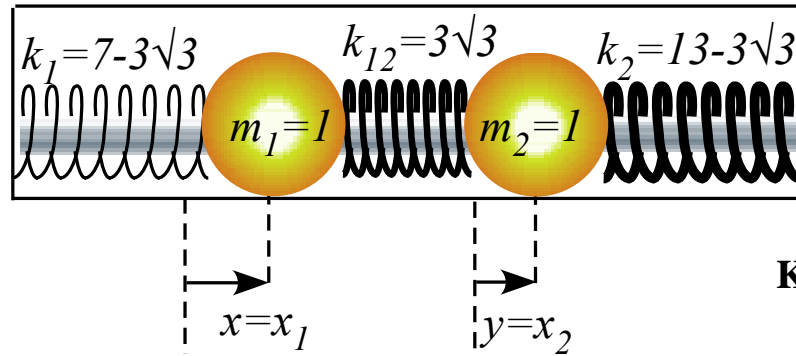
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2D harmonic oscillator equations

Lagrangian and matrix forms and Reciprocity symmetry

2D harmonic oscillator equation eigensolutions

Geometric method

Matrix-algebraic eigensolutions with example $M = \begin{pmatrix} 4 & 1 \\ 3 & 2 \end{pmatrix}$

Secular equation

Hamilton-Cayley equation and projectors

Idempotent projectors (how eigenvalues \Rightarrow eigenvectors)

Operator orthonormality and Completeness (Idempotent means: $P \cdot P = P$)

Spectral Decompositions

Functional spectral decomposition

Orthonormality vs. Completeness vis-a`-vis Operator vs. State

Lagrange functional interpolation formula

Diagonalizing Transformations (D-Ttran) from projectors

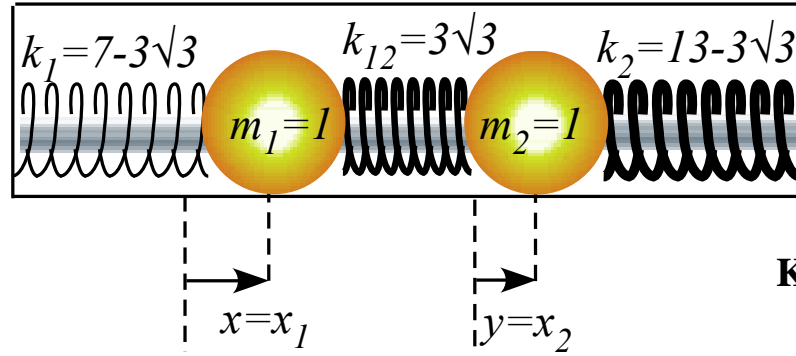
2D-HO eigensolution example with bilateral (B-Type) symmetry

Mixed mode beat dynamics and fixed $\pi/2$ phase

2D-HO eigensolution example with asymmetric (A-Type) symmetry

➔ Initial state projection, mixed mode beat dynamics with variable phase ←

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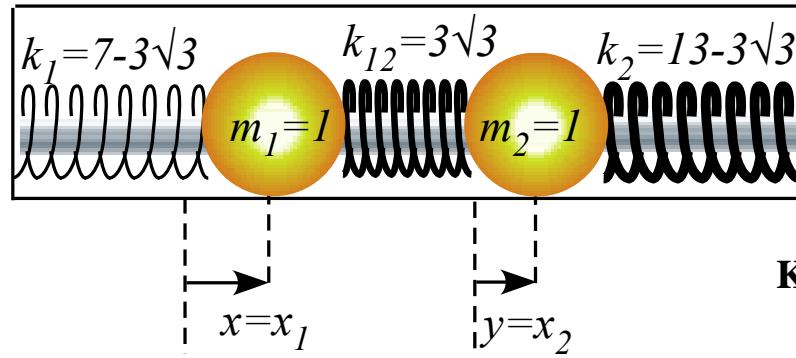
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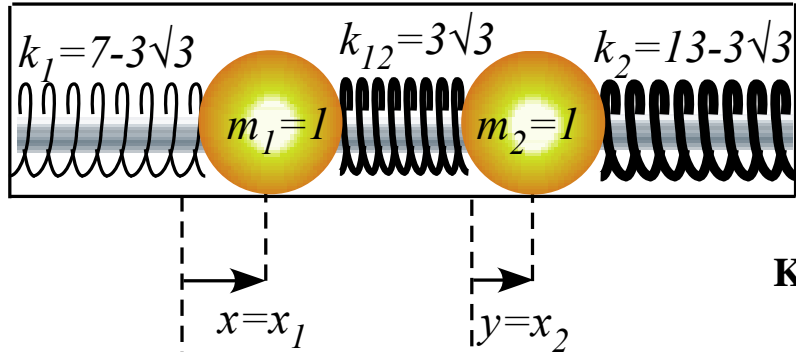
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(Note projection onto eigen-axes)

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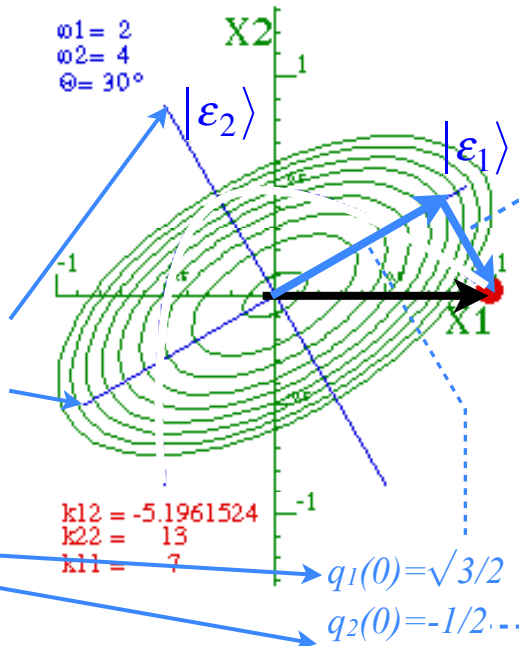
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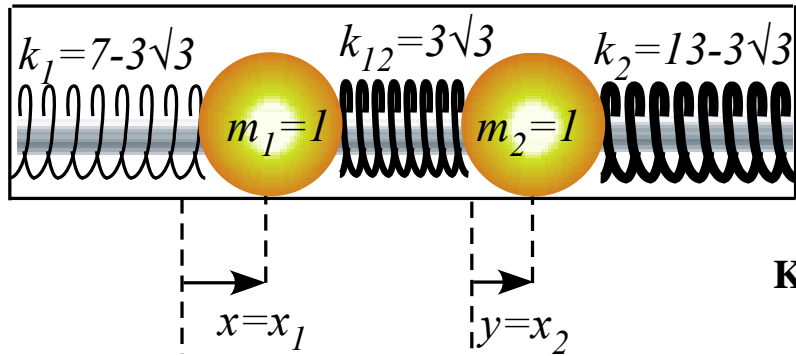
$$\begin{pmatrix} q_1(t) = \frac{\sqrt{3}}{2} \cos 2t, & q_2(t) = -\frac{1}{2} \cos 4t \end{pmatrix}$$

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Eigenbra vectors: $\langle\epsilon_1| = \begin{pmatrix} \sqrt{3}/2 & 1/2 \end{pmatrix}, \quad \langle\epsilon_2| = \begin{pmatrix} -1/2 & \sqrt{3}/2 \end{pmatrix}$

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$$\mathbf{1} \cdot \mathbf{x}(0) = (\mathbf{P}_1 + \mathbf{P}_2) \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} \sqrt{3}/2 \\ 1/2 \end{pmatrix} \otimes \begin{pmatrix} \sqrt{3}/2 & 1/2 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} + \begin{pmatrix} -1/2 \\ \sqrt{3}/2 \end{pmatrix} \otimes \begin{pmatrix} -1/2 & \sqrt{3}/2 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$= \begin{pmatrix} \sqrt{3}/2 \\ 1/2 \end{pmatrix} \left(\frac{\sqrt{3}}{2}\right) + \begin{pmatrix} -1/2 \\ \sqrt{3}/2 \end{pmatrix} \left(-\frac{1}{2}\right)$$

$$\begin{pmatrix} q_1(t) = \frac{\sqrt{3}}{2} \cos 2t, & q_2(t) = -\frac{1}{2} \cos 4t \end{pmatrix}$$

Using $\cos 4t = 2 \cos^2 2t - 1$ derives a parabolic trajectory!

$$q_2(t) = -\frac{1}{2} 2 \cos^2 2t + \frac{1}{2} = -\frac{4}{3} [q_1(t)]^2 + \frac{1}{2}$$

$$\mathbf{x}(0) = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

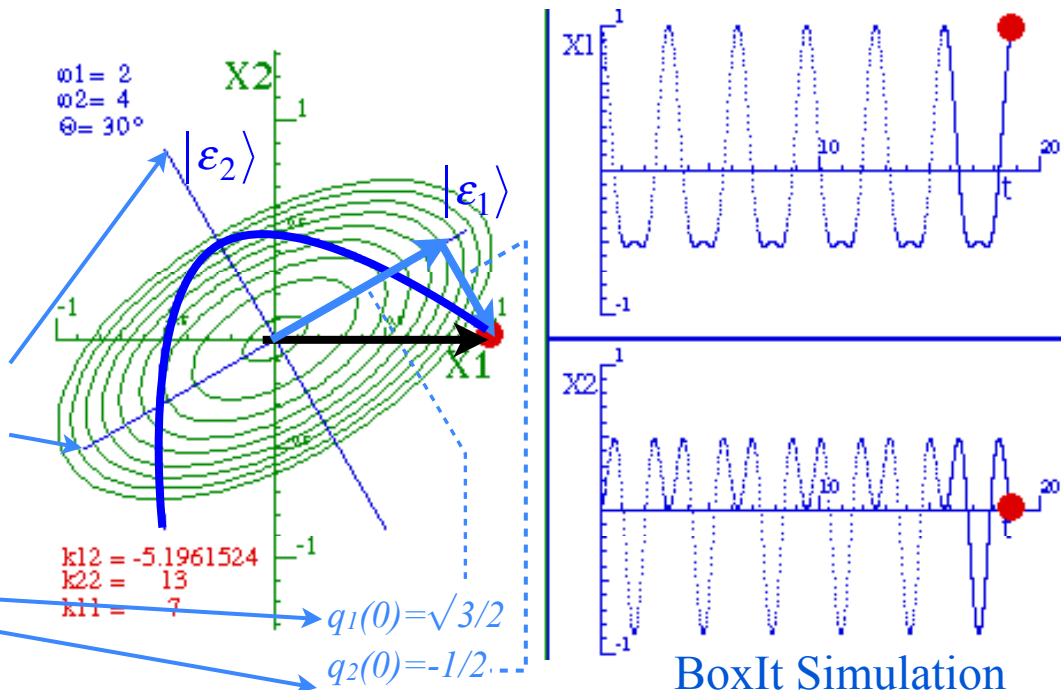
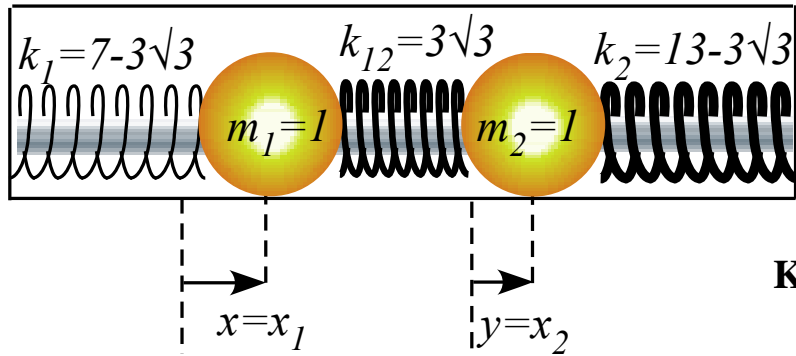


Fig. 3.3.6 Normal coordinate axes, coupled oscillator trajectories and equipotential ($V=\text{const.}$) ovals for an integral 1:2 eigenfrequency ratio ($\omega_0(\epsilon_1)=2.0, \omega_0(\epsilon_2)=4.0$) and zero initial velocity.

Spectral decomposition of 2D-HO mode dynamics for lower symmetry



$$\mathbf{K} = \begin{pmatrix} K_{11} & K_{12} \\ K_{12} & K_{22} \end{pmatrix} = \begin{pmatrix} k_1 + k_{12} & -k_{12} \\ -k_{12} & k_2 + k_{12} \end{pmatrix} = \begin{pmatrix} 7 & -3\sqrt{3} \\ -3\sqrt{3} & 13 \end{pmatrix}$$

The \mathbf{K} secular equation $K^2 - \text{Trace}(\mathbf{K})K + \text{Det}(\mathbf{K}) = K^2 - 20K + 64 = 0 = (K - 4)(K - 16)$

$\text{Det}(\mathbf{K}) = 7 \cdot 13 - 27 = 91 - 27 = 64$
 $\text{Trace}(\mathbf{K}) = 7 + 13 = 20$

Eigenvalues K_k and squared eigenfrequencies $\omega_0(\epsilon_k)^2$ $K_1 = \omega_0^2(\epsilon_1) = 4, \quad K_2 = \omega_0^2(\epsilon_2) = 16,$

Eigen-projectors \mathbf{P}_k

$$\mathbf{P}_1 = \frac{\begin{pmatrix} K_{11} - K_2 & K_{12} \\ K_{12} & K_{22} - K_2 \end{pmatrix}}{K_1 - K_2} = \frac{\begin{pmatrix} 7 - 16 & -3\sqrt{3} \\ -3\sqrt{3} & 13 - 16 \end{pmatrix}}{4 - 16} = \frac{\begin{pmatrix} 9 & +3\sqrt{3} \\ +3\sqrt{3} & 3 \end{pmatrix}}{12}$$

$$= \frac{\begin{pmatrix} 3 & \sqrt{3} \\ \sqrt{3} & 1 \end{pmatrix}}{4} = \begin{pmatrix} \sqrt{3}/2 \\ 1/2 \end{pmatrix} \begin{pmatrix} \sqrt{3}/2 & 1/2 \end{pmatrix} = |\epsilon_1\rangle\langle\epsilon_1|$$

$$\mathbf{P}_2 = \frac{\begin{pmatrix} K_{11} - K_1 & K_{12} \\ K_{12} & K_{22} - K_1 \end{pmatrix}}{K_2 - K_1} = \frac{\begin{pmatrix} 7 - 4 & -3\sqrt{3} \\ -3\sqrt{3} & 13 - 4 \end{pmatrix}}{16 - 4} = \frac{\begin{pmatrix} 3 & -3\sqrt{3} \\ -3\sqrt{3} & 9 \end{pmatrix}}{12}$$

$$= \frac{\begin{pmatrix} 1 & -\sqrt{3} \\ -\sqrt{3} & 3 \end{pmatrix}}{4} = \begin{pmatrix} -1/2 \\ \sqrt{3}/2 \end{pmatrix} \begin{pmatrix} -1/2 & \sqrt{3}/2 \end{pmatrix} = |\epsilon_2\rangle\langle\epsilon_2|$$

Eigenbra vectors: $\langle\epsilon_1| = \begin{pmatrix} \sqrt{3}/2 & 1/2 \end{pmatrix}, \quad \langle\epsilon_2| = \begin{pmatrix} -1/2 & \sqrt{3}/2 \end{pmatrix}$

Spectral decomposition of initial state $\mathbf{x}(0) = (1, 0)$:

$$\mathbf{1} \cdot \mathbf{x}(0) = (\mathbf{P}_1 + \mathbf{P}_2) \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} \sqrt{3}/2 \\ 1/2 \end{pmatrix} \otimes \begin{pmatrix} \sqrt{3}/2 & 1/2 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} + \begin{pmatrix} -1/2 \\ \sqrt{3}/2 \end{pmatrix} \otimes \begin{pmatrix} -1/2 & \sqrt{3}/2 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$= \begin{pmatrix} \sqrt{3}/2 \\ 1/2 \end{pmatrix} \begin{pmatrix} \sqrt{3}/2 \\ 1/2 \end{pmatrix} + \begin{pmatrix} -1/2 \\ \sqrt{3}/2 \end{pmatrix} \begin{pmatrix} -1/2 \\ \sqrt{3}/2 \end{pmatrix}$$

(Note projection of $\mathbf{x}(0)$ onto eigen-axes)

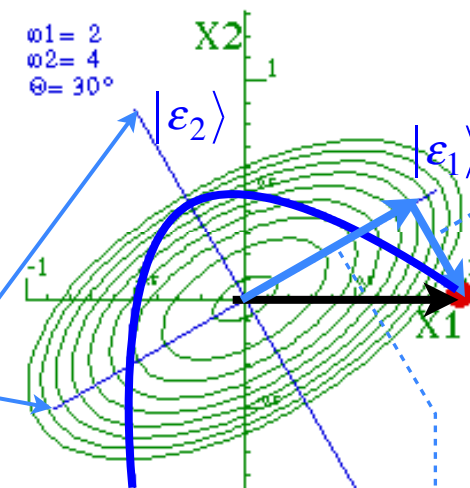
$$\begin{pmatrix} q_1(t) = \frac{\sqrt{3}}{2} \cos 2t, & q_2(t) = -\frac{1}{2} \cos 4t \end{pmatrix}$$

Using $\cos 4t = 2 \cos^2 2t - 1$ derives a parabolic trajectory!

$$q_2(t) = -\frac{1}{2} (2 \cos^2 2t - 1) = -\frac{4}{3} [q_1(t)]^2 + \frac{1}{2}$$

$$\mathbf{x}(0) = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

Example of a Tschebycheff Polynomial order 2



BoxIt Simulation

Pafnuty Chebyshev



Pafnuty Lvovich Chebyshev was a Russian mathematician. His name can be alternatively transliterated as Chebychev, Chebysheff, Chebyshev, Tchebychev or Tchebycheff, or Tschebyshev or Tschebyscheff. Wikipedia

Born: May 16, 1821, Borovsk
 Died: December 8, 1894, Saint Petersburg