Lecture 9 Revised 12.21.2012 9.18.2012

Geometry of Dual Quadratic Forms: Lagrange vs Hamilton (Ch. 11 and Ch. 12 of Unit 1)

Introduction to dual matrix operator geometry

Review of dual IHO elliptic orbits (Lecture 7-8)

Construction by Phasor-pair projection

Construction by Kepler anomaly projection

Operator geometric sequences and eigenvectors

Rescaled description of matrix operator geometry

Vector calculus of tensor operation

Introduction to Lagrangian-Hamiltonian duality

Review of partial differential relations

Chain rule and order symmetry

Duality relations of Lagrangian and Hamiltonian ellipse

Introducing the 1^{st} (partial $\frac{\partial ?}{\partial ?}$) differential equations of mechanics

Introduction to dual matrix operator geometry

Review of dual IHO elliptic orbits (Lecture 7-8)

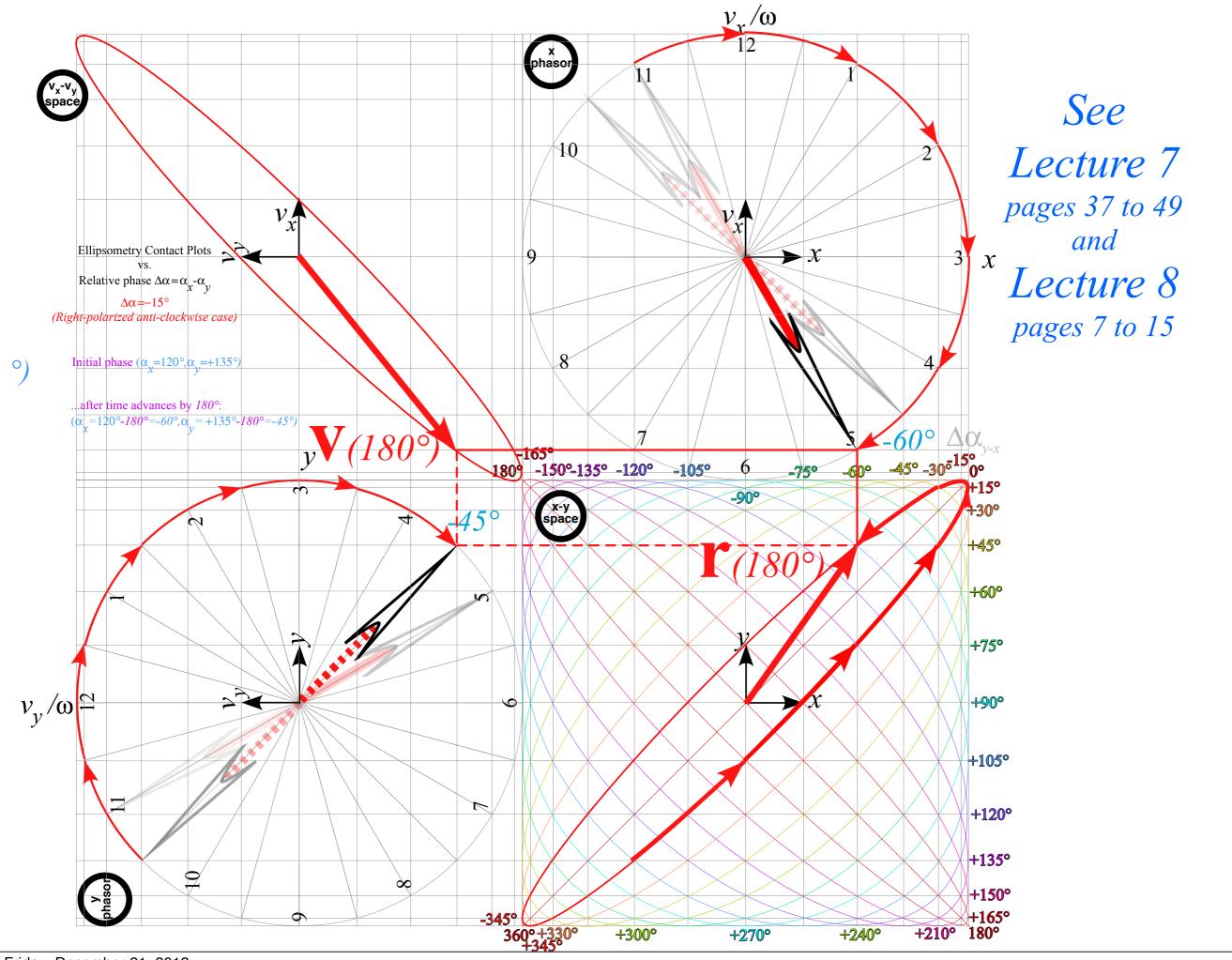
Construction by Phasor-pair projection

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Introduction to dual matrix operator geometry

Review of dual IHO elliptic orbits (Lecture 7-8)

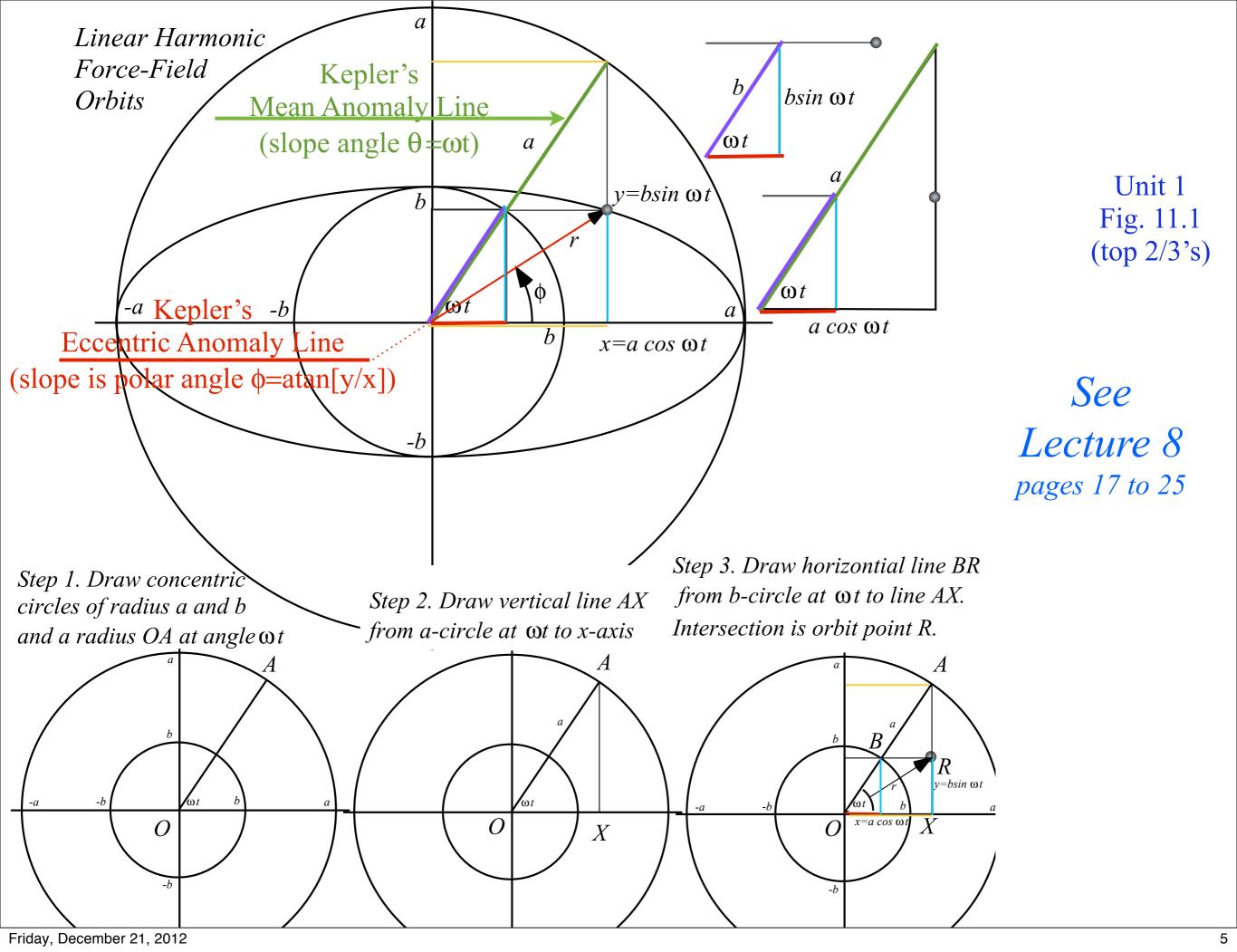
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Quadratic forms and tangent contact geometry of their ellipses

A matrix Q that generates an ellipse by $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ is called positive-definite (if $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}$ always > 0)

$$\begin{pmatrix} x & y \end{pmatrix} \bullet \begin{pmatrix} \frac{1}{a^2} & 0 \\ 0 & \frac{1}{b^2} \end{pmatrix} \bullet \begin{pmatrix} x \\ y \end{pmatrix} = 1 = \begin{pmatrix} x & y \end{pmatrix} \bullet \begin{pmatrix} \frac{x}{a^2} \\ \frac{y}{b^2} \end{pmatrix} = \frac{x^2}{a^2} + \frac{y^2}{b^2}$$

A inverse matrix Q^{-1} generates an ellipse by $\mathbf{p} \cdot Q^{-1} \cdot \mathbf{p} = 1$ called inverse or dual ellipse:

$$\mathbf{p} \bullet \mathbf{Q}^{-1} \bullet \mathbf{p} = 1$$

$$\left(\begin{array}{ccc} p_x & p_y \\ 0 & b^2 \end{array} \right) \bullet \left(\begin{array}{ccc} a^2 & 0 \\ p_y \end{array} \right) = 1 = \left(\begin{array}{ccc} p_x & p_y \\ p_y \end{array} \right) \bullet \left(\begin{array}{ccc} a^2 p_x \\ b^2 p_y \end{array} \right) = a^2 p_x^2 + b^2 p_y^2$$

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Quadratic forms and tangent contact geometry of their ellipses

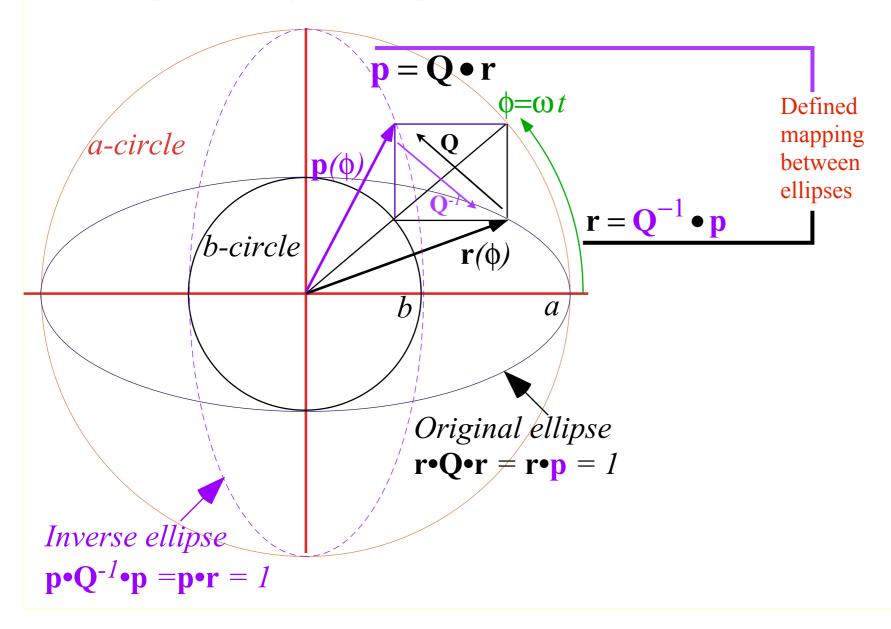
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Defined mapping between ellipses

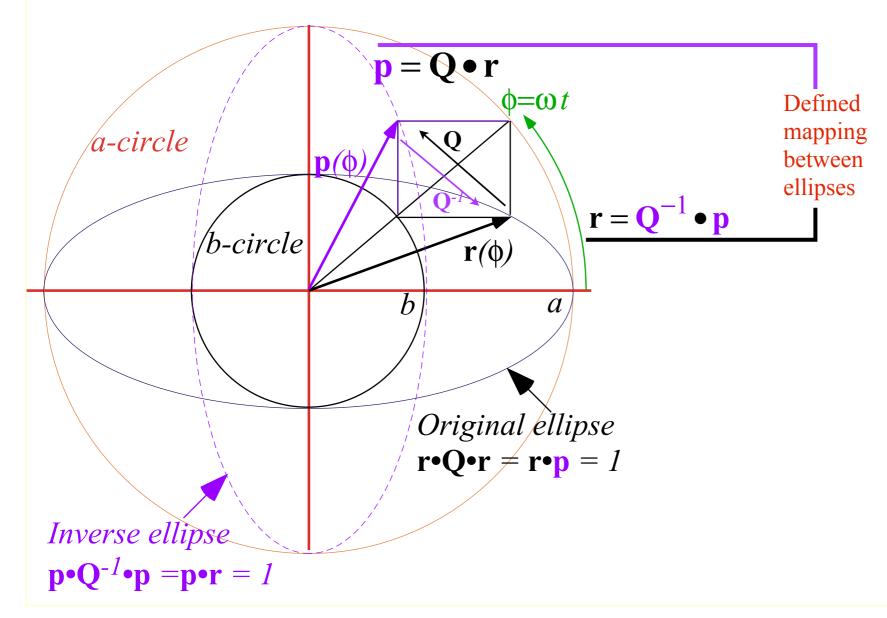
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based on Unit 1 Fig. 11.6

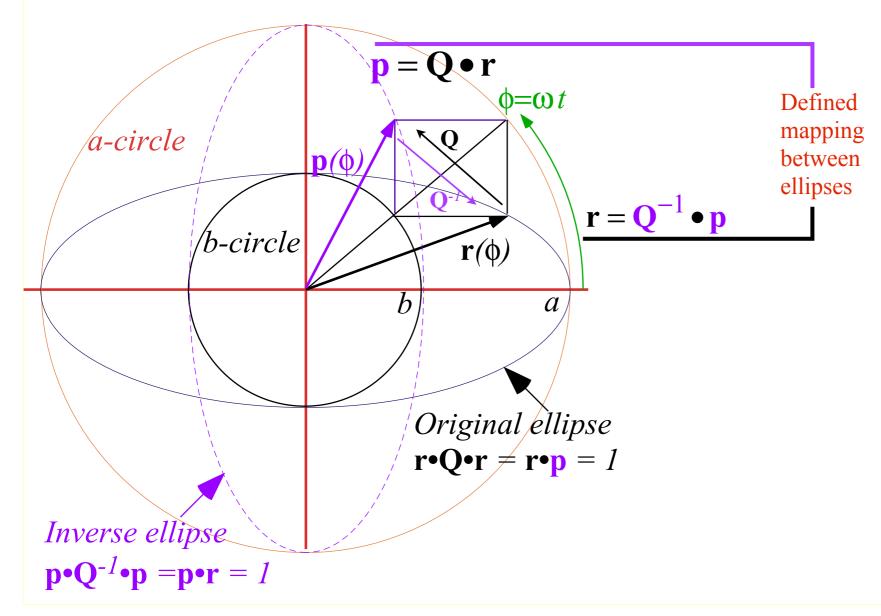


based on Unit 1 Fig. 11.6



Quadratic form $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ has mutual duality relations with inverse form $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1 = \mathbf{p} \cdot \mathbf{r}$

based on Unit 1 Fig. 11.6

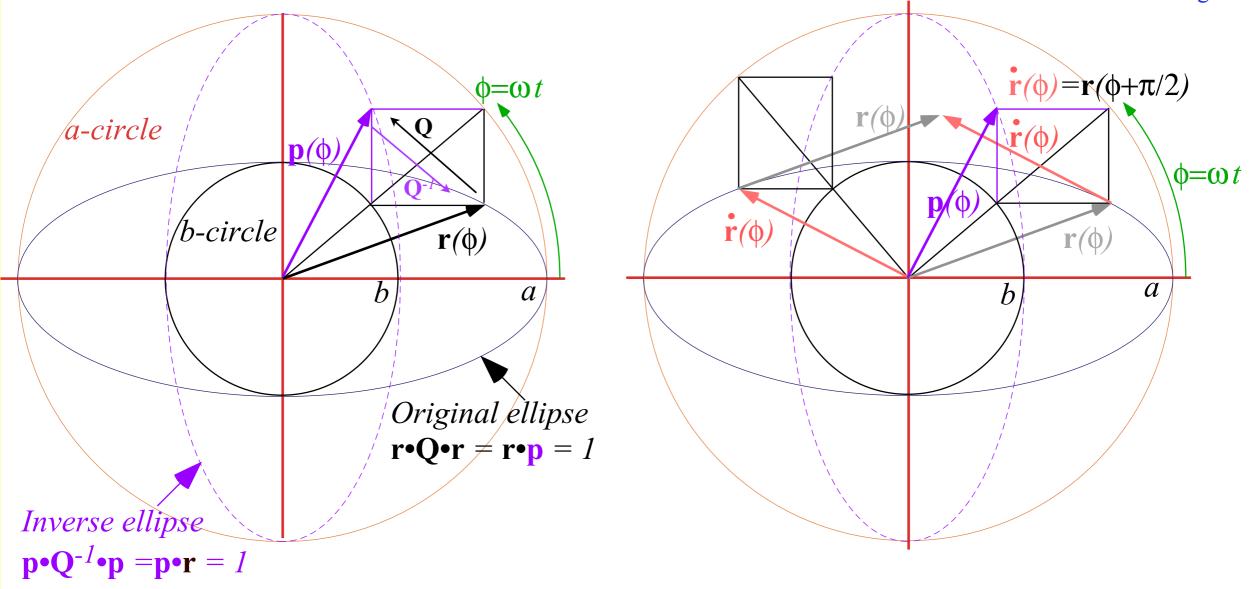


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$$\mathbf{p} = \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \bullet \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix} = \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \text{ where: } \begin{cases} x = r_x = a\cos\phi = a\cos\omega t \\ y = r_y = b\sin\phi = b\sin\omega t \end{cases} \text{ so: } \mathbf{p} \cdot \mathbf{r} = 1$$

(b) Ellipse tangents

based on Unit 1 Fig. 11.6

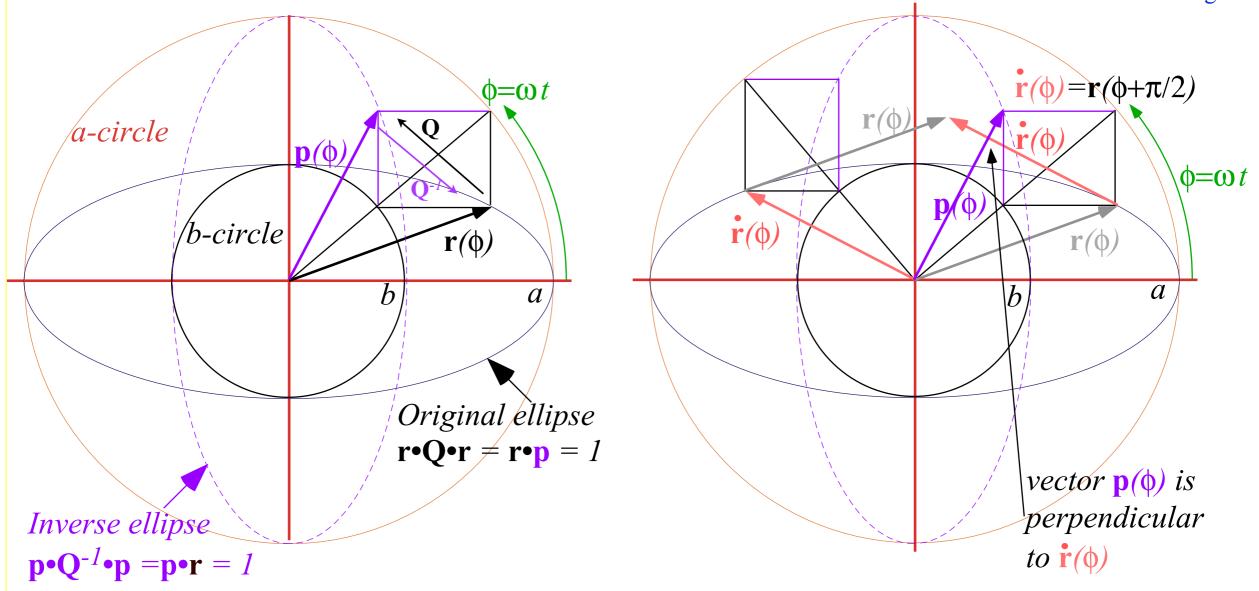


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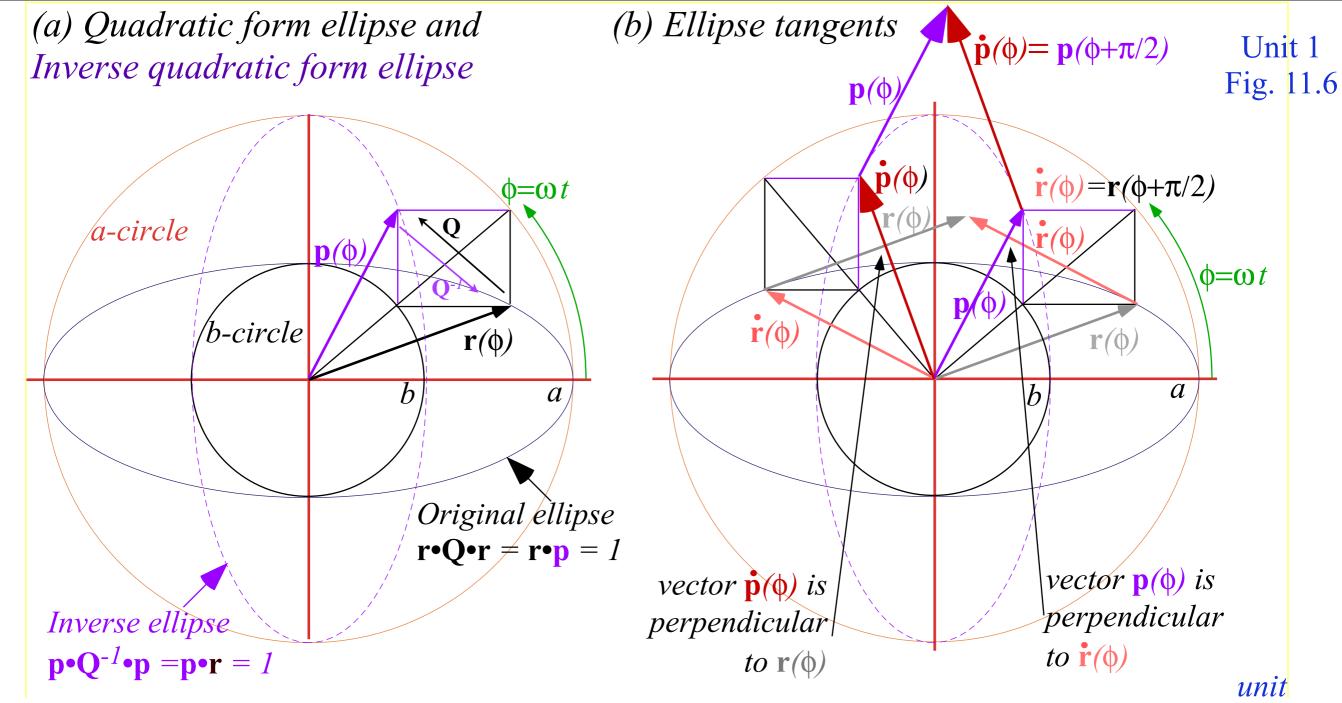


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 ${f p}$ is perpendicular to velocity ${f v}={f \dot r}$, a mutual orthogonality

$$\begin{vmatrix}
\dot{\mathbf{r}} \bullet \mathbf{p} = 0 \\
\dot{\mathbf{r}}_{x} & \dot{r}_{y}
\end{vmatrix} \bullet \begin{pmatrix}
p_{x} \\
p_{y}
\end{pmatrix} = \begin{pmatrix}
-a\sin\phi & b\cos\phi
\end{pmatrix} \bullet \begin{pmatrix}
(1/a)\cos\phi \\
(1/b)\sin\phi
\end{pmatrix}$$
where:
$$\begin{vmatrix}
\dot{r}_{x} = -a\sin\phi \\
\dot{r}_{y} = b\cos\phi$$
and:
$$\begin{aligned}
p_{x} = (1/a)\cos\phi \\
p_{y} = (1/b)\sin\phi
\end{aligned}$$



Quadratic form $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$ has mutual duality relations with inverse form $\mathbf{p} \cdot \mathbf{Q}^{-1} \cdot \mathbf{p} = 1$

mutual projection

$$\mathbf{p} = \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \bullet \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix} = \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \text{ where: } \begin{cases} x = r_x = a\cos\phi = a\cos\omega t \\ y = r_y = b\sin\omega t \end{cases}$$

so: $\mathbf{p} \cdot \mathbf{r} = 1$

 ${\bf p}$ is perpendicular to velocity ${\bf v}={\bf \dot r}$, a mutual orthogonality. So is ${\bf r}$ perpendicular to ${\bf \dot p}$:

$$\begin{vmatrix}
\dot{\mathbf{r}} \cdot \mathbf{p} = 0 \\
\dot{\mathbf{r}}_{x} & \dot{r}_{y}
\end{vmatrix} = \begin{pmatrix} -a\sin\phi & b\cos\phi \\ p_{y} \end{pmatrix} = \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \text{ where: } \begin{vmatrix} \dot{r}_{x} = -a\sin\phi \\ \dot{r}_{y} = b\cos\phi \end{vmatrix} \text{ and: } \begin{aligned} p_{x} = (1/a)\cos\phi \\ p_{y} = (1/b)\sin\phi \end{aligned}$$

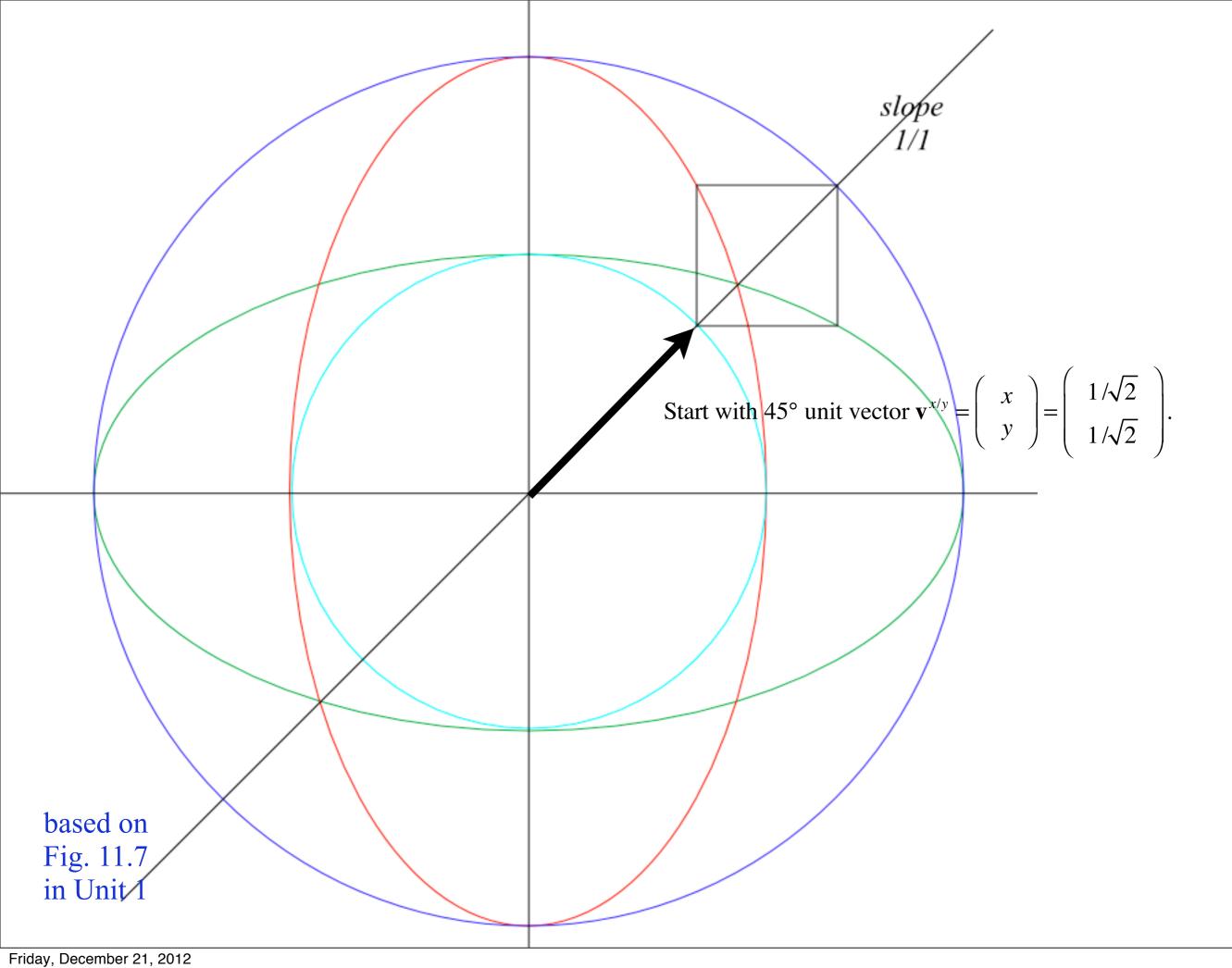
Introduction to dual matrix operator geometry

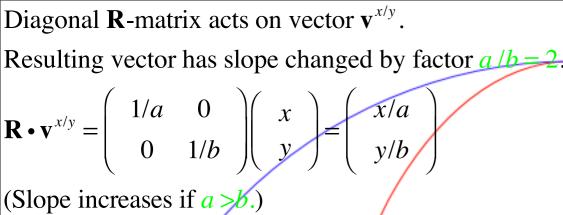
Review of dual IHO elliptic orbits (Lecture 7-8)

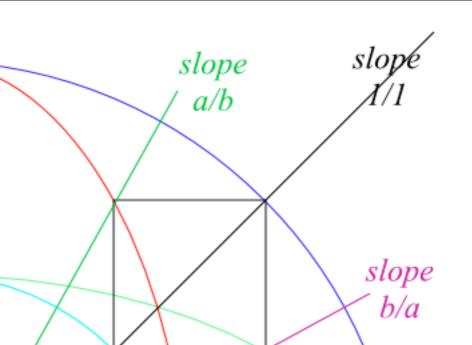
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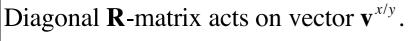




Diagonal \mathbf{R}^{-1} -matrix acts on vector $\mathbf{v}^{x/y}$. Resulting vector has slope changed by factor b/a.

$$\mathbf{R}^{-1} \bullet \mathbf{v}^{x/y} = \begin{pmatrix} a & 0 \\ 0 & b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x \cdot a \\ y \cdot b \end{pmatrix}$$

(Slope decreases if b < a.)



Resulting vector has slope changed by factor a/b = 2.

$$\mathbf{R} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a & 0 \\ 0 & 1/b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a \\ y/b \end{pmatrix}$$

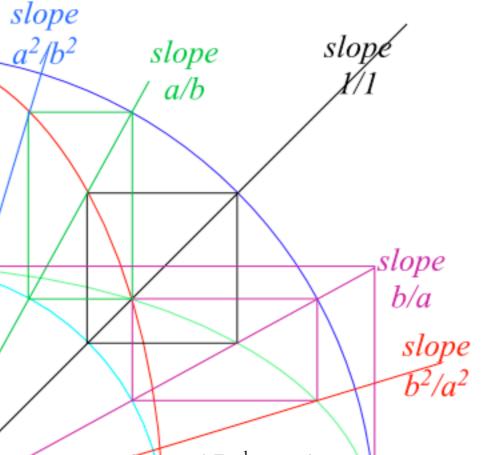
(It increases if a > b.)

Diagonal ($\mathbb{R}^2 = \mathbb{Q}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a^2/b^2 = 4$.

$$\mathbf{Q} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix}$$

(It increases if a > b.)



Diagonal \mathbf{R}^{-1} -matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor b/a=1/2.

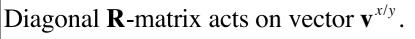
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Diagonal ($\mathbf{R}^{-2} = \mathbf{Q}^{-1}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $b^2/a^2=1/4$.

$$\mathbf{Q}^{-1} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} a^2 & 0 \\ 0 & b^2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x \cdot a^2 \\ y \cdot b^2 \end{pmatrix}$$

based on Fig. 11.7 in Unit 1



Resulting vector has slope changed by factor a/b = 2. $a^3/b^3 a^2/b^2$

$$\mathbf{R} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a & 0 \\ 0 & 1/b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a \\ y/b \end{pmatrix}$$

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Diagonal ($\mathbb{R}^2 = \mathbb{Q}$)-matrix acts on vector $\mathbf{v}^{x/y}$

Resulting vector has slope changed by factor $a^2/b^2 = 4$.

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(It increases if a > b.)

Either process can go on forever...

Diagonal ($\mathbb{R}^{2n} = \mathbb{Q}^n$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a^{2n}/b^{2n} \neq 4^n$.

Either process can go on forever...

slopeslope

slope

/a/b

Diagonal ($\mathbf{R}^{-2n} = \mathbf{Q}^{-n}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

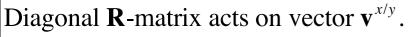
Resulting vector has slope changed by factor $b^{2n}/a^{2n} = 4^{-n}$.

slope

b/a

 $slope \\ b^2/a^2 \\ slope \\ b^3/a^3$

based on Fig. 11.7 in Unit 1



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(It increases if a > b.)

 $|y\rangle$

slopeslope

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Resulting vector has slope changed by factor $a^{2n}/b^{2n} = 4^n$.

...Finally, the result approaches *EIGENVECTOR* $|y\rangle = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$

of ∞ -slope which is "immune" to **R**, **Q** or **Q**ⁿ:

$$\mathbf{R}|y\rangle = (1/b)|y\rangle \qquad \mathbf{Q}^n|y\rangle = (1/b^2)^n|y\rangle$$

Either process can go on forever...

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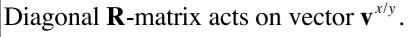
slope

...Finally, the result approaches *EIGENVECTOR* $|x\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$

of 0-slope which is "immune" to \mathbf{R}^{-1} , \mathbf{Q}^{-1} or \mathbf{Q}^{-n} :

$$\mathbf{R}^{-1}|\mathbf{x}\rangle = (\mathbf{a})|\mathbf{x}\rangle \qquad \mathbf{Q}^{-n}|\mathbf{x}\rangle = (\mathbf{a}^2)^n|\mathbf{x}\rangle$$

slope



Resulting vector has slope changed by factor a/b = 2.

$$\mathbf{R} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a & 0 \\ 0 & 1/b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a \\ y/b \end{pmatrix}$$

(It increases if a > b.)

 $|y\rangle$

slopeslope

EIGENVECTOR

slope

Diagonal ($\mathbb{R}^2 = \mathbb{Q}$)-matrix acts on vector $\mathbf{v}^{x/y}$.

Resulting vector has slope changed by factor $a^2/b^2 =$

$$\mathbf{Q} \cdot \mathbf{v}^{x/y} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix}$$

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slope

b/a

...Finally, the result approaches *EIGENVECTOR* $|x\rangle = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$

of 0-slope which is "immune" to \mathbf{R}^{-1} , \mathbf{Q}^{-1} or \mathbf{Q}^{-n} :

$$\mathbf{R}|y\rangle = (1/b)|y\rangle \qquad \mathbf{Q}^{n}|y\rangle = (1/b^{2})^{n}|y\rangle \qquad Eigensolution \qquad \mathbf{R}^{-1}|x\rangle = (a)|x\rangle \qquad \mathbf{Q}^{-n}|x\rangle = (a^{2})^{n}|x\rangle$$

$$Eigenvalues \qquad Relations \qquad Eigenvalues$$

$$\mathbf{R}^{-1}|x\rangle = (a)|x\rangle \qquad \mathbf{Q}^{-n}|x\rangle = (a^2)^n|x\rangle$$
Eigenvalues

Introduction to dual matrix operator geometry

Review of dual IHO elliptic orbits (Lecture 7-8)

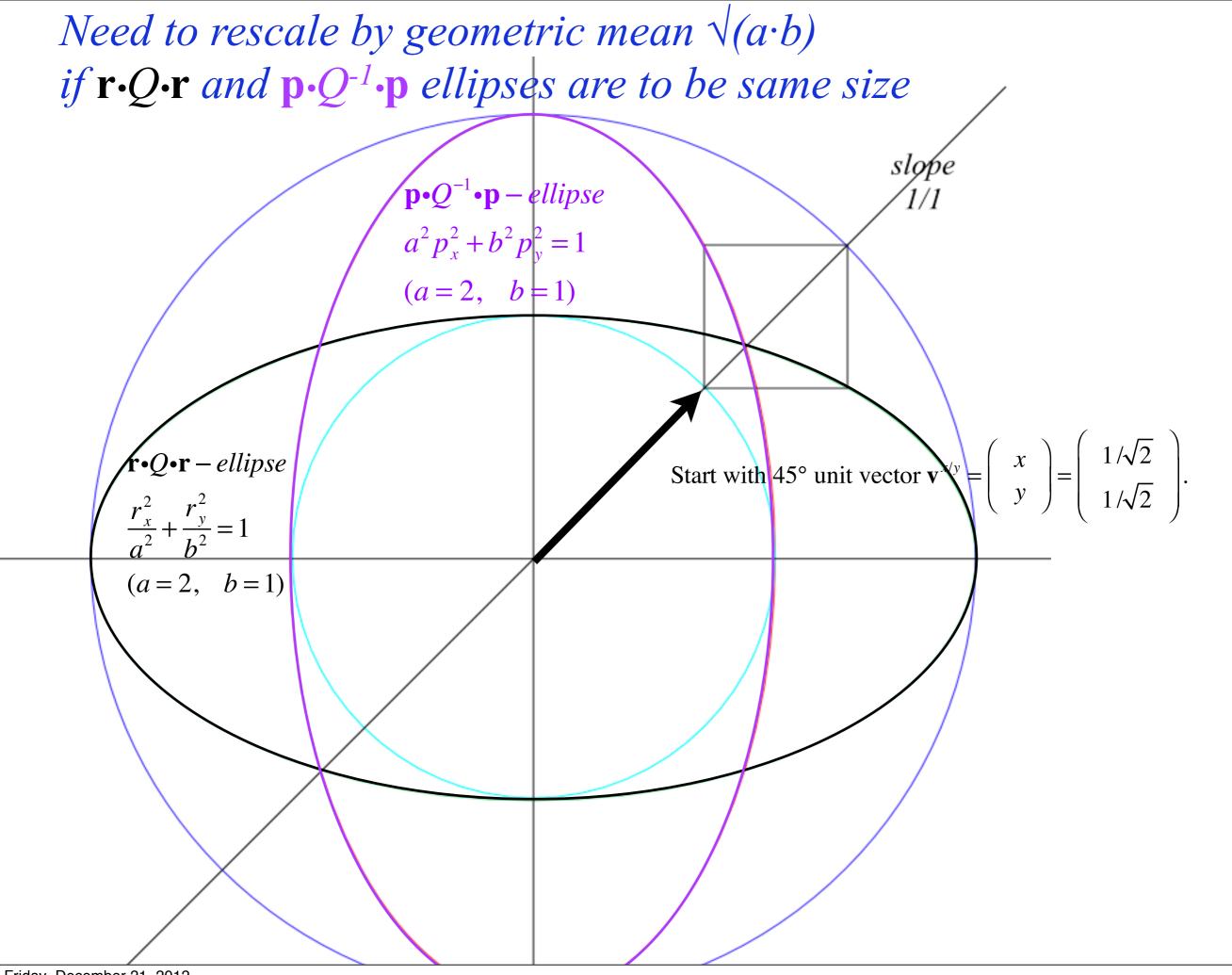
Construction by Phasor-pair projection

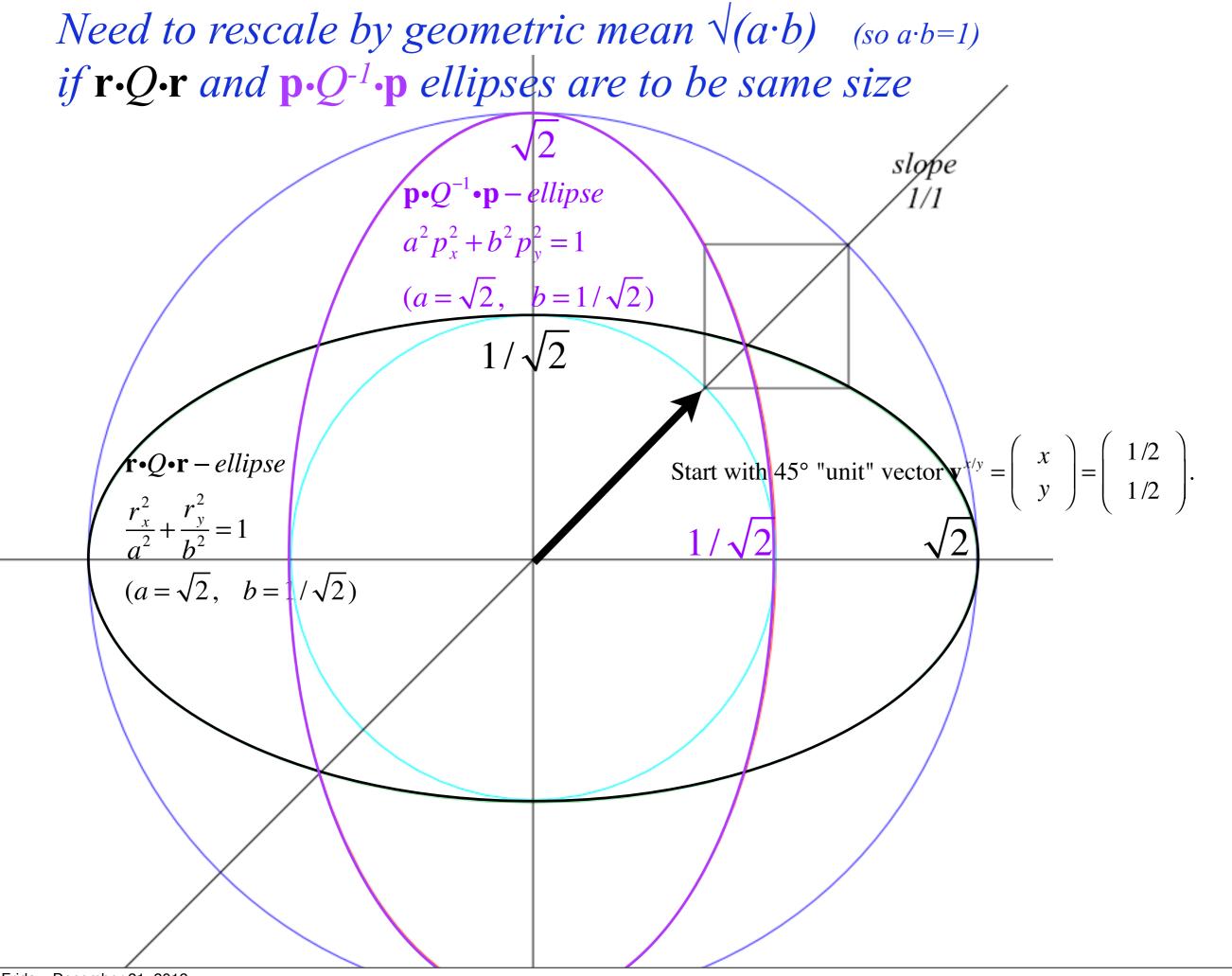
Construction by Kepler anomaly projection

Operator geometric sequences and eigenvectors

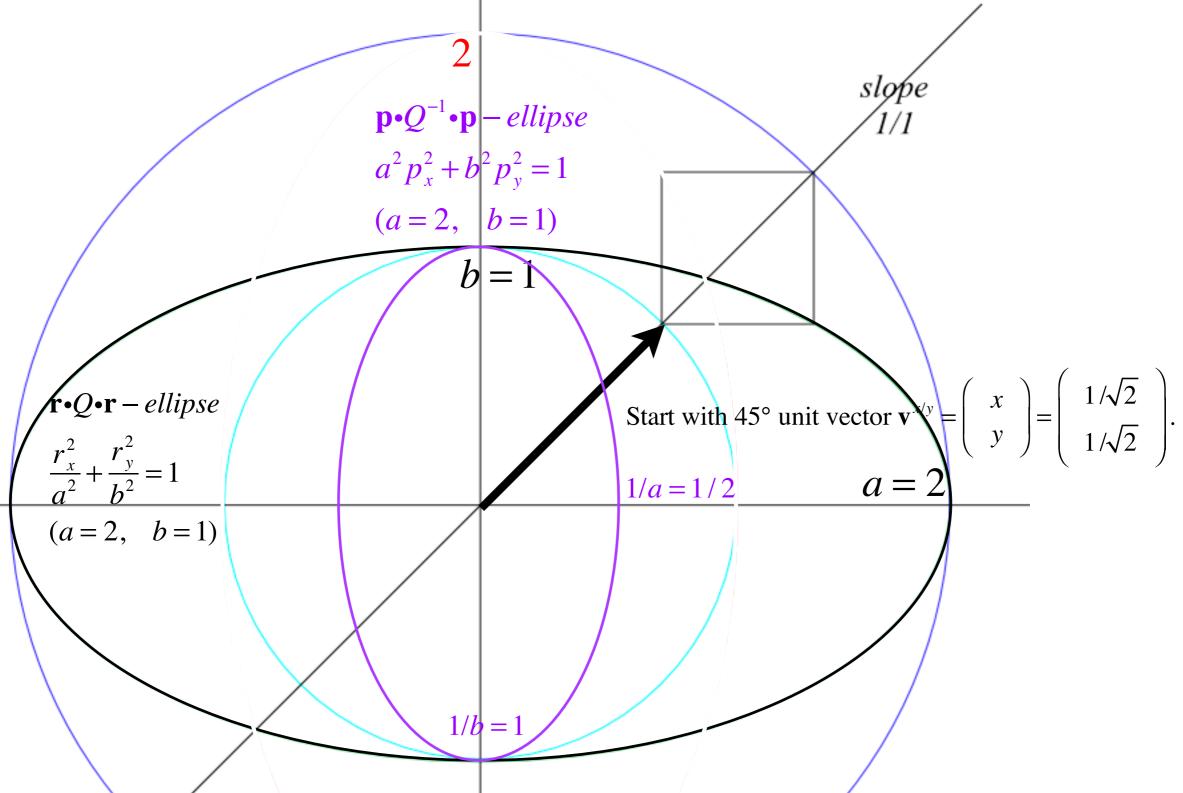
Rescaled description of matrix operator geometry

Vector calculus of tensor operation

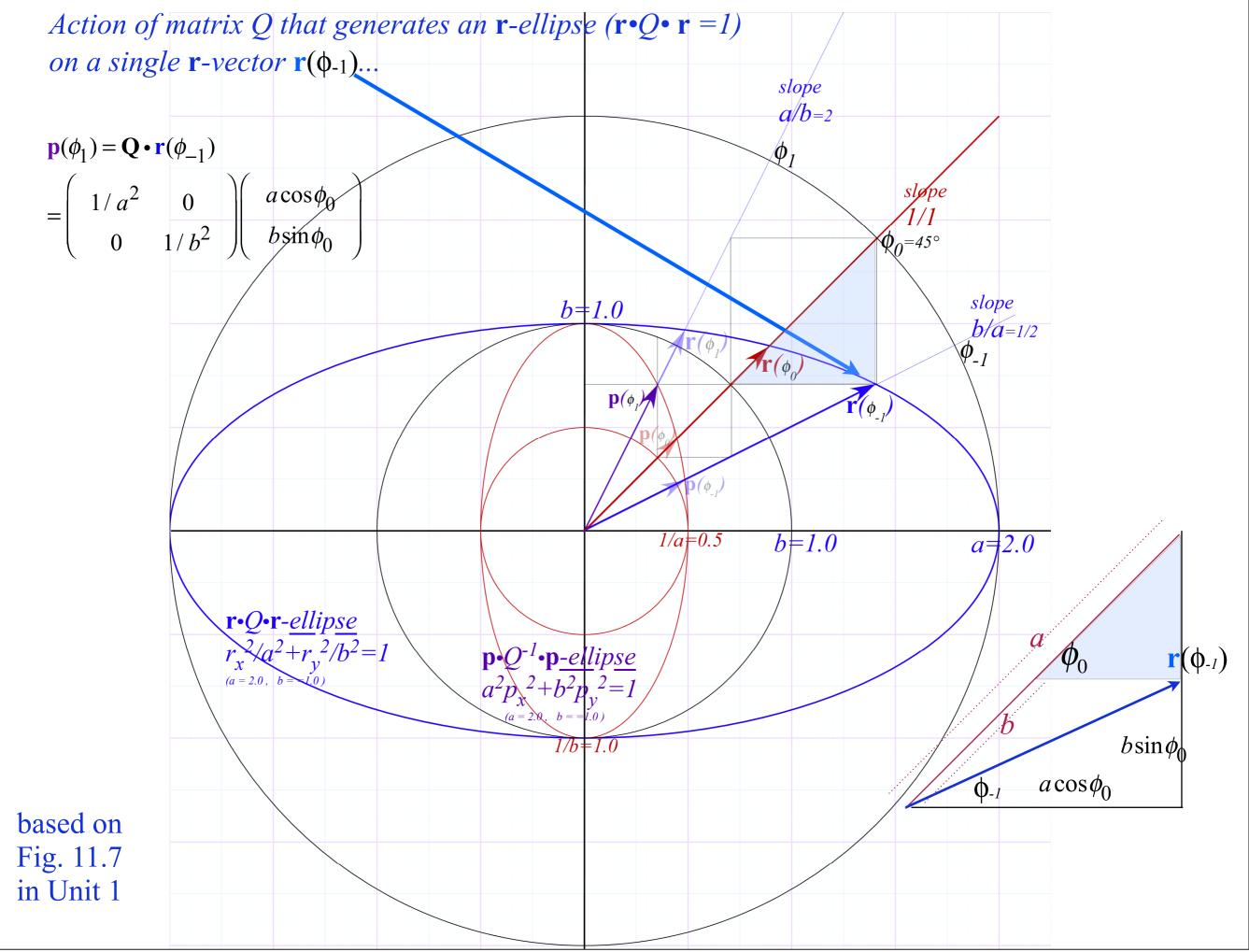


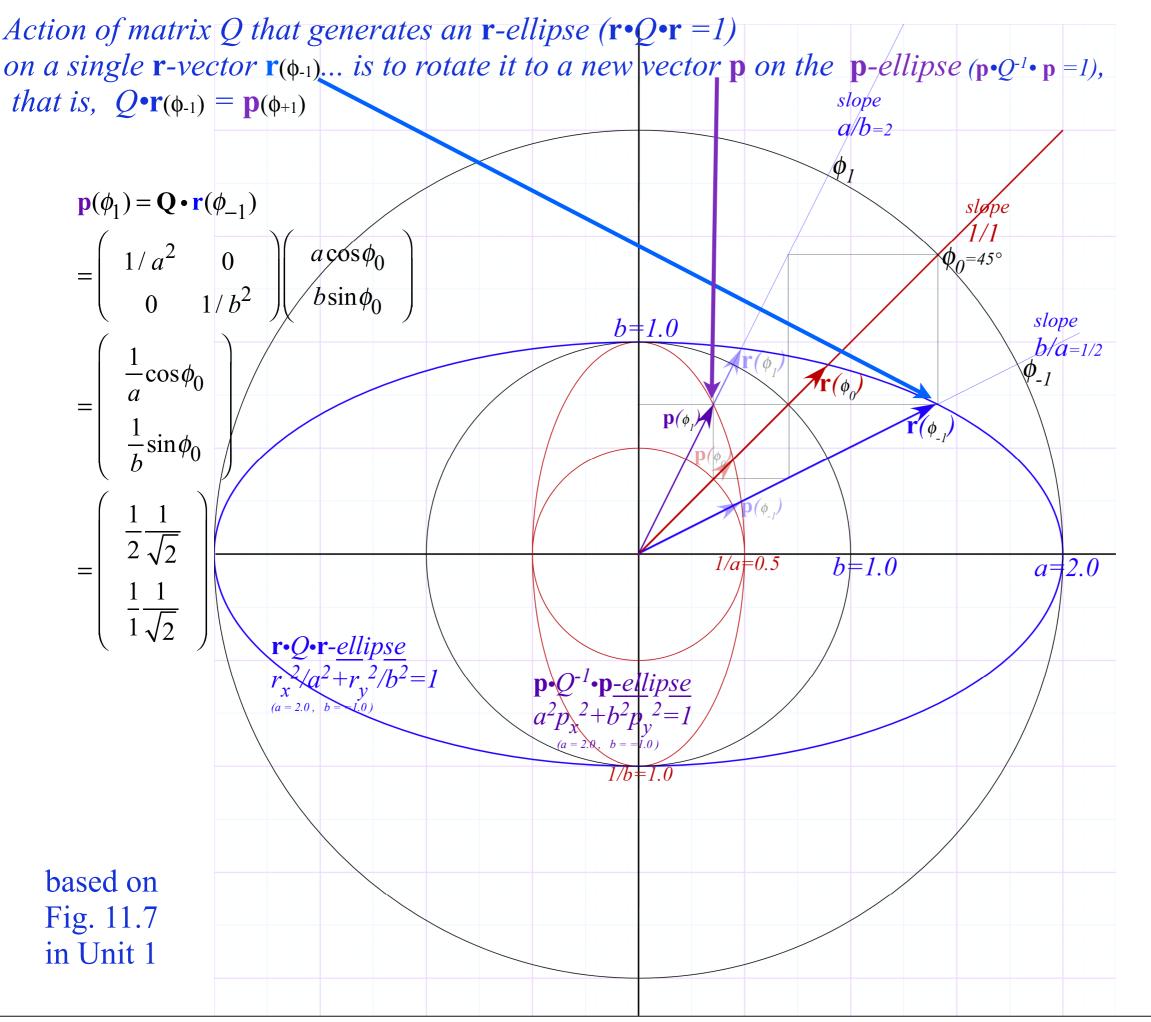


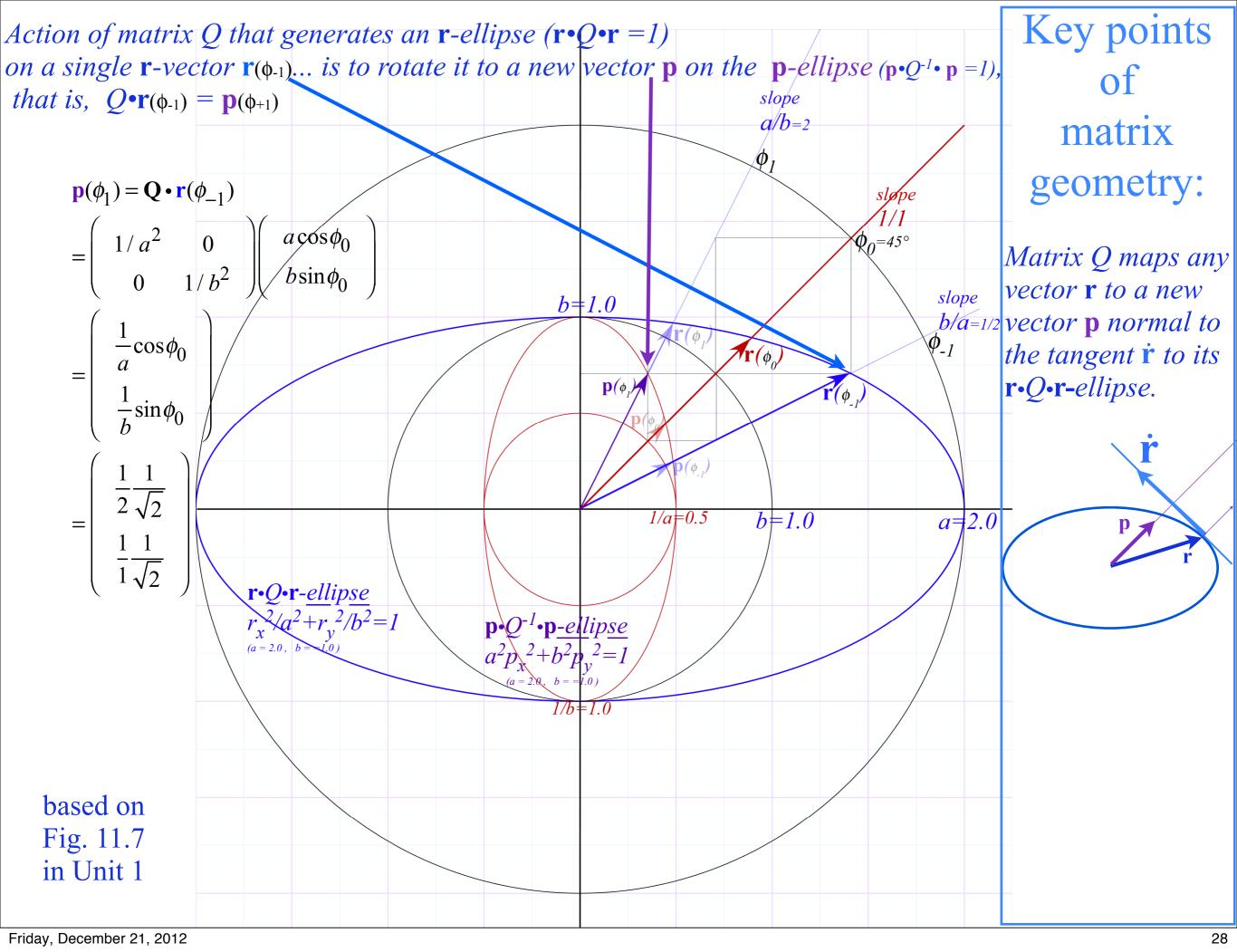


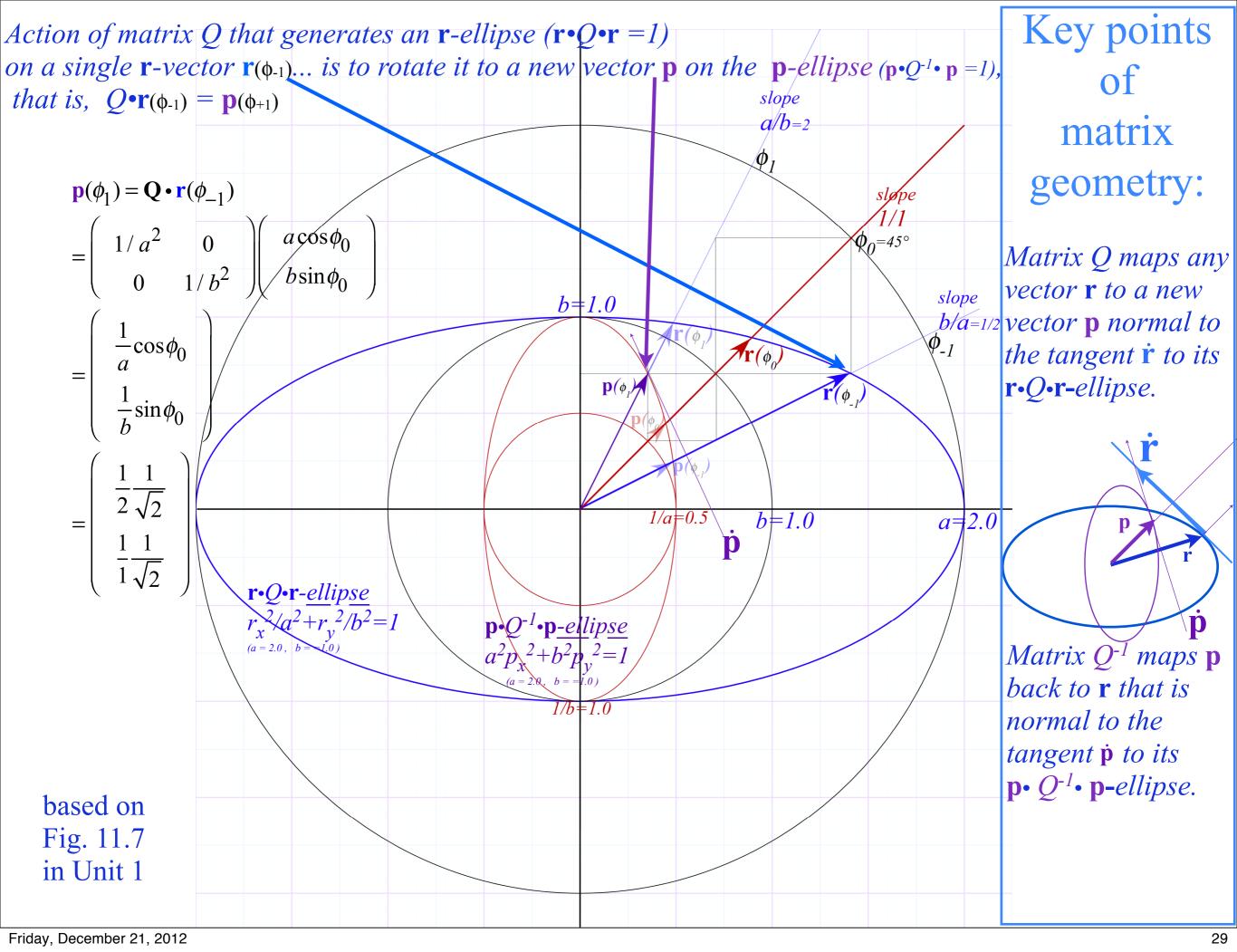


This is a clearer choice. It separates **r** and **p** into different spaces









Introduction to dual matrix operator geometry

Review of dual IHO elliptic orbits (Lecture 7-8)

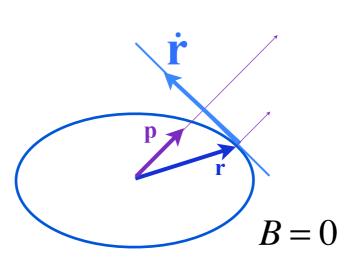
Construction by Phasor-pair projection

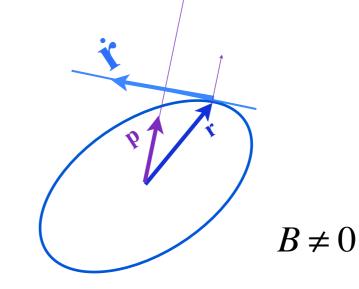
Construction by Kepler anomaly projection

Operator geometric sequences and eigenvectors

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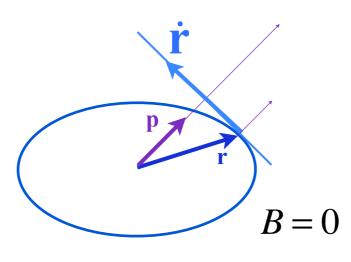


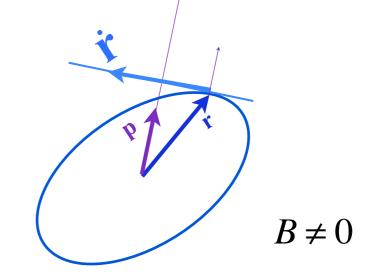


Derive matrix "normal-to-ellipse" geometry by vector calculus:

Let matrix
$$Q = \begin{pmatrix} A & B \\ B & D \end{pmatrix}$$

define the ellipse
$$1 = \mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A & B \\ B & D \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{pmatrix} = A \cdot x^2 + 2B \cdot xy + D \cdot y^2 = 1$$





Derive matrix "normal-to-ellipse" geometry by vector calculus:

Let matrix
$$Q = \begin{pmatrix} A & B \\ B & D \end{pmatrix}$$

define the ellipse
$$1 = \mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A & B \\ B & D \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{pmatrix} = A \cdot x^2 + 2B \cdot xy + D \cdot y^2 = 1$$

with

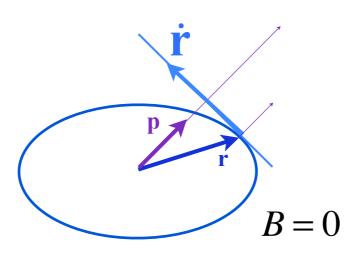
Compare operation by Q on vector **r**

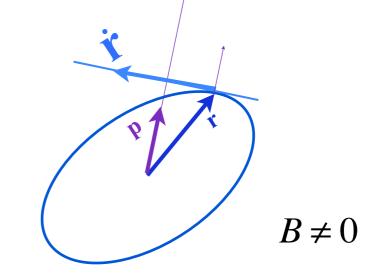
$$\begin{pmatrix} A & B \\ B & D \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{pmatrix}$$

vector derivative or gradient of r•Q•r

$$\frac{\partial}{\partial \mathbf{r}}(\mathbf{r} \cdot Q \cdot \mathbf{r}) = \nabla(\mathbf{r} \cdot Q \cdot \mathbf{r})$$

$$\begin{pmatrix} \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \end{pmatrix} (A \cdot x^2 + 2B \cdot xy + D \cdot y^2) = \begin{pmatrix} 2A \cdot x + 2B \cdot y \\ 2B \cdot x + 2D \cdot y \end{pmatrix}$$





Derive matrix "normal-to-ellipse" geometry by vector calculus:

Let matrix
$$Q = \begin{pmatrix} A & B \\ B & D \end{pmatrix}$$

define the ellipse
$$1 = \mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A & B \\ B & D \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{pmatrix} = A \cdot x^2 + 2B \cdot xy + D \cdot y^2 = 1$$

with

Compare operation by Q on vector **r**

$$\begin{pmatrix} A & B \\ B & D \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{pmatrix}$$

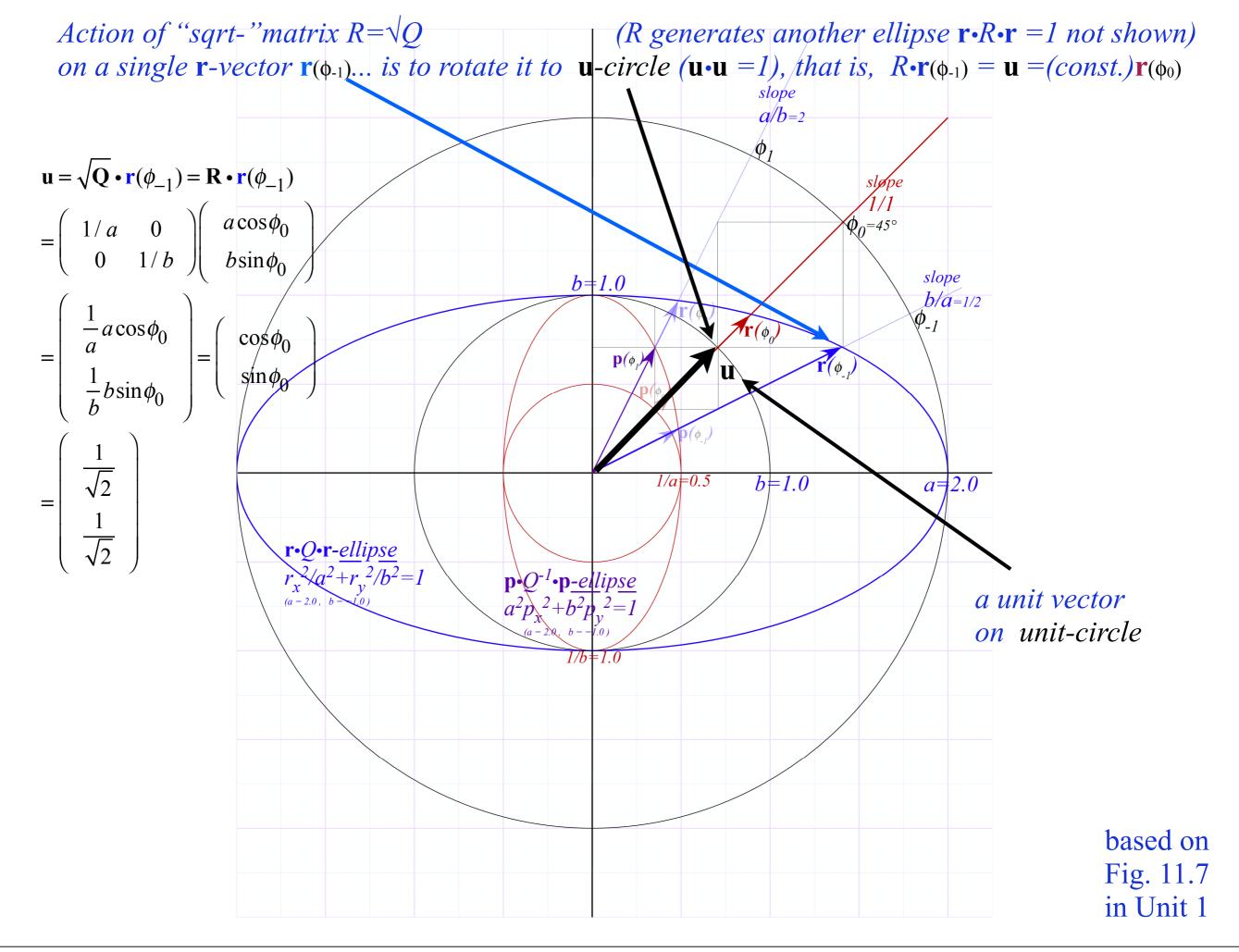
vector derivative or gradient of r•Q•r

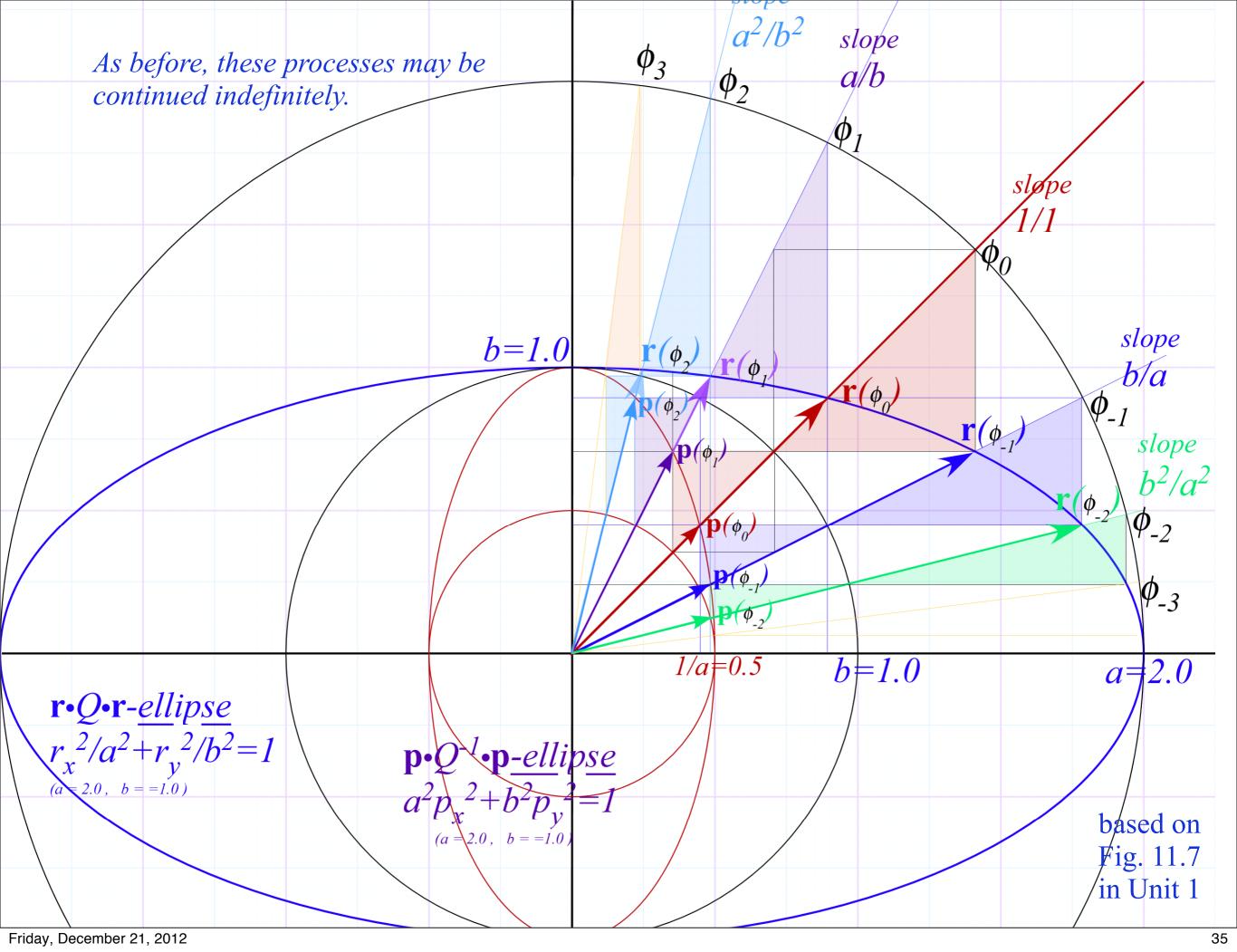
$$\frac{\partial}{\partial \mathbf{r}} (\mathbf{r} \cdot Q \cdot \mathbf{r}) = \nabla (\mathbf{r} \cdot Q \cdot \mathbf{r})$$

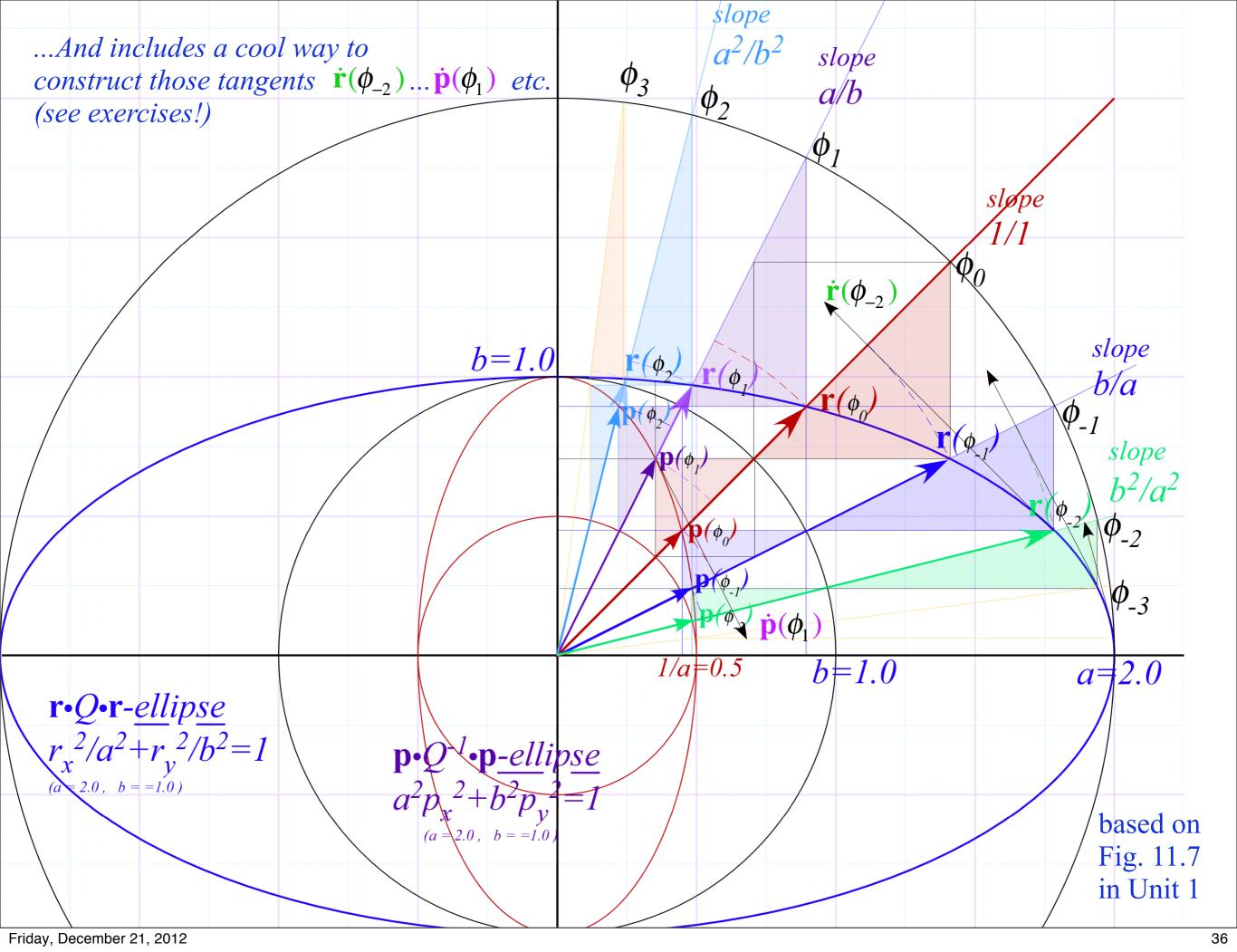
$$\begin{pmatrix} \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \end{pmatrix} (A \cdot x^2 + 2B \cdot xy + D \cdot y^2) = \begin{pmatrix} 2A \cdot x + 2B \cdot y \\ 2B \cdot x + 2D \cdot y \end{pmatrix}$$

Very simple result:

$$\frac{\partial}{\partial \mathbf{r}} \left(\frac{\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}}{2} \right) = \nabla \left(\frac{\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}}{2} \right) = \mathbf{Q} \cdot \mathbf{r}$$







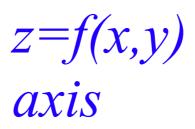
Introduction to Lagrangian-Hamiltonian duality

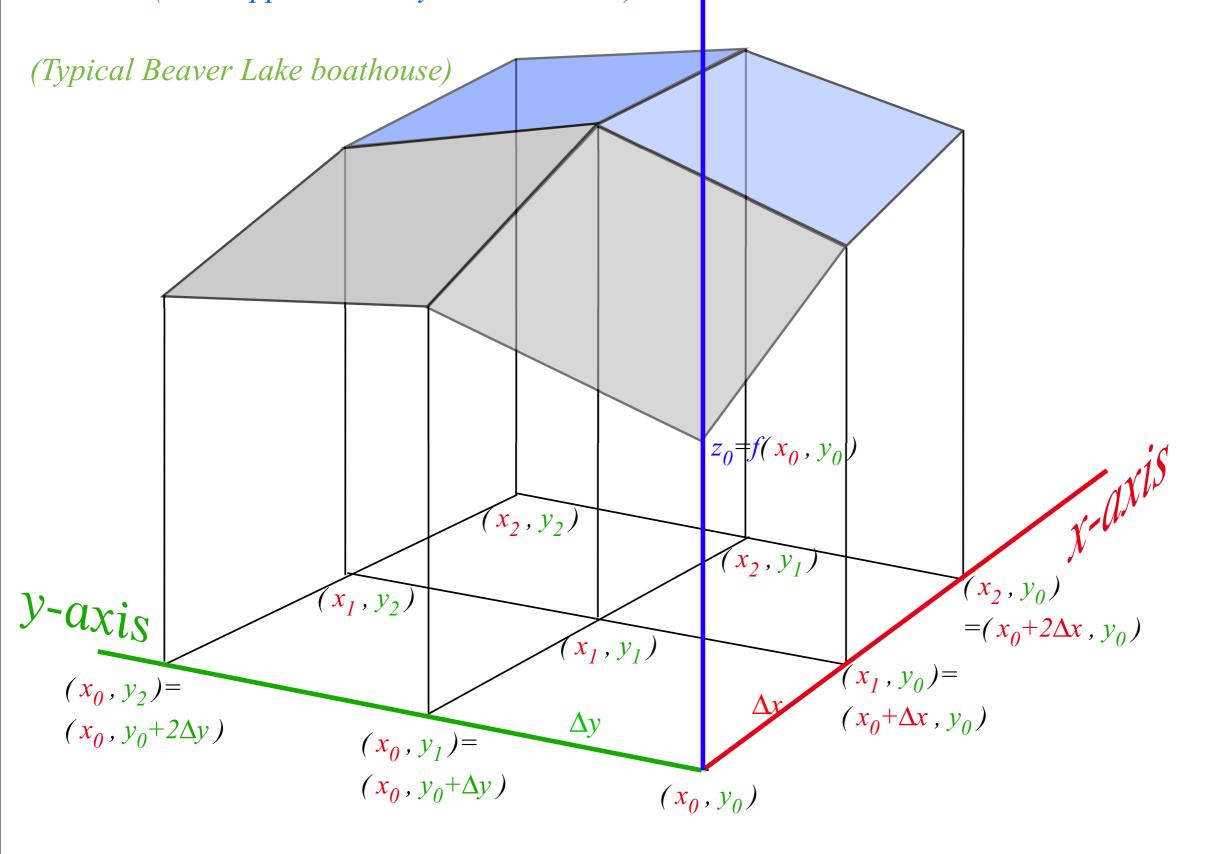
Review of partial differential relations

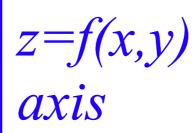
Chain rule and order symmetry

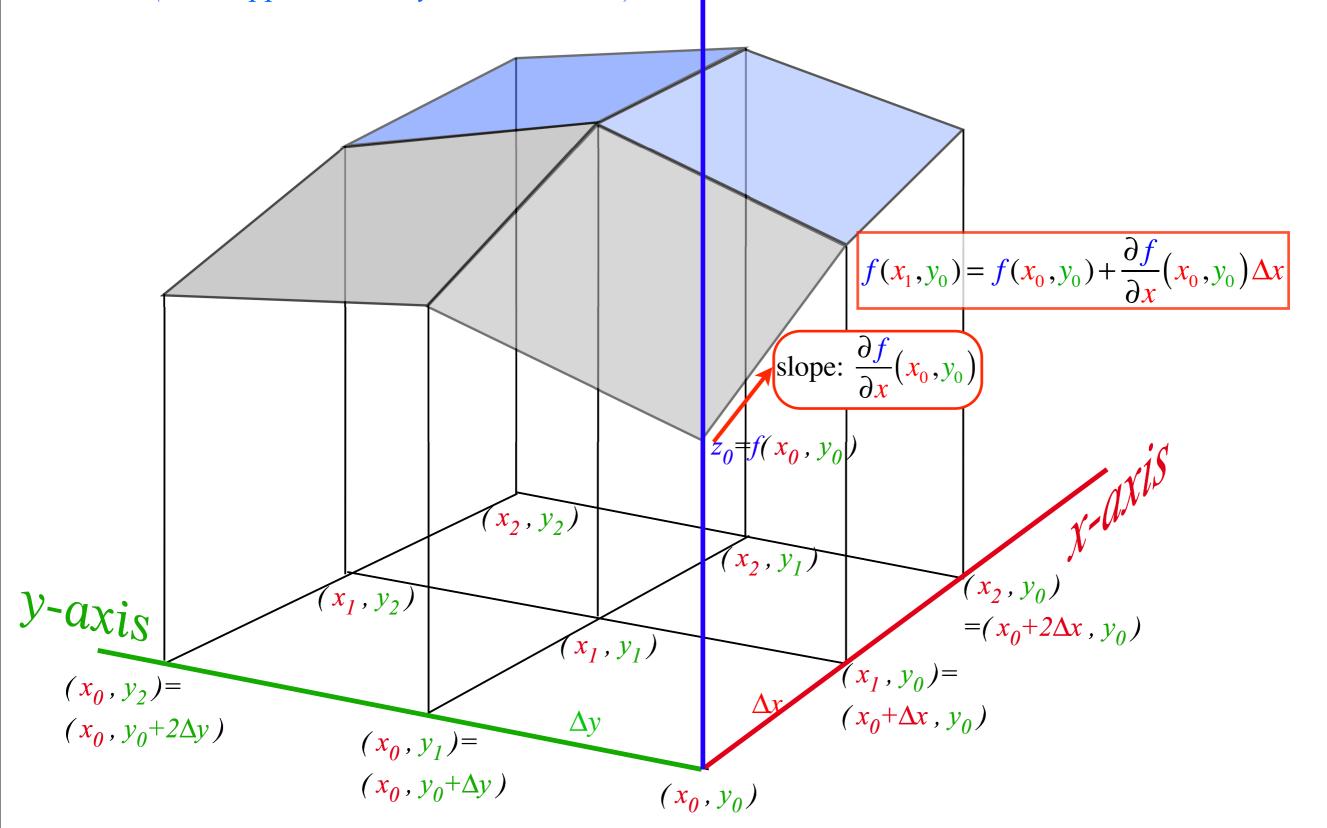
Duality relations of Lagrangian and Hamiltonian ellipse

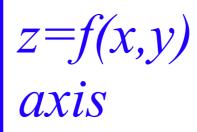
Introducing the 1^{st} (partial $\frac{\partial?}{\partial?}$) differential equations of mechanics

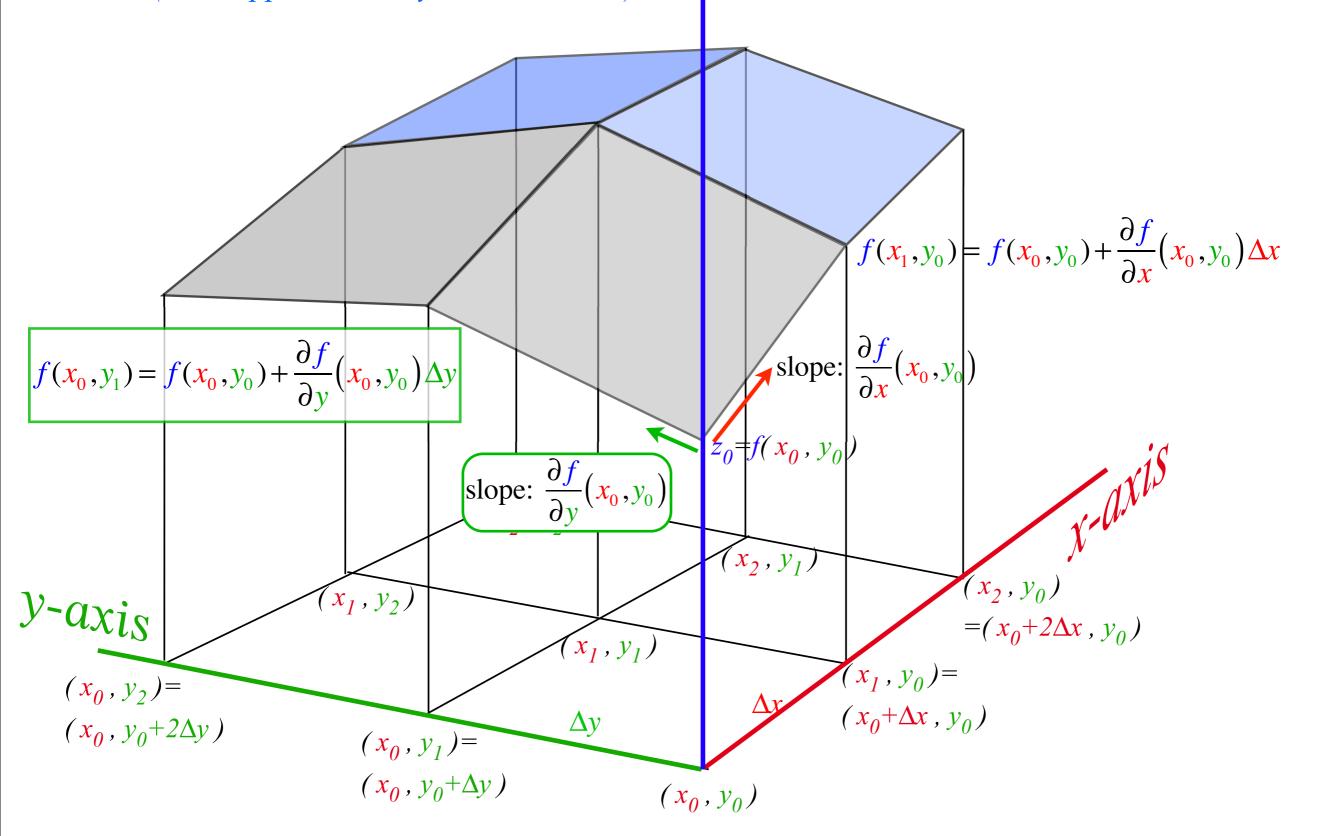




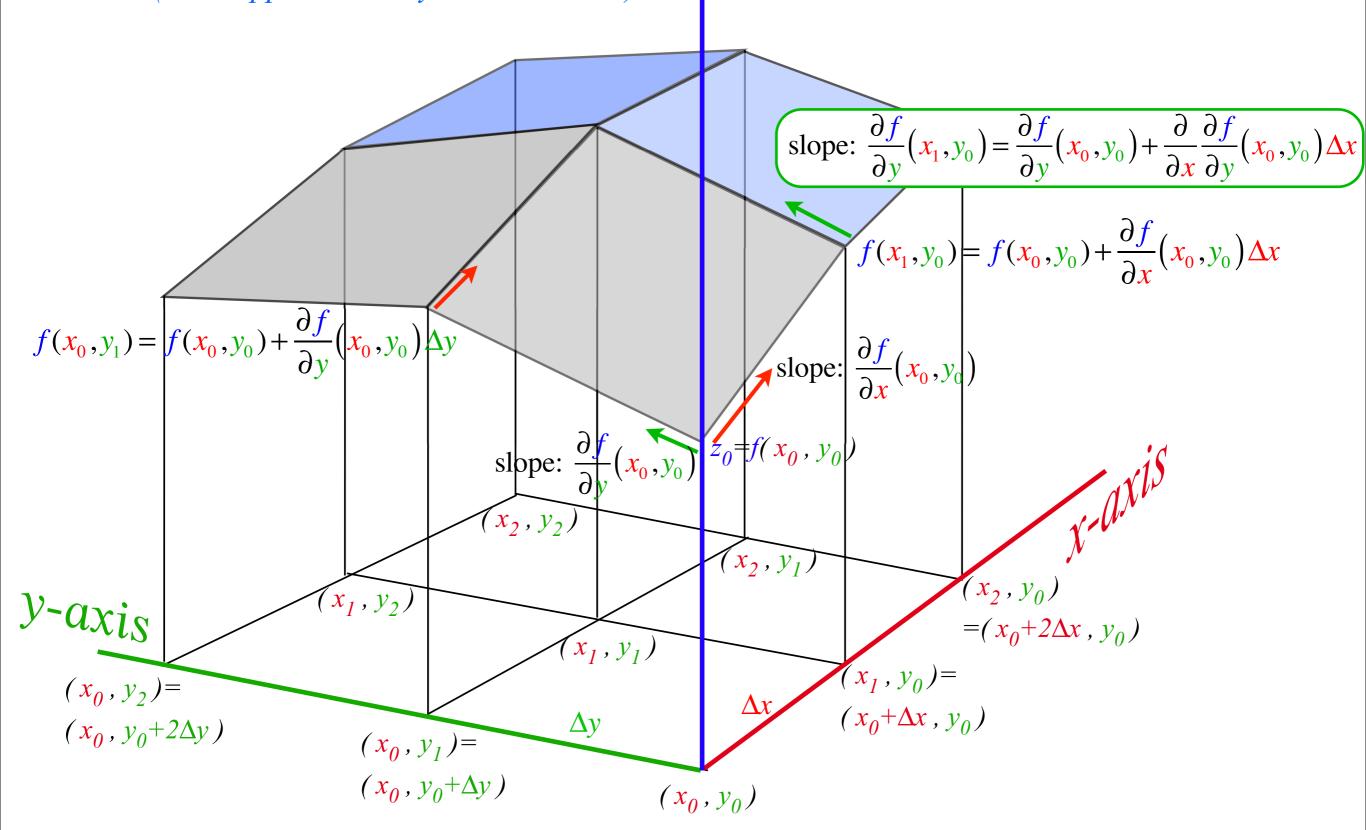




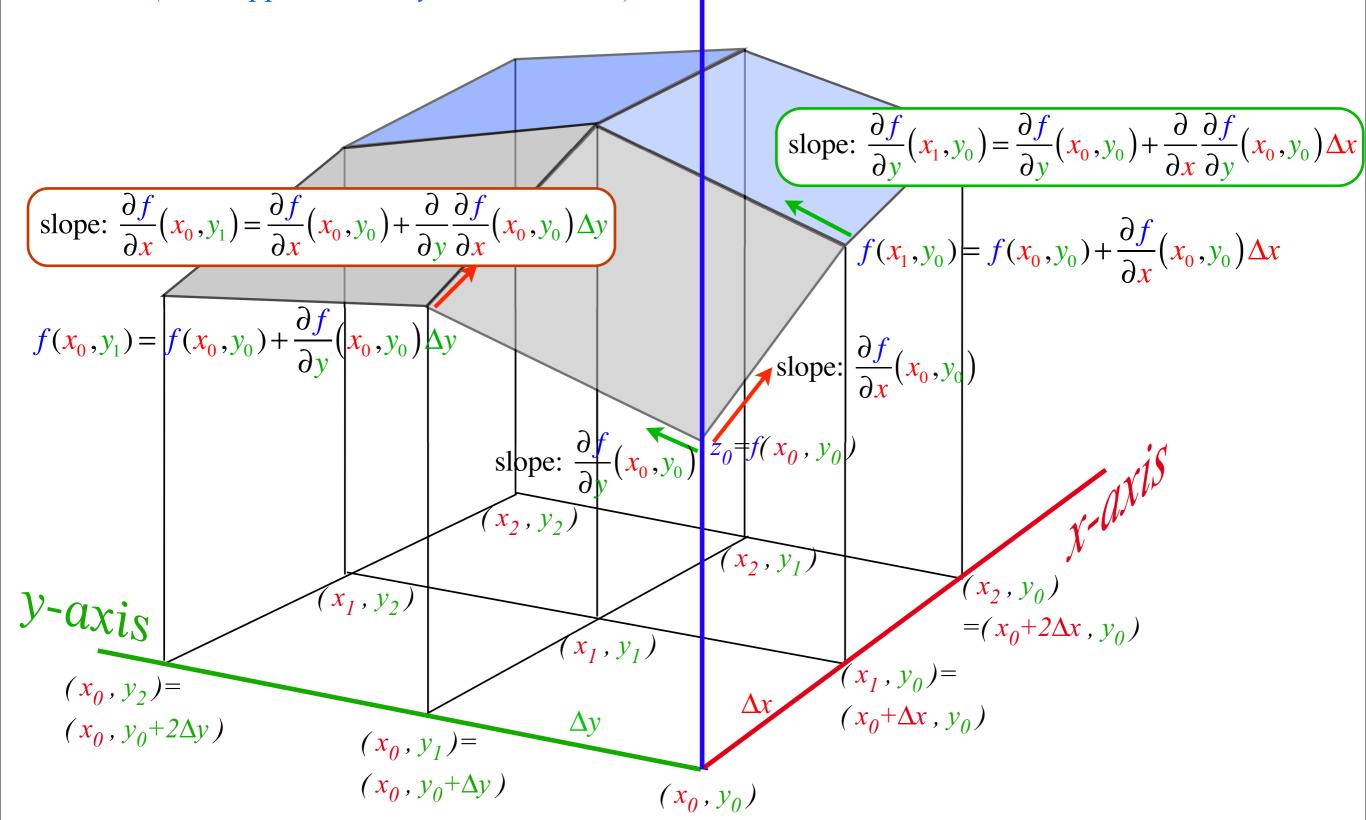


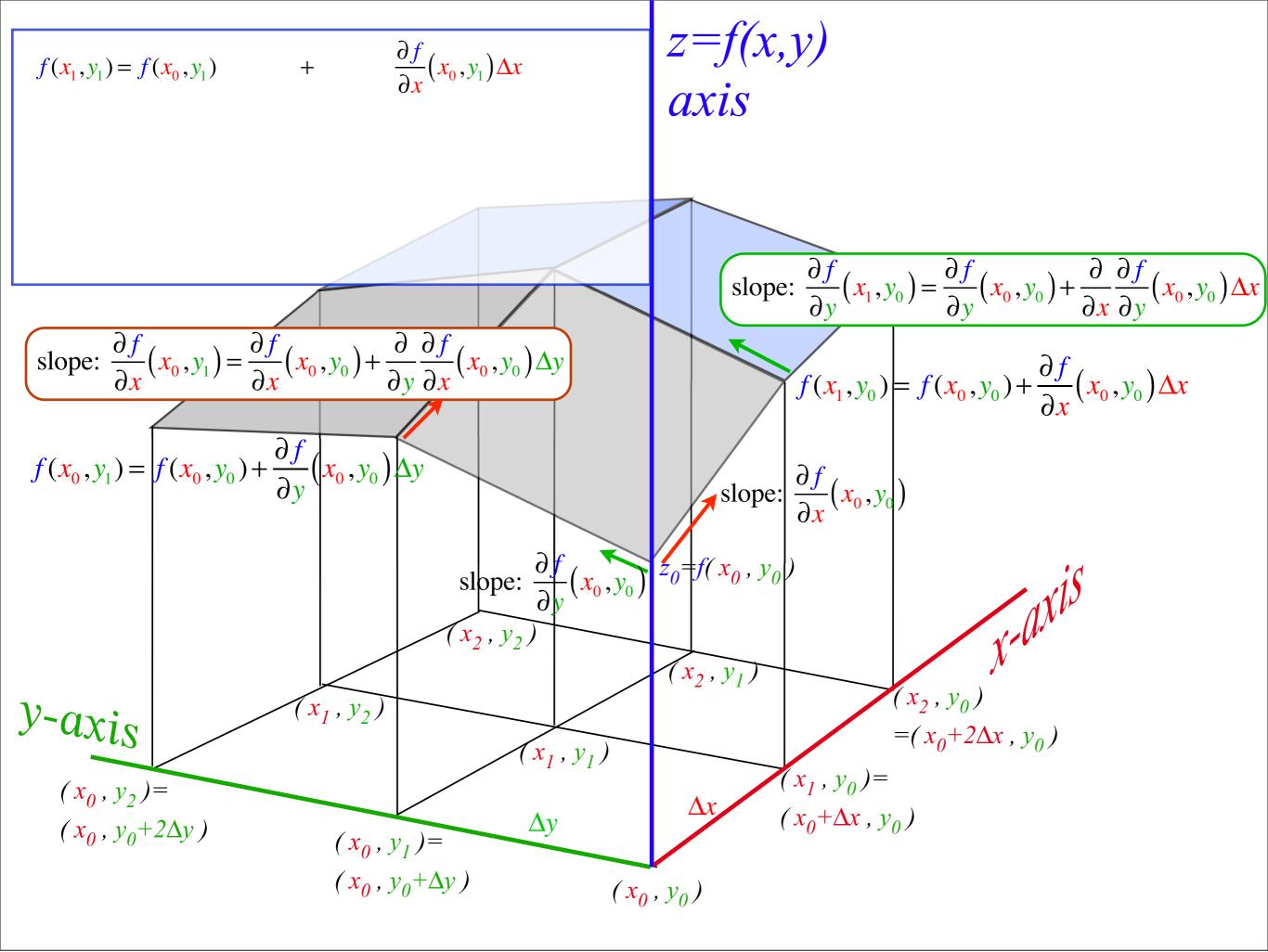


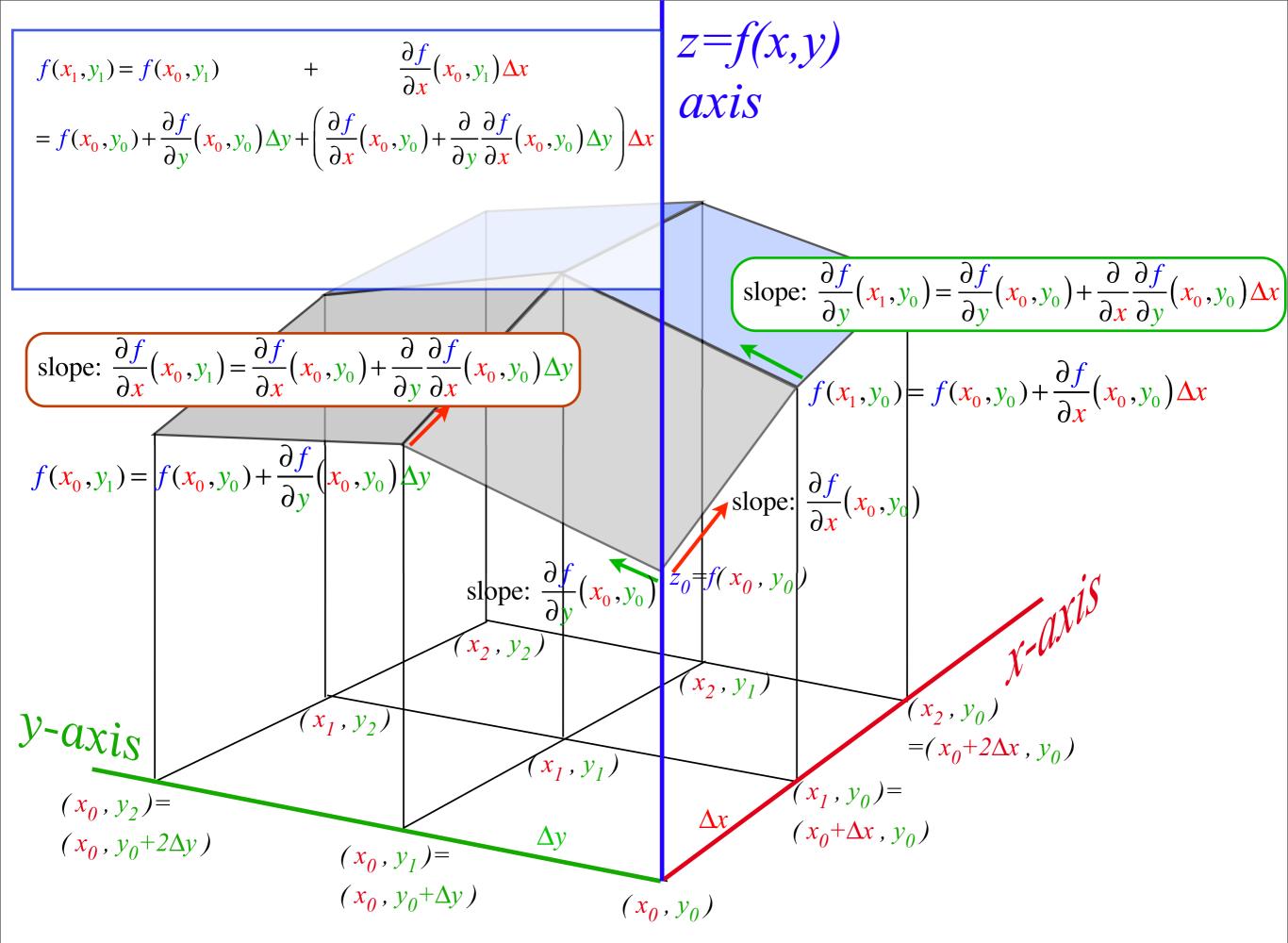
$$z=f(x,y)$$
axis

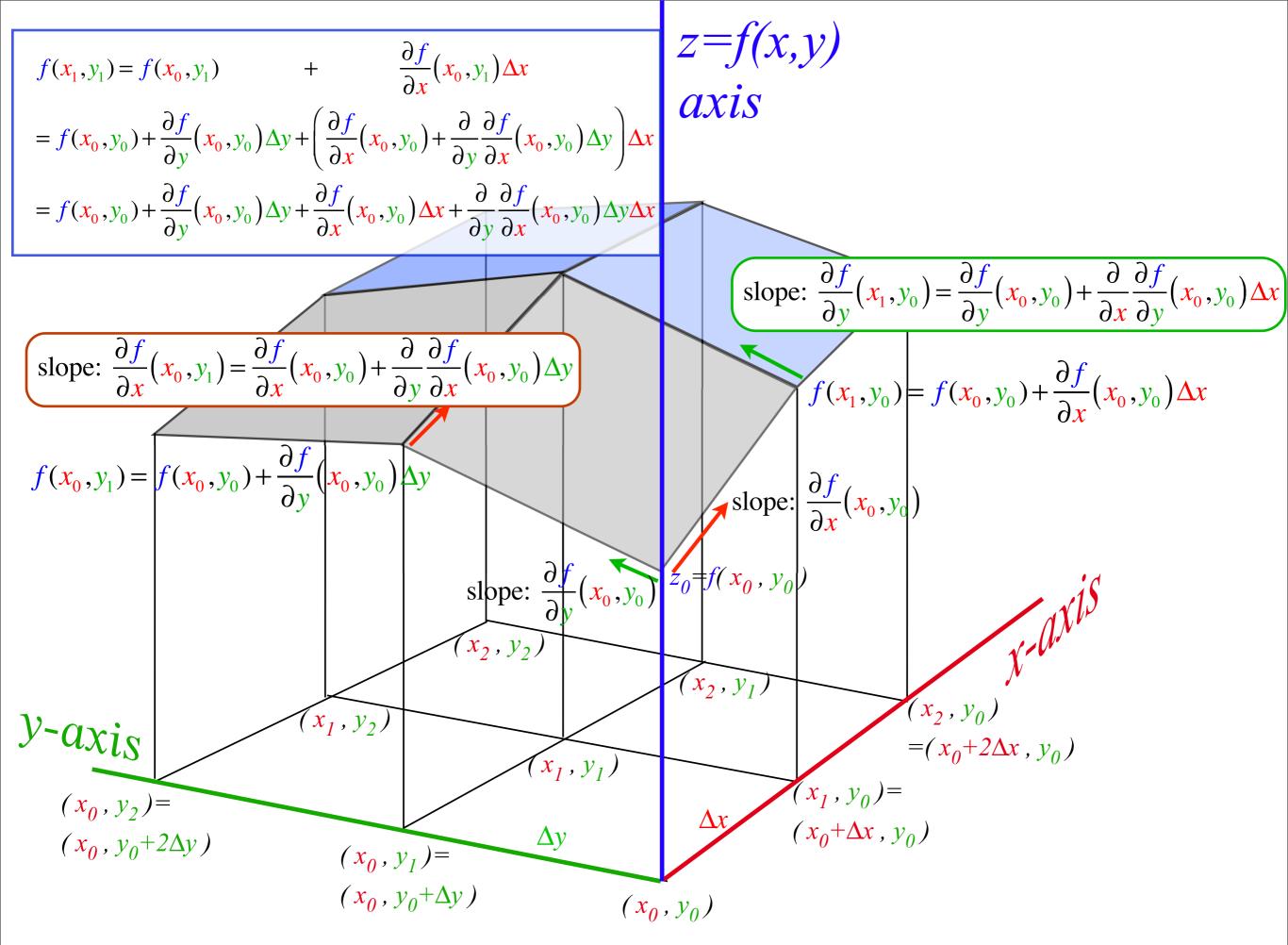


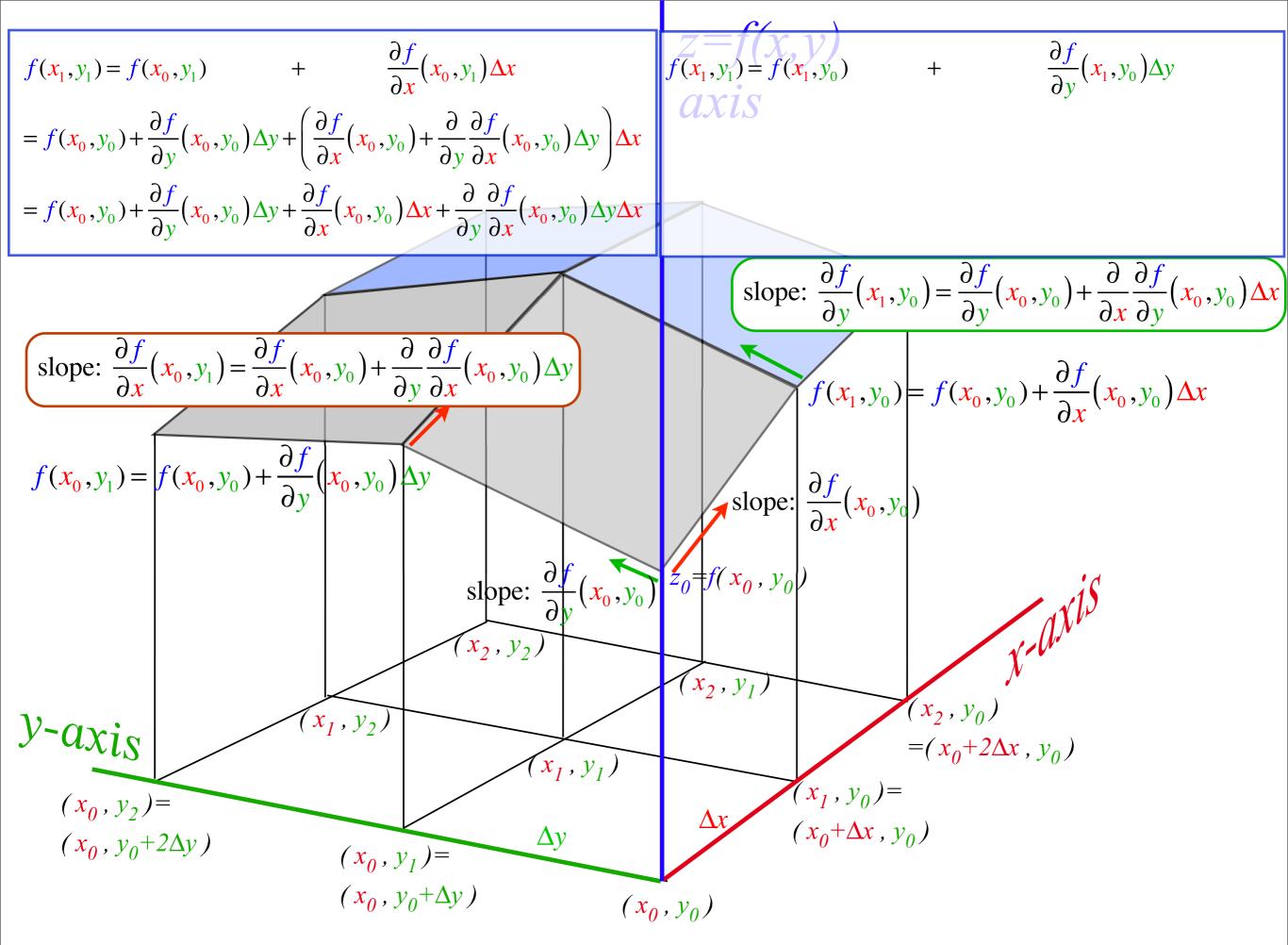
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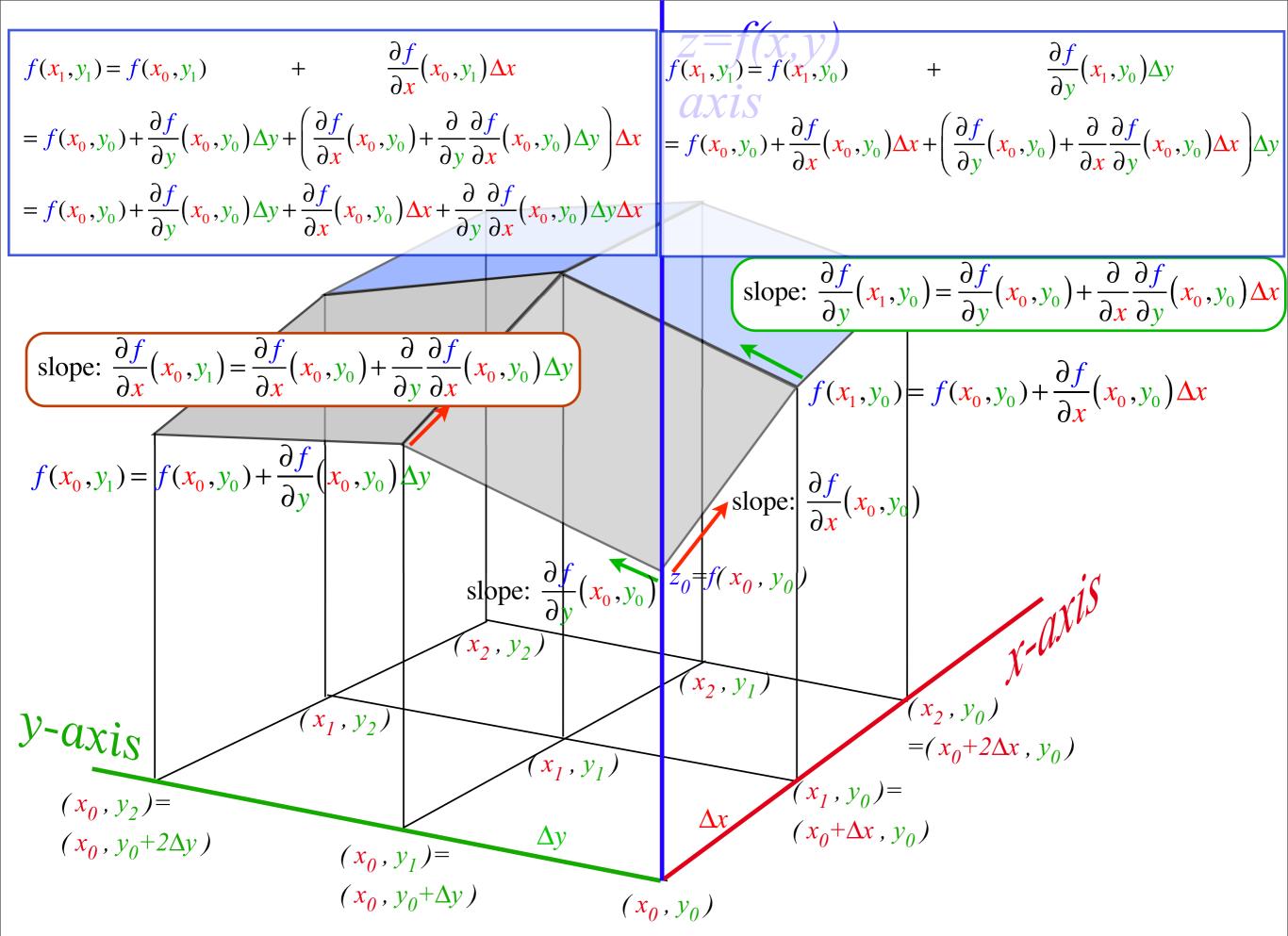


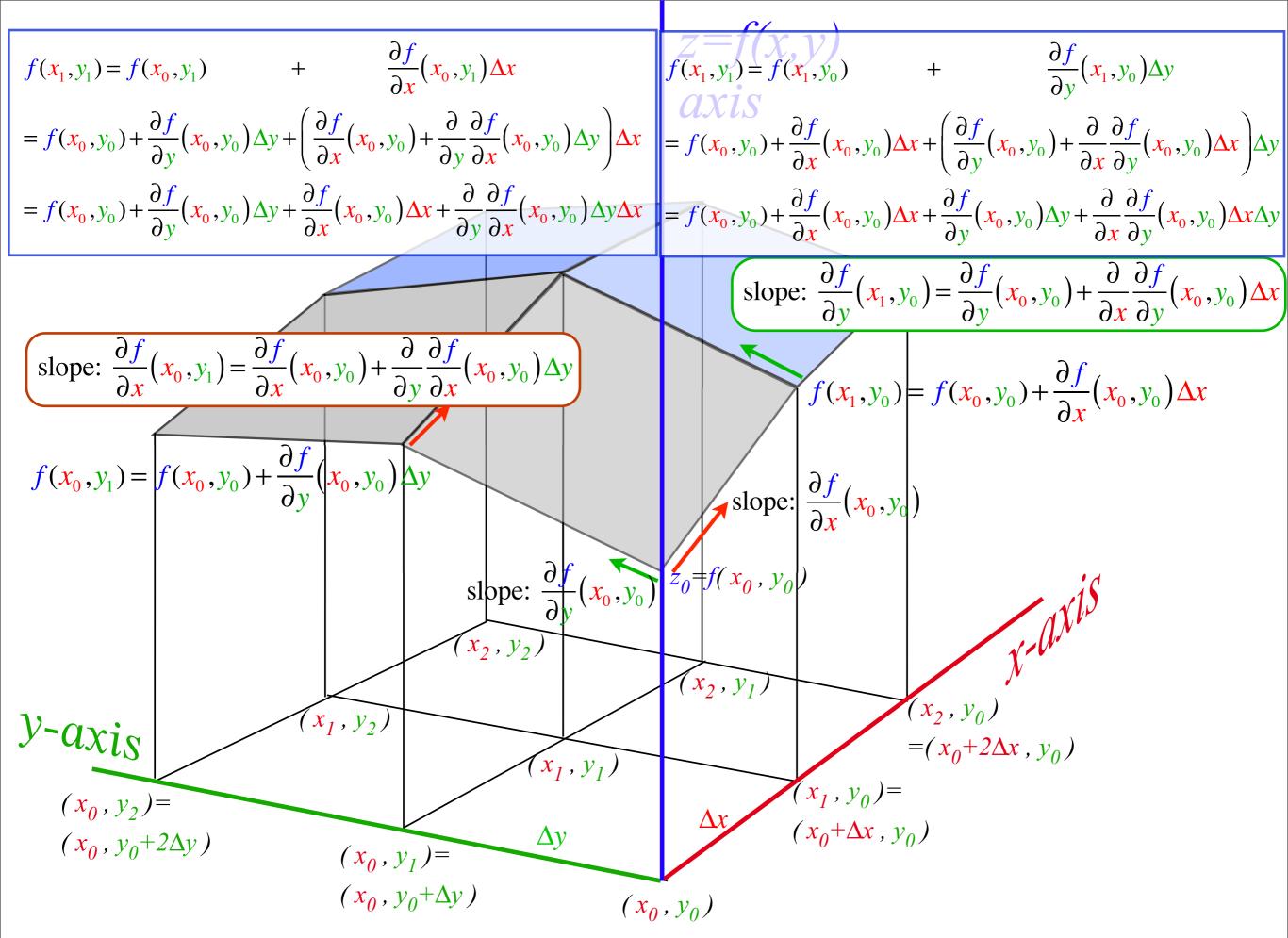












Introduction to Lagrangian-Hamiltonian duality

Review of partial differential relations

Chain rule and order symmetry

Duality relations of Lagrangian and Hamiltonian ellipse Introducing the 1^{st} (partial $\frac{\partial?}{\partial?}$) differential equations of mechanics

$$f(x_1, y_1) = f(x_0, y_0) + \frac{\partial f}{\partial x}(x_0, y_0) \Delta x + \frac{\partial f}{\partial y}(x_0, y_0) \Delta y + \frac{\partial f}{\partial y}(x_0, y_0) \Delta x \Delta y$$

$$= f(x_0, y_0) + \frac{\partial f}{\partial y}(x_0, y_0) \Delta y + \frac{\partial f}{\partial x}(x_0, y_0) \Delta x + \frac{\partial f}{\partial x}(x_0, y_0) \Delta x \Delta y$$

$$f(x_{1}, y_{1}) = f(x_{0}, y_{0}) + \frac{\partial f}{\partial x}(x_{0}, y_{0}) \Delta x + \frac{\partial f}{\partial y}(x_{0}, y_{0}) \Delta y + \frac{\partial f}{\partial y}(x_{0}, y_{0}) \Delta x \Delta y$$

$$= f(x_{0}, y_{0}) + \frac{\partial f}{\partial y}(x_{0}, y_{0}) \Delta y + \frac{\partial f}{\partial x}(x_{0}, y_{0}) \Delta x + \frac{\partial f}{\partial x}(x_{0}, y_{0}) \Delta x \Delta y$$

1. Chain rules

$$\begin{aligned} \left[f(\mathbf{x}_{1}, y_{1}) - f(\mathbf{x}_{0}, y_{0}) \right] &= df = \frac{\partial f}{\partial \mathbf{x}} (\mathbf{x}_{0}, y_{0}) d\mathbf{x} + \frac{\partial f}{\partial y} (\mathbf{x}_{0}, y_{0}) dy \dots_{(keep \ 1^{st} - order \ terms \ only!)} \\ \frac{df}{dt} &= \frac{\partial f}{\partial \mathbf{x}} (\mathbf{x}_{0}, y_{0}) \frac{d\mathbf{x}}{dt} + \frac{\partial f}{\partial y} (\mathbf{x}_{0}, y_{0}) \frac{d\mathbf{y}}{dt} \\ \dot{f} &= \frac{\partial f}{\partial \mathbf{x}} \dot{\mathbf{x}} + \frac{\partial f}{\partial y} \dot{\mathbf{y}} \qquad (shorthand \ notation) \end{aligned}$$

$$f(x_{1}, y_{1}) = f(x_{0}, y_{0}) + \frac{\partial f}{\partial x}(x_{0}, y_{0}) \Delta x + \frac{\partial f}{\partial y}(x_{0}, y_{0}) \Delta y + \frac{\partial f}{\partial y}(x_{0}, y_{0}) \Delta x \Delta y$$

$$= f(x_{0}, y_{0}) + \frac{\partial f}{\partial y}(x_{0}, y_{0}) \Delta y + \frac{\partial f}{\partial x}(x_{0}, y_{0}) \Delta x + \frac{\partial f}{\partial x}(x_{0}, y_{0}) \Delta x + \frac{\partial f}{\partial x}(x_{0}, y_{0}) \Delta y \Delta x$$

1. Chain rules

$$\begin{aligned} \left[f(\mathbf{x}_{1}, y_{1}) - f(\mathbf{x}_{0}, y_{0}) \right] &= df = \frac{\partial f}{\partial x} (\mathbf{x}_{0}, y_{0}) dx + \frac{\partial f}{\partial y} (\mathbf{x}_{0}, y_{0}) dy \dots_{(keep \ 1^{st} - order \ terms \ only!)} \\ \frac{df}{dt} &= \frac{\partial f}{\partial x} (\mathbf{x}_{0}, y_{0}) \frac{dx}{dt} + \frac{\partial f}{\partial y} (\mathbf{x}_{0}, y_{0}) \frac{dy}{dt} \\ \dot{f} &= \frac{\partial f}{\partial x} \dot{x} + \frac{\partial f}{\partial y} \dot{y} \qquad (shorthand \ notation) = \partial_{x} f \dot{x} + \partial_{y} f \dot{y} \end{aligned}$$

2. Symmetry of partial deriv. ordering

(pay attention to the 2^{nd} – order terms, too!)

$$\frac{\partial}{\partial y} \frac{\partial f}{\partial x} = \frac{\partial}{\partial x} \frac{\partial f}{\partial y} \quad \text{or:} \quad \frac{\partial^2 f}{\partial y \partial x} = \frac{\partial^2 f}{\partial x \partial y} \quad \text{or:} \quad \partial_y \partial_x f = \partial_x \partial_y f$$

(shorthand notation)

$$f(x_{1}, y_{1}) = f(x_{0}, y_{0}) + \frac{\partial f}{\partial x}(x_{0}, y_{0}) \Delta x + \frac{\partial f}{\partial y}(x_{0}, y_{0}) \Delta y + \frac{\partial f}{\partial y}(x_{0}, y_{0}) \Delta x \Delta y$$

$$= f(x_{0}, y_{0}) + \frac{\partial f}{\partial y}(x_{0}, y_{0}) \Delta y + \frac{\partial f}{\partial x}(x_{0}, y_{0}) \Delta x + \frac{\partial f}{\partial x}(x_{0}, y_{0}) \Delta x + \frac{\partial f}{\partial x}(x_{0}, y_{0}) \Delta y \Delta x$$

1. Chain rules

$$\begin{aligned} \left[f(\mathbf{x}_{1}, y_{1}) - f(\mathbf{x}_{0}, y_{0}) \right] &= df = \frac{\partial f}{\partial x} (\mathbf{x}_{0}, y_{0}) dx + \frac{\partial f}{\partial y} (\mathbf{x}_{0}, y_{0}) dy \dots_{(keep \ 1^{st} - order \ terms \ only!)} \\ \frac{df}{dt} &= \frac{\partial f}{\partial x} (\mathbf{x}_{0}, y_{0}) \frac{dx}{dt} + \frac{\partial f}{\partial y} (\mathbf{x}_{0}, y_{0}) \frac{dy}{dt} \\ \dot{f} &= \frac{\partial f}{\partial x} \dot{x} + \frac{\partial f}{\partial y} \dot{y} \qquad (shorthand \ notation) = \partial_{x} f \dot{x} + \partial_{y} f \dot{y} \end{aligned}$$

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(shorthand notation)

$$Let: \vec{\nabla} = \begin{pmatrix} \partial_x & \partial_y \end{pmatrix} \quad so: \vec{\nabla}f \cdot \mathbf{dr} = \begin{pmatrix} \partial_x f & \partial_y f \end{pmatrix} \cdot \begin{pmatrix} dx \\ dy \end{pmatrix} = \partial_x f \, dx + \partial_y f \, dy = df$$

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Three ways to express energy: Consider kinetic energy (KE) first

1. Lagrangian is explicit function of velocity: $\mathbf{v} = \begin{pmatrix} v_1 \\ v_2 \end{pmatrix}$

$$L(v_k...) = \frac{1}{2} (m_1 v_1^2 + m_2 v_2^2 + ...) = L(\mathbf{v}...) = \frac{1}{2} \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + ... = \frac{1}{2} \begin{pmatrix} v_1 & v_2 \end{pmatrix} \begin{pmatrix} m_1 & 0 \\ 0 & m_2 \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} + ...$$

2. "Estrangian" is explicit function of R-rescaled velocity:

or: "speedinum" $V = \mathbf{V} = \mathbf{V} \cdot \mathbf{v}$ or: $\begin{pmatrix} V_1 \\ V_2 \end{pmatrix} = \begin{pmatrix} \sqrt{m_1} & 0 \\ 0 & \sqrt{m_2} \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix}$

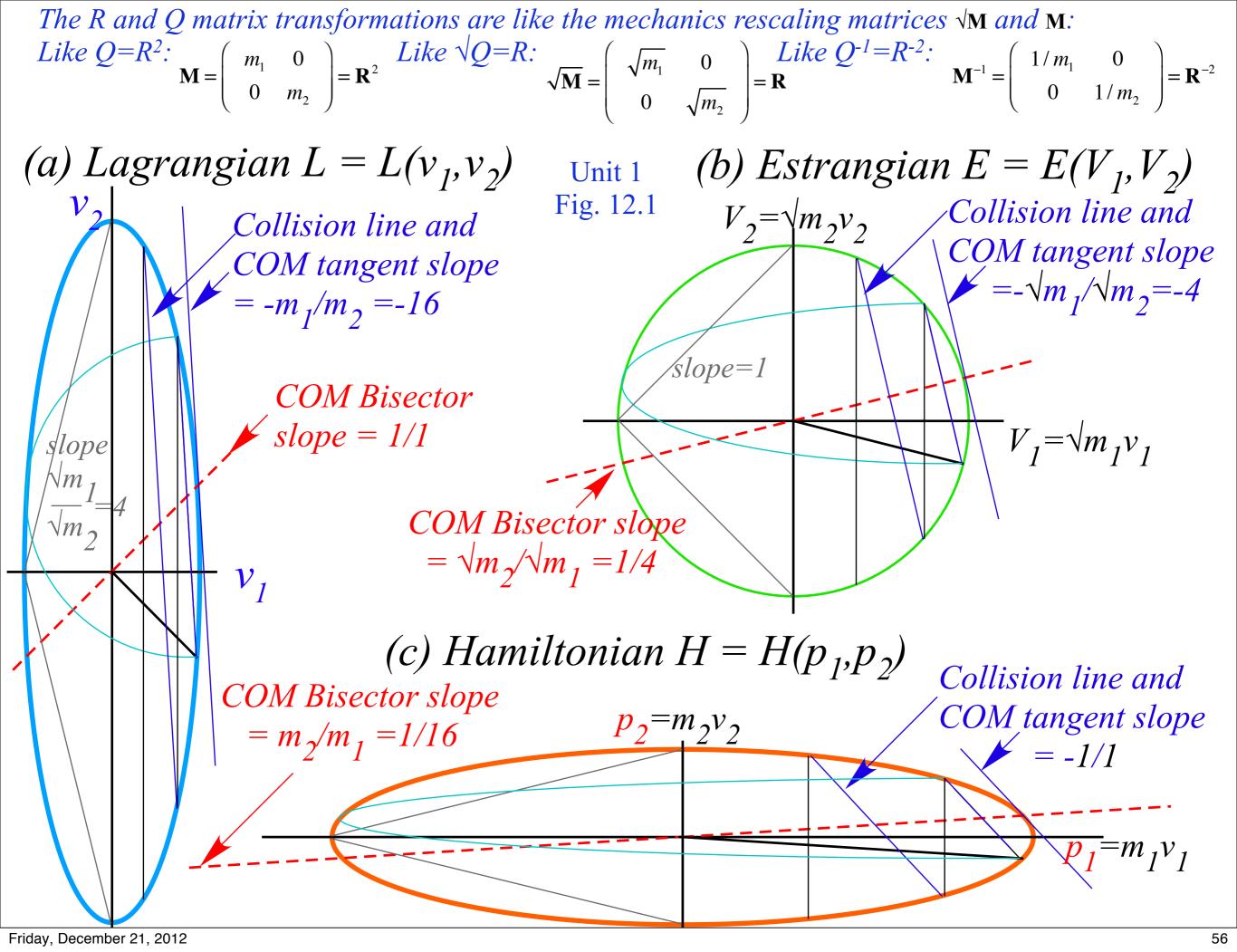
$$E(V_k...) = \frac{1}{2} (V_1^2 + V_2^2 + ...) = E(\mathbf{V}...) = \frac{1}{2} \mathbf{V} \cdot \mathbf{1} \cdot \mathbf{V} + ... = \frac{1}{2} \begin{pmatrix} V_1 & V_2 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} V_1 \\ V_2 \end{pmatrix} + ...$$

3. **Hamiltonian** is explicit function of **M=R**²-rescaled velocity:

or: momentum
$$p$$
 $\mathbf{p} = \mathbf{M} \cdot \mathbf{v}$ or: $\begin{pmatrix} p_1 \\ p_2 \end{pmatrix} = \begin{pmatrix} m_1 & 0 \\ 0 & m_2 \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} = \begin{pmatrix} m_1 v_1 \\ m_2 v_2 \end{pmatrix}$

$$H(p_k...) = \frac{1}{2}(\frac{p_1^2}{m_1} + \frac{p_2^2}{m_2} + ...) = H(\mathbf{p}...) = \frac{1}{2}\mathbf{p} \cdot \mathbf{M}^{-1} \cdot \mathbf{p} + ... = \frac{1}{2}\begin{pmatrix} p_1 & p_2 \end{pmatrix} \begin{pmatrix} 1/m_1 & 0 \\ 0 & 1/m_2 \end{pmatrix} \begin{pmatrix} p_1 \\ p_2 \end{pmatrix} + ...$$

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Duality relations of Lagrangian and Hamiltonian ellipse

Introducing the 1^{st} (partial $\frac{\partial?}{\partial?}$) differential equations of mechanics

Introducing the (partial ³?) differential equations of mechanics

Starts out with simple demands for explicit-dependence, "loyalty" or "fealty to the colors"

Lagrangian and Estrangian have <u>no</u> explicit dependence on momentum p

$$\frac{\partial L}{\partial p_k} \equiv 0 \equiv \frac{\partial E}{\partial p_k}$$

Hamiltonian and Estrangian have <u>no</u> explicit dependence on velocity v

$$\frac{\partial \mathbf{H}}{\partial \mathbf{v}_{k}} \equiv 0 \equiv \frac{\partial \mathbf{E}}{\partial \mathbf{v}_{k}}$$

Lagrangian and Hamiltonian have <u>no</u> explicit dependence on speedinum V

$$\frac{\partial L}{\partial V_k} \equiv 0 \equiv \frac{\partial H}{\partial V_k}$$

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Introducing the (partial $\frac{\partial r}{\partial r}$) differential equations of mechanics

Starts out with simple demands for explicit-dependence, "loyalty" or "fealty to the colors"

Lagrangian and Estrangian have <u>no</u> explicit dependence on momentum p

$$\frac{\partial L}{\partial p_k} \equiv 0 \equiv \frac{\partial E}{\partial p_k}$$

Hamiltonian and Estrangian have <u>no</u> explicit dependence on velocity v

$$\frac{\partial \mathbf{H}}{\partial \mathbf{v}_{k}} \equiv 0 \equiv \frac{\partial \mathbf{E}}{\partial \mathbf{v}_{k}}$$

Lagrangian and Hamiltonian have <u>no</u> explicit dependence On speedinum V

$$\frac{\partial L}{\partial V_k} \equiv 0 \equiv \frac{\partial H}{\partial V_k}$$

Such non-dependencies hold in spite of "under-the-table" matrix and partial-differential connections

$$\nabla_{v} L = \frac{\partial L}{\partial \mathbf{v}} = \frac{\partial}{\partial \mathbf{v}} \frac{\mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v}}{2}$$
$$= \mathbf{M} \cdot \mathbf{v} = \mathbf{p}$$

$$\begin{pmatrix} \frac{\partial L}{\partial v_1} \\ \frac{\partial L}{\partial v_2} \end{pmatrix} = \begin{pmatrix} m_1 & 0 \\ 0 & m_2 \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} = \begin{pmatrix} p_1 \\ p_2 \end{pmatrix}$$

$$\nabla_{p} H = \mathbf{v} = \frac{\partial H}{\partial \mathbf{p}} = \frac{\partial}{\partial \mathbf{p}} \frac{\mathbf{p} \cdot \mathbf{M}^{-1} \cdot \mathbf{p}}{2}$$
$$= \mathbf{M}^{-1} \cdot \mathbf{p} = \mathbf{v}$$

$$\begin{pmatrix} \frac{\partial L}{\partial v_1} \\ \frac{\partial L}{\partial v_2} \end{pmatrix} = \begin{pmatrix} m_1 & 0 \\ 0 & m_2 \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} = \begin{pmatrix} p_1 \\ p_2 \end{pmatrix} \qquad \qquad \begin{pmatrix} \frac{\partial H}{\partial p_1} \\ \frac{\partial H}{\partial p_2} \end{pmatrix} = \begin{pmatrix} m_1^{-1} & 0 \\ 0 & m_2^{-1} \end{pmatrix} \begin{pmatrix} p_1 \\ p_2 \end{pmatrix} = \begin{pmatrix} v_1 \\ v_2 \end{pmatrix}$$

Introducing the (partial ³/₂) differential equations of mechanics

Starts out with simple demands for explicit-dependence, "loyalty" or "fealty to the colors"

Lagrangian and Estrangian have <u>no</u> explicit dependence on momentum p

$$\frac{\partial L}{\partial p_k} \equiv 0 \equiv \frac{\partial E}{\partial p_k}$$

Hamiltonian and Estrangian have <u>no</u> explicit dependence on velocity v

$$\frac{\partial \mathbf{H}}{\partial \mathbf{v}_{k}} \equiv 0 \equiv \frac{\partial \mathbf{E}}{\partial \mathbf{v}_{k}}$$

Lagrangian and Hamiltonian have <u>no</u> explicit dependence on speedinum V

$$\frac{\partial L}{\partial V_k} \equiv 0 \equiv \frac{\partial H}{\partial V_k}$$

Such non-dependencies hold in spite of "under-the-table" matrix and partial-differential connections

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$$= \mathbf{M} \cdot \mathbf{v} = \mathbf{p}$$

$$\begin{pmatrix} \frac{\partial L}{\partial v_1} \\ \frac{\partial L}{\partial v_2} \end{pmatrix} = \begin{pmatrix} m_1 & 0 \\ 0 & m_2 \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} = \begin{pmatrix} p_1 \\ p_2 \end{pmatrix}$$

Lagrange's 1st equation(s)

$$\frac{\partial L}{\partial v_k} = p_k$$
 or: $\frac{\partial L}{\partial \mathbf{v}} = \mathbf{p}$

$$\nabla_{p} H = \mathbf{v} = \frac{\partial H}{\partial \mathbf{p}} = \frac{\partial}{\partial \mathbf{p}} \frac{\mathbf{p} \cdot \mathbf{M}^{-1} \cdot \mathbf{p}}{2}$$
$$= \mathbf{M}^{-1} \cdot \mathbf{p} = \mathbf{v}$$

$$\begin{pmatrix}
\frac{\partial H}{\partial p_1} \\
\frac{\partial H}{\partial p_2}
\end{pmatrix} = \begin{pmatrix}
m_1^{-1} & 0 \\
0 & m_2^{-1}
\end{pmatrix} \begin{pmatrix}
p_1 \\
p_2
\end{pmatrix} = \begin{pmatrix}
v_1 \\
v_2
\end{pmatrix}$$

Hamilton's 1st equation(s)

$$\frac{\partial H}{\partial p_k} = v_k \quad \text{or:} \quad \frac{\partial H}{\partial \mathbf{p}} = \mathbf{v}$$

(Forget Estrangian for now)

