## Lecture 12 Thur. 2.25.2016

# Kepler Geometry of IHO (Isotropic Harmonic Oscillator) Elliptical Orbits (Ch. 8 and Ch. 9 of Unit 1)

"Sophomore-Physics-Earth" models: 3 key energy "steps" and 4 key energy "levels"
"Crushed-Earth" models: 3 key energy "steps" and 4 key energy "levels"
Earth matter vs nuclear matter vs Schwarzchild singular matter:
Introducing the "neutron starlet" ("fingertip physics")
Fantasizing a completely crushed "Black-Hole-Earth"
Introducing Isotropic Harmonic Oscillator (IHO) energy and frequency relations
Constructing 2D-IHO orbits using Kepler anomaly plots
Mean-anomaly and eccentric-anomaly geometry with web-app animation
Calculus and vector geometry of IHO orbits
Constructing 2D-IHO orbits using orbital phasor-clock plots
Phasor geometry of coordinate (x,y) and velocity (V<sub>x</sub>, V<sub>y</sub>) space with web-app animation

Kepler "laws" (Some that apply to all central (isotropic) F(r) force fields)(Derived here)Angular momentum invariance of IHO:  $F(r) = -k \cdot r$  with  $U(r) = k \cdot r^2/2$ (Derived here)Angular momentum invariance of Coulomb:  $F(r) = -GMm/r^2$  with U(r) = -GMm/r(Derived later)Total energy E = KE + PE invariance of IHO:  $F(r) = -k \cdot r$ (Derived here)Total energy E = KE + PE invariance of Coulomb:  $F(r) = -GMm/r^2$ (Derived later)Total energy E = KE + PE invariance of Coulomb:  $F(r) = -GMm/r^2$ (Derived later)

 $Link \Rightarrow \underline{BoxIt \ simulation \ of \ IHO \ orbits}$   $\underline{Link} \rightarrow \underline{IHO \ orbital \ time \ rates \ of \ change}$   $\underline{Link} \rightarrow \underline{IHO \ Exegesis \ Plot}$ 

## Geometry of idealized "Sophomore-physics Earth"

Coulomb field <u>outside</u>Isotropic Harmonic Oscillator (IHO) field <u>inside</u>Contact-geometry of potential curve(s)

"Crushed-Earth" models: 3 key energy "steps" and 4 key energy "levels" Earth matter vs nuclear matter:

Introducing the "neutron starlet" and "Black-Hole-Earth"





Thursday, March 3, 2016

















Thursday, March 3, 2016



















"Sophomore-Physics-Earth" models: 3 key energy "steps" and 4 key energy "levels" "Crushed-Earth" models: 3 key energy "steps" and 4 key energy "levels" Earth matter vs nuclear matter vs Schwarzchild singular matter: Introducing the "neutron starlet" ("fingertip physics") Fantasizing a completely crushed "Black-Hole-Earth"



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*Earth matter* Earth mass:  $M_{\oplus} = 5.9722 \times 10^{24} kg \approx 6.0 \cdot 10^{24} kg$ . Density  $\rho_{\oplus} = ??$ Earth radius:  $R_{\oplus} = 6.371 \cdot 10^6 m \approx 6.4 \cdot 10^6 m$  Earth volume:  $(4\pi/3)R_{\oplus}^3 \approx 4 \cdot 262 \cdot 10^{18} \sim 10^{21} m^3$ 

 $(6.4)^3 \sim 262$  and  $(4\pi/3)^2 62 = 1089 \sim 10^3$ 

$$\rho_{\oplus} = \frac{M_{\oplus}}{(4\pi/3)R_{\oplus}^3}$$

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Density of solid Fe= $7.9 \cdot 10^3$ kg/m<sup>3</sup> Density of liquid Fe= $6.9 \cdot 10^3$ kg/m<sup>3</sup>

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Nuclear matter Nucleon mass =  $1.67 \cdot 10^{-27} kg \sim 2 \cdot 10^{-27} kg$ . ("fingertip physics") Say a nucleus of atomic weight 50 has a radius of 3 fm, or 50 nucleons each with a mass  $2 \cdot 10^{-27} kg$ .

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Nuclear matter Nucleon mass =  $1.67 \cdot 10^{-27} kg$ . ~  $2 \cdot 10^{-27} kg$ . ("fingertip physics") Say a nucleus of atomic weight 50 has a radius of 3 fm, or 50 nucleons each with a mass  $2 \cdot 10^{-27} kg$ .  $4\pi/_3 3^3 = 36\pi = 113 \sim 10^2$ That's  $100 \cdot 10^{-27} = 10^{-25} kg$  packed into a volume of  $4\pi/_3 r^3 = 4\pi/_3 (3 \cdot 10^{-15})^3 m^3$  or about  $10^{-43} m^3$ .

$$\rho_{\oplus} = \frac{M_{\oplus}}{(4\pi/3)R_{\oplus}^3}$$

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Earth radius crushed by a factor of  $0.5 \cdot 10^{-5}$  to  $R_{crush\oplus} \simeq 300m$  would approach neutron-star density.



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Introducing the "Neutron starlet" 1 cm<sup>3</sup> of nuclear matter: mass =  $10^{12}$  kg.

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#### Introducing the "Black Hole Earth" Suppose Earth is crushed so that its



surface escape velocity is the speed of light  $c \cong 3.0.10^8$  m/s.  $c \equiv 299,792,458$  m/s (EXACTLY)  $V_{escape} = \sqrt{(2GM/R_{\otimes})}$ 

(from p. 49)  $G = 6.67384(80) \cdot 10^{-11} Nm^2/C^2 \sim (2/3)10^{-10}$  "Sophomore-Physics-Earth" models: 3 key energy "steps" and 4 key energy "levels"
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Introducing Isotropic Harmonic Oscillator (IHO) energy and frequency relations Constructing 2D-IHO orbits using Kepler anomaly plots Mean-anomaly and eccentric-anomaly geometry with web-app animation Calculus and vector geometry of IHO orbits
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Isotropic Harmonic Oscillator makes balls in parallel tunnel track each other



Isotropic Harmonic Oscillator makes balls in parallel tunnels track each other...



...even if track length is just  $g = 1m \text{ so } d \sim (1/12) \text{ micron}$ 

They all take about 84 minutes to go from right to left and back, again.

#### Isotropic Harmonic Oscillator makes balls in parallel tunnels track each other...



...even if track length is just g = 1m so d = (1/12) micron

The all take about 84 minutes to go from right to left and back, again.

Most neutron starlet (Corbits are centered ellipses



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Unit 1

Fig. 9.1

(top 2/3's)



Introducing Isotropic Harmonic Oscillator (IHO) energy and frequency relations
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(a) Orbits  
Calculus of IHO orbits  
To make velocity vector **v**  
just rotate by 
$$\pi/2$$
 or 90°  
the mean-anomaly  $\phi$  of position vector **r**  
reducing  
radius vector :  $\mathbf{r} = \begin{pmatrix} r_x \\ r_y \end{pmatrix} = \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} a \cos \omega t \\ b \sin \omega t \end{pmatrix} = \begin{pmatrix} a \cos \omega t \\ b \sin \omega t \end{pmatrix} = \begin{pmatrix} a \cos (\phi + \frac{\pi}{2}) \\ b \sin (\phi + \frac{\pi}{2}) \end{pmatrix}$  (for  $\omega = 1$ )  
mean-anomaly  $\phi$  of position vector **r**  
rotated by  $\pi/2$  or 90° is m.a. of vector **v**  
velocity vector :  $\mathbf{v} = \begin{pmatrix} v_x \\ v_y \end{pmatrix} = \begin{pmatrix} -a\omega \sin \omega t \\ b\omega \cos \omega t \end{pmatrix} = \frac{d\mathbf{r}}{dt} = \dot{\mathbf{r}} = \begin{pmatrix} a \cos (\phi + \frac{\pi}{2}) \\ b \sin (\phi + \frac{\pi}{2}) \end{pmatrix}$  (for  $\omega = 1$ )  
 $another \pi/2$  is m.a. of vector **v** rotated by  
another \pi/2 is m.a. of vector **a**  
acceleration or force vector :  $\frac{\mathbf{F}}{m} = \mathbf{a} = \begin{pmatrix} a_x \\ a_y \end{pmatrix} = \begin{pmatrix} -a\omega^2 \cos \omega t \\ -b\omega^2 \sin \omega t \end{pmatrix} = \frac{d\mathbf{v}}{dt} = \dot{\mathbf{v}} = \ddot{\mathbf{r}} = \frac{d^2\mathbf{r}}{dt^2} = \begin{pmatrix} a \cos(\phi + \frac{2\pi}{2}) \\ b \sin(\phi + \frac{2\pi}{2}) \end{pmatrix}$  ...and so forth...  
jerk or change of acceleration :  $\mathbf{j} = \begin{pmatrix} J_x \\ J_y \end{pmatrix} = \begin{pmatrix} +a\omega^3 \sin \omega t \\ -b\omega^3 \cos \omega t \end{pmatrix} = \frac{d\mathbf{a}}{dt} = \dot{\mathbf{a}} = \ddot{\mathbf{v}} = \ddot{\mathbf{r}} = \frac{d^3\mathbf{r}}{dt^3} = \begin{pmatrix} a \cos(\phi + \frac{3\pi}{2}) \\ b \sin(\phi + \frac{3\pi}{2}) \end{pmatrix}$  ...and so forth...

(a) Orbits  
Calculus of IHO orbits  
To make velocity vector **v**  
just rotate by 
$$\pi/2$$
 or 90°  
the mean-anomaly  $\phi$  of position vector **r**  
 $r(0)$   
 $r$ 

(a) Orbits  
To make velocity vector v  
istrotate by 
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the mean-anomaly  $\phi$  of position vector r  
 $r(0)$   
 $r(0)$   



*Geometry of Kepler anomalies for vectors*  $[\mathbf{r}(\phi), \mathbf{v}(\phi), \mathbf{a}(\phi), \mathbf{j}(\phi), ]$  *in coordinate (x,y) space rendered by animation web-apps BoxIt and RelaWavity.* 



Geometry of Kepler anomalies for vectors  $[\mathbf{r}(\phi), \mathbf{v}(\phi)]$  in coordinate (x, y) space rendered by animation web-apps BoxIt and RelaWavity described below after p.70.

http://www.uark.edu/ua/modphys/testing/markup/RelaWavityWeb.html?plotType=112

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Geometry of Kepler anomalies for vectors  $[\mathbf{r}(\phi), \mathbf{v}(\phi)]$  in coordinate (x, y) space rendered by animation web-apps BoxIt and RelaWavity described below after p.70.

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Geometry of dual ellipse Kepler anomalies for vectors  $[\mathbf{r}(\phi), \mathbf{p}(\phi)]$  and  $d_{dt}[\mathbf{r}(\phi), \mathbf{p}(\phi), ]$  in coordinate (x, y) space rendered by animation web-app in RelaWavity and described in Lect. 12-advanced.



Geometry of Kepler anomalies for vectors  $[\mathbf{r}(\phi)]$  in coordinate (x,y) space and 2-particle  $(x_1,x_2)$  space rendered by animation web-apps BoxIt.

**BoxIt Web Simulation** 



Geometry of Kepler anomalies for vectors  $[\mathbf{r}(\phi), \mathbf{v}(\phi), \mathbf{a}(\phi), \mathbf{j}(\phi), \mathbf{j}$ 



Geometry of vectors  $[\mathbf{r}(\phi), \mathbf{p}(\phi)]$  and quantum spin S-space BoxIt Web Simulation - B-Type Motion and 2-particle  $(x_1, x_2)$  space rendered by animation web-apps BoxIt.

Thursday, March 3, 2016

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<u>RelaWavity Web Simulation</u> <u>Ellipsometry</u>

BoxIt simulation of U(2) orbits http://www.uark.edu/ua/modphys/markup/BoxItWeb.html

















Review of IHO orbital phasor "clock" dynamics in uniform-body with "home-made movie" examples



























Thursday, March 3, 2016







# Kepler "laws" (Some that apply to all central (isotropic) F(r) force fields)Angular momentum invariance of IHO: $F(r)=-k \cdot r$ with $U(r)=k \cdot r^2/2$ (Derived here)Angular momentum invariance of Coulomb: $F(r)=-GMm/r^2$ with U(r)=-GMm/r(Derived in Unit 5)Total energy E=KE+PE invariance of IHO: $F(r)=-k \cdot r$ (Derived here)Total energy E=KE+PE invariance of Coulomb: $F(r)=-GMm/r^2$ (Derived in Unit 5)(Derived here)(Derived in Unit 5)Total energy E=KE+PE invariance of Coulomb: $F(r)=-GMm/r^2$ (Derived in Unit 5)



1. Area of triangle  $\measuredangle_{\mathbf{r}}^{\mathbf{v}} = \mathbf{r} \times \mathbf{v}/2$  is constant

$$\mathbf{r} \times \mathbf{v} = r_x v_y - r_y v_x = a \cos \omega t \cdot (b\omega \cos \omega t) - b \sin \omega t \cdot (-a\omega \sin \omega t) = ab \cdot \omega (\cos^2 \omega t + \sin^2 \omega t)$$

$$\begin{pmatrix} r_x \\ r_y \end{pmatrix} = \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} a \cos \omega t \\ b \sin \omega t \end{pmatrix}$$

$$\begin{pmatrix} v_x \\ v_y \end{pmatrix} = \begin{pmatrix} -a\omega \sin \omega t \\ b\omega \cos \omega t \end{pmatrix}$$



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for IHO

#### 2. Angular momentum $\mathbf{L} = m\mathbf{r} \times \mathbf{v}$ is conserved

$$L = m |\mathbf{r} \times \mathbf{v}| = m \left( r_x v_y - r_y v_x \right) = m \cdot ab \cdot \omega$$
 for IHO

$$|\mathbf{r} \times \mathbf{v}| = r \cdot v \cdot \sin \lambda_r \quad \mathbf{v} \quad \mathbf{r}$$
$$|\mathbf{r} \cdot \mathbf{v}| = r \cdot v \cdot \cos \lambda_r$$

Some Kepler's "laws" that apply to any central (isotropic) force F(r)...and certainly apply to the IHO:  $F(r) = -k \cdot r$  with  $U(r) = k \cdot r^2/2$  (Recall p.19:  $k = G \frac{4\pi}{3} m \rho_{\oplus}$ ) Unit 1 t = 0  $b r v = a \odot 0$   $t = \pi/3 \odot r$   $t = \pi/2 \odot r$   $r v = b \odot 0$ 1. Area of triangle  $\measuredangle_r^v = \mathbf{r} \times \mathbf{v}/2$  is constant  $\mathbf{r} \times \mathbf{v} = r_x v_y - r_y v_x = a \cos \omega t \cdot (b\omega \cos \omega t) - a \sin \omega t \cdot (-b\omega \sin \omega t) = ab \cdot \omega$  for IHO 2. Angular momentum  $\mathbf{L} = m\mathbf{r} \times \mathbf{v}$  is conserved

$$L = m |\mathbf{r} \times \mathbf{v}| = m \left( r_x v_y - r_y v_x \right) = m \cdot ab \cdot \omega$$
 for IHO

3. Equal area is swept by radius vector in each equal time interval T

$$A_{T} = \int_{0}^{T} \frac{\mathbf{r} \times d\mathbf{r}}{2} = \int_{0}^{T} \frac{\mathbf{r} \times \frac{d\mathbf{r}}{dt}}{2} dt = \int_{0}^{T} \frac{\mathbf{r} \times \mathbf{v}}{2} dt = \frac{L}{2m} \int_{0}^{T} dt = \frac{L}{2m} T$$

$$for IHO$$

$$|by 2| = r \cdot dr \cdot \sin \frac{dr}{dr}$$

Some Kepler's "laws" that apply to any central (isotropic) force F(r)...and certainly apply to the IHO:  $F(r) = -k \cdot r$  with  $U(r) = k \cdot r^2/2$  (Recall p.19:  $k = G \frac{4\pi}{3} m \rho_{\oplus}$ ) Unit 1 t = 0  $b r = \pi/3 \omega$   $t = \pi/3 \omega$   $r = \pi/2 \omega$   $r = b \omega$   $r = b \omega$ 

1. Area of triangle  $\measuredangle_{\mathbf{r}}^{\mathbf{v}} = \mathbf{r} \times \mathbf{v}/2$  is constant

 $\mathbf{r} \times \mathbf{v} = r_x v_y - r_y v_x = a \cos \omega t \cdot (b\omega \cos \omega t) - a \sin \omega t \cdot (-b\omega \sin \omega t) = ab \cdot \omega$ 2. Angular momentum  $L = m\mathbf{r} \times \mathbf{v}$  is conserved

$$L = m\mathbf{r} \times \mathbf{v} = m\left(r_x v_y - r_y v_x\right) = m \cdot ab \cdot \boldsymbol{\omega} = m \cdot ab \cdot \frac{2\pi}{\tau}$$
 for IHO

3. Equal area is swept by radius vector in each equal time interval T

$$A_{T} = \int_{0}^{T} \frac{\mathbf{r} \times d\mathbf{r}}{2} = \int_{0}^{T} \frac{\mathbf{r} \times \frac{d\mathbf{r}}{dt}}{2} dt = \int_{0}^{T} \frac{\mathbf{r} \times \mathbf{v}}{2} dt = \frac{L}{2m} \int_{0}^{T} dt = \frac{L}{2m} T$$
 for IHO  
In one period:  $\tau = \frac{1}{v} = \frac{2\pi}{\omega} = \frac{2mA_{\tau}}{L}$  the area is:  $A_{\tau} = \frac{L\tau}{2m} (= ab \cdot \pi \text{ for ellipse orbit})$ 

Some Kepler's "laws" that apply to any central (isotropic) force F(r)...and certainly apply to the IHO:  $F(r) = -k \cdot r$  with  $U(r) = k \cdot r^2/2$ (Recall p.19:  $k=G\frac{\pi n}{3}m\rho_{\oplus}$ ) Unit 1 Fig. 9.8 t = $t = \pi/3\omega$  $t = \pi/2\omega$  $v = a \omega$ b  $v=b \omega$ 

1. Area of triangle  $\measuredangle_{\mathbf{r}}^{\mathbf{v}} = \mathbf{r} \times \mathbf{v}/2$  is constant

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In one period:  $\tau = \frac{1}{\upsilon} = \frac{2\pi}{\omega} = \frac{2mA_{\tau}}{L}$  the area is:  $A_{\tau} = \frac{L\tau}{2m} (= ab \cdot \pi \text{ for ellipse orbit})$ 

(Recall from Lecture 7:  $\omega = \sqrt{k/m} = \sqrt{G\rho_{\oplus} 4\pi/3}$ )

(1)

Kepler "laws" (Some that apply to all central (isotropic) F(r) force fields)<br/>
Angular momentum invariance of IHO:  $F(r)=-k\cdot r$  with  $U(r)=k\cdot r^2/2$  (Derived here)Angular momentum invariance of Coulomb:  $F(r)=-GMm/r^2$  with U(r)=-GMm/r (Derived in Unit 5)<br/>
Total energy E=KE+PE invariance of IHO:  $F(r)=-k\cdot r$  (Derived here)<br/>
Total energy E=KE+PE invariance of Coulomb:  $F(r)=-GMm/r^2$  (Derived in Unit 5)

Some Kepler's "laws" that apply to any central (isotropic) force F(r)Apply to IHO:  $F(r)=-k \cdot r$  with  $U(r)=k \cdot r^2/2$  and Coulomb:  $F(r)=-GMm/r^2$  with  $U(r)=-GMm \cdot r$ 



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 $m \cdot a^{-1/2} b \sqrt{GM_{\oplus}}$ 

any central

F(r)

IHO and

Coulomb
Some Kepler's "laws" that apply to any central (isotropic) force F(r)Apply to IHO:  $F(r) = -k \cdot r$  with  $U(r) = k \cdot r^2/2$  and Coulomb:  $F(r) = -GMm/r^2$  with  $U(r) = -GMm \cdot r$ 



Kepler "laws" (Some that apply to all central (isotropic) F(r) force fields)<br/>
Angular momentum invariance of IHO:  $F(r) = -k \cdot r$  with  $U(r) = k \cdot r^2/2$ <br/>
Angular momentum invariance of Coulomb:  $F(r) = -GMm/r^2$  with U(r) = -GMm/r(Derived here)<br/>
(Derived in Unit 5)<br/>
(Derived here)<br/>
Total energy E = KE + PE invariance of Coulomb:  $F(r) = -GMm/r^2$ (Derived in Unit 5)<br/>
(Derived here)<br/>
(Derived here)<br/>
(Derived in Unit 5)

Kepler laws involve  $\measuredangle$ -momentum conservation in isotropic force F(r)Now consider orbital energy conservation of the IHO:  $F(r)=-k \cdot r$  with  $U(r)=k \cdot r^2/2$ Total energy=KE + PE is constant

$$KE + PE = \frac{1}{2}\mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + \frac{1}{2}\mathbf{r} \cdot \mathbf{K} \cdot \mathbf{r}$$

$$= \frac{1}{2} \begin{pmatrix} v_x & v_y \\ v_x \end{pmatrix} \cdot \begin{pmatrix} m & 0 \\ 0 & m \end{pmatrix} \cdot \begin{pmatrix} v_x \\ v_y \end{pmatrix} + \begin{pmatrix} r_x & r_y \\ r_x \end{pmatrix} \cdot \begin{pmatrix} k & 0 \\ 0 & k \end{pmatrix} \cdot \begin{pmatrix} r_x \\ r_y \end{pmatrix}$$

$$= \frac{1}{2}mv_x^2 + \frac{1}{2}mv_y^2 + \frac{1}{2}kr_x^2 + \frac{1}{2}kr_y^2$$

$$= \frac{1}{2}m(-a\omega\sin\omega t)^2 + \frac{1}{2}m(b\omega\cos\omega t)^2 + \frac{1}{2}k(a\cos\omega t)^2 + \frac{1}{2}k(b\sin\omega t)^2$$

$$\begin{pmatrix} v_x \\ v_y \end{pmatrix} = \begin{pmatrix} -a\omega\sin\omega t \\ b\omega\cos\omega t \end{pmatrix}$$

$$\begin{pmatrix} r_x \\ r_y \end{pmatrix} = \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} a\cos\omega t \\ b\sin\omega t \end{pmatrix}$$

Kepler laws involve  $\measuredangle$ -momentum conservation in isotropic force F(r)Now consider orbital energy conservation of the IHO:  $F(r) = -k \cdot r$  with  $U(r) = k \cdot r^2/2$ Total IHO energy=KE + PE is constant

$$KE + PE = \frac{1}{2}\mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + \frac{1}{2}\mathbf{r} \cdot \mathbf{K} \cdot \mathbf{r}$$

$$= \frac{1}{2} \begin{pmatrix} v_x & v_y \end{pmatrix} \bullet \begin{pmatrix} m & 0 \\ 0 & m \end{pmatrix} \bullet \begin{pmatrix} v_x \\ v_y \end{pmatrix} + \begin{pmatrix} r_x & r_y \end{pmatrix} \bullet \begin{pmatrix} k & 0 \\ 0 & k \end{pmatrix} \bullet \begin{pmatrix} r_x \\ r_y \end{pmatrix}$$

$$= \frac{1}{2}mv_x^2 + \frac{1}{2}mv_y^2 + \frac{1}{2}kr_x^2 + \frac{1}{2}kr_y^2$$

$$= \frac{1}{2}m(-a\omega\sin\omega t)^2 + \frac{1}{2}m(b\omega\cos\omega t)^2 + \frac{1}{2}k(a\cos\omega t)^2 + \frac{1}{2}k(b\sin\omega t)^2$$

$$= \frac{1}{2}ma^2\omega^2(\sin^2\omega t) + \frac{1}{2}mb^2\omega^2(\cos^2\omega t)^2 + \frac{1}{2}ka^2(\cos^2\omega t) + \frac{1}{2}kb^2(\sin^2\omega t)$$

$$= \frac{1}{2}m\omega^2(a^2 + b^2) \qquad \text{Given } : k = m\omega^2$$

Kepler laws involve  $\measuredangle$ -momentum conservation in isotropic force F(r)Now consider orbital energy conservation of the IHO:  $F(r) = -k \cdot r$  with  $U(r) = k \cdot r^2/2$ Total IHO energy=KE + PE is constant

$$\begin{split} & KE + PE = \frac{1}{2} \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + \frac{1}{2} \mathbf{r} \cdot \mathbf{K} \cdot \mathbf{r} \\ &= \frac{1}{2} \begin{pmatrix} v_x & v_y \end{pmatrix} \bullet \begin{pmatrix} m & 0 \\ 0 & m \end{pmatrix} \bullet \begin{pmatrix} v_x \\ v_y \end{pmatrix} + \begin{pmatrix} r_x & r_y \end{pmatrix} \bullet \begin{pmatrix} k & 0 \\ 0 & k \end{pmatrix} \bullet \begin{pmatrix} r_x \\ r_y \end{pmatrix} \\ &= \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 + \frac{1}{2} k r_x^2 + \frac{1}{2} k r_y^2 \\ &= \frac{1}{2} m (-a\omega \sin \omega t)^2 + \frac{1}{2} m (b\omega \cos \omega t)^2 + \frac{1}{2} k (a \cos \omega t)^2 + \frac{1}{2} k (b \sin \omega t)^2 \\ &= \frac{1}{2} m a^2 \omega^2 (\sin^2 \omega t) + \frac{1}{2} m b^2 \omega^2 (\cos^2 \omega t)^2 + \frac{1}{2} k a^2 (\cos^2 \omega t) + \frac{1}{2} k b^2 (\sin^2 \omega t) \\ &= \frac{1}{2} m \omega^2 (a^2 + b^2) \qquad Given : k = m \omega^2 \\ &= KE + PE = \frac{1}{2} m \omega^2 (a^2 + b^2) = \frac{1}{2} k (a^2 + b^2) \quad \text{since: } \omega = \sqrt{\frac{k}{m}} = \sqrt{G\rho_{\oplus} 4\pi/3} \quad \text{or: } m \omega^2 = k \end{split}$$

Kepler "laws" (Some that apply to all central (isotropic) F(r) force fields)<br/>Angular momentum invariance of IHO:  $F(r)=-k \cdot r$  with  $U(r)=k \cdot r^2/2$ (Derived here)<br/>(Derived here)<br/>(Derived in Unit 5)<br/>(Derived here)Angular momentum invariance of Coulomb:  $F(r)=-GMm/r^2$  with U(r)=-GMm/r(Derived in Unit 5)<br/>(Derived here)Total energy E=KE+PE invariance of Coulomb:  $F(r)=-GMm/r^2$ (Derived in Unit 5)<br/>(Derived here)Total energy E=KE+PE invariance of Coulomb:  $F(r)=-GMm/r^2$ (Derived in Unit 5)<br/>(Derived in Unit 5)

Kepler laws involve  $\measuredangle$ -momentum conservation in isotropic force F(r)Now consider orbital energy conservation of the IHO:  $F(r)=-k \cdot r$  with  $U(r)=k \cdot r^2/2$ Total IHO energy=KE + PE is constant

$$\begin{aligned} KE + PE &= \frac{1}{2} \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + \frac{1}{2} \mathbf{r} \cdot \mathbf{K} \cdot \mathbf{r} \\ &= \frac{1}{2} \begin{pmatrix} v_x & v_y \end{pmatrix} \cdot \begin{pmatrix} m & 0 \\ 0 & m \end{pmatrix} \cdot \begin{pmatrix} v_x \\ v_y \end{pmatrix} + \begin{pmatrix} r_x & r_y \end{pmatrix} \cdot \begin{pmatrix} k & 0 \\ 0 & k \end{pmatrix} \cdot \begin{pmatrix} r_x \\ r_y \end{pmatrix} \\ &= \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 + \frac{1}{2} k r_x^2 + \frac{1}{2} k r_y^2 \\ &= \frac{1}{2} m (-a\omega \sin \omega t)^2 + \frac{1}{2} m (b\omega \cos \omega t)^2 + \frac{1}{2} k (a \cos \omega t)^2 + \frac{1}{2} k (b \sin \omega t)^2 \\ &= \frac{1}{2} m a^2 \omega^2 (\sin^2 \omega t) + \frac{1}{2} m b^2 \omega^2 (\cos^2 \omega t)^2 + \frac{1}{2} k a^2 (\cos^2 \omega t) + \frac{1}{2} k b^2 (\sin^2 \omega t) \\ &= \frac{1}{2} m \omega^2 (a^2 + b^2) = \frac{1}{2} k (a^2 + b^2) \quad \text{since: } \omega = \sqrt{\frac{k}{m}} = \sqrt{G \rho_{\oplus} 4\pi/3} \quad \text{or: } m \omega^2 = k \end{aligned}$$
We'll see that the Coul, orbits are simpler: (*like the period...*not a function of b)

Kepler laws involve  $\measuredangle$ -momentum conservation in isotropic force F(r)Now consider orbital energy conservation of the IHO:  $F(r) = -k \cdot r$  with  $U(r) = k \cdot r^2/2$ Total IHO energy=KE + PE is constant

$$\begin{split} & KE + PE = \frac{1}{2} \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + \frac{1}{2} \mathbf{r} \cdot \mathbf{K} \cdot \mathbf{r} \\ &= \frac{1}{2} \begin{pmatrix} v_x & v_y \end{pmatrix} \mathbf{\bullet} \begin{pmatrix} m & 0 \\ 0 & m \end{pmatrix} \mathbf{\bullet} \begin{pmatrix} v_x \\ v_y \end{pmatrix} \mathbf{+} \begin{pmatrix} r_x & r_y \end{pmatrix} \mathbf{\bullet} \begin{pmatrix} k & 0 \\ 0 & k \end{pmatrix} \mathbf{\bullet} \begin{pmatrix} r_x \\ r_y \end{pmatrix} \\ &= \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 + \frac{1}{2} k r_x^2 + \frac{1}{2} k r_y^2 \\ &= \frac{1}{2} m (-a\omega \sin \omega t)^2 + \frac{1}{2} m (b\omega \cos \omega t)^2 + \frac{1}{2} k (a \cos \omega t)^2 + \frac{1}{2} k (b \sin \omega t)^2 \\ &= \frac{1}{2} m a^2 \omega^2 (\sin^2 \omega t) + \frac{1}{2} m b^2 \omega^2 (\cos^2 \omega t)^2 + \frac{1}{2} k a^2 (\cos^2 \omega t) + \frac{1}{2} k b^2 (\sin^2 \omega t) \\ &= \frac{1}{2} m \omega^2 (a^2 + b^2) = \frac{1}{2} k (a^2 + b^2) \quad \text{since: } \omega = \sqrt{\frac{k}{m}} = \sqrt{G\rho_{\oplus} 4\pi / 3} \quad \text{or: } m\omega^2 = k \\ \text{We'll see that the Coul. orbits are simpler: } (like the period...not a function of b) \\ &E = KE + PE = \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 - \frac{k}{r} = \frac{1}{2} m v_x^2 + \frac{1}{2} m v_y^2 - \frac{GM_{\oplus}m}{r} = -\frac{GM_{\oplus}m}{a} \end{split}$$

Thursday, March 3, 2016



