Lecture 9 Tue. 9.18.2012

## Geometry of Dual Quadratic Forms: Lagrange vs Hamilton (Ch. 11 and Ch. 12 of Unit 1)

## Introduction to dual matrix operator geometry

Review of dual IHO elliptic orbits (Lecture 7-8) Construction by Phasor-pair projection Construction by Kepler anomaly projection Operator geometric sequences and eigenvectors Rescaled description of matrix operator geometry Vector calculus of tensor operation

# Introduction to Lagrangian-Hamiltonian duality

Review of partial differential relations Chain rule and order symmetry Duality relations of Lagrangian and Hamiltonian ellipse Introducing the 1<sup>st</sup> (partial  $\frac{\partial?}{\partial?}$ ) differential equations of mechanics





# Quadratic forms and tangent contact geometry of their ellipses

A matrix *Q* that generates an ellipse by  $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r} = 1$  is called positive-definite (if  $\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}$  always >0)

$$\mathbf{r} \bullet \mathbf{Q} \bullet \mathbf{r} = 1$$

$$\begin{pmatrix} x & y \end{pmatrix} \bullet \begin{pmatrix} \frac{1}{a^2} & 0 \\ 0 & \frac{1}{b^2} \end{pmatrix} \bullet \begin{pmatrix} x \\ y \end{pmatrix} = 1 = \begin{pmatrix} x & y \end{pmatrix} \bullet \begin{pmatrix} \frac{x}{a^2} \\ \frac{y}{b^2} \end{pmatrix} = \frac{x^2}{a^2} + \frac{y^2}{b^2}$$

A inverse matrix  $Q^{-1}$  generates an ellipse by  $\mathbf{p} \cdot Q^{-1} \cdot \mathbf{p} = 1$  called inverse or dual ellipse:

$$p \bullet \mathbf{Q}^{-1} \bullet \mathbf{p} = 1$$

$$\begin{pmatrix} p_x & p_y \\ 0 & b^2 \end{pmatrix} \bullet \begin{pmatrix} p_x \\ p_y \end{pmatrix} = 1 = \begin{pmatrix} p_x & p_y \\ b^2 p_y \end{pmatrix} \bullet \begin{pmatrix} a^2 p_x \\ b^2 p_y \end{pmatrix} = a^2 p_x^2 + b^2 p_y^2$$

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Defined mapping between ellipses

A inverse matrix  $Q^{-1}$  generates an ellipse by  $\mathbf{p} \cdot Q^{-1} \cdot \mathbf{p} = 1$  called inverse or dual ellipse:

$$\begin{pmatrix} p_x & p_y \end{pmatrix} \bullet \begin{pmatrix} a^2 & 0 \\ 0 & b^2 \end{pmatrix} \bullet \begin{pmatrix} p_x \\ p_y \end{pmatrix} = 1 = \begin{pmatrix} p_x & p_y \\ p_x \end{pmatrix} \bullet \begin{pmatrix} a^2 p_x \\ b^2 p_y \end{pmatrix} = a^2 p_x^2 + b^2 p_y^2$$

# (a) Quadratic form ellipse and *Inverse quadratic form ellipse*

based on Unit 1 Fig. 11.6



#### (a) Quadratic form ellipse and Inverse quadratic form ellipse





#### (a) Quadratic form ellipse and Inverse quadratic form ellipse





$$\mathbf{p} = \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \bullet \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix} = \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \text{ where: } \begin{aligned} x = r_x = a\cos\phi = a\cos\phi t \\ y = r_y = b\sin\phi = b\sin\omega t \end{aligned} \text{ so: } \mathbf{p} \cdot \mathbf{r} = h$$



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$$\mathbf{p} = \mathbf{Q} \cdot \mathbf{r} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \bullet \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix} = \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \text{ where: } \begin{aligned} x = r_x = a\cos\phi = a\cos\phi t \\ y = r_y = b\sin\phi t \end{aligned} \text{ so: } \mathbf{p} \cdot \mathbf{r} = I \\ \mathbf{p} \text{ is perpendicular to velocity } \mathbf{v} = \mathbf{\dot{r}}, a \text{ mutual orthogonality} \\ \mathbf{\dot{r}} \bullet \mathbf{p} = 0 = \begin{pmatrix} \dot{r}_x & \dot{r}_y \end{pmatrix} \bullet \begin{pmatrix} p_x \\ p_y \end{pmatrix} = \begin{pmatrix} -a\sin\phi & b\cos\phi \end{pmatrix} \bullet \begin{pmatrix} (1/a)\cos\phi \\ (1/b)\sin\phi \end{pmatrix} \text{ where: } \begin{aligned} \dot{r}_x = -a\sin\phi \\ \dot{r}_y = b\cos\phi \end{aligned} \text{ and: } \begin{aligned} p_x = (1/a)\cos\phi \\ p_y = (1/b)\sin\phi \end{aligned}$$











Diagonal **R**-matrix acts on vector 
$$\mathbf{v}^{try}$$
.  
Resulting vector has slope changed by factor  $a/b = 2$ .  
 $\mathbf{R} \cdot \mathbf{v}^{xy} = \begin{pmatrix} 1/a & 0 \\ 0 & 1/b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a \\ y/b \end{pmatrix}$ .  
(It increases if  $a > b$ .)  
Diagonal ( $\mathbf{R}^2 = \mathbf{Q}$ )-matrix acts on vector  $\mathbf{v}^{try}$ .  
Resulting vector has slope changed by factor  $a^2/b^2 = 1$ .  
 $\mathbf{Q} \cdot \mathbf{v}^{xy} = \begin{pmatrix} Ma^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix}$ .  
Either process can go on forever...  
Diagonal ( $\mathbf{R}^{2n} = \mathbf{Q}^n$ )-matrix acts on vector  $\mathbf{v}^{try}$ .  
Resulting vector has slope changed by factor  $a^{2n}/b^2 = 1$ .  
Either process can go on forever...  
Diagonal ( $\mathbf{R}^{2n} = \mathbf{Q}^n$ )-matrix acts on vector  $\mathbf{v}^{try}$ .  
Resulting vector has slope changed by factor  $a^{2n}/b^{2n} = 4^n$ .  
...Finally, the result approaches *EIGENVECTOR*  $|x\rangle = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$   
of  $\infty$ -slope which is "immune" to  $\mathbf{R}$ ,  $\mathbf{Q}$  or  $\mathbf{Q}^n$ :  
 $\mathbf{R} |y| = (1/b)|y\rangle$   $\mathbf{Q}^n |y| = (1/b^2)^n |y\rangle$ 

Diagonal **R**-matrix acts on vector 
$$\mathbf{v}^{sy}$$
.  
Resulting vector has slope changed by factor  $a/b = 2$ .  
**R** •  $\mathbf{v}^{sy} = \begin{pmatrix} 1/a & 0 \\ 0 & 1/b \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a \\ y/b \end{pmatrix}$   
(It increases if  $a > b$ .)  
Diagonal ( $\mathbf{R}^2 = \mathbf{Q}$ )-matrix acts on vector  $\mathbf{v}^{sy}$ .  
Resulting vector has slope changed by factor  $a^2/b^2$   
 $\mathbf{Q} \cdot \mathbf{v}^{sy} = \begin{pmatrix} 1/a^2 & 0 \\ 0 & 1/b^2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x/a^2 \\ y/b^2 \end{pmatrix}$   
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of  $\infty$ -slope which is "immune" to  $\mathbf{R}$ ,  $\mathbf{Q}$  or  $\mathbf{Q}^n$ :  
 $\mathbf{R}|y| = (1/b)|y\rangle$   $\mathbf{Q}^n|y| = (1/b^2)^n|y\rangle$  *Eigensolution*  
 $\mathbf{R}^{-1}|x| = (a)|x\rangle$   $\mathbf{Q}^{-1}|x| = (a^2)^n|x\rangle$ 

















Derive matrix "normal-to-ellipse" geometry by vector calculus: Let matrix  $Q = \begin{pmatrix} A & B \\ B & D \end{pmatrix}$ define the ellipse  $1 = \mathbf{r} \cdot Q \cdot \mathbf{r} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A & B \\ B & D \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{pmatrix} = A \cdot x^2 + 2B \cdot xy + D \cdot y^2 = 1$ 



Derive matrix "normal-to-ellipse" geometry by vector calculus: Let matrix  $Q = \begin{pmatrix} A & B \\ B & D \end{pmatrix}$ define the ellipse  $1 = \mathbf{r} \cdot Q \cdot \mathbf{r} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A & B \\ B & D \end{pmatrix} \cdot \begin{pmatrix} x \\ y \end{pmatrix} = \begin{pmatrix} x & y \end{pmatrix} \cdot \begin{pmatrix} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{pmatrix} = A \cdot x^2 + 2B \cdot xy + D \cdot y^2 = 1$ 

*Compare operation by Q on vector* **r** 

with ve

$$\left(\begin{array}{cc} A & B \\ B & D \end{array}\right) \bullet \left(\begin{array}{c} x \\ y \end{array}\right) = \left(\begin{array}{c} A \cdot x + B \cdot y \\ B \cdot x + D \cdot y \end{array}\right)$$

vector derivative or gradient of 
$$\mathbf{r} \cdot Q \cdot \mathbf{r}$$
  
 $\frac{\partial}{\partial \mathbf{r}} (\mathbf{r} \cdot Q \cdot \mathbf{r}) = \nabla (\mathbf{r} \cdot Q \cdot \mathbf{r})$   
 $\begin{pmatrix} \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \end{pmatrix} (A \cdot x^2 + 2B \cdot xy + D \cdot y^2) = \begin{pmatrix} 2A \cdot x + 2B \cdot y \\ 2B \cdot x + 2D \cdot y \end{pmatrix}$ 



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*Compare operation by Q on vector* **r** 

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$$\frac{\partial}{\partial \mathbf{r}} (\mathbf{r} \cdot Q \cdot \mathbf{r}) = \nabla (\mathbf{r} \cdot Q \cdot \mathbf{r})$$

$$\begin{pmatrix} \frac{\partial}{\partial x} \\ \frac{\partial}{\partial y} \end{pmatrix} (A \cdot x^2 + 2B \cdot xy + D \cdot y^2) = \begin{pmatrix} 2A \cdot x + 2B \cdot y \\ 2B \cdot x + 2D \cdot y \end{pmatrix}$$

*Very simple result:* 

$$\frac{\partial}{\partial \mathbf{r}} \left( \frac{\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}}{2} \right) = \nabla \left( \frac{\mathbf{r} \cdot \mathbf{Q} \cdot \mathbf{r}}{2} \right) = \mathbf{Q} \cdot \mathbf{r}$$







# Introduction to Lagrangian-Hamiltonian duality

Review of partial differential relations Chain rule and order symmetry Duality relations of Lagrangian and Hamiltonian ellipse Introducing the 1<sup>st</sup> (partial  $\frac{\partial?}{\partial?}$ ) differential equations of mechanics

















![](_page_45_Figure_0.jpeg)

![](_page_46_Figure_0.jpeg)

![](_page_47_Figure_0.jpeg)

## Introduction to Lagrangian-Hamiltonian duality

Review of partial differential relations Chain rule and order symmetry Duality relations of Lagrangian and Hamiltonian ellipse Introducing the 1<sup>st</sup> (partial  $\frac{\partial}{\partial 2}$ ) differential equations of mechanics

$$f(x_1, y_1) = f(x_0, y_0) + \frac{\partial f}{\partial x}(x_0, y_0) \Delta x + \frac{\partial f}{\partial y}(x_0, y_0) \Delta y + \frac{\partial f}{\partial y} \frac{\partial f}{\partial x}(x_0, y_0) \Delta x \Delta y$$
$$= f(x_0, y_0) + \frac{\partial f}{\partial y}(x_0, y_0) \Delta y + \frac{\partial f}{\partial x}(x_0, y_0) \Delta x + \frac{\partial f}{\partial x} \frac{\partial f}{\partial y}(x_0, y_0) \Delta y \Delta x$$

$$f(x_1, y_1) = f(x_0, y_0) + \frac{\partial f}{\partial x}(x_0, y_0) \Delta x + \frac{\partial f}{\partial y}(x_0, y_0) \Delta y + \frac{\partial g}{\partial y} \frac{\partial f}{\partial x}(x_0, y_0) \Delta x \Delta y$$
$$= f(x_0, y_0) + \frac{\partial f}{\partial y}(x_0, y_0) \Delta y + \frac{\partial f}{\partial x}(x_0, y_0) \Delta x + \frac{\partial g}{\partial x} \frac{\partial f}{\partial y}(x_0, y_0) \Delta y \Delta x$$

1. Chain rules

$$\begin{bmatrix} f(x_1, y_1) - f(x_0, y_0) \end{bmatrix} = df = \frac{\partial f}{\partial x} (x_0, y_0) dx + \frac{\partial f}{\partial y} (x_0, y_0) dy \dots_{(keep \ 1^{st} - order \ terms \ only!)}$$
$$\frac{df}{dt} = \frac{\partial f}{\partial x} (x_0, y_0) \frac{dx}{dt} + \frac{\partial f}{\partial y} (x_0, y_0) \frac{dy}{dt}$$
$$\dot{f} = \frac{\partial f}{\partial x} \dot{x} + \frac{\partial f}{\partial y} \dot{y} \qquad (shorthand \ notation)$$

$$f(x_1, y_1) = f(x_0, y_0) + \frac{\partial f}{\partial x}(x_0, y_0) \Delta x + \frac{\partial f}{\partial y}(x_0, y_0) \Delta y + \frac{\partial f}{\partial y} \frac{\partial f}{\partial x}(x_0, y_0) \Delta x \Delta y$$
$$= f(x_0, y_0) + \frac{\partial f}{\partial y}(x_0, y_0) \Delta y + \frac{\partial f}{\partial x}(x_0, y_0) \Delta x + \frac{\partial f}{\partial x} \frac{\partial f}{\partial y}(x_0, y_0) \Delta y \Delta x$$

1. Chain rules

$$f(x_{1}, y_{1}) - f(x_{0}, y_{0})] = df = \frac{\partial f}{\partial x}(x_{0}, y_{0})dx + \frac{\partial f}{\partial y}(x_{0}, y_{0})dy \dots_{(keep \ 1^{st} - order \ terms \ only!)}$$
$$\frac{df}{dt} = \frac{\partial f}{\partial x}(x_{0}, y_{0})\frac{dx}{dt} + \frac{\partial f}{\partial y}(x_{0}, y_{0})\frac{dy}{dt}$$
$$\dot{f} = \frac{\partial f}{\partial x}\dot{x} + \frac{\partial f}{\partial y}\dot{y} \qquad (shorthand \ notation) = \partial_{x} f \dot{x}$$

. Symmetry of partial deriv. ordering  

$$\frac{\partial}{\partial y}\frac{\partial f}{\partial x} = \frac{\partial}{\partial x}\frac{\partial f}{\partial y} \quad \text{or:} \quad \frac{\partial^2 f}{\partial y \partial x} = \frac{\partial^2 f}{\partial x \partial y} \quad \text{or:} \quad \partial_y \partial_x f = \partial_x \partial_y f$$
(pay attention to the second s

 $(pay attention to the 2^{nd} - order terms, too!)$ 

 $+ \partial_{y} f \dot{y}$ 

(shorthand notation)

$$f(x_1, y_1) = f(x_0, y_0) + \frac{\partial f}{\partial x}(x_0, y_0) \Delta x + \frac{\partial f}{\partial y}(x_0, y_0) \Delta y + \frac{\partial f}{\partial y} \frac{\partial f}{\partial x}(x_0, y_0) \Delta x \Delta y$$
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1. Chain rules

2. Symmetry of partial deriv. ordering  

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(shorthand notation)

$$Let: \vec{\nabla} = \begin{pmatrix} \partial_x & \partial_y \end{pmatrix} \quad so: \vec{\nabla}f \cdot \mathbf{dr} = \begin{pmatrix} \partial_x f & \partial_y f \end{pmatrix} \cdot \begin{pmatrix} dx \\ dy \end{pmatrix} = \partial_x f \, dx + \partial_y f \, dy = df$$

Tuesday, September 18, 2012

 $(pay attention to the 2^{nd} - order terms, too!)$ 

## Introduction to Lagrangian-Hamiltonian duality Review of partial differential relations Chain rule and order symmetry → Duality relations of Lagrangian and Hamiltonian ellipse Introducing the 1<sup>st</sup> (partial <sup>∂?</sup>/<sub>∂2</sub>) differential equations of mechanics

Three ways to express energy: Consider kinetic energy (KE) first

1. Lagrangian is explicit function of velocity: 
$$\mathbf{v} = \begin{pmatrix} v_1 \\ v_2 \end{pmatrix}$$
  
 $L(v_k...) = \frac{1}{2} (m_1 v_1^2 + m_2 v_2^2 + ...) = L(\mathbf{v}...) = \frac{1}{2} \mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v} + ... = \frac{1}{2} \begin{pmatrix} v_1 & v_2 \end{pmatrix} \begin{pmatrix} m_1 & 0 \\ 0 & m_2 \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} + ...$ 

2. "Estrangian" is explicit function of **R**-rescaled velocity:  
(or l'Estrangian Or: "speedinum" 
$$V = \mathbf{R} \cdot \mathbf{v}$$
 or:  $\begin{pmatrix} V_1 \\ V_2 \end{pmatrix} = \begin{pmatrix} \sqrt{m_1} & 0 \\ 0 & \sqrt{m_2} \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix}$   
 $E(V_k \dots) = \frac{1}{2} (V_1^2 + V_2^2 + \dots) = E(\mathbf{V} \dots) = \frac{1}{2} \mathbf{V} \cdot \mathbf{1} \cdot \mathbf{V} + \dots = \frac{1}{2} \begin{pmatrix} V_1 & V_2 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \begin{pmatrix} V_1 \\ V_2 \end{pmatrix} + \dots$ 

3. Hamiltonian is explicit function of 
$$\mathbf{M} = \mathbf{R}^2$$
-rescaled velocity:  
or: momentum  $p$   $\mathbf{p} = \mathbf{M} \cdot \mathbf{v}$  or:  $\begin{pmatrix} p_1 \\ p_2 \end{pmatrix} = \begin{pmatrix} m_1 & 0 \\ 0 & m_2 \end{pmatrix} \begin{pmatrix} \mathbf{v}_1 \\ \mathbf{v}_2 \end{pmatrix} = \begin{pmatrix} m_1 \mathbf{v}_1 \\ m_2 \mathbf{v}_2 \end{pmatrix}$   
 $H(p_k \dots) = \frac{1}{2}(\frac{p_1^2}{m_1} + \frac{p_2^2}{m_2} + \dots) = H(\mathbf{p} \dots) = \frac{1}{2}\mathbf{p} \cdot \mathbf{M}^{-1} \cdot \mathbf{p} + \dots = \frac{1}{2}\begin{pmatrix} p_1 & p_2 \end{pmatrix} \begin{pmatrix} 1/m_1 & 0 \\ 0 & 1/m_2 \end{pmatrix} \begin{pmatrix} p_1 \\ p_2 \end{pmatrix} + \dots$ 

![](_page_55_Figure_0.jpeg)

 Introduction to Lagrangian-Hamiltonian duality

 Review of partial differential relations

 Chain rule and order symmetry

 Duality relations of Lagrangian and Hamiltonian ellipse

 Introducing the 1<sup>st</sup> (partial  $\frac{\partial^2}{\partial \gamma}$ ) differential equations of mechanics

# Introducing the (partial $\frac{\partial 2}{\partial 2}$ ) differential equations of mechanics

Starts out with simple demands for explicit-dependence, "loyalty" or "fealty to the colors"

*Lagrangian* and *Estrangian* have <u>no</u> explicit dependence on *momentum* **p** 

$$\frac{\partial \boldsymbol{L}}{\partial \boldsymbol{p}_k} \equiv 0 \equiv \frac{\partial \boldsymbol{E}}{\partial \boldsymbol{p}_k}$$

*Hamiltonian* and *Estrangian* have <u>no</u> explicit dependence on *velocity* v

$$\frac{\partial H}{\partial v_k} \equiv 0 \equiv \frac{\partial E}{\partial v_k}$$

Lagrangian and Hamiltonian have <u>no</u> explicit dependence ON speedinum V

$$\frac{\partial L}{\partial V_k} \equiv 0 \equiv \frac{\partial H}{\partial V_k}$$

# Introducing the (partial $\frac{\partial 2}{\partial 2}$ ) differential equations of mechanics

Starts out with simple demands for explicit-dependence, "loyalty" or "fealty to the colors"

Lagrangian and Estrangian have <u>no</u> explicit dependence on momentum **p**  *Hamiltonian* and *Estrangian* have <u>no</u> explicit dependence on *velocity* v

Lagrangian and Hamiltonian have <u>no</u> explicit dependence ON speedinum V

$$\frac{\partial L}{\partial p_k} \equiv 0 \equiv \frac{\partial E}{\partial p_k} \qquad \qquad \frac{\partial H}{\partial v_k} \equiv 0 \equiv \frac{\partial E}{\partial v_k} \qquad \qquad \frac{\partial L}{\partial V_k} \equiv 0 \equiv \frac{\partial H}{\partial V_k}$$

Such non-dependencies hold in spite of "under-the-table" matrix and partial-differential connections

$$\nabla_{v}L = \frac{\partial L}{\partial \mathbf{v}} = \frac{\partial}{\partial \mathbf{v}} \frac{\mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v}}{2} \qquad \nabla_{p}H = \mathbf{v} = \frac{\partial H}{\partial \mathbf{p}} = \frac{\partial}{\partial \mathbf{p}} \frac{\mathbf{p} \cdot \mathbf{M}^{-1} \cdot \mathbf{p}}{2} = \mathbf{M} \cdot \mathbf{v} = \mathbf{p} \qquad = \mathbf{M}^{-1} \cdot \mathbf{p} = \mathbf{v}$$
$$\begin{pmatrix} \frac{\partial L}{\partial v_{1}} \\ \frac{\partial L}{\partial v_{2}} \end{pmatrix} = \begin{pmatrix} m_{1} & 0 \\ 0 & m_{2} \end{pmatrix} \begin{pmatrix} v_{1} \\ v_{2} \end{pmatrix} = \begin{pmatrix} p_{1} \\ p_{2} \end{pmatrix} \qquad \begin{pmatrix} \frac{\partial H}{\partial p_{1}} \\ \frac{\partial H}{\partial p_{2}} \end{pmatrix} = \begin{pmatrix} m_{1}^{-1} & 0 \\ 0 & m_{2}^{-1} \end{pmatrix} \begin{pmatrix} p_{1} \\ p_{2} \end{pmatrix} = \begin{pmatrix} v_{1} \\ v_{2} \end{pmatrix}$$

# Introducing the (partial $\frac{\partial 2}{\partial 2}$ ) differential equations of mechanics

Starts out with simple demands for explicit-dependence, "loyalty" or "fealty to the colors"

Lagrangian and Estrangian have <u>no</u> explicit dependence on momentum **p**  *Hamiltonian* and *Estrangian* have <u>no</u> explicit dependence on *velocity* v

Lagrangian and Hamiltonian have <u>no</u> explicit dependence ON speedinum V

$$\frac{\partial L}{\partial p_k} \equiv 0 \equiv \frac{\partial E}{\partial p_k} \qquad \qquad \frac{\partial H}{\partial v_k} \equiv 0 \equiv \frac{\partial E}{\partial v_k} \qquad \qquad \frac{\partial L}{\partial V_k} \equiv 0 \equiv \frac{\partial H}{\partial V_k}$$

Such non-dependencies hold in spite of "under-the-table" matrix and partial-differential connections

$$\nabla_{v}L = \frac{\partial L}{\partial \mathbf{v}} = \frac{\partial}{\partial \mathbf{v}} \frac{\mathbf{v} \cdot \mathbf{M} \cdot \mathbf{v}}{2} \qquad \nabla_{p}H = \mathbf{v} = \frac{\partial H}{\partial \mathbf{p}} = \frac{\partial}{\partial \mathbf{p}} \frac{\mathbf{p} \cdot \mathbf{M}^{-1} \cdot \mathbf{p}}{2} \qquad (Forget Estrangian for now)$$
$$= \mathbf{M} \cdot \mathbf{v} = \mathbf{p} \qquad = \mathbf{M}^{-1} \cdot \mathbf{p} = \mathbf{v}$$
$$\begin{pmatrix} \frac{\partial L}{\partial v_{1}} \\ \frac{\partial L}{\partial v_{2}} \end{pmatrix} = \begin{pmatrix} m_{1} & 0 \\ 0 & m_{2} \end{pmatrix} \begin{pmatrix} v_{1} \\ v_{2} \end{pmatrix} = \begin{pmatrix} p_{1} \\ p_{2} \end{pmatrix} \qquad \begin{pmatrix} \frac{\partial H}{\partial p_{1}} \\ \frac{\partial H}{\partial p_{2}} \end{pmatrix} = \begin{pmatrix} m_{1}^{-1} & 0 \\ 0 & m_{2}^{-1} \end{pmatrix} \begin{pmatrix} p_{1} \\ p_{2} \end{pmatrix} = \begin{pmatrix} v_{1} \\ v_{2} \end{pmatrix} \qquad Hamilton's 1^{st} equation(s)$$
$$\frac{\partial L}{\partial v_{k}} = p_{k} \quad \text{or:} \quad \frac{\partial L}{\partial \mathbf{v}} = \mathbf{p}$$

Tuesday, September 18, 2012

![](_page_60_Figure_0.jpeg)