

Lecture 4

Thur. 9.5.2013

Kinetic Derivation of 1D Potentials and Force Fields

(Ch. 6, and Ch. 7 of Unit 1)

Review of $(V_1, V_2) \rightarrow (y_1, y_2)$ relations High mass ratio $M_1/m_2 = 49$

Force “field” or “pressure” due to many small bounces

Force defined as momentum transfer rate

The 1D-Isothermal force field $F(y) = \text{const.}/y$ and the 1D-Adiabatic force field $F(y) = \text{const.}/y^3$

Potential field due to many small bounces

Example of 1D-Adiabatic potential $U(y) = \text{const.}/y^2$

Physicist’s Definition $F = -\Delta U / \Delta y$ vs. Mathematician’s Definition $F = +\Delta U / \Delta y$

Example of 1D-Isothermal potential $U(y) = \text{const.} \ln(y)$

“Monster Mash” classical segue to Heisenberg action relations

Example of very very large M_1 ball-wall(s) crushing a poor little m_2

How m_2 keeps its action

An interesting wave analogy: The “Tiny-Big-Bang” [Harter, J. Mol. Spec. 210, 166-182 (2001)], [Harter, Li IMSS (2012)]

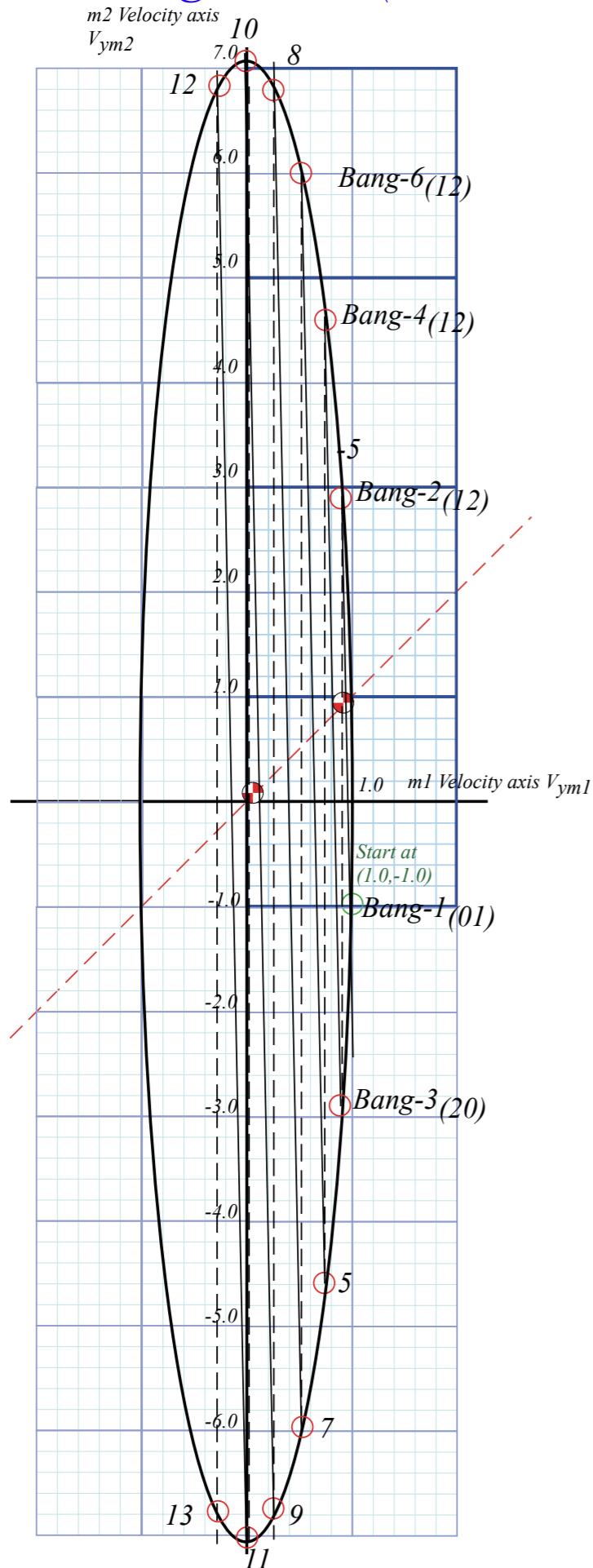
A lesson in geometry of fractions and fractals: Ford Circles and Farey Sums

[Lester R. Ford, Am. Math. Monthly 45, 586(1938)] [John Farey, Phil. Mag.(1816)]

Review of $(V_1, V_2) \rightarrow (y_1, y_2)$ relations

→ *High mass ratio $M_1/m_2 = 49$*

Geometric “Integration” (Converting Velocity data to Space-time trajectory)



Example with masses: $m_1=49$ and $m_2=1$

Kinetic Energy Ellipse

$$KE = \frac{1}{2} m_1 v_1^2 + \frac{1}{2} m_2 v_2^2 = \frac{49}{2} + \frac{1}{2} = 25$$

$$1 = \frac{v_1^2}{2KE/m_1} + \frac{v_2^2}{2KE/m_2} = \frac{x_1^2}{a_1^2} + \frac{x_2^2}{a_2^2}$$

Ellipse radius 1

$$\begin{aligned} a_1 &= \sqrt{2KE/m_1} \\ &= \sqrt{2KE/49} \\ &= \sqrt{50/49} \\ &= 1.01 \end{aligned}$$

Ellipse radius 2

$$\begin{aligned} a_2 &= \sqrt{2KE/m_2} \\ &= \sqrt{2KE/1} \\ &= \sqrt{50/1} \\ &= 7.07 \end{aligned}$$

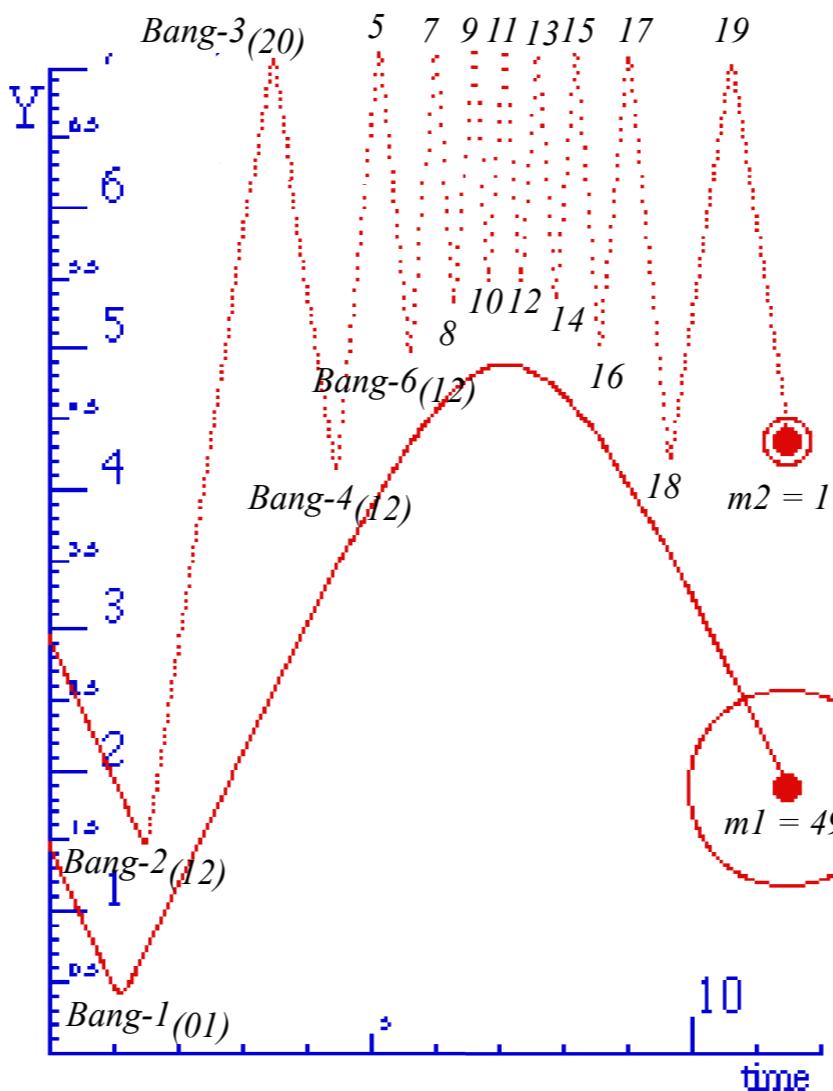


Fig. 5.1
in Unit 1

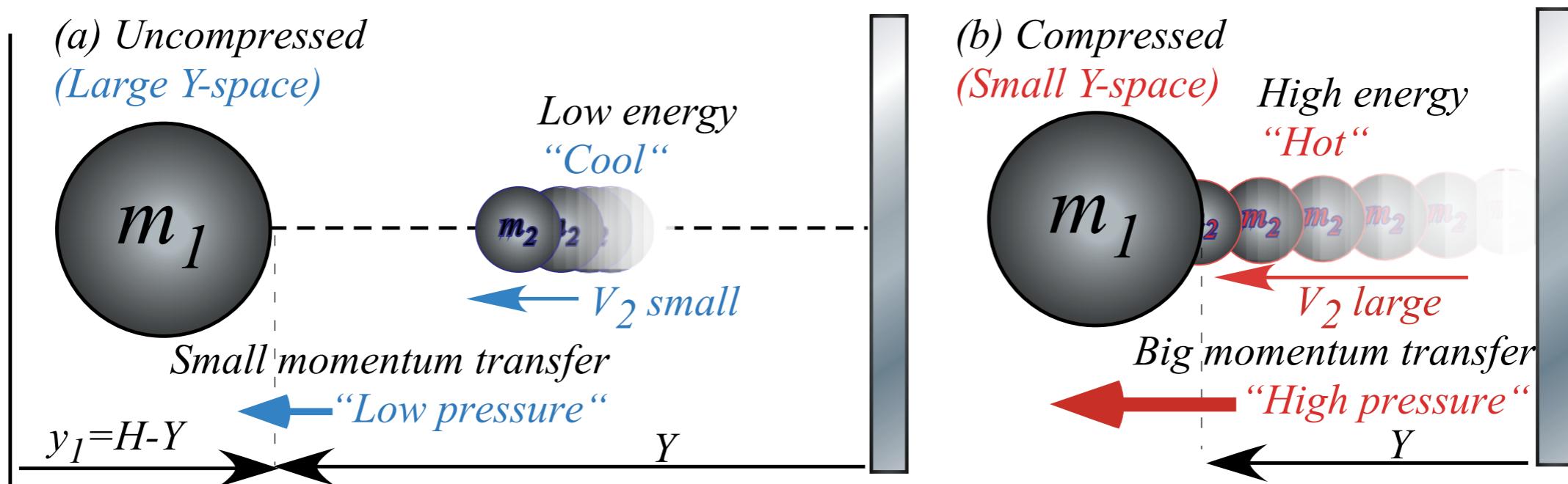
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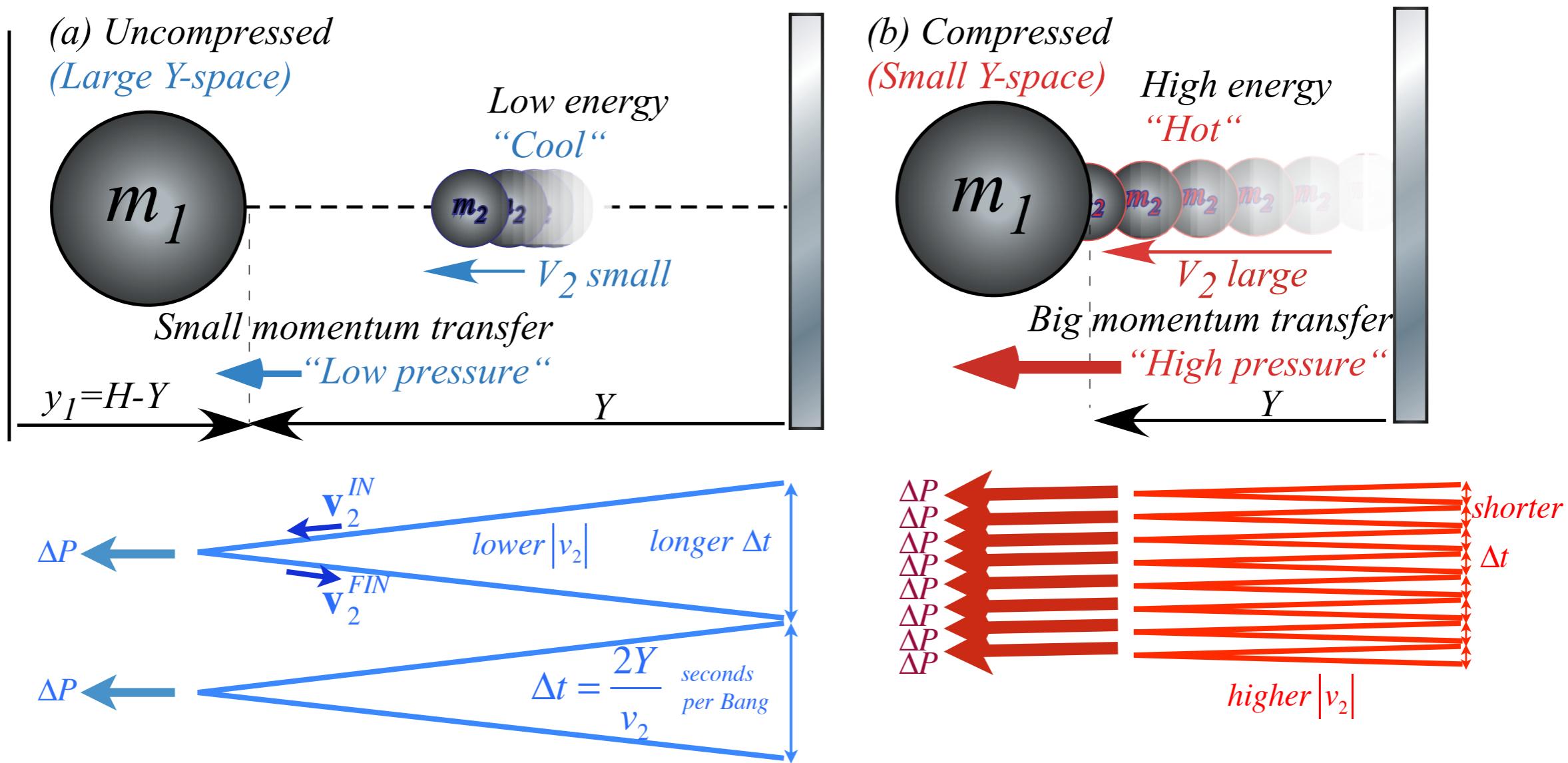
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Unit 1
Fig. 6.1



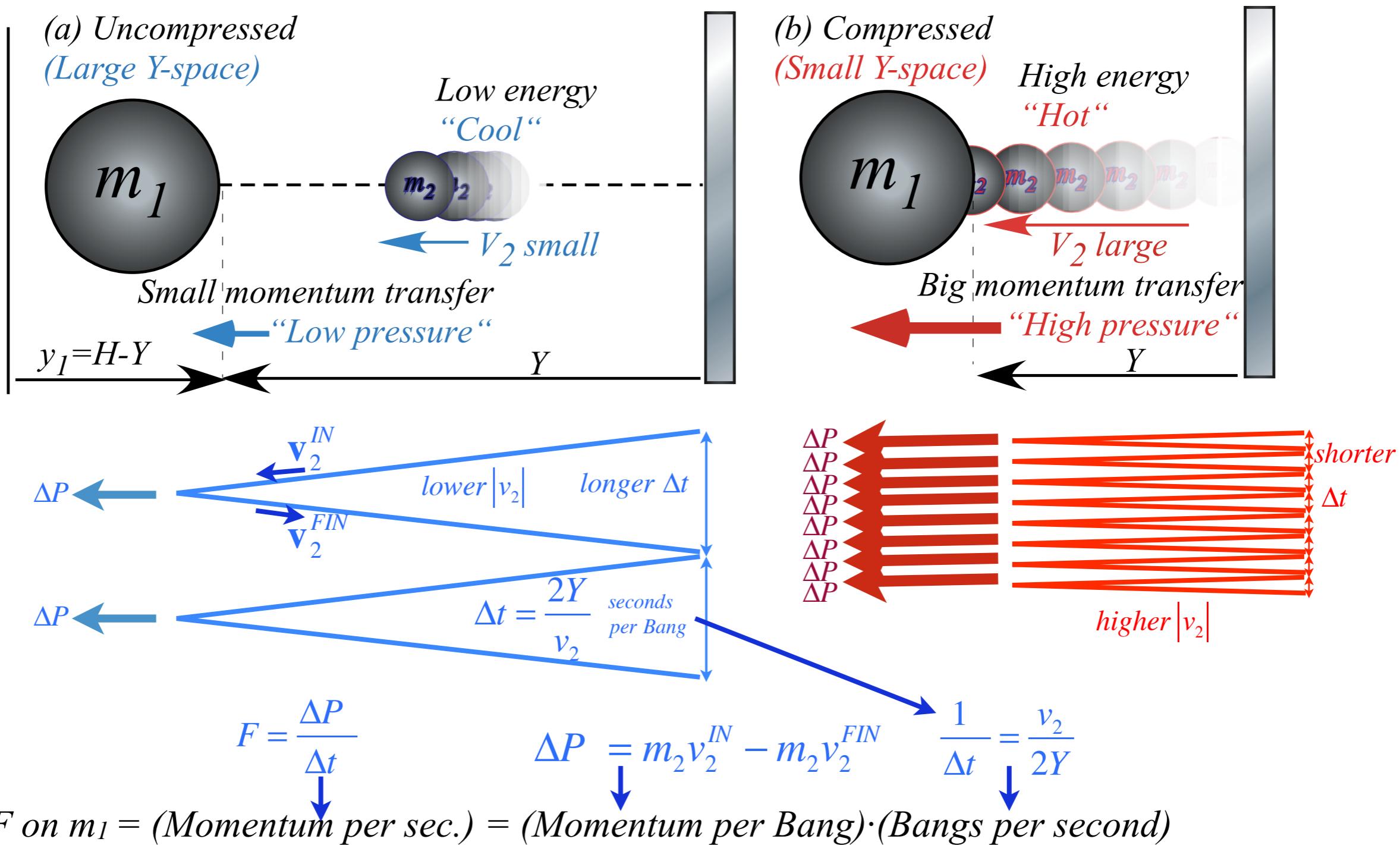
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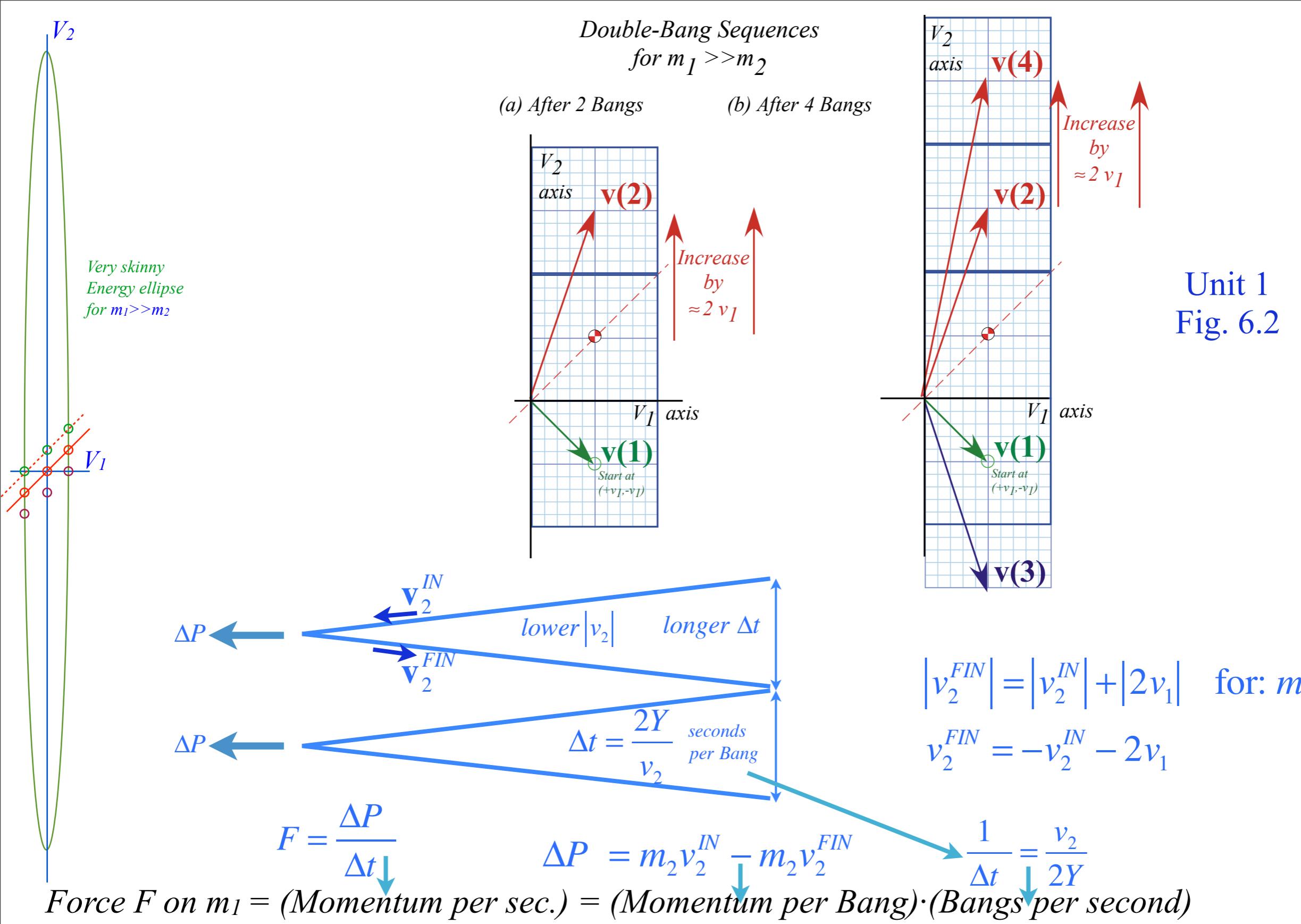
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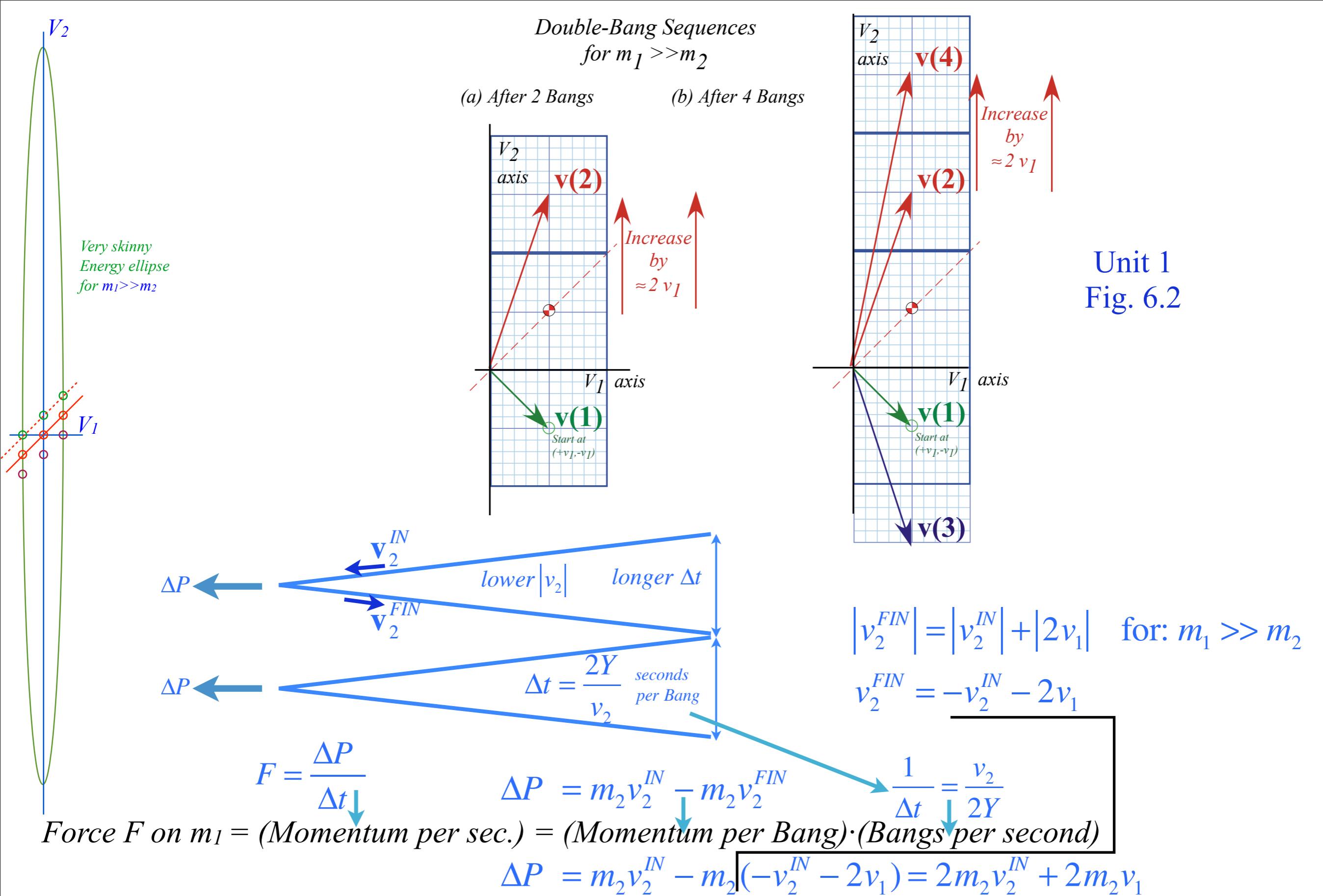


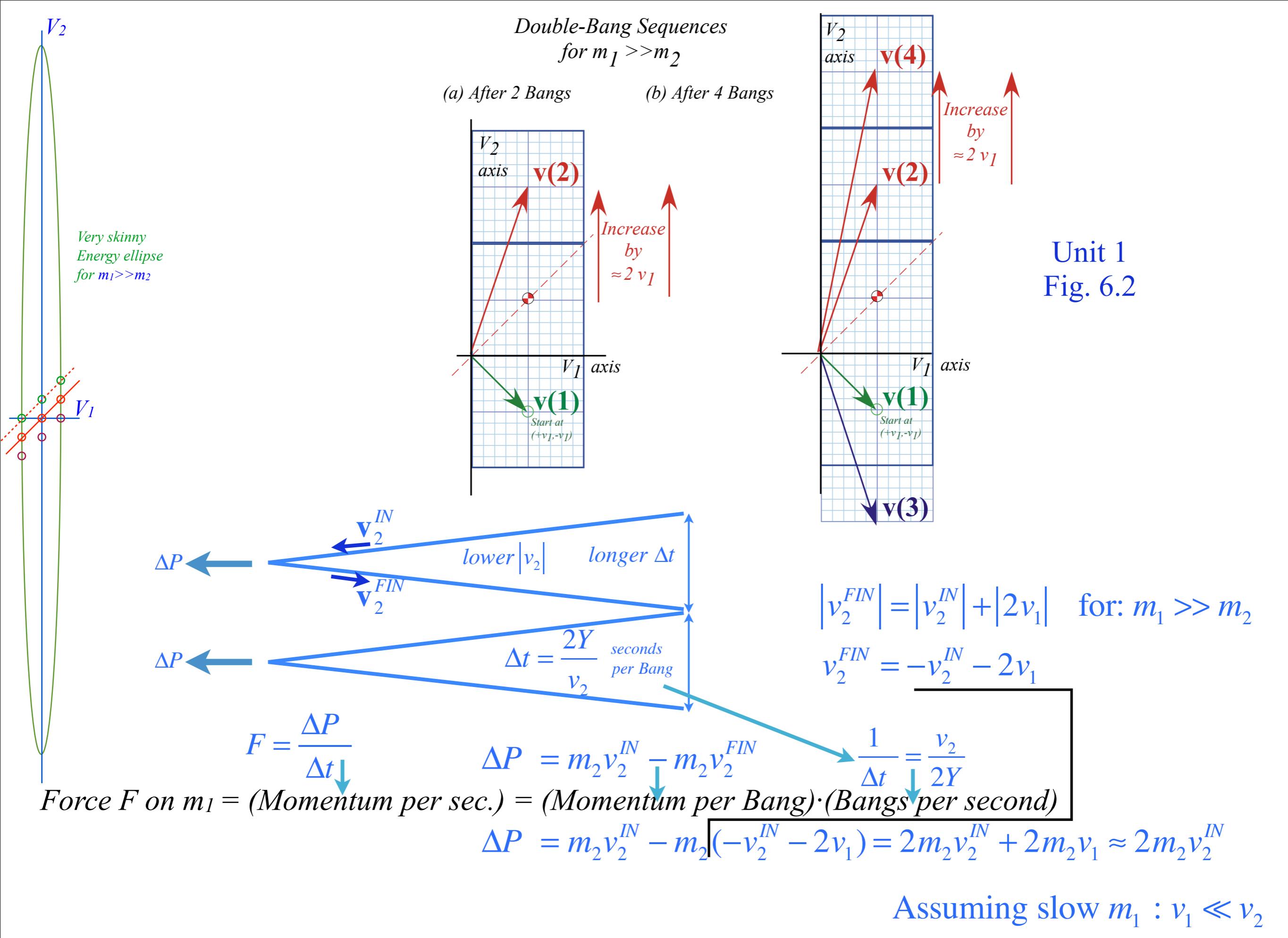
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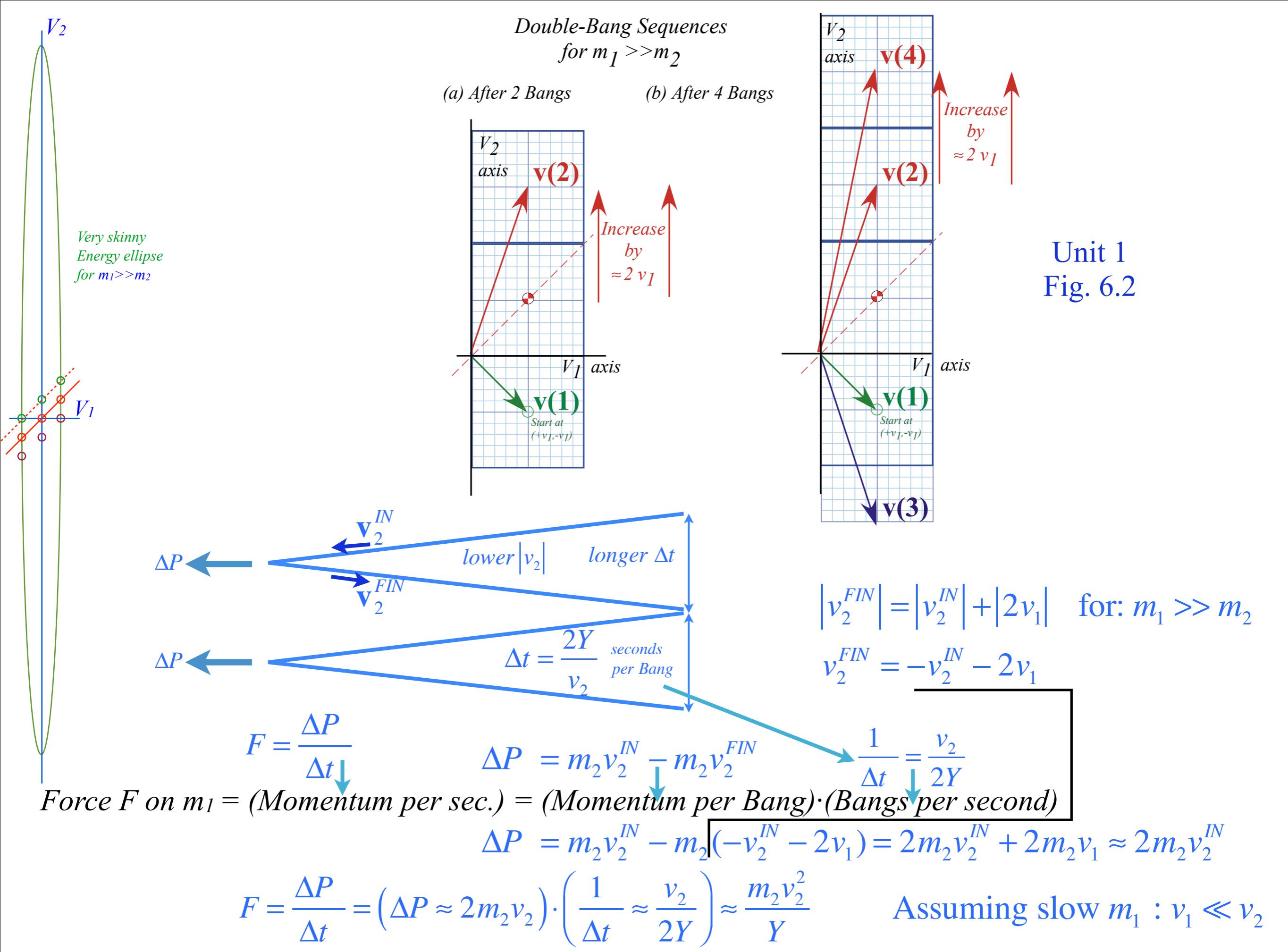








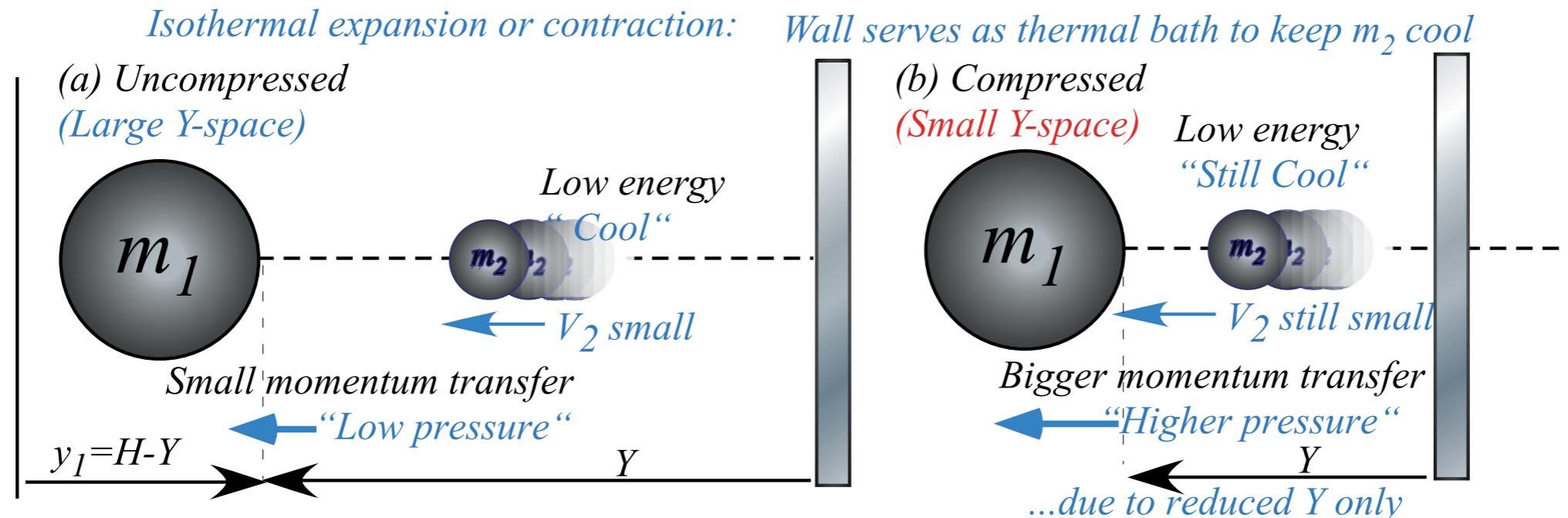
Unit 1
Fig. 6.2



$$F = \frac{\Delta P}{\Delta t} = (\Delta P \approx 2m_2 v_2) \cdot \left(\frac{1}{\Delta t} \approx \frac{v_2}{2Y} \right) \approx \frac{m_2 v_2^2}{Y}$$

1D-Isothermal Force Law (assume v_2 is constant for all Y):

$$F = \frac{m_2 v_2^2}{Y} = \frac{\text{const.}}{Y}$$



Force “field” or “pressure” due to many small bounces

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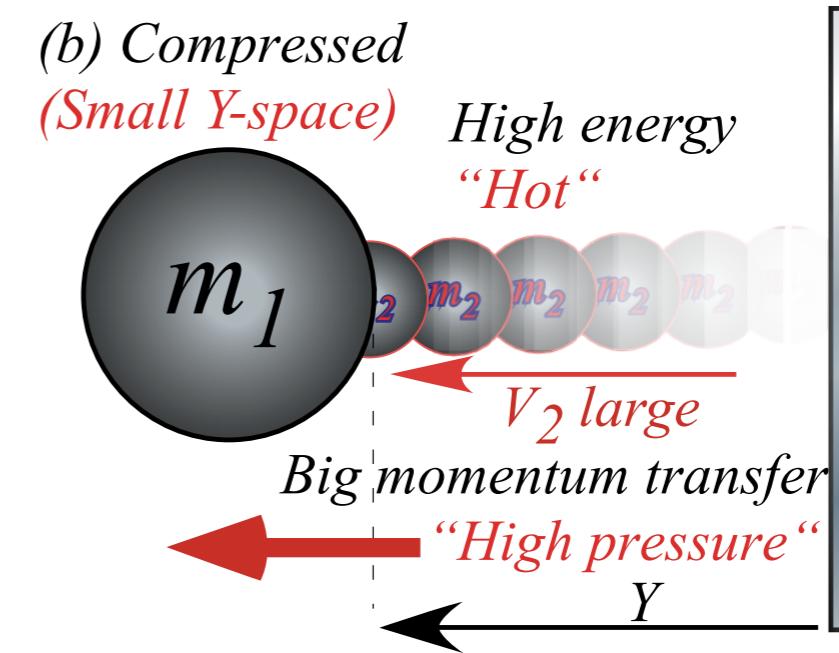
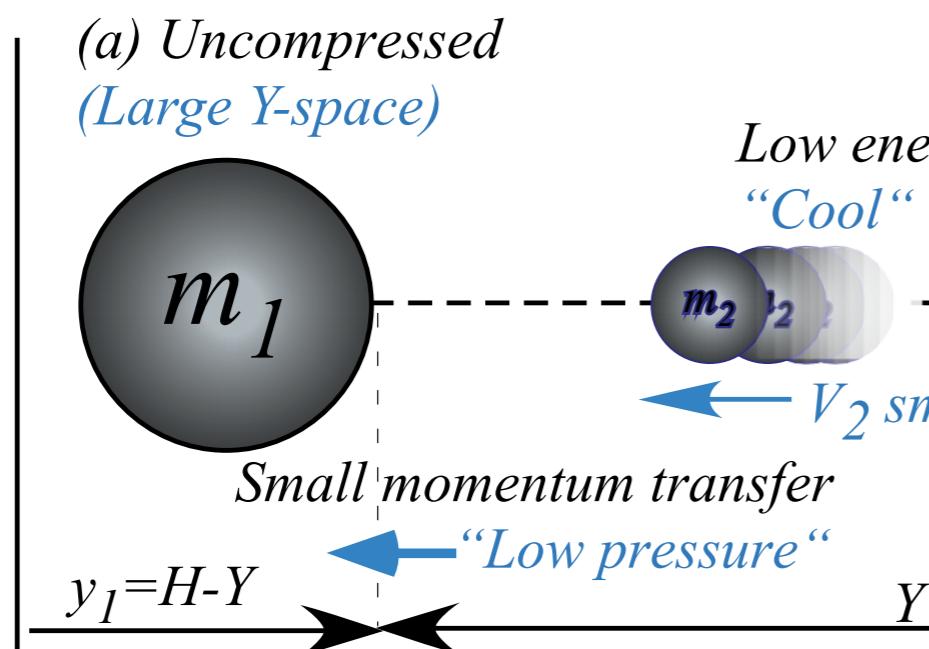
$$F = \frac{m_2 v_2^2}{Y} = \frac{\text{const.}}{Y}$$

However, if ceiling is elastic, v_2 isn't constant if m_1 changes bounce range Y : $\frac{dy_1}{dt} \equiv v_1 = -\frac{dY}{dt}$

When m_1 collides with m_2 it adds twice its velocity ($2v_1$) to v_2 . This occurs at “bang-rate” $B=v_2/2Y$.

$$\frac{dv_2}{dt} = 2v_1 B = 2v_1 \frac{v_2}{2Y} = -2 \frac{dY}{dt} \frac{v_2}{2Y}$$

Wall not given time to give or take KE



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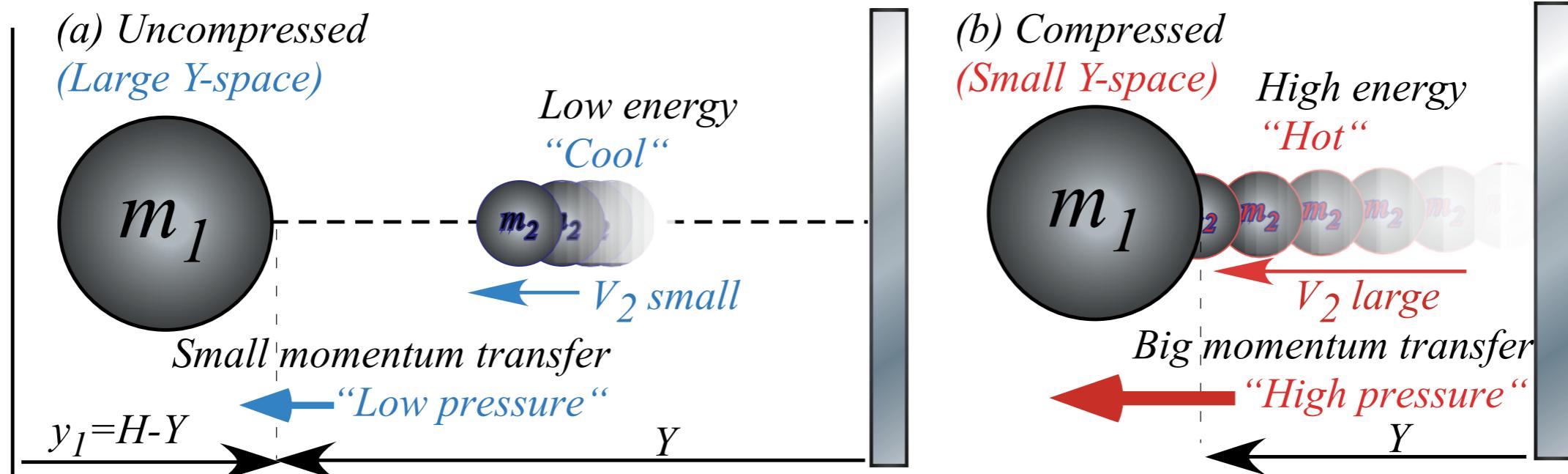
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Differential equation results and has logarithmic integral. $\int \frac{dx}{x} = \ln x + C = \log_e x + \log_e e^C = \log_e(e^C x)$

$$\frac{dv_2}{v_2} = -\frac{dY}{Y} \quad \text{integrates to: } \ln v_2 = -\ln Y + C \quad \text{or: } \ln v_2 = \ln \frac{\text{const.}}{Y} \quad \text{or: } v_2 = \frac{\text{const.}}{Y}$$

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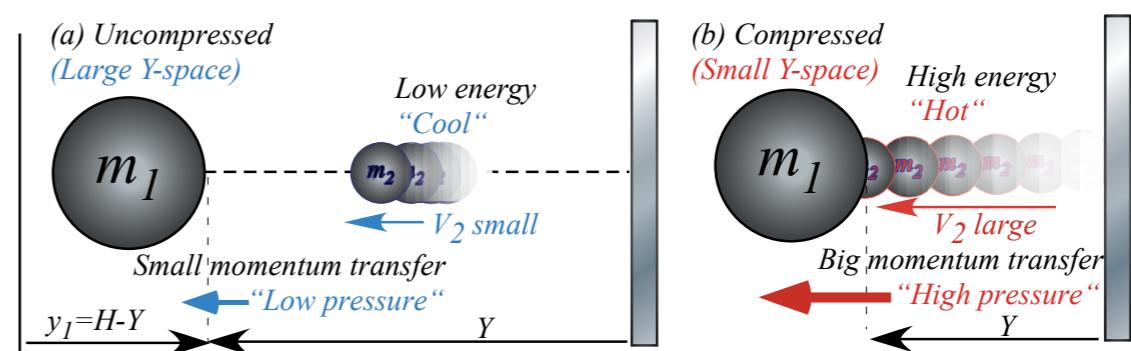
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Force law with this variable v_2 is called *adiabatic* or not-*adiabatic* or not-gradual.

1D-Adiabatic Force Law (assume v_2 varies: $v_2 = \frac{\text{const.}}{Y} = \frac{v_2^{IN} Y(t=0)}{Y}$):

$$F = \frac{m_2 (v_2^{IN} Y(t=0))^2}{Y^3} = \frac{\text{const.}}{Y^3}$$



Potential field due to many small bounces

→ Example of 1D-Adiabatic potential $U(y)=\text{const.}/y^2$

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Example of 1D-Isothermal potential $U(y)=\text{const.} \ln(y)$

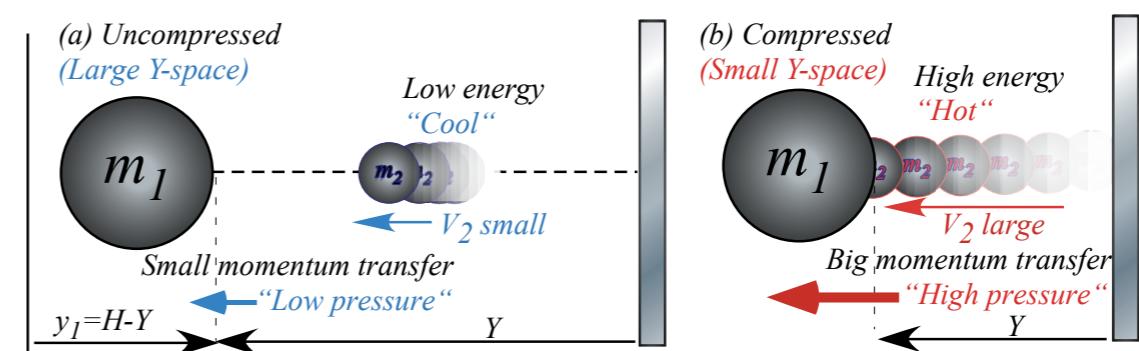
Big mass- m_1 ball feeling “potential-field” or “gradient” due to small ($m_2 \ll m_1$) rapidly ($v_2 \gg v_1$) bouncing ball

In adiabatic case where $v_2 = \frac{\text{const.}}{Y}$ the total energy E is strictly conserved.

$$\text{const.} = E = \frac{1}{2}m_1v_1^2 + \frac{1}{2}m_2v_2^2 = \frac{1}{2}m_1v_1^2 + \frac{1}{2}m_2\left(\frac{\text{const.}}{Y}\right)^2$$

Define for big mass m_1 : *Kinetic energy* $KE(v_1)$ vs *Potential energy* $PE(Y) = U(Y)$

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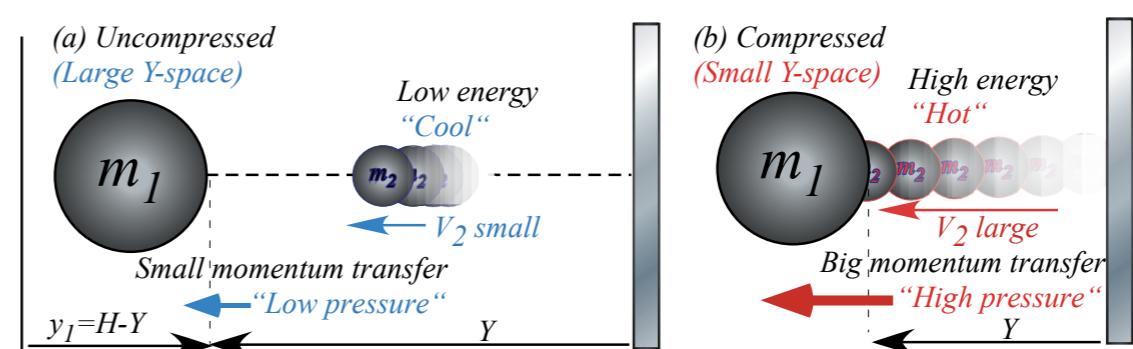
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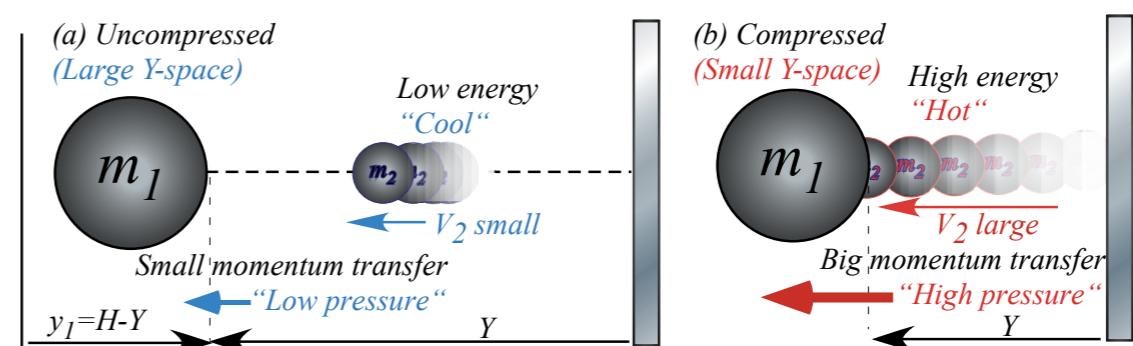
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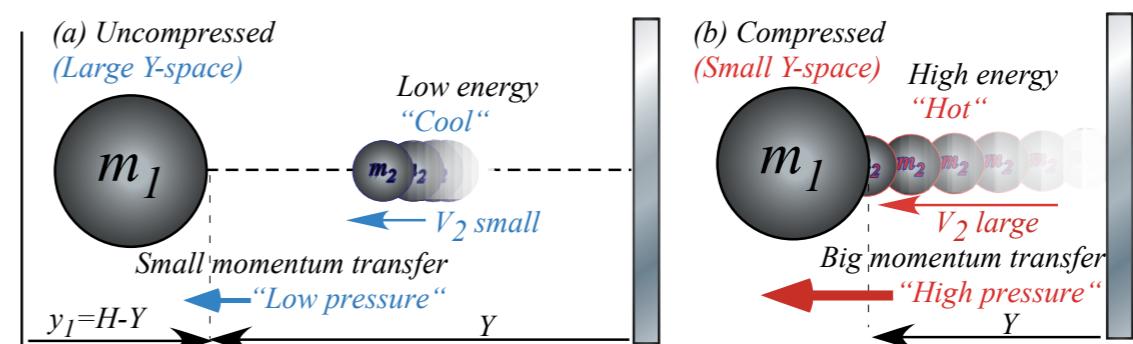
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$$\text{or else : } \mathbf{F} \cdot \frac{dY}{dt} = \frac{dp}{dt} \cdot V = \frac{d(mV)}{dt} \cdot V = \frac{d(mV^2)/2}{dt} = \frac{dU}{dt}$$



Potential field due to many small bounces

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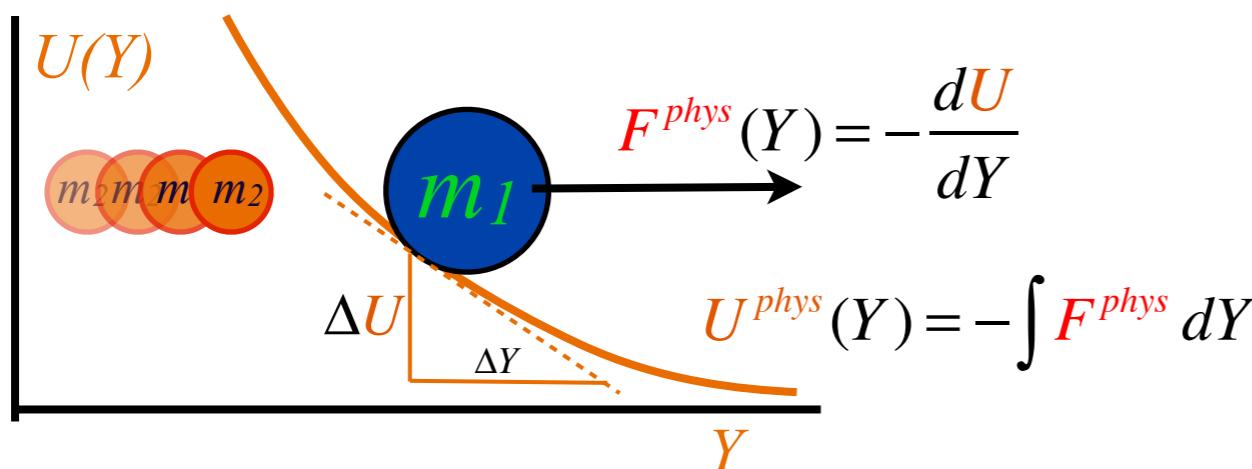
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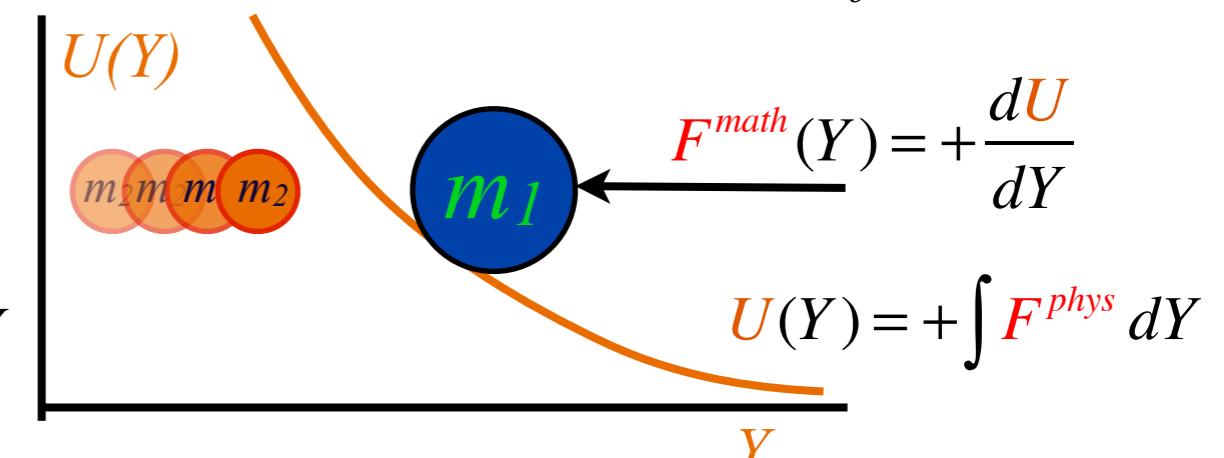
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The “Physicist” View of Force



The “Mathematician” View of Force



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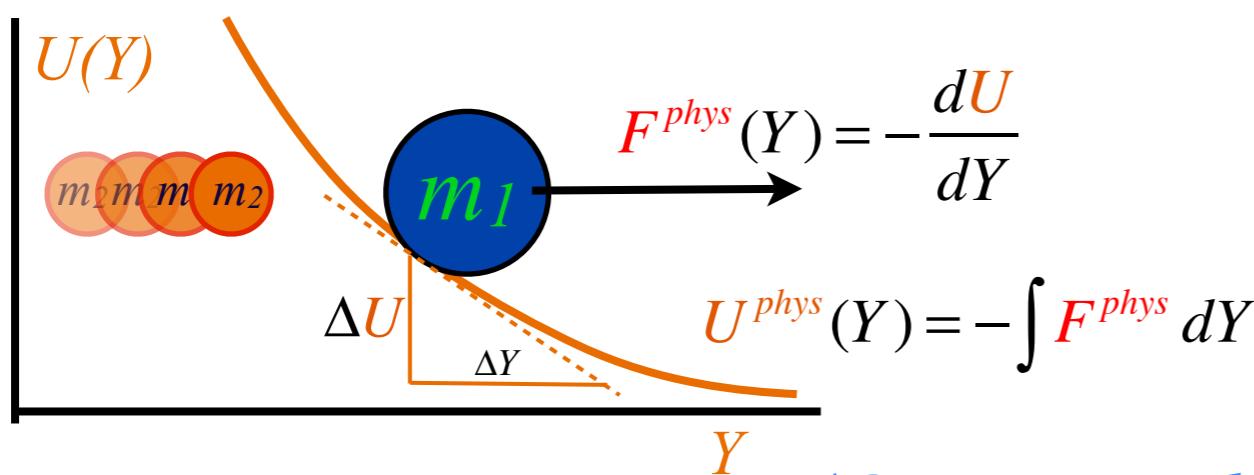
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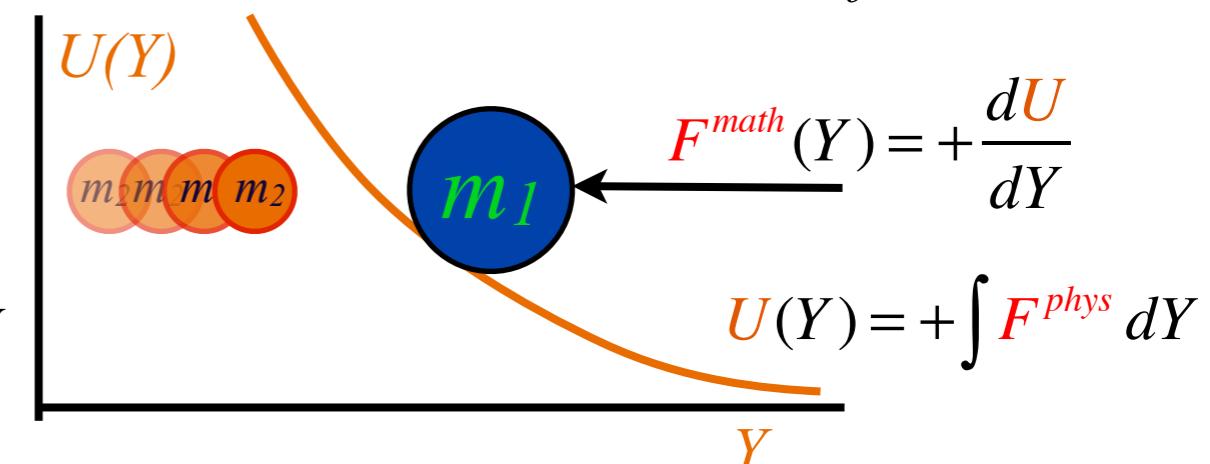
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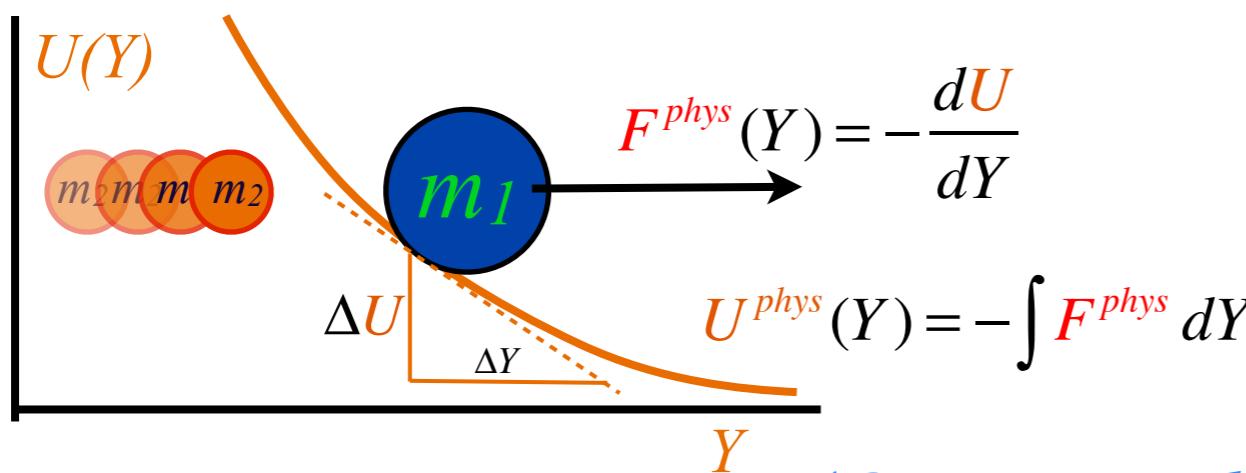
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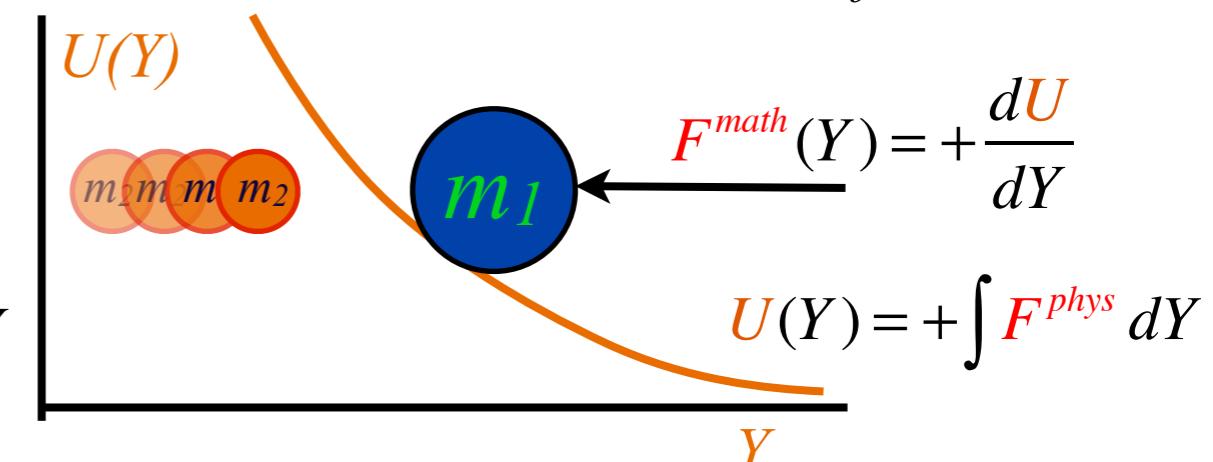
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The “Physicist” View of Force



The “Mathematician” View of Force



(OK, But, does it work?)

$$F^{\text{phys}} = m_2 \frac{(\text{const.})^2}{Y^3}$$

consistent
with :

$$F^{\text{phys}} = -\frac{\Delta U}{\Delta Y} = -\frac{d}{dY} \frac{1}{2}m_2\left(\frac{\text{const.}}{Y}\right)^2 = m_2 \frac{(\text{const.})^2}{Y^3}$$

(Hurrah!)

Potential field due to many small bounces

Example of 1D-Adiabatic potential $U(y)=\text{const.}/y^2$

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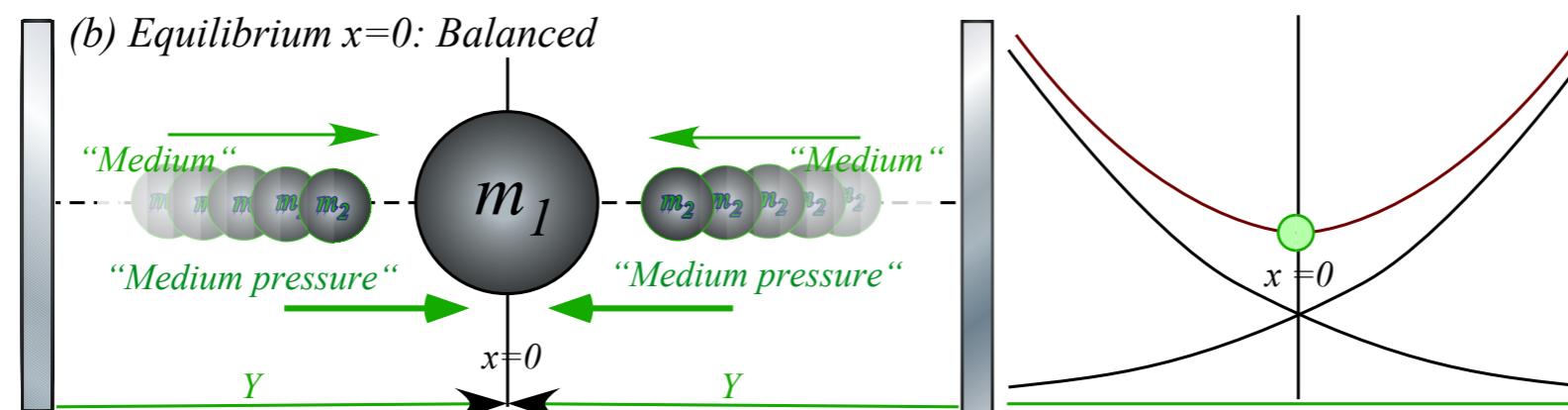
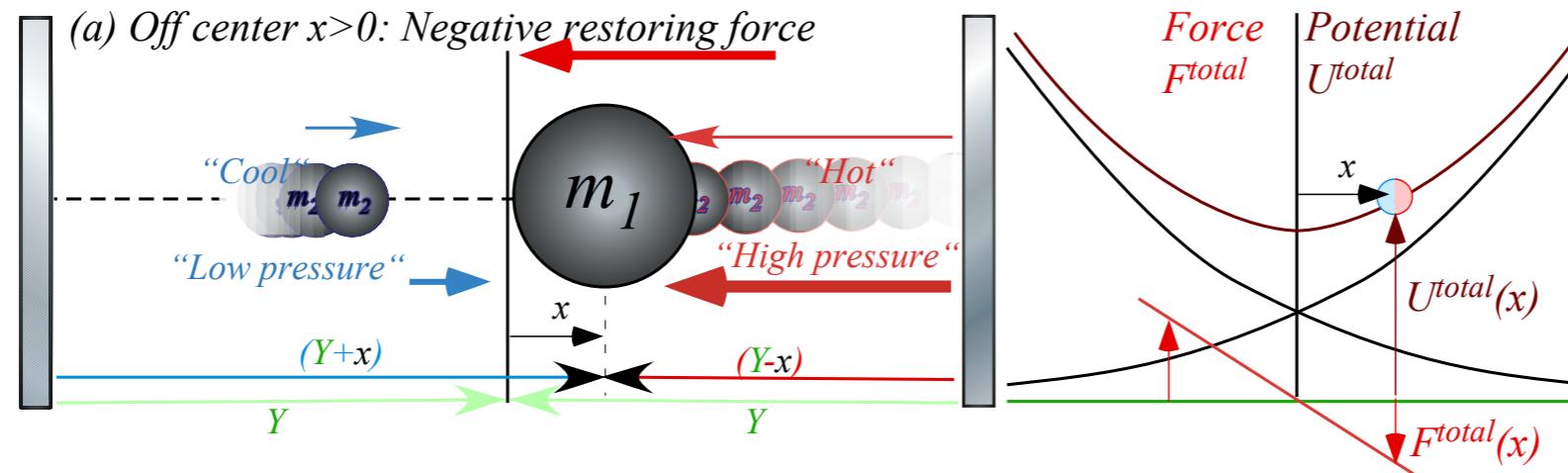
1D-Isothermal Force Law (assume v_2 is constant for all Y):

$$F = \frac{m_2 v_2^2}{Y} = \frac{\text{const.}}{Y}$$

$$F^{phys} = \frac{m_2 v_2^2}{Y} = -\frac{\Delta U}{\Delta Y}$$

implies :

$$U = -m_2 v_2^2 \ln(Y)$$



Two opposing 1D-Isothermal Force fields

$$F^{total} = \frac{f}{1+x} - \frac{f}{1-x} = f[1-x+x^2-x^3\dots] - f[1+x+x^2+x^3\dots] = -2f \cdot x - 2f \cdot x^3 - \dots$$

$$F^{HO} = -k \cdot x = -\frac{\partial U^{HO}}{\partial x}$$

$$U^{HO} = \frac{1}{2} k \cdot x^2 = - \int F^{HO} dx$$

HO ↘ frequency: $\omega = \sqrt{\frac{k}{m_1}} = 2\pi\nu$

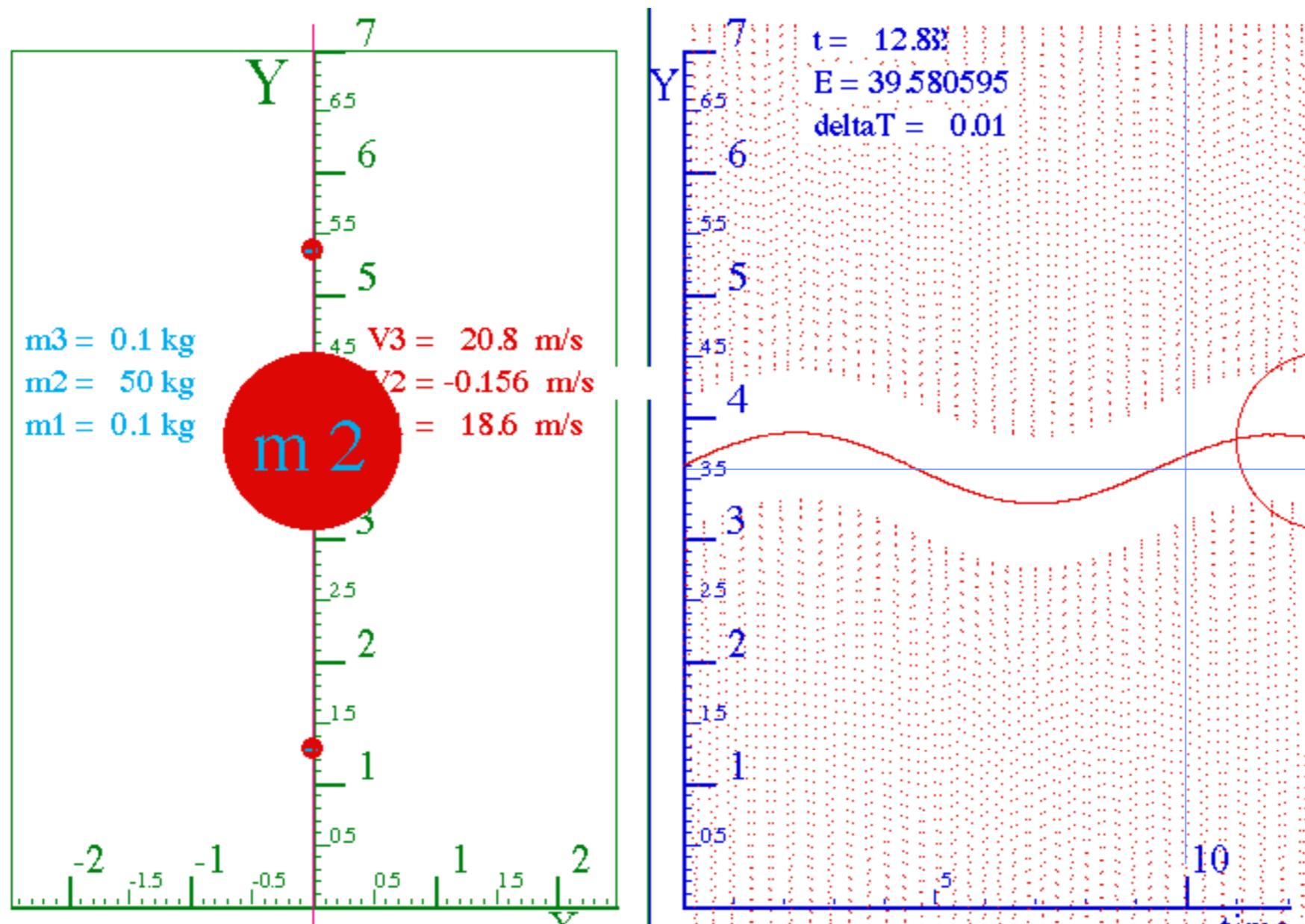
Unit 1
Fig. 6.2

Anharmonic oscillator terms...

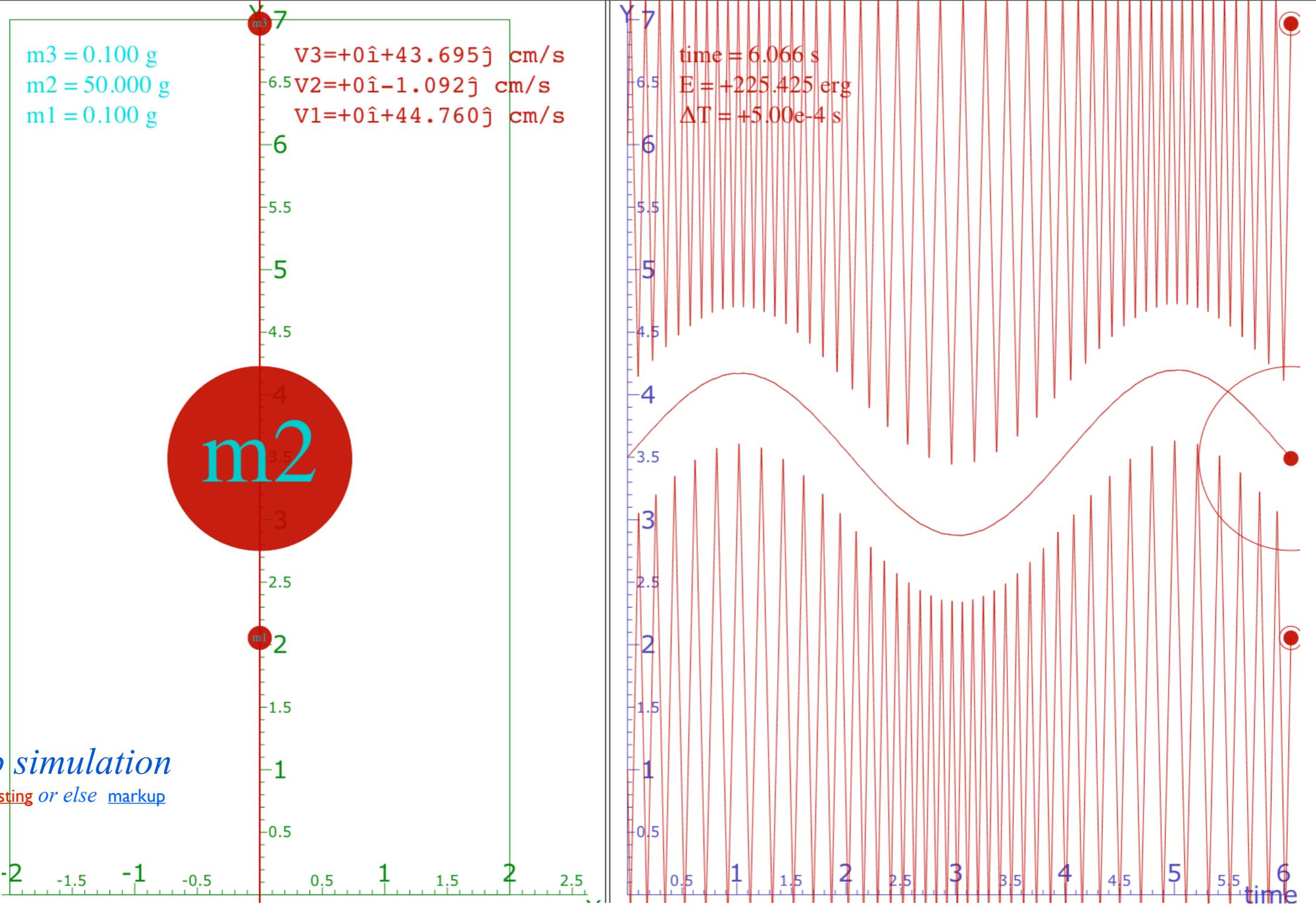
Harmonic oscillator term

Unit 1
Fig. 6.3

Simulation of
the adiabatic case



See Homework problem 1.6.1: Compute frequency and/or period for both isoT and adiabatic cases



Initial x1 =	<input type="text" value="0.75"/>	y Min =	<input type="text" value="0"/>	Adiabatic force scenarios											
Max x PE plot =	<input type="text" value="0.5"/>	T Max =	<input type="text" value="6"/>	Quasi-harmonic oscillation (m1:m2 = 100:1)											
F-Vector scale =	<input type="text" value="0.003"/>	V2y Max =	<input type="text" value="3"/>	Quasi-harmonic oscillation (m1:m2 = 50:1)											
Error step =	<input type="text" value="0.000"/>	V2y Min =	<input type="text" value="-2"/>	Quasi-harmonic oscillation (m1:m2 = 25:1)											
m1=	<input type="text" value="1"/>	x10 [^]	<input type="text" value="-1"/>	x10 [^]	X1 ₀ =	<input type="text" value="1"/>	x10 [^]	<input type="text" value="0"/>	x10 [^]	V1 ₀ =	<input type="text" value="-4.5"/>	x10 [^]	<input type="text" value="1"/>	x10 [^]	{cm/s}
m2=	<input type="text" value="5"/>	x10 [^]	<input type="text" value="1"/>	x10 [^]	X2 ₀ =	<input type="text" value="3.5"/>	x10 [^]	<input type="text" value="0"/>	x10 [^]	V2 ₀ =	<input type="text" value="1"/>	x10 [^]	<input type="text" value="0"/>	x10 [^]	{cm/s}
m3=	<input type="text" value="1"/>	x10 [^]	<input type="text" value="-1"/>	x10 [^]	X3 ₀ =	<input type="text" value="6"/>	x10 [^]	<input type="text" value="0"/>	x10 [^]	V3 ₀ =	<input type="text" value="4.5"/>	x10 [^]	<input type="text" value="1"/>	x10 [^]	{cm/s}

“Monster Mash” classical segue to Heisenberg action relations

→ *Example of very very large M_1 ball-walls crushing a poor little m_2*

How m_2 keeps its action

An interesting wave analogy: The “Tiny-Big-Bang” [Harter, J. Mol. Spec. 210, 166-182 (2001)], [Harter, Li IMSS (2012)]

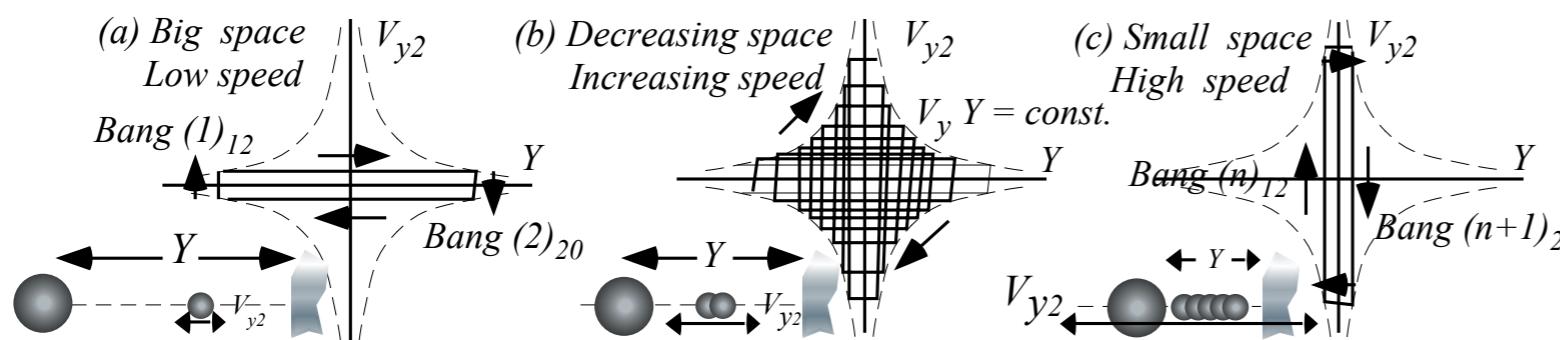
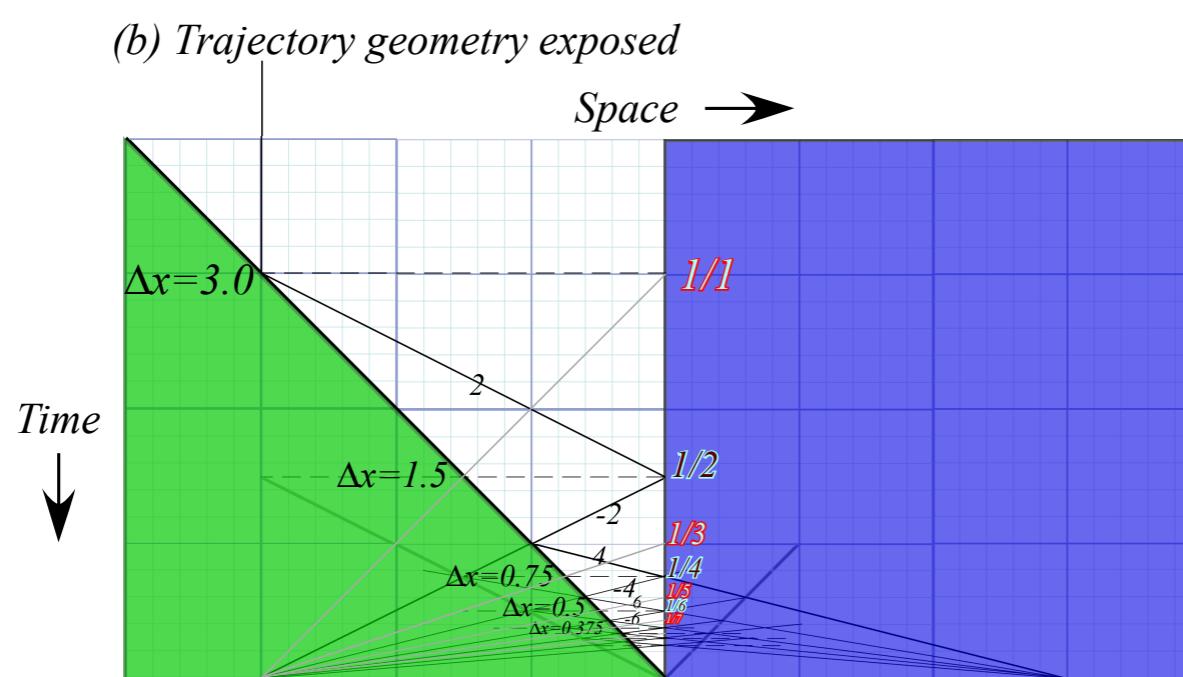
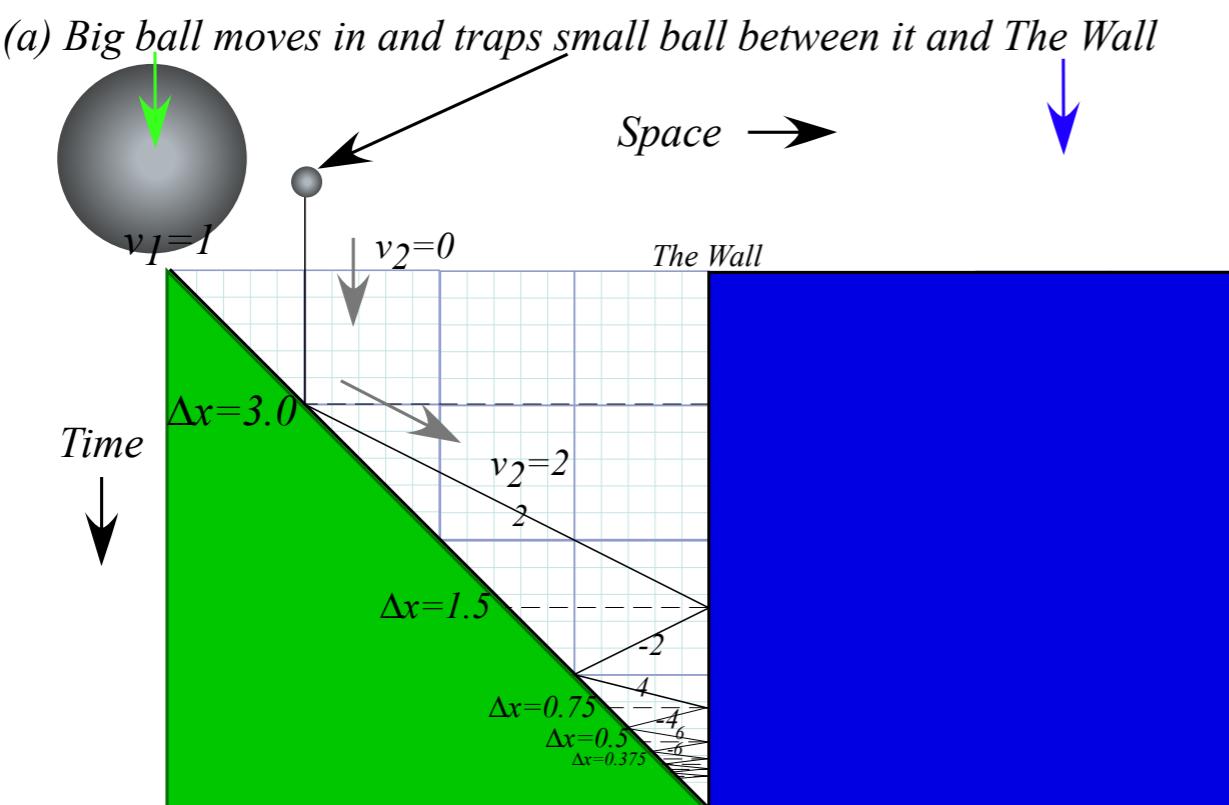
A lesson in geometry of fractions and fractals: Ford Circles and Farey Sums

[Lester. R. Ford, Am. Math. Monthly 45, 586(1938)] [John Farey, Phil. Mag.(1816)]

The Classical “Monster Mash”

Classical introduction to

Heisenberg “Uncertainty” Relations

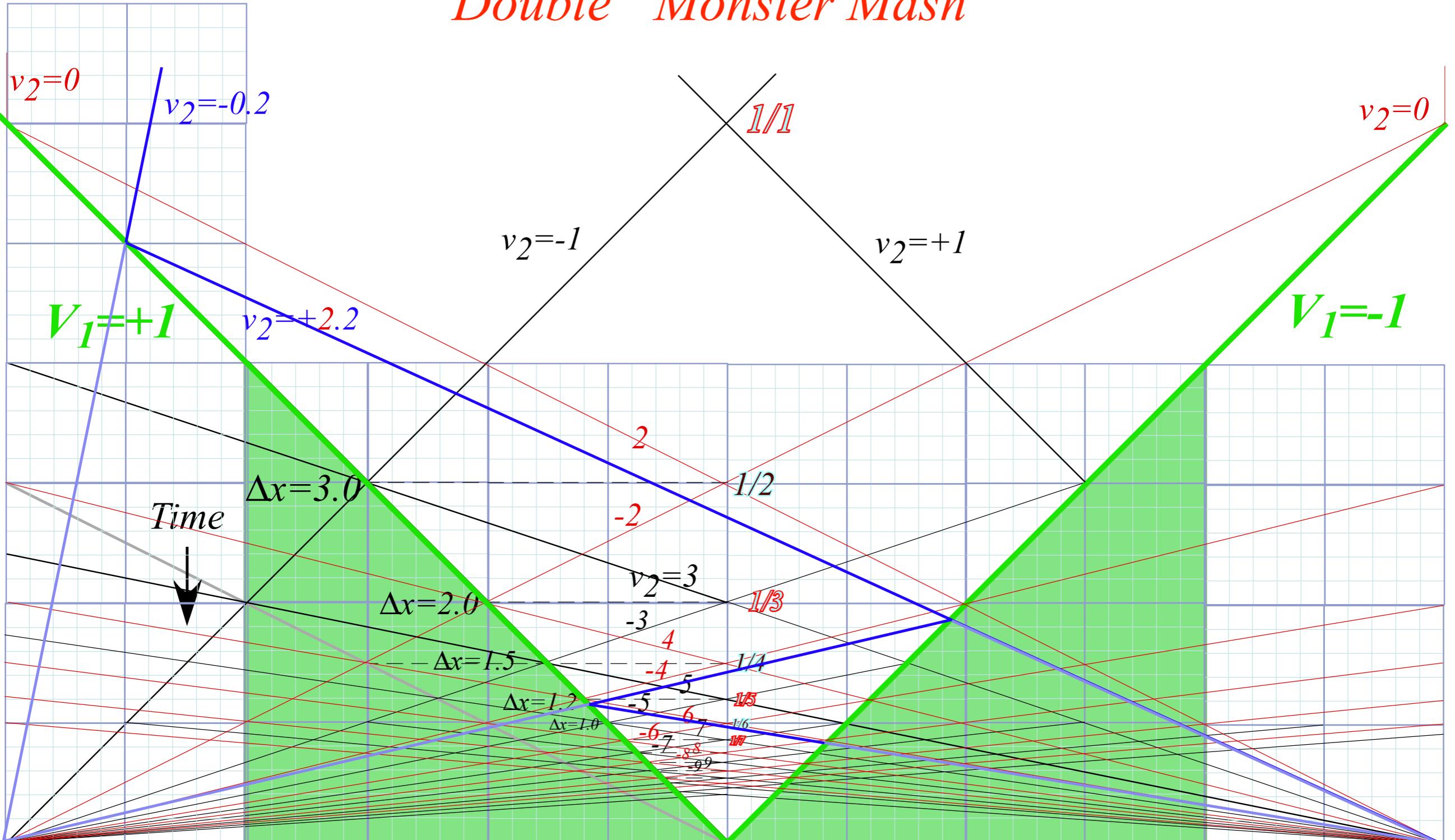


$$v_2 = \frac{\text{const.}}{Y} \quad \text{or: } Y \cdot v_2 = \text{const.}$$

is analogous to: $\Delta x \cdot \Delta p = N \cdot \hbar$

Unit 1
Fig. 6.4

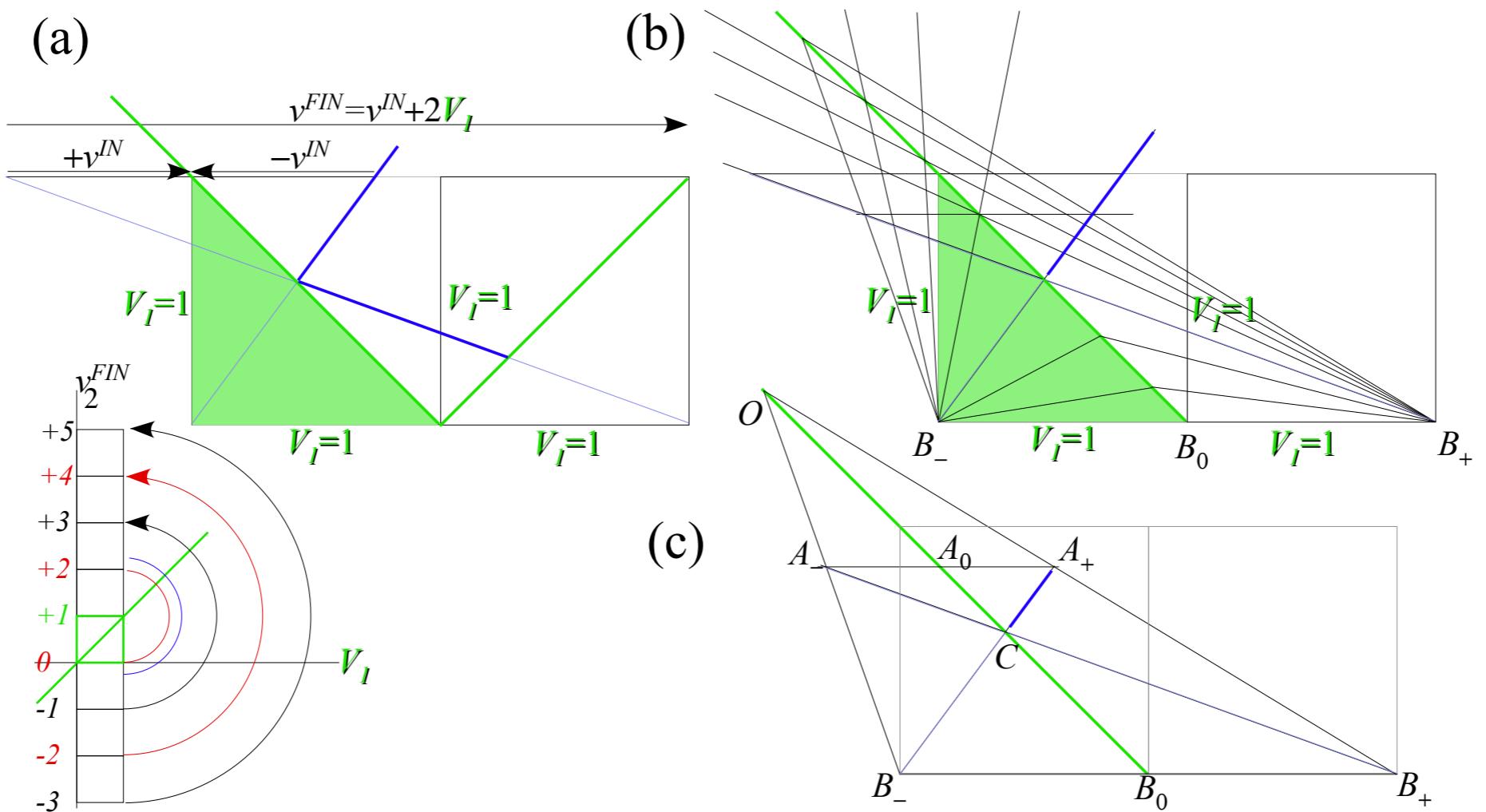
Double “Monster Mash”



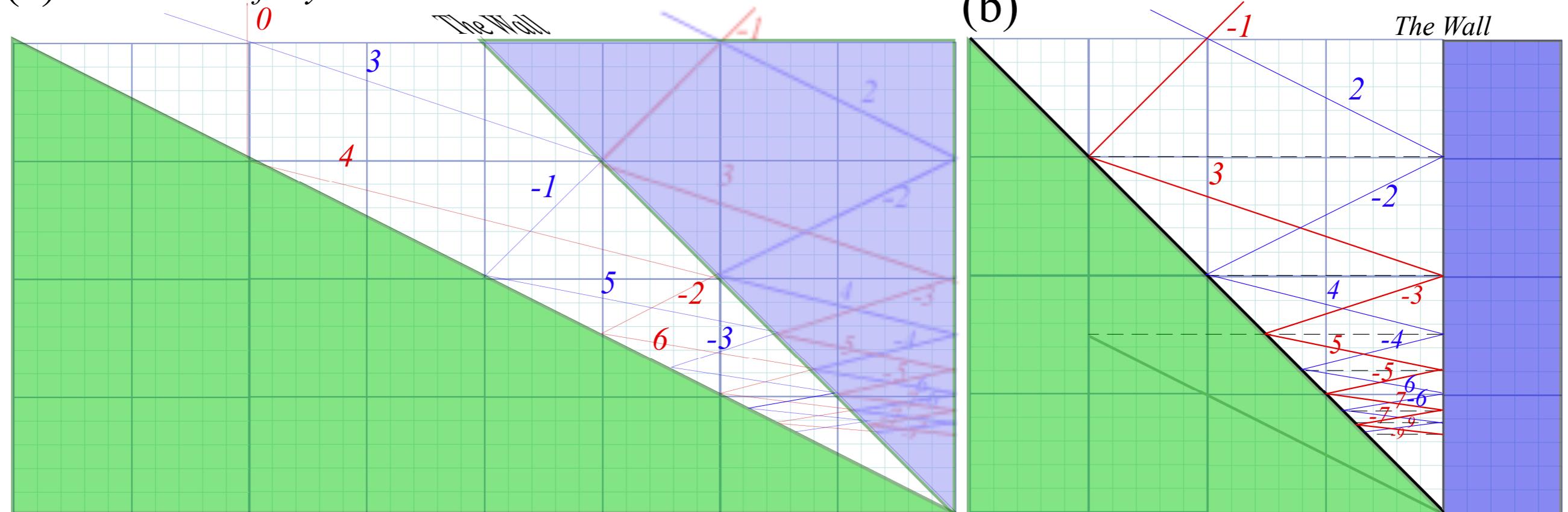
Unit 1
Fig. 6.5

See Homework problem 1.6.2: Construct related spacetime case

Unit 1
Fig. 6.6
and
Fig. 6.7



(a) Galilean shift by $V=1$



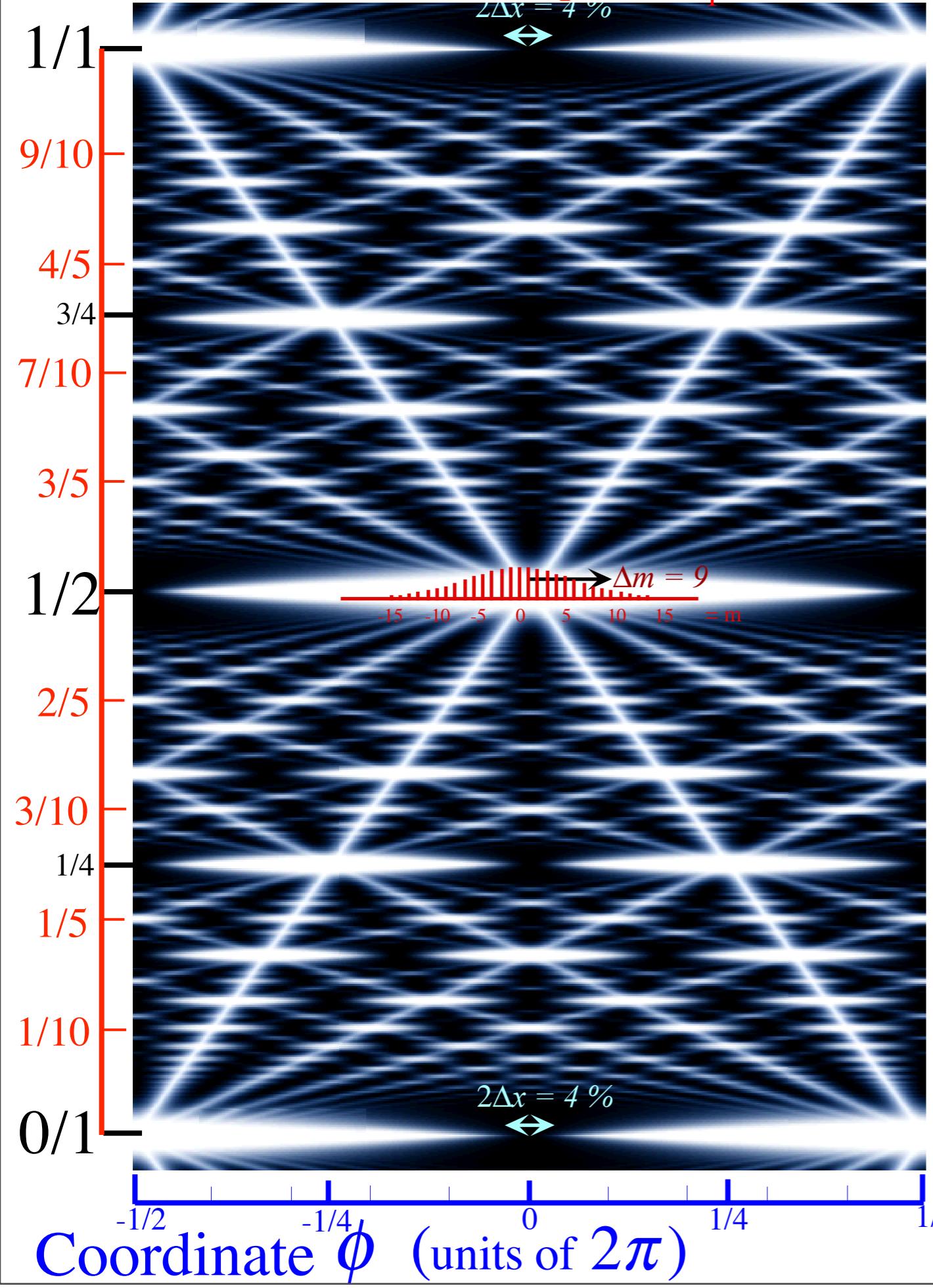
“Monster Mash” classical segue to Heisenberg action relations

Example of very very large M_1 ball-walls crushing a poor little m_2

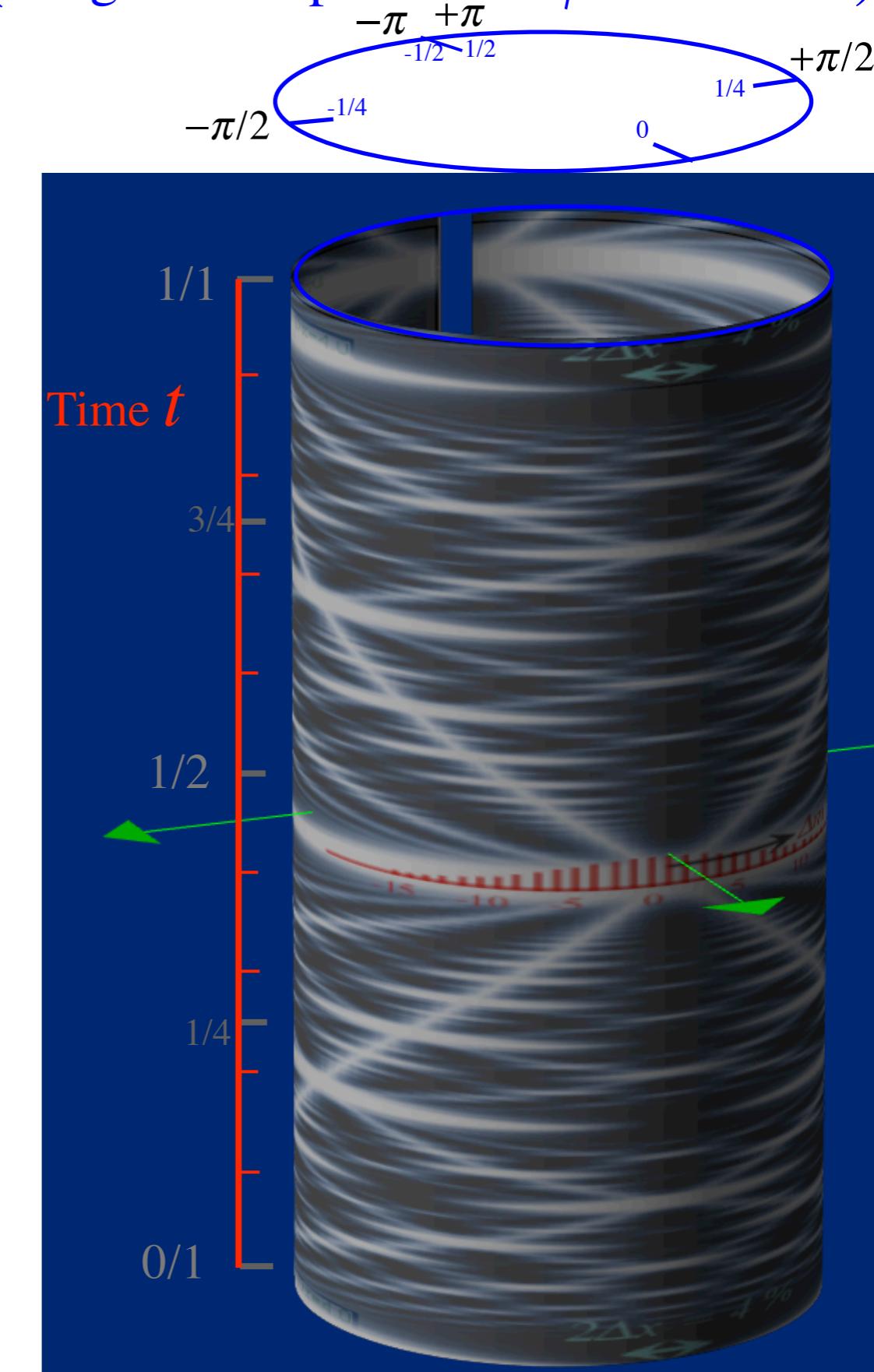
How m_2 keeps its action

- *An interesting wave analogy: The “Tiny-Big-Bang”* [Harter, J. Mol. Spec. 210, 166-182 (2001)], [Harter, Li IMSS (2012)]
A lesson in geometry of fractions and fractals: Ford Circles and Farey Sums
[Lester. R. Ford, Am. Math. Monthly 45, 586(1938)] [John Farey, Phil. Mag.(1816)]

Time t (units of fundamental period τ_1)



(Imagine "wrap-around" ϕ -coordinate)



Web simulation

<http://www.uark.edu/ua/modphys/markup/WaveltWeb.html>

Try [testing or else markup](#)

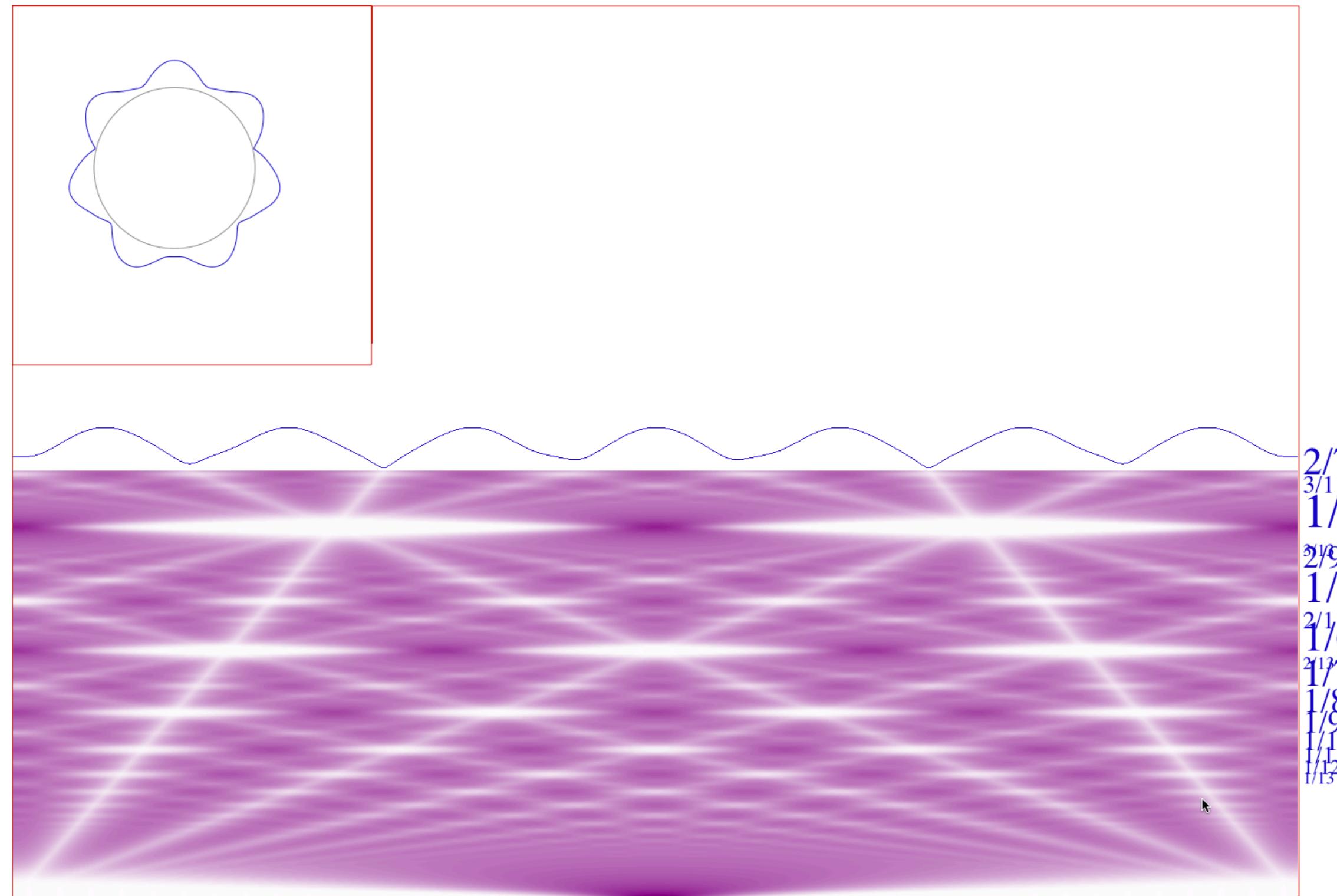
Click here....

[Launch](#) [Fourier Control](#) [Scenarios](#) [Pause](#) [Set T=0](#) [Zero Amps](#) T-Scale=

...then here....

Twelve (n=12) oscillator
Twelve (n=12) oscillator
Twelve (n=12) oscillator
C(n) Character Table
Quantum Carpet

time = 0.29T



Web simulation

Try [testing](#) or [else](#) markup

[Click here....](#)

<http://www.uark.edu/ua/modphys/markup/WaveltWeb.html>

[Local Control](#)

[Fourier Control](#)

[Scenarios](#)

[Pause](#)

[Set T=0](#)

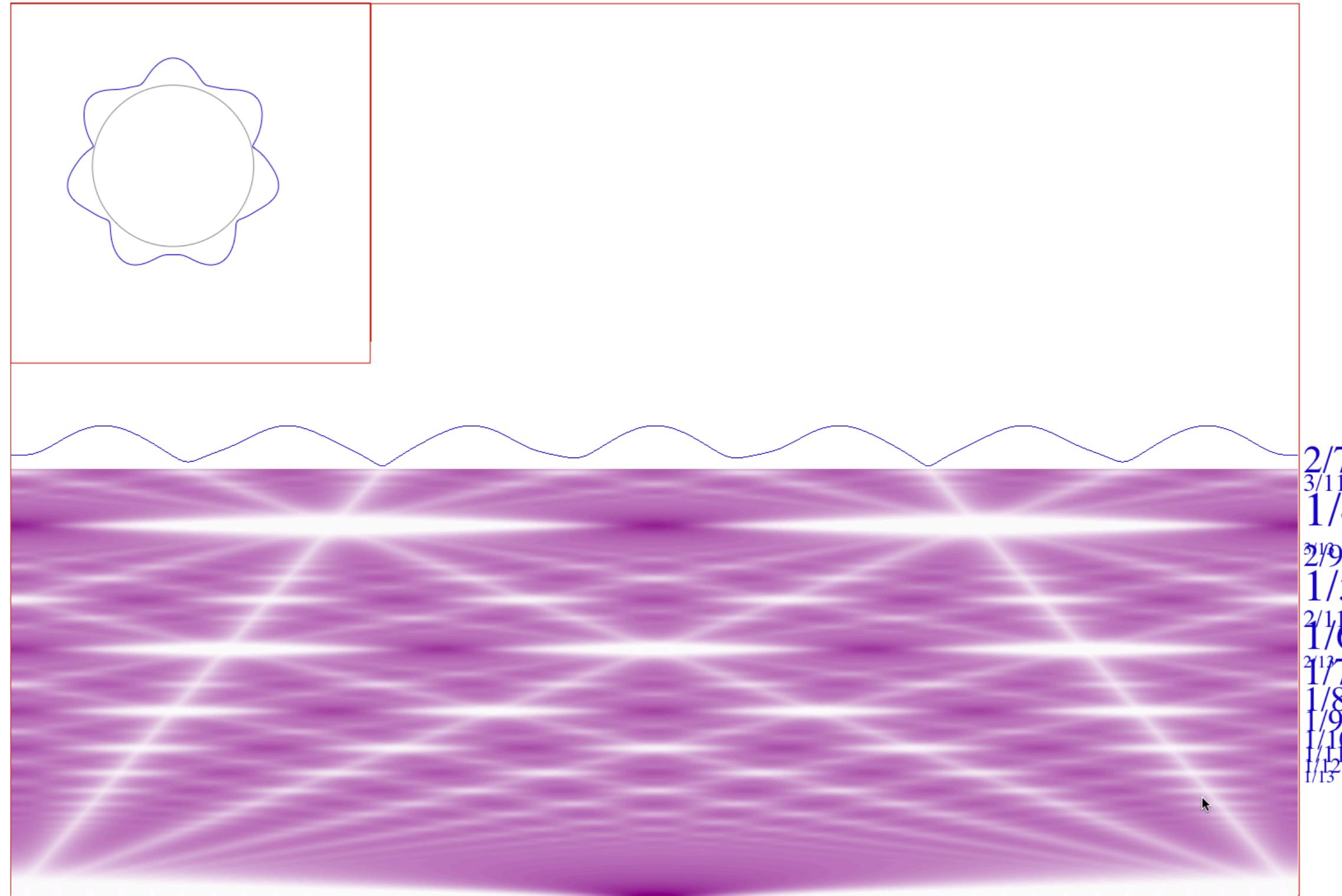
[Zero Amps](#)

T-Scale=



[...then here....](#)

time = 0.29T



[Launch](#)[Fourier Control](#)[Scenarios](#)[Pause](#)[Set T=0](#)[Zero Amps](#)

T-Scale= 1

*Set this and then click here....*Type [Quantum Carpet](#)Time Behavior [Pause at End](#)

Time Start (% Period) = 0

Time End (% Period)= 60

Del-x Width (% L) = 4

Excitation (Max n) = 20

Left (% L) = 0

Right (% L)= 100

n-Mean (% Max n)= 0

Peak1 Mean (% L)= 50

OverAll Scale = 1

Peak2 Mean (% L)= 0

Peak2 Amp (% Peak1)= 0

Draw Ring m/n Labels m-Boxcar Draw m-Bars m-Bars Max = 30

Aspect Ratio {W/H} = 1.5

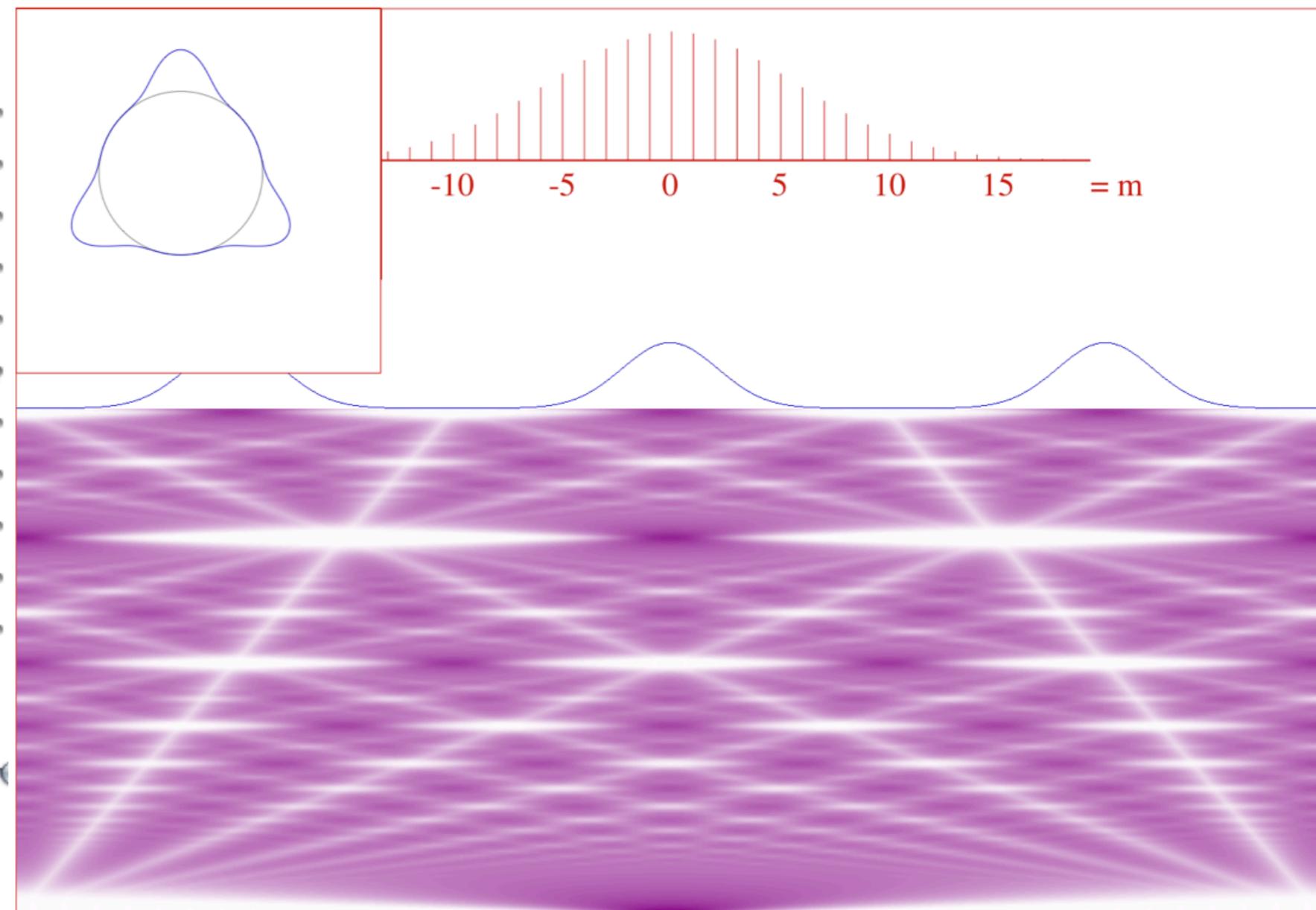
Red Level = 128

Green Level = 0

Blue Level = 128

Alpha Level = 1

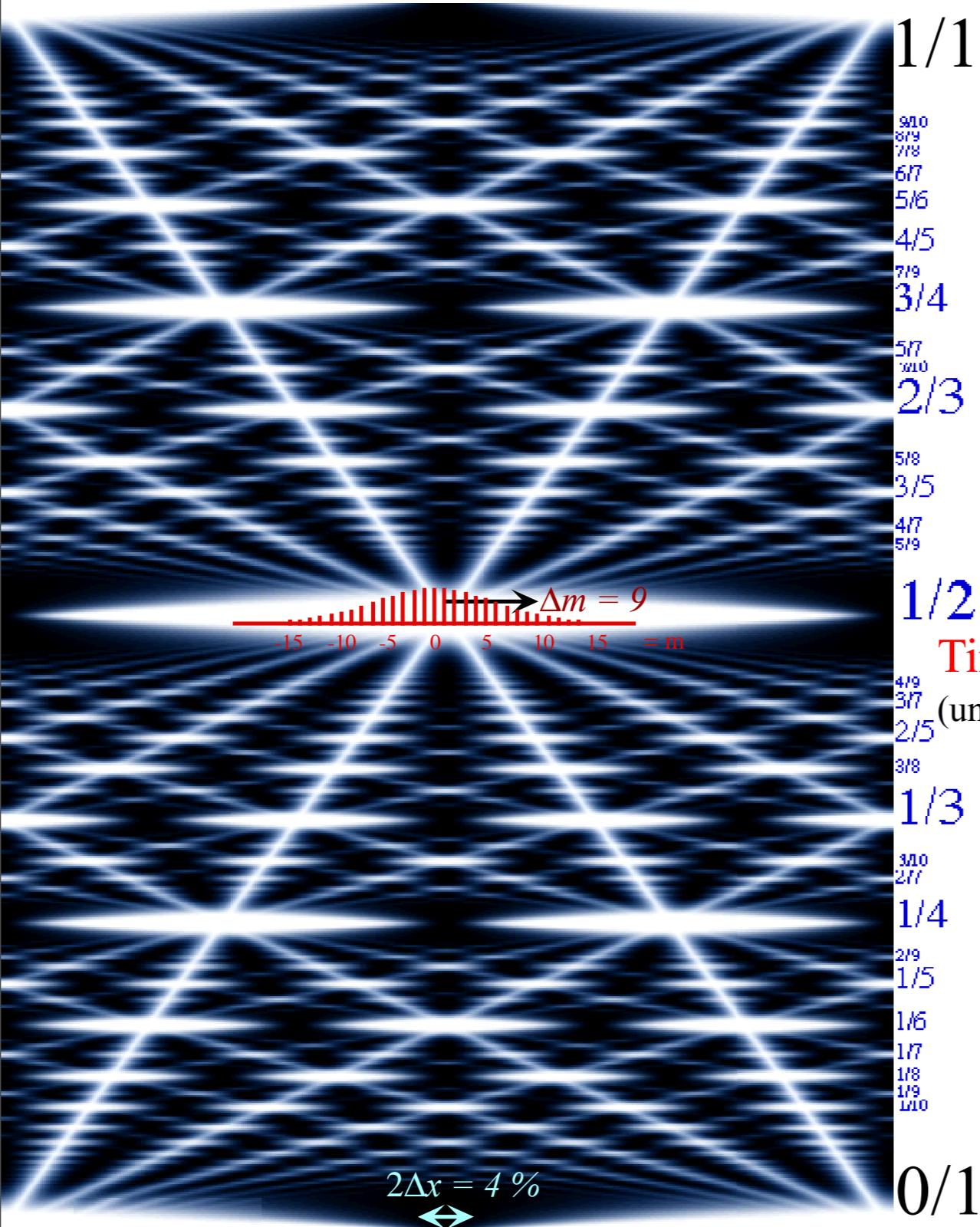
Definition Level = 0.5



1/3
2/9
3/11
1/4
2/9
1/5
2/16
1/7
1/8
1/9
1/10
1/12
1/13

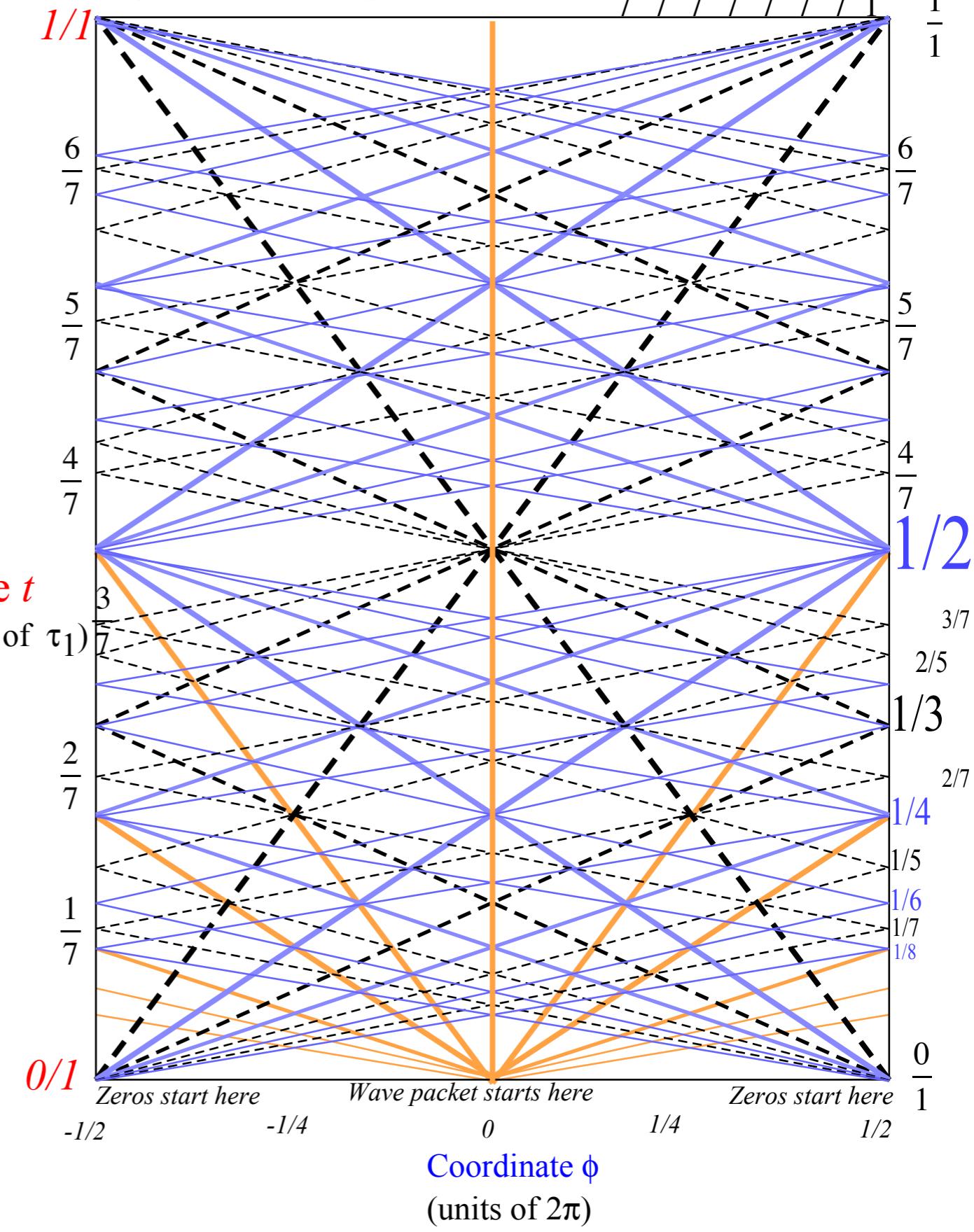
N-level-system and revival-beat wave dynamics

(9 or 10-levels ($0, \pm 1, \pm 2, \pm 3, \pm 4, \dots, \pm 9, \pm 10, \pm 11 \dots$) excited)



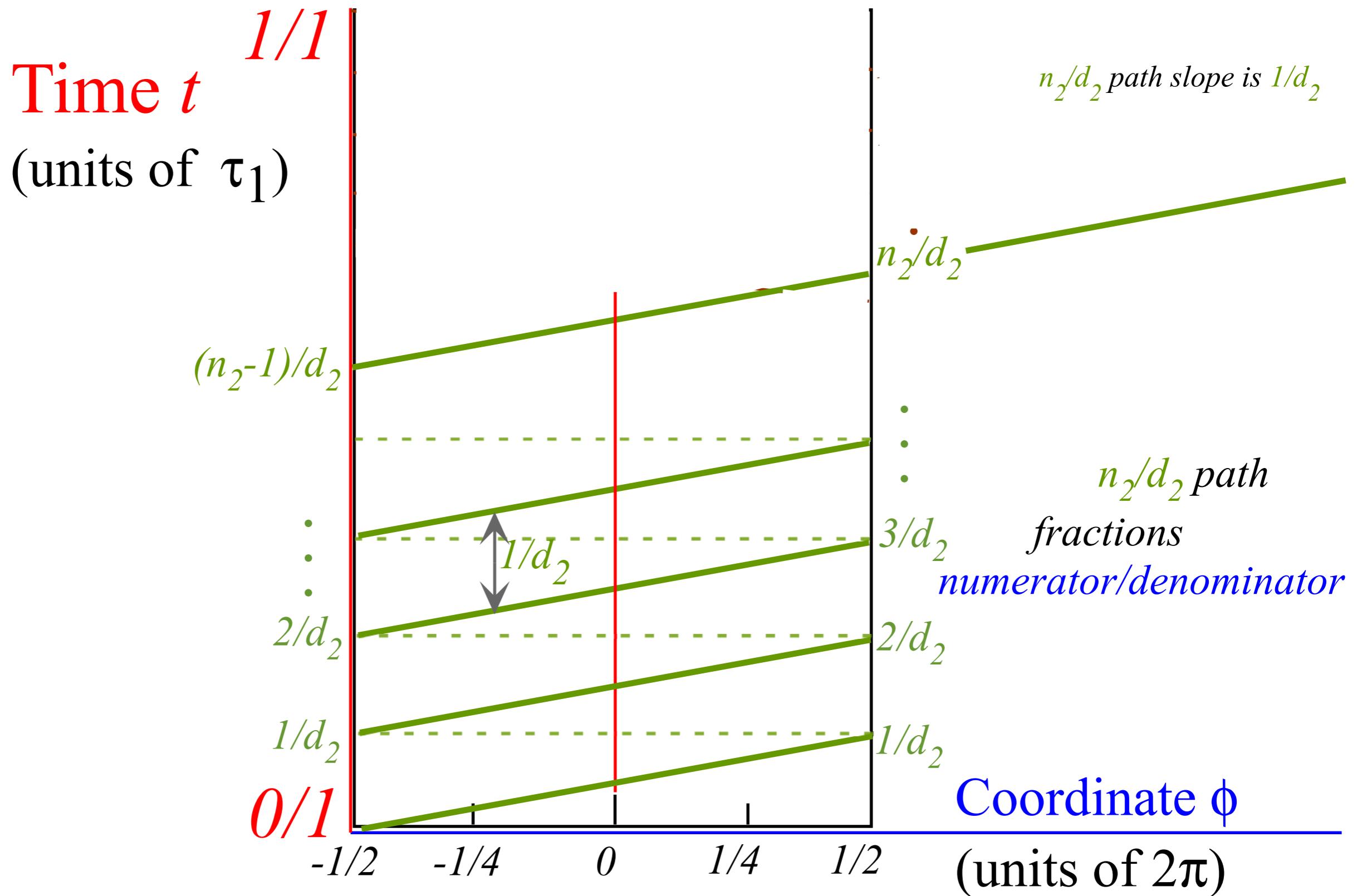
Zeros (clearly) and “particle-packets” (faintly) have paths labeled by fraction sequences like:

$$\frac{0}{7}, \frac{1}{7}, \frac{2}{7}, \frac{3}{7}, \frac{4}{7}, \frac{5}{7}, \frac{6}{7}, \frac{1}{1}$$



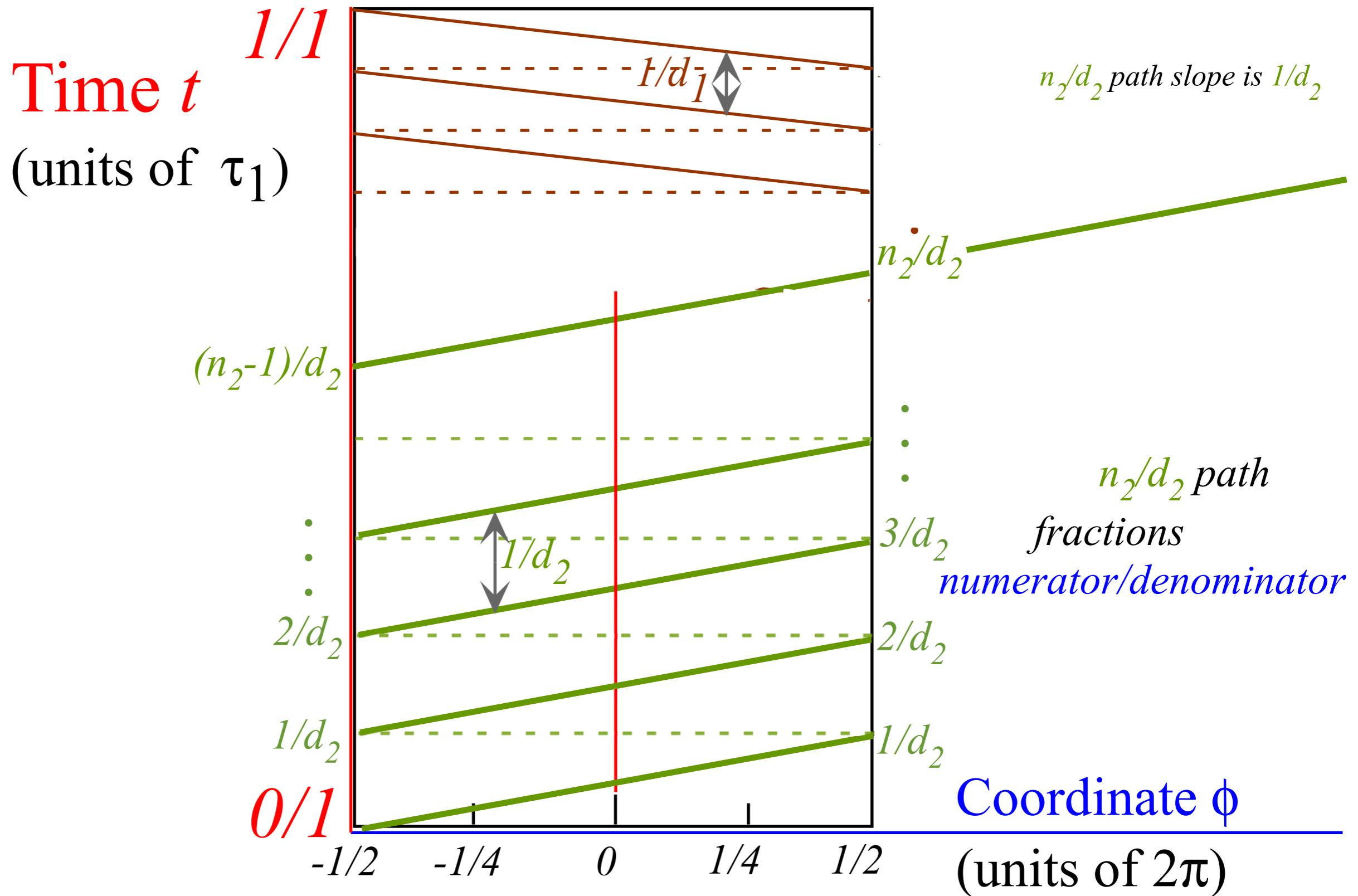
Farey Sum algebra of revival-beat wave dynamics

Label by *numerators N* and *denominators D* of rational fractions N/D



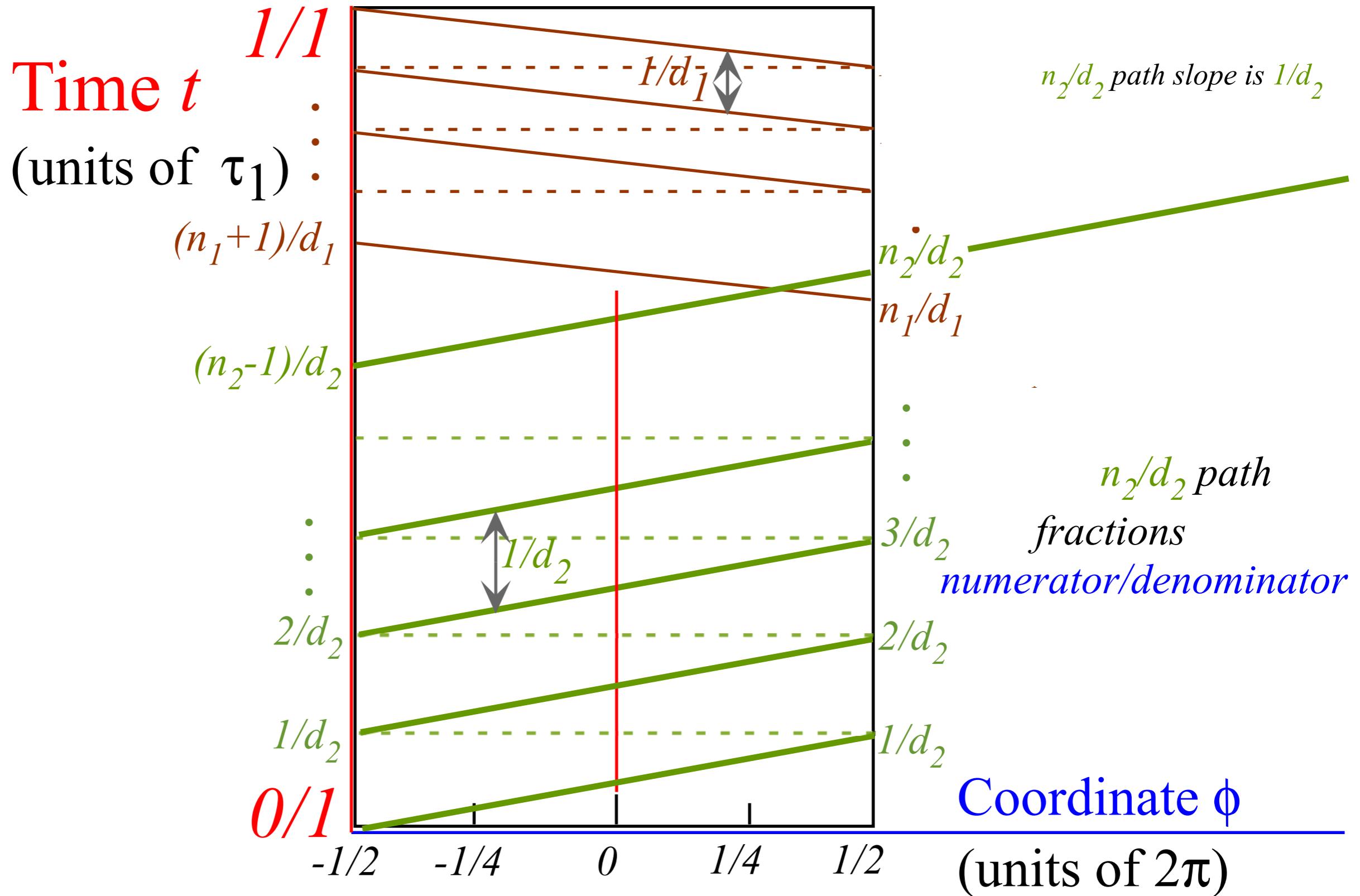
Farey Sum algebra of revival-beat wave dynamics

Label by numerators N and denominators D of rational fractions N/D



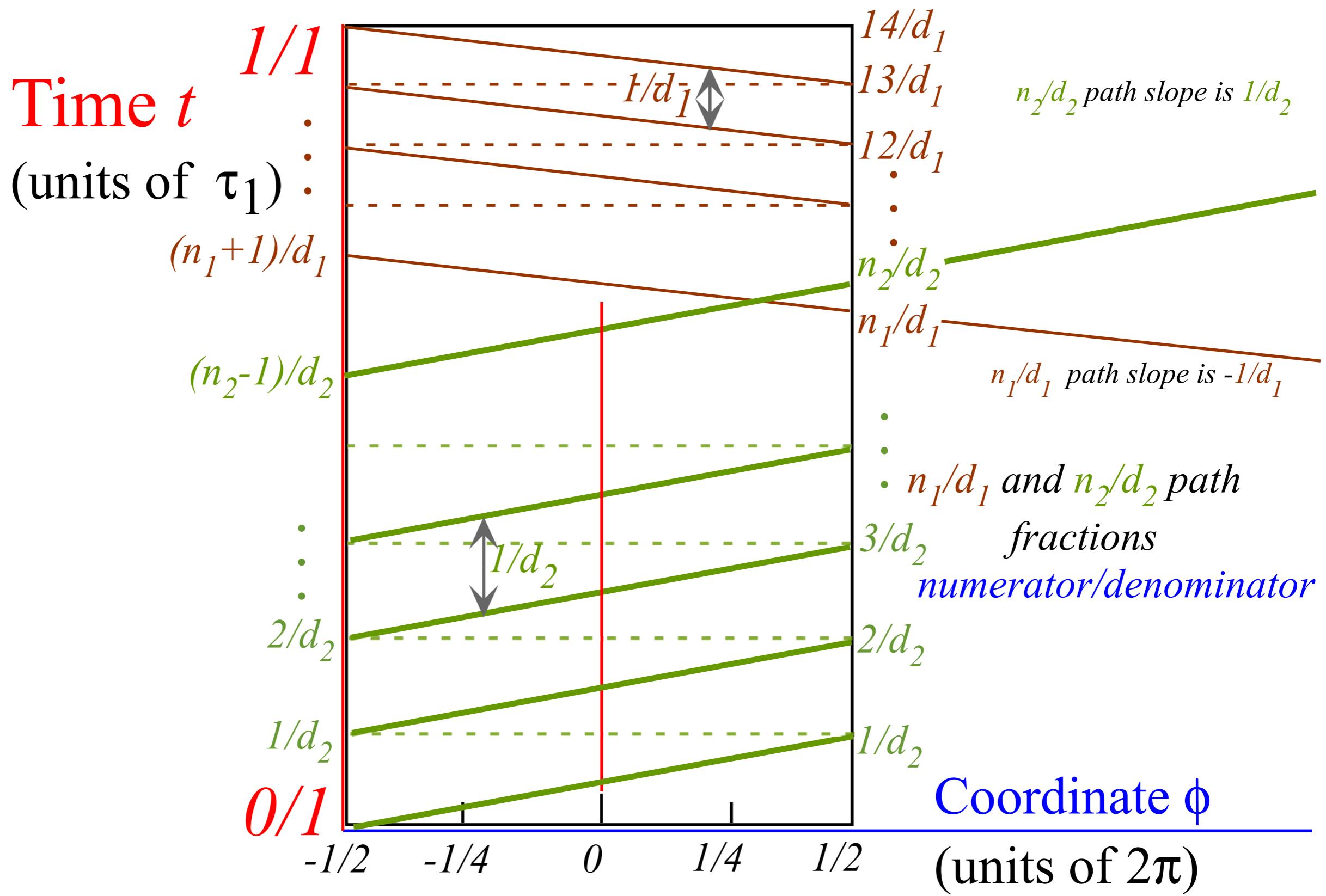
Farey Sum algebra of revival-beat wave dynamics

Label by numerators N and denominators D of rational fractions N/D



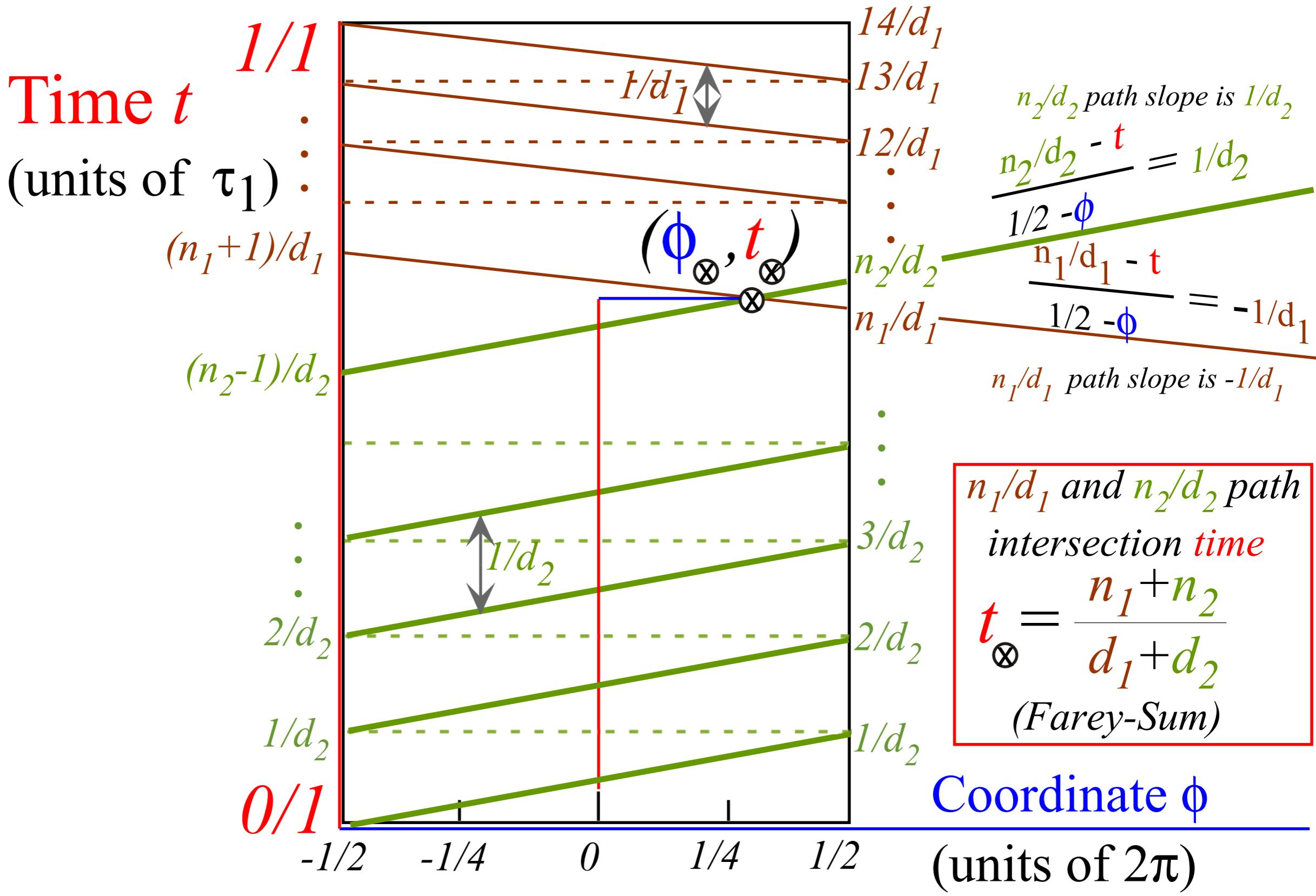
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Label by numerators N and denominators D of rational fractions N/D



Farey Sum algebra of revival-beat wave dynamics

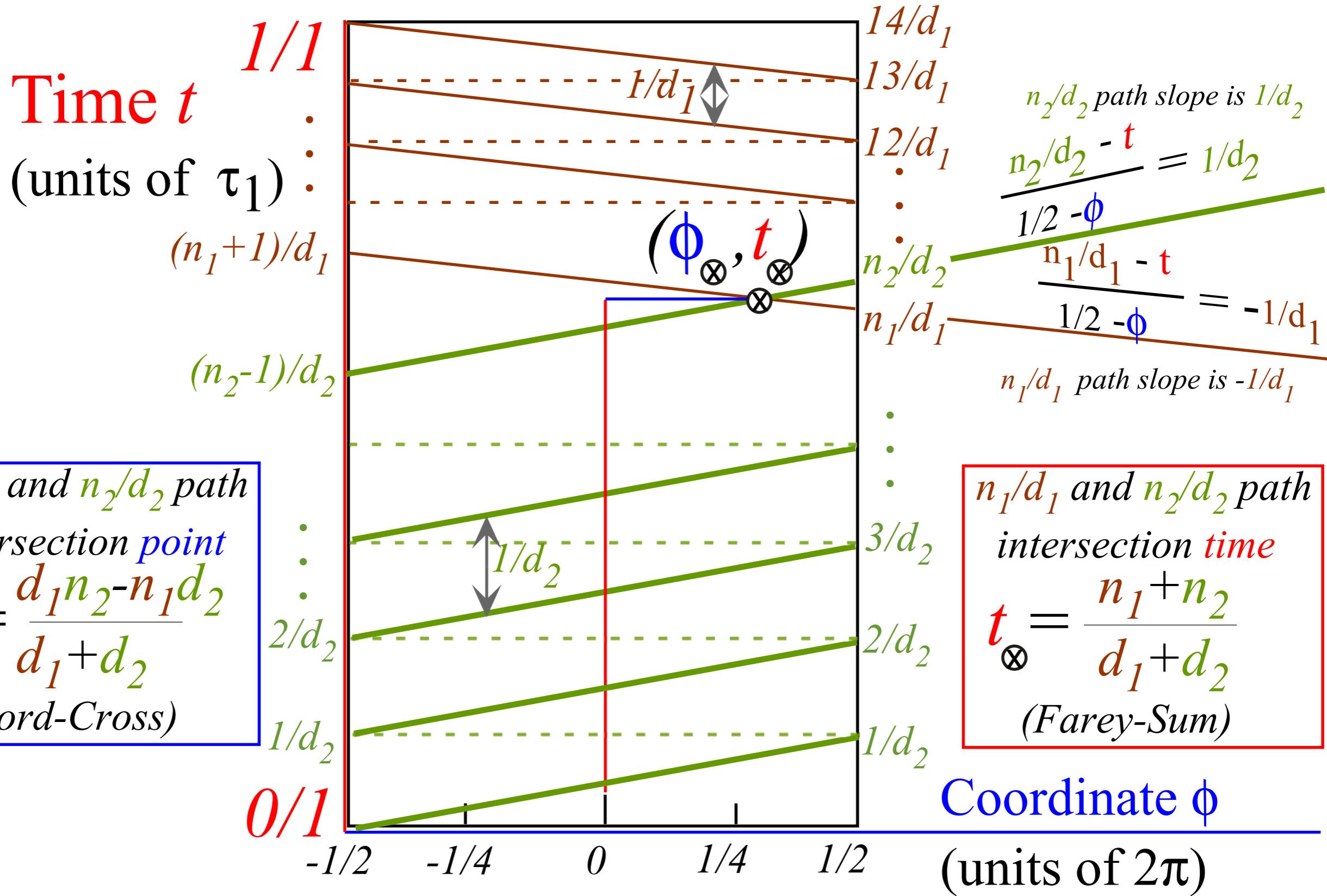
Label by numerators N and denominators D of rational fractions N/D



[John Farey, Phil. Mag. (1816)]

Farey Sum algebra of revival-beat wave dynamics

Label by numerators N and denominators D of rational fractions N/D



[Lester R. Ford, Am. Math. Monthly 45, 586(1938)]

[John Farey, Phil. Mag.(1816)]

“Monster Mash” classical segue to Heisenberg action relations

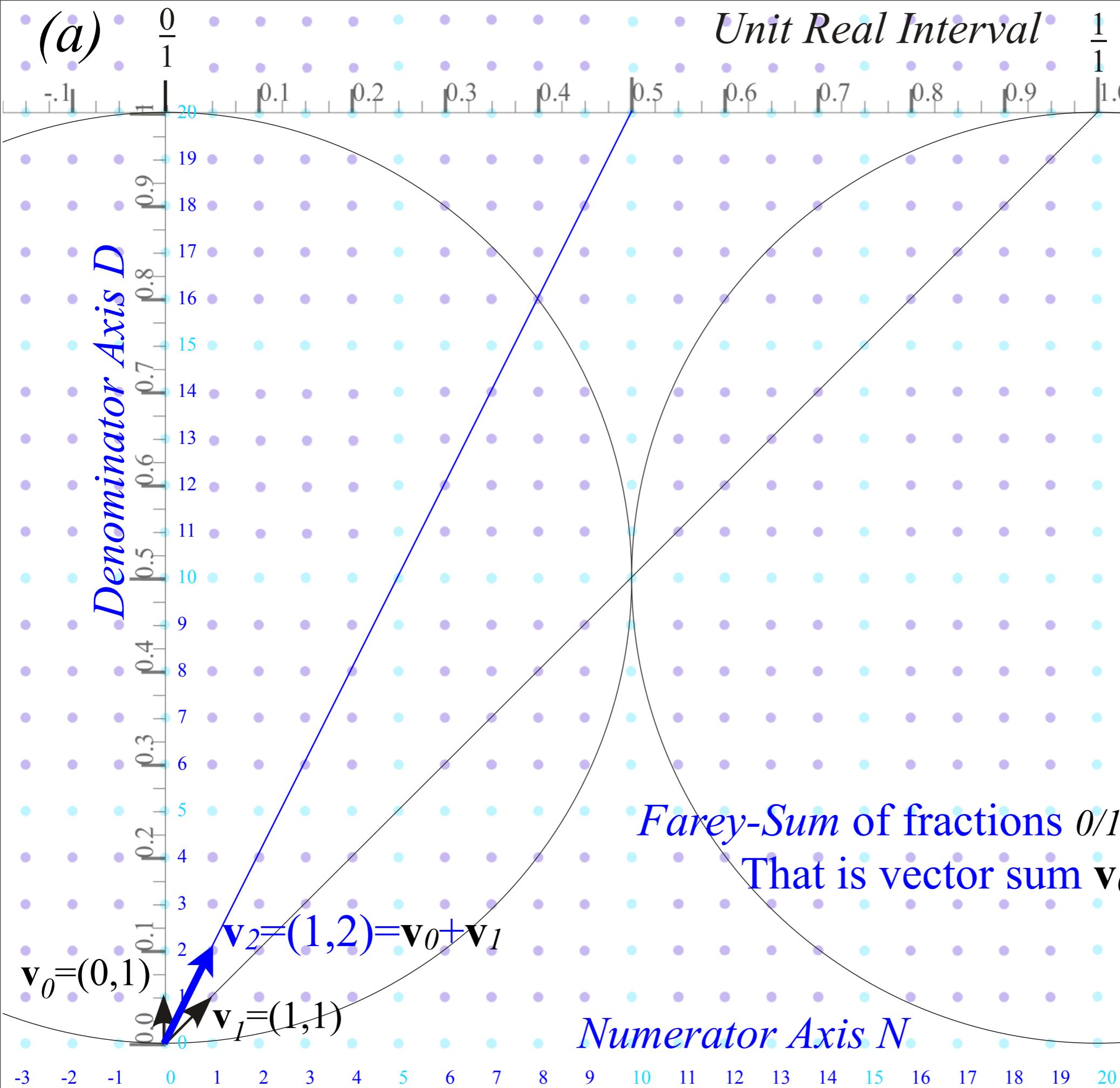
Example of very very large M_1 ball-walls crushing a poor little m_2

How m_2 keeps its action

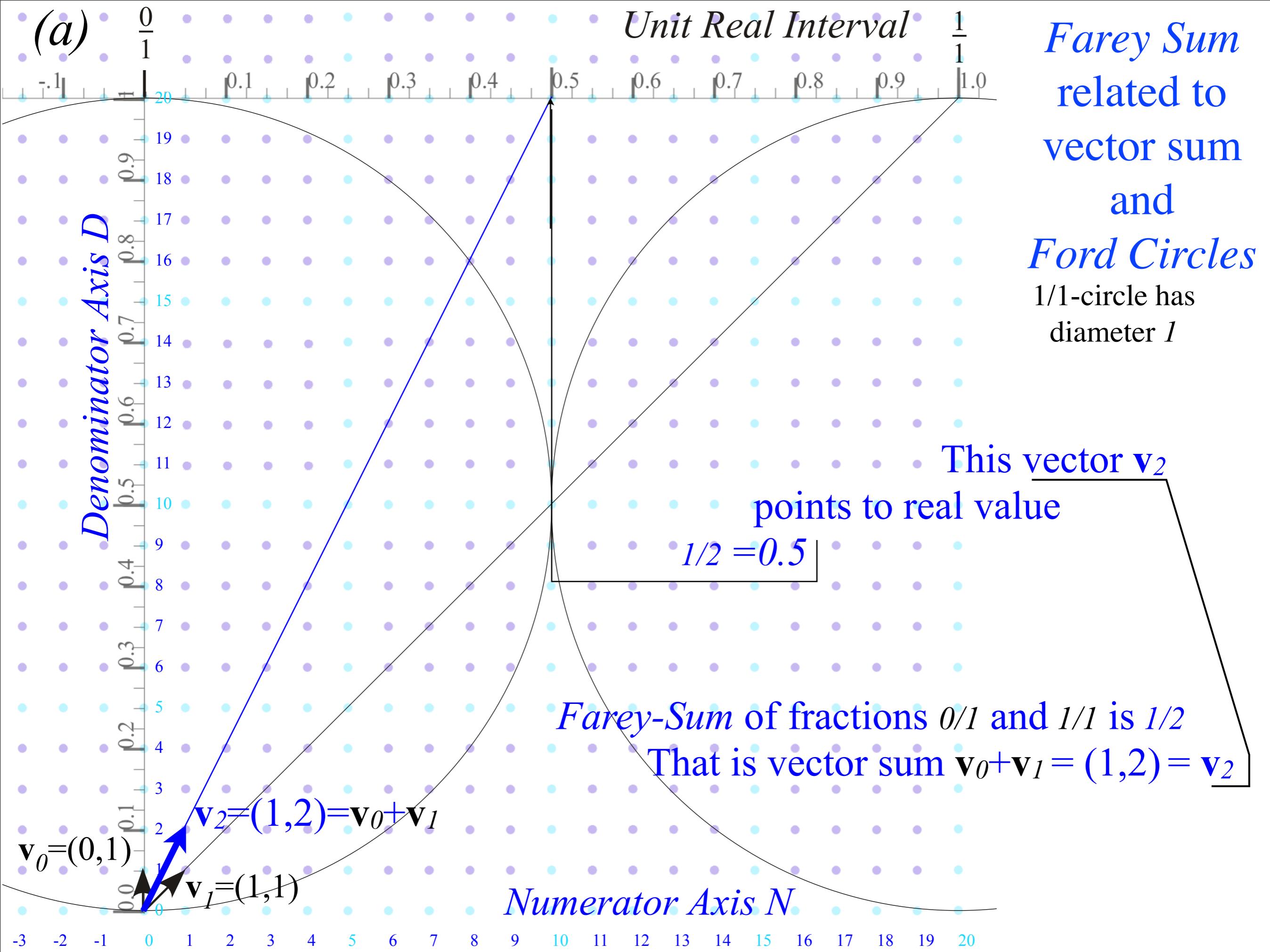
An interesting wave analogy: The “Tiny-Big-Bang” [Harter, J. Mol. Spec. 210, 166-182 (2001)], [Harter, Li IMSS (2012)]

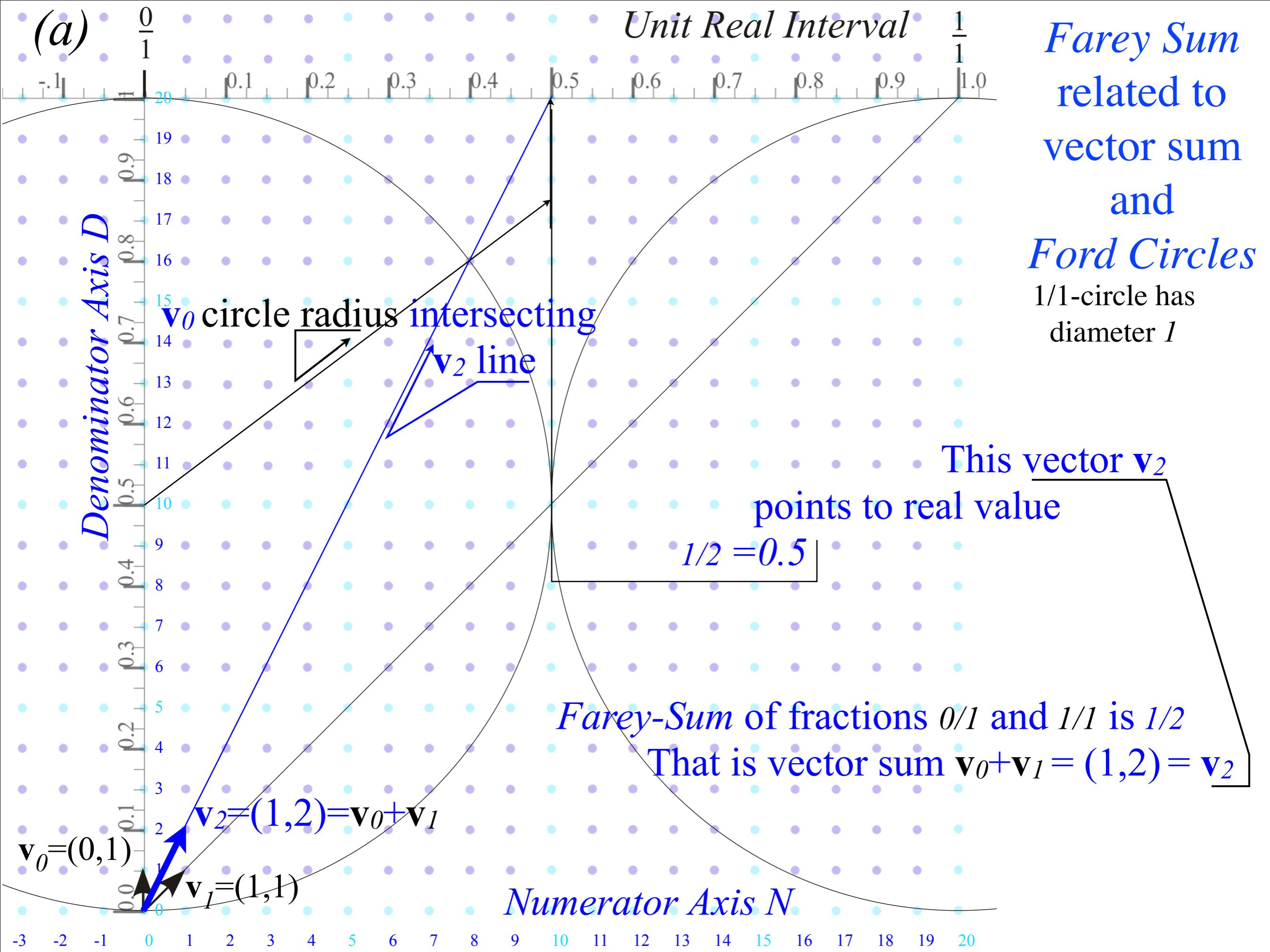
A lesson in geometry of fractions and fractals: Ford Circles and Farey Sums

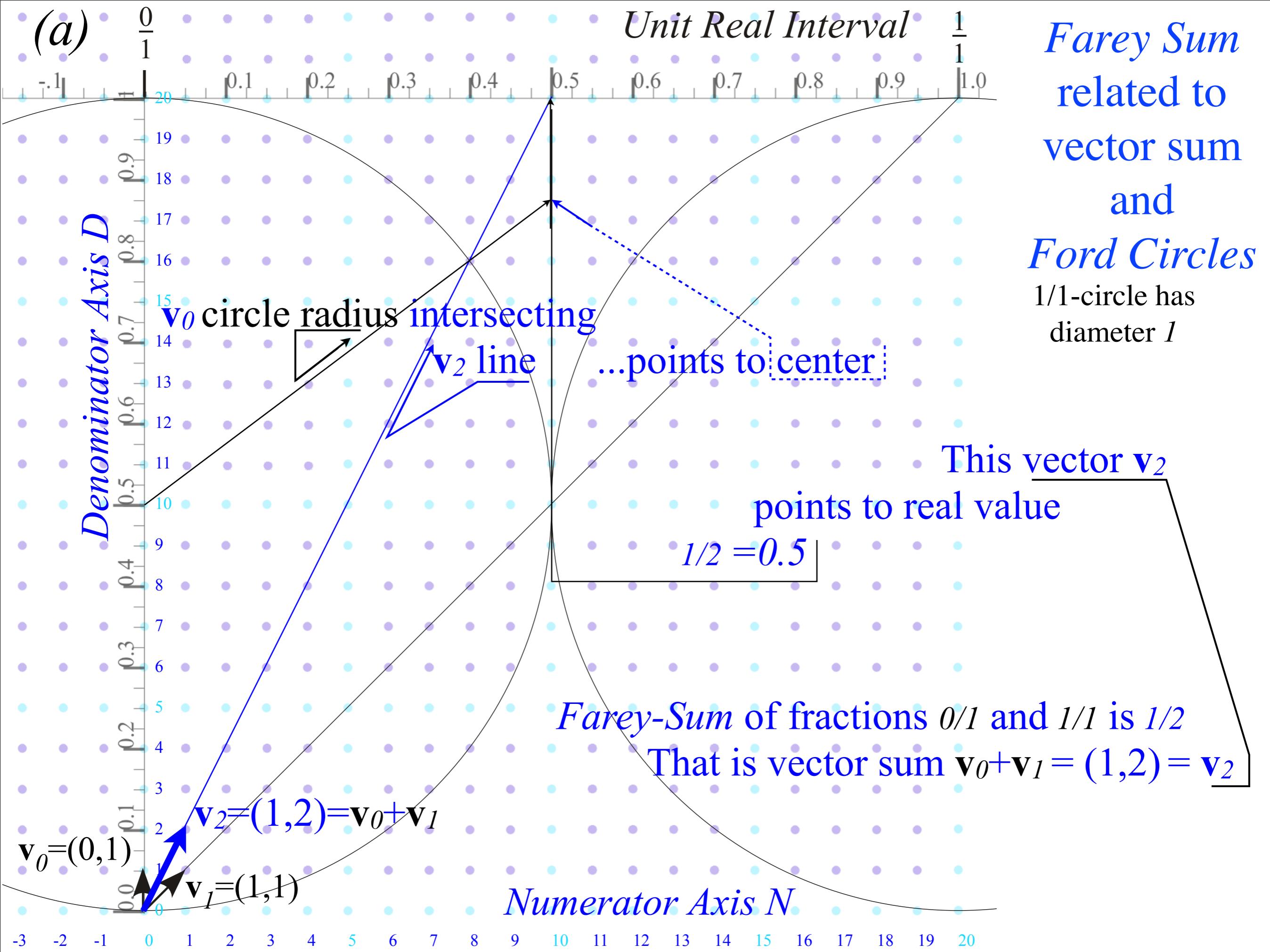
[Lester R. Ford, Am. Math. Monthly 45, 586(1938)] [John Farey, Phil. Mag.(1816)]

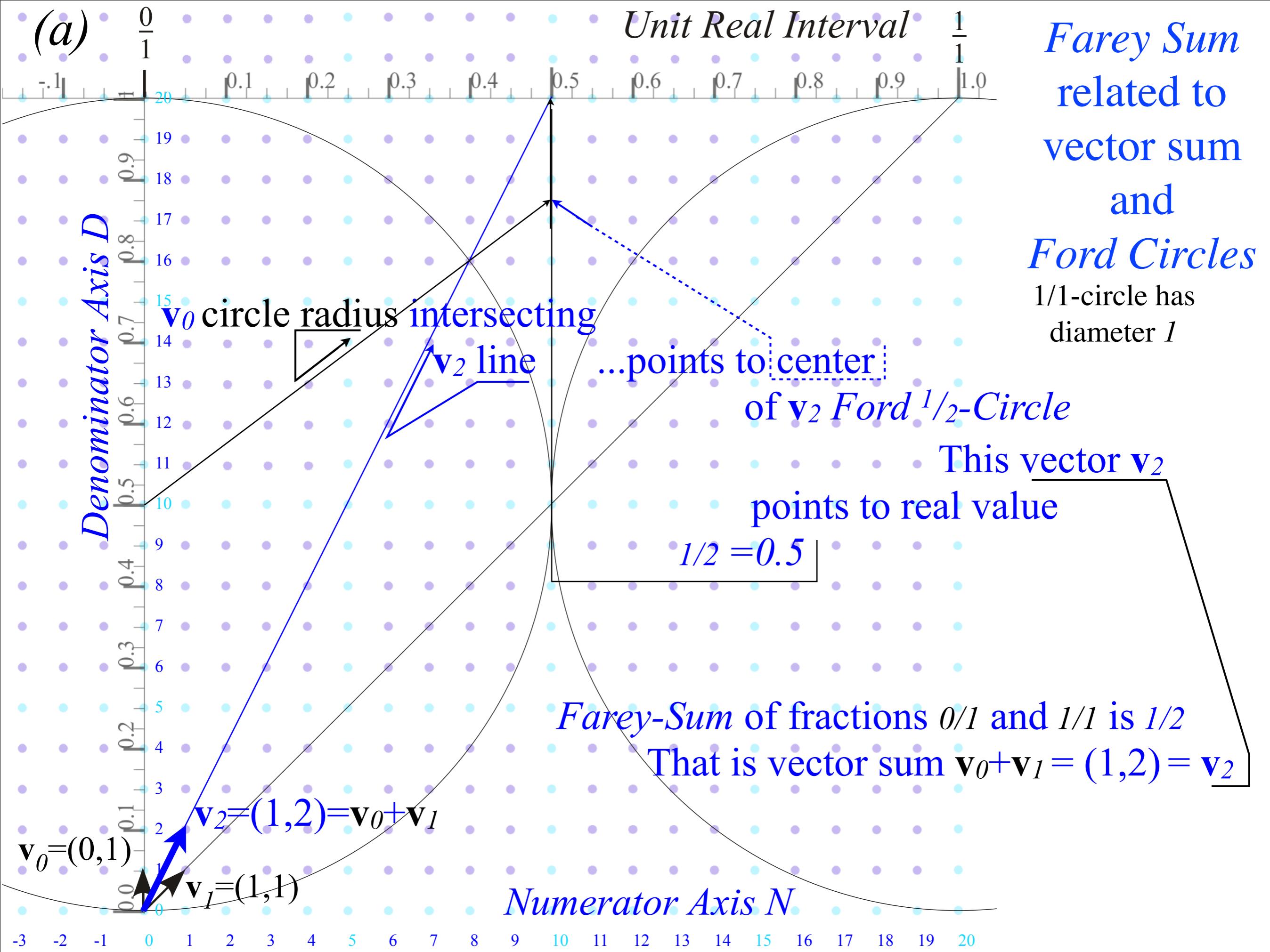


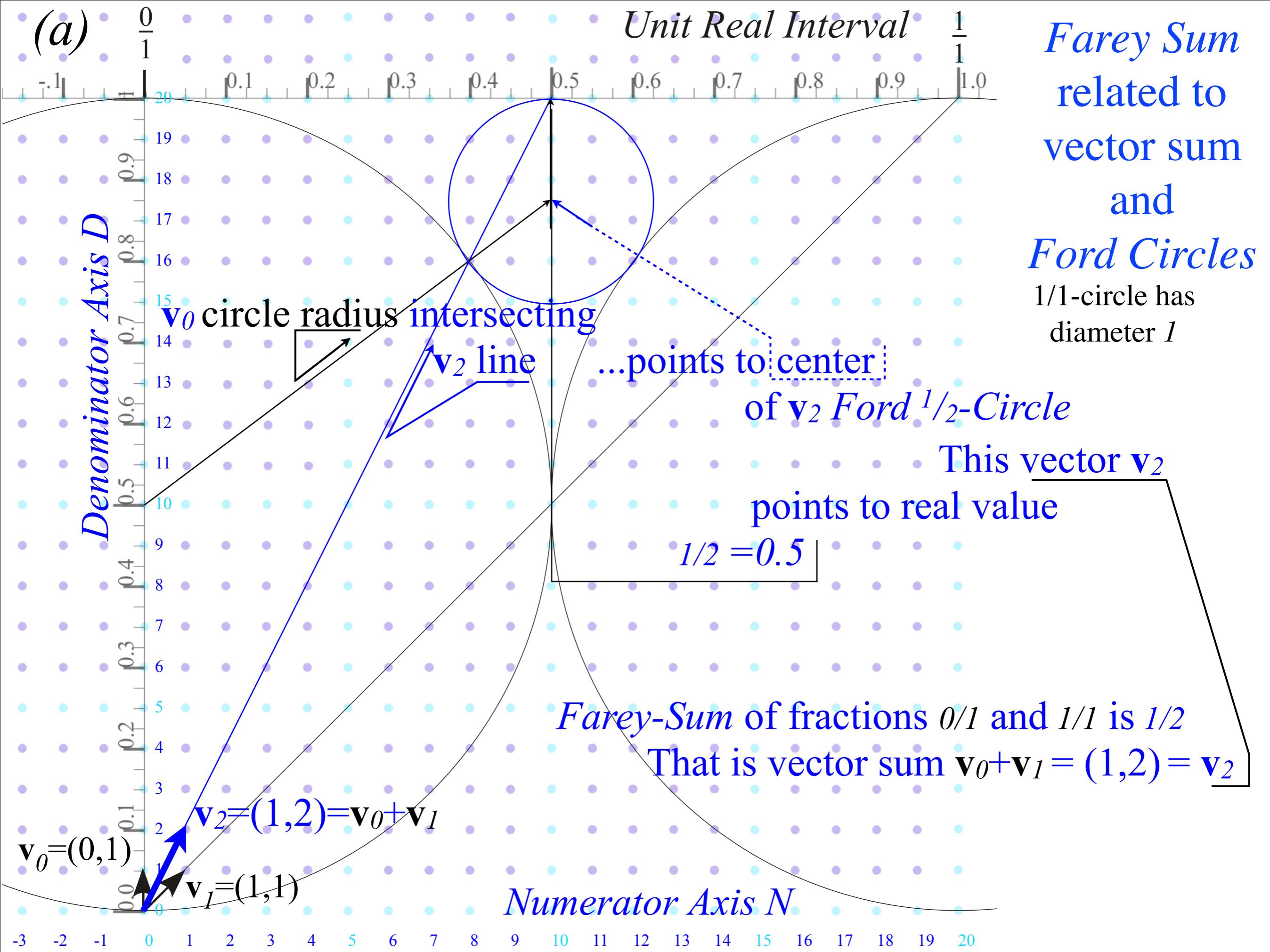
Farey Sum
related to
vector sum
and
Ford Circles
1/1-circle has
diameter 1

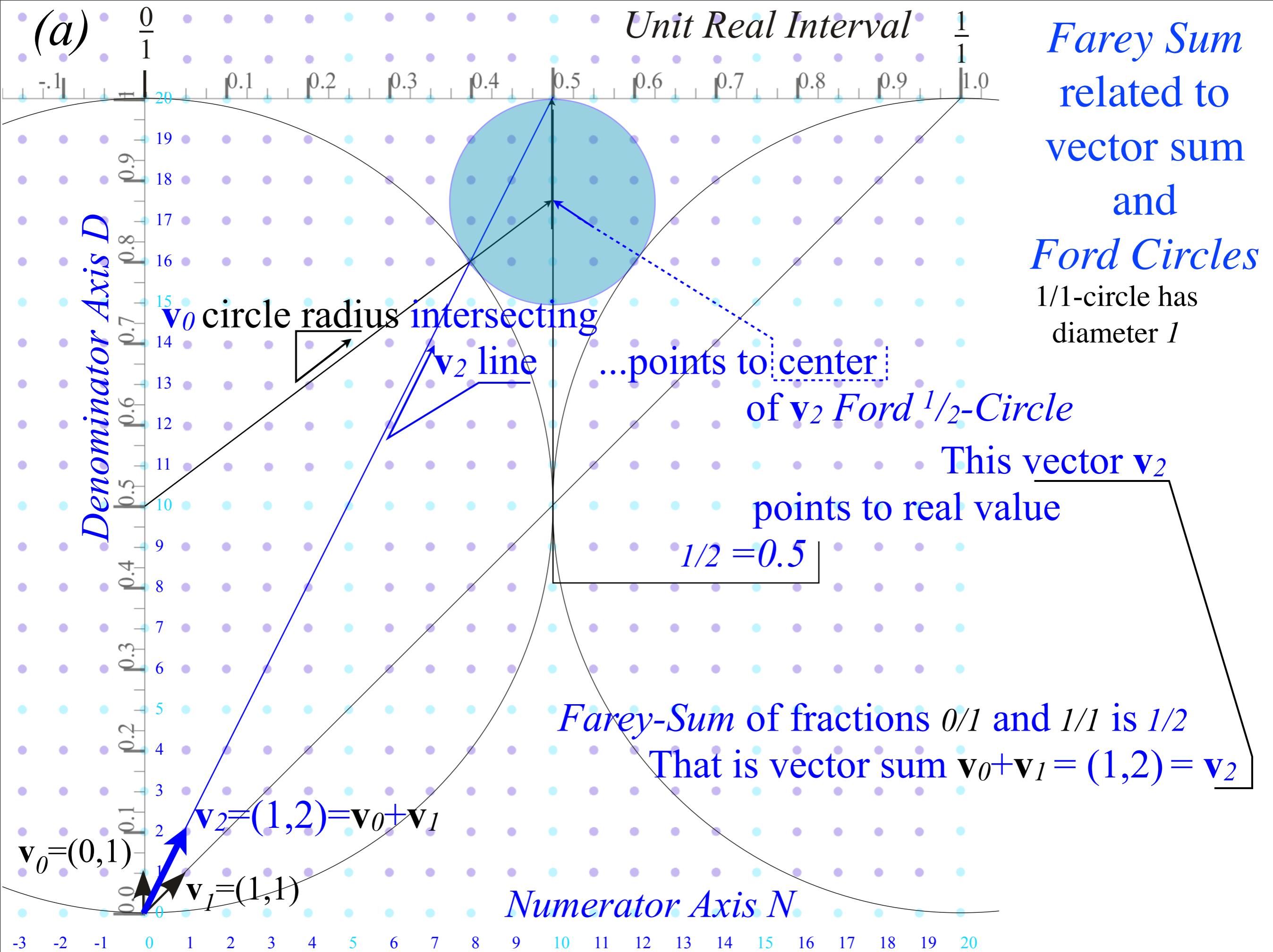


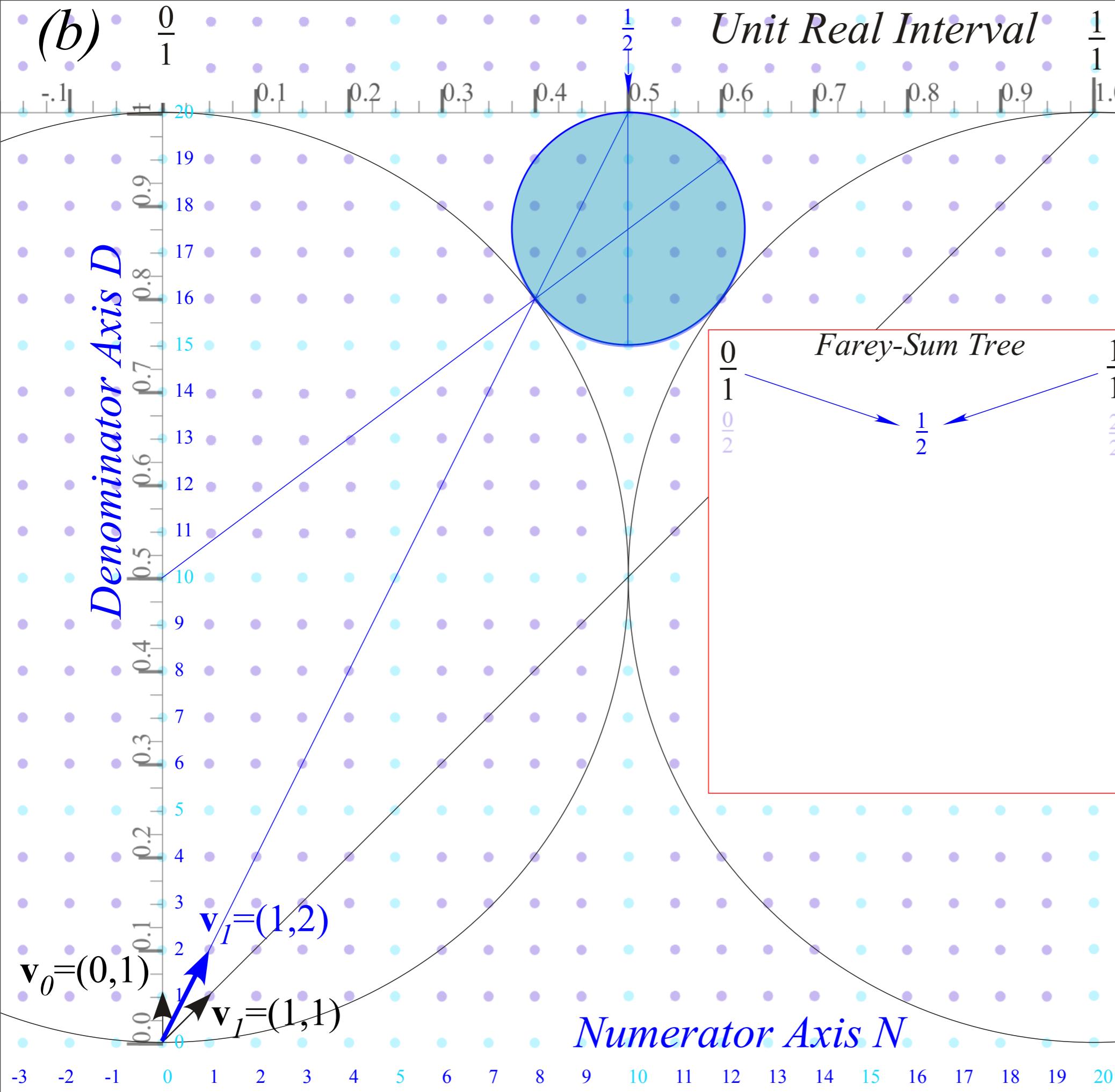








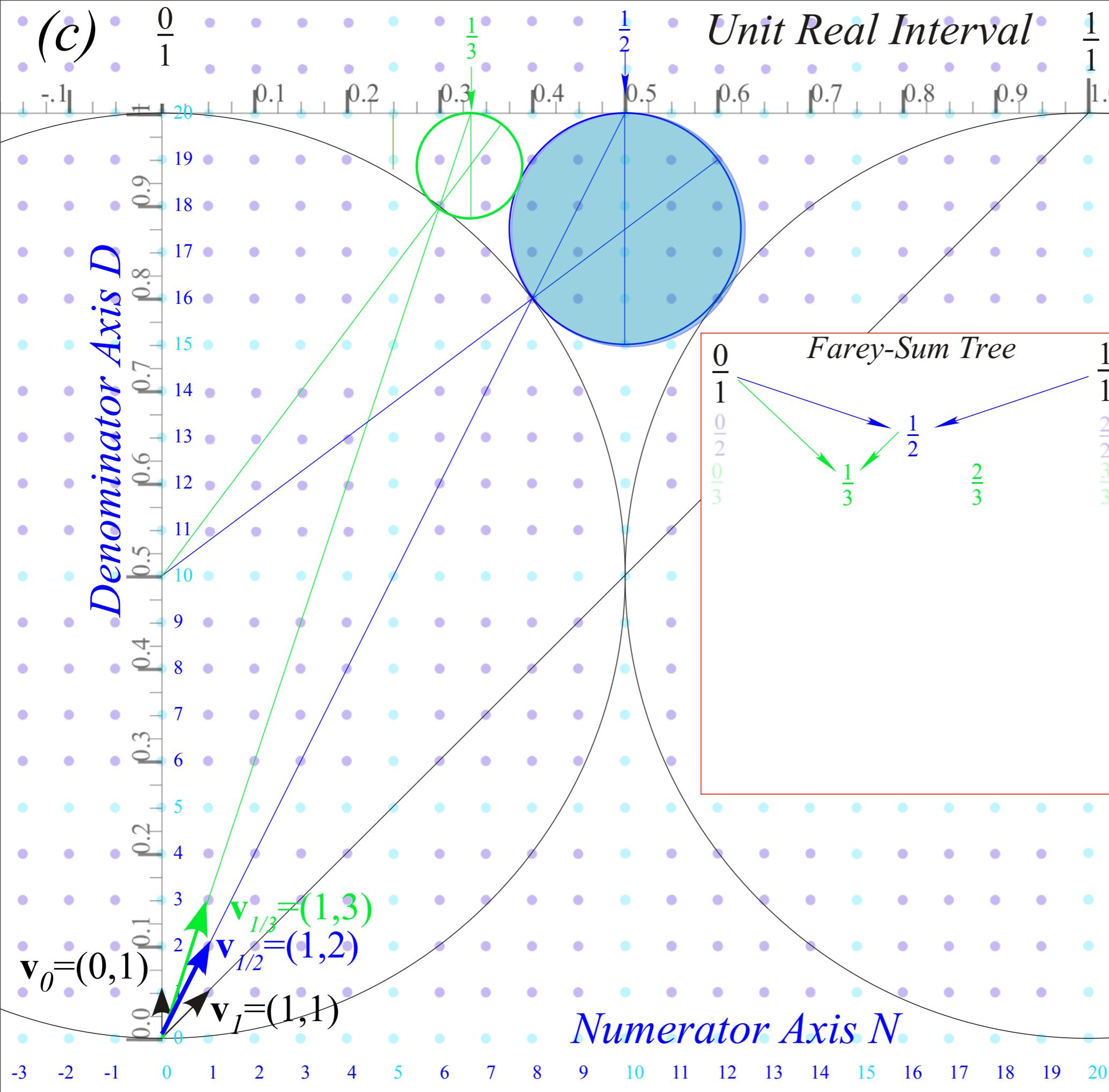


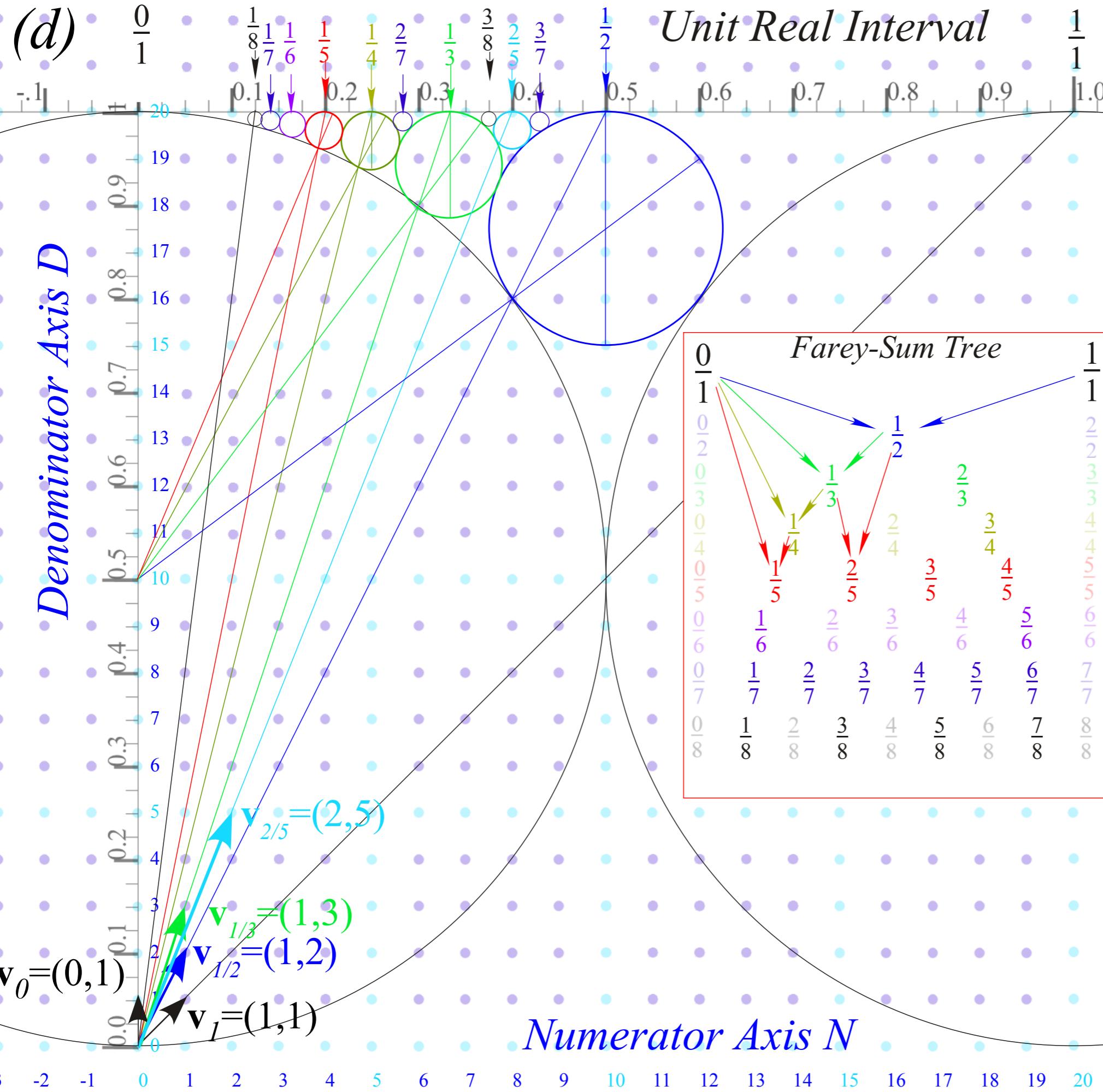


Farey Sum
related to
vector sum
and
Ford Circles

1/1-circle has diameter 1

1/2-circle has diameter $1/2^2=1/4$





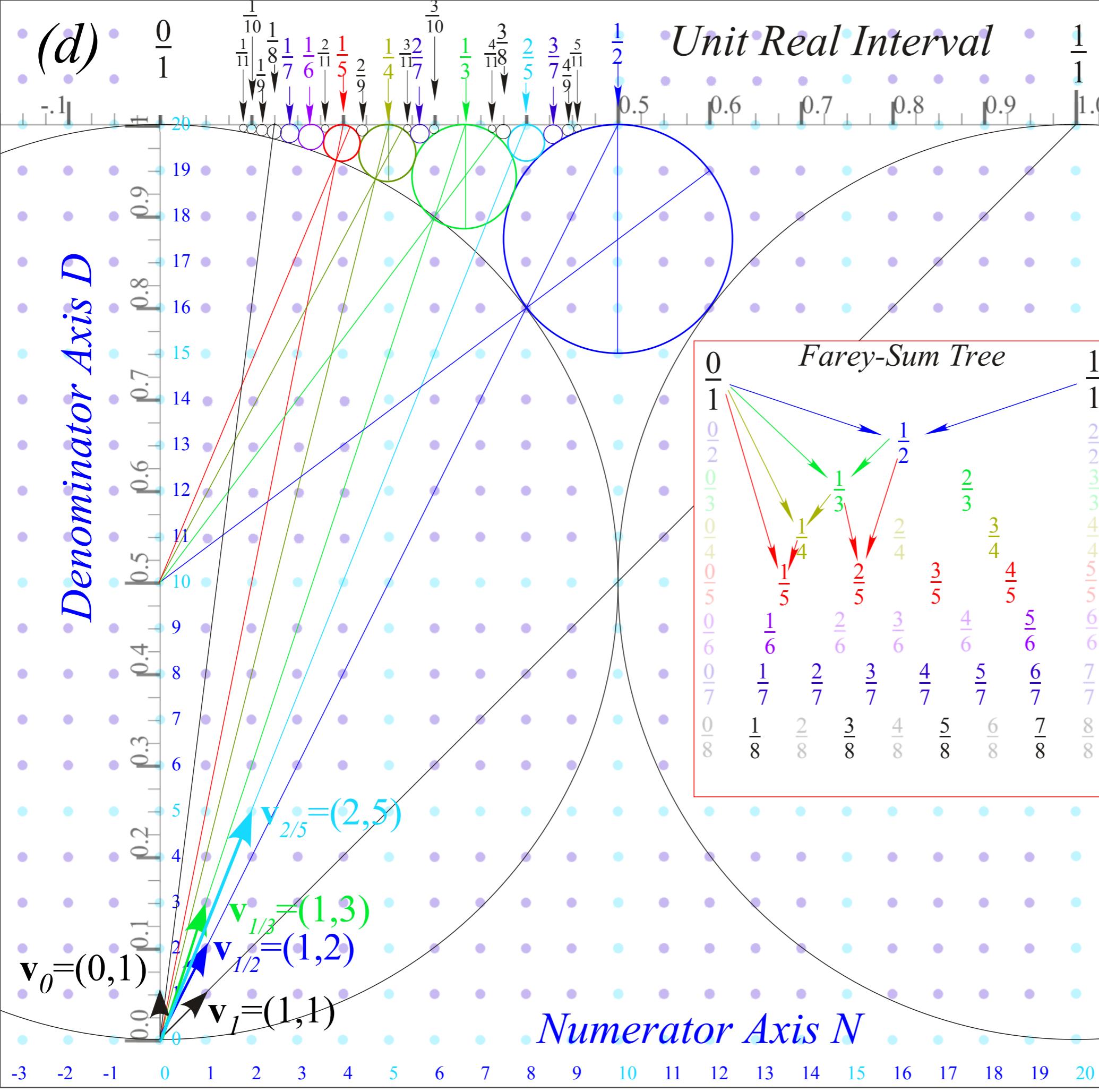
Farey Sum
related to
vector sum
and
Ford Circles

1/2-circle has
diameter $1/2^2 = 1/4$

1/3-circles have
diameter $1/3^2 = 1/9$

n/d-circles have
diameter $1/d^2$

*Farey Sum
related to
vector sum
and
Ford Circles*

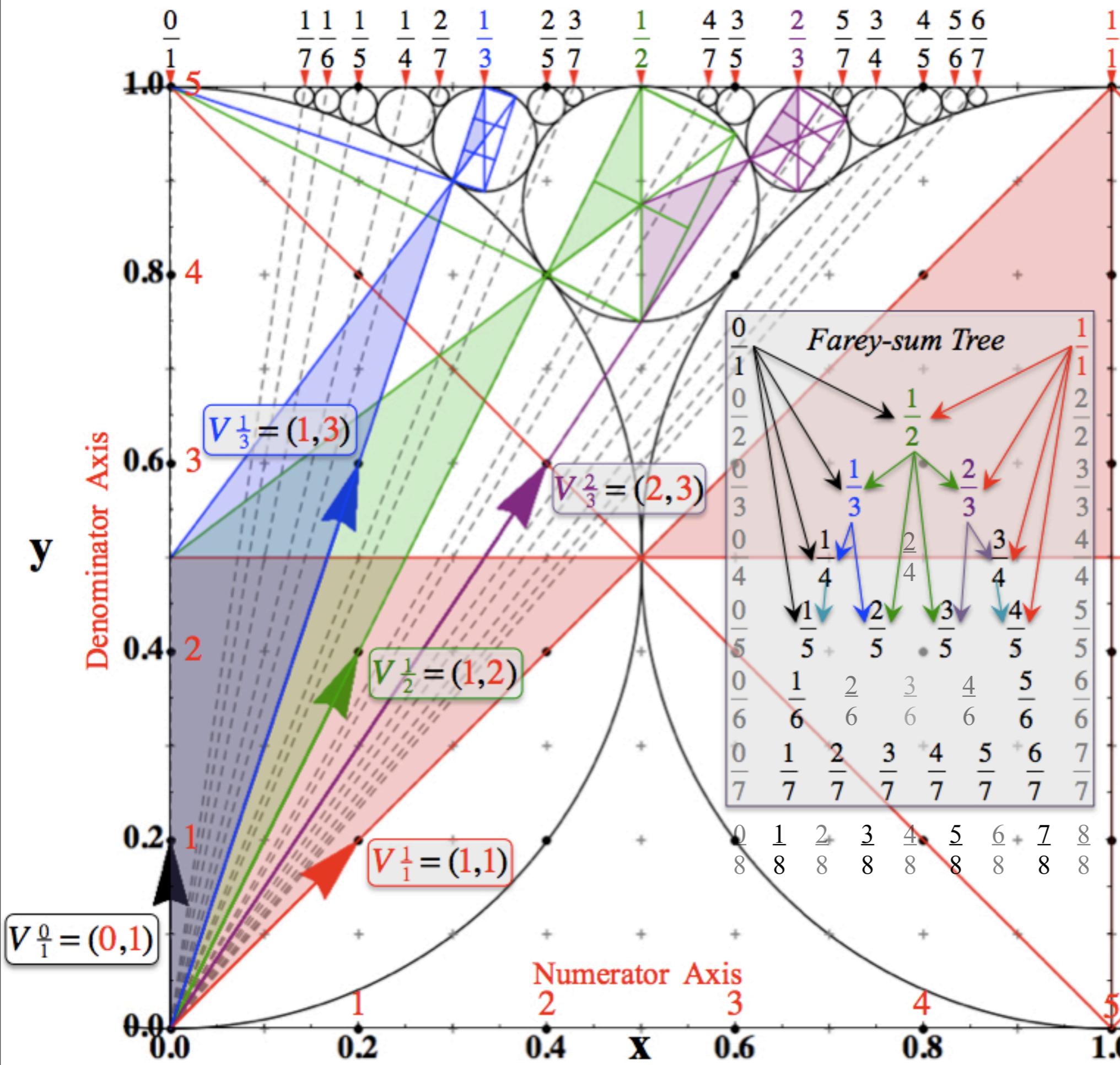


1/2-circle has diameter $1/2^2 = 1/4$

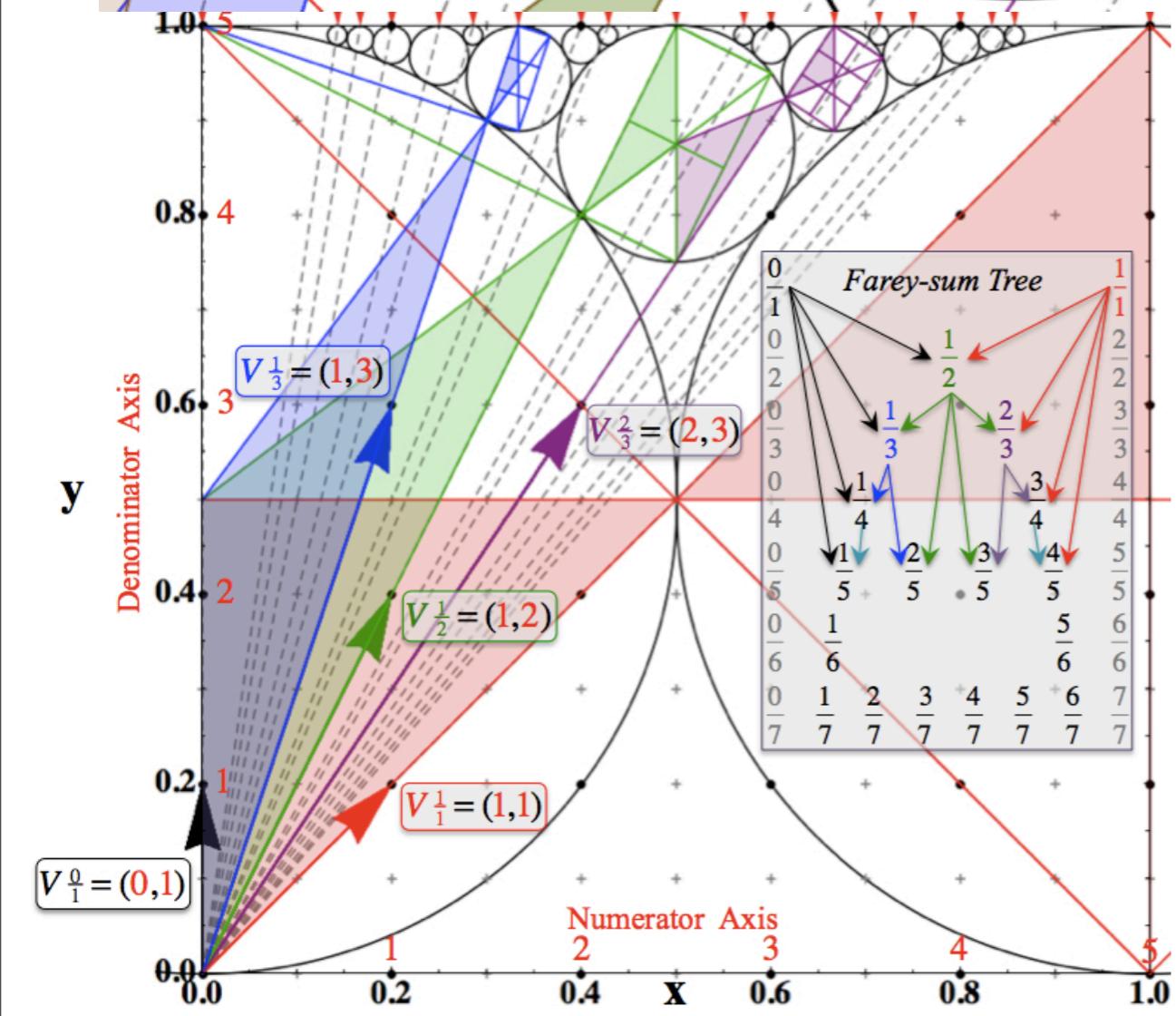
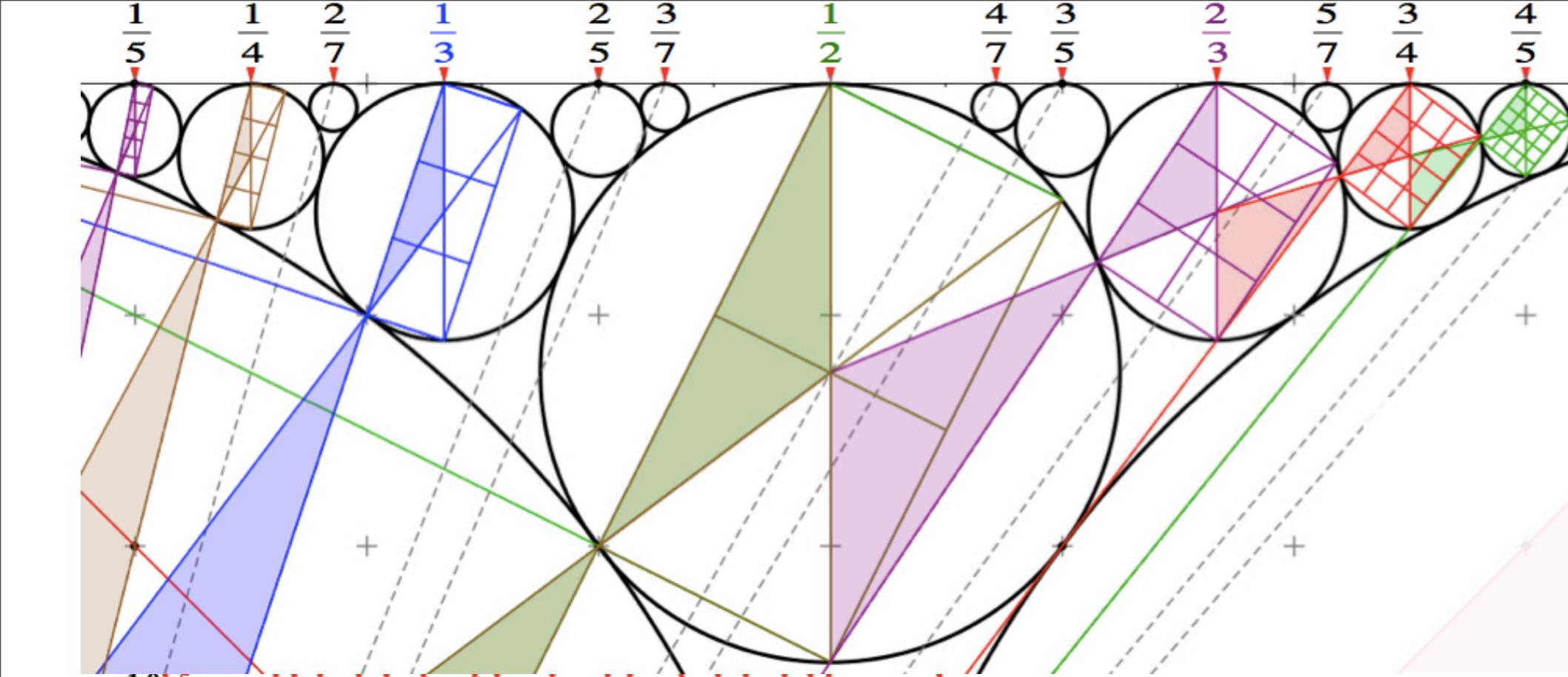
1/3-circles have diameter $1/3^2 = 1/9$

n/d -circles have diameter $1/d^2$

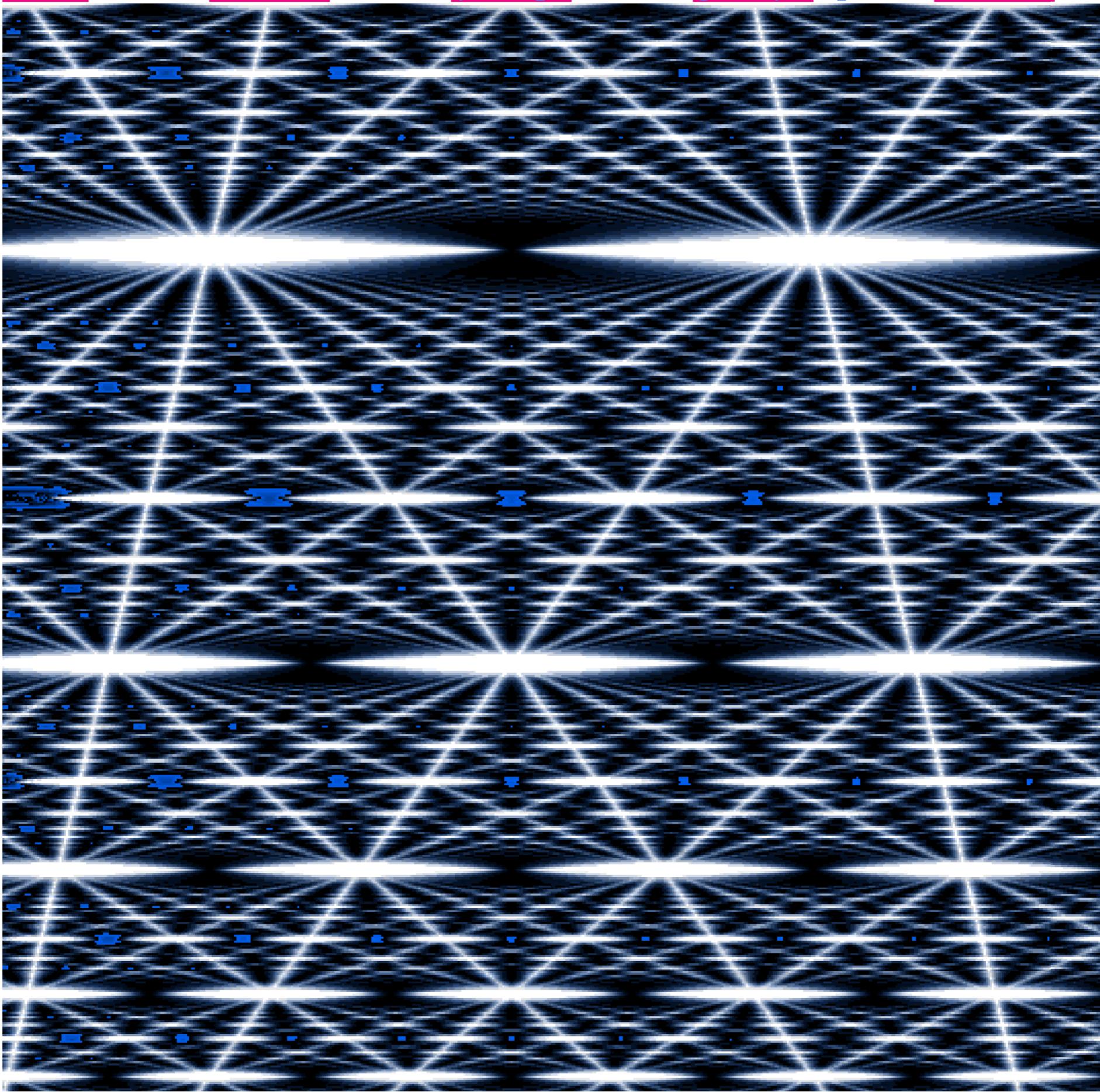
Thales
Rectangles
provide
analytic geometry
of
fractal structure



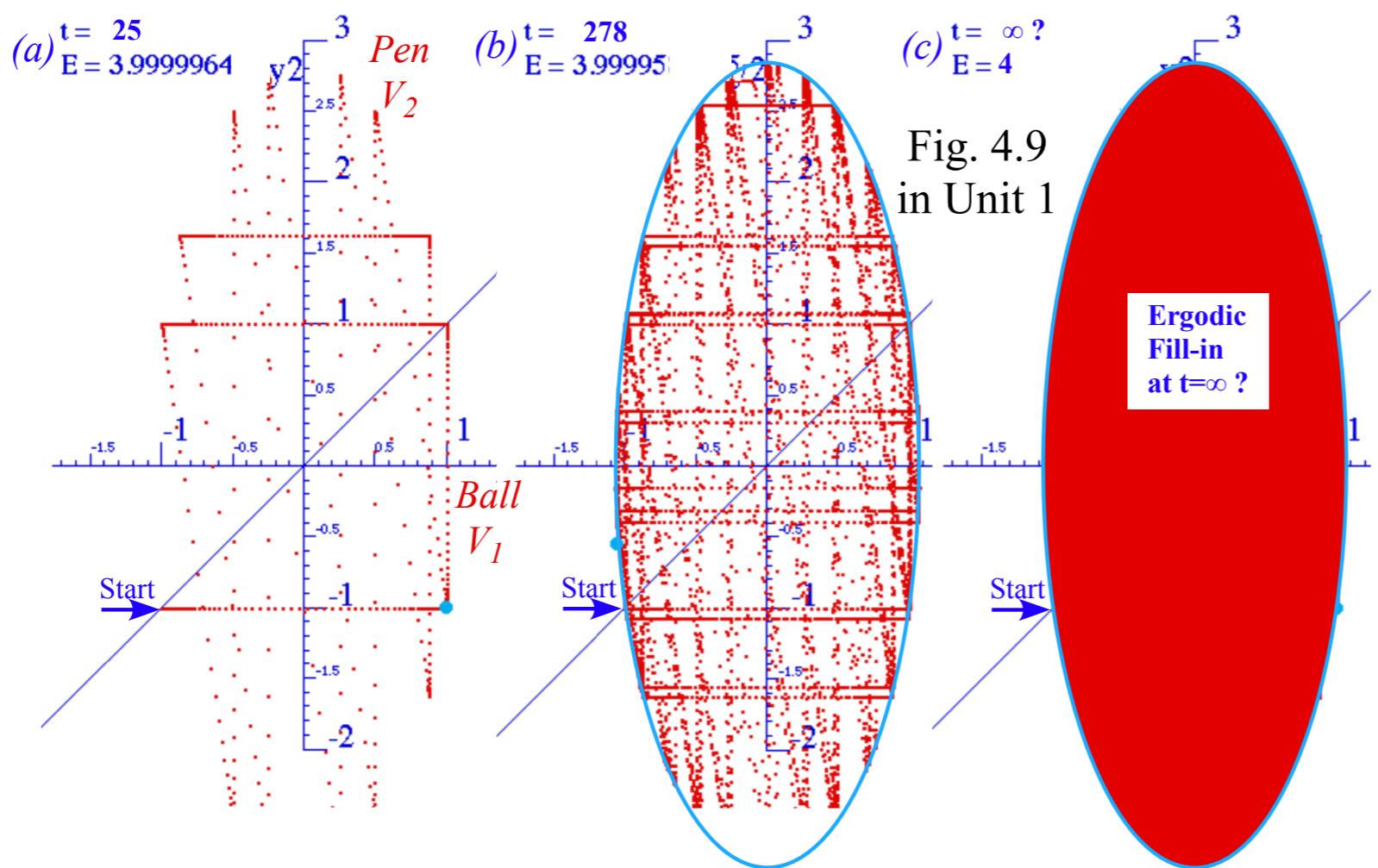
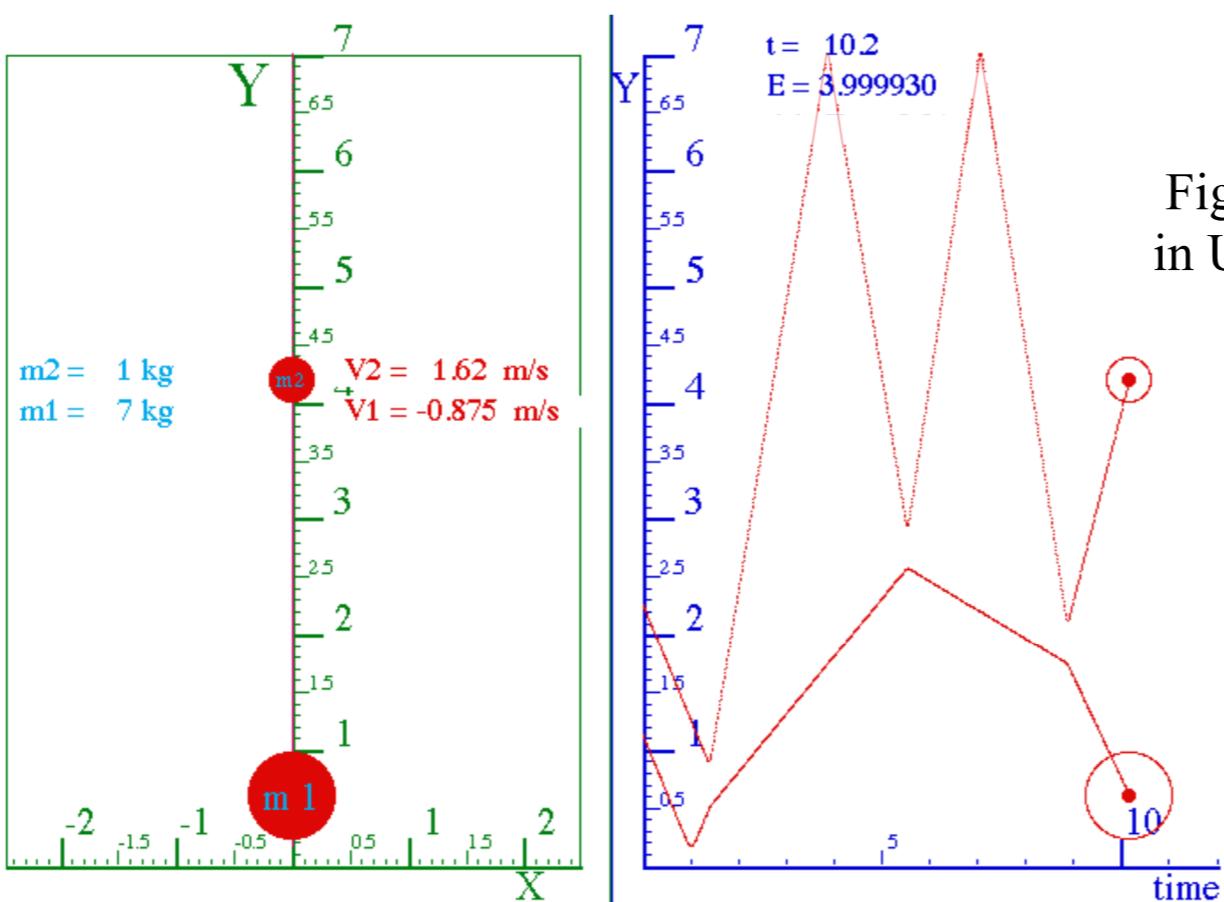
“Quantized”
Thales
Rectangles
provide
analytic geometry
of
fractal structure



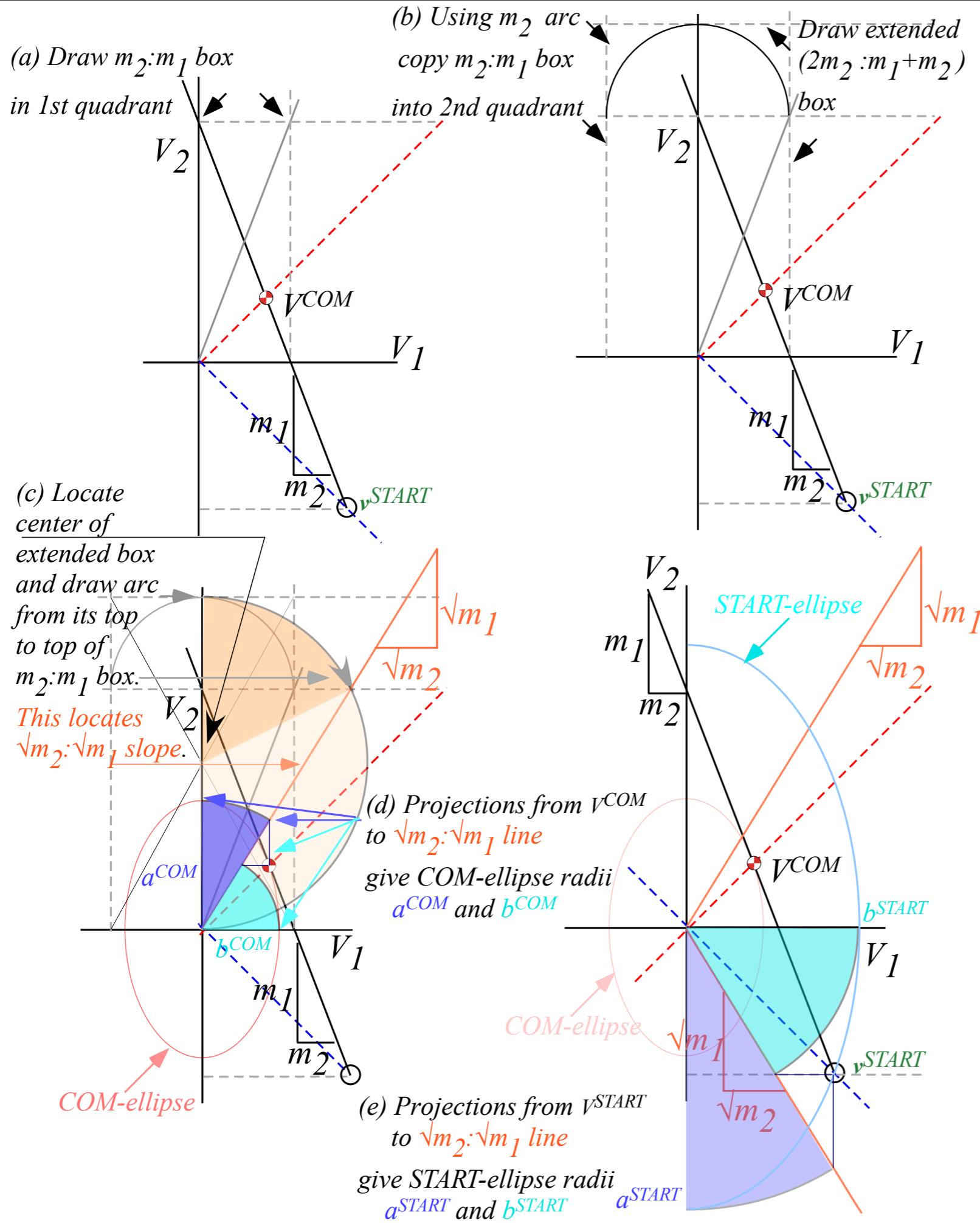
(Quantum computer simulation)
That makes an ∞ -ly deep “3D-Magic-Eye” picture



Geometric “Integration” (Converting Velocity data to Spacetime)



Unit 1
Fig. 8.4a-d



This is a construction of the energy ellipse in a Lagrangian (v_1, v_2) plot given the initial (v_1, v_2) .

The Estrangian (V_1, V_2) plot makes the (v_1, v_2) plot and this construction obsolete.

(Easier to just draw circle through initial (V_1, V_2) .)