Lecture 26 Mon. 11.25.2019

R=

KE/PE

eccentricity vector

Geometry and Symmetry of Coulomb Orbital Dynamics (Ch. 2-4 of Unit 5)

Rutherford scattering and hyperbolic orbit geometry Backward vs forward scattering angles and orbit construction example Parabolic "kite" and orbital envelope geometry $h_{=}$ Differential and total scattering cross-sections *impact parameter Eccentricity vector* $\boldsymbol{\varepsilon}$ *and* $(\boldsymbol{\varepsilon}, \boldsymbol{\lambda})$ *-geometry of orbital mechanics* Projection ε •r geometry of ε -vector and orbital radius r *Review and connection to usual orbital algebra (previous lecture) Projection* $\varepsilon \cdot \mathbf{p}$ *geometry of* ε *-vector and momentum* $\mathbf{p} = m\mathbf{v}$ General geometric orbit construction using ε -vector and (γ, R) -parameters Derivation of ε -construction by analytic geometry Coulomb orbit algebra of ε -vector and Kepler dynamics of momentum $\mathbf{p}=m\mathbf{v}$ *Example of complete* (**r**,**p**)*-geometry of elliptical orbit Connection formulas for* (γ, \mathbf{R}) *-parameters with* (a, b) *and* (ε, λ) elevation angle

Excerpts from Lect. 27

This Lecture's Reference Link Listing

<u>Web Resources - front page</u> <u>UAF Physics UTube channel</u> Quantum Theory for the Computer Age

Principles of Symmetry, Dynamics, and Spectroscopy

Classical Mechanics with a Bang!

Modern Physics and its Classical Foundations

2017 Group Theory for QM 2018 Adv CM 2018 AMOP 2019 Advanced Mechanics

Lecture #22-26

In reverse order

Coullt Web App Simulations: p19, p32, p72, p73, p92, R=-0.375, R=+0.5, Rutherford

OscillatorPE Web App: <u>IHO Scenario 2</u>, <u>Coulomb Scenario 3</u> RelaWavity Web App/Simulator/Calculator: <u>Elliptical - IHO orbits</u>

Jerklt Web App: 2-, 2+, Amp50Omega147-, Amp50Omega296, Amp50Omega602, Gap(1) MolVibes Web App: C3vN3 Wavelt Web App: Dim = <u>3 w/Wave Components</u>; Static Char Table: 6, 12, 12(b), 16, 36, 256 Quantum Carpet with N=20: Gaussian, Boxcar Quantum Revivals of Morse Oscillators and Farey-Ford Geometry - Li-Harter-CPL-2015 QTCA Unit 5 Ch14 2013 Lester, R. Ford, Am. Math. Monthly 45,586(1938) John Farey, Phil. Mag.(1816) Wolfram Harter, J. Mol. Spec. 210, 166-182 (2001) Harter, Li IMSS (2013) Li. Harter, Chem. Phys. Letters (2015) Advanced Atomic and Molecular Optical Physics 2018 Class #9, pages: 5, 61 **BoxIt Web Simulations** Pure A-Type A=4.9, B=0, C=0, & D=4.0 Pure B-Type: A=4.0, B=-0.2, C=0, & D=4.0 Pure C-Type A,D=4.055, B=0, C=0.1 Mixed AB-Type w/Cosine Mixed AB Type A=4.0, BU2=0.866..., CU2=0, & D=1.0 w/Stokes & Freq rats Mixed AB Type A=5.086 B=-0.27 C=0 D=2.024 w/Stokes plot Mixed ABC Type A=4.833 B=0.2403 C=0.4162 D=4.277 w/Stokes plot Recent mixed ABC Type A=0.325 B=0.375 C=0.825 D=0.05 w/Stokes plot

Select, exciting, and/or related Research

<u>This Indestructible NASA Camera Revealed Hidden Patterns on Jupiter</u> - seeker-yt-2019 <u>What did NASA's New Horizons discover around Pluto? - Astrum-yt-2018</u> <u>Synthetic Chiral Light for Efficient Control of Chiral Light–Matter Interaction - Ayuso-np-2019</u>

In development, but close to role out. More Advanced QM and classical references will *soon* be available through our: <u>References Page</u> Would be great to have our <u>Apache SOLR</u> Search & Index system up for a bigger Bang!)

Classical Mechanics with a Bang! 2018 Lectures <u>8</u>, <u>9</u>, <u>23 page 93</u> Text <u>Unit 6</u>, <u>page=27</u> <u>ColorU2 for the Web</u> - in development Group Theory for Quantum Mechanics - 2017 Lectures: <u>6</u>, <u>7</u>, <u>8</u>, and the <u>combined 9-10</u> Quantum Theory for the Computer Age <u>Unit 3 Ch.7-10</u>, <u>page=90</u> <u>Spectral Decomposition with Repeated Eigenvalues - 2017 GTQM - Lecture 5</u>

Web based 3D & XR (x∈{A,M,V}, R=Reality) <u>https://www.babylonjs.com/</u> Web based 3D graphics <u>WebGL API (Graphics Layer modeled after OpenGL)</u>

Recent In-House draft Articles:

Springer handbook on Molecular Symmetry and Dynamics - Ch_32 - Molecular Symmetry AMOP Ch 0 Space-Time Symmetry - 2019 Seminar at Rochester Institute of Optics, Auxiliary slides, June 19, 2018

Quantum_Computing - (Current) State of the Art - Reimer-www-2019 Geometric Algebra- A Guided Tour through Space and Time - Reimer-www-2019 Wildlife Monitoring Identification and Behavioral Study - Section 1 - Reimer-www-2019 Wildlife Monitoring Identification and Behavioral Study - Section 2 - Reimer-www-2019

Quantum Computing (QC) and Geometric Algebra (GA) references:

<u>*Quantum_Supremacy_Using_a_Programmable_Superconducting_Processor__Arute-n-2019</u></u> <u><i>Quantum Computing for Computer Scientists - Helwer-mr-yt-2018*, Slides Quantum Computing and Workforce, Curriculum, and App Devel - Roetteler-MS-2019</u></u>

Quantum_Computing - (Current) State of the Art - Reimer-www-2019 **Excerpts** (Page 44-47 in *Preliminary Draft*) for a GA take on the Complex Numbers Geometric Algebra- A Guided Tour through Space and Time - Reimer-www-2019 GA & QC references (Page 11-16 in Preliminary Draft)

Continued for 3 more pages

This Lecture's Reference Link Listing

<u>Web Resources - front page</u> <u>UAF Physics UTube channel</u> Quantum Theory for the Computer Age

Principles of Symmetry, Dynamics, and Spectroscopy

Classical Mechanics with a Bang!

Modern Physics and its Classical Foundations

Lectures #12 through #21

In reverse order

2017 Group Theory for QM 2018 Adv CM 2018 AMOP 2019 Advanced Mechanics

Wiki on Pafnuty Chebyshev Nobelprize.org 2005 Physics Award

BoxIt Web Simulations:

A-Type w/Cosine, A-Type w/Freq ratios, AB-Type w/Cosine, AB-Type 2:1 Freq ratio

OscillIt Web Simulations:

Default/Generic, Weakly Damped #18, Forced : Way below resonance,On resonance Way above resonance,Underdamped Complex Response Plot

Coullt Web Simulations:

<u>Stark-Coulomb : Bound-state motion in parabolic coordinates</u> <u>Molecular Ion : Bound-state motion in hyperbolic coordinates</u> <u>Synchrotron Motion, Synchrotron Motion #2</u> <u>Mechanical Analog to EM Motion (YouTube video)</u> iBall demo - Quasi-periodicity (YouTube video)

Trebuchet Web Simulations:

Default/Generic URL, Montezuma's Revenge, Seige of Kenilworth, "Flinger", Position Space (Course), Position Space (Fine) Wacky Waving Solid Metal Arm Flailing Chaos Pendulum - Scooba_Steeve-yt-2015 Triple Double-Pendulum - Cohen-yt-2008 Punkin Chunkin - TheArmchairCritic-2011 Jersey Team Claims Title in Punkin Chunkin - sussexcountyonline-1999 Shooting range for medieval siege weapons. Anybody knows? - twcenter.net/forums The Trebuchet - Chevedden-SciAm-1995 NOVA Builds a Trebuchet

Recent Articles of Interest:

<u>A_Semi-Classical_Approach_to_the_Calculation_of_Highly_Excited_Rotational_Energies for</u> ...
 <u>Asymmetric-Top_Molecules_-_Schmiedt-pccp-2017</u>
 Tunable and broadband coherent perfect absorption by ultrathin blk phos metasurfaces - Guo-josab-2019
 Vortex Detection in Vector Fields Using Geometric Algebra - Pollock-aaca-2013.pdf

Pirelli Relativity Challenge (Introduction level) - Visualizing Waves:

Using Earth as a clock, Tesla's AC Phasors, Phasors using complex numbers. CM wBang Unit 1 - Chapter 10, pdf_page=135 Calculus of exponentials, logarithms, and complex fields, RelaWavity Web Simulation - Unit Circle and Hyperbola (Mixed labeling) Smith Chart, Invented by Phillip H. Smith (1905-1987)

Select, exciting, and related Research

Clifford Algebra And The Projective Model Of Homogeneous Metric Spaces -Foundations - Sokolov-x-2013 Geometric Algebra 3 - Complex Numbers - MacDonald-yt-2015 Biquaternion - Complexified Quaternion - Roots of -1 - Sangwine-x-2015 An Introduction to Clifford Algebras and Spinors - Vaz-Rocha-op-2016 Unified View on Complex Numbers and Quaternions- Bongardt-wcmms-2015 Complex Functions and the Cauchy-Riemann Equations - complex2 - Friedman-columbia-2019 An sp-hybridized Molecular Carbon Allotrope- cyclo-18-carbon - Kaiser-s-2019 An Atomic-Scale View of Cyclocarbon Synthesis - Maier-s-2019 Discovery Of Topological Weyl Fermion Lines And Drumhead Surface States in a Room Temperature Magnet - Belopolski-s-2019 "Weyl"ing away Time-reversal Symmetry - Neto-s-2019 Non-Abelian Band Topology in Noninteracting Metals - Wu-s-2019 What Industry Can Teach Academia - Mao-s-2019 RoVib- quantum state resolution of the C60 fullerene - Changala-Ye-s-2019 (Alt) A Degenerate Fermi Gas of Polar molecules - DeMarco-s-2019

An assist from *Physics Girl* (YouTube Channel):

How to Make VORTEX RINGS in a Pool Crazy pool vortex - pg-yt-2014 Fun with Vortex Rings in the Pool - pg-yt-2014

Running Reference Link Listing

Lectures #11 through #7

In reverse order

Eric J Heller Gallery:

Main portal, Consonance and Dissonance II, Bessel 21, Chladni

The Semiclassical Way to Molecular Spectroscopy - Heller-acs-1981 Quantum_dynamical_tunneling_in_bound_states_-_Davis-Hellerjcp-1981

Pendulum Web Simulation Cycloidulum Web Simulation

Links to previous lecture: <u>Page=74</u>, <u>Page=75</u>, <u>Page=79</u>

Pendulum Web Sim

Cycloidulum Web Sim

JerkIt Web Simulations: Basic/Generic: Inverted, FVPlot

CMwithBang Lecture 8, page=20

WWW.sciencenewsforstudents.org: Cassini - Saturnian polar vortex

"RelaWavity" Web Simulations:
<u>2-CW laser wave, Lagrangian vs Hamiltonian,</u> <u>Physical Terms Lagrangian L(u) vs Hamiltonian H(p)</u>
<u>Coullt Web Simulation of the Volcanoes of Io</u>
BohrIt Multi-Panel Plot:
Relativistically shifted Time-Space plots of 2 CW light waves

BoxIt Web Simulations:

<u>Generic/Default</u> <u>Most Basic A-Type</u> <u>Basic A-Type w/reference lines</u> <u>Basic A-Type A-Type with Potential energy</u> <u>A-Type with Potential energy and Stokes Plot</u> <u>A-Type w/3 time rates of change</u> <u>A-Type w/3 time rates of change with Stokes Plot</u> <u>B-Type (A=1.0, B=-0.05, C=0.0, D=1.0)</u>

RelaWavity Web Elliptical Motion Simulations:

Orbits with b/a=0.125 Orbits with b/a=0.5 Orbits with b/a=0.7 Exegesis with b/a=0.125 Exegesis with b/a=0.5 Exegesis with b/a=0.7 Contact Ellipsometry

Coullt Web Simulations: Basic/Generic

Exploding Starlet Volcanoes of Io (Color Quantized)

JerkIt Web Simulations:

<u>Basic/Generic</u> Catcher in the Eye - IHO with Linear Hooke perturbation - Force-potential-Velocity Plot

OscillatorPE Web Simulation:

Coulomb-Newton-Inverse_Square, Hooke-Isotropic Harmonic, Pendulum-Circular Constraint

AMOP Ch 0 Space-Time Symmetry - 2019 Seminar at Rochester Institute of Optics, Aux. slides-2018

NASA Astronomy Picture of the Day -<u>Io: The Prometheus Plume (Just Image)</u> <u>NASA Galileo - Io's Alien Volcanoes</u> <u>New Horizons - Volcanic Eruption Plume on Jupiter's moon IO</u> <u>NASA Galileo - A Hawaiian-Style Volcano on Io</u>

<u>Pirelli Site: Phasors animimation</u> <u>CMwithBang Lecture #6, page=70 (9.10.18)</u>

Select, exciting, and related Research & Articles of Interest:

Burning a hole in reality—design for a new laser may be powerful enough to pierce space-time - Sumner-KOS-2019 Trampoline mirror may push laser pulse through fabric of the Universe - Lee-ArsTechnica-2019 Achieving Extreme Light Intensities using Optically Curved Relativistic Plasma Mirrors - Vincenti-prl-2019 <u>A Soft Matter Computer for Soft Robots - Garrad-sr-2019</u> <u>Correlated Insulator Behaviour at Half-Filling in Magic-Angle Graphene Superlattices - cao-n-2018</u> <u>Sorting ultracold atoms in a three-dimensional optical lattice in a</u> realization of Maxwell's Demon - Kumar-n-2018 Synthetic three-dimensional atomic structures assembled atom by atom - Barredo-n-2018 Older ones: Wave-particle duality of C60 molecules - Arndt-Itn-1999 Optical Vortex Knots - One Photon _ At A Time - Tempone-Wiltshire-Sr-2018 Baryon Deceleration by Strong Chromofields in Ultrarelativistic ,

<u>Baryon_Deceleration_by_Strong_Chromofields_in_Ottrarelativistic_</u>, <u>Nuclear_Collisions - Mishustin-PhysRevC-2007</u>, <u>APS Link & Abstract</u> Hadronic Molecules - Guo-x-2017

Hidden-charm pentaquark and tetraquark states - Chen-pr-2016

Running Reference Link Listing

Lectures #6 through #1

In reverse order

RelaWavity Web Simulation: Contact EllipsometryBoxIt Web Simulation: Elliptical Motion (A-Type)CMwBang Course: Site Title PagePirelli Relativity Challenge: Describing Wave Motion With Complex PhasorsUAF Physics UTube channel

Velocity Amplification in Collision Experiments Involving Superballs - Harter, 1971 <u>MIT OpenCourseWare: High School/Physics/Impulse and Momentum</u> <u>Hubble Site: Supernova - SN 1987A</u>

BounceItIt Web Animation - Scenarios:

49:1 y vs t, 49:1 V2 vs V1, 1:500:1 - 1D Gas Model w/ faux restorative force (Cool), 1:500:1 - 1D Gas (Warm), 1:500:1 - 1D Gas Model (Cool, Zoomed in),
Farey Sequence - Wolfram
Fractions - Ford-AMM-1938
Monstermash BounceItIt Animations: 1000:1 - V2 vs V1, 1000:1 with t vs x - Minkowski Plot
Quantum Revivals of Morse Oscillators and Farey-Ford Geometry - Li-Harter-2013
Quantum Revivals of Morse Oscillators and Farey-Ford Geometry - Li-Harter-cpl-2015
Quant. Revivals of Morse Oscillators and Farey-Ford Geom. - Harter-Li-CPL-2015 (Publ.)
Velocity Amplification in Collision Experiments Involving Superballs-Harter-1971
WaveIt Web Animation - Scenarios: Quantum Carpet, Quantum Carpet wMBars, Quantum Carpet BCar, Quantum Carpet BCar_wMBars
Wave Node Dynamics and Revival Symmetry in Quantum Rotors - Harter-JMS-2001 (Publ.)

<u>AJP article on superball dynamics</u> <u>AAPT Summer Reading List</u> <u>Scitation.org - AIP publications</u> <u>HarterSoft Youtube Channel</u>

BounceIt Web Animation - Scenarios:

Generic Scenario: <u>2-Balls dropped no Gravity (7:1) - V vs V Plot (Power=4)</u> 1-Ball dropped w/Gravity=0.5 w/Potential Plot: <u>Power=1, Power=4</u> <u>7:1 - V vs V Plot: Power=1</u> <u>3-Ball Stack (10:3:1) w/Newton plot (y vs t) - Power=4</u> <u>3-Ball Stack (10:3:1) w/Newton plot (y vs t) - Power=1 w/Gaps</u> <u>4-Ball Stack (27:9:3:1) w/Newton plot (y vs t) - Power=4</u> <u>4-Newton's Balls (1:1:1:1) w/Newtonian plot (y vs t) - Power=4</u> <u>5-Ball Totally Inelastic (1:1:1:1:1) w/Gaps: Newtonian plot (t vs x), V6 vs V5 plot</u> <u>5-Ball Totally Inelastic Pile-up w/ 5-Stationary-Balls - Minkowski plot (t vs x1) w/Gaps</u>

BounceIt Dual plots

 $m_{1}:m_{2} = 3:1$ $v_{2} vs v_{1} and V_{2} vs V_{1}, (v_{1}, v_{2}) = (1, 0.1), (v_{1}, v_{2}) = (1, 0)$ $y_{2} vs y_{1} plots: (v_{1}, v_{2}) = (1, 0.1), (v_{1}, v_{2}) = (1, 0), (v_{1}, v_{2}) = (1, -1)$ Estrangian plot $V_{2} vs V_{1}: (v_{1}, v_{2}) = (0, 1), (v_{1}, v_{2}) = (1, -1)$ $m_{1}:m_{2} = 4:1$ $v_{2} vs v_{1}, v_{2} vs y_{1}$ $m_{1}:m_{2} = 100:1, (v_{1}, v_{2}) = (1, 0): V_{2} vs V_{1} Estrangian plot, y_{2} vs y_{1} plot$ With g=0 and 70:10 mass ratio With non zero g, velocity dependent damping and mass ratio of 70:35 $M_{1}=49, M_{2}=1 with Newtonian time plot$ $M_{1}=49, M_{2}=1 with V_{2} vs V_{1} plot$ Example with friction Low force constant with drag displaying a Pass-thru, Fall-Thru, Bounce-Off $m_{1}:m_{2}= 3:1 and (v_{1}, v_{2}) = (1, 0) Comparison with Estrangian$

X2 paper: <u>Velocity Amplification in Collision Experiments Involving Superballs - Harter, et. al. 1971 (pdf)</u> Car Collision Web Simulator: <u>https://modphys.hosted.uark.edu/markup/CMMotionWeb.html</u> Superball Collision Web Simulator: <u>https://modphys.hosted.uark.edu/markup/BounceItWeb.html</u>; with Scenarios: <u>1007</u> <u>BounceIt web simulation with g=0 and 70:10 mass ratio</u> <u>With non zero g, velocity dependent damping and mass ratio of 70:35</u> Elastic Collision Dual Panel Space vs Space: <u>Space vs Time (Newton)</u>, <u>Time vs. Space(Minkowski)</u> Inelastic Collision Dual Panel Space vs Space: <u>Space vs Time (Newton)</u>, <u>Time vs. Space(Minkowski)</u> Matrix Collision Simulator: M_1 =49, M_2 =1 V_2 vs V_1 plot <<Under Construction>> Rutherford scattering and hyperbolic orbit geometry Backward vs forward scattering angles and orbit construction example Parabolic "kite" and orbital envelope geometry Differential and total scattering cross-sections

Eccentricity vector $\boldsymbol{\varepsilon}$ and (ε, λ) -geometry of orbital mechanics Projection $\boldsymbol{\varepsilon} \cdot \mathbf{r}$ geometry of $\boldsymbol{\varepsilon}$ -vector and orbital radius \mathbf{r} Review and connection to usual orbital algebra (previous lecture) Projection $\boldsymbol{\varepsilon} \cdot \mathbf{p}$ geometry of $\boldsymbol{\varepsilon}$ -vector and momentum $\mathbf{p} = m\mathbf{v}$

General geometric orbit construction using ε -vector and (γ, R) -parameters Derivation of ε -construction by analytic geometry

Coulomb orbit algebra of ε -vector and Kepler dynamics of momentum $\mathbf{p}=m\mathbf{v}$ Example of complete (\mathbf{r},\mathbf{p}) -geometry of elliptical orbit

Connection formulas for (γ, \mathbf{R}) *-parameters with* (a, b) *and* (ε, λ)





Rutherford scattering of α^{+2} particles from Au^{+79} nucleus at O Assume "Dead-On" closest approach 2a. (E=k/2a) $a\sim 10^{-11}m >> 7.3 \cdot 10^{-15}m$

Pick an "impact parameter" line y = b. Draw circle of radius a around center point C=(-a,b) tangent to y-axis. Draw "focus-locus" line OCF.



Rutherford scattering of α^{+2} particles from Au^{+79} nucleus at O Assume "Dead-On" closest approach 2a. (E=k/2a) $a\sim 10^{-11}m >> 7.3 \cdot 10^{-15}m$

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Copy angle \angle BCF (equal to $\Phi/2$) to make angle \angle FCB' (also equal to $\Phi/2$) Resulting line CB' is outgoing asymptote at scattering angle Θ .



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Locate secondary focus O' by drawing circle around point C of diameter CO thru point O. Diameter O'CO is $2a\varepsilon$. Hyperbolic orbit points P now found using constant 2a=PO-PO'



Rutherford scattering of α^{+2} particles from Au^{+79} nucleus at O Assume "Dead-On" closest approach 2a. (E=k/2a) $a\sim 10^{-11}m >> 7.3 \cdot 10^{-15}m$

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Rutherford scattering and hyperbolic orbit geometry

Backward vs forward scattering angles and orbit construction example
 Parabolic "kite" and orbital envelope geometry
 Differential and total scattering cross-sections

Eccentricity vector ε and (ε, λ) -geometry of orbital mechanics Projection ε •**r** geometry of ε -vector and orbital radius **r** Review and connection to usual orbital algebra (previous lecture) Projection ε •**p** geometry of ε -vector and momentum **p**=m**v**

General geometric orbit construction using ε -vector and (γ, R) -parameters Derivation of ε -construction by analytic geometry

Coulomb orbit algebra of ε -vector and Kepler dynamics of momentum $\mathbf{p}=m\mathbf{v}$ Example of complete (\mathbf{r},\mathbf{p}) -geometry of elliptical orbit

Connection formulas for (γ, \mathbf{R}) *-parameters with* (a, b) *and* (ε, λ)

















Rutherford scattering and hyperbolic orbit geometry

Backward vs forward scattering angles and orbit construction example



- *Parabolic "kite" and orbital envelope geometry* Differential and total scattering cross-sections

Eccentricity vector $\boldsymbol{\varepsilon}$ *and* ($\boldsymbol{\varepsilon}, \boldsymbol{\lambda}$)*-geometry of orbital mechanics* Projection ε •r geometry of ε -vector and orbital radius r *Review and connection to usual orbital algebra (previous lecture) Projection* $\varepsilon \cdot \mathbf{p}$ *geometry of* ε *-vector and momentum* $\mathbf{p} = m\mathbf{v}$

General geometric orbit construction using ε -vector and (γ, \mathbf{R}) -parameters Derivation of ε -construction by analytic geometry

Coulomb orbit algebra of ε -vector and Kepler dynamics of momentum $\mathbf{p}=m\mathbf{v}$ *Example of complete* (**r**,**p**)*-geometry of elliptical orbit*

Connection formulas for (γ, \mathbf{R}) *-parameters with* (a, b) *and* (ε, λ)

Rutherford scattering geometry





Chapter 1 Orbit Families and Action

Families of particle orbits are drawn in a varying color which represents the classical action or Hamiltoon's characteristic function SH dq.(Sometimes SH is called 'reduced action'.) The color is chosen by first calculating c = SH modulo h-bar (You can change Planck's consta its default value h/2π = 1.0) The chromatic value c assigns the hue by its position on the color wheel (e.g.; c=0 is red, c=0.2 is a yellow, c=0. green, etc.).

Chapter 2 Rutherford Scattering

A parallel beam of iso-energetic alpha particles undergo Rutherford scattering from a coulomb field of a nucleus as calculated in the It is also the ideal pattern of paths followed by intergalactic hydrogen in perturbed by the solar wind.

Chapter 3 Coulomb Field (H atom)

Orbits in an attractive Coulomb field are calculated here. You may select the initial position (x(0),y(0)) by moving the mouse to a des launch point, and then select the initial momentum (px(0), py(0)) by pressing the mouse button and dragging.

Chapter 4 Molecular Ion Orbits

Orbits around two fixed nuclei are calculated here. A set of elliptic coordinates are drawn in the background. After running a few tra you may notice that their caustics conform to one or two of the elliptic coordinate lines.

H = ∫p ant from 5 is a	Volcanoes of Io (Paths=180, No color quant.) Parabolic Fountain (Uniform) Space Bomb (Coulomb) Exploding Starlet (IHO) Synchrotron Motion (Crossed E & B fields)
se demos.	Rutherford scattering 2-Electron Orbits
ired	Atomic Orbits
ajectories	Molecular Ion Orbits
	Oscillator Scattering (2-Particle Orbits) (2-Particle Collision)











tangent slope=-5/2 slo

slope=1/2





Rutherford scattering and hyperbolic orbit geometry Backward vs forward scattering angles and orbit construction example Parabolic "kite" and orbital envelope geometry Differential and total scattering cross-sections

Eccentricity vector $\boldsymbol{\varepsilon}$ and $(\boldsymbol{\varepsilon}, \boldsymbol{\lambda})$ -geometry of orbital mechanics Projection $\boldsymbol{\varepsilon} \cdot \mathbf{r}$ geometry of $\boldsymbol{\varepsilon}$ -vector and orbital radius \mathbf{r} Review and connection to usual orbital algebra (previous lecture) Projection $\boldsymbol{\varepsilon} \cdot \mathbf{p}$ geometry of $\boldsymbol{\varepsilon}$ -vector and momentum $\mathbf{p} = m\mathbf{v}$

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Connection formulas for (γ, \mathbf{R}) *-parameters with* (a, b) *and* (ε, λ)



Also: Approximate model of deep-space H-atom scattering from solar wind as our Sun travels around galaxy. Lyman- α shock wave found just inside Mars orbital radius 2a~1.2Au.

Fig. 5.3.2 Family of iso-energetic Rutherford scattering orbits with varying impact parameter.

Incremental window $d\sigma = b \cdot db$ normal to beam axis at $x = -\infty$ scatters to area $dA = R^2 \sin \Theta d\Theta d\varphi = R^2 d\Omega$ onto a sphere at $R = +\infty$ where is called the *incremental solid angle* $d\Omega = \sin \Theta d\Theta d\varphi$



Also: Approximate model of deep-space H-atom scattering from solar wind as our Sun travels around galaxy. Lyman- α shock wave found just inside Mars orbital radius 2a~1.2Au.

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Ratio $\frac{d\sigma}{d\Omega} = \frac{b \, db \, d\phi}{\sin \Theta d\Theta d\phi} = \frac{b}{\sin \Theta} \frac{db}{d\Theta}$ is called the *differential scattering cross-section (DSC)*



H-atom scattering from solar wind as our Sun travels around galaxy. *Lyman*- α *shock wave found just inside Mars* orbital radius 2a~1.2Au.

Incremental window $d\sigma = b \cdot db$ normal to beam axis at $x = -\infty$ scatters to area $dA = R^2 \sin \Theta d\Theta d\phi = R^2 d\Omega$



Also: Approximate model of deep-space *H-atom scattering* from solar wind as our Sun travels around galaxy. *Lyman*- α *shock wave found just inside Mars* orbital radius 2a~1.2Au.

Fig. 5.3.2 Family of iso-energetic Rutherford scattering orbits with varying impact parameter.

Incremental window $d\sigma = b \cdot db$ normal to beam axis at $x = -\infty$ scatters to area $dA = R^2 \sin \Theta d\Theta d\phi = R^2 d\Omega$ onto a sphere at $R = +\infty$ where is called the *incremental solid angle* $d\Omega = \sin \Theta d\Theta d\phi$

Ratio
$$\frac{d\sigma}{d\Omega} = \frac{b \, db \, d\phi}{\sin \Theta d\Theta d\phi} = \frac{b}{\sin \Theta} \frac{db}{d\Theta}$$
 is called the *differential scattering cross-section (DSC)*
Geometry: $b = a \cot \frac{\Theta}{2} = \frac{k}{2E} \cot \frac{\Theta}{2}$
with: $\frac{db}{d\Theta} = \frac{-a}{2} \csc^2 \frac{\Theta}{2} = \frac{-a}{2\sin^2 \frac{\Theta}{2}}$
(Never forget!: $a = \frac{-k}{2E}$ or: $E = \frac{-k}{2a}$)



model of deep-space *H-atom scattering* from solar wind as our Sun travels around galaxy. *Lyman*- α *shock wave found just inside Mars* orbital radius 2a~1.2Au.

Incremental window $d\sigma = b \cdot db$ normal to beam axis at $x = -\infty$ scatters to area $dA = R^2 \sin \Theta d\Theta d\phi = R^2 d\Omega$



model of deep-space *H-atom scattering* from solar wind as our Sun travels around galaxy. *Lyman*- α *shock wave* found just inside Mars orbital radius 2a~1.2Au.

Fig. 5.3.2 Family of iso-energetic Rutherford scattering orbits with varying impact parameter.

Incremental window $d\sigma = b \cdot db$ normal to beam axis at $x = -\infty$ scatters to area $dA = R^2 \sin \Theta d\Theta d\phi = R^2 d\Omega$


Eccentricity vector $\boldsymbol{\varepsilon}$ and (ε, λ) -geometry of orbital mechanics Projection $\boldsymbol{\varepsilon} \cdot \mathbf{r}$ geometry of $\boldsymbol{\varepsilon}$ -vector and orbital radius \mathbf{r} Review and connection to usual orbital algebra (previous lecture) Projection $\boldsymbol{\varepsilon} \cdot \mathbf{p}$ geometry of $\boldsymbol{\varepsilon}$ -vector and momentum $\mathbf{p}=m\mathbf{v}$

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Eccentricity vector $\boldsymbol{\varepsilon}$ *and* ($\boldsymbol{\varepsilon}, \boldsymbol{\lambda}$) *geometry of orbital mechanics*

Isotropic field V=V(r) guarantees conservation *angular momentum vector* **L**

 $\mathbf{L} = \mathbf{r} \times \mathbf{p} = m \, \mathbf{r} \times \dot{\mathbf{r}}$

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Coulomb V = -k/r also conserves *eccentricity vector* ε

 $\mathbf{\hat{\varepsilon}} = \hat{\mathbf{r}} - \frac{\mathbf{p} \times \mathbf{L}}{km} = \frac{\mathbf{r}}{r} - \frac{\mathbf{p} \times (\mathbf{r} \times \mathbf{p})}{km}$

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(...for sake of comparison...) IHO $V = (k/2)r^2$ also conserves Stokes vector **S** $S_A = \frac{1}{2}(x_1^2 + p_1^2 - x_2^2 - p_2^2)$ $S_B = x_1 p_1 + x_2 p_2$ $S_C = x_1 p_2 - x_2 p_1$

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Eccentricity vector ε and (ε, λ) -geometry of orbital mechanics Projection ε -**r** geometry of ε -vector and orbital radius **r** Review and connection to usual orbital algebra (previous lecture) Projection ε -**p** geometry of ε -vector and momentum $\mathbf{p}=m\mathbf{v}$

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Consider dot product of ε with a radial vector **r**:

 $\mathbf{\varepsilon} \bullet \mathbf{r} = \frac{\mathbf{r} \bullet \mathbf{r}}{r} - \frac{\mathbf{r} \bullet \mathbf{p} \times \mathbf{L}}{km} = r - \frac{\mathbf{r} \times \mathbf{p} \bullet \mathbf{L}}{km} = r - \frac{\mathbf{L} \bullet \mathbf{L}}{km}$

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 $\mathbf{A} = km \cdot \varepsilon$ is known as the *Laplace-Hamilton-Gibbs-Runge-Lenz vector*.

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... or of ε with momentum vector **p**:

$$\mathbf{\varepsilon} \bullet \mathbf{p} = \frac{\mathbf{p} \bullet \mathbf{r}}{r} - \frac{\mathbf{p} \bullet \mathbf{p} \times \mathbf{L}}{km} = \mathbf{p} \bullet \hat{\mathbf{r}} = p_r$$

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Eccentricity vector ε and (ε,λ)-geometry of orbital mechanics
Projection ε•r geometry of ε-vector and orbital radius r
Review and connection to usual orbital algebra (previous lecture)
Projection ε•p geometry of ε-vector and momentum p=mv

General geometric orbit construction using ε -vector and (γ, R) -parameters Derivation of ε -construction by analytic geometry

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NOTE: Lengths of vectors **p** and -**p** are not drawn so they correctly show that momentum $\mathbf{p}=m\mathbf{v}$ grows as radial distance $r=|\mathbf{r}|$ falls. (To be shown on p. <u>95-104</u>)



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Dot product of ε with momentum vector **p**:

 $\boldsymbol{\varepsilon} \bullet \mathbf{p} = \frac{\mathbf{p} \bullet \mathbf{r}}{r} - \frac{\mathbf{p} \bullet \mathbf{p} \times \mathbf{r}}{km}$ $= \mathbf{p} \bullet \hat{\mathbf{r}} = p_r = \boldsymbol{\varepsilon} p_x$

This says: "Projection p_r of \mathbf{p} onto radial \mathbf{r} or \mathbf{r}' lines equals eccentricity $\boldsymbol{\varepsilon}$ times projection p_x of \mathbf{p} onto orbit major axis : $(\hat{\mathbf{x}} = \hat{\boldsymbol{\varepsilon}})$ "

Focal geometry demands: "Momentum \mathbf{p} must bisect angle $\measuredangle_{\mathbf{r}}^{\mathbf{r}}$, between radial \mathbf{r} or \mathbf{r}' lines."

> *Hyperbola has eccentricity* $\varepsilon > 1$ (*Here* : $\varepsilon = 5/4 = 1.25$)

 $p_r = \epsilon p_x$

 $p_{\rm x}$

b

ae

a

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Next several pages give step-by-step constructions of ε -vector and Coulomb orbit and trajectory physics

Fig. 5.4.2 Construction of eccentricity vector ε *and*

General geometric orbit construction using ε -vector and (γ, \mathbf{R}) -parameters

Pick launch point P (radius vector r) and elevation angle γ from radius (momentum initial p direction) inital momentum elevation angle γ r

General geometric orbit construction using ε -vector and (γ, \mathbf{R}) -parameters



General geometric orbit construction using ε -vector and (γ, \mathbf{R}) -parameters *Copy double angle 2* γ *(* \angle **FPQ***) onto* \angle **PFT** Copy F-center circle around launch point P Pick launch point P *Extend* ∠PFT *chord* PT *to make R***-ratio scale line** *Copy elevation angle* γ (\angle FPP') *onto* \angle P'PQ (radius vector **r**) and elevation angle γ from radius Extend resulting line QPQ' to make focus locus (momentum initial **p** direction) inital momentum р *wpied* elevation angle γ inital momentum D p elevation angle γ *copied* F **P**' inital momentum p elevation angle γ FOCUS Reason for focus locu Line **r** from 1st focus **F**/"reflects r off line **p** (or **P'P**) toward 2nd focus **F** somewhere so incident-angle γ equals reflected-angle γ









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Play embedded movie with controls above.

or follow this link to try your hand at ε-construction using the CoulIt Web App

Just click and drag in main window to set new initial conditions. The Lenz vector will display as part of an overlay.

Rutherford scattering and hyperbolic orbit geometry Backward vs forward scattering angles and orbit construction example Parabolic "kite" and orbital envelope geometry Differential and total scattering cross-sections

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Connection formulas for (γ, \mathbf{R}) *-parameters with* (a, b) *and* (ε, λ)

Finding time derivatives of orbital coordinates r, ϕ , x, y, and eventually velocity **v** or momentum **p**=m**v**

Polar angle
$$\phi$$
 using: $L = mr^2 \frac{d\phi}{dt} = mr^2 \dot{\phi}$

Radius r: (p. <u>48</u> to p.<u>53</u>) $r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi}$

Finding time derivatives of orbital coordinates r, ϕ , x, y, and eventually velocity **v** or momentum **p**=m**v**

Radius r:

 $r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi}$

Polar angle
$$\phi$$
 using: $L = mr^2 \frac{d\phi}{dt} = mr^2 \dot{\phi}$
 $\dot{\phi} = \frac{L}{mr^2} = \frac{L}{m} \frac{1}{r^2}$

Radius r:

$$r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi}$$

$$polar angle \phi using: L = mr^2 \frac{d\phi}{dt} = mr^2 \dot{\phi}$$

$$\dot{\phi} = \frac{L}{mr^2} = \frac{L}{m} \frac{1}{r^2} = \frac{L}{m} \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos \phi)^2$$

using:
$$\frac{1}{r^2} = \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos\phi)^2$$

Radius r:

$$r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi}$$

$$\dot{r} = \frac{dr}{dt} = \frac{L^2}{km} \frac{-\frac{d}{dt}(-\varepsilon \cos \phi)}{(1 - \varepsilon \cos \phi)^2}$$

$$\dot{r} = \frac{dr}{dt} = \frac{L^2}{km} \frac{-\frac{d}{dt}(-\varepsilon \cos \phi)}{(1 - \varepsilon \cos \phi)^2}$$

$$using: \frac{1}{r^2} = \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos \phi)^2$$

Radius r:Polar angle
$$\phi$$
 using: $L = mr^2 \frac{d\phi}{dt} = mr^2 \dot{\phi}$ $r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2/km}{1 - \varepsilon \cos \phi}$ $\dot{\phi} = \frac{L}{mr^2} = \frac{L}{m} \frac{1}{r^2} = \frac{L}{m} \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos \phi)^2$ $\dot{r} = \frac{dr}{dt} = \frac{L^2}{km} \frac{-\frac{d}{dt}(-\varepsilon \cos \phi)}{(1 - \varepsilon \cos \phi)^2}$ $\dot{\phi} = \frac{L}{mr} = \frac{L}{m} \frac{1}{r}$

$$\begin{aligned} \text{Radius } r: & \text{Polar angle } \phi \text{ using: } L = mr^2 \frac{d\phi}{dt} = mr^2 \dot{\phi} \\ r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi} & \dot{\phi} = \frac{L}{mr^2} = \frac{L}{m} \frac{1}{r^2} = \frac{L}{m} \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos \phi)^2 \\ \dot{r} = \frac{dr}{dt} = \frac{L^2}{km} \frac{-\frac{d}{dt} (-\varepsilon \cos \phi)}{(1 - \varepsilon \cos \phi)^2} & r\dot{\phi} = \frac{L}{mr} = \frac{L}{m} \frac{1}{r} = \frac{L}{m} \left(\frac{km}{L^2}\right) (1 - \varepsilon \cos \phi) = \frac{k}{L} (1 - \varepsilon \cos \phi) \\ using: \frac{1}{r} = \left(\frac{km}{L^2}\right) (1 - \varepsilon \cos \phi) \end{aligned}$$

$$\begin{aligned} Radius r: & Polar angle \phi using: \ L = mr^2 \frac{d\phi}{dt} = mr^2 \dot{\phi} \\ r = \frac{\lambda}{1 - \varepsilon \cos\phi} = \frac{L^2/km}{1 - \varepsilon \cos\phi} & \dot{\phi} = \frac{L}{mr^2} = \frac{L}{m} \frac{1}{r^2} = \frac{L}{m} \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos\phi)^2 \\ \dot{r} = \frac{dr}{dt} = \frac{L^2}{km} \frac{-\frac{d}{dt}(-\varepsilon \cos\phi)}{(1 - \varepsilon \cos\phi)^2} & r\dot{\phi} = \frac{L}{mr} = \frac{L}{m} \frac{1}{r} = \frac{L}{m} \left(\frac{km}{L^2}\right) (1 - \varepsilon \cos\phi) = \frac{k}{L} (1 - \varepsilon \cos\phi) \\ r\dot{\phi} = \frac{L}{mr} = \frac{L}{m} \frac{1}{r} = \frac{L}{m} \left(\frac{km}{L^2}\right) (1 - \varepsilon \cos\phi) = \frac{k}{L} (1 - \varepsilon \cos\phi)^2 \\ using: \frac{1}{r^2} = \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos\phi)^2 \end{aligned}$$

$$\begin{aligned} \text{Radius } r: & \text{Polar angle } \phi \text{ using: } L = mr^2 \frac{d\phi}{dt} = mr^2 \phi \\ r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2/km}{1 - \varepsilon \cos \phi} & \phi = \frac{L}{mr^2} = \frac{L}{m} \frac{1}{r^2} = \frac{L}{m} \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos \phi)^2 \\ \dot{r} = \frac{dr}{dt} = \frac{L^2}{km} \frac{-\frac{d}{dt}(-\varepsilon \cos \phi)}{(1 - \varepsilon \cos \phi)^2} & r\phi = \frac{L}{mr} = \frac{L}{m} \frac{1}{r} = \frac{L}{m} \left(\frac{km}{L^2}\right) (1 - \varepsilon \cos \phi) = \frac{k}{L} (1 - \varepsilon \cos \phi)^2 \\ \dot{r} = \frac{L^2}{km} \frac{-\varepsilon \sin \phi \phi}{(1 - \varepsilon \cos \phi)^2} & using: \frac{1}{r^2} = \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos \phi)^2 \\ \dot{r} = -\frac{L^2}{km} \left(\frac{km}{L^2}\right)^2 r^2 \dot{\phi} \varepsilon \sin \phi & using: \frac{1}{(1 - \varepsilon \cos \phi)^2} = \left(\frac{km}{L^2}\right)^2 r^2 \end{aligned}$$

$$\begin{aligned} \text{Radius } r: & \text{Polar angle } \phi \text{ using: } L = mr^2 \frac{d\phi}{dt} = mr^2 \dot{\phi} \\ r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi} & \dot{\phi} = \frac{L}{mr^2} = \frac{L}{m} \frac{1}{r^2} = \frac{L}{m} \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos \phi)^2 \\ \dot{r} = \frac{dr}{dt} = \frac{L^2}{\sqrt{m}} \frac{-\frac{d}{dt} (-\varepsilon \cos \phi)}{(1 - \varepsilon \cos \phi)^2} & r\dot{\phi} = \frac{L}{mr} = \frac{L}{mr} = \frac{L}{m} \frac{1}{r} = \frac{L}{m} \left(\frac{km}{L^2}\right) (1 - \varepsilon \cos \phi) = \frac{k}{L} (1 - \varepsilon \cos \phi)^2 \\ \dot{r} = -\frac{L^2}{km} \left(\frac{km}{L^2}\right)^2 r^2 \dot{\phi} \varepsilon \sin \phi & using: \frac{1}{(1 - \varepsilon \cos \phi)^2} = \left(\frac{km}{L^2}\right)^2 r^2 \\ \dot{r} = -\frac{k}{L^2} mr^2 \dot{\phi} \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi & again using: L = mr^2 \dot{\phi} \end{aligned}$$

$$\begin{array}{c} Radius r: \\ r = \frac{\lambda}{1 - \varepsilon \cos\phi} = \frac{L^2/km}{1 - \varepsilon \cos\phi} & \phi = \frac{L}{1 - \varepsilon \cos\phi} \\ \dot{r} = \frac{\lambda}{dt} = \frac{L^2}{1 - \varepsilon \cos\phi} & \phi = \frac{L}{mr^2} = \frac{L}{m} \frac{1}{r^2} = \frac{L}{m} \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos\phi)^2 \\ \dot{r} = \frac{dr}{dt} = \frac{L^2}{km} \frac{-\frac{d}{dt}(-\varepsilon \cos\phi)}{(1 - \varepsilon \cos\phi)^2} & r\phi = \frac{L}{mr} = \frac{L}{m} \frac{1}{r} = \frac{L}{m} \left(\frac{km}{L^2}\right)^2 (1 - \varepsilon \cos\phi)^2 \\ \dot{r} = \frac{L^2}{km} \frac{-\varepsilon \sin\phi\phi}{(1 - \varepsilon \cos\phi)^2} & r\phi = \frac{L}{mr} = \frac{L}{m} \frac{1}{r} = \frac{L}{m} \left(\frac{km}{L^2}\right) (1 - \varepsilon \cos\phi) = \frac{k}{L} (1 - \varepsilon \cos\phi)^2 \\ \frac{L^2}{km} \left(\frac{km}{L^2}\right)^2 \dot{r}^2 \dot{\phi} \varepsilon \sin\phi & using: \frac{1}{(1 - \varepsilon \cos\phi)^2} = \left(\frac{km}{L^2}\right)^2 r^2 \\ \frac{L^2}{km} \frac{km}{L^2} = \frac{k}{r} \frac{m^2 \dot{\phi} \varepsilon \sin\phi}{L^2} & using: \frac{1}{(1 - \varepsilon \cos\phi)^2} = \left(\frac{km}{L^2}\right)^2 r^2 \\ \frac{L^2}{cartesian} x = r \cos\phi: \\ \dot{x} = \frac{dx}{dt} = -\dot{r} \cos\phi - \sin\phi r\phi & \dot{y} = \frac{dy}{dt} = -\dot{r} \sin\phi + \cos\phi r\phi \end{array}$$

$$\begin{aligned} \text{Radius } r: \\ r &= \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi} \\ \dot{r} &= \frac{\lambda}{dt} = \frac{L^2}{km} - \frac{d}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} &= \frac{dr}{dt} = \frac{L^2}{km} - \frac{d}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} &= \frac{dr}{dt} = \frac{L^2}{km} - \frac{d}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} &= \frac{L^2}{km} - \frac{\varepsilon \sin \phi \phi}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} &= \frac{L^2}{km} - \frac{\varepsilon \sin \phi \phi}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} &= \frac{L^2}{km} - \frac{\varepsilon \sin \phi \phi}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} &= \frac{L^2}{km} - \frac{\varepsilon \sin \phi \phi}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} &= \frac{L^2}{km} - \frac{\varepsilon \sin \phi \phi}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} &= \frac{L^2}{km} - \frac{\varepsilon \sin \phi \phi}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} &= \frac{L^2}{km} - \frac{L^2}{km} - \frac{k}{L^2} - \frac{\varepsilon \sin \phi}{k} \\ \dot{r} &= \frac{k}{L} - \frac{\varepsilon \sin \phi}{k} \\ \dot{r} &= \frac{k}{L} \\ \dot{r} &=$$

Radius r:

$$r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2/km}{1 - \varepsilon \cos \phi}$$

$$r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2/km}{1 - \varepsilon \cos \phi}$$

$$r = \frac{dr}{dt} = \frac{L^2}{km} - \frac{d}{(1 - \varepsilon \cos \phi)^2}$$

$$r = \frac{dr}{dt} = \frac{L^2}{km} - \frac{d}{(1 - \varepsilon \cos \phi)^2}$$

$$r = \frac{L^2}{km} - \frac{\varepsilon \sin \phi}{(1 - \varepsilon \cos \phi)^2}$$

$$r = -\frac{L^2}{km} - \frac{k}{L^2} - \frac{\varepsilon \sin \phi}{L^2}$$

$$r = -\frac{L^2}{km} - \frac{k}{L^2} - \frac{k}{L$$

$$\begin{aligned} \text{Radius } r: \\ r = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi} \\ \dot{r} = \frac{\lambda}{1 - \varepsilon \cos \phi} = \frac{L^2 / km}{1 - \varepsilon \cos \phi} \\ \dot{r} = \frac{dr}{dt} = \frac{L^2}{km} - \frac{d}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} = \frac{dr}{dt} = \frac{L^2}{km} - \frac{c \sin \phi \phi}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} = \frac{L^2}{km} - \frac{c \sin \phi \phi}{(1 - \varepsilon \cos \phi)^2} \\ \dot{r} = -\frac{L^2}{km} \left(\frac{km}{L^2}\right)^2 r^2 \phi \varepsilon \sin \phi \\ \dot{r} = -\frac{k}{L^2} mr^2 \phi \varepsilon \sin \phi = -\frac{k}{L} \varepsilon \sin \phi \\ \dot{r} = -\frac{k}{L} \cos \phi - \varepsilon \\ \dot{r} = -\frac{k}{L} \cos$$

Rutherford scattering and hyperbolic orbit geometry Backward vs forward scattering angles and orbit construction example Parabolic "kite" and orbital envelope geometry Differential and total scattering cross-sections

Eccentricity vector $\boldsymbol{\varepsilon}$ and (ε, λ) -geometry of orbital mechanics Projection $\boldsymbol{\varepsilon} \cdot \mathbf{r}$ geometry of $\boldsymbol{\varepsilon}$ -vector and orbital radius \mathbf{r} Review and connection to usual orbital algebra (previous lecture) Projection $\boldsymbol{\varepsilon} \cdot \mathbf{p}$ geometry of $\boldsymbol{\varepsilon}$ -vector and momentum $\mathbf{p} = m\mathbf{v}$

General geometric orbit construction using ε -vector and (γ, R) -parameters Derivation of ε -construction by analytic geometry

Coulomb orbit algebra of ε-vector and Kepler dynamics of momentum p=mv
 Example of complete (r,p)-geometry of elliptical orbit

Connection formulas for (γ, \mathbf{R}) *-parameters with* (a, b) *and* (ε, λ)













0 > R = KE/PE > -1 scale subtends angle 2γ with length $2r \sin\gamma$ as is derived before on p. 66-70. Note similarity of (**R**,**r**)-triangle in **r**-circle of radius r to that in **p**-circle of diameter p above. Rutherford scattering and hyperbolic orbit geometry Backward vs forward scattering angles and orbit construction example Parabolic "kite" and orbital envelope geometry Differential and total scattering cross-sections

Eccentricity vector $\boldsymbol{\varepsilon}$ and (ε, λ) -geometry of orbital mechanics Projection $\boldsymbol{\varepsilon} \cdot \mathbf{r}$ geometry of $\boldsymbol{\varepsilon}$ -vector and orbital radius \mathbf{r} Review and connection to usual orbital algebra (previous lecture) Projection $\boldsymbol{\varepsilon} \cdot \mathbf{p}$ geometry of $\boldsymbol{\varepsilon}$ -vector and momentum $\mathbf{p} = m\mathbf{v}$

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• Connection formulas for (γ, \mathbf{R}) -parameters with (a, b) and (ε, λ)

Algebra of ε -construction geometry The eccentricity parameter relates ratios $R = \frac{KE}{PE}$ and $\frac{b^2}{a^2}$

$$\varepsilon^{2} = 1 + 4R(R+1)\sin^{2}\gamma$$
$$= 1 - \frac{b^{2}}{a^{2}} \quad \text{for ellipse} \quad (\varepsilon < 1)$$
$$= 1 + \frac{b^{2}}{a^{2}} \quad \text{for hyperbola } (\varepsilon > 1)$$

Three pairs of parameters for Coulomb orbits: 1.Cartesian (a,b), 2.Physics (E,L), 3.Polar (ε , λ) Now we relate a 4th pair: 4.Initial (γ ,**R**) Algebra of ε -construction geometry The eccentricity parameter relates ratios $R = \frac{KE}{PE}$ and $\frac{b^2}{a^2}$

 $\varepsilon^2 = 1 + 4R(R+1)\sin^2\gamma$

and $\frac{b^2}{a^2}$ 1.Cartesian (a,b), 2.Physics (E,L), 3.Polar (ε, λ) Now we relate a 4th pair: 4.Initial (γ, R)

Three pairs of parameters for Coulomb orbits:

$$= 1 - \frac{b^2}{a^2} \quad \text{for ellipse} \quad (\varepsilon < 1) \quad \text{where:} \quad 4R(R+1)\sin^2\gamma = -\frac{b^2}{a^2} = \varepsilon^2 - 1$$
$$= 1 + \frac{b^2}{a^2} \quad \text{for hyperbola} \ (\varepsilon > 1) \quad \text{where:} \quad 4R(R+1)\sin^2\gamma = +\frac{b^2}{a^2} = \varepsilon^2 - 1$$

Algebra of ε -construction geometry The eccentricity parameter relates ratios $R = \frac{KE}{PE}$ and $\frac{b^2}{a^2}$ Three pairs of parameters for Coulomb orbits: 1.Cartesian (a,b), 2.Physics (E,L), 3.Polar (ε , λ) Now we relate a 4th pair: 4.Initial (γ ,**R**)

$$\varepsilon^{2} = 1 + 4R(R+1)\sin^{2}\gamma$$

$$= 1 - \frac{b^{2}}{a^{2}} \text{ for ellipse } (\varepsilon < 1) \text{ where: } 4R(R+1)\sin^{2}\gamma = -\frac{b^{2}}{a^{2}} = \varepsilon^{2} - 1 \text{ implying: } R(R+1) < 0$$

$$= 1 + \frac{b^{2}}{a^{2}} \text{ for hyperbola } (\varepsilon > 1) \text{ where: } 4R(R+1)\sin^{2}\gamma = +\frac{b^{2}}{a^{2}} = \varepsilon^{2} - 1 \text{ implying: } R(R+1) > 0$$
Three pairs of parameters for Coulomb orbits:Three pairs of parameters for Coulomb orbits:I. Cartesian (a,b), 2.Physics (E,L), 3.Polar (ε , λ)Now we relate a 4th pair: 4.Initial (γ , R) $\varepsilon^2 = 1 + 4R(R+1)\sin^2\gamma$ $= 1 - \frac{b^2}{a^2}$ for ellipse ($\varepsilon < 1$) where: $4R(R+1)\sin^2\gamma = -\frac{b^2}{a^2} = \varepsilon^2 - 1$ implying: R(R+1) < 0 (or: $-R^2 > R$) $= 1 + \frac{b^2}{a^2}$ for hyperbola ($\varepsilon > 1$) where: $4R(R+1)\sin^2\gamma = +\frac{b^2}{a^2} = \varepsilon^2 - 1$ implying: R(R+1) > 0 (or: $-R^2 < R$)
(or: -1 > R > 0)

Algebra of ε -construction geometry The eccentricty parameter relates ratios $R = \frac{KE}{PE}$ and $\frac{b^2}{a^2}$ Three pairs of parameters for Coulomb orbits: 1. Cartesian (a,b), 2. Physics (E,L), 3. Polar (ε, λ) Now we relate a 4th pair: 4. Initial (γ, R) $\varepsilon^2 = 1+4R(R+1)\sin^2\gamma$ $= 1 - \frac{b^2}{a^2}$ for ellipse $(\varepsilon < 1)$ where: $4R(R+1)\sin^2\gamma = -\frac{b^2}{a^2} = \varepsilon^2 - 1$ implying: R(R+1) < 0 (or: $-R^2 > R$) (or: 0 > R > -1) $= 1 + \frac{b^2}{a^2}$ for hyperbola $(\varepsilon > 1)$ where: $4R(R+1)\sin^2\gamma = +\frac{b^2}{a^2} = \varepsilon^2 - 1$ implying: R(R+1) > 0 (or: $-R^2 < R$) $(or: -R^2 < R)$ $(or: -R^2 < R)$

Three pairs of parameters for Coulomb orbits: Algebra of ε -construction geometry 1. Cartesian (a,b), 2. Physics (E,L), 3. Polar (ε , λ) The *eccentricty* parameter relates ratios $R = \frac{KE}{PE}$ and $\frac{b^2}{a^2}$ Now we relate a 4th pair: 4. Initial (γ, \mathbf{R}) $\varepsilon^2 = 1 + 4R(R+1)\sin^2\gamma$ $=1-\frac{b^2}{a^2} \quad \text{for ellipse} \quad (\varepsilon < 1) \text{ where:} \quad 4R(R+1)\sin^2\gamma = -\frac{b^2}{a^2} = \varepsilon^2 - 1 \text{ implying: } R(R+1) < 0 \quad (\text{or: } -R^2 > R) \\ (\text{or: } 0 > R > -1) \quad (\text{or: } 0 > R > -1) \end{cases}$ $= 1 + \frac{b^2}{a^2} \text{ for hyperbola } (\varepsilon > 1) \text{ where: } 4R(R+1)\sin^2\gamma = + \frac{b^2}{a^2} = \varepsilon^2 - 1 \text{ implying: } R(R+1) > 0 \qquad (\text{or: } -R^2 < R) \\ (\text{or: } -1 > R > 0) \end{cases}$ Total $\frac{-k}{2a} = E = energy = KE + PE$ relates ratio $R = \frac{KE}{PE}$ to individual radii a, b, and λ . $\frac{-k}{2a} = E = KE + PE = \frac{R}{PE} + PE = (\frac{R}{R} + 1)PE = (\frac{R}{R} + 1)\frac{-k}{r} \text{ or: } \frac{1}{2a} = (\frac{R}{R} + 1)\frac{1}{r} = (\frac{R}{R} + 1)$ $\left(a = \frac{r}{2(R+1)} = \left(\frac{1}{2(R+1)} \text{ assuming unit initial radius } (r \equiv 1).\right)$

Three pairs of parameters for Coulomb orbits: Algebra of ε -construction geometry 1. Cartesian (a,b), 2. Physics (E,L), 3. Polar (ε , λ) The *eccentricty* parameter relates ratios $R = \frac{KE}{PE}$ and $\frac{b^2}{a^2}$ Now we relate a 4th pair: 4. Initial (γ, \mathbf{R}) $\varepsilon^2 = 1 + 4R(R+1)\sin^2\gamma$ $=1-\frac{b^2}{a^2} \quad \text{for ellipse} \quad (\varepsilon < 1) \quad \text{where:} \quad 4R(R+1)\sin^2\gamma = -\frac{b^2}{a^2} = \varepsilon^2 - 1 \quad \text{implying:} \quad R(R+1) < 0 \qquad (\text{or: } -R^2 > R) \\ (\text{or: } 0 > R > -1) \quad (\text{or: } 0 > = 1 + \frac{b^2}{a^2} \text{ for hyperbola } (\varepsilon > 1) \text{ where: } 4R(R+1)\sin^2\gamma = + \frac{b^2}{a^2} = \varepsilon^2 - 1 \text{ implying: } R(R+1) > 0 \quad (\text{or: } -R^2 < R) \text{ (or: } -R^2 < R) \text{$ Total $\frac{-k}{2a} = E = energy = KE + PE$ relates ratio $R = \frac{KE}{PE}$ to individual radii a, b, and λ . $\frac{-k}{2a} = E = KE + PE = \frac{R}{PE} + PE = (\frac{R}{R} + 1)PE = (\frac{R}{R} + 1)\frac{-k}{r} \text{ or: } \frac{1}{2a} = (\frac{R}{R} + 1)\frac{1}{r} = (\frac{R}{R} + 1)$ $\left| a = \frac{r}{2(R+1)} = \left(\frac{1}{2(R+1)} \text{ assuming unit initial radius } (r \equiv 1). \right) \right|$ $4R(R+1)\sin^2\gamma = \mp \frac{b^2}{a^2} \text{ implies: } 2\sqrt{\mp R(R+1)}\sin\gamma = \frac{b}{a} \text{ or: } b = 2a\sqrt{\mp R(R+1)}\sin\gamma$ $b = r \sqrt{\frac{\mp R}{R+1}} \sin \gamma \left(= \sqrt{\frac{\mp R}{R+1}} \sin \gamma \text{ assuming unit initial radius } (r \equiv 1) \right)$

Three pairs of parameters for Coulomb orbits: Algebra of ε -construction geometry 1. Cartesian (a,b), 2. Physics (E,L), 3. Polar (ε , λ) The *eccentricty* parameter relates ratios $R = \frac{KE}{PE}$ and $\frac{b^2}{a^2}$ Now we relate a 4th pair: 4. Initial (γ, \mathbf{R}) $\varepsilon^2 = 1 + 4R(R+1)\sin^2\gamma$ $=1-\frac{b^2}{a^2} \quad \text{for ellipse} \quad (\varepsilon < 1) \quad \text{where:} \quad 4R(R+1)\sin^2\gamma = -\frac{b^2}{a^2} = \varepsilon^2 - 1 \quad \text{implying:} \quad R(R+1) < 0 \qquad (\text{or: } -R^2 > R) \\ (\text{or: } 0 > R > -1) \quad (\text{or: } 0 > = 1 + \frac{b^2}{a^2} \text{ for hyperbola } (\varepsilon > 1) \text{ where: } 4R(R+1)\sin^2\gamma = + \frac{b^2}{a^2} = \varepsilon^2 - 1 \text{ implying: } R(R+1) > 0 \quad (\text{or: } -R^2 < R) \text{ (or: } -R^2 < R) \text{ (or: } -R^2 < R) \text{ (or: } -1 > R > 0)$ Total $\frac{-k}{2a} = E = energy = KE + PE$ relates ratio $R = \frac{KE}{PE}$ to individual radii a, b, and λ . $\frac{-k}{2a} = E = KE + PE = \frac{R}{PE} + PE = (\frac{R}{R} + 1)PE = (\frac{R}{R} + 1)\frac{-k}{r} \text{ or: } \frac{1}{2a} = (\frac{R}{R} + 1)\frac{1}{r} = (\frac{R}{R} + 1)$ $a = \frac{r}{2(R+1)} = \left(\frac{1}{2(R+1)} \text{ assuming unit initial radius } (r \equiv 1).\right)$ $4R(R+1)\sin^2\gamma = \pm \frac{b^2}{a^2} \text{ implies: } 2\sqrt{\pm R(R+1)}\sin\gamma = \frac{b}{a} \text{ or: } b = 2a\sqrt{\pm R(R+1)}\sin\gamma$ $b = r \sqrt{\frac{\mp R}{R+1}} \sin \gamma \left(= \sqrt{\frac{\mp R}{R+1}} \sin \gamma \text{ assuming unit initial radius } (r \equiv 1) \right)$

Latus radius is similarly related:

$$\lambda = \frac{b^2}{a} = \mp 2r R \sin^2 \gamma$$





Eccentricity vector $\boldsymbol{\varepsilon}$ *and* ($\boldsymbol{\varepsilon}, \boldsymbol{\lambda}$)*-geometry of orbital mechanics* Analytic geometry derivation of ε -construction *Connection formulas for* (a,b) *and* (ε,λ) *with* (γ, R) Detailed ruler & compass construction of ε -vector and orbits $(R=-0.375 \ elliptic \ orbit)$ (R=+0.5 hyperbolic orbit)Properties of Coulomb trajectory families and envelopes Graphical ε -development of orbits Launch angle fixed-Varied launch energy - Launch energy fixed-Varied launch angle Launch optimization and orbit family envelopes



Excerpts from Lect. 27



for orbits in a repulsive field. In either case a radial line marked ±10°,

 $\pm 20^{\circ}$, ..., is the focus locus for an orbit with the initial velocity an

angle ±10°, ±20°, . . . , above the nadir line.

Fig. 11. Sample computer trajectories. (a) Family of hyperbolic orbits with R = 1. (b) Family of hyperbolic orbits with R < 1.



Coulomb envelope geometry











Fig. 5.4.5 in Unit 5 of CMwBANG!





<u>Coullt Web Simulation Repulsive</u> <u>Coulomb Burst - Tight</u>



Rutherford scattering (roughly!)



Eccentricity vector ε and (ε, λ) -geometry of orbital mechanics Analytic geometry derivation of ε -construction Connection formulas for (a,b) and (ε,λ) with (γ, \mathbf{R}) Detailed ruler & compass construction of ε -vector and orbits $(R=-0.375 \ elliptic \ orbit)$ $(R=+0.5 \ hyperbolic \ orbit)$ Properties of Coulomb trajectory families and envelopes \blacktriangleright Graphical ε -development of orbits Launch angle fixed-Varied launch energy Launch energy fixed-Varied launch angle \blacktriangleright Launch optimization and orbit family envelopes



Fig. 11. Sample computer trajectories. (a) Family of hyperbolic orbits with R = 1. (b) Family of hyperbolic orbits with R < 1.

(b)

Lenz Vector...analog computersAJP 44 4 (1974)

Fig. 7. Coordinate grids for orbital analog computers. (a) Elliptical orbit scale (0 > R > -1). This can be used with the apparatus in Figs. 8 or 9. Radial lines marked $\pm 10^{\circ}, \pm 20^{\circ}, \ldots$, are each the focus locus for orbits with an initial velocity $\pm 10^{\circ}, \pm 20^{\circ}, \ldots$, above the horizon line. The circle marked $20^{\circ}, 40^{\circ}, \ldots, 340^{\circ}$ can be taken as the Earth's surface, or any circle inside this one can be taken to be the surface of any celestial body. The *R* values apply correctly in either case, while the velocity values are marked for the former case only. (b) Hyperbolic orbit scale $(0 < R < \infty)$ and $(-\infty < R < -1)$. This can only be used with the apparatus shown in Fig. 9. Outer circles locate foci for orbits of particles attracted to the force center, while inner circles locate foci for orbits in a repulsive field. In either case a radial line marked $\pm 10^{\circ}$, $\pm 20^{\circ}, \ldots$, is the focus locus for an orbit with the initial velocity an angle $\pm 10^{\circ}, \pm 20^{\circ}, \ldots$, above the nadir line.

The Lenz vector and orbital analog computers*

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A single geometrical diagram based on the Lenz vector shows the qualitative and quantitative features of all three types of Coulomb orbits. A simple analog computer can be made for an overhead projector by using this theory, and a number of interesting effects can be efficiently demonstrated.

(1)

I. INTRODUCTION: THE ECCENTRICITY VECTOR

Occasionally, a geometrical construction and accompanying picture is worth a great deal more to the physicist or the physics student than pages of equations and solutions, especially now that computer graphics are so available. Since Newton's time the geometrical approach has come to be regarded as more clumsy than other methods of thought, and some very pretty pictures and proofs of physical phenomena have undoubtedly been lost. An example of such a construction involving Rutherford scattering was discussed in a recent article¹ by my students and myself, and the following is an improvement of this which describes general Coulombic orbit mechanics.

The generalization we shall describe below is based partly on a more recently discovered quantity called the Lenz-Runge vector^{2,3} or the "eccentricity" vector ϵ defined by Eq. (1). There **r** is the position vector of the orbiting particle, **L** is its angular momentum, **p** is its linear momentum, *m* is its mass, and *k* is the gravitational (or electrostatic) coefficient:

$$\epsilon = \mathbf{r}/r - \mathbf{L} \times \mathbf{p}/km.$$

Lately this quantity has received a flurry of attention in group theoretical studies of the hydrogen atom⁴; however, we shall use only its geometrical and classical properties.

In particular, the main property of ϵ is that it is a constant vector for any particle moving according to a Coulomb field. Vector ϵ points along the major axis of ellipse, parabola, or hyperbola, whichever is the appropriate orbit of the particle. Furthermore, the magnitude ϵ of this vector is the eccentricity of the orbit.

To show that this is consistent with the usual formulation, we take the dot product of this vector ϵ with the position **r** as in Eq. (2). This then reduces to the following equation (3) of a conic section in polar coordinates, which is the general orbit equation⁵:

$$\epsilon r \cos \theta = \epsilon \cdot \mathbf{r} = r - \mathbf{L} \times \mathbf{p} \cdot \mathbf{r} / km$$

$$= r + \mathbf{L} \cdot \mathbf{L} / km,$$

$$r = -(L^2 / km)(1 - \epsilon \cos \theta)^{-1}.$$
(3)

In Sec. II a simple geometric construction using these properties is shown to describe qualitatively and quantitatively the Coulomb orbits for all three cases: namely, the attractive case (k < 0) with positive energy, with negative energy, and the repulsive case (k > 0).

S(star)

Fig. 1. Initial position and momentum must be given before construction of the resulting orbit is possible.

Finally, it is shown how this construction leads to an analog computer of orbits that can be made for a few dollars to fit onto an overhead projector. This device can be appreciated by elementary classes (even *large* elementary classes if you use the right projector) when they know only a little about conic sections, since the more tedious mathematics is built into the device.

II. COMPUTING ORBITS BY RULER AND COMPASS

We start by simply listing three steps of an orbit construction while demonstrating their application to a particular case of a satellite orbiting a star. Then a general proof of the steps will be provided along with further discussion and applications.

Suppose you are given the initial position and velocity of a satellite relative to some very massive star. If these quantities are given in a pictorial form which shows the angle γ between momentum vector vector $\mathbf{p} = m\mathbf{v}$ and the radius line PS in Fig. 1, and if the magnitude of \mathbf{v} is given by the ratio R = T/V of the kinetic energy $(T = mv^2/2)$ to the potential energy (V = k/r), then the construction below proceeds immediately. Otherwise, these quantities must be calculated before proceeding. (In the potential energy of the star's gravity we have k = -GMm, where M is the star mass and G is the universal constant of gravitation.) Note that R is minus the squared ratio of initial velocity to the escape velocity in



Fig. 2. Doubling the angle between the momentum and the position vectors gives a line QP which must contain the orbit focus. MARKING PEN TIP WITH GROOVE

Fig. 8. Orbital computer design: cheaper model that computes elliptical orbits only using the scale of Fig. 7(a). A Plexiglas sheet that is about $\frac{1}{3}$ in. thick has a string hole at the orbit's center. A transparency (Xerox, 3-M, etc.) of Fig. 7(a) is taped in position on the underside. (*Caution:* One should avoid marking pens that permanently mark plastic.)

8 go with Fig. 7(a) and can be assembled in a few minutes by using odds and ends. A more sophisticated apparatus is shown in Fig. 9. The simple apparatus of Fig. 8 produces the well-known elliptical orbits and trajectories of planets and satellites, but not the hyperbolic trajectories characteristic of higher-than-escape velocity meteors or of the repulsive Coulomb force problems. The second apparatus (Fig. 9) is designed to handle all cases, provided that the appropriate focal point scale is inserted.

The operation of either plotter begins with the positioning of second focal point S' according to the scale on the plotting board. Then the marking pen is poked into a small identation at P and held while the strings to S and S' are tightened. Finally, you slide the pen out along the board in such a way that the strings stay tight and the desired trajectory is drawn.

The apparatus in Fig. 8 will thus make an ellipse since the sum of distances SP and S'P is constant. The apparatus in Fig. 9 does the same when the spool brakes are



Fig. 9. Orbital computer design: a more elaborate model that computes general Coulomb orbits. The ellipse drawing mode is obtained when, first, the clamp is opened to allow the string to slide and then the spool brakes are tightened after the initial adjustment has been made with the use of the scale in Fig. 7(a). The hyperbola drawing mode is obtained when the string clamp and spool clutch are tightened but the brakes are released. The spools must turn together after the initial adjustment has been made with the use of the scale in Fig. 7(b). One hand can maintain the string tension while drawing the orbit, and the other hand can control the paying out of string. (Alternately, springs on the spool axis accomplish the same thing.)





Fig. 10. Sample computer trajectory problem. One finds the minimumenergy trajectory for a given range ρ . Initial velocity v and θ follow easily from the geometry of the computer scale in Fig. 7(a).

tightened and the string clamp on the marking pen is loosened.

The apparatus in Fig. 9 will produce a hyperbola if the *difference* between distances SP and S'P is constant. This is accomplished by tightening the string clamp and the clutch so that two string spools reel equal amounts of string in or out and the constant difference in length is maintained.

IV. SOME USES FOR COMPUTERS

For the computer to be set up to draw elliptical orbits, there is one important question that can be answered immediately: To throw with *minimum* initial velocity a free falling spacecraft between two fixed points near the earth and a distance ρ apart, what initial angle θ and speed vare needed? (We imagine a fixed coordinate system here, not a rotating one.)

Measuring this range by a great circle angle ρ , we see that the focus of the orbit must be on a line through the Earth's center, making an angle $\rho/2$ with the launch point P. The smallest R circle [recall Fig. 7(a)] intersecting this line is the one tangent to it and represents the solution to the problem. Indeed, the algebraic solution to this problem follows from the diagram in Fig. 10 and is given there. Note that the angle θ approaches 45° as the range becomes small compared to the radius of the Earth.

Note that one may change the radius of the starting point P by simply reinterpreting the scale of the computer. For example, if the starting point is located at a height of, say, four times the Earth's radius, then the velocities marked on this scale are all divided by the square root of this factor, in this case, by 2.

With the computer set up to draw hyperbolic orbits and the appropriate scales available, there are a number of interesting problems to examine. For example, the attractive-field positive-energy scale allows one to exhibit the paths of meteorites. Given the impact direction and speed, one can extrapolate to find its origin.

The hyperbolic computer setup can be used to demonstrate Rutherford scattering for either the repulsive field (see Ref. 1) or the attractive field. At the same time