

Lecture 20
Mon. 11.04 thru Wed. 11.06 2019

Introduction to classical oscillation and resonance

(Ch. 1 of Unit 4)

1D forced-damped-harmonic oscillator equations and Green's function solutions

Linear harmonic oscillator equation of motion.

*Linear **damped**-harmonic oscillator equation of motion.*

Frequency retardation and amplitude damping

Figure of oscillator merit (the 5% solution $3/\Gamma$ and other numbers)

*Linear **forced**-**damped**-harmonic oscillator equation of motion.*

Phase lag and amplitude resonance amplification

Figure of resonance merit: (angular) Quality factor $q = \omega_0/2\Gamma$

*Properties of **Green's function** solutions and their mathematical/physical behavior*

Transient solutions vs. Steady State solutions

*Complete **Green's Solution** for the **FDHO** (**Forced-Damped-Harmonic Oscillator**)*

Quality factors: Beat, lifetimes, and uncertainty

*Approximate Lorentz-**Green's Function** for high quality **FDHO** (Quantum propagator)*

Common Lorentzian (a.k.a. Witch of Agnesi)

Smith Charts (Graph paper)

This Lecture's Reference Link Listing

[Web Resources - front page](#)

[Quantum Theory for the Computer Age](#)

[2017 Group Theory for QM](#)

[UAF Physics UTube channel](#)

[Principles of Symmetry, Dynamics, and Spectroscopy](#)

[2018 Adv CM](#)

[Classical Mechanics with a Bang!](#)

[2018 AMOP](#)

[Modern Physics and its Classical Foundations](#)

[2019 Advanced Mechanics](#)

Lectures #12 through #20

In reverse order

[Smith Chart, Invented by Phillip H. Smith \(1905-1987\)](#)

[Nobelprize.org](#)

[2005 Physics Award](#)

Intra-lecture links: [Page=18](#), [Page=35](#), [Page=47](#)

Oscillt Web Simulations:

[Default/Generic, Weakly Damped #18](#),

Forced : [Way below resonance](#), [On resonance](#)

[Way above resonance](#), [Underdamped](#)

[Complex Response Plot](#)

CoulIt Web Simulations:

[Stark-Coulomb : Bound-state motion in parabolic coordinates](#)

[Molecular Ion : Bound-state motion in hyperbolic coordinates](#)

[Synchrotron Motion, Synchrotron Motion #2](#)

[Mechanical Analog to EM Motion \(YouTube video\)](#)

[iBall demo - Quasi-periodicity \(YouTube video\)](#)

Trebuchet Web Simulations:

[Default/Generic URL](#), [Montezuma's Revenge](#), [Seige of Kenilworth](#),
["Flinger"](#),

[Position Space \(Course\)](#), [Position Space \(Fine\)](#)

[Wacky Waving Solid Metal Arm Flailing Chaos Pendulum - Scooba_Steve-yt-2015](#)

[Triple Double-Pendulum - Cohen-yt-2008](#)

[Punkin Chunkin - TheArmchairCritic-2011](#)

[Jersey Team Claims Title in Punkin Chunkin - sussexcountyonline-1999](#)

[Shooting range for medieval siege weapons. Anybody knows? - twcenter.net/forums](#)

[The Trebuchet - Chevedden-SciAm-1995](#)

[NOVA Builds a Trebuchet](#)

Recent Articles of Interest:

[Springer handbook on Molecular Symmetry and Dynamics - Ch 32 - Molecular Symmetry](#)

[Synthetic Chiral Light for Efficient Control of Chiral Light-Matter Interaction - Ayuso-np-2019](#)

[A Semi-Classical Approach to the Calculation of Highly Excited Rotational Energies for ...](#)

[Asymmetric-Top Molecules - Schmiedt-pccp-2017](#)

[Quantum Chaos - An Introduction - Stockmann-cup-2006, Review by E. Heller](#)

[Tunable and broadband coherent perfect absorption by ultrathin blk phos metasurfaces - Guo-josab-2019](#)

[Quantum Supremacy Using a Programmable Superconducting Processor - Arute-n-2019](#)

[Vortex Detection in Vector Fields Using Geometric Algebra - Pollock-aaca-2013.pdf](#)

Pirelli Relativity Challenge (Introduction level) - Visualizing Waves:

[Using Earth as a clock](#),

[Tesla's AC Phasors](#),

[Phasors using complex numbers](#).

[CM wBang Unit 1 - Chapter 10, pdf_page=135](#)

[Calculus of exponentials, logarithms, and complex fields](#),

[RelaWavity Web Simulation - Unit Circle and Hyperbola \(Mixed labeling\)](#)

Select, exciting, and related Research

[Clifford Algebra And The Projective Model Of Homogeneous Metric Spaces - Foundations - Sokolov-x-2013](#)

[Geometric Algebra 3 - Complex Numbers - MacDonald-yt-2015](#)

[Biquaternion -Complexified Quaternion- Roots of -1 - Sangwine-x-2015](#)

[An Introduction to Clifford Algebras and Spinors - Vaz-Rocha-op-2016](#)

[Unified View on Complex Numbers and Quaternions- Bongardt-wcmms-2015](#)

[Complex Functions and the Cauchy-Riemann Equations - complex2 - Friedman-columbia-2019](#)

[An sp-hybridized Molecular Carbon Allotrope- cyclo-18-carbon - Kaiser-s-2019](#)

[An Atomic-Scale View of Cyclocarbon Synthesis - Maier-s-2019](#)

[Discovery Of Topological Weyl Fermion Lines And Drumhead Surface States in a Room Temperature Magnet - Belopolski-s-2019](#)

["Weyl"ing away Time-reversal Symmetry - Neto-s-2019](#)

[Non-Abelian Band Topology in Noninteracting Metals - Wu-s-2019](#)

[What Industry Can Teach Academia - Mao-s-2019](#)

[RoVib- quantum state resolution of the C60 fullerene - Changala-Ye-s-2019 \(Alt\)](#)

[A Degenerate Fermi Gas of Polar molecules - DeMarco-s-2019](#)

An assist from *Physics Girl* (YouTube Channel):

[How to Make VORTEX RINGS in a Pool](#)

[Crazy pool vortex - pg-yt-2014](#)

[Fun with Vortex Rings in the Pool - pg-yt-2014](#)

Excerpts (Page 44-47 in *Preliminary Draft*) from the

[Geometric Algebra- A Guided Tour through Space and Time - Reimer-www-2019](#)

Running Reference Link Listing

Lectures #11 through #7

In reverse order

Eric J Heller Gallery:

[Main portal](#), [Consonance and Dissonance II](#), [Bessel 21](#), [Chladni](#)

[The Semiclassical Way to Molecular Spectroscopy - Heller-acs-1981](#)
[Quantum dynamical tunneling in bound states - Davis-Heller-jcp-1981](#)

[Pendulum Web Simulation](#)

[Cycloidulum Web Simulation](#)

Links to previous lecture: [Page=74](#), [Page=75](#), [Page=79](#)

[Pendulum Web Sim](#)

[Cycloidulum Web Sim](#)

JerKIt Web Simulations: [Basic/Generic: Inverted](#), [FVPlot](#)

[CMwithBang Lecture 8](#), page=20

[WWW.sciencenewsforstudents.org: Cassini - Saturnian polar vortex](#)

“RelaWavity” Web Simulations:

[2-CW laser wave](#), [Lagrangian vs Hamiltonian](#),

[Physical Terms Lagrangian L\(u\) vs Hamiltonian H\(p\)](#)

[CouIt Web Simulation of the Volcanoes of Io](#)

[BohrIt Multi-Panel Plot:](#)

[Relativistically shifted Time-Space plots of 2 CW light waves](#)

BoxIt Web Simulations:

[Generic/Default](#)

[Most Basic A-Type](#)

[Basic A-Type w/reference lines](#)

[Basic A-Type A-Type with Potential energy](#)

[A-Type with Potential energy and Stokes Plot](#)

[A-Type w/3 time rates of change](#)

[A-Type w/3 time rates of change with Stokes Plot](#)

[B-Type \(A=1.0, B=-0.05, C=0.0, D=1.0\)](#)

RelaWavity Web Elliptical Motion Simulations:

[Orbits with b/a=0.125](#)

[Orbits with b/a=0.5](#)

[Orbits with b/a=0.7](#)

[Exegesis with b/a=0.125](#)

[Exegesis with b/a=0.5](#)

[Exegesis with b/a=0.7](#)

[Contact Ellipsometry](#)

CouIt Web Simulations:

[Basic/Generic](#)

[Exploding Starlet](#)

[Volcanoes of Io \(Color Quantized\)](#)

JerKIt Web Simulations:

[Basic/Generic](#)

[Catcher in the Eye - IHO with Linear Hooke perturbation - Force-potential-Velocity Plot](#)

OscillatorPE Web Simulation:

[Coulomb-Newton-Inverse_Square](#),

[Hooke-Isotropic Harmonic](#),

[Pendulum-Circular_Constraint](#)

[AMOP Ch 0 Space-Time Symmetry - 2019](#)

[Seminar at Rochester Institute of Optics, Aux. slides-2018](#)

NASA Astronomy Picture of the Day -

[Io: The Prometheus Plume \(Just Image\)](#)

[NASA Galileo - Io's Alien Volcanoes](#)

[New Horizons - Volcanic Eruption Plume on Jupiter's moon IO](#)

[NASA Galileo - A Hawaiian-Style Volcano on Io](#)

[Pirelli Site: Phasors animation](#)

[CMwithBang Lecture #6](#), page=70 (9.10.18)

Select, exciting, and related Research & Articles of Interest:

[Burning a hole in reality—design for a new laser may be powerful enough to pierce space-time - Sumner-KOS-2019](#)

[Trampoline mirror may push laser pulse through fabric of the Universe - Lee-ArsTechnica-2019](#)

[Achieving Extreme Light Intensities using Optically Curved Relativistic Plasma Mirrors - Vincenti-prl-2019](#)

[A Soft Matter Computer for Soft Robots - Garrad-sr-2019](#)

[Correlated Insulator Behaviour at Half-Filling in Magic-Angle Graphene Superlattices - cao-n-2018](#)

[Sorting ultracold atoms in a three-dimensional optical lattice in a realization of Maxwell's Demon - Kumar-n-2018](#)

[Synthetic three-dimensional atomic structures assembled atom by atom - Barredo-n-2018](#)

Older ones:

[Wave-particle duality of C60 molecules - Arndt-ltn-1999](#)

[Optical Vortex Knots - One Photon At A Time - Tempone-Wiltshire-Sr-2018](#)

[Baryon Deceleration by Strong Chromofields in Ultrarelativistic](#)

[Nuclear Collisions - Mishustin-PhysRevC-2007](#), APS Link & Abstract

[Hadronic Molecules - Guo-x-2017](#)

[Hidden-charm pentaquark and tetraquark states - Chen-pr-2016](#)

Running Reference Link Listing

Lectures #6 through #1

In reverse order

[RelaWavity Web Simulation: Contact Ellipsometry](#)

[BoxIt Web Simulation: Elliptical Motion \(A-Type\)](#)

[CMwBang Course: Site Title Page](#)

[Pirelli Relativity Challenge: Describing Wave Motion With Complex Phasors](#)

[UAF Physics UTube channel](#)

[Velocity Amplification in Collision Experiments Involving Superballs - Harter, 1971](#)

[MIT OpenCourseWare: High School/Physics/Impulse and Momentum](#)

[Hubble Site: Supernova - SN 1987A](#)

[BounceIt Web Animation - Scenarios:](#)

[49:1 y vs t, 49:1 V2 vs V1, 1:500:1 - 1D Gas Model w/ faux restorative force \(Cool\),](#)

[1:500:1 - 1D Gas \(Warm\), 1:500:1 - 1D Gas Model \(Cool, Zoomed in\),](#)

[Farey Sequence - Wolfram](#)

[Fractions - Ford-AMM-1938](#)

[Monstermash BounceIt Animations:](#)

[1000:1 - V2 vs V1, 1000:1 with t vs x - Minkowski Plot](#)

[Quantum Revivals of Morse Oscillators and Farey-Ford Geometry - Li-Harter-2013](#)

[Quantum Revivals of Morse Oscillators and Farey-Ford Geometry - Li-Harter-cpl-2015](#)

[Quant. Revivals of Morse Oscillators and Farey-Ford Geom. - Harter-Li-CPL-2015 \(Publ.\)](#)

[Velocity Amplification in Collision Experiments Involving Superballs-Harter-1971](#)

[WaveIt Web Animation - Scenarios:](#)

[Quantum Carpet, Quantum Carpet wMBars,](#)

[Quantum Carpet BCar, Quantum Carpet BCar wMBars](#)

[Wave Node Dynamics and Revival Symmetry in Quantum Rotors - Harter-JMS-2001](#)

[Wave Node Dynamics and Revival Symmetry in Quantum Rotors - Harter-jms-2001 \(Publ.\)](#)

[AJP article on superball dynamics](#)

[AAPT Summer Reading List](#)

[Scitation.org - AIP publications](#)

[HarterSoft Youtube Channel](#)

[BounceIt Web Animation - Scenarios:](#)

[Generic Scenario: 2-Balls dropped no Gravity \(7:1\) - V vs V Plot \(Power=4\)](#)

[1-Ball dropped w/Gravity=0.5 w/Potential Plot: Power=1, Power=4](#)

[7:1 - V vs V Plot: Power=1](#)

[3-Ball Stack \(10:3:1\) w/Newton plot \(y vs t\) - Power=4](#)

[3-Ball Stack \(10:3:1\) w/Newton plot \(y vs t\) - Power=1](#)

[3-Ball Stack \(10:3:1\) w/Newton plot \(y vs t\) - Power=1 w/Gaps](#)

[4-Ball Stack \(27:9:3:1\) w/Newton plot \(y vs t\) - Power=4](#)

[4-Newton's Balls \(1:1:1:1\) w/Newtonian plot \(y vs t\) - Power=4 w/Gaps](#)

[6-Ball Totally Inelastic \(1:1:1:1:1\) w/Gaps: Newtonian plot \(t vs x\), V6 vs V5 plot](#)

[5-Ball Totally Inelastic Pile-up w/ 5-Stationary-Balls - Minkowski plot \(t vs x1\) w/Gaps](#)

[1-Ball Totally Inelastic Pile-up w/ 5-Stationary-Balls - Vx2 vs Vx1 plot w/Gaps](#)

[BounceIt Dual plots](#)

[m1:m2 = 3:1](#)

[v2 vs v1 and V2 vs V1, \(v1, v2\)=\(1, 0.1\), \(v1, v2\)=\(1, 0\)](#)

[y2 vs y1 plots: \(v1, v2\)=\(1, 0.1\), \(v1, v2\)=\(1, 0\), \(v1, v2\)=\(1, -1\)](#)

[Estrangian plot V2 vs V1: \(v1, v2\)=\(0, 1\), \(v1, v2\)=\(1, -1\)](#)

[m1:m2 = 4:1](#)

[v2 vs v1, y2 vs y1](#)

[m1:m2 = 100:1, \(v1, v2\)=\(1, 0\): V2 vs V1 Estrangian plot, y2 vs y1 plot](#)

[With g=0 and 70:10 mass ratio](#)

[With non zero g, velocity dependent damping and mass ratio of 70:35](#)

[M1=49, M2=1 with Newtonian time plot](#)

[M1=49, M2=1 with V2 vs V1 plot](#)

[Example with friction](#)

[Low force constant with drag displaying a Pass-thru, Fall-Thru, Bounce-Off](#)

[m1:m2= 3:1 and \(v1, v2\) = \(1, 0\) Comparison with Estrangian](#)

X2 paper: [Velocity Amplification in Collision Experiments Involving Superballs - Harter, et. al. 1971 \(pdf\)](#)

Car Collision Web Simulator: <https://modphys.hosted.uark.edu/markup/CMMotionWeb.html>

Superball Collision Web Simulator: <https://modphys.hosted.uark.edu/markup/BounceItWeb.html>; with Scenarios: [1007](#)

[BounceIt web simulation with g=0 and 70:10 mass ratio](#)

[With non zero g, velocity dependent damping and mass ratio of 70:35](#)

Elastic Collision Dual Panel Space vs Space: [Space vs Time \(Newton\)](#), [Time vs. Space\(Minkowski\)](#)

Inelastic Collision Dual Panel Space vs Space: [Space vs Time \(Newton\)](#), [Time vs. Space\(Minkowski\)](#)

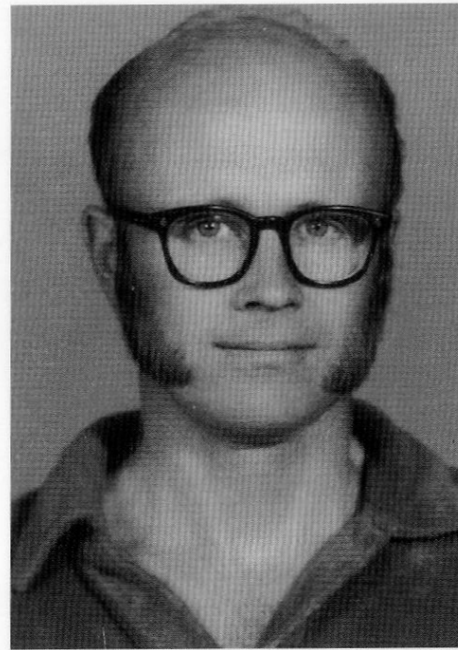
Matrix Collision Simulator: [M1=49, M2=1 V2 vs V1 plot <<Under Construction>>](#)

More Advanced QM and classical references will soon be available through our: [**Mechanics References Page**](#)

(Now in Development)

*Without resonance...
...we are all blind, deaf, and dumb.*

Anonymous



PAULINIA, BRASIL 1976

THE SPEED OF LIGHT IS
299,792,458 METERS PER SECOND!

-- The Purest Light and a Resonance Hero – Ken Evenson (1932-2002) --

Ken Evenson

When travelers punch up their GPS coordinates they owe a debt of gratitude to an under sung hero who, alongside his colleagues and students, often toiled 18 hour days deep inside a laser laboratory lit only by the purest light in the universe.

Ken was an “Indiana Jones” of modern physics. While he may never have been called “Montana Ken,” such a name would describe a real life hero from Bozeman, Montana, whose extraordinary accomplishments in many ways surpass the fictional characters in cinematic thrillers like *Raiders of the Lost Arc*.

Indeed, there were some exciting real life moments shared by his wife Vera, one together with Ken in a canoe literally inches from the hundred-foot drop-off of Brazil’s largest waterfall. But, such outdoor exploits, of which Ken had many, pale in the light of an in-the-lab brilliance and courage that profoundly enriched the world.

Ken is one of few researchers and perhaps the only physicist to be twice listed in the *Guinness Book of Records*. The listings are not for jungle exploits but for his lab’s highest frequency measurement and for a speed of light determination that made c many times more precise due to his lab’s pioneering work with John Hall in laser resonance and metrology[†].

The meter-kilogram-second (mks) system of units underwent a redefinition largely because of these efforts. Thereafter, the speed of light c was set to $299,792,458\text{ms}^{-1}$. The meter was defined in terms of c , instead of the other way around since his time precision had so far trumped that for distance. Without such resonance precision, the Global Positioning System (GPS), the first large-scale wave space-time coordinate system, would not be possible.

Ken’s courage and persistence at the Time and Frequency Division of the Boulder Laboratories in the National Bureau of Standards (now the National Institute of Standards and Technology or NIST) are legendary as are his railings against boneheaded administrators who seemed bent on thwarting his best efforts. Undaunted, Ken’s lab painstakingly exploited the resonance properties of metal-insulator diodes, and succeeded in literally counting the waves of near-infrared radiation and eventually visible light itself.

Those who knew Ken miss him terribly. But, his indelible legacy resonates today as ultra-precise atomic and molecular wave and pulse quantum optics continue to advance and provide heretofore unimaginable capability. Our quality of life depends on their metrology through the Quality and Finesse of the resonant oscillators that are the heartbeats of our technology.

Before being taken by Lou Gehrig’s disease, Ken began ultra-precise laser spectroscopy of unusual molecules such as HO_2 , the radical cousin of the more common H_2O . Like Ken, such radical molecules affect us as much or more than better known ones. But also like Ken, they toil in obscurity, illuminated only by the purest light in the universe.

In 2005 the Nobel Prize in physics was awarded to Glauber, Hall, and Hensch^{††} for laser optics and metrology.

[†] K. M. Evenson, J.S. Wells, F.R. Peterson, B.L. Danielson, G.W. Day, R.L. Barger and J.L. Hall, Phys. Rev. Letters 29, 1346(1972).

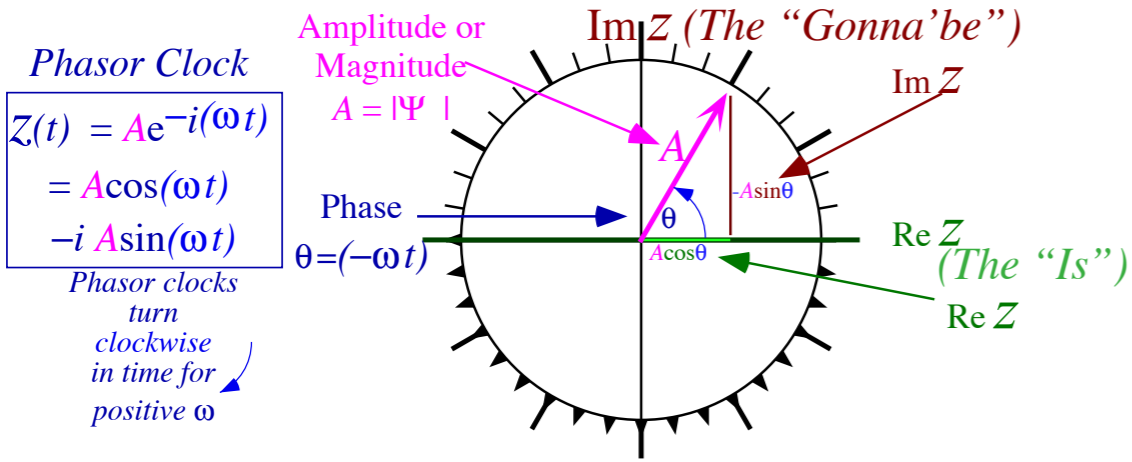
^{††} *The Nobel Prize in Physics, 2005*. <http://nobelprize.org/>

Linear forced-damped-harmonic oscillator equation of motion.

$$F_{total}(t) = m \frac{d^2 z}{dt^2} = F_{damping} + F_{restore} + F_{stimulus}$$

$$\frac{d^2 z}{dt^2} = \frac{F_{damping}}{m} + \frac{F_{restore}}{m} + \frac{F_{stimulus}}{m}$$

Stimulating acceleration $a_{stimulus} = a(t)$ due to stimulating force $F_{stimulus}(t)$ (Typically **E**-field)



$$\frac{d^2 z}{dt^2} + 2\Gamma \frac{dz}{dt} + \omega_0^2 z = a_{stimulus} = \frac{e}{m} E(t)$$

Coordinate $z=z(t)$ is the response coordinate for a particle of mass m and charge e

driven by external **stimulating force** $F_{stimulus}(t) = eE(t)$

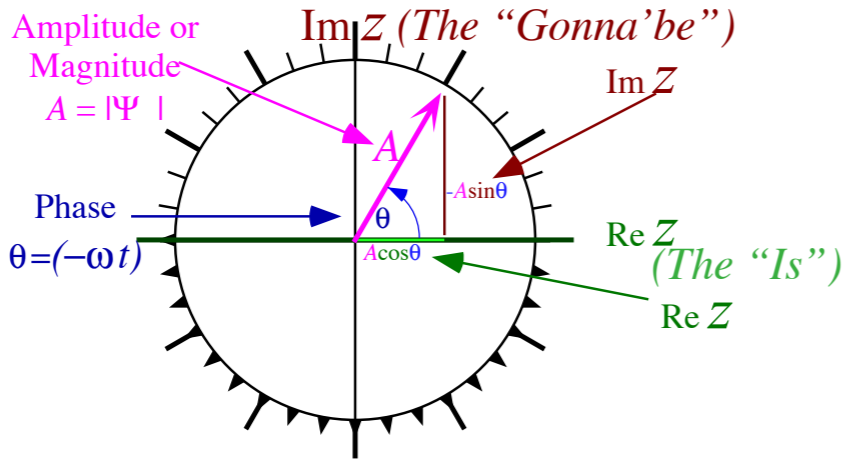
held back by a **harmonic (linear) restoring force** $F_{restore} = -kz, (k = \omega_0^2 m),$

retarded by **frictional damping force** $F_{damping} = -b \frac{dz}{dt}, (b = 2\Gamma m)$

Linear

harmonic oscillator equation of motion.

Phasor Clock
 $Z(t) = Ae^{-i(\omega t)}$
 $= A\cos(\omega t)$
 $-i A\sin(\omega t)$
 Phasor clocks
 turn
 clockwise
 in time for
 positive ω



$$F_{total}(t) = m \frac{d^2 z}{dt^2} = \frac{F_{restore}}{m} \frac{d^2 z}{dt^2} + \omega_0^2 z = 0$$

Coordinate $z=z(t)$ is the response coordinate for a particle of mass m and charge e

held back by a harmonic (linear) restoring force $\longrightarrow F_{restore} = -kz, (k = \omega_0^2 m),$

Linear

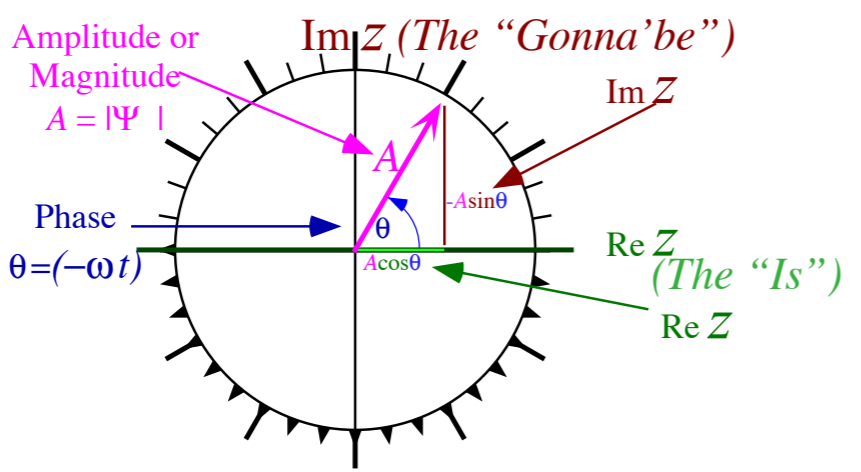
harmonic oscillator equation of motion.

$$F_{total}(t) = m \frac{d^2 z}{dt^2} = F_{restore}$$

$$\frac{d^2 z}{dt^2} = \frac{F_{restore}}{m}$$

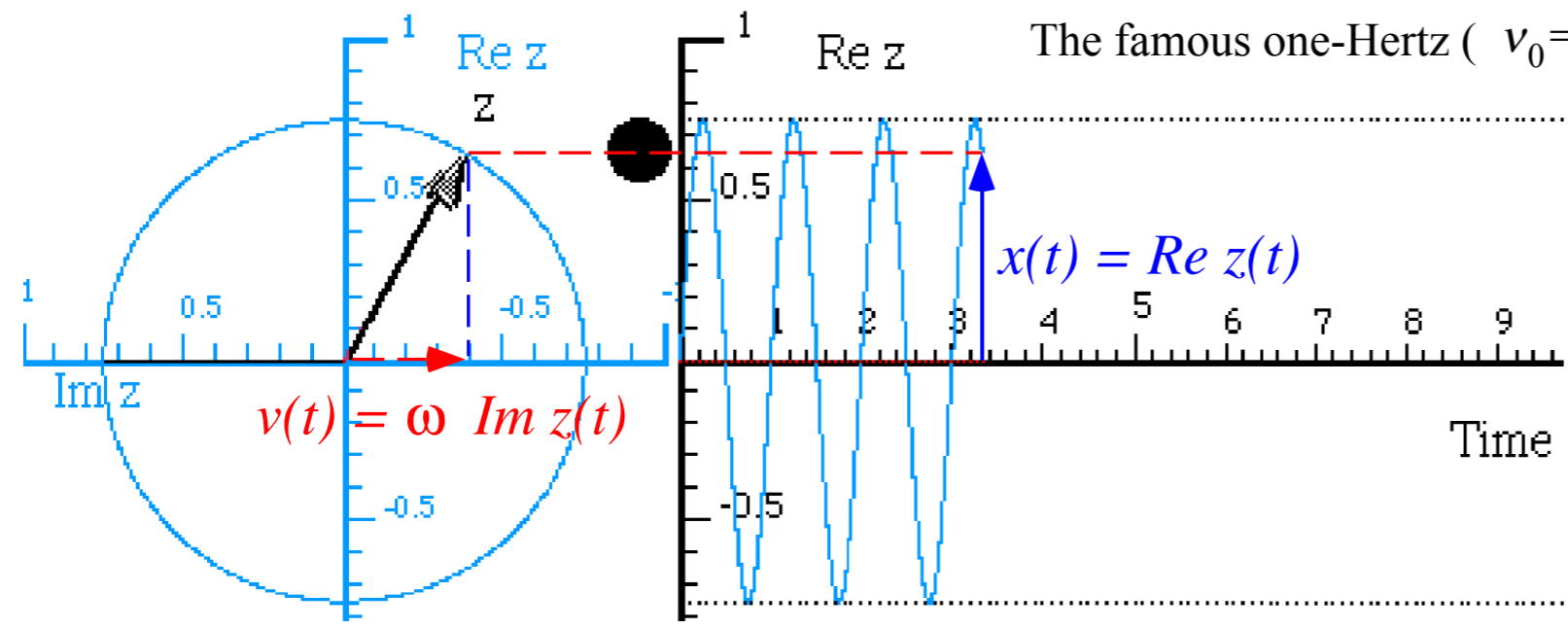
$$\frac{d^2 z}{dt^2} + \omega_0^2 z = 0$$

Phasor Clock
 $Z(t) = Ae^{-i(\omega t)}$
 $= A\cos(\omega t)$
 $-i A\sin(\omega t)$
 Phasor clocks turn clockwise in time for positive ω



Coordinate $z=z(t)$ is the response coordinate for a particle of mass m and charge e

held back by a harmonic (linear) restoring force $\longrightarrow F_{restore} = -kz, (k = \omega_0^2 m),$



The famous one-Hertz ($\nu_0=1/s.$ or: $\omega_0 = 2\pi = 6.2832\text{rad/s.}$) oscillator.

[OscillIt Web Simulation](#)

Fig. 3.2.2 Phasor z and corresponding coordinate versus time plot for $\omega_0=2\pi$ and $\Gamma=0$ <https://modphys.hosted.uark.edu/markup/OscillItWeb.html>

OscillIt Web Simulation (Generic):

Linear *damped-harmonic oscillator equation of motion.*

$$F_{total}(t) = m \frac{d^2 z}{dt^2} = F_{damping} + F_{restore}$$

$$\frac{d^2 z}{dt^2} = \frac{F_{damping}}{m} + \frac{F_{restore}}{m}$$

$$\frac{d^2 z}{dt^2} + 2\Gamma \frac{dz}{dt} + \omega_0^2 z = 0$$

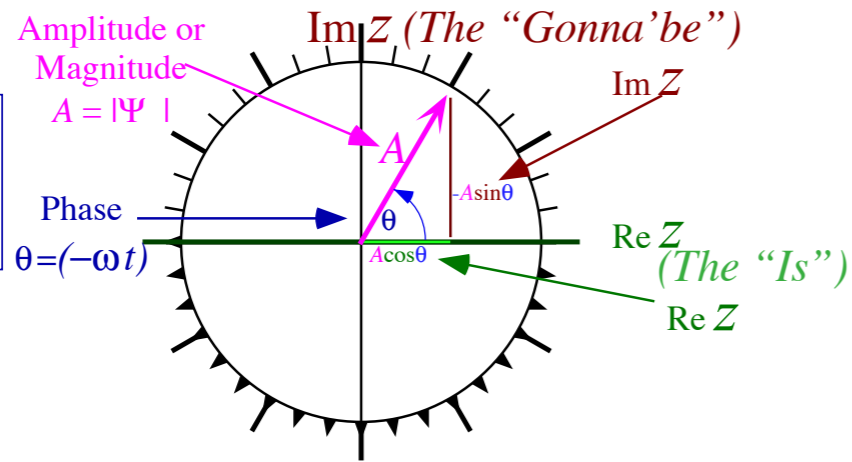
Phasor Clock

$$z(t) = Ae^{-i(\omega t)}$$

$$= A \cos(\omega t)$$

$$-i A \sin(\omega t)$$

Phasor clocks
turn
clockwise
in time for
positive ω



Coordinate $z=z(t)$ is the response coordinate for a particle of mass m and charge e

held back by a **harmonic (linear) restoring force** $\longrightarrow F_{restore} = -kz, \quad (k = \omega_0^2 m),$

retarded by **frictional damping force** $\longrightarrow F_{damping} = -b \frac{dz}{dt}, \quad (b = 2\Gamma m)$

Linear damped-harmonic oscillator equation of motion.

$$F_{total}(t) = m \frac{d^2 z}{dt^2} = F_{damping} + F_{restore}$$

$$\frac{d^2 z}{dt^2} = \frac{F_{damping}}{m} + \frac{F_{restore}}{m}$$

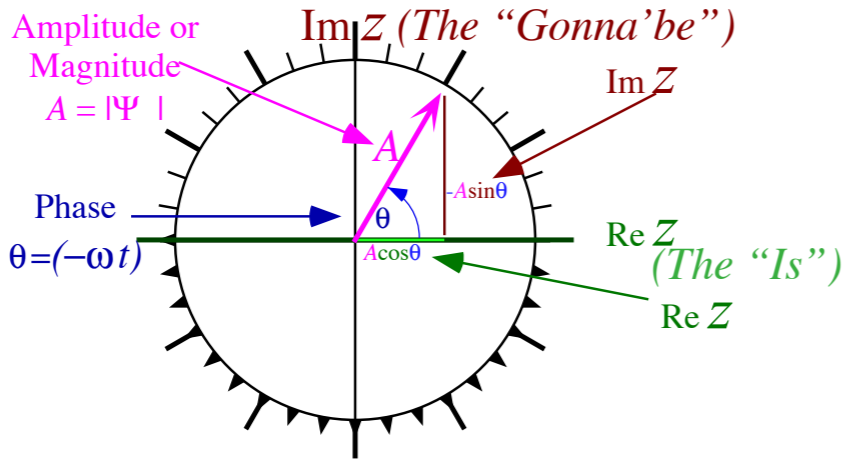
$$\frac{d^2 z}{dt^2} + 2\Gamma \frac{dz}{dt} + \omega_0^2 z = 0$$

Trick:
Set: $z = z(t) = Ae^{-i\omega t}$

$$\left[(-i\omega)^2 + 2\Gamma(-i\omega) + \omega_0^2 \right] e^{-i\omega t} = 0$$

$$\omega^2 + 2i\Gamma\omega - \omega_0^2 = 0$$

Phasor Clock
 $Z(t) = Ae^{-i(\omega t)}$
 $= A\cos(\omega t)$
 $-i A\sin(\omega t)$
 Phasor clocks
 turn
 clockwise
 in time for
 positive ω



Coordinate $z = z(t)$ is the response coordinate for a particle of mass m and charge e

held back by a harmonic (linear) restoring force

$$F_{restore} = -kz$$

retarded by frictional damping force

$$F_{damping} = -b \frac{dz}{dt}$$

Linear damped-harmonic oscillator equation of motion.

$$F_{total}(t) = m \frac{d^2 z}{dt^2} = F_{damping} + F_{restore}$$

$$\frac{d^2 z}{dt^2} = \frac{F_{damping}}{m} + \frac{F_{restore}}{m}$$

$$\frac{d^2 z}{dt^2} + 2\Gamma \frac{dz}{dt} + \omega_0^2 z = 0$$

Trick:

Set: $z = z(t) = Ae^{-i\omega t}$

$$\left[(-i\omega)^2 + 2\Gamma(-i\omega) + \omega_0^2 \right] e^{-i\omega t} = 0$$

$$\omega^2 + 2i\Gamma\omega - \omega_0^2 = 0$$

Solve for: $\omega = \omega_{\pm}$

$$\omega_{\pm} = \frac{-2i\Gamma \pm \sqrt{-4\Gamma^2 + 4\omega_0^2}}{2}$$

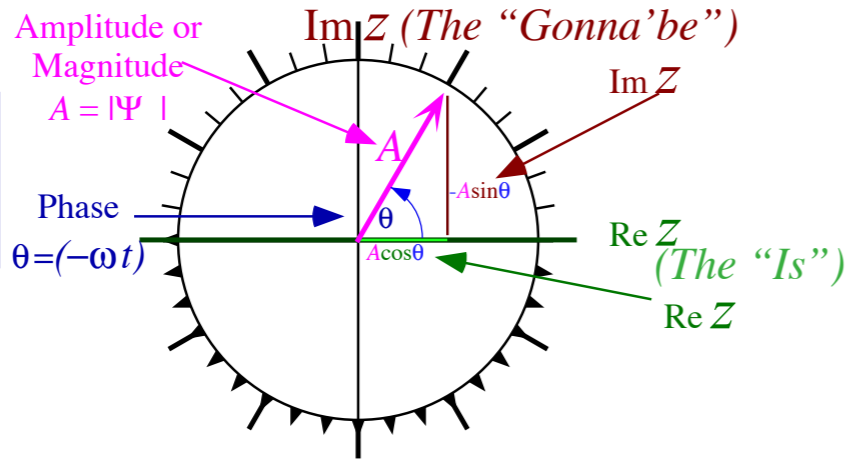
Phasor Clock

$$z(t) = Ae^{-i(\omega t)}$$

$$= A\cos(\omega t)$$

$$-i A\sin(\omega t)$$

Phasor clocks
turn
clockwise
in time for
positive ω



Coordinate $z = z(t)$ is the response coordinate for a particle of mass m and charge e

held back by a harmonic (linear) restoring force

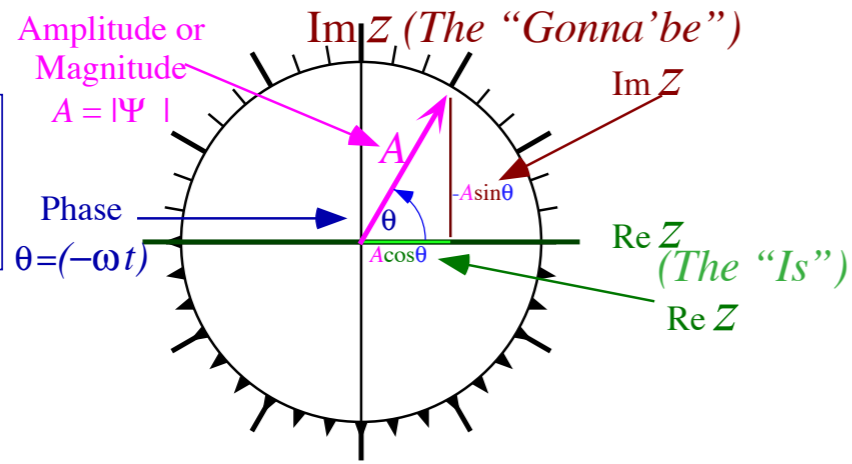
$$F_{restore} = -kz$$

retarded by frictional damping force

$$F_{damping} = -b \frac{dz}{dt}$$

Linear damped-harmonic oscillator equation of motion.

Phasor Clock
 $Z(t) = Ae^{-i(\omega t)}$
 $= A\cos(\omega t)$
 $-i A\sin(\omega t)$
 Phasor clocks
 turn
 clockwise
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$$F_{total}(t) = m \frac{d^2 z}{dt^2} = F_{damping} + F_{restore}$$

$$\frac{d^2 z}{dt^2} = \frac{F_{damping}}{m} + \frac{F_{restore}}{m}$$

$$\frac{d^2 z}{dt^2} + 2\Gamma \frac{dz}{dt} + \omega_0^2 z = 0$$

Trick:
 Set: $z = z(t) = Ae^{-i\omega t}$

$$\left[(-i\omega)^2 + 2\Gamma(-i\omega) + \omega_0^2 \right] e^{-i\omega t} = 0$$

$$\omega^2 + 2i\Gamma\omega - \omega_0^2 = 0$$

Solve for: $\omega = \omega_{\pm}$

$$\omega_{\pm} = \frac{-2i\Gamma \pm \sqrt{-4\Gamma^2 + 4\omega_0^2}}{2}$$

$$= -i\Gamma \pm \sqrt{\omega_0^2 - \Gamma^2}$$

Coordinate $z = z(t)$ is the response coordinate for a particle of mass m and charge e

held back by a harmonic (linear) restoring force

$$F_{restore} = -kz$$

retarded by frictional damping force

$$F_{damping} = -b \frac{dz}{dt}$$

Linear *damped-harmonic oscillator equation of motion.*

$$F_{total}(t) = m \frac{d^2 z}{dt^2} = F_{damping} + F_{restore}$$

$$\frac{d^2 z}{dt^2} = \frac{F_{damping}}{m} + \frac{F_{restore}}{m}$$

$$\frac{d^2 z}{dt^2} + 2\Gamma \frac{dz}{dt} + \omega_0^2 z = 0$$

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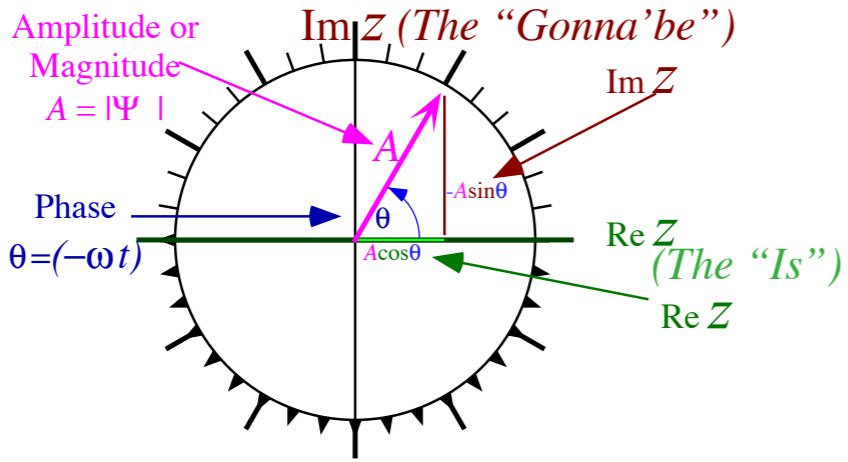
$$\omega_{\pm} = \frac{-2i\Gamma \pm \sqrt{-4\Gamma^2 + 4\omega_0^2}}{2}$$

$$= -i\Gamma \pm \sqrt{\omega_0^2 - \Gamma^2}$$

Solution:

$$\begin{aligned} z(t) &= e^{-i\left(-i\Gamma \pm \sqrt{\omega_0^2 - \Gamma^2}\right)t} \\ &= e^{\left(-\Gamma \pm i\sqrt{\omega_0^2 - \Gamma^2}\right)t} \\ &= e^{-\Gamma t} e^{\pm i\sqrt{\omega_0^2 - \Gamma^2}t} \\ &= e^{-\Gamma t} e^{\pm i\omega_{\Gamma}t} \end{aligned}$$

Phasor Clock
 $Z(t) = Ae^{-i(\omega t)}$
 $= A\cos(\omega t)$
 $-iA\sin(\omega t)$
Phasor clocks turn clockwise in time for positive ω



Coordinate $z = z(t)$ is the response coordinate for a particle of mass m and charge e

held back by a harmonic (linear) restoring force

$$F_{restore} = -kz$$

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$$= -i\Gamma \pm \sqrt{\omega_0^2 - \Gamma^2}$$

Solution:

$$z(t) = e^{-i\left(-i\Gamma \pm \sqrt{\omega_0^2 - \Gamma^2}\right)t}$$

$$= e^{\left(-\Gamma \pm i\sqrt{\omega_0^2 - \Gamma^2}\right)t}$$

$$= e^{-\Gamma t} e^{\pm i\sqrt{\omega_0^2 - \Gamma^2}t}$$

$$= e^{-\Gamma t} e^{\pm i\omega_{\Gamma}t}$$

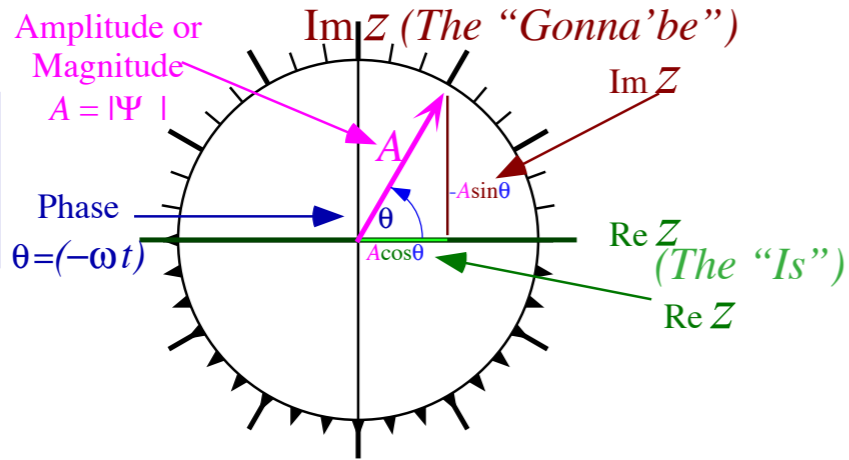
Phasor Clock

$$z(t) = Ae^{-i(\omega t)}$$

$$= A\cos(\omega t)$$

$$-iA\sin(\omega t)$$

Phasor clocks turn clockwise in time for positive ω



Coordinate $z = z(t)$ is the response coordinate for a particle of mass m and charge e

held back by a **harmonic (linear) restoring force**

$$F_{restore} = -kz$$

retarded by **frictional damping force**

$$F_{damping} = -b \frac{dz}{dt}$$

It oscillates at an angular frequency ω_{Γ} reduced slightly by .05% from ω_0 due to damping $\Gamma = 0.2$.

$$\omega_{\Gamma} = \sqrt{\omega_0^2 - \Gamma^2} = \omega_0 - \frac{1}{2}(\Gamma^2 / \omega_0) + \dots = 6.2831853 - 0.003183 + \dots = 6.280002 + \dots = 6.280001$$

Linear *damped-harmonic oscillator equation of motion.*

$$F_{total}(t) = m \frac{d^2 z}{dt^2} = F_{damping} + F_{restore}$$

$$\frac{d^2 z}{dt^2} = \frac{F_{damping}}{m} + \frac{F_{restore}}{m}$$

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Set: $z = z(t) = Ae^{-i\omega t}$

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$$\omega^2 + 2i\Gamma\omega - \omega_0^2 = 0$$

Solve for: $\omega = \omega_{\pm}$

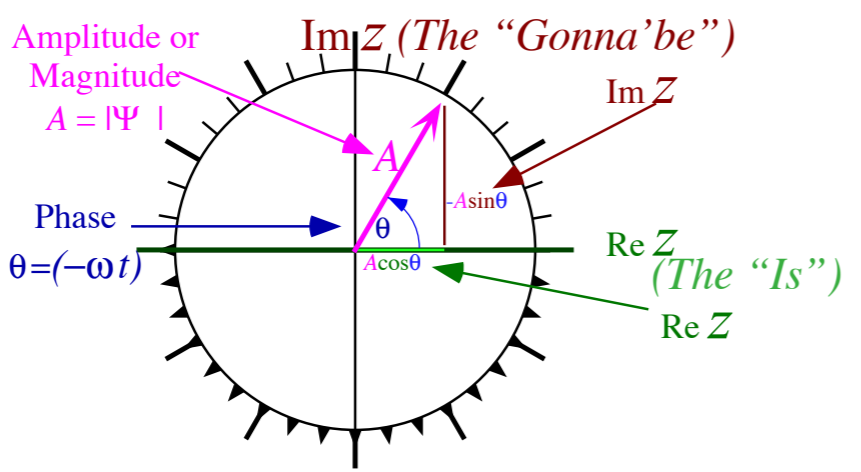
$$\omega_{\pm} = \frac{-2i\Gamma \pm \sqrt{-4\Gamma^2 + 4\omega_0^2}}{2}$$

$$= -i\Gamma \pm \sqrt{\omega_0^2 - \Gamma^2}$$

Solution:

$$\begin{aligned} z(t) &= e^{-i\left(-i\Gamma \pm \sqrt{\omega_0^2 - \Gamma^2}\right)t} \\ &= e^{\left(-\Gamma \pm i\sqrt{\omega_0^2 - \Gamma^2}\right)t} \\ &= e^{-\Gamma t} e^{\pm i\sqrt{\omega_0^2 - \Gamma^2}t} \\ &= e^{-\Gamma t} e^{\pm i\omega_{\Gamma}t} \end{aligned}$$

Phasor Clock
 $Z(t) = Ae^{-i(\omega t)}$
 $= A\cos(\omega t)$
 $-iA\sin(\omega t)$
Phasor clocks turn clockwise in time for positive ω



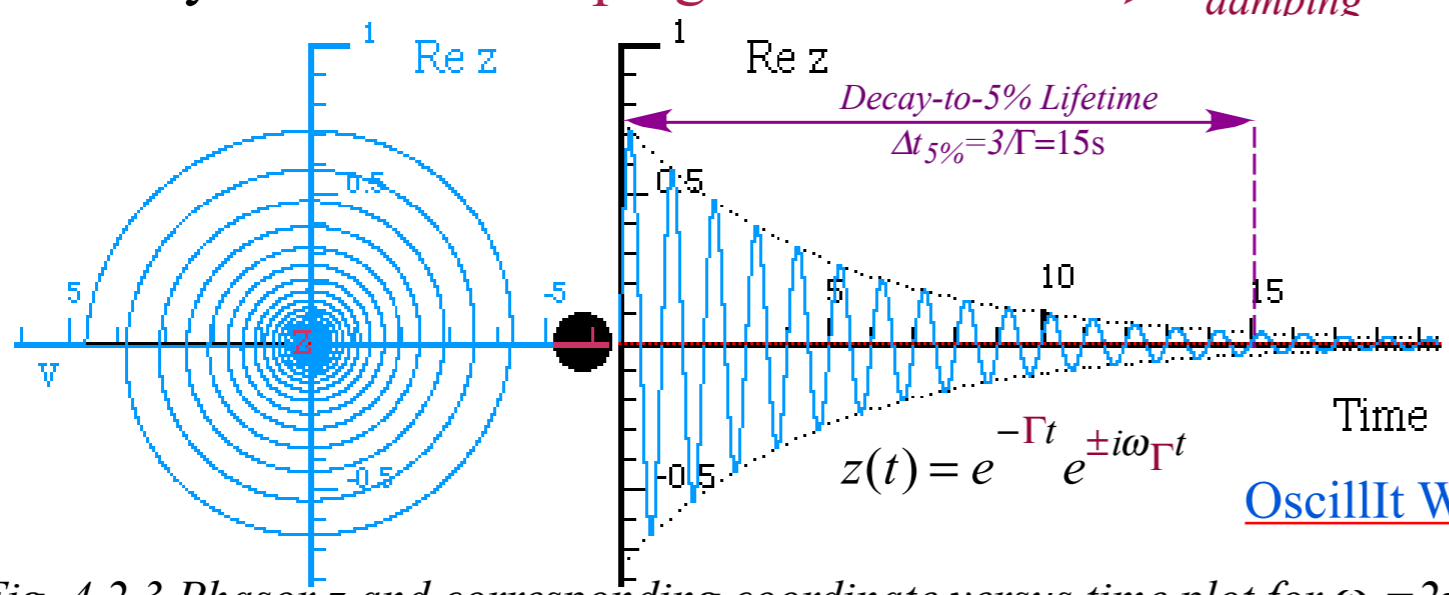
Coordinate $z = z(t)$ is the response coordinate for a particle of mass m and charge e

held back by a harmonic (linear) restoring force

$$F_{restore} = -kz$$

retarded by frictional damping force

$$F_{damping} = -b \frac{dz}{dt}$$



[OscillIt Web Simulation](#)

Fig. 4.2.3 Phasor z and corresponding coordinate versus time plot for $\omega_0 = 2\pi$ and $\Gamma = 0.2$

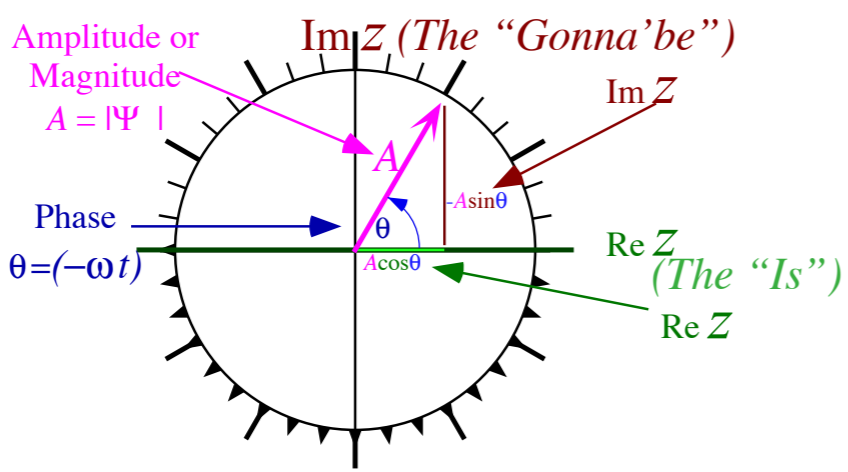
Linear *damped-harmonic oscillator equation of motion.*

$$F_{total}(t) = m \frac{d^2 z}{dt^2} = F_{damping} + F_{restore}$$

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$$\frac{d^2 z}{dt^2} + 2\Gamma \frac{dz}{dt} + \omega_0^2 z = 0$$

Phasor Clock
 $Z(t) = Ae^{-i(\omega t)}$
 $= A \cos(\omega t)$
 $-i A \sin(\omega t)$
 Phasor clocks turn clockwise in time for positive ω



Coordinate $z=z(t)$ is the response coordinate for a particle of mass m and charge e

held back by a **harmonic (linear) restoring force**

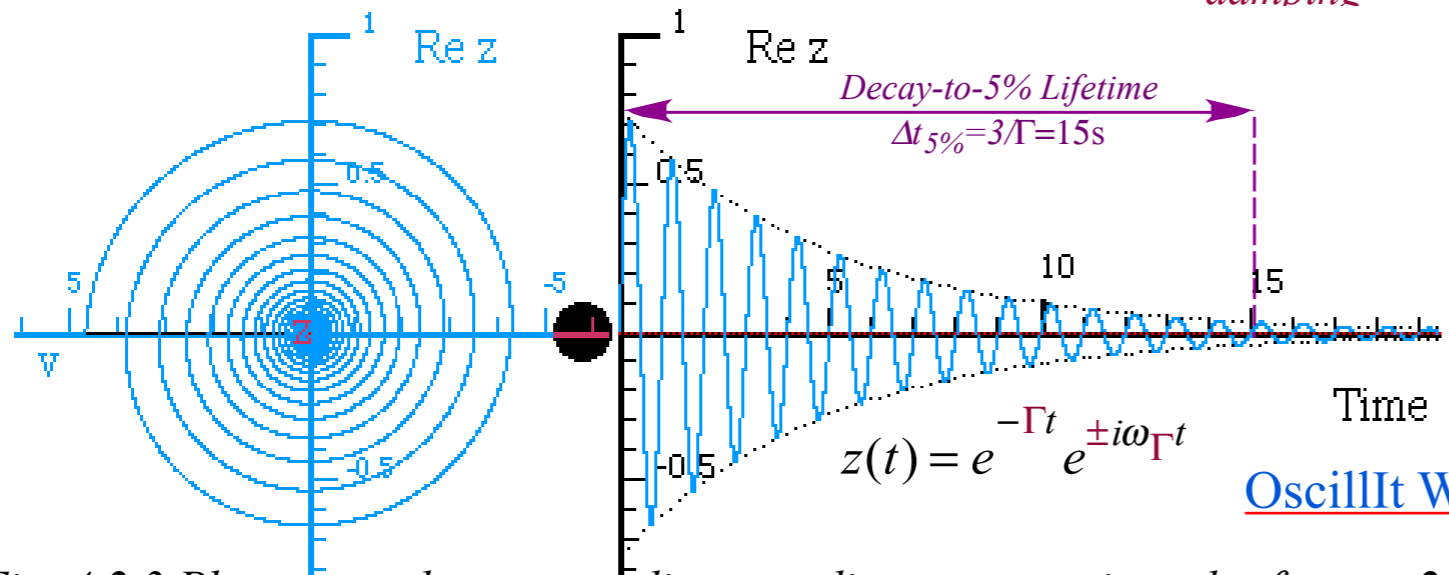
$$F_{restore} = -kz$$

retarded by **frictional damping force**

$$F_{damping} = -b \frac{dz}{dt}$$

Oscillator Figures of Merit:

Time required to reduce amplitude to 5%



Easy-to-recall 5% approximation:

$$e^{-3} \cong 0.05$$

$$t_{5\%} = \frac{3}{\Gamma} = \frac{3}{0.2} = 15$$

[OscillIt Web Simulation](#)

Fig. 4.2.3 Phasor z and corresponding coordinate versus time plot for $\omega_0=2\pi$ and $\Gamma=0.2$

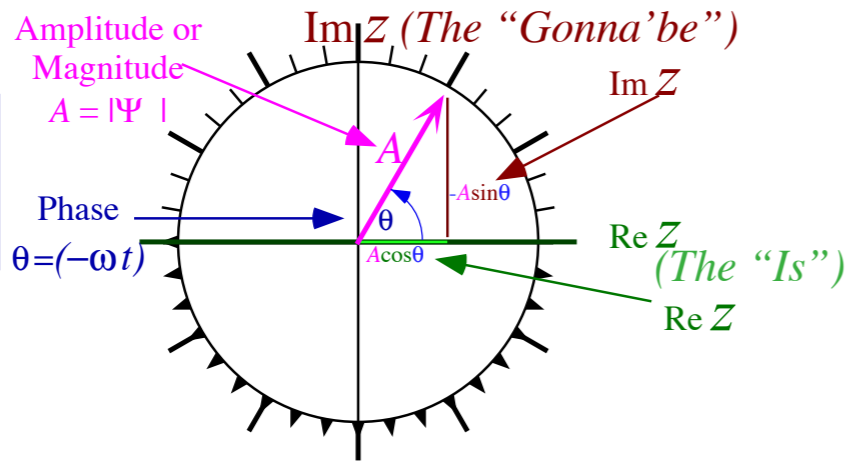
Linear *damped-harmonic oscillator equation of motion.*

$$F_{total}(t) = m \frac{d^2 z}{dt^2} = F_{damping} + F_{restore}$$

$$\frac{d^2 z}{dt^2} = \frac{F_{damping}}{m} + \frac{F_{restore}}{m}$$

$$\frac{d^2 z}{dt^2} + 2\Gamma \frac{dz}{dt} + \omega_0^2 z = 0$$

Phasor Clock
 $Z(t) = Ae^{-i(\omega t)}$
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 Phasor clocks turn clockwise in time for positive ω



Coordinate $z=z(t)$ is the response coordinate for a particle of mass m and charge e

held back by a harmonic (linear) restoring force \longrightarrow

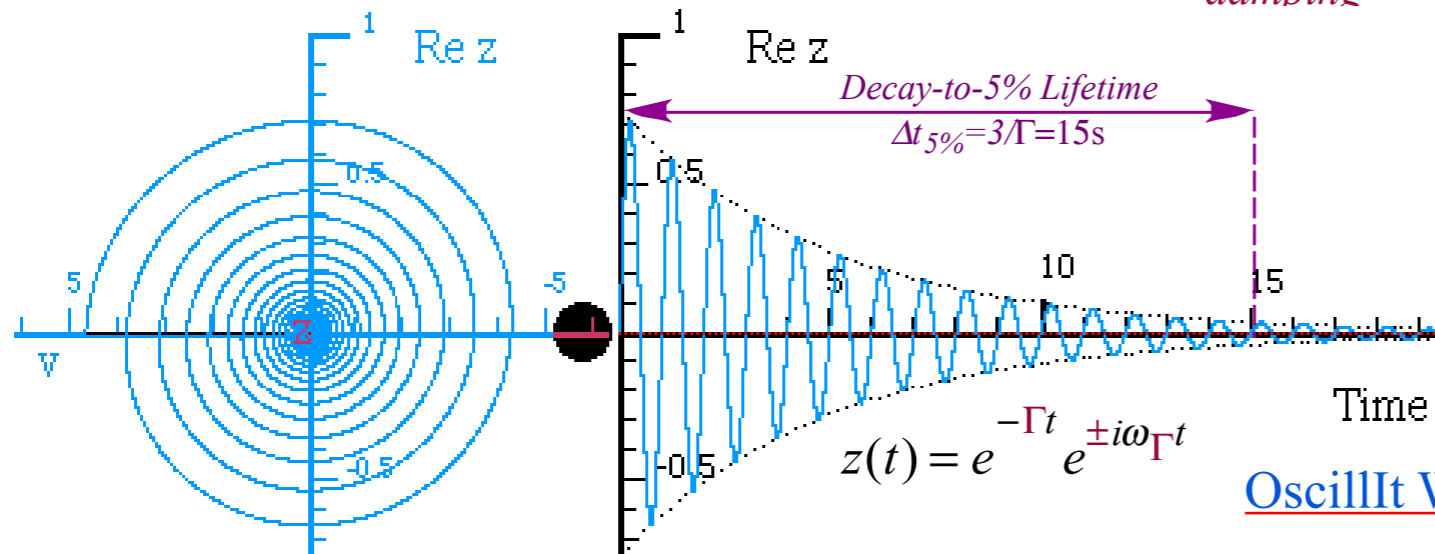
$$F_{restore} = -kz$$

retarded by frictional damping force \longrightarrow

$$F_{damping} = -b \frac{dz}{dt}$$

Oscillator Figures of Merit:

Time required to reduce amplitude to 5% (or 4.321%)



Easy-to-recall 5% approximation: $e^{-3} \cong 0.05$ More precise one: $e^{-\pi} \cong 0.04321$

$$t_{5\%} = \frac{3}{\Gamma} = \frac{3}{0.2} = 15 \quad t_{4.321\%} = \frac{\pi}{\Gamma} = \frac{\pi}{0.2} = 15.708$$

[OscillIt Web Simulation](#)

Fig. 4.2.3 Phasor z and corresponding coordinate versus time plot for $\omega_0=2\pi$ and $\Gamma=0.2$

Linear damped-harmonic oscillator equation of motion.

$$F_{total}(t) = m \frac{d^2 z}{dt^2} = F_{damping} + F_{restore}$$

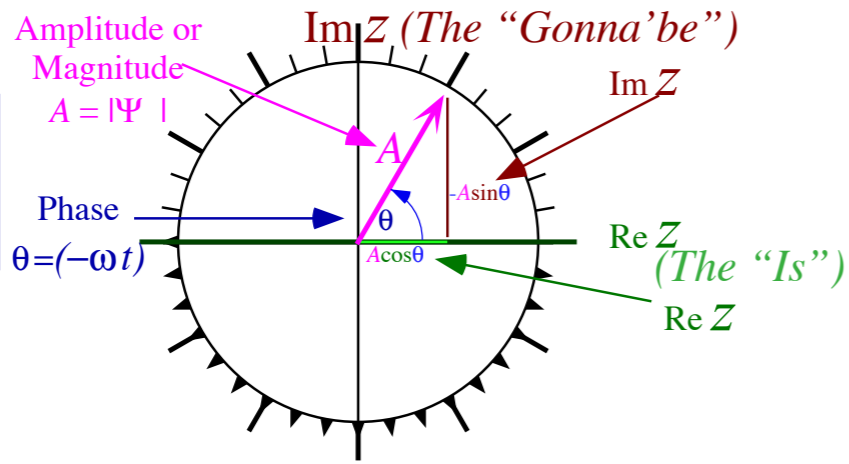
$$\frac{d^2 z}{dt^2} = \frac{F_{damping}}{m} + \frac{F_{restore}}{m}$$

$$\frac{d^2 z}{dt^2} + 2\Gamma \frac{dz}{dt} + \omega_0^2 z = 0$$

[Click to see p.35.](#)

[Click to see p.47.](#)

Phasor Clock
 $Z(t) = Ae^{-i(\omega t)}$
 $= A \cos(\omega t)$
 $-i A \sin(\omega t)$
 Phasor clocks turn clockwise in time for positive ω



Coordinate $z=z(t)$ is the response coordinate for a particle of mass m and charge e

held back by a harmonic (linear) restoring force

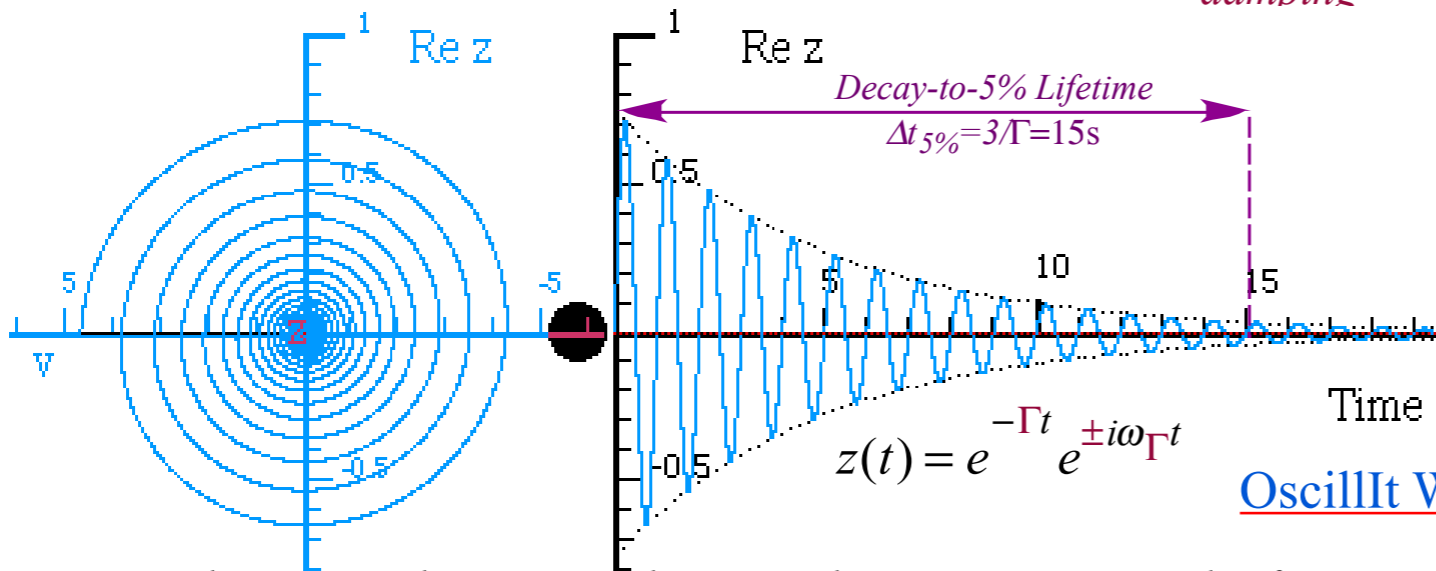
$$F_{restore} = -kz$$

retarded by frictional damping force

$$F_{damping} = -b \frac{dz}{dt}$$

Oscillator Figures of Merit:

Number N of oscillations to reduce amplitude to 5% (or 4.321%)



Easy-to-recall 5% approximation: $e^{-3} \cong 0.05$ More precise one: $e^{-\pi} \cong 0.04321$

$$N_{5\%} = \frac{\omega_{\Gamma} \cdot t_{5\%}}{2\pi} = \frac{3\omega_{\Gamma}}{2\pi\Gamma} \sim \frac{\omega_{\Gamma}}{2\Gamma}$$

$$t_{4.321\%} = \frac{\pi}{\Gamma} = \frac{\pi}{0.2} = 15.708$$

[OscillIt Web Simulation](#)

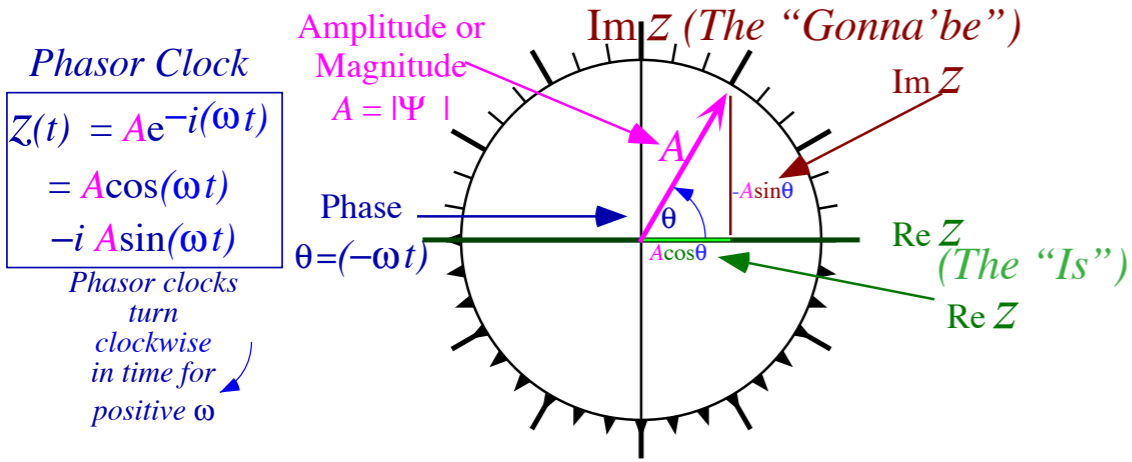
Fig. 4.2.3 Phasor z and corresponding coordinate versus time plot for $\omega_0=2\pi$ and $\Gamma=0.2$

Linear forced-damped-harmonic oscillator equation of motion.

$$F_{total}(t) = m \frac{d^2 z}{dt^2} = F_{damping} + F_{restore} + F_{stimulus}$$

$$\frac{d^2 z}{dt^2} = \frac{F_{damping}}{m} + \frac{F_{restore}}{m} + \frac{F_{stimulus}}{m}$$

Stimulating acceleration $a_{stimulus} = a(t)$ due to stimulating force $F_{stimulus}(t)$ (Typically **E**-field)



$$\frac{d^2 z}{dt^2} + 2\Gamma \frac{dz}{dt} + \omega_0^2 z = a_{stimulus} = \frac{e}{m} E(t)$$

Coordinate $z=z(t)$ is the response coordinate for a particle of mass m and charge e

driven by external **stimulating force** $F_{stimulus}(t) = eE(t)$

held back by a **harmonic (linear) restoring force** $F_{restore} = -kz, (k = \omega_0^2 m),$

retarded by **frictional damping force** $F_{damping} = -b \frac{dz}{dt}, (b = 2\Gamma m)$

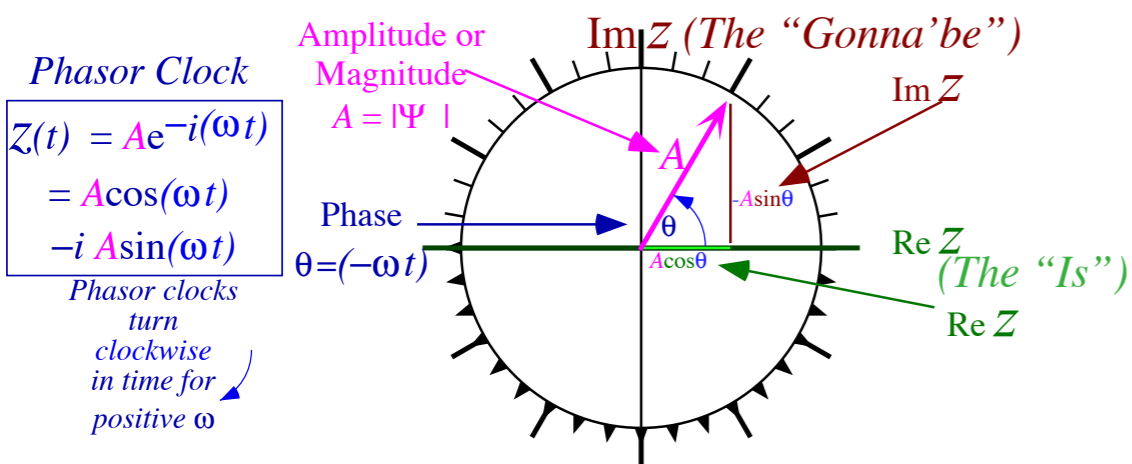
Linear forced-damped-harmonic oscillator equation of motion.

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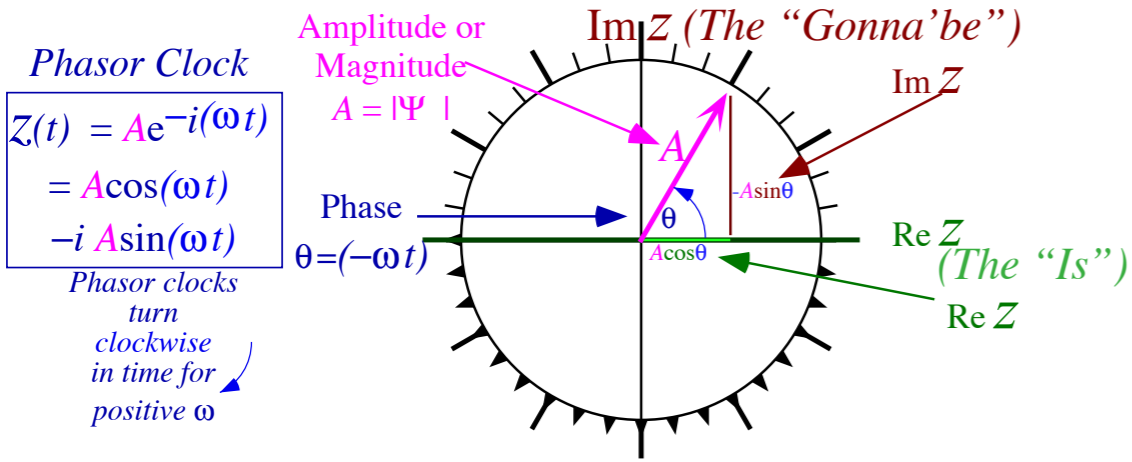
Solving for $z_{stimulus}(t)$ given $a_{stimulus}$:

Linear forced-damped-harmonic oscillator equation of motion.

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$$\frac{d^2 z}{dt^2} + 2\Gamma \frac{dz}{dt} + \omega_0^2 z = a_{stimulus} = \frac{e}{m} E(t)$$

Solving for $z_{stimulus}(t)$ given $a_{stimulus}$:

$$\left(\frac{d^2}{dt^2} + 2\Gamma \frac{d}{dt} + \omega_0^2 \right) z = a_{stimulus}$$

$$z = \frac{1}{\frac{d^2}{dt^2} + 2\Gamma \frac{d}{dt} + \omega_0^2} a_{stimulus}$$

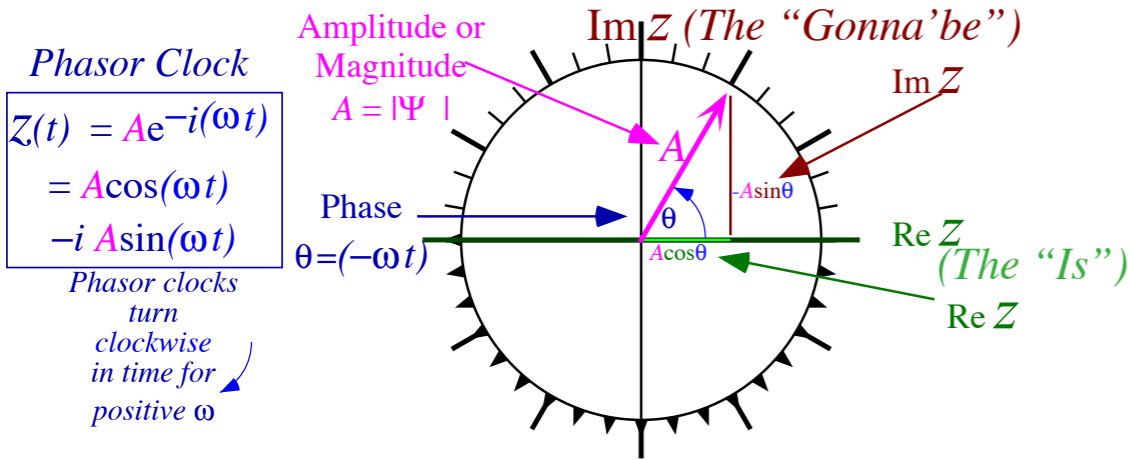
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Pretty crazy?

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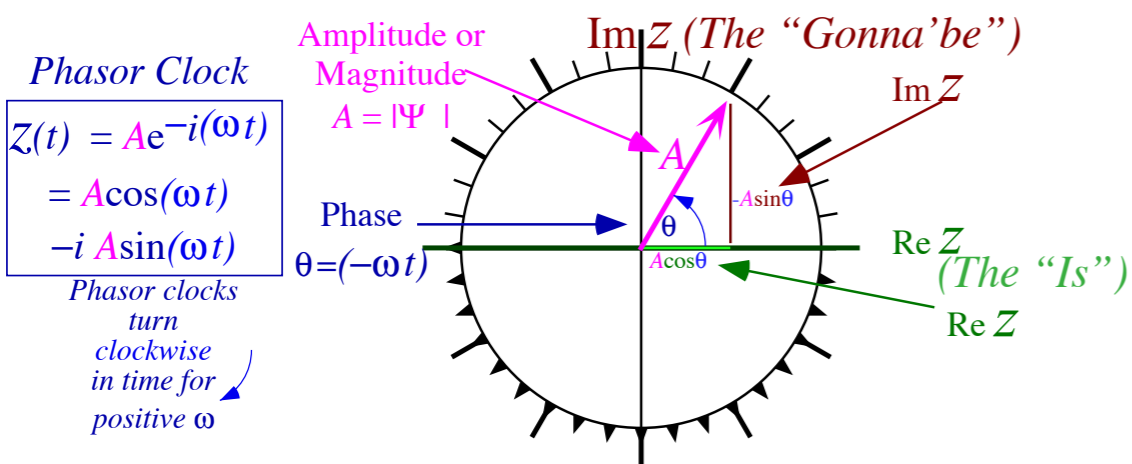
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$$\left(\frac{d^2}{dt^2} + 2\Gamma \frac{d}{dt} + \omega_0^2 \right) z = a_{stimulus}$$

Pretty crazy? But not so crazy if

$$a_{stimulus}(t) = |a_{stimulus}| e^{-i\omega_{stimulus} t} = |a_s| e^{-i\omega_s t}$$

$$z = \frac{1}{\frac{d^2}{dt^2} + 2\Gamma \frac{d}{dt} + \omega_0^2} a_{stimulus}$$

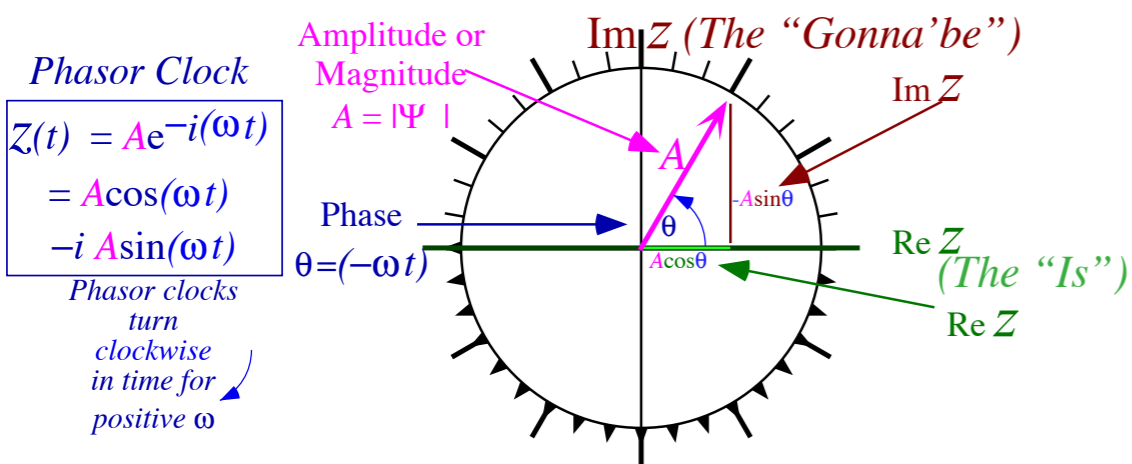
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$$z_s e^{-i\omega_s t} = \frac{1}{\omega_0^2 - \omega_s^2 - i2\Gamma\omega_s} a_s e^{-i\omega_s t}$$

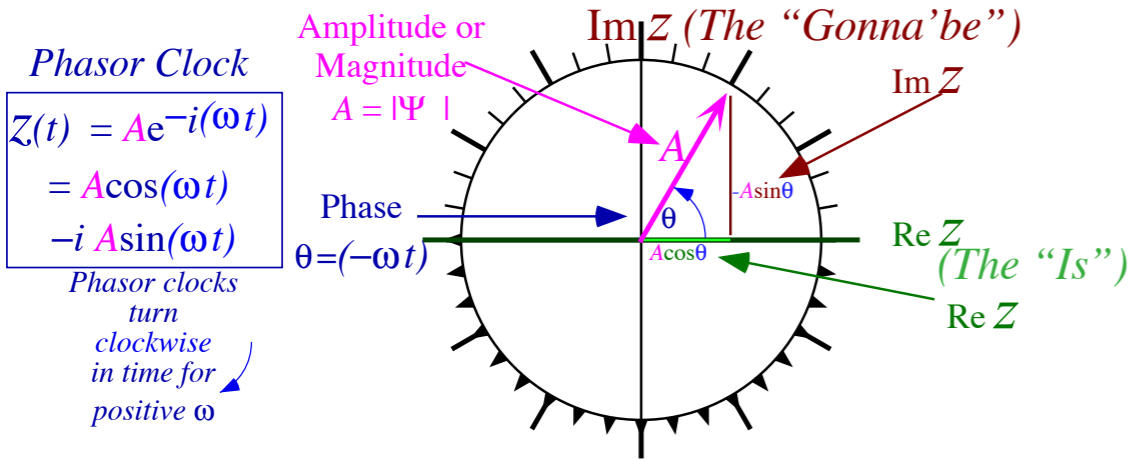
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$$z_{stimulus} = \frac{1}{-\omega_s^2 - i2\Gamma\omega_s + \omega_0^2} a_s e^{-i\omega_s t}$$

$$z_s e^{-i\omega_s t} = \frac{1}{\omega_0^2 - \omega_s^2 - i2\Gamma\omega_s} a_s e^{-i\omega_s t}$$

$$z_s = G_{\omega_0}(\omega_s) \cdot a_s$$

Linear forced-damped-harmonic oscillator equation of motion.

$$F_{total}(t) = m \frac{d^2 z}{dt^2} = F_{damping} + F_{restore} + F_{stimulus}$$

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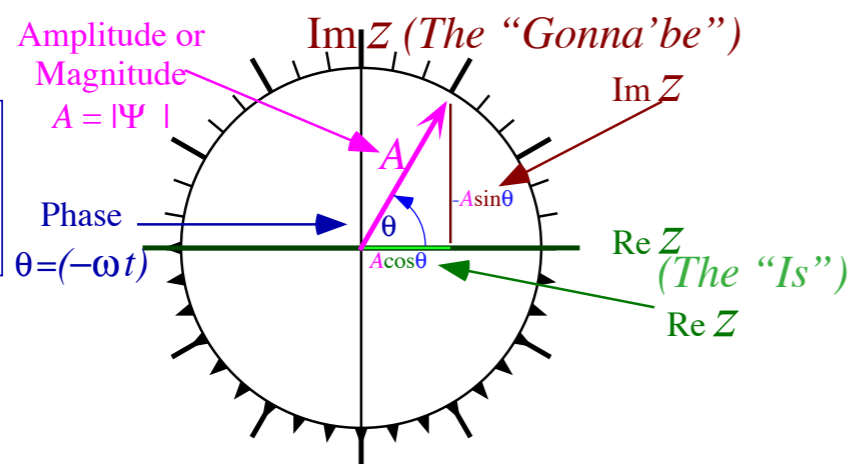
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Phasor Clock

$$z(t) = A e^{-i(\omega t)} \\ = A \cos(\omega t) - i A \sin(\omega t)$$

Phasor clocks turn clockwise in time for positive ω



Solving for $z_{stimulus}(t)$ given $a_{stimulus}$:

$$\left(\frac{d^2}{dt^2} + 2\Gamma \frac{d}{dt} + \omega_0^2 \right) z = a_{stimulus}$$

Pretty crazy? But not so crazy if

$$a_{stimulus}(t) = |a_{stimulus}| e^{-i\omega_{stimulus} t} = |a_s| e^{-i\omega_s t}$$

$$z_{stimulus} = \frac{1}{-\omega_s^2 - i2\Gamma\omega_s + \omega_0^2} a_s e^{-i\omega_s t}$$

$$z_s e^{-i\omega_s t} = \frac{1}{\omega_0^2 - \omega_s^2 - i2\Gamma\omega_s} a_s e^{-i\omega_s t}$$

$$z_s = G_{\omega_0}(\omega_s) \cdot a_s$$



George Green (14 July 1793 – 31 May 1841)

Green's Function for the F-D-H Oscillator (FDHO)

Green's Function for the **FDHO** (**F**orced-**D**amped-**H**armonic Oscillator)

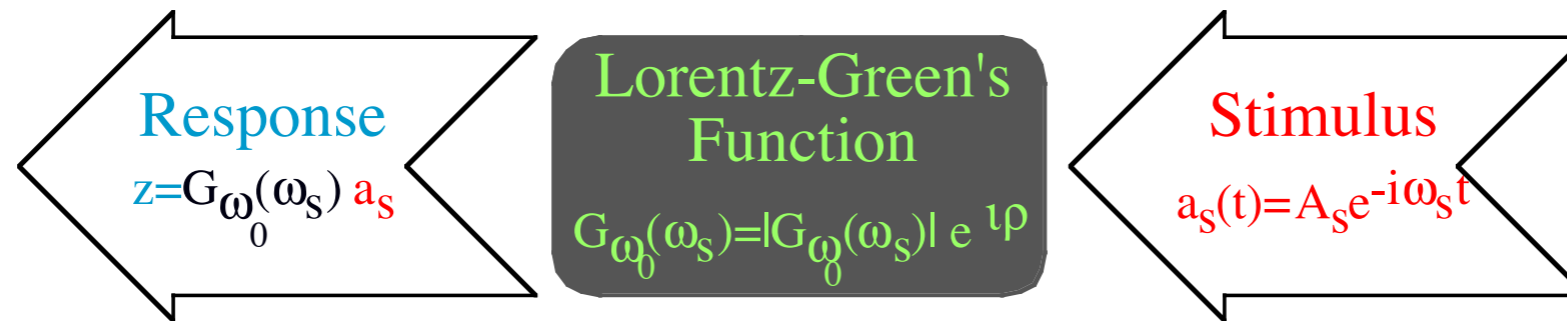


Fig. 4.2.4 Black-box diagram of oscillator response to monochromatic stimulus

$$G_{\omega_0}(\omega_s) = \frac{1}{\omega_0^2 - \omega_s^2 - i2\Gamma\omega_s} = \text{Re } G_{\omega_0}(\omega_s) + i \text{Im } G_{\omega_0}(\omega_s)$$

Real and imaginary parts of the *rectangular form* of G :

Hendrik A. Lorentz



July 18, 1853. - February 4, 1928

Green's Function for the **FDHO** (**F**orced-**D**amped-**H**armonic Oscillator)

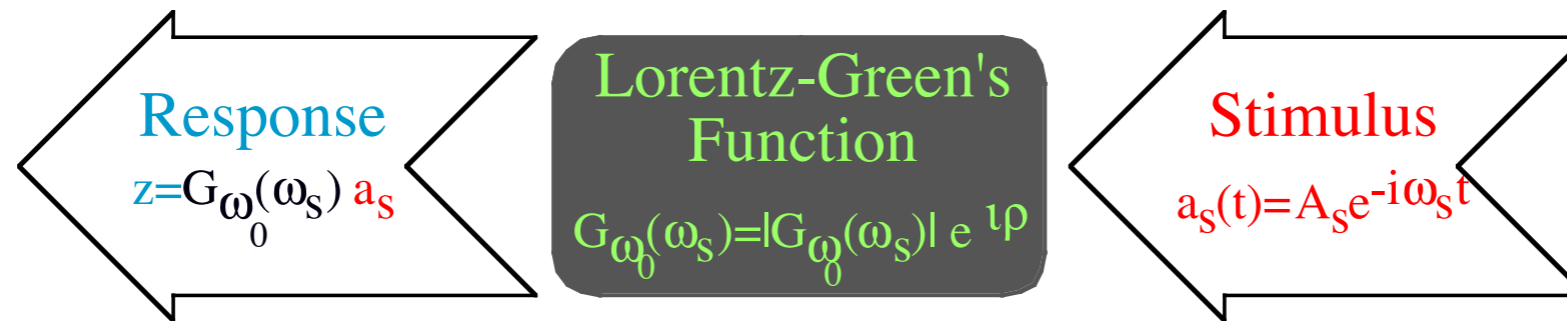


Fig. 4.2.4 Black-box diagram of oscillator response to monochromatic stimulus

$$G_{\omega_0}(\omega_s) = \frac{1}{\omega_0^2 - \omega_s^2 - i2\Gamma\omega_s} = \text{Re } G_{\omega_0}(\omega_s) + i \text{Im } G_{\omega_0}(\omega_s)$$

Real and imaginary parts of the *rectangular form* of G : $\frac{1}{x-iy} = \frac{1}{x-iy} \frac{x+iy}{x+iy} = \frac{x+iy}{x^2+y^2}$

Green's Function for the **FDHO** (**F**orced-**D**amped-**H**armonic Oscillator)

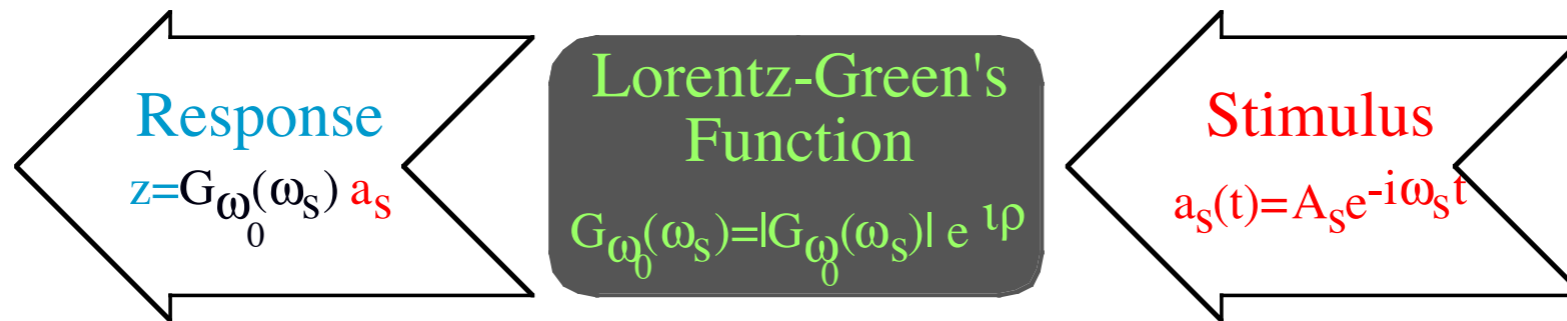


Fig. 4.2.4 Black-box diagram of oscillator response to monochromatic stimulus

$$G_{\omega_0}(\omega_s) = \frac{1}{\omega_0^2 - \omega_s^2 - i2\Gamma\omega_s} = \text{Re } G_{\omega_0}(\omega_s) + i \text{Im } G_{\omega_0}(\omega_s)$$

Real and imaginary parts of the *rectangular form* of G : $\frac{1}{x-iy} = \frac{1}{x-iy} \frac{x+iy}{x+iy} = \frac{x+iy}{x^2+y^2} = \frac{x}{x^2+y^2} + i \frac{y}{x^2+y^2}$

$$\text{Re } G_{\omega_0}(\omega_s) = \frac{\omega_0^2 - \omega_s^2}{(\omega_0^2 - \omega_s^2)^2 + (2\Gamma\omega_s)^2}$$

$$\text{Im } G_{\omega_0}(\omega_s) = \frac{2\Gamma\omega_s}{(\omega_0^2 - \omega_s^2)^2 + (2\Gamma\omega_s)^2}$$

Green's Function for the **FDHO** (**F**orced-**D**amped-**H**armonic Oscillator)

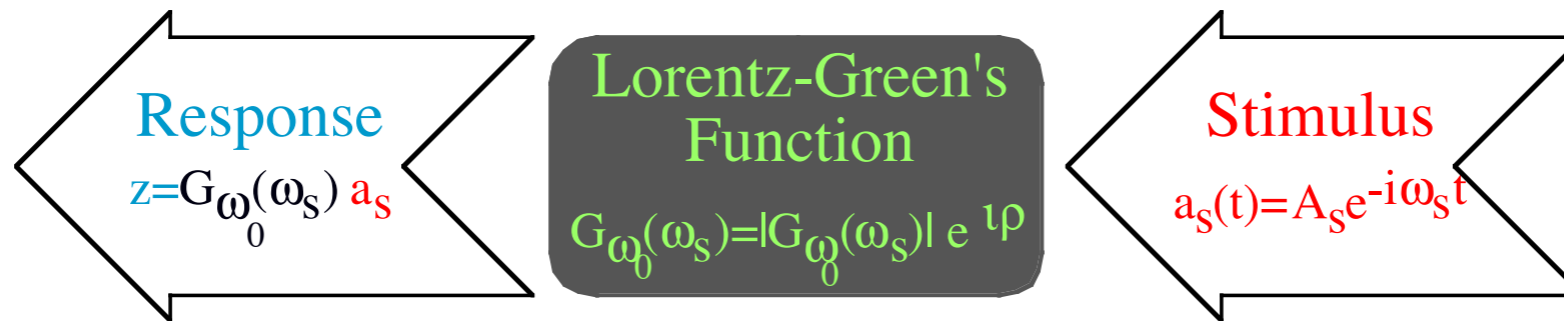


Fig. 4.2.4 Black-box diagram of oscillator response to monochromatic stimulus

$$G_{\omega_0}(\omega_s) = \frac{1}{\omega_0^2 - \omega_s^2 - i2\Gamma\omega_s} = \text{Re } G_{\omega_0}(\omega_s) + i \text{Im } G_{\omega_0}(\omega_s) = |G_{\omega_0}(\omega_s)| e^{i\rho}$$

Real and imaginary parts of the *rectangular form* of G :

$$\text{Re } G_{\omega_0}(\omega_s) = \frac{\omega_0^2 - \omega_s^2}{(\omega_0^2 - \omega_s^2)^2 + (2\Gamma\omega_s)^2}$$

$$\text{Im } G_{\omega_0}(\omega_s) = \frac{2\Gamma\omega_s}{(\omega_0^2 - \omega_s^2)^2 + (2\Gamma\omega_s)^2}$$

Magnitude $|G_{\omega_0}(\omega_s)|$ and polar angle ρ of the *polar form* of G :

$$|G_{\omega_0}(\omega_s)| = \frac{1}{\sqrt{(\omega_0^2 - \omega_s^2)^2 + (2\Gamma\omega_s)^2}}$$

$$\rho = \tan^{-1} \left(\frac{2\Gamma\omega_s}{\omega_0^2 - \omega_s^2} \right)$$

Green's Function for the FDHO (Forced-Damped-Harmonic Oscillator)

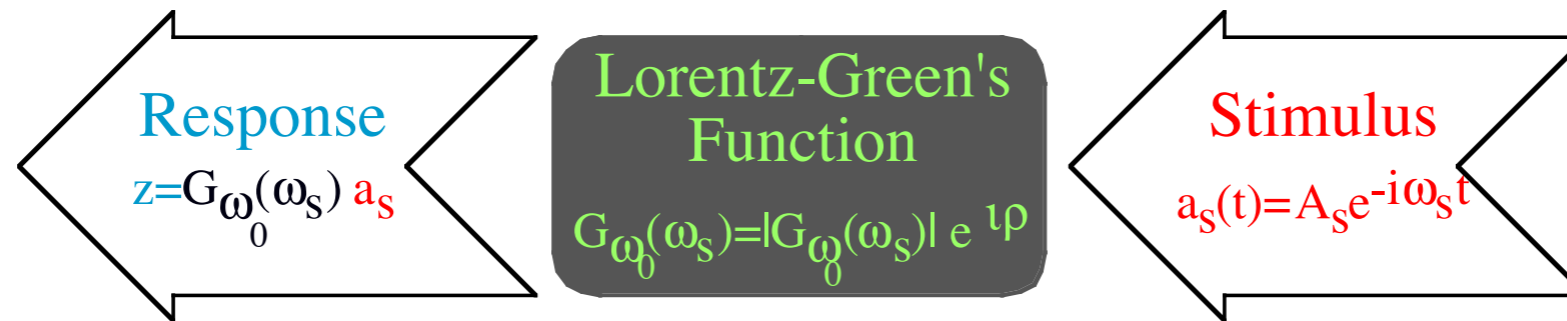


Fig. 4.2.4 Black-box diagram of oscillator response to monochromatic stimulus

$$G_{\omega_0}(\omega_s) = \frac{1}{\omega_0^2 - \omega_s^2 - i2\Gamma\omega_s} = \text{Re } G_{\omega_0}(\omega_s) + i \text{Im } G_{\omega_0}(\omega_s) = |G_{\omega_0}(\omega_s)| e^{i\rho}$$

Real and imaginary parts of the *rectangular form* of G :

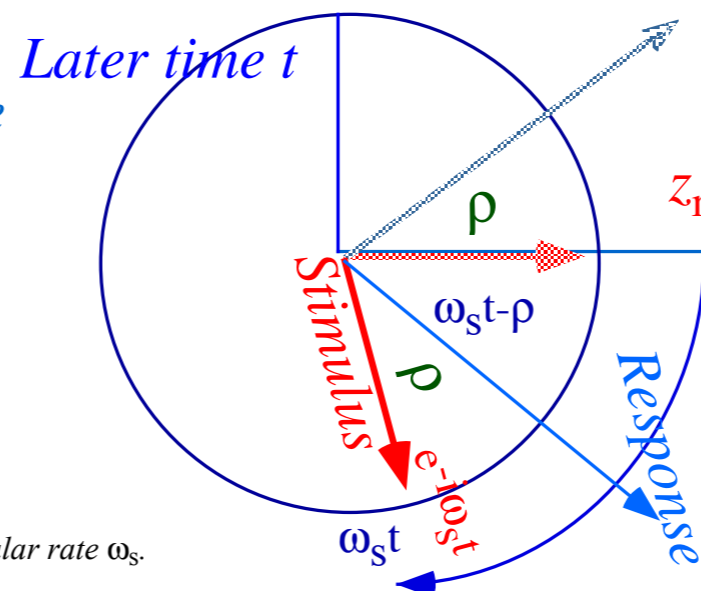
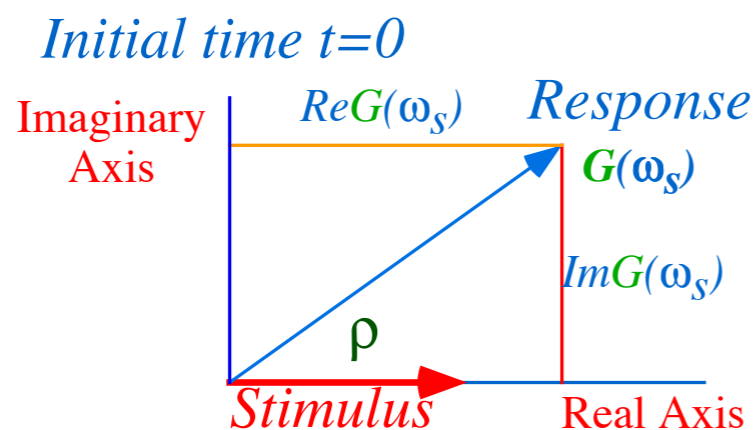
$$\text{Re } G_{\omega_0}(\omega_s) = \frac{\omega_0^2 - \omega_s^2}{(\omega_0^2 - \omega_s^2)^2 + (2\Gamma\omega_s)^2}$$

$$\text{Im } G_{\omega_0}(\omega_s) = \frac{2\Gamma\omega_s}{(\omega_0^2 - \omega_s^2)^2 + (2\Gamma\omega_s)^2}$$

Magnitude $|G_{\omega_0}(\omega_s)|$ and *polar angle* ρ of the *polar form* of G :

$$|G_{\omega_0}(\omega_s)| = \frac{1}{\sqrt{(\omega_0^2 - \omega_s^2)^2 + (2\Gamma\omega_s)^2}}$$

$$\rho = \tan^{-1}\left(\frac{2\Gamma\omega_s}{\omega_0^2 - \omega_s^2}\right)$$



polar angle ρ is the *phase lag angle* ρ

$$z_{\text{response}}(t) = |G_{\omega_0}(\omega_s)| a(0) e^{-i(\omega_s t - \rho)}$$

Fig. 4.2.5 Oscillator response and stimulus phasors rotate rigidly at angular rate ω_s .

Green's Function for the FDHO (Forced-Damped-Harmonic Oscillator)

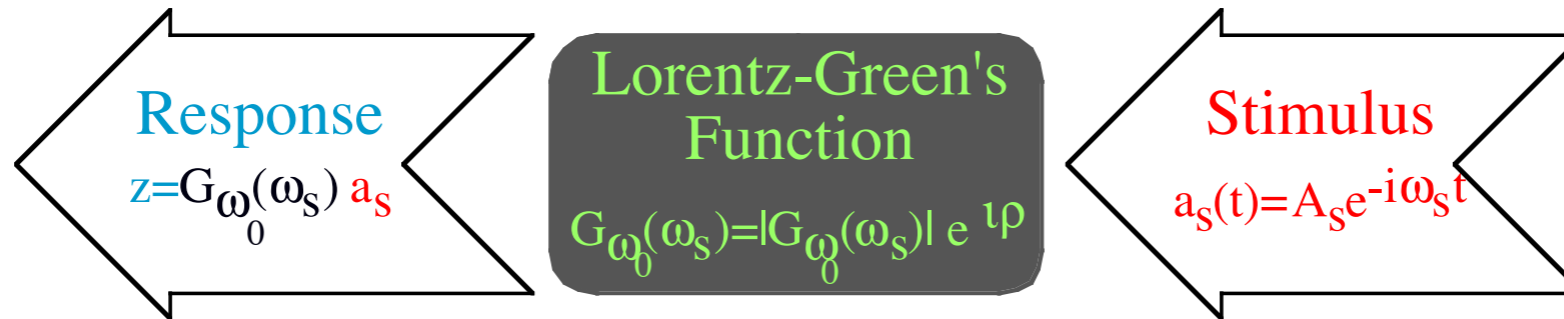


Fig. 4.2.4 Black-box diagram of oscillator response to monochromatic stimulus

$$G_{\omega_0}(\omega_s) = \frac{1}{\omega_0^2 - \omega_s^2 - i2\Gamma\omega_s} = \text{Re } G_{\omega_0}(\omega_s) + i \text{Im } G_{\omega_0}(\omega_s) = |G_{\omega_0}(\omega_s)| e^{i\rho}$$

Real and imaginary parts of the rectangular form of G :

$$\text{Re } G_{\omega_0}(\omega_s) = \frac{\omega_0^2 - \omega_s^2}{(\omega_0^2 - \omega_s^2)^2 + (2\Gamma\omega_s)^2}$$

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Magnitude $|G_{\omega_0}(\omega_s)|$ and *polar angle* ρ of the *polar form* of G :

$$|G_{\omega_0}(\omega_s)| = \frac{1}{\sqrt{(\omega_0^2 - \omega_s^2)^2 + (2\Gamma\omega_s)^2}}$$

$$\rho = \tan^{-1}\left(\frac{2\Gamma\omega_s}{\omega_0^2 - \omega_s^2}\right)$$

polar angle ρ is the *phase lag angle* ρ

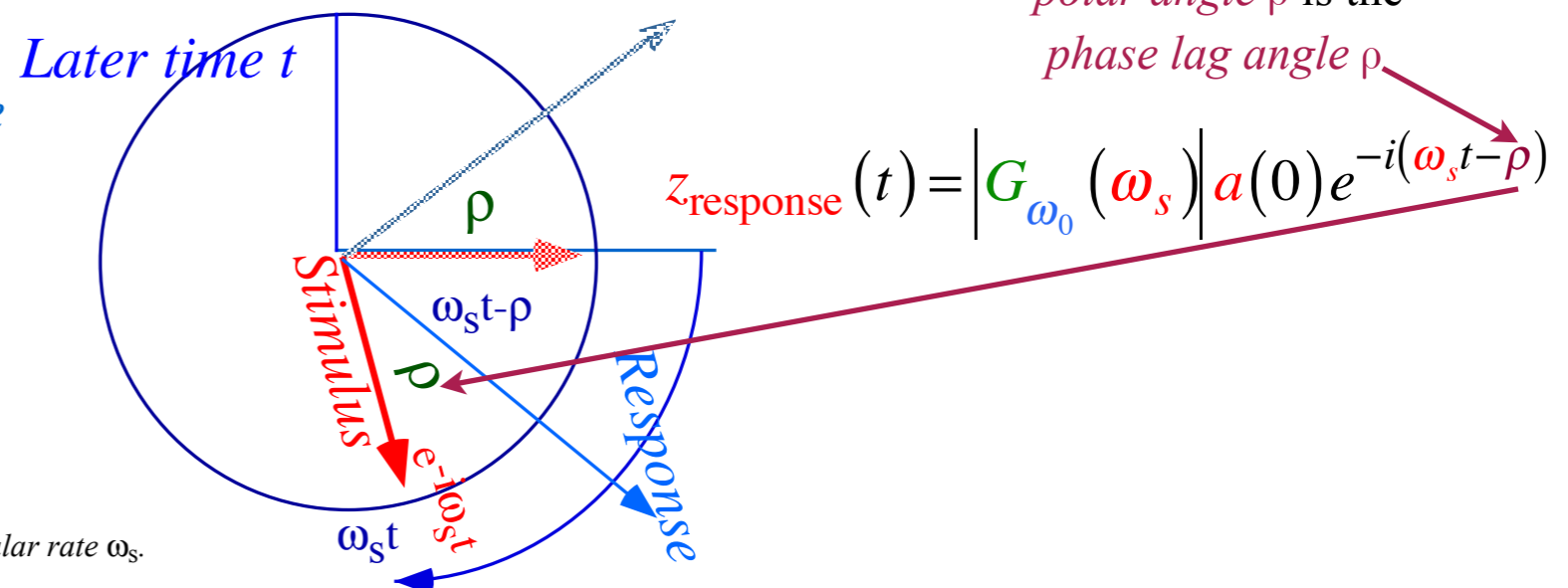
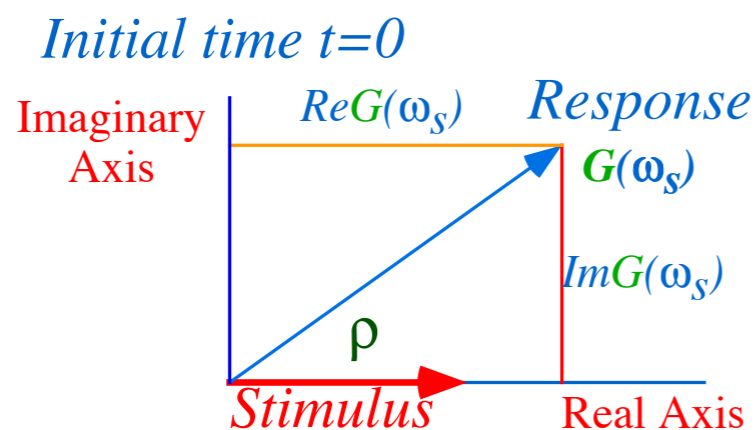
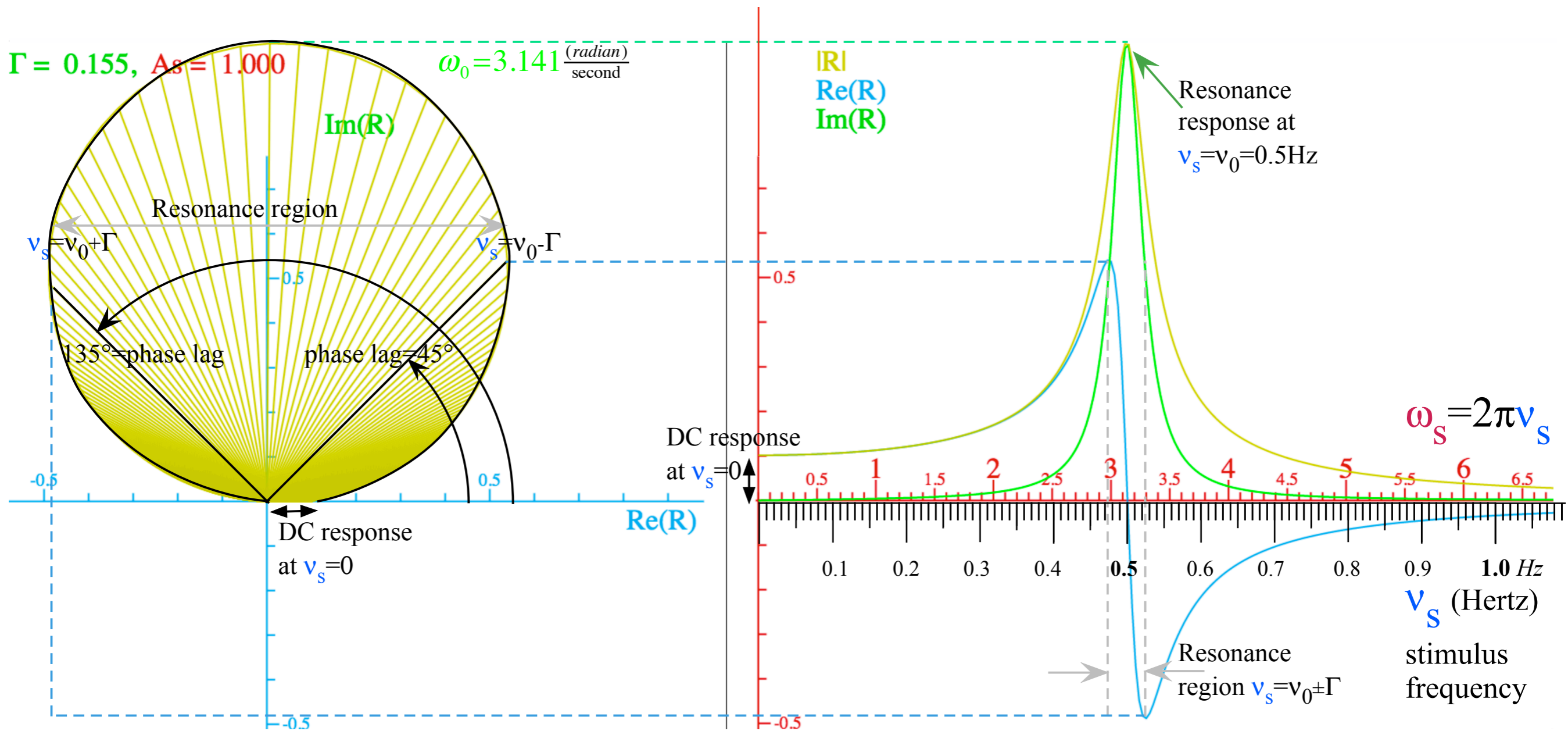


Fig. 4.2.5 Oscillator response and stimulus phasors rotate rigidly at angular rate ω_s .

Lorentz-Green's function for $\nu_0 = 0.5 \text{ Hz}$ or $\omega_0 = \pi \frac{\text{(radian)}}{\text{second}}$



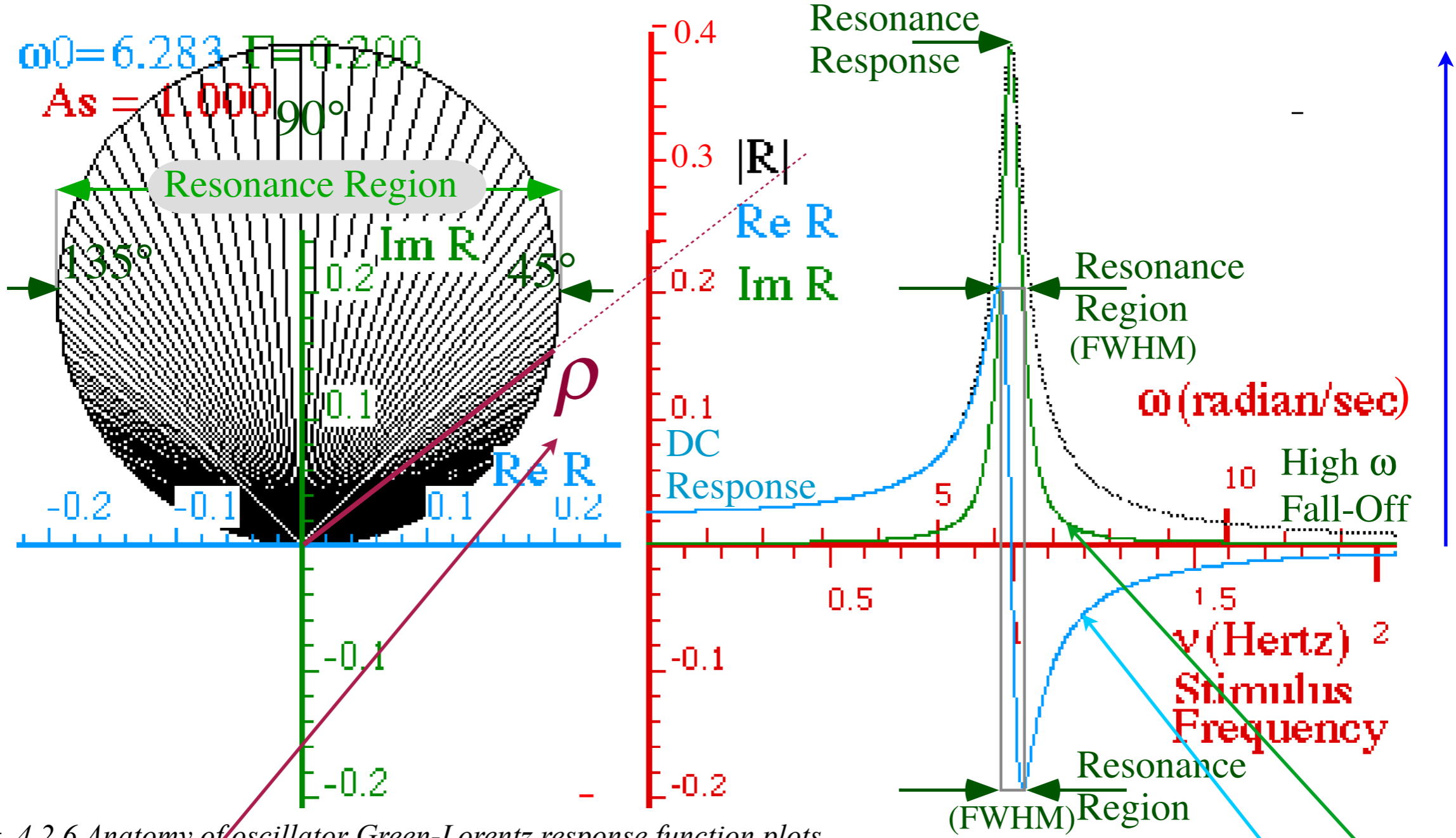


Fig. 4.2.6 Anatomy of oscillator Green-Lorentz response function plots

Phase lag angle

$$\rho = \tan^{-1} \left(\frac{2\Gamma\omega_s}{\omega_0^2 - \omega_s^2} \right)$$

$$\text{Re } G_{\omega_0}(\omega_s) = \frac{\omega_0^2 - \omega_s^2}{(\omega_0^2 - \omega_s^2)^2 + (2\Gamma\omega_s)^2}$$

$$\text{Im } G_{\omega_0}(\omega_s) = \frac{2\Gamma\omega_s}{(\omega_0^2 - \omega_s^2)^2 + (2\Gamma\omega_s)^2}$$

Real part

Imaginary part

$$Q \text{ factor} = \frac{\nu_0}{2\Gamma} = \frac{q}{2\pi}$$

Amplitude Amplification Factor

$$AAF = \frac{\text{Resonant response}}{\text{DC response}} = \frac{|G_{\omega_0}(\omega_s = \omega_0)|}{|G_{\omega_0}(0)|} = \frac{1/(2\Gamma\omega_0)}{1/\omega_0^2} = \frac{\omega_0}{2\Gamma} \equiv q \quad (\text{angular quality factor})$$

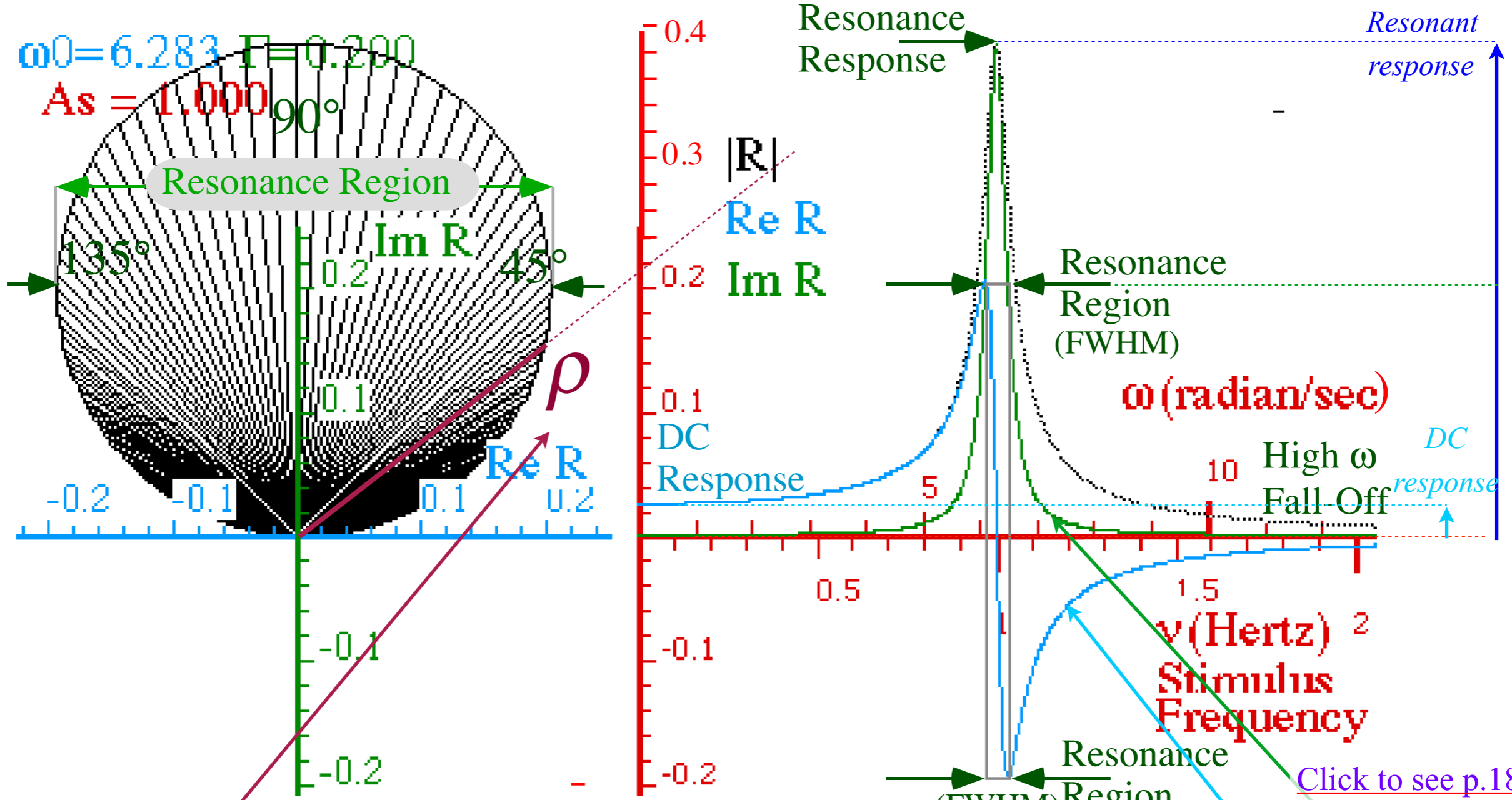


Fig. 4.2.6 Anatomy of oscillator Green-Lorentz response function plots

Phase lag angle

$$\rho = \tan^{-1} \left(\frac{2\Gamma\omega_s}{\omega_0^2 - \omega_s^2} \right)$$

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$$\text{Im } G_{\omega_0}(\omega_s) = \frac{2\Gamma\omega_s}{(\omega_0^2 - \omega_s^2)^2 + (2\Gamma\omega_s)^2}$$

[Click to see p.18.](#)

[Click to see p.47.](#)

Real part

Imaginary part

$$Q \text{ factor} = \frac{\nu_0}{2\Gamma} = \frac{q}{2\pi}$$

Amplitude Amplification Factor

$$AAF = \frac{\text{Resonant response}}{\text{DC response}} = \frac{|G_{\omega_0}(\omega_s = \omega_0)|}{|G_{\omega_0}(0)|} = \frac{1/(2\Gamma\omega_0)}{1/\omega_0^2} = \frac{\omega_0}{2\Gamma} \equiv q \quad (\text{angular quality factor})$$

OscillIt Web Simulation:
Lorentz Function w/
Gamma=0.2

OscillIt Web Simulation:
Lorentz Function w/
Gamma=0.1

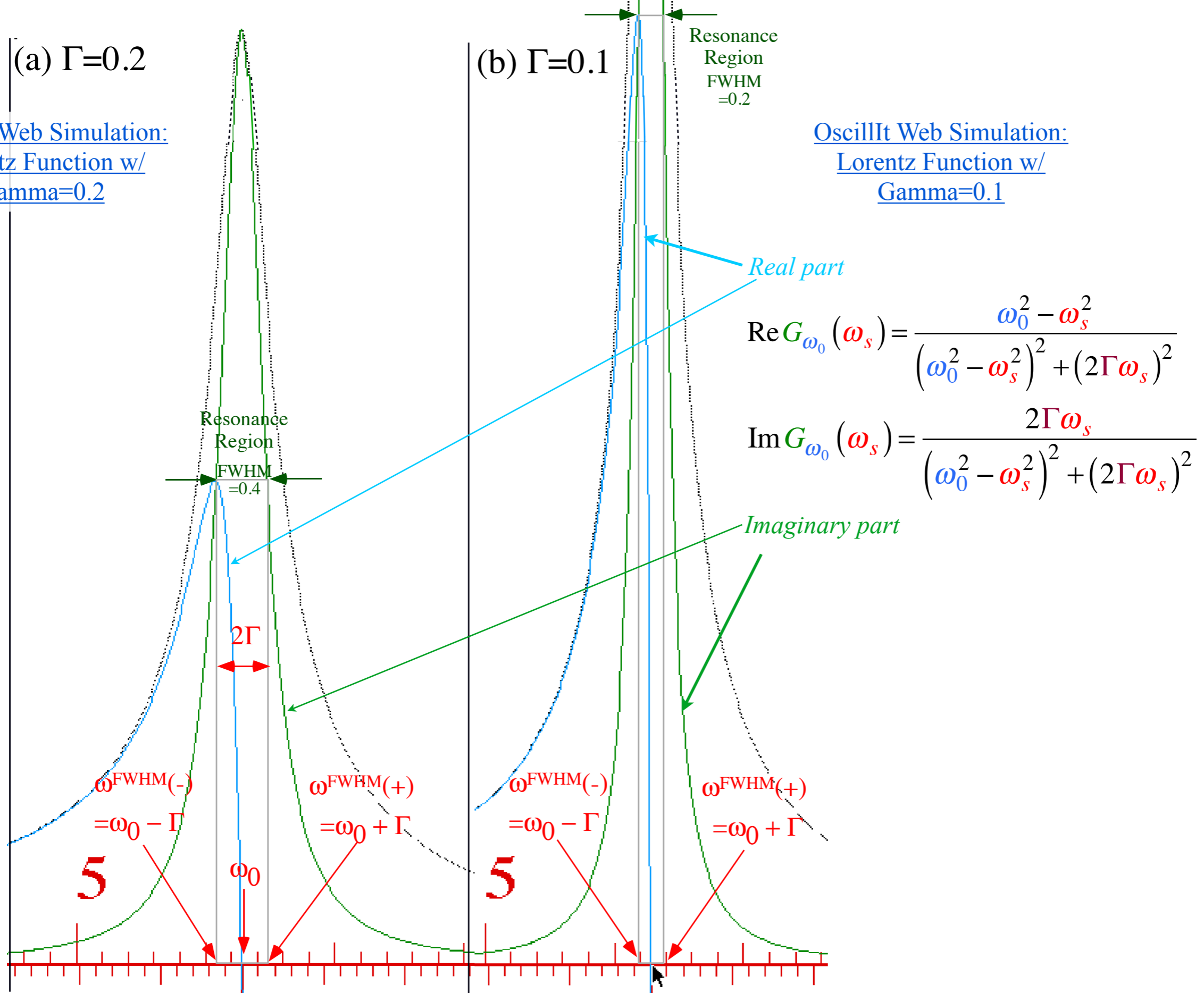


Fig. 4.2.7 Comparing Lorentz-Green resonance region for (a) $\Gamma=0.2$ and (b) $\Gamma=0.1$.

Maximum and minimum points of $\text{Re}G(\omega)$ and inflection points of $\text{Im}G(\omega)$ are near region boundaries $\omega^{\text{FWHM}(\pm)} = \omega_0 \pm \Gamma$.

Complete *Green's Solution* for the *FDHO* (*Forced-Damped-Harmonic Oscillator*)

$$z(t) = z_{\text{transient}}(t) + z_{\text{response}}(t) \equiv z_{\text{decaying}}(t) + z_{\text{steady state}}(t)$$

$$= Ae^{-\Gamma t} e^{-i\omega_{\Gamma} t} + G_{\omega_0}(\omega_s) a(0) e^{-i\omega_s t}$$

$$= \underbrace{Ae^{-\Gamma t} e^{-i\omega_{\Gamma} t}}_{\text{Known as "homogeneous" solution (no force)}} + \underbrace{\left| G_{\omega_0}(\omega_s) \right| a(0) e^{-i(\omega_s t - \rho)}}_{\text{Known as "inhomogeneous" solution}}$$

Let's you set initial values or boundary conditions

Not function of initial values. Marches to stimulus only.

Complete *Green's Solution* for the *FDHO* (*Forced-Damped-Harmonic Oscillator*)

$$z(t) = z_{\text{transient}}(t) + z_{\text{response}}(t) \equiv z_{\text{decaying}}(t) + z_{\text{steady state}}(t)$$

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Let's you set initial values or boundary conditions

Known as *Transient* solution since it dies-off as time advances past initial conditions

Not function of initial values. Marches to stimulus only.

Known as *Steady State* solution since it is present as long as stimulus is.

Complete *Green's Solution* for the *FDHO* (*Forced-Damped-Harmonic Oscillator*)

$$z(t) = z_{\text{transient}}(t) + z_{\text{response}}(t) \equiv z_{\text{decaying}}(t) + z_{\text{steady state}}(t)$$

$$= Ae^{-\Gamma t} e^{-i\omega_{\Gamma} t} + G_{\omega_0}(\omega_s) a(0) e^{-i\omega_s t}$$

$$= Ae^{-\Gamma t} e^{-i\omega_{\Gamma} t} + \left| G_{\omega_0}(\omega_s) \right| a(0) e^{-i(\omega_s t - \rho)}$$

Known as “homogeneous” solution (no force)
Let's you set initial values or boundary conditions

Known as *Transient* solution since it dies-off as time advances past initial conditions

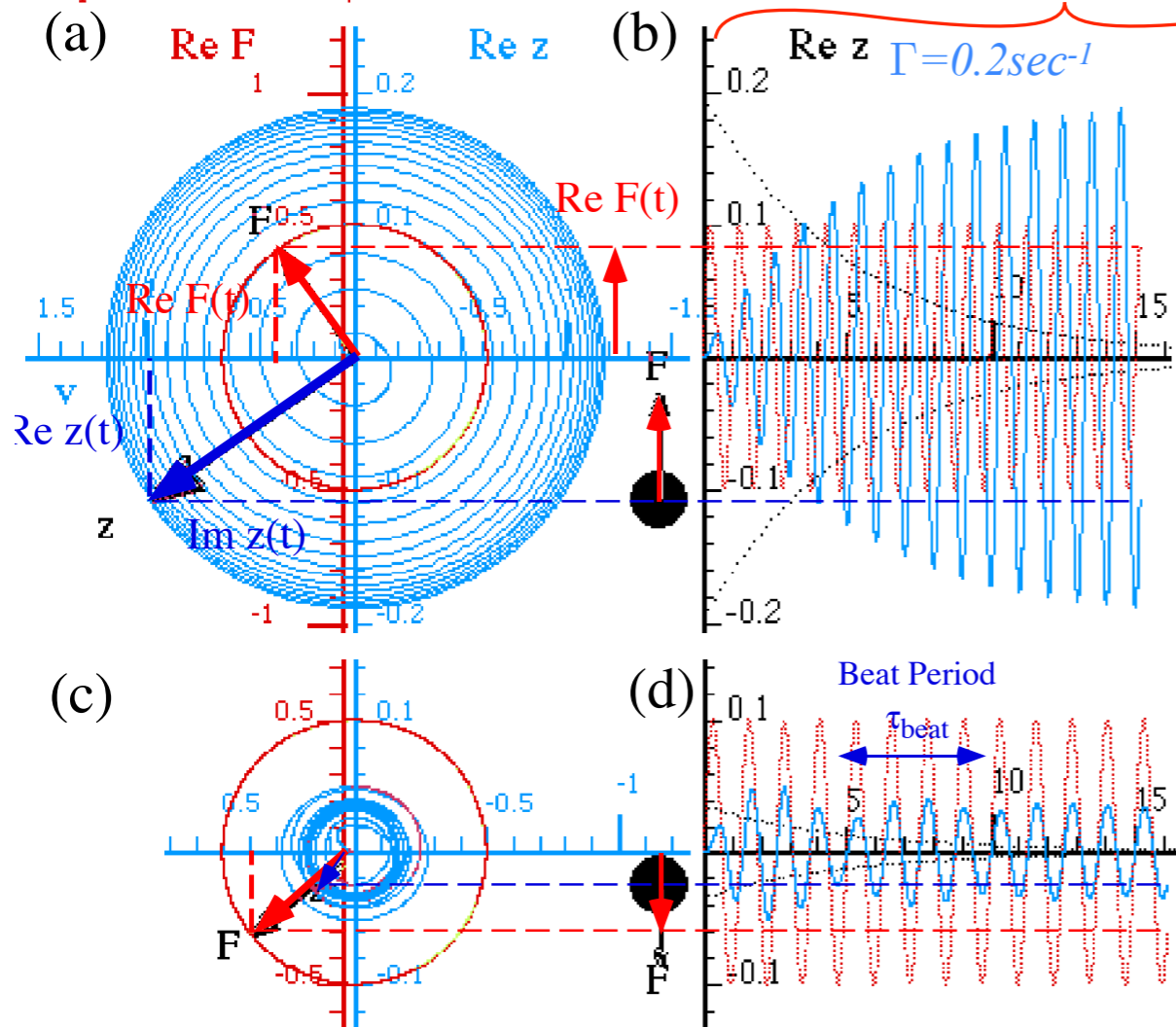
Known as “inhomogeneous” solution
Not function of initial values. Marches to stimulus only.

Known as *Steady State* solution since it is present as long as stimulus is.

Stimulus: $A_s = 0.5000$ $\omega = 6.2832$
Response: $R = 0.1989$ $\rho = 1.5708$

About $t = 3/\Gamma = 15 \text{ sec}$

About $t = \text{forever}$



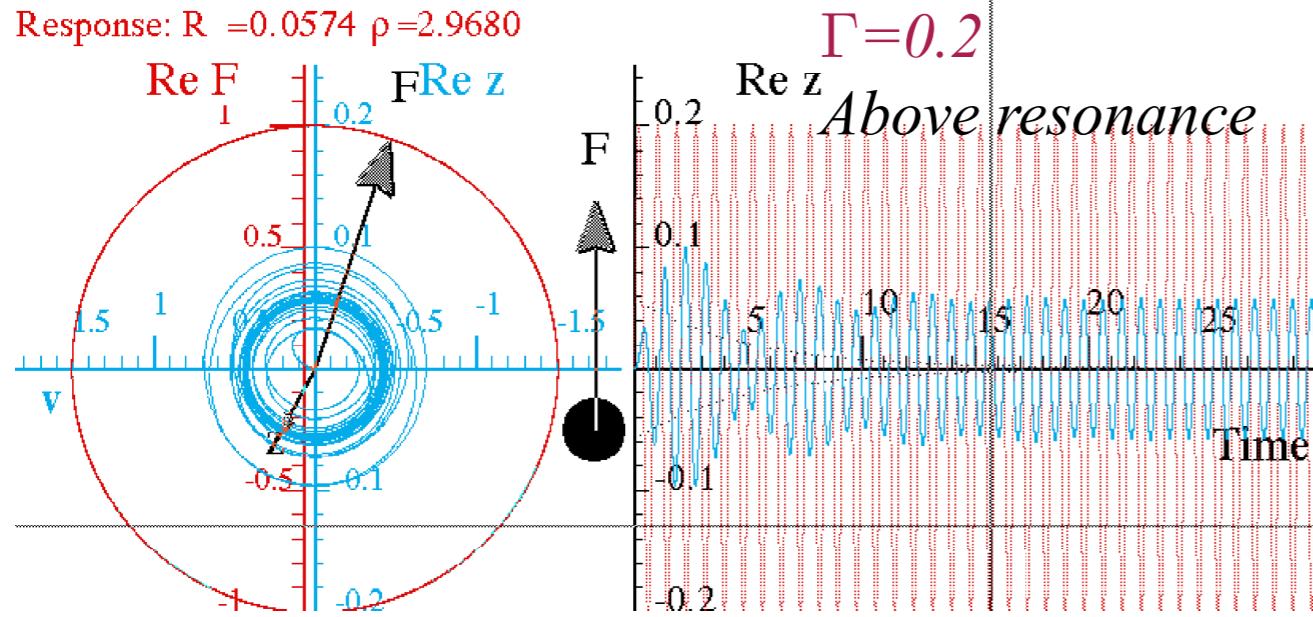
OscillIt (On Resonance) Simulation

Fig. 4.2.8 On Resonance (a) Response z -phasor lags $\rho = 90^\circ$ behind stimulus F -phasor. ($\omega_s = \omega_0 = 2\pi$, $\omega_0 = 2\pi$, and $\Gamma = 0.2$). (b) Time plots of $\text{Re } z(t)$ and $\text{Re } F(t)$

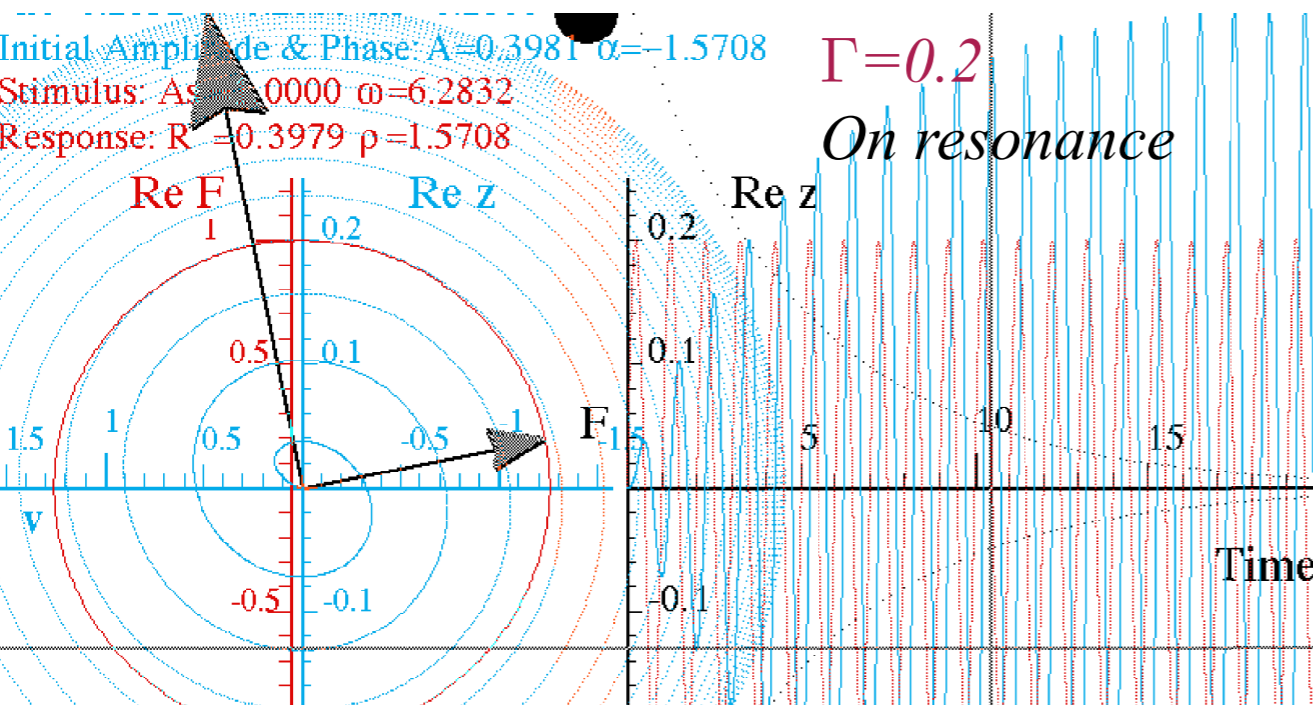
Fig. 4.2.8 Below Resonance (c) Response z -phasor lags $\rho = 8.05^\circ$ behind stimulus F -phasor. ($\omega_s = 5.03$, $\omega_0 = 2\pi$, and $\Gamma = 0.2$). (d) Time plots of $\text{Re } z(t)$ and $\text{Re } F(t)$. Beats are barely visible.

OscillIt (Way Below Resonance) Simulation

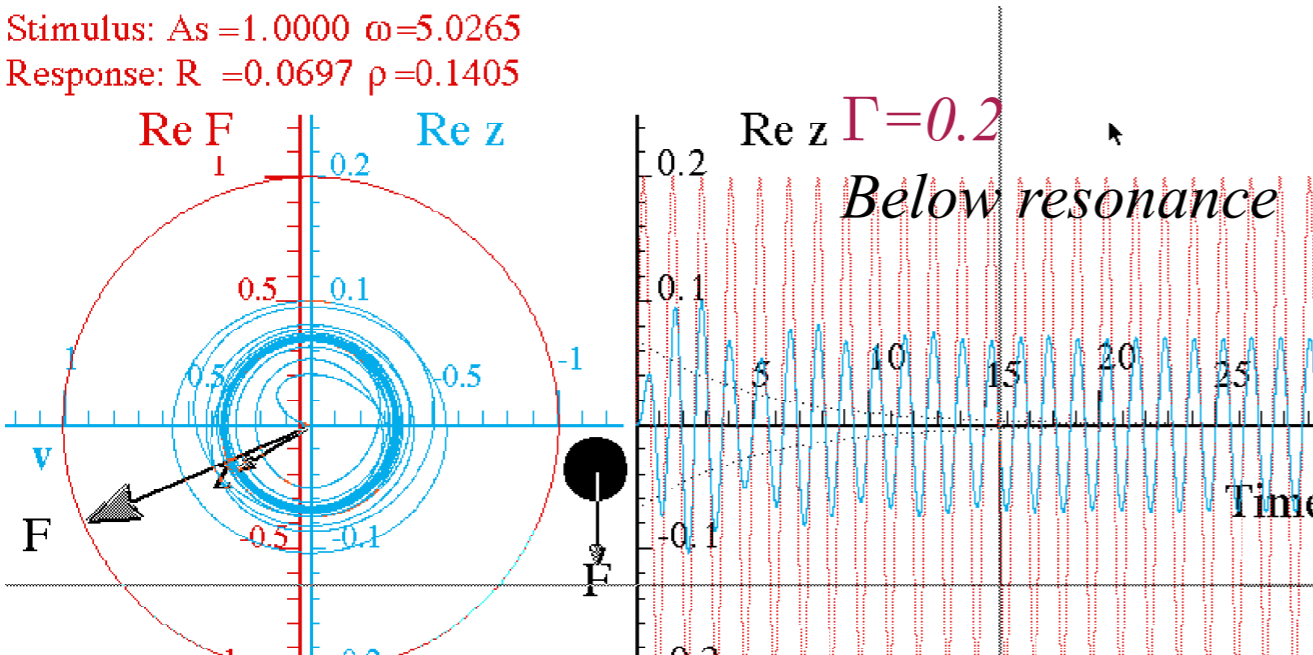
Stimulus: $A_s = 1.0000$ $\omega = 7.5265$
Response: $R = 0.0574$ $\rho = 2.9680$



[OscillIt \(Way Above Resonance\) Simulation](#)

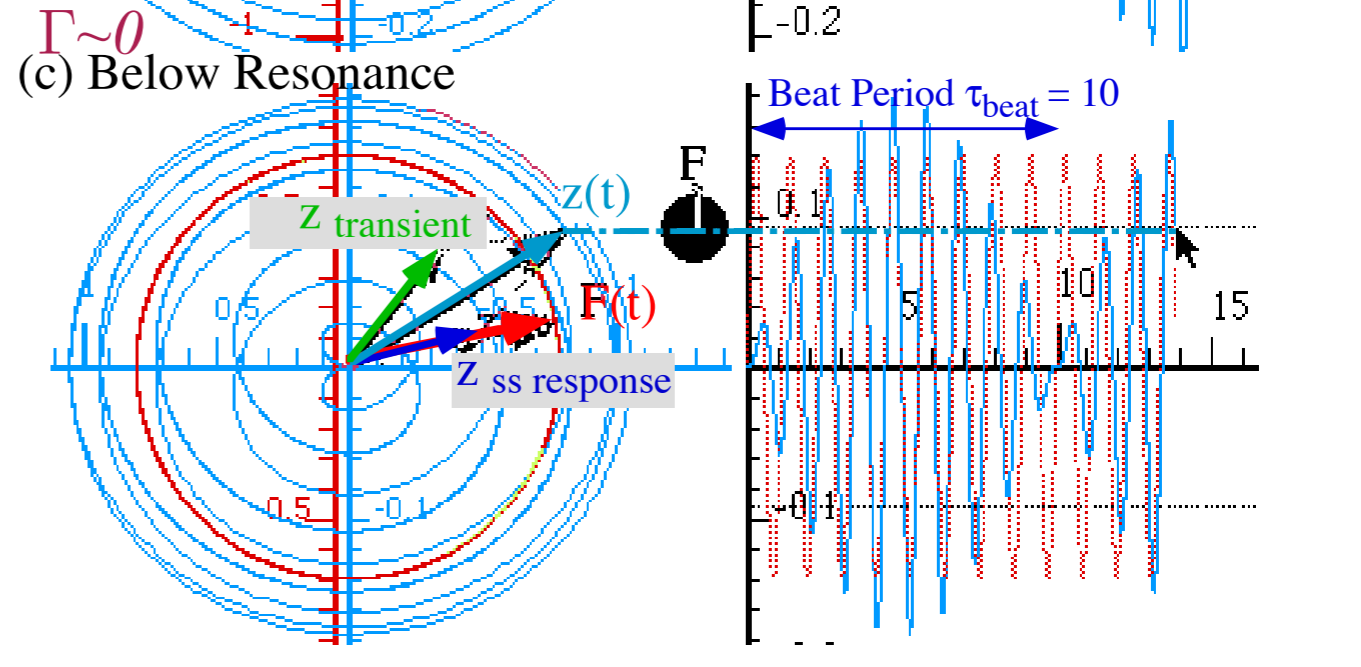
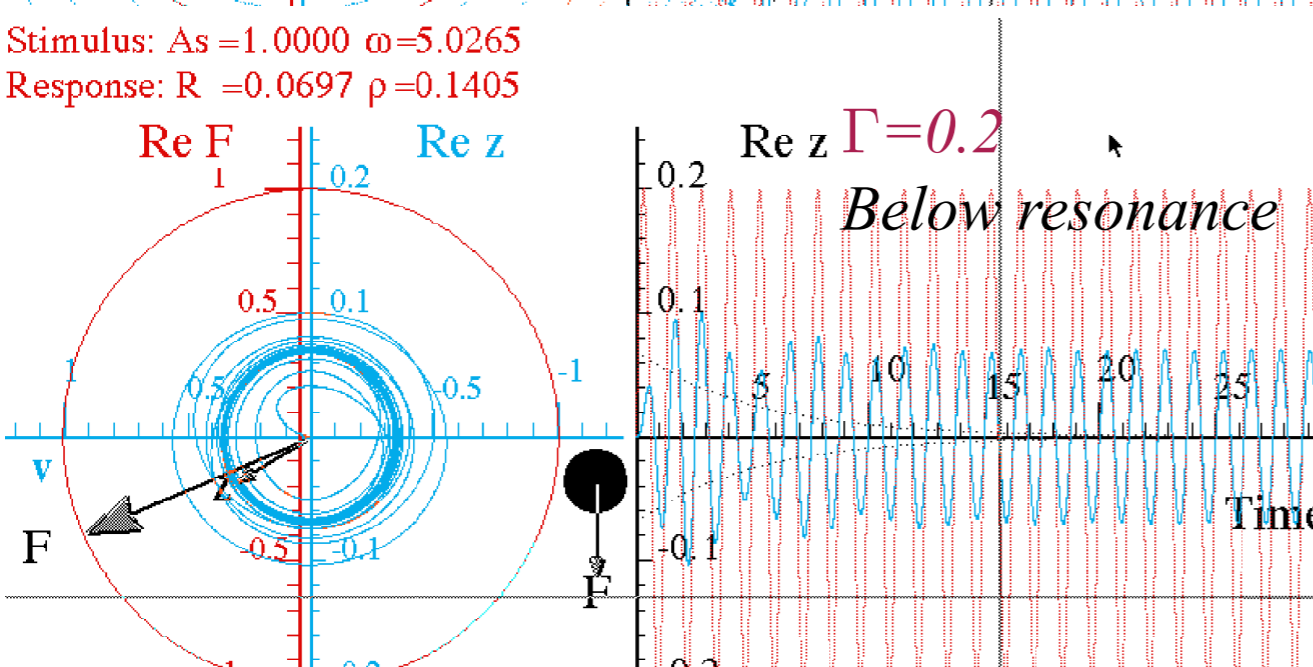
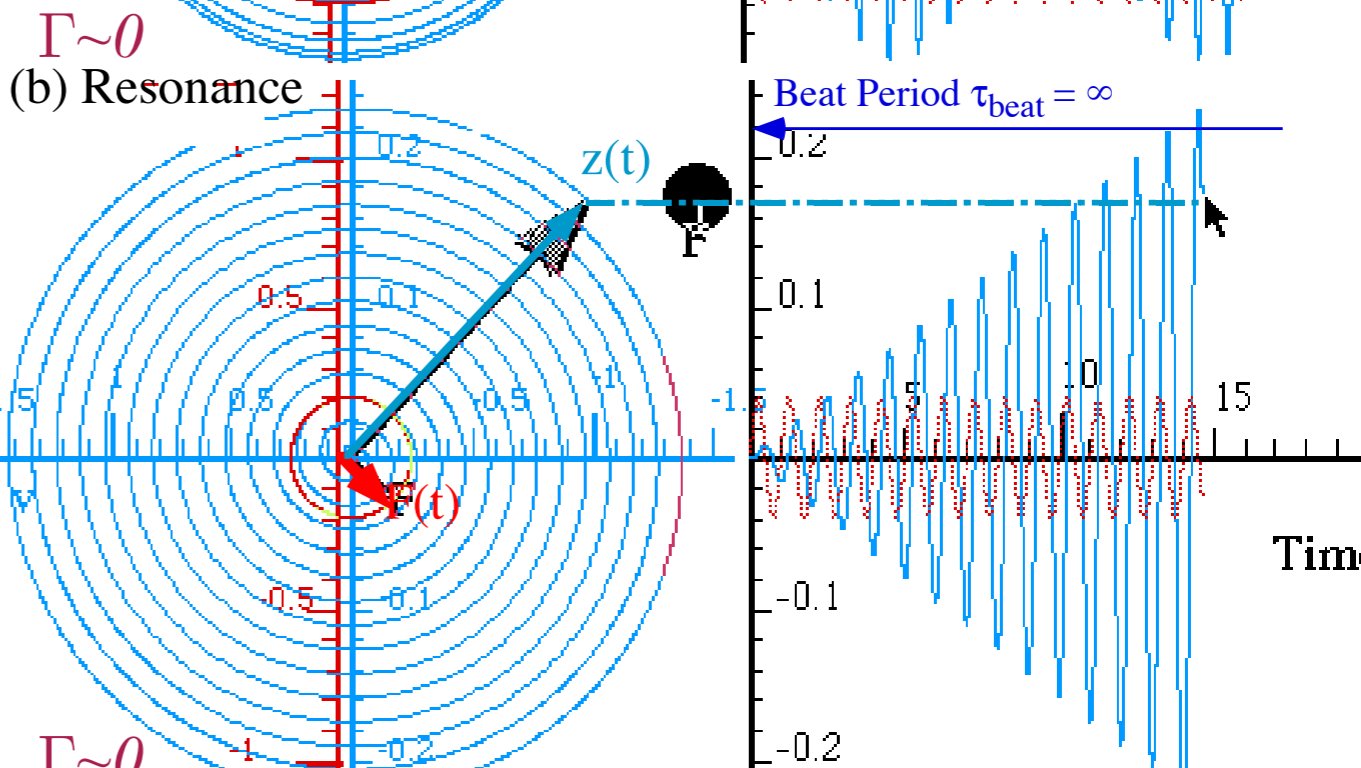
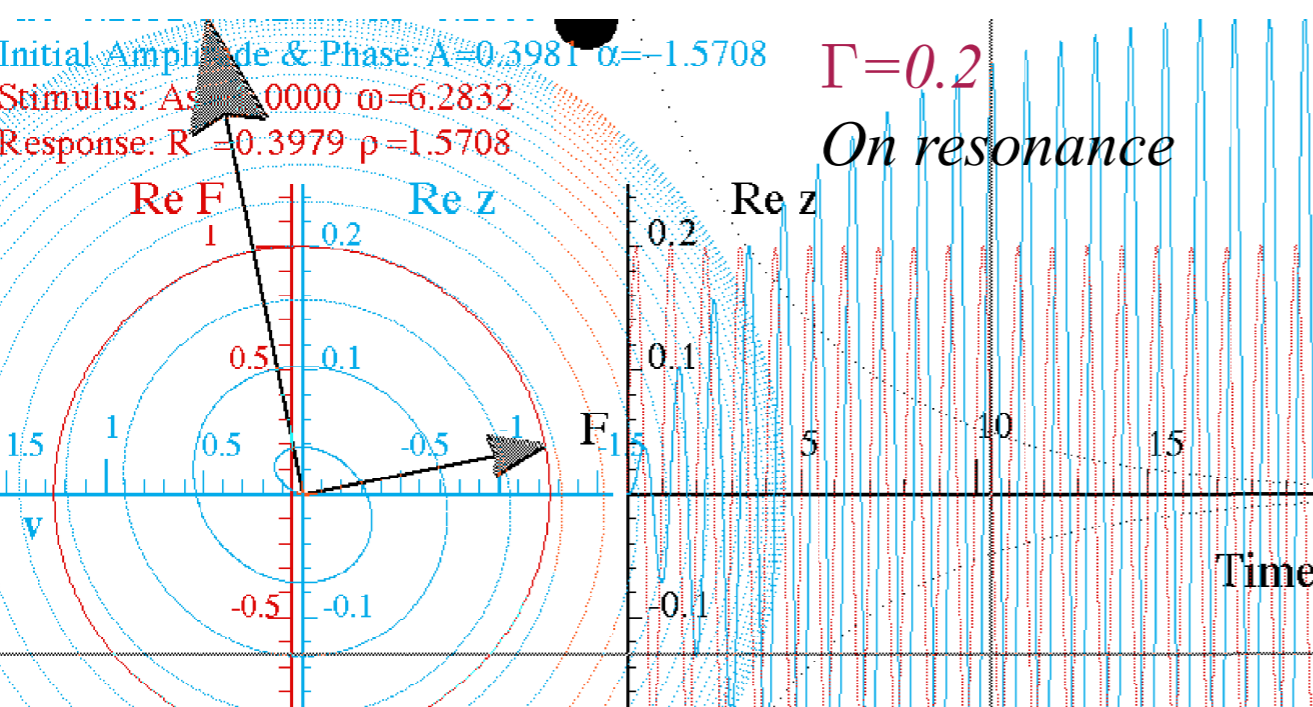
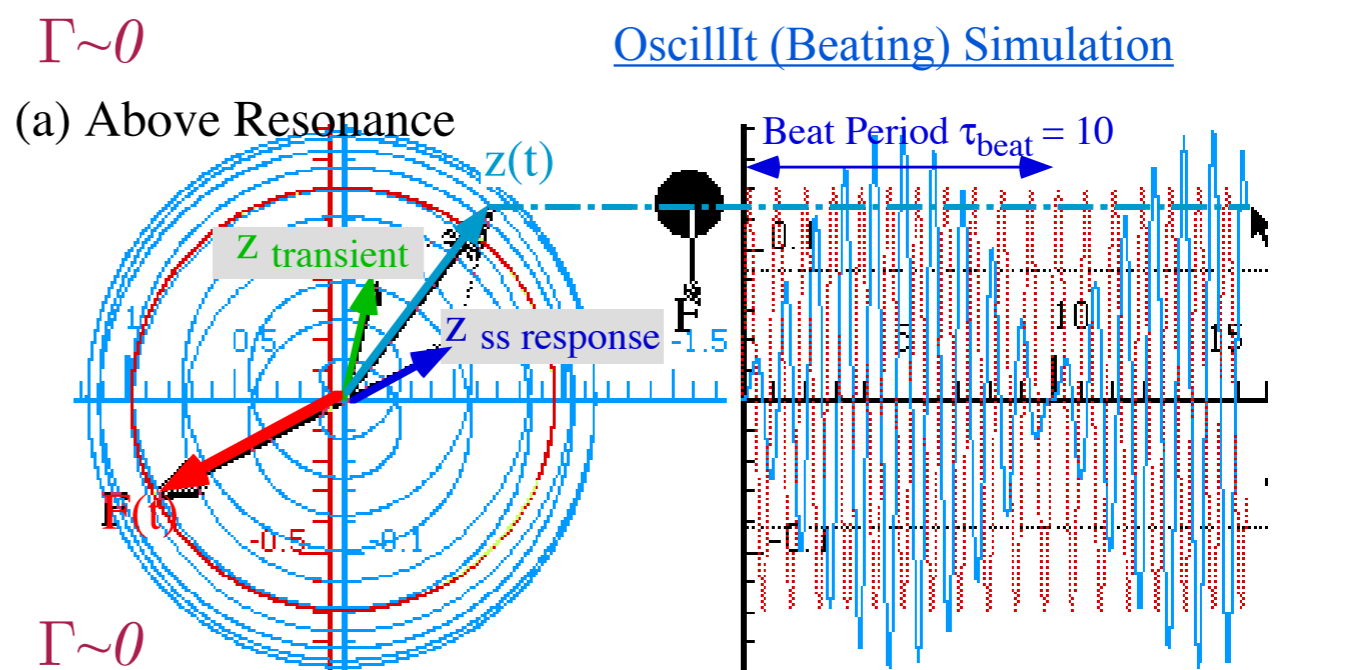
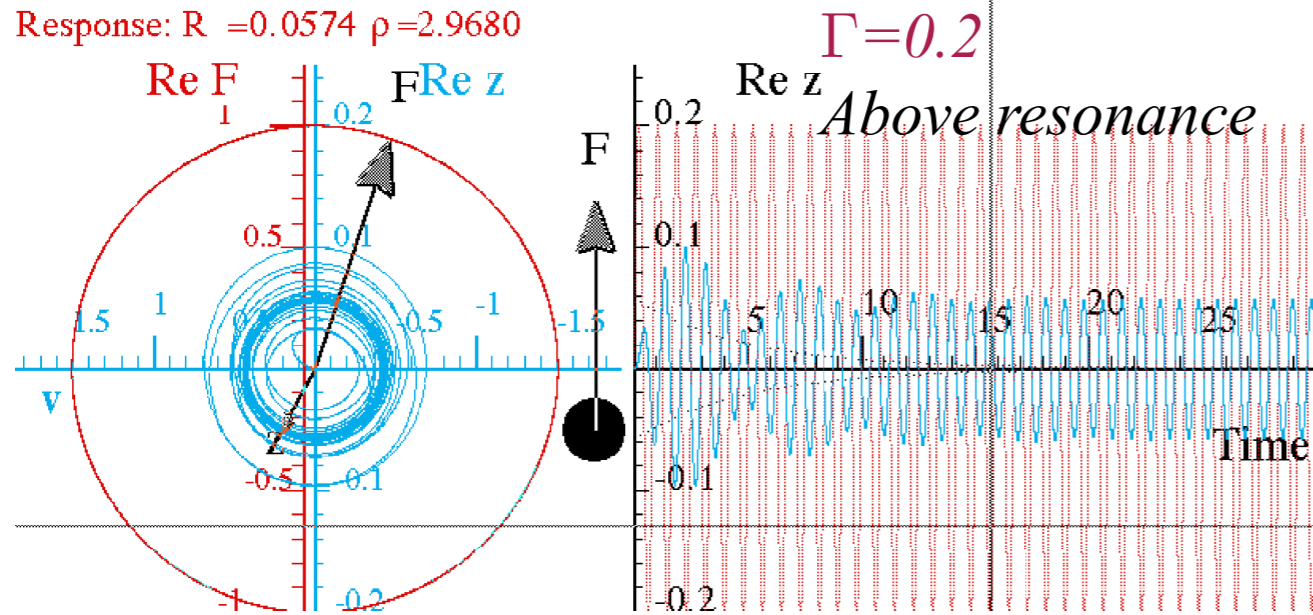


[OscillIt \(On Resonance\) Simulation](#)



[OscillIt \(Way Below Resonance\) Simulation](#)

Stimulus: $A_s = 1.0000$ $\omega = 7.5265$
 Response: $R = 0.0574$ $\rho = 2.9680$



Lorentz-Green's Function for high quality *FDHO*

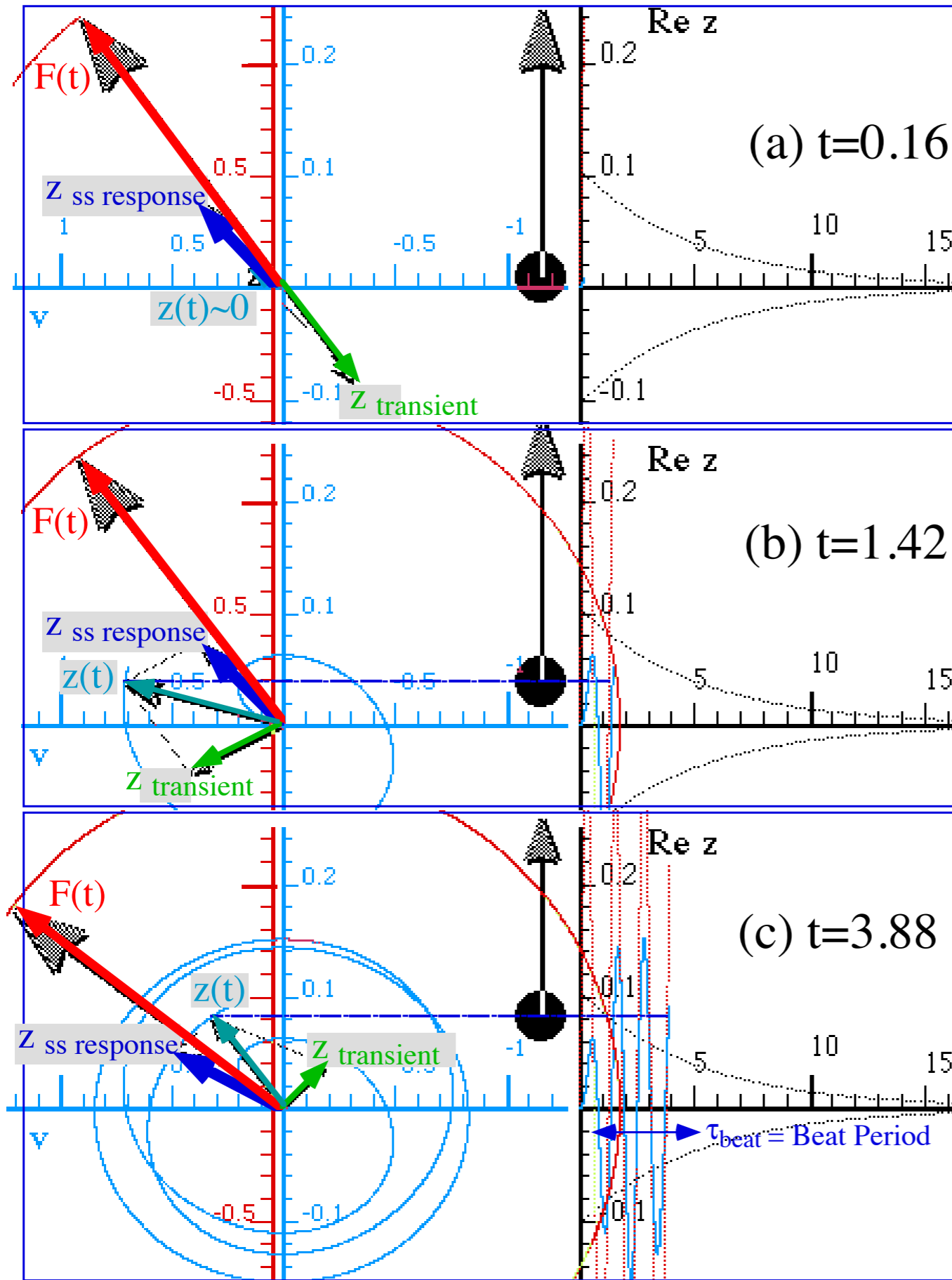


Fig. 4.2.9 Beat formation.

Transient phasor $Z_{transient}$ catches up with F -phasor and passes it.

[OscillIt \(Beating\) Simulation](#)

Oscillator figures of merit: quality factors Q and $q=2\pi Q$

$$AAF = \frac{\text{Resonant response}}{\text{DC response}} = \frac{|G_{\omega_0}(\omega_s = \omega_0)|}{|G_{\omega_0}(0)|} = \frac{1/(2\Gamma\omega_0)}{1/\omega_0^2} = \frac{\omega_0}{2\Gamma} \equiv q \quad (\text{angular quality factor})$$

$$\text{Amplification factor } q = \omega_0/2\Gamma$$

Natural oscillation frequency is approximately $\nu_0 = \omega_0/2\pi$ (for $\omega_0 \gg \Gamma$ we have $\omega_0 \sim \omega_\Gamma$).

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Natural oscillation frequency is approximately $\nu_0 = \omega_0/2\pi$ (for $\omega_0 \gg \Gamma$ we have $\omega_0 \sim \omega_\Gamma$).

$$\left(\begin{array}{l} t_{5\%} = 3/\Gamma = \text{Lifetime} \\ \text{for decaying oscillator} \\ \text{to lose 95\% of} \\ \text{amplitude} \end{array} \right) \text{times} \left(\nu_0 = \frac{\omega_0}{2\pi} \right) = \begin{array}{l} \text{number } n_{5\%} \\ \text{of oscillations} \\ \text{in a } t_{5\%} \text{ Lifetime} \end{array}$$

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$$n_{5\%} = t_{5\%} \nu_0 = \frac{3}{\Gamma} \cdot \frac{\omega_0}{2\pi} \cong \frac{\omega_0}{2\Gamma} = q$$

The “Heartbeat Count”
measure of lifetime

Oscillator figures of merit: quality factors Q and $q=2\pi Q$

$$AAF = \frac{\text{Resonant response}}{\text{DC response}} = \frac{|G_{\omega_0}(\omega_s = \omega_0)|}{|G_{\omega_0}(0)|} = \frac{1/(2\Gamma\omega_0)}{1/\omega_0^2} = \frac{\omega_0}{2\Gamma} \equiv q \quad (\text{angular quality factor})$$

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$$\left(\begin{array}{l} t_{5\%} = 3/\Gamma = \text{Lifetime} \\ \text{for decaying oscillator} \\ \text{to lose 95\% of} \\ \text{amplitude} \end{array} \right) \text{times} \left(\nu_0 = \frac{\omega_0}{2\pi} \right) = \begin{array}{l} \text{number } n_{5\%} \\ \text{of oscillations} \\ \text{in a } t_{5\%} \text{ Lifetime} \end{array}$$

$$n_{5\%} = t_{5\%} \nu_0 = \frac{3}{\Gamma} \cdot \frac{\omega_0}{2\pi} \cong \frac{\omega_0}{2\Gamma} = q$$

The “Heartbeat Count”
measure of lifetime

Energy decay
(proportional to the square of oscillator amplitude):

$$\left(e^{\Gamma t} \right)^2 = e^{-2\Gamma t} \qquad dE = -2\Gamma E$$

Oscillator figures of merit: quality factors Q and $q=2\pi Q$

$$AAF = \frac{\text{Resonant response}}{\text{DC response}} = \frac{|G_{\omega_0}(\omega_s = \omega_0)|}{|G_{\omega_0}(0)|} = \frac{1/(2\Gamma\omega_0)}{1/\omega_0^2} = \frac{\omega_0}{2\Gamma} \equiv q \quad (\text{angular quality factor})$$

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Natural oscillation frequency is approximately $\nu_0 = \omega_0/2\pi$ (for $\omega_0 \gg \Gamma$ we have $\omega_0 \sim \omega_\Gamma$).

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$$n_{5\%} = t_{5\%} \nu_0 = \frac{3}{\Gamma} \cdot \frac{\omega_0}{2\pi} \cong \frac{\omega_0}{2\Gamma} = q$$

The “Heartbeat Count”
measure of lifetime

Energy decay
(proportional to the square of oscillator amplitude):

$$\left(e^{\Gamma t} \right)^2 = e^{-2\Gamma t} \quad dE = -2\Gamma E$$

Relative amount
of energy lost
each cycle period

$$= \tau_0 \left(\frac{-dE}{E} \right) = \frac{2\Gamma}{\nu_0} \equiv \frac{1}{Q} = \frac{2\pi}{q}$$

$$\left(\tau_0 = \frac{1}{\nu_0} \right)$$

$$Q = (\text{Standard amplitude quality factor}) = \frac{q}{2\pi}$$

[Click to see p.18.](#)

[Click to see p.35.](#)

Oscillator figures of merit: Uncertainty 1/q

To see a beat we need $\tau_{\text{half-beat}}$ to be less than $\tau_{5\%}$ or $3/\Gamma$. (Here we approximate $\pi \sim 3.0$, again.)

$$\pi / |\omega_s - \omega_0| < 3 / \Gamma$$

$$|\omega_s - \omega_0| > \Gamma$$

This means ω -detuning error is greater than or equal to the decay rate Γ .

Any detuning less than Γ is virtually undetectable.

Total ω uncertainty is $\pm\Gamma$ or twice Γ (that is: FWHM $\Delta\omega = 2\Gamma$). Linear frequency uncertainty is:

The *relative frequency uncertainty* $\frac{2\Gamma}{\omega_0} = \frac{\Delta\omega}{\omega_0} = \frac{1}{q} = \frac{\Delta\nu}{\nu_0}$ $\Delta\nu = \Delta\omega / 2\pi = \Gamma / \pi$

is the *inverse* of the *angular quality factor* q .

$$Q = (\text{Standard amplitude quality factor}) = \frac{q}{2\pi}$$

If we think of the 5% or 4.321% lifetime of a musical note as its time uncertainty Δt , then:

$$\Delta t \Delta \nu = 3 / \pi \approx 1$$

$$\Delta t = t_{5\%} = 3 / \Gamma$$

$$\Delta t = t_{4.321\%} = \pi / \Gamma$$

Very precise measures of imprecision

Approximate Lorentz-Green's Function for high quality *FDHO* (Quantum propagator)

$$G_{\omega_0}(\omega_s) = \frac{1}{\omega_0^2 - \omega_s^2 - i2\Gamma\omega_s} \xrightarrow{\omega_s \rightarrow \omega_0} \frac{1}{2\omega_s} \frac{1}{\omega_0 - \omega_s - i\Gamma} \approx \frac{1}{2\omega_0} \frac{1}{\Delta - i\Gamma} = \frac{1}{2\omega_0} L(\Delta - i\Gamma)$$

Complex detuning-decay $\delta = \Delta - i\Gamma$ variable δ is defined with the real detuning $\Delta = \omega_0 - \omega_s$

Approximate Lorentz-Green's Function for high quality *FDHO* (Quantum propagator)

$$G_{\omega_0}(\omega_s) = \frac{1}{\omega_0^2 - \omega_s^2 - i2\Gamma\omega_s} \xrightarrow{\omega_s \rightarrow \omega_0} \frac{1}{2\omega_s} \frac{1}{\omega_0 - \omega_s - i\Gamma} \approx \frac{1}{2\omega_0} \frac{1}{\Delta - i\Gamma} = \frac{1}{2\omega_0} L(\Delta - i\Gamma)$$

Complex detuning-decay $\delta = \Delta - i\Gamma$ variable δ is defined with the real detuning $\Delta = \omega_0 - \omega_s$

$$L(\Delta - i\Gamma) = \frac{1}{\Delta - i\Gamma} = \text{Re } L + i \text{Im } L = \frac{\Delta}{\Delta^2 + \Gamma^2} + i \frac{\Gamma}{\Delta^2 + \Gamma^2} = |L|^2 \Delta + i |L|^2 \Gamma$$

Approximate Lorentz-Green's Function for high quality *FDHO* (Quantum propagator)

$$G_{\omega_0}(\omega_s) = \frac{1}{\omega_0^2 - \omega_s^2 - i2\Gamma\omega_s} \xrightarrow{\omega_s \rightarrow \omega_0} \frac{1}{2\omega_s} \frac{1}{\omega_0 - \omega_s - i\Gamma} \approx \frac{1}{2\omega_0} \frac{1}{\Delta - i\Gamma} = \frac{1}{2\omega_0} L(\Delta - i\Gamma)$$

Complex detuning-decay $\delta = \Delta - i\Gamma$ variable δ is defined with the real detuning $\Delta = \omega_0 - \omega_s$

$$\begin{aligned} L(\Delta - i\Gamma) &= \frac{1}{\Delta - i\Gamma} = \text{Re } L + i \text{Im } L = \frac{\Delta}{\Delta^2 + \Gamma^2} + i \frac{\Gamma}{\Delta^2 + \Gamma^2} = |L|^2 \Delta + i |L|^2 \Gamma \\ &= |L| e^{i\rho} = |L| \cos \rho + i |L| \sin \rho = \frac{\cos \rho}{\sqrt{\Delta^2 + \Gamma^2}} + i \frac{\sin \rho}{\sqrt{\Delta^2 + \Gamma^2}} \quad \text{where: } |L| = \frac{1}{\sqrt{\Delta^2 + \Gamma^2}} \end{aligned}$$

Approximate Lorentz-Green's Function for high quality *FDHO* (Quantum propagator)

$$G_{\omega_0}(\omega_s) = \frac{1}{\omega_0^2 - \omega_s^2 - i2\Gamma\omega_s} \xrightarrow{\omega_s \rightarrow \omega_0} \frac{1}{2\omega_s} \frac{1}{\omega_0 - \omega_s - i\Gamma} \approx \frac{1}{2\omega_0} \frac{1}{\Delta - i\Gamma} = \frac{1}{2\omega_0} L(\Delta - i\Gamma)$$

Complex detuning-decay $\delta = \Delta - i\Gamma$ variable δ is defined with the *real detuning* $\Delta = \omega_0 - \omega_s$

$$L(\Delta - i\Gamma) = \frac{1}{\Delta - i\Gamma} = \text{Re } L + i \text{Im } L = \frac{\Delta}{\Delta^2 + \Gamma^2} + i \frac{\Gamma}{\Delta^2 + \Gamma^2} = |L|^2 \Delta + i |L|^2 \Gamma$$

$$= |L| e^{i\rho} = |L| \cos \rho + i |L| \sin \rho = \frac{\cos \rho}{\sqrt{\Delta^2 + \Gamma^2}} + i \frac{\sin \rho}{\sqrt{\Delta^2 + \Gamma^2}} \quad \text{where: } |L| = \frac{1}{\sqrt{\Delta^2 + \Gamma^2}}$$

Ideal Lorentz-Green's functions

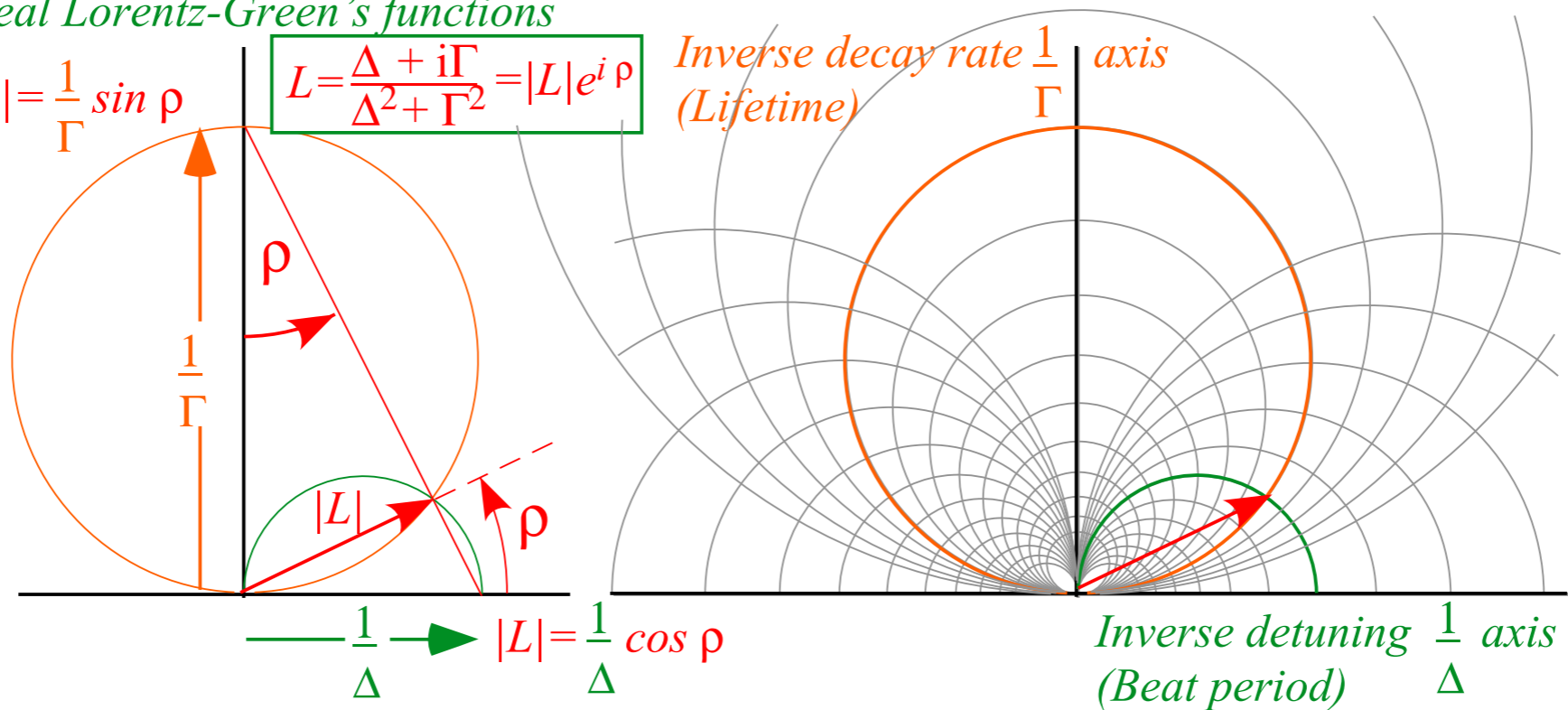
$$L = \frac{\Delta + i\Gamma}{\Delta^2 + \Gamma^2} = |L| e^{i\rho}$$

Inverse decay rate $\frac{1}{\Gamma}$ axis
(Lifetime)

Smith plots

$$|L| = \frac{1}{\Gamma} \sin \rho$$

$$|L| = \frac{1}{\Delta} \cos \rho$$



Approximate Lorentz-Green's Function for high quality *FDHO* (Quantum propagator)

$$G_{\omega_0}(\omega_s) = \frac{1}{\omega_0^2 - \omega_s^2 - i2\Gamma\omega_s} \xrightarrow{\omega_s \rightarrow \omega_0} \frac{1}{2\omega_s} \frac{1}{\omega_0 - \omega_s - i\Gamma} \approx \frac{1}{2\omega_0} \frac{1}{\Delta - i\Gamma} = \frac{1}{2\omega_0} L(\Delta - i\Gamma)$$

Complex detuning-decay $\delta = \Delta - i\Gamma$ variable δ is defined with the real detuning $\Delta = \omega_0 - \omega_s$

$$L(\Delta - i\Gamma) = \frac{1}{\Delta - i\Gamma} = \text{Re } L + i \text{Im } L = \frac{\Delta}{\Delta^2 + \Gamma^2} + i \frac{\Gamma}{\Delta^2 + \Gamma^2} = |L|^2 \Delta + i |L|^2 \Gamma$$

$$= |L| e^{i\rho} = |L| \cos \rho + i |L| \sin \rho = \frac{\cos \rho}{\sqrt{\Delta^2 + \Gamma^2}} + i \frac{\sin \rho}{\sqrt{\Delta^2 + \Gamma^2}} \quad \text{where: } |L| = \frac{1}{\sqrt{\Delta^2 + \Gamma^2}}$$

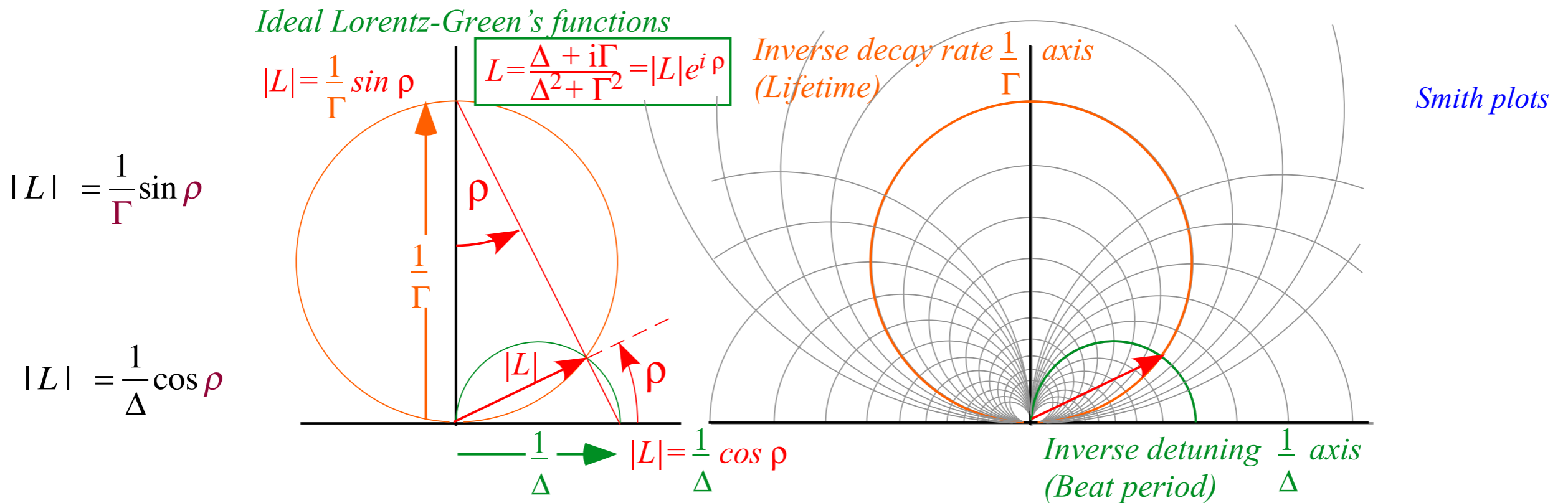


Fig. 4.2.13 Ideal Lorentzian in inverse rate space. (Smith life-time $1/\Gamma$ vs. beat-period $1/\Delta$ coordinates)

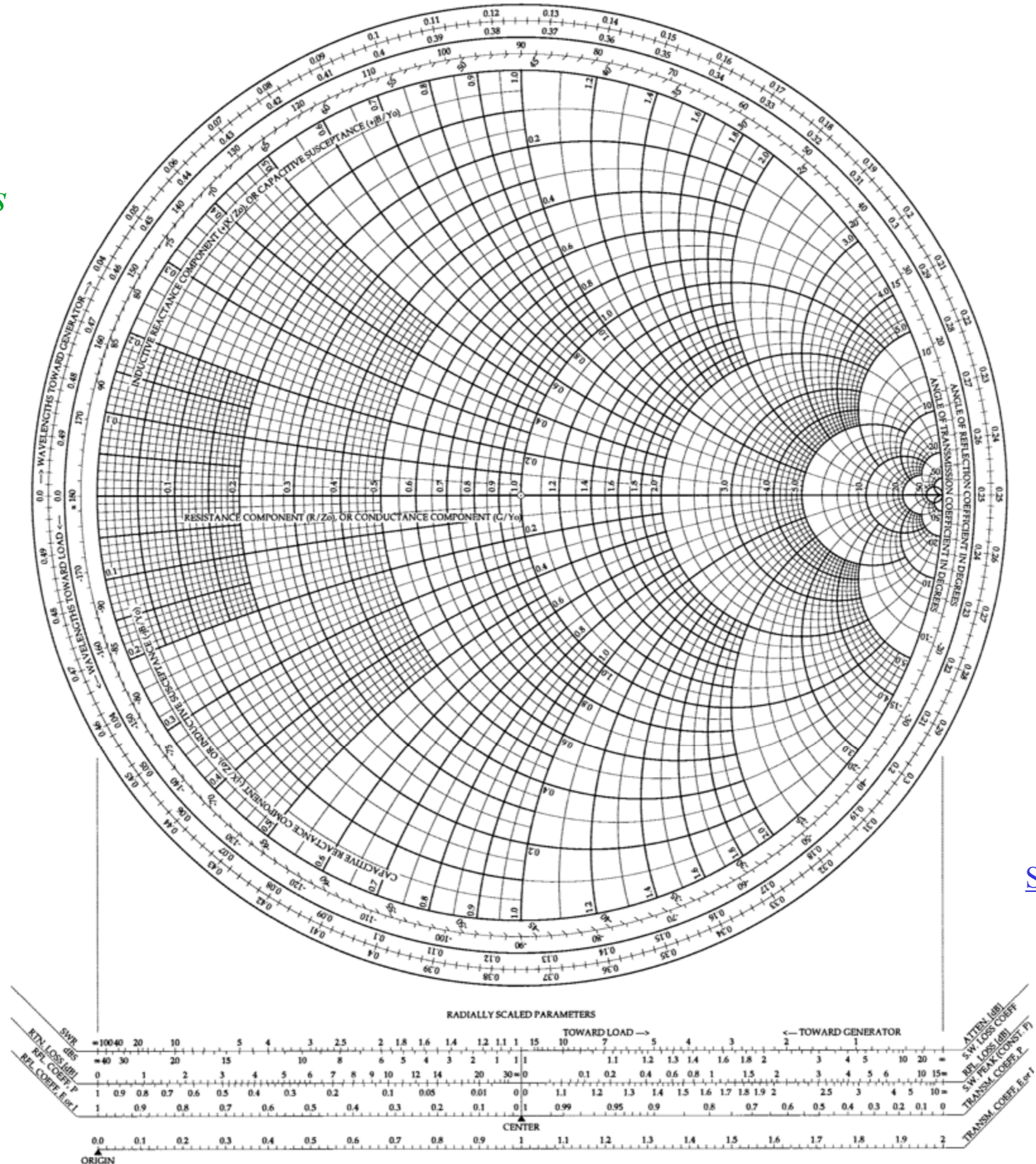
Constant Δ and Γ curves in Fig. 4.2.13 are orthogonal circles of $1/z$ -dipolar coordinates. Recall Fig. 1.10.11.

SMITH CHART (Invented by Phillip H. Smith 1905-1987)

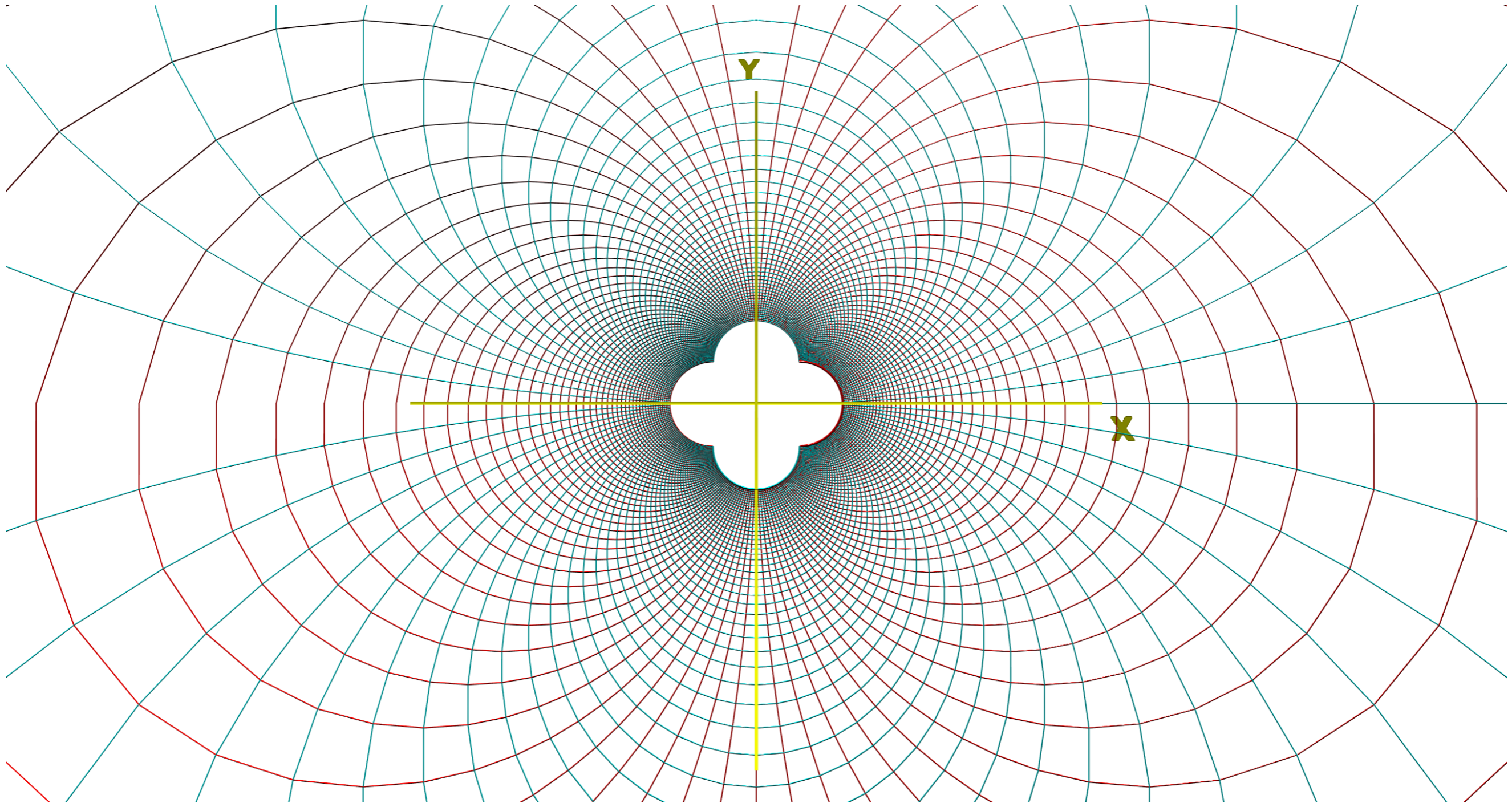
An FDHO Green's
Function
Slide rule

A plot of
 $f(z) = 1/z$

For wavy
"Ohm's Laws"
 $V = I \cdot Z$
 $I = V/Z$



Smith plot: Graph paper



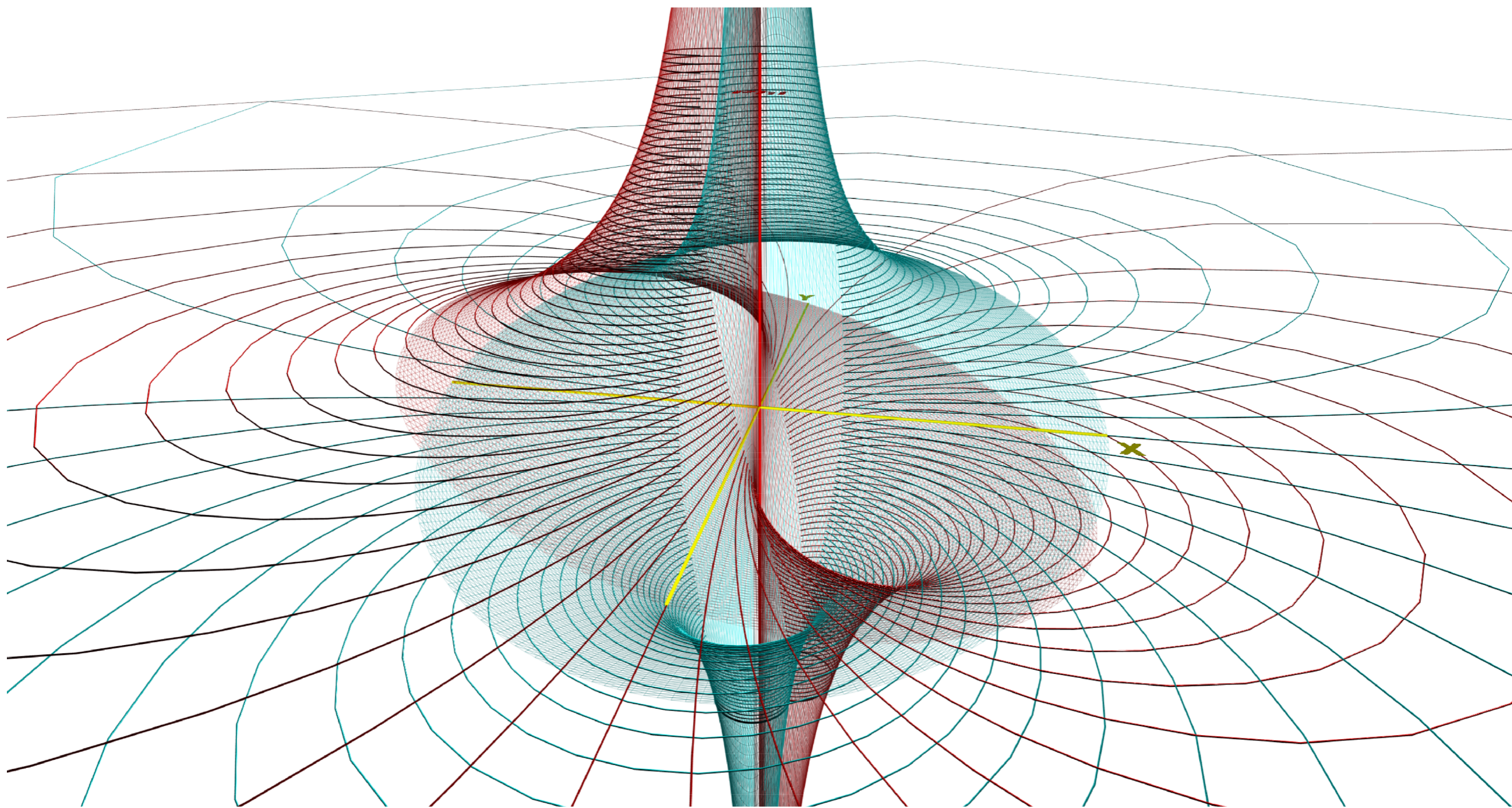
https://modphys.hosted.uark.edu/video/AnalyIt_0-3.webm

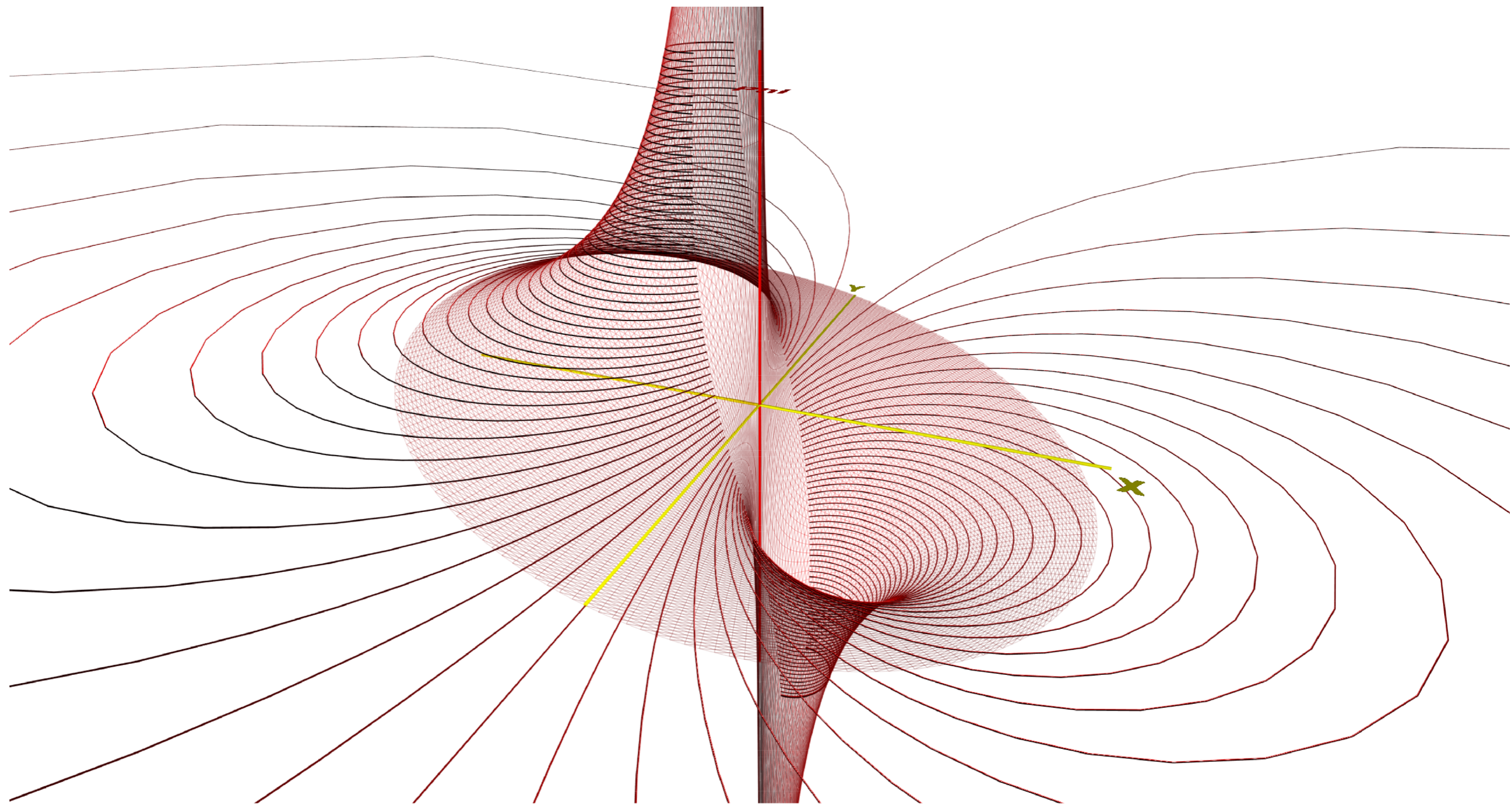
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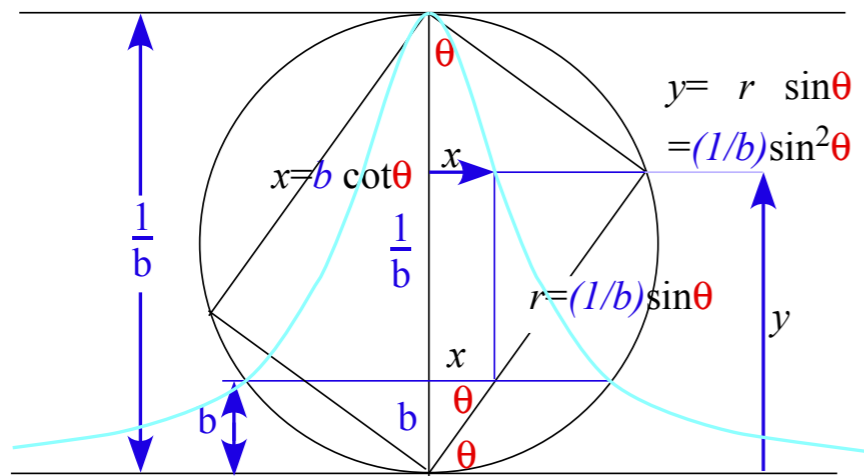


The Common Lorentzian (a.k.a. The Witch of Agnesi)

Maria Gaetana Agnesi



Born May 16, 1718
Died January 9, 1799 (aged 80)
Residence Italy
Nationality Italy
Fields Mathematics



$$x^2 = b^2 \cot^2 \theta = b^2 \frac{\cos^2 \theta}{\sin^2 \theta} = b^2 \frac{1 - \sin^2 \theta}{\sin^2 \theta} = \frac{b^2}{\sin^2 \theta} - b^2$$

$$x^2 + b^2 = \frac{b^2}{\sin^2 \theta} = \frac{b}{y}$$

$$y = \frac{b}{x^2 + b^2}$$

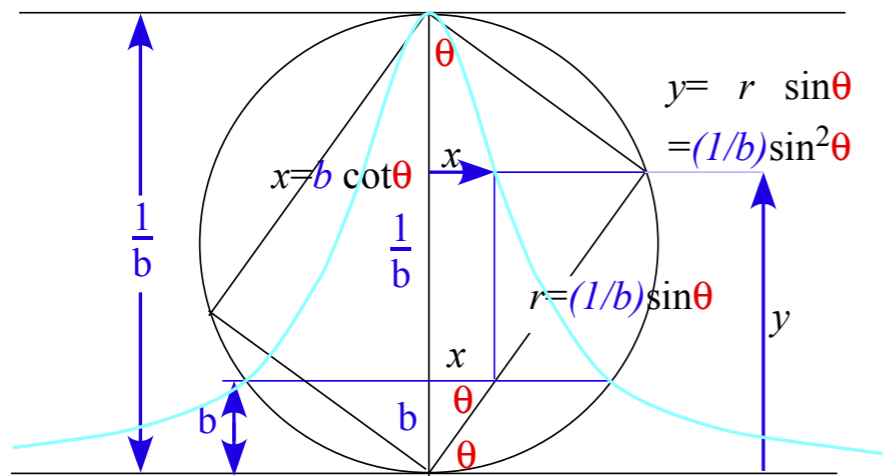
Common Lorentzian function I.
(imaginary "absorbive" part)

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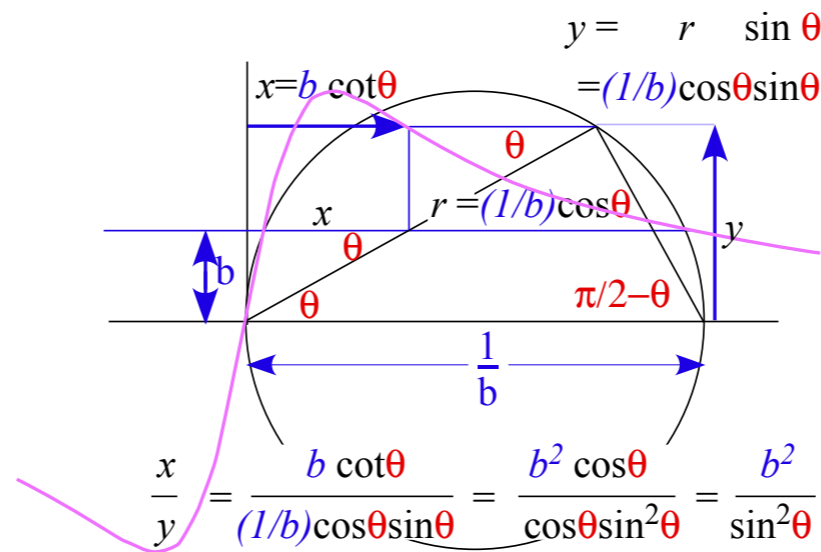
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$$x^2 + b^2 = \frac{b^2}{\sin^2 \theta} = \frac{b}{y} \quad \boxed{y = \frac{b}{x^2 + b^2}}$$

*Common Lorentzian function I.
(imaginary "absorbive" part)*



$$\frac{x}{y} = \frac{b \cot \theta}{(1/b) \cos \theta \sin \theta} = \frac{b^2 \cos \theta}{\cos \theta \sin^2 \theta} = \frac{b^2}{\sin^2 \theta}$$

$$x^2 + b^2 = \frac{b^2}{\sin^2 \theta} = \frac{x}{y} \quad \boxed{y = \frac{x}{x^2 + b^2}}$$

*Common Lorentzian function II.
(real "refractory" part)*

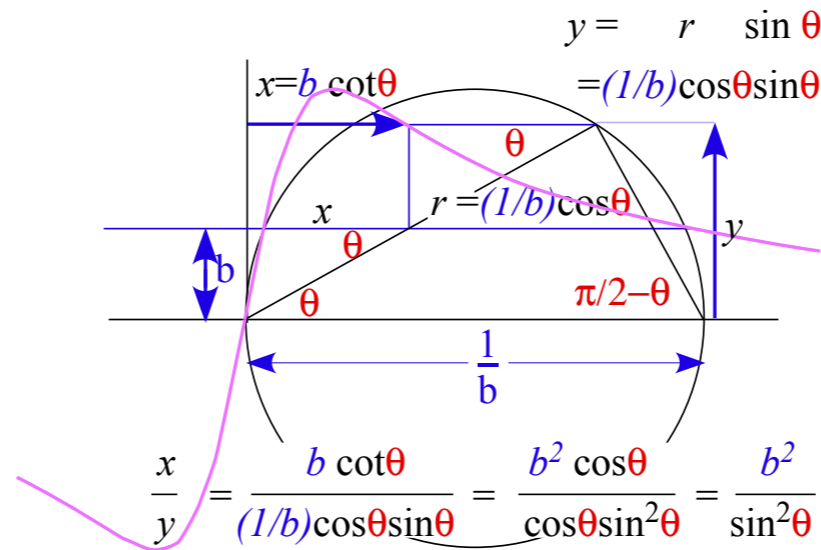
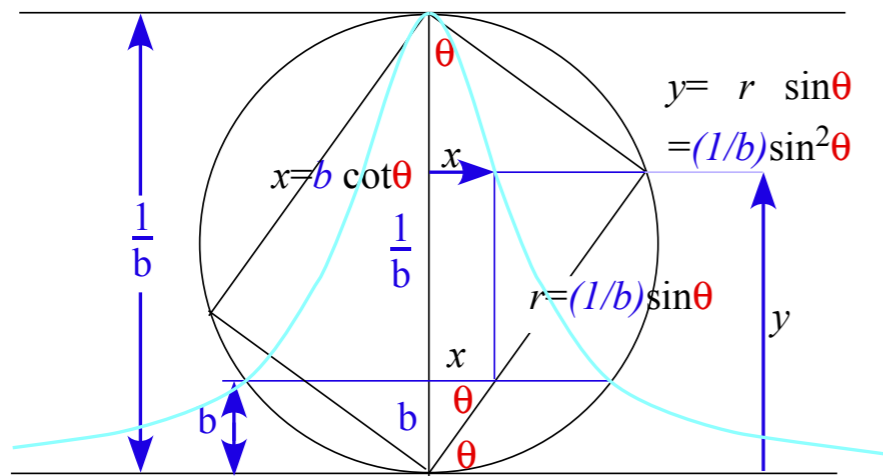


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$$x^2 + b^2 = \frac{b^2}{\sin^2 \theta} = \frac{b}{y}$$

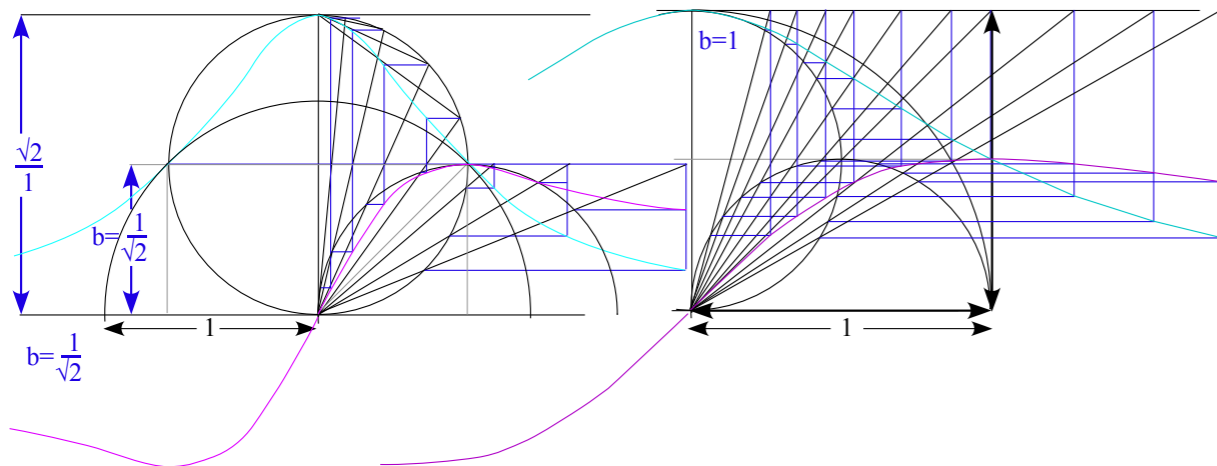
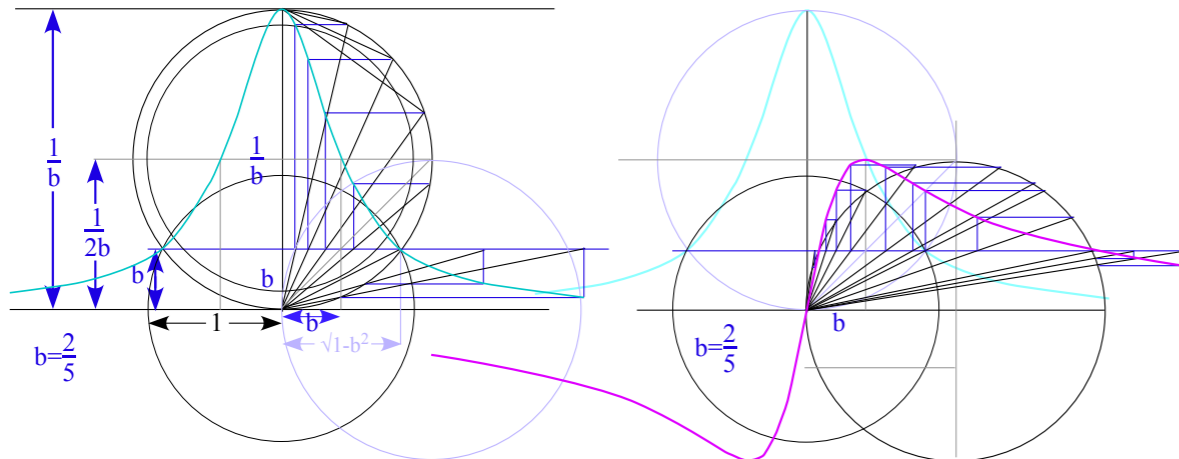
$$y = \frac{b}{x^2 + b^2}$$

*Common Lorentzian function I.
(imaginary "absorbive" part)*

$$x^2 + b^2 = \frac{b^2}{\sin^2 \theta} = \frac{x}{y}$$

$$y = \frac{x}{x^2 + b^2}$$

*Common Lorentzian function II.
(real "refractory" part)*

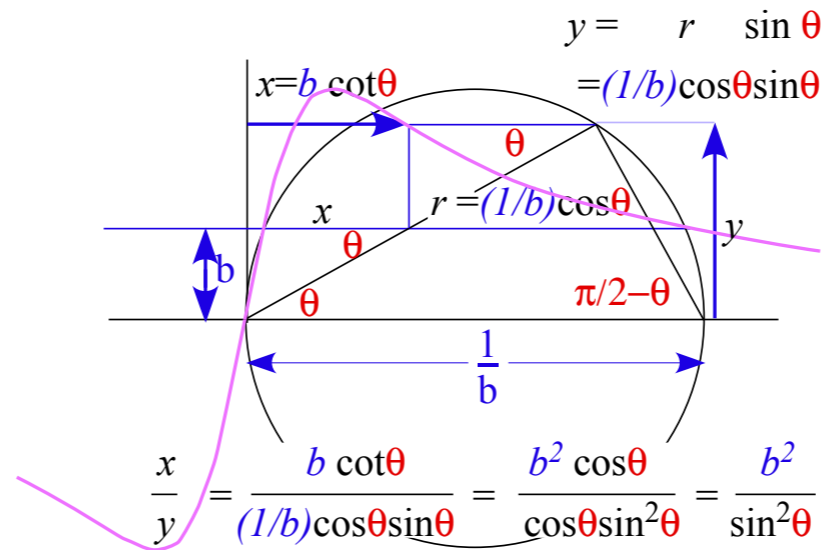
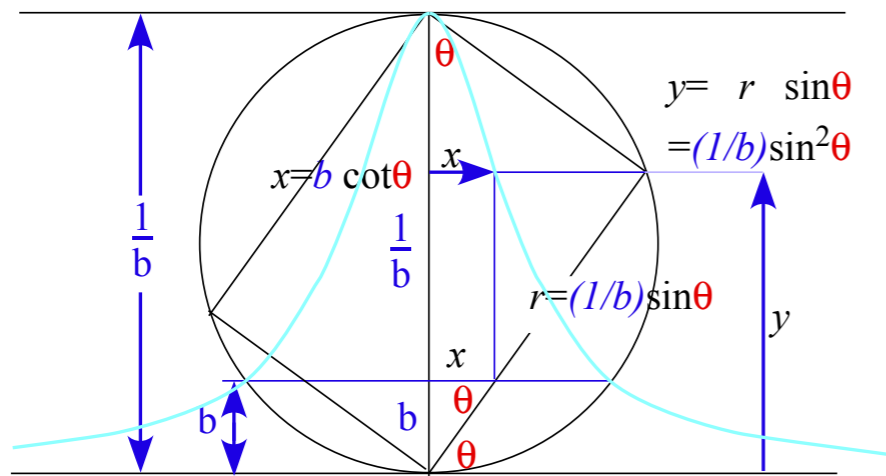


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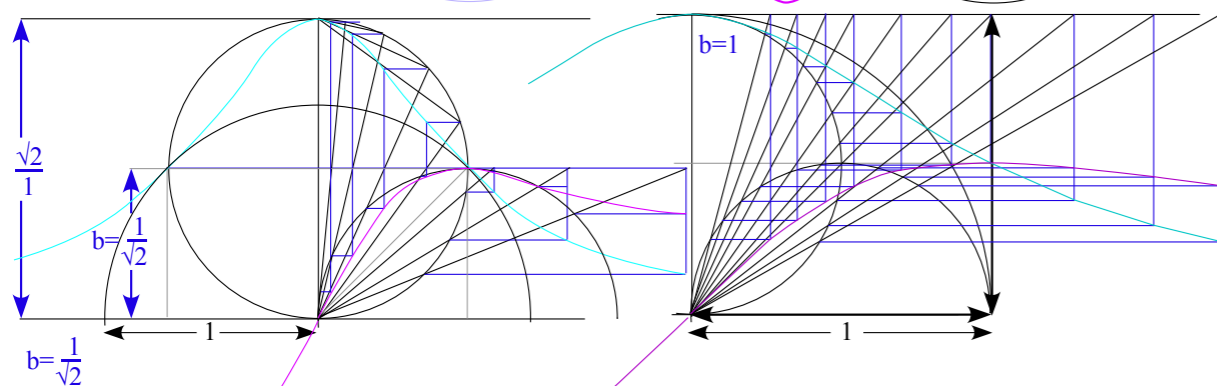
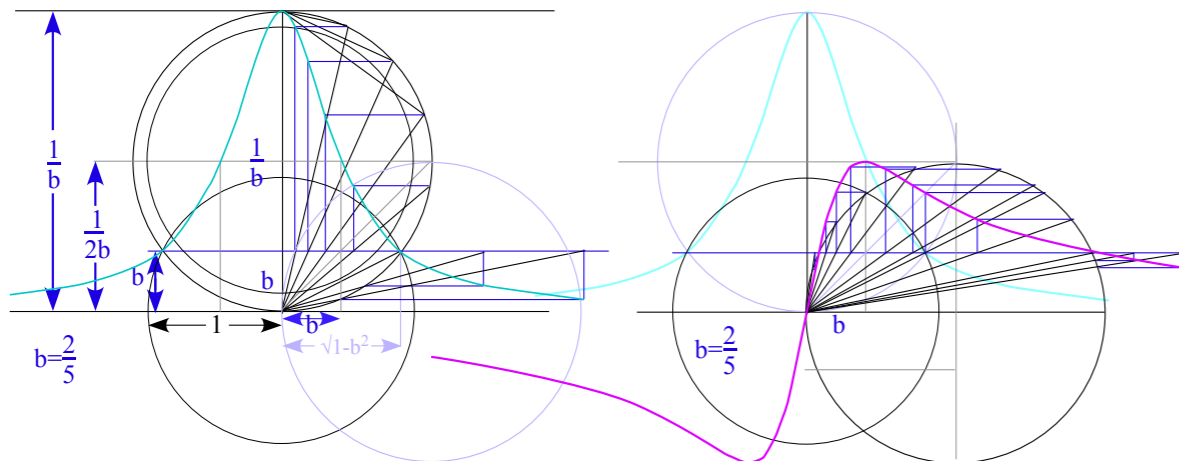
$$x^2 = b^2 \cot^2 \theta = b^2 \frac{\cos^2 \theta}{\sin^2 \theta} = b^2 \frac{1 - \sin^2 \theta}{\sin^2 \theta} = \frac{b^2}{\sin^2 \theta} - b^2$$

$$x^2 + b^2 = \frac{b^2}{\sin^2 \theta} = \frac{b}{y} \quad y = \frac{b}{x^2 + b^2}$$

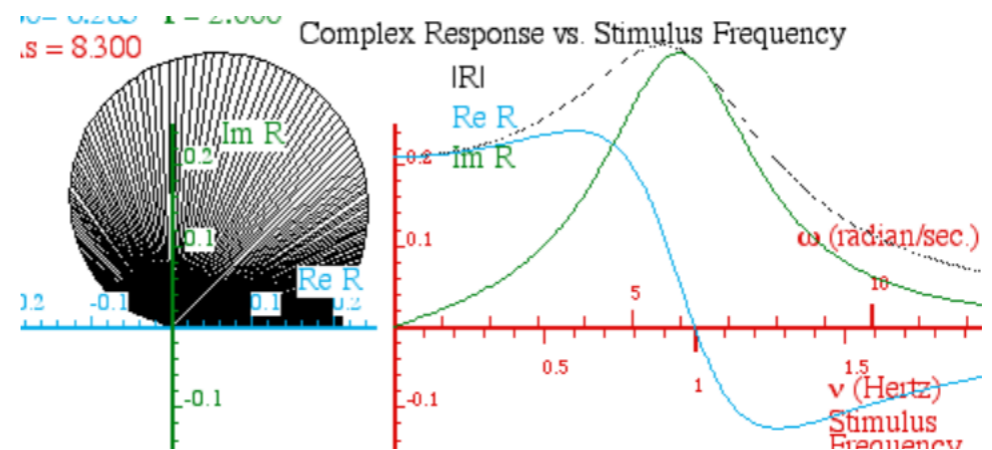
Common Lorentzian function I.
(imaginary "absorbptive" part)

$$x^2 + b^2 = \frac{b^2}{\sin^2 \theta} = \frac{x}{y} \quad y = \frac{x}{x^2 + b^2}$$

Common Lorentzian function II.
(real "refractory" part)



Underlined below are links to the OscillIt Web Simulations
Compare ideal Lorentzians ($\Gamma=0.2$)
with a very non-ideal one ($\Gamma=2$)

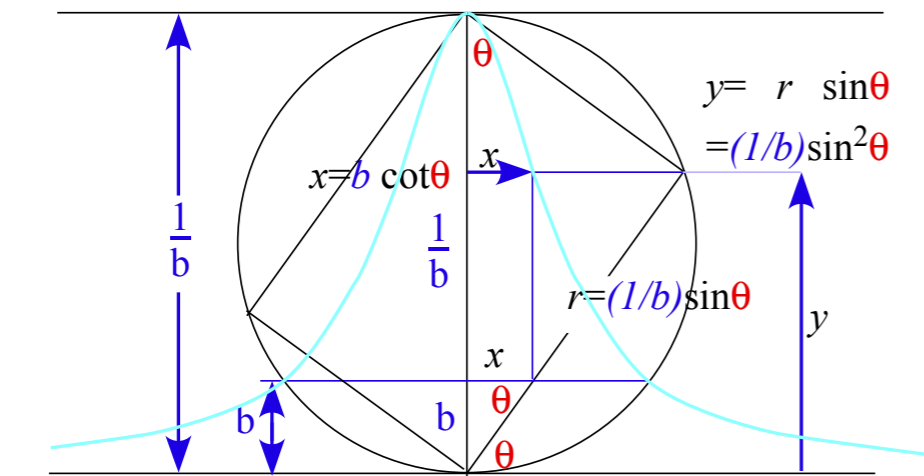


The Common Lorentzian (a.k.a. The Witch of Agnesi)

Maria Gaetana Agnesi



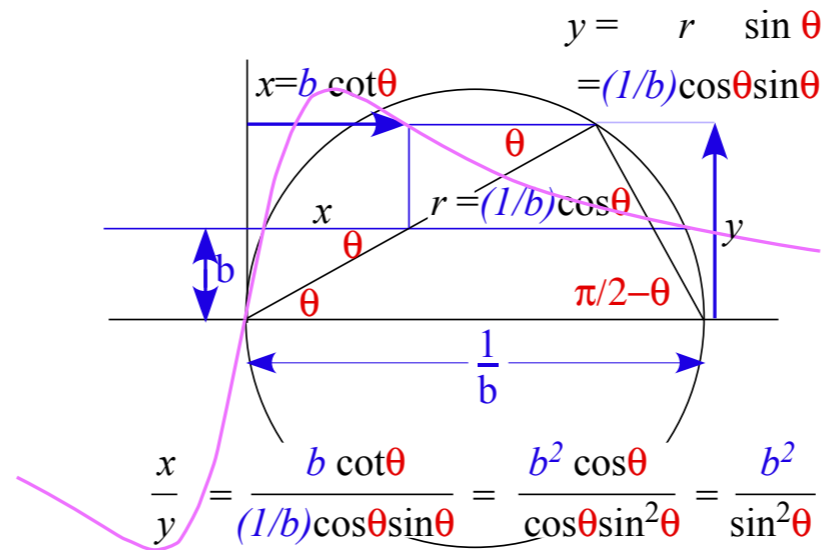
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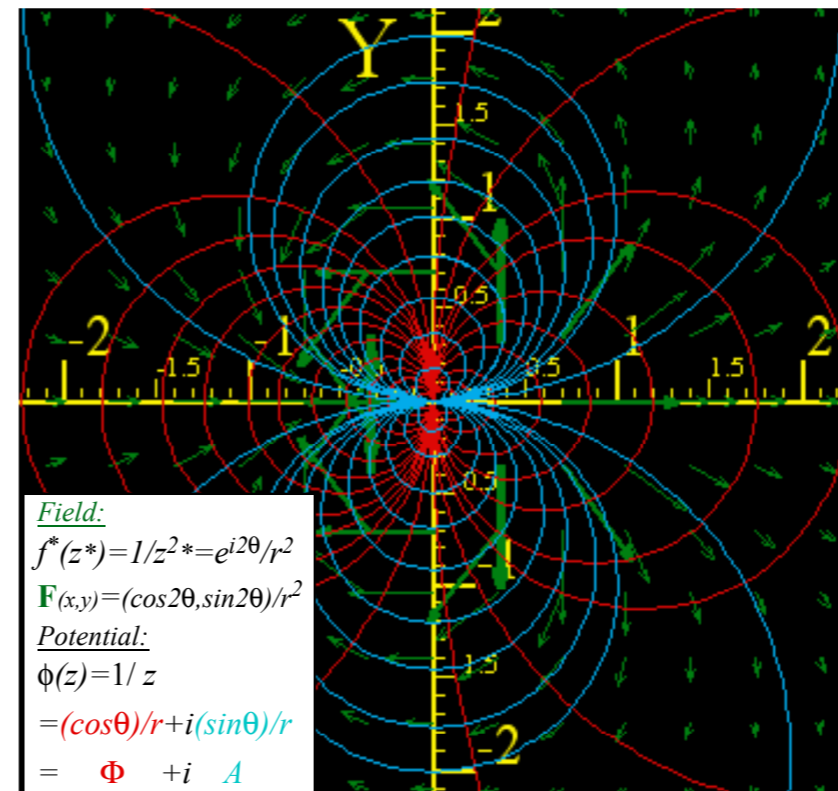
Common Lorentzian function I.
(imaginary "absorbive" part)



$$\frac{x}{y} = \frac{b \cot \theta}{(1/b) \cos \theta \sin \theta} = \frac{b^2 \cos \theta}{\cos \theta \sin^2 \theta} = \frac{b^2}{\sin^2 \theta}$$

$$x^2 + b^2 = \frac{b^2}{\sin^2 \theta} = \frac{x}{y} \quad y = \frac{x}{x^2 + b^2}$$

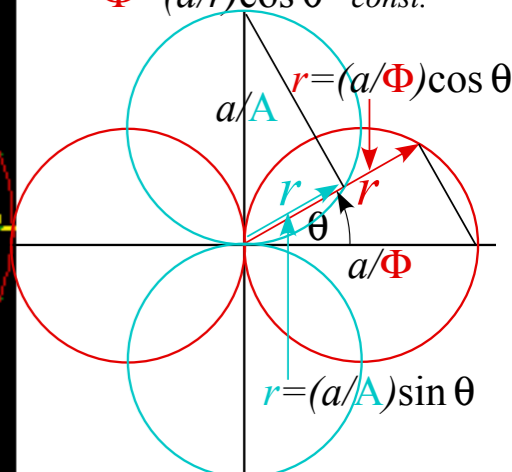
Common Lorentzian function II.
(real "refractory" part)



Field:
 $f^*(z^*) = 1/z^{*2} = e^{i2\theta}/r^2$
 $\mathbf{F}(x,y) = (\cos 2\theta, \sin 2\theta)/r^2$
 Potential:
 $\phi(z) = 1/z$
 $= (\cos \theta)/r + i(\sin \theta)/r$
 $= \Phi + i A$

Scalar potentials

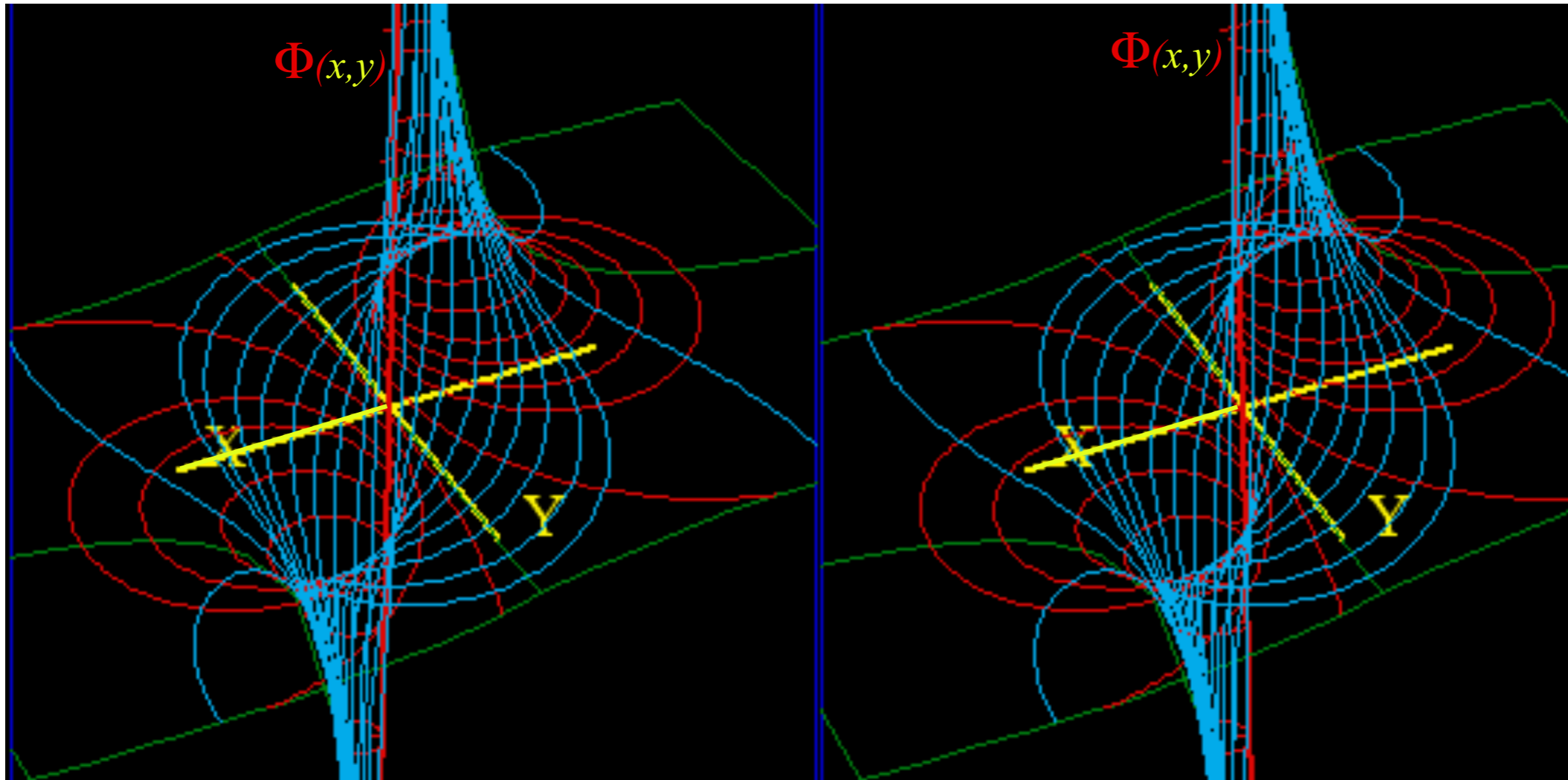
$$\Phi = (a/r) \cos \theta = \text{const.}$$



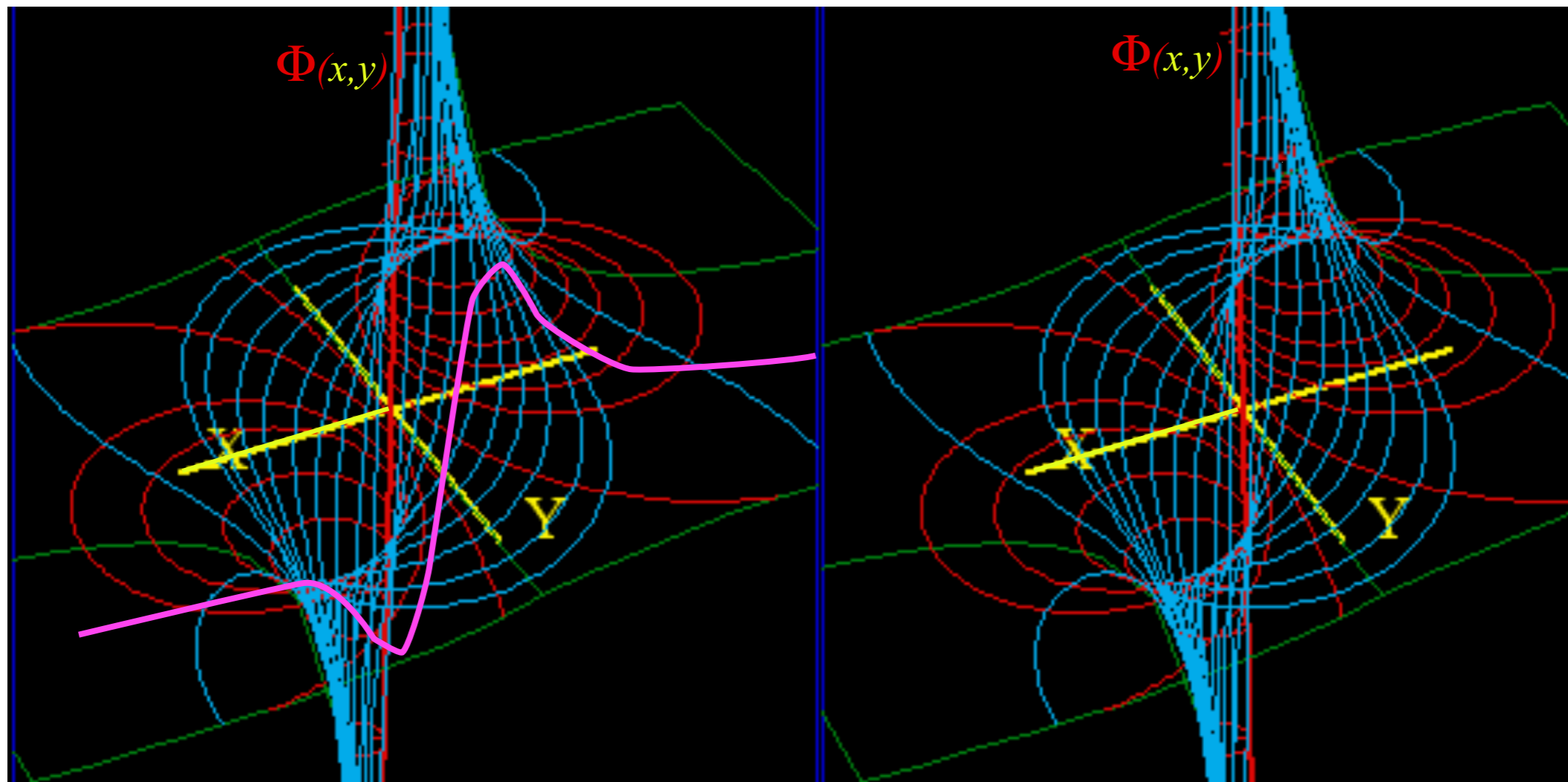
Vector potentials

$$A = (a/r) \sin \theta = \text{const.}$$

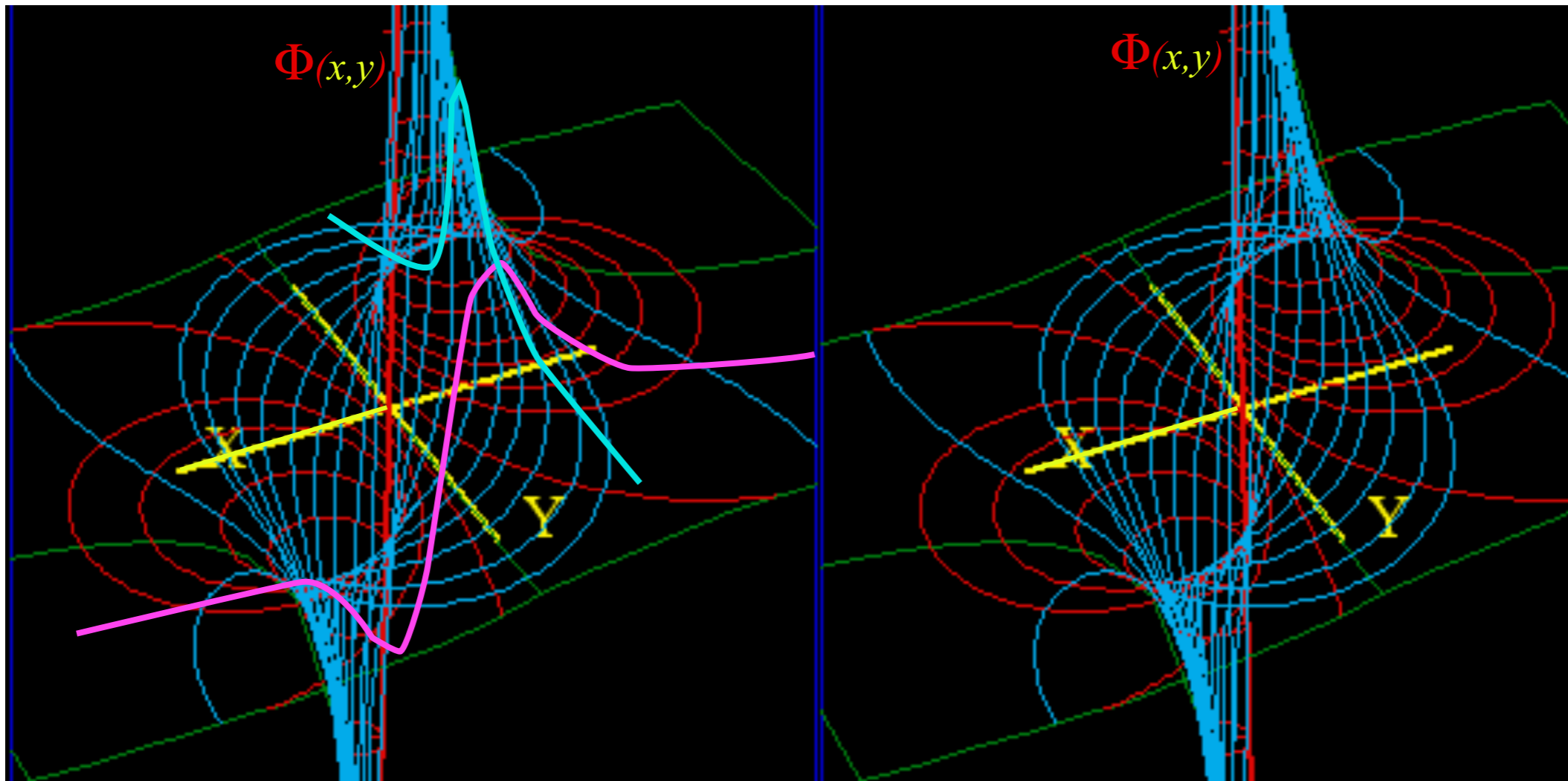
Fig. 10.11 Dipole \mathbf{F} -field $f(z) = 1/z^2$ and scalar potential ($\Phi = \text{const.}$)-circles orthogonal to ($A = \text{const.}$)-circles.



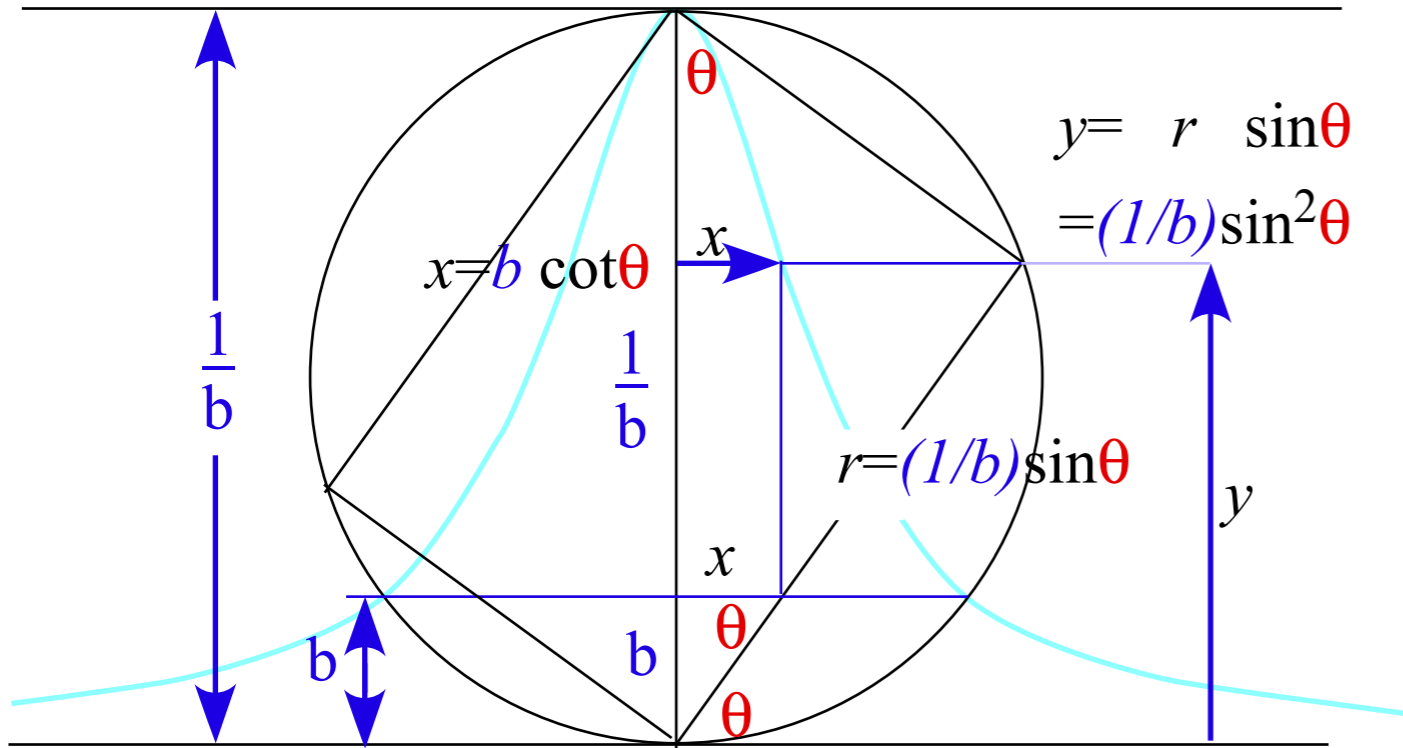
From: Fig. 1.10.12



From: Fig. 1.10.12



From: Fig. 1.10.12

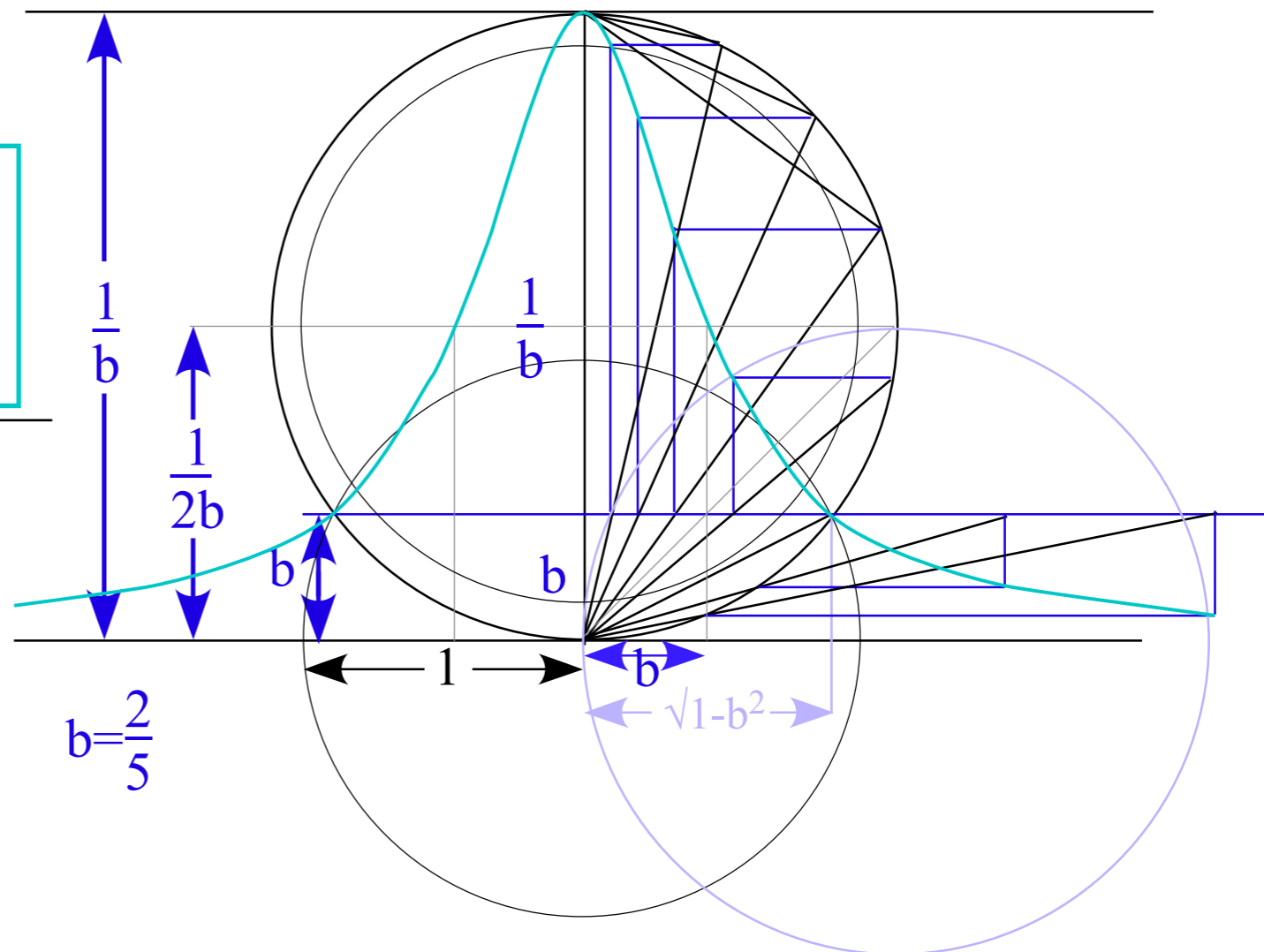


$$x^2 = b^2 \cot^2 \theta = b^2 \frac{\cos^2 \theta}{\sin^2 \theta} = b^2 \frac{1 - \sin^2 \theta}{\sin^2 \theta} = \frac{b^2}{\sin^2 \theta} - b^2$$

$$x^2 + b^2 = \frac{b^2}{\sin^2 \theta} = \frac{b}{y}$$

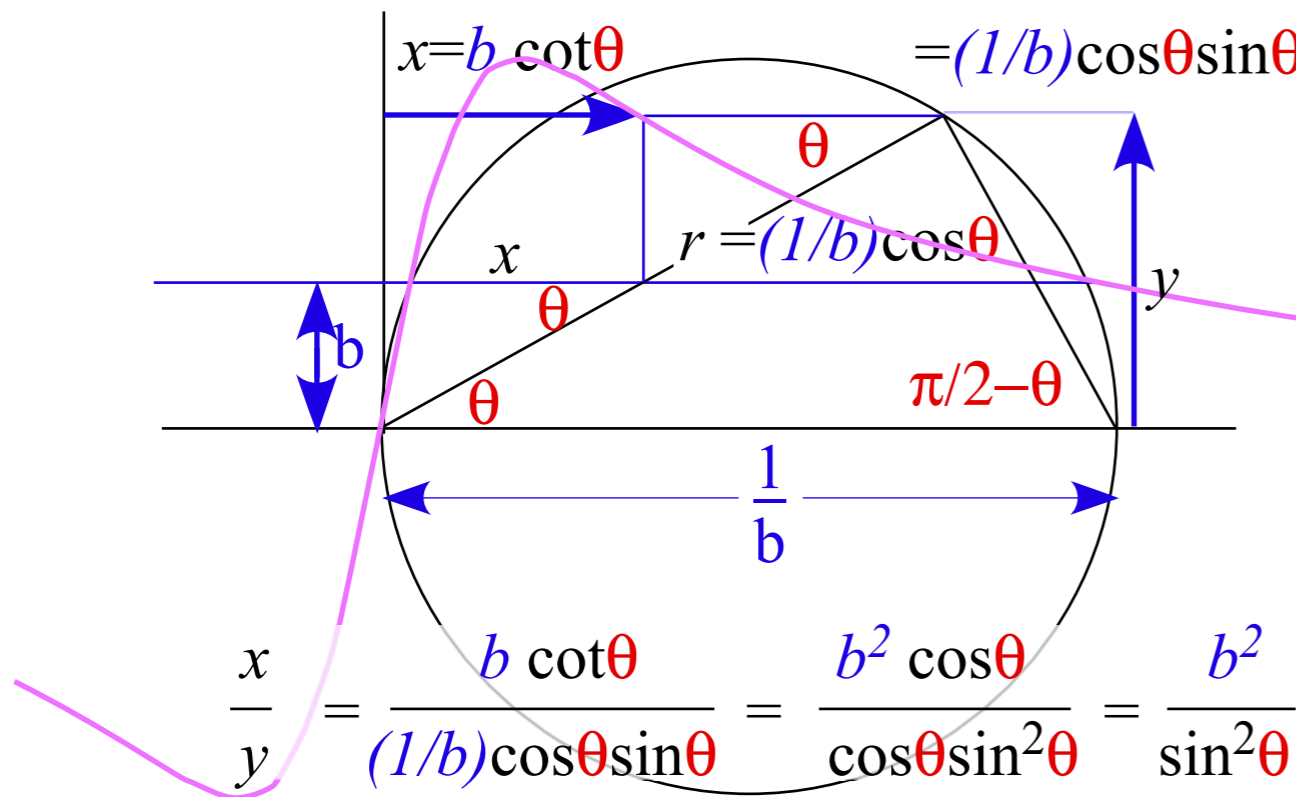
$$y = \frac{b}{x^2 + b^2}$$

*Common Lorentzian function I.
(imaginary "absorbative" part)*



$$y = r \sin \theta$$

$$= (1/b) \cos \theta \sin \theta$$



$$x^2 + b^2 = \frac{b^2}{\sin^2 \theta} = \frac{x}{y}$$

$$y = \frac{x}{x^2 + b^2}$$
 Common Lorentzian function II.
 (real "refractory" part)

