Lecture 13.5
Mon. 10.14.2019

Complex Variables, Series, and Field Coordinates III.

(Ch. 10 of Unit 1)

1. The Story of e (A Tale of Great $Interest$)
How good are those power series?
Taylor-Maclaurin series, imaginary interest, and complex exponentials
2. What good are complex exponentials?
Easy trig
Easy 2D vector analysis
Easy oscillator phase analysis
Easy rotation and "dot" or "cross" products
3. Easy 2D vector calculus
Easy 2D vector derivatives
Easy 2D source-free field theory
Easy 2D vector field-potential theory
4. Riemann-Cauchy relations (What's analytic? What's not?)
Lect. 13 ended here
Easy 2D curvilinear coordinate discovery
Easy 2D circulation and flux integrals
Lect 13.5 starts here p.30
Easy 2D monopole, dipole, and 2^n-pole analysis
(Preview of Unit 4.)
Easy 2^n-multipole field and potential expansion
Easy stereo-projection visualization
Cauchy integrals, Laurent-Maclaurin series
5. Mapping and Non-analytic 2D source field analysis

1. Complex numbers provide "automatic trigonometry"

2. Complex numbers add like vectors.

3. Complex exponentials $Ae^{-i\omega t}$ track position and velocity using Phasor Clock.

4. Complex products provide 2D rotation operations.

5. Complex products provide 2D "dot"(•) and "cross"(x) products.

(Review of topics in Lect.12)

6. Complex derivative contains "divergence"($\nabla \cdot \mathbf{F}$) and "curl"($\nabla \times \mathbf{F}$) of 2D vector field

7. Invent source-free 2D vector fields [$\nabla \cdot \mathbf{F}=0$ and $\nabla \times \mathbf{F}=0$]

8. Complex potential ϕ contains "scalar"($\mathbf{F}=\nabla\Phi$) and "vector"($\mathbf{F}=\nabla \times \mathbf{A}$) potentials

The half-n'-half results: (Riemann-Cauchy Derivative Relations)

9. Complex potentials define 2D Orthogonal Curvilinear Coordinates (OCC) of field

10. Complex integrals $\int f(z)dz$ count 2D "circulation"($\int \mathbf{F}\cdot d\mathbf{r}$) and "flux"($\int \mathbf{F} \times d\mathbf{r}$)

11. Complex integrals define 2D monopole fields and potentials

12. Complex derivatives give 2D dipole fields

13. More derivatives give 2D 2^N-pole fields…

14. ...and 2^N-pole multipole expansions of fields and potentials...

15. ...and Laurent Series...

16. ...and non-analytic source analysis.

17. …and mapping…

This Lecture's Reference Link Listing

Web Resources - front page
UAF Physics UTube channel

Quantum Theory for the Computer Age
Principles of Symmetry, Dynamics, and Spectroscopy
Classical Mechanics with a Bang!
Modern Physics and its Classical Foundations

2017 Group Theory for QM
2018 Adv CM
2018 AMOP
2019 Advanced Mechanics

Lectures #12 through #13.2

In reverse order

Pirelli Relativity Challenge (Introduction level) - *Visualizing Waves:*
Using Earth as a clock,
Tesla's AC Phasors ,
Phasors using complex numbers.
CM wBang Unit 1 - Chapter 10: Calculus of exponentials, logarithms, and complex fields, pdf_page=135
RelaWavity Web Simulation - Unit Circle and Hyperbola (Mixed labeling)
Smith Chart, Invented by Phillip H. Smith (1905-1987)

Select, exciting, and related Research & Articles of Interest

(Many of these may be just beyond this course, but are included to lend added insight):

Clifford_Algebra_And_The_Projective_Model_Of_Homogeneous_Metric_Spaces - Foundations - Sokolov-x-2013
Geometric Algebra 3 - Complex Numbers - MacDonald-yt-2015
Biquaternion -Complexified Quaternion- Roots of -1 - Sangwine-x-2015
An_Introduction_to_Clifford_Algebras_and_Spinors_-_Vaz-Rocha-op-2016
Unified View on Complex Numbers and Quaternions- Bongardt-wcmms-2015
Complex Functions and the Cauchy-Riemann Equations - complex2 - Friedman-columbia-2019

Smith Chart, Invented by Phillip H. Smith (1905-1987)

An assist from *Physics Girl* (YouTube Channel):

Posted this year:
How to Make VORTEX RINGS in a Pool

Crazy pool vortex (new inclusion with more background)
Crazy pool vortex - pg-yt-2014

Posting with the best visuals:
Fun with Vortex Rings in the Pool - pg-yt-2014

She covers it beautifully!

Excerpts (Page 44-47 in *Preliminary Draft*) from the
Geometric Algebra- A Guided Tour through Space and Time - Reimer-www-2019

Past Articles of Interest:

An_sp-hybridized_Molecular_Carbon_Allotrope-_cyclo-18-carbon_-_Kaiser-s-2019
An_Atomic-Scale_View_of_Cyclocarbon_Synthesis_-_Maier-s-2019
Discovery_Of_Topological_Weyl_Fermion_Lines_And_Drumhead_Surface_States_in_a_Room_Temperature_Magnet_-_Belopolski-s-2019
"Weyl"ing_away_Time-reversal_Symmetry_-_Neto-s-2019
Non-Abelian_Band_Topology_in_Noninteracting_Metals_-_Wu-s-2019
What_Industry_Can_Teach_Academia_-_Mao-s-2019
Rovibrational_quantum_state_resolution_of_the_C60_fullerene_-_Changala-Ye-s-2019 (Alt)
A_Degenerate_Fermi_Gas_of_Polar_molecules_-_DeMarco-s-2019

Running Reference Link Listing

Lectures #11 through #7

In reverse order

Eric J Heller Gallery:

Main portal, Consonance and Dissonance II, Bessel 21, Chladni

The Semiclassical Way to Molecular Spectroscopy - Heller-acs-1981
Quantum_dynamical_tunneling_in_bound_states_-_Davis-Heller-jcp-1981

Pendulum Web Simulation
Cycloidulum Web Simulation

Links to previous lecture: *Page=74*, *Page=75*, *Page=79*

Pendulum Web Sim

Cycloidulum Web Sim

***JerkIt* Web Simulations:** *Basic/Generic: Inverted*, *FVPlot*

CMwithBang Lecture 8, page=20

WWW.sciencenewsforstudents.org: Cassini - Saturnian polar vortex

"RelaWavity" Web Simulations:
- 2-CW laser wave, Lagrangian vs Hamiltonian,
- Physical Terms Lagrangian L(u) vs Hamiltonian H(p)

CoulIt Web Simulation of the Volcanoes of Io

BohrIt Multi-Panel Plot:
- Relativistically shifted Time-Space plots of 2 CW light waves

BoxIt Web Simulations:
- *Generic/Default*
- *Most Basic A-Type*
- *Basic A-Type w/reference lines*
- *Basic A-Type A-Type with Potential energy*
- *A-Type with Potential energy and Stokes Plot*
- *A-Type w/3 time rates of change*
- *A-Type w/3 time rates of change with Stokes Plot*
- *B-Type (A=1.0, B=-0.05, C=0.0, D=1.0)*

RelaWavity Web Elliptical Motion Simulations:
- *Orbits with b/a=0.125*
- *Orbits with b/a=0.5*
- *Orbits with b/a=0.7*
- *Exegesis with b/a=0.125*
- *Exegesis with b/a=0.5*
- *Exegesis with b/a=0.7*
- *Contact Ellipsometry*

***CoulIt* Web Simulations:**
- *Basic/Generic*
- *Exploding Starlet*
- *Volcanoes of Io (Color Quantized)*

***JerkIt* Web Simulations:**
- *Basic/Generic*
- *Catcher in the Eye - IHO with Linear Hooke perturbation - Force-potential-Velocity Plot*

***OscillatorPE* Web Simulation:**
- *Coulomb-Newton-Inverse_Square*,
- *Hooke-Isotropic Harmonic*,
- *Pendulum-Circular_Constraint*

AMOP Ch 0 Space-Time Symmetry - 2019
Seminar at Rochester Institute of Optics, Aux. slides-2018

NASA Astronomy Picture of the Day -
Io: The Prometheus Plume *(Just Image)*
NASA Galileo - *Io's Alien Volcanoes*
New Horizons - *Volcanic Eruption Plume on Jupiter's moon IO*
NASA Galileo - A Hawaiian-Style Volcano on Io

Pirelli Site: Phasors animimation
CMwithBang Lecture #6, page=70 (9.10.18)

Select, exciting, and related Research & Articles of Interest:

Burning a hole in reality—design for a new laser may be powerful enough to pierce space-time - Sumner-KOS-2019
Trampoline mirror may push laser pulse through fabric of the Universe - Lee-ArsTechnica-2019
Achieving_Extreme_Light_Intensities_using_Optically_Curved_Relativistic_Plasma_Mirrors_-_Vincenti-prl-2019

A_Soft_Matter_Computer_for_Soft_Robots_-_Garrad-sr-2019
Correlated_Insulator_Behaviour_at_Half-Filling_in_Magic-Angle_Graphene_Superlattices_-_cao-n-2018

*Sorting_ultracold_atoms_in_a_three-dimensional_optical_lattice_in_a_*
realization_of_Maxwell's_Demon - Kumar-n-2018
Synthetic_three-dimensional_atomic_structures_assembled_atom_by_atom - Barredo-n-2018

Older ones:
Wave-particle_duality_of_C60_molecules - Arndt-ltn-1999
Optical_Vortex_Knots_-_One_Photon__At_A_Time - Tempone-Wiltshire-Sr-2018
Baryon_Deceleration_by_Strong_Chromofields_in_Ultrarelativistic_,
Nuclear_Collisions - Mishustin-PhysRevC-2007, *APS Link & Abstract*
Hadronic Molecules - Guo-x-2017
Hidden-charm_pentaquark_and_tetraquark_states - Chen-pr-2016

Running Reference Link Listing

Lectures #6 through #1

In reverse order

RelaWavity Web Simulation: Contact Ellipsometry
BoxIt Web Simulation: Elliptical Motion (A-Type)
CMwBang Course: Site Title Page
Pirelli Relativity Challenge: Describing Wave Motion With Complex Phasors
UAF Physics UTube channel

Velocity Amplification in Collision Experiments Involving Superballs - Harter, 1971
MIT OpenCourseWare: High School/Physics/Impulse and Momentum
Hubble Site: Supernova - SN 1987A

BounceItIt Web Animation - Scenarios:
49:1 y vs t, *49:1 V2 vs V1*, *1:500:1 - 1D Gas Model w/ faux restorative force (Cool)*,
1:500:1 - 1D Gas (Warm), *1:500:1 - 1D Gas Model (Cool, Zoomed in)*,
Farey Sequence - Wolfram
Fractions - Ford-AMM-1938
Monstermash BounceItIt Animations:
1000:1 - V2 vs V1, *1000:1 with t vs x - Minkowski Plot*
Quantum Revivals of Morse Oscillators and Farey-Ford Geometry - Li-Harter-2013
Quantum_Revivals_of_Morse_Oscillators_and_Farey-Ford_Geometry - Li-Harter-cpl-2015
Quant. Revivals of Morse Oscillators and Farey-Ford Geom. - Harter-Li-CPL-2015 (Publ.)
Velocity_Amplification_in_Collision_Experiments_Involving_Superballs-Harter-1971
WaveIt Web Animation - Scenarios:
Quantum_Carpet, *Quantum_Carpet_wMBars*,
Quantum_Carpet_BCar, *Quantum_Carpet_BCar_wMBars*
Wave Node Dynamics and Revival Symmetry in Quantum Rotors - Harter-JMS-2001
Wave Node Dynamics and Revival Symmetry in Quantum Rotors - Harter-jms-2001 (Publ.)

BounceIt Web Animation - Scenarios:
Generic Scenario: 2-Balls dropped no Gravity (7:1) - V vs V Plot (Power=4)
1-Ball dropped w/Gravity=0.5 w/Potential Plot: Power=1, Power=4
7:1 - V vs V Plot: Power=1
3-Ball Stack (10:3:1) w/Newton plot (y vs t) - Power=4
3-Ball Stack (10:3:1) w/Newton plot (y vs t) - Power=1
3-Ball Stack (10:3:1) w/Newton plot (y vs t) - Power=1 w/Gaps
4-Ball Stack (27:9:3:1) w/Newton plot (y vs t) - Power=4
4-Newton's Balls (1:1:1:1) w/Newtonian plot (y vs t) - Power=4 w/Gaps
6-Ball Totally Inelastic (1:1:1:1:1:1) w/Gaps: Newtonian plot (t vs x), V6 vs V5 plot
5-Ball Totally Inelastic Pile-up w/ 5-Stationary-Balls - Minkowski plot (t vs x1) w/Gaps
1-Ball Totally Inelastic Pile-up w/ 5-Stationary-Balls - Vx2 vs Vx1 plot w/Gaps

BounceIt Dual plots
m_1:m_2 = 3:1
v_2 vs v_1 and V_2 vs V_1, *(v_1, v_2)=(1, 0.1)*, *(v_1, v_2)=(1, 0)*
y_2 vs y_1 plots: *(v_1, v_2)=(1, 0.1)*, *(v_1, v_2)=(1, 0)*, *(v_1, v_2)=(1, -1)*
Estrangian plot *V_2 vs V_1:* *(v_1, v_2)=(0, 1)*, *(v_1, v_2)=(1, -1)*
m_1:m_2 = 4:1
v2 vs v1, *y2 vs y1*
m_1:m_2 = 100:1, *(v_1, v_2)=(1, 0):* *V2 vs V1 Estrangian plot*, *y2 vs y1 plot*

With g=0 and 70:10 mass ratio
With non zero g, velocity dependent damping and mass ratio of 70:35
M_1=49, M_2=1 with Newtonian time plot
M_1=49, M_2=1 with V_2 vs V_1 plot
Example with friction
Low force constant with drag displaying a Pass-thru, *Fall-Thru*, *Bounce-Off*
m1:m2= 3:1 and (v1, v2) = (1, 0) Comparison with Estrangian

AJP article on superball dynamics
AAPT Summer Reading List
Scitation.org - AIP publications
HarterSoft Youtube Channel

X2 paper: *Velocity Amplification in Collision Experiments Involving Superballs - Harter, et. al. 1971 (pdf)*
Car Collision Web Simulator: https://modphys.hosted.uark.edu/markup/CMMotionWeb.html
Superball Collision Web Simulator: https://modphys.hosted.uark.edu/markup/BounceItWeb.html; with Scenarios: 1007
BounceIt web simulation with g=0 and 70:10 mass ratio
With non zero g, velocity dependent damping and mass ratio of 70:35
Elastic Collision Dual Panel Space vs Space: Space vs Time (Newton), Time vs. Space(Minkowski)
Inelastic Collision Dual Panel Space vs Space: Space vs Time (Newton), Time vs. Space(Minkowski)
Matrix Collision Simulator: M_1=49, M_2=1 V_2 vs V_1 plot <<*Under Construction*>>

More Advanced QM and classical references will *soon* be available through our: ***Mechanics References Page***
(Now in Development)

6. Complex derivative contains "divergence"($\nabla\cdot\mathbf{F}$) and "curl"($\nabla\times\mathbf{F}$) of 2D vector field

Relation of (z,z^*) to $(x=\text{Re}z, y=\text{Im}z)$ defines a z-derivative $\frac{df}{dz}$ and "star" z^*-derivative. $\frac{df}{dz^*}$

$$z = x + iy \qquad x = \frac{1}{2}(z + z^*)$$

$$z^* = x - iy \qquad y = \frac{1}{2i}(z - z^*)$$

$$\frac{df}{dz} = \frac{\partial x}{\partial z}\frac{\partial f}{\partial x} + \frac{\partial y}{\partial z}\frac{\partial f}{\partial y} = \frac{1}{2}\frac{\partial f}{\partial x} - \frac{i}{2}\frac{\partial f}{\partial y}$$

$$\frac{df}{dz^*} = \frac{\partial x}{\partial z^*}\frac{\partial f}{\partial x} + \frac{\partial y}{\partial z^*}\frac{\partial f}{\partial y} = \frac{1}{2}\frac{\partial f}{\partial x} + \frac{i}{2}\frac{\partial f}{\partial y}$$

$$\frac{d}{dz} = \frac{1}{2}\frac{\partial}{\partial x} - \frac{i}{2}\frac{\partial}{\partial y}$$

$$\frac{d}{dz^*} = \frac{1}{2}\frac{\partial}{\partial x} + \frac{i}{2}\frac{\partial}{\partial y}$$

Derivative chain-rule shows real part of $\frac{df}{dz}$ has 2D divergence $\nabla\bullet\mathbf{f}$ and imaginary part has curl $\nabla\times\mathbf{f}$.

$$\frac{df}{dz} = \frac{d}{dz}(f_x + i f_y) = \frac{1}{2}\left(\frac{\partial f}{\partial x} - i\frac{\partial f}{\partial y}\right)(f_x + i f_y) = \frac{1}{2}\underbrace{\left(\frac{\partial f_x}{\partial x} + \frac{\partial f_y}{\partial y}\right)} + \frac{i}{2}\underbrace{\left(\frac{\partial f_y}{\partial x} - \frac{\partial f_x}{\partial y}\right)} = \frac{1}{2}\nabla\bullet\mathbf{f} + \frac{i}{2}|\nabla\times\mathbf{f}|_{Z\perp(x,y)}$$

7. Invent source-free 2D vector fields [$\nabla\cdot\mathbf{F}=0$ and $\nabla\times\mathbf{F}=0$]

We can invent *source-free 2D vector fields* that are both *zero-divergence* and *zero-curl*. Take any function $f(z)$, conjugate it (change all i's to $-i$) to give $f^*(z^*)$ for which $\boxed{\frac{df^*}{dz} = 0}$

For example: if $f(z)=a\cdot z$ then $f^*(z^*)=a\cdot z^*=a(x-iy)$ is not function of z so it has *zero z-derivative*.

$\mathbf{F}=(F_x,F_y)=(f^*_x,f^*_y)=(a\cdot x,-a\cdot y)$ *has zero divergence:* $\nabla\bullet\mathbf{F}=0$ *and has zero curl:* $|\nabla\times\mathbf{F}|=0$.

$$\nabla\bullet\mathbf{F} = \frac{\partial F_x}{\partial x} + \frac{\partial F_y}{\partial y} = \frac{\partial(ax)}{\partial x} + \frac{\partial F(-ay)}{\partial y} = 0 \qquad |\nabla\times\mathbf{F}|_{Z\perp(x,y)} = \frac{\partial F_y}{\partial x} - \frac{\partial F_x}{\partial y} = \frac{\partial(-ay)}{\partial x} - \frac{\partial F(ax)}{\partial y} = 0$$

A *DFL* field $\mathbf{F}$ (*Divergence-Free-Laminar*)

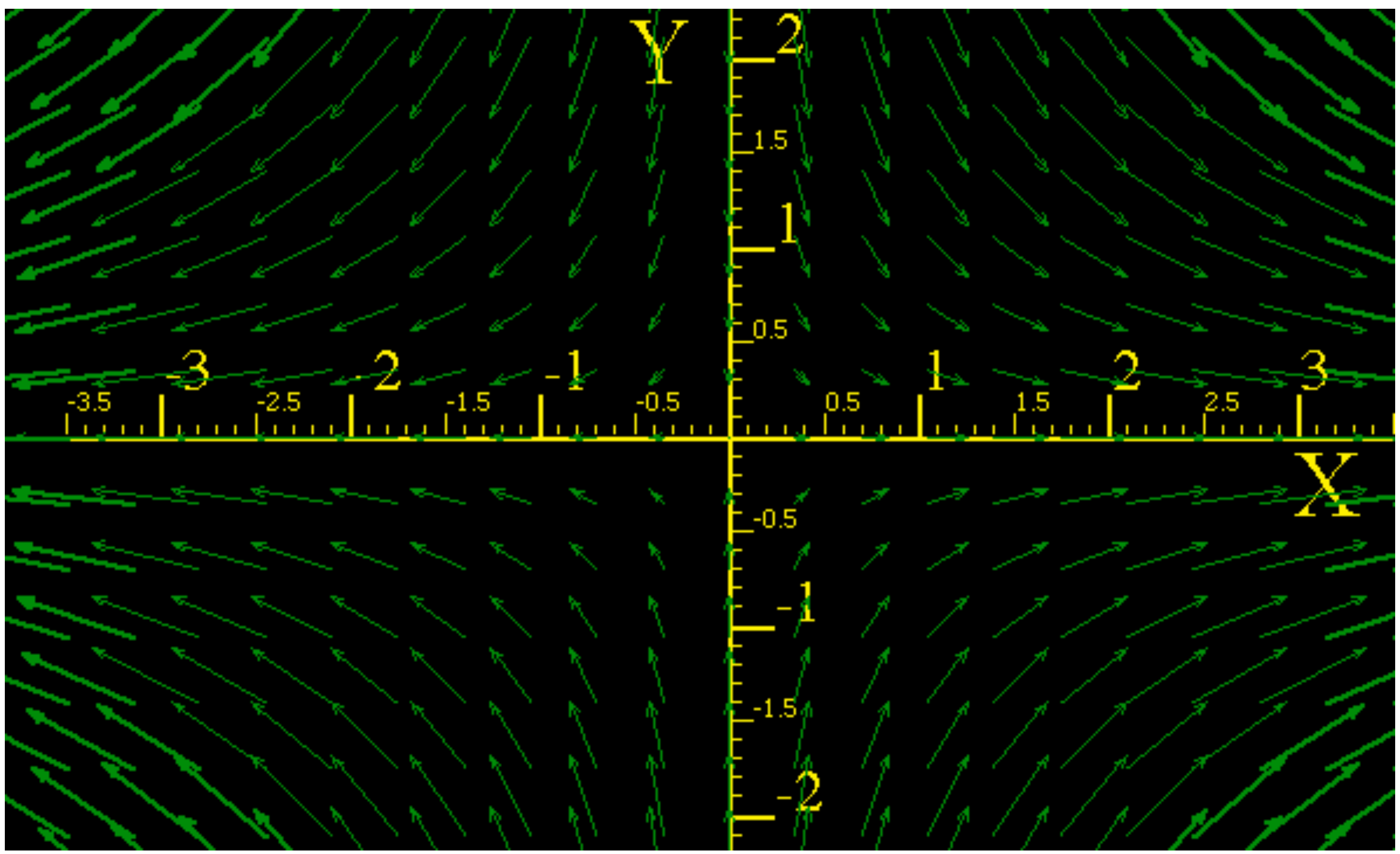
7. Invent source-free 2D vector fields [$\nabla\cdot\mathbf{F}=0$ and $\nabla\times\mathbf{F}=0$]

We can invent *source-free 2D vector fields* that are both *zero-divergence* and *zero-curl*.
Take any function *f(z)*, conjugate it (change all *i*'s to –*i*) to give $f^*(z^*)$ for which $\boxed{\frac{df^*}{dz}=0}$.

For example: if $f(z)=a\cdot z$ then $f^*(z^*)=a\cdot z^*=a(x-iy)$ is not function of *z* so it has *zero z-derivative*.

$\mathbf{F}=(F_x,F_y)=(f^*_x,f^*_y)=(a\cdot x,-a\cdot y)$ has *zero divergence*: $\nabla\bullet\mathbf{F}=0$ and has *zero curl*: $|\nabla\times\mathbf{F}|=0$.

$$\nabla\bullet\mathbf{F}=\frac{\partial F_x}{\partial x}+\frac{\partial F_y}{\partial y}=\frac{\partial(ax)}{\partial x}+\frac{\partial F(-ay)}{\partial y}=0 \qquad |\nabla\times\mathbf{F}|_{Z\perp(x,y)}=\frac{\partial F_y}{\partial x}-\frac{\partial F_x}{\partial y}=\frac{\partial(-ay)}{\partial x}-\frac{\partial F(ax)}{\partial y}=0$$



*precursor to
Unit 1
Fig. 10.7*

$\mathbf{F}=(f^*_x,f^*_y)=(a\cdot x,-a\cdot y)$ is a *divergence-free laminar (DFL)* field.

8. Complex potential ϕ contains "scalar"($\mathbf{F}=\nabla\Phi$) and "vector"($\mathbf{F}=\nabla x\mathbf{A}$) potentials

Any *DFL* field $\mathbf{F}$ is a gradient of a *scalar potential field* Φ or a curl of a *vector potential field* $\mathbf{A}$.

$$\mathbf{F}=\nabla\Phi \qquad\qquad \mathbf{F}=\nabla\times\mathbf{A}$$

A *complex potential* $\phi(z)=\Phi(x,y)+i\mathbf{A}(x,y)$ exists whose z-derivative is $f(z)=d\phi/dz$.

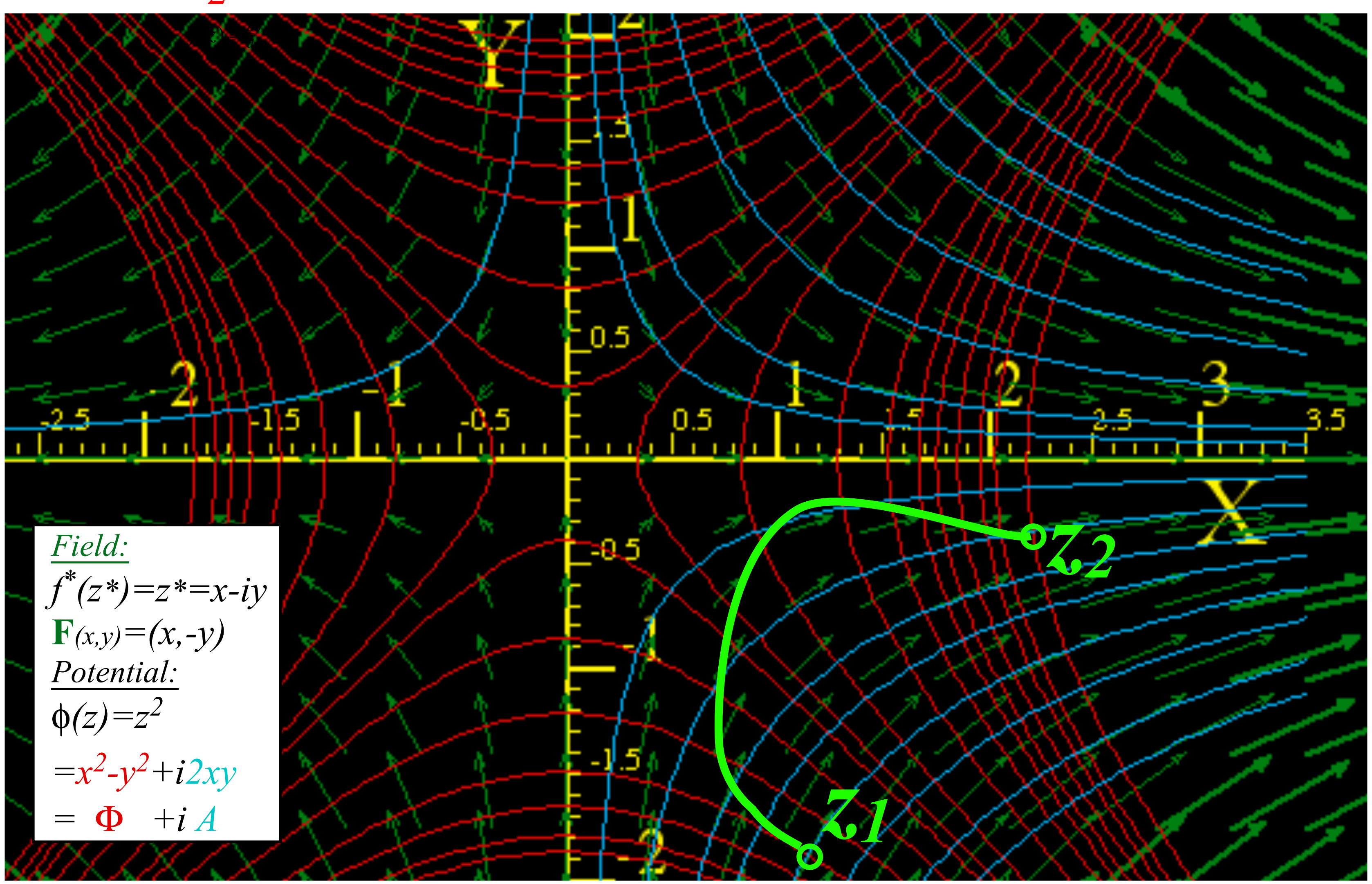
Its complex conjugate $\phi^*(z^*)=\Phi(x,y)-i\mathbf{A}(x,y)$ has z^*-derivative $f^*(z^*)=d\phi^*/dz^*$ giving *DFL* field $\mathbf{F}$.

To find $\phi=\Phi+i\mathbf{A}$ integrate $f(z)=a\cdot z$ to get ϕ and isolate real (Re $\phi=\Phi$) and imaginary (Im $\phi=\mathbf{A}$) parts.

$$f(z)=\frac{d\phi}{dz} \quad\Rightarrow\quad \phi=\overbrace{\Phi}^{}+i\,\mathbf{A}=\int f\cdot dz=\int az\cdot dz=\tfrac{1}{2}\,az^2=\tfrac{1}{2}\,a(x+iy)^2$$

$$=\underbrace{\tfrac{1}{2}\,a(x^2-y^2)}_{\Phi}+i\,\underbrace{axy}_{\mathbf{A}}$$

BONUS!
Get a free coordinate system!



Unit 1
Fig. 10.7

The $(\Phi,\mathbf{A})$ grid is a GCC coordinate system*:

$q^1=\Phi=(x^2-y^2)/2=$ *const.*

$q^2=\mathbf{A}=(xy)=$ *const.*

*Actually it's OCC.

*8. (contd.) Complex potential ϕ contains "scalar"(**F=**$\nabla\Phi$) and "vector"(**F=**∇x**A**) potentials*

...and either one (or half-n'-half!) works just as well.

Derivative $\frac{d\phi^*}{dz^*}$ has 2D gradient $\nabla\Phi = \begin{pmatrix} \frac{\partial\Phi}{\partial x} \\ \frac{\partial\Phi}{\partial y} \end{pmatrix}$ of scalar $\mathbf{\Phi}$ and curl $\nabla\times\mathbf{A} = \begin{pmatrix} \frac{\partial \mathbf{A}}{\partial y} \\ -\frac{\partial \mathbf{A}}{\partial y} \end{pmatrix}$ of vector **A** (*and they're equal!*)

The half-n'-half result

$$\frac{d}{dz^*}\phi^* = \frac{d}{dz^*}(\mathbf{\Phi} - i\mathbf{A}) = \frac{1}{2}\left(\frac{\partial}{\partial x} + i\frac{\partial}{\partial y}\right)(\mathbf{\Phi} - i\mathbf{A}) = \frac{1}{2}\left(\frac{\partial\mathbf{\Phi}}{\partial x} + i\frac{\partial\mathbf{\Phi}}{\partial y}\right) + \frac{1}{2}\left(\frac{\partial\mathbf{A}}{\partial y} - i\frac{\partial\mathbf{A}}{\partial x}\right) = \frac{1}{2}\nabla\mathbf{\Phi} + \frac{1}{2}\nabla\times\mathbf{A}$$

Note, *mathematician definition* of force field **F**=+$\nabla\mathbf{\Phi}$ replaces usual physicist's definition **F**=-$\nabla\mathbf{\Phi}$

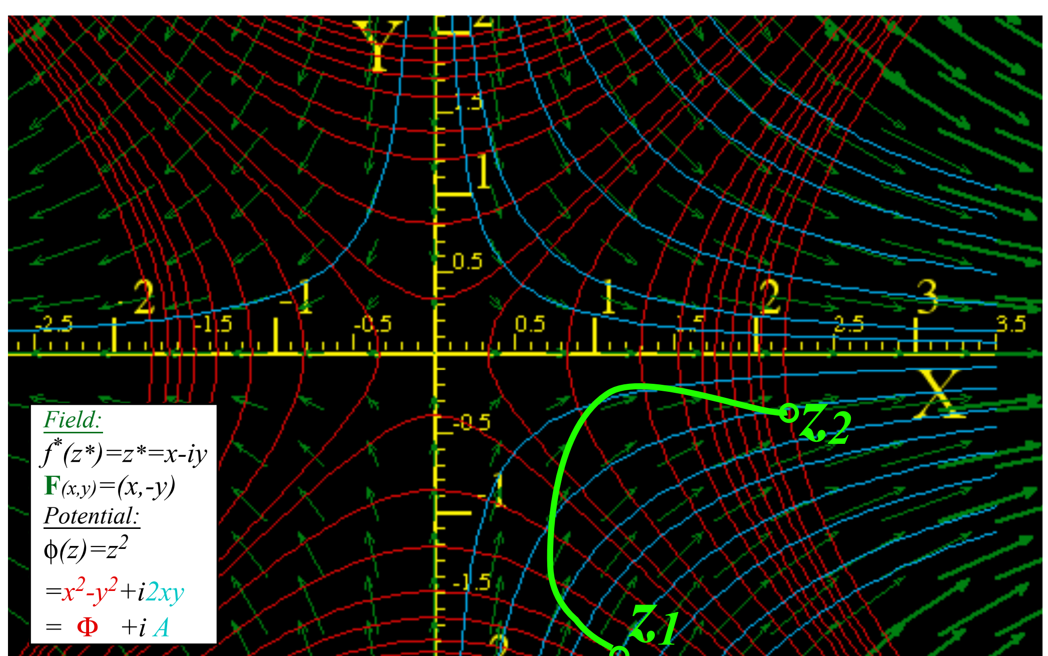
Given ϕ:
$$\boxed{\begin{aligned} \phi &= \mathbf{\Phi} + i\,\mathbf{A} \\ &= \tfrac{1}{2}a(x^2 - y^2) + i\,axy \end{aligned}}$$

The half-n'-half result

find:
$$\nabla\mathbf{\Phi} = \begin{pmatrix} \frac{\partial\mathbf{\Phi}}{\partial x} \\ \frac{\partial\mathbf{\Phi}}{\partial y} \end{pmatrix} = \begin{pmatrix} \frac{\partial}{\partial x}\frac{a}{2}(x^2 - y^2) \\ \frac{\partial}{\partial y}\frac{a}{2}(x^2 - y^2) \end{pmatrix} = \begin{pmatrix} ax \\ -ay \end{pmatrix} = \mathbf{F}$$

or find:
$$\nabla\times\mathbf{A} = \begin{pmatrix} \frac{\partial\mathbf{A}}{\partial y} \\ -\frac{\partial\mathbf{A}}{\partial x} \end{pmatrix} = \begin{pmatrix} \frac{\partial}{\partial y}axy \\ -\frac{\partial}{\partial x}axy \end{pmatrix} = \begin{pmatrix} ax \\ -ay \end{pmatrix} = \mathbf{F}$$

Scalar *static potential lines* $\mathbf{\Phi}$=*const.* and vector *flux potential lines* **A**=*const.* define *DFL field-net*.



The half-n'-half results are called Riemann-Cauchy Derivative Relations

$$\frac{\partial\mathbf{\Phi}}{\partial x} = \frac{\partial\mathbf{A}}{\partial y} \quad \text{is:} \quad \frac{\partial \mathrm{Re} f(z)}{\partial x} = \frac{\partial \mathrm{Im} f(z)}{\partial y}$$

$$\frac{\partial\mathbf{\Phi}}{\partial y} = -\frac{\partial\mathbf{A}}{\partial x} \quad \text{is:} \quad \frac{\partial \mathrm{Re} f(z)}{\partial y} = -\frac{\partial \mathrm{Im} f(z)}{\partial x}$$

Review (z,z) to (x,y) transformation relations* (Review of topics in Lect.12)

$$z = x + iy \qquad x = \tfrac{1}{2}(z + z^*)$$
$$z^* = x - iy \qquad y = \tfrac{1}{2i}(z - z^*)$$

$$\frac{df}{dz} = \frac{\partial x}{\partial z}\frac{\partial f}{\partial x} + \frac{\partial y}{\partial z}\frac{\partial f}{\partial y} = \frac{1}{2}\frac{\partial f}{\partial x} + \frac{1}{2i}\frac{\partial f}{\partial y} = \frac{1}{2}\left(\frac{\partial}{\partial x} - i\frac{\partial}{\partial y}\right)f$$
$$\frac{df}{dz^*} = \frac{\partial x}{\partial z^*}\frac{\partial f}{\partial x} + \frac{\partial y}{\partial z^*}\frac{\partial f}{\partial y} = \frac{1}{2}\frac{\partial f}{\partial x} - \frac{1}{2i}\frac{\partial f}{\partial y} = \frac{1}{2}\left(\frac{\partial}{\partial x} + i\frac{\partial}{\partial y}\right)f$$

Criteria for a field function $f = f_x(x,y) + i f_y(x,y)$ to be an ***analytic function f(z)*** *of $z=x+iy$:*
First, f(z) must <u>not</u> be a function of $z^*=x-iy$, that is: $\frac{df}{dz^*} = 0$
This implies f(z) satisfies differential equations known as the ***Riemann-Cauchy conditions***

$$\frac{df}{dz^*} = 0 = \frac{1}{2}\left(\frac{\partial}{\partial x} + i\frac{\partial}{\partial y}\right)(f_x + i f_y) = \frac{1}{2}\left(\frac{\partial f_x}{\partial x} - \frac{\partial f_y}{\partial y}\right) + \frac{i}{2}\left(\frac{\partial f_y}{\partial x} + \frac{\partial f_x}{\partial y}\right) \text{ implies : } \boxed{\frac{\partial f_x}{\partial x} = \frac{\partial f_y}{\partial y} \quad \text{and : } \quad \frac{\partial f_y}{\partial x} = -\frac{\partial f_x}{\partial y}}$$

$$\frac{df}{dz} = \frac{1}{2}\left(\frac{\partial}{\partial x} - i\frac{\partial}{\partial y}\right)(f_x + i f_y) = \frac{1}{2}\left(\frac{\partial f_x}{\partial x} + \frac{\partial f_y}{\partial y}\right) + \frac{i}{2}\left(\frac{\partial f_y}{\partial x} - \frac{\partial f_x}{\partial y}\right) = \frac{\partial f_x}{\partial x} + i\frac{\partial f_y}{\partial x} = \frac{\partial f_y}{\partial y} - i\frac{\partial f_x}{\partial y} = \frac{\partial}{\partial x}(f_x + i f_y) = \frac{\partial}{\partial iy}(f_x + i f_y)$$

Criteria for a field function $f = f_x(x,y) + i f_y(x,y)$ to be an ***analytic function f(z*)*** *of $z^*=x-iy$:*
First, f(z) must <u>not</u> be a function of $z=x+iy$, that is:* $\frac{df}{dz} = 0$
This implies f(z) satisfies differential equations we call* ***Anti-Riemann-Cauchy conditions***

$$\frac{df}{dz} = 0 = \frac{1}{2}\left(\frac{\partial}{\partial x} - i\frac{\partial}{\partial y}\right)(f_x + i f_y) = \frac{1}{2}\left(\frac{\partial f_x}{\partial x} + \frac{\partial f_y}{\partial y}\right) + \frac{i}{2}\left(\frac{\partial f_y}{\partial x} - \frac{\partial f_x}{\partial y}\right) = \text{implies : } \boxed{\frac{\partial f_x}{\partial x} = -\frac{\partial f_y}{\partial y} \quad \text{and : } \quad \frac{\partial f_y}{\partial x} = \frac{\partial f_x}{\partial y}}$$

$$\frac{df}{dz^*} = \frac{1}{2}\left(\frac{\partial}{\partial x} + i\frac{\partial}{\partial y}\right)(f_x + i f_y) = \frac{1}{2}\left(\frac{\partial f_x}{\partial x} - \frac{\partial f_y}{\partial y}\right) + \frac{i}{2}\left(\frac{\partial f_y}{\partial x} + \frac{\partial f_x}{\partial y}\right) = \frac{\partial f_x}{\partial x} + i\frac{\partial f_y}{\partial x} = -\frac{\partial f_y}{\partial y} + i\frac{\partial f_x}{\partial y} = \frac{\partial}{\partial x}(f_x + i f_y) = -\frac{\partial}{\partial iy}(f_x + i f_y)$$

What's analytic? (...and what's not?)

Example: Q: Is $f(x,y) = 2x + i4y$ an analytic function of $z=x+iy$?

Well, test it using definitions: $z = x + iy$ and: $z^* = x - iy$
or: $x = (z+z^*)/2$ and: $y =-i(z-z^*)/2$

$$f(x,y) = 2x + i4y = 2\,(z+z^*)/2 + i4(-i(z-z^*)/2)$$
$$= z+z^* + (2z-2z^*)$$
$$= 3z-z^*$$

A: **NO!** *It's a function of z and z* so not analytic for either.*

Example 2: Q: Is $r(x,y) = x^2 + y^2$ an analytic function of $z=x+iy$?

A: **NO!** $r(xy)=z^*z$ *is a function of z and z* so not analytic for either.*

Example 3: Q: Is $s(x,y) = x^2-y^2 + 2ixy$ an analytic function of $z=x+iy$?

A: **YES!** $s(xy)=(x+iy)^2 =z^2$ *is analytic function of z.* (Yay!)

4. Riemann-Cauchy conditions *What's analytic? (...and what's <u>not</u>?)*

Easy 2D circulation and flux integrals

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Easy stereo-projection visualization

*9. Complex integrals ∫ f(z)dz count 2D "circulation"(∫**F•dr**) and "flux"(∫**F**x**dr**)*

Integral of $f(z)$ between point z_1 and point z_2 is potential difference $\Delta\phi = \phi(z_2) - \phi(z_1)$

$$\Delta\phi = \phi(z_2) - \phi(z_1) = \int_{z_1}^{z_2} f(z)dz = \underbrace{\Phi(x_2,y_2) - \Phi(x_1,y_1)}_{\Delta\Phi} + i\,[\underbrace{\mathbf{A}(x_2,y_2) - \mathbf{A}(x_1,y_1)}_{\Delta\mathbf{A}}]$$

$$\Delta\phi = \Delta\Phi + i\,\Delta\mathbf{A}$$

In *DFL*-field **F**, $\Delta\phi$ is independent of the integration path $z(t)$ connecting z_1 and z_2.

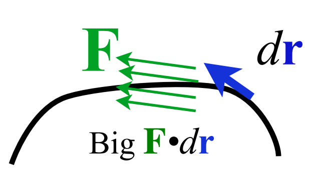
$$\int f(z)dz = \int \left(f^*(z^*)\right)^* dz = \int \left(f^*(z^*)\right)^* (dx + i\,dy) = \int \left(f_x^* + i\,f_y^*\right)^* (dx + i\,dy) = \int \left(f_x^* - i\,f_y^*\right)(dx + i\,dy)$$

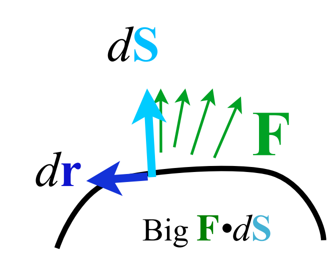
$$= \int (f_x^*\,dx + f_y^*\,dy) + i\int (f_x^*\,dy - f_y^*\,dx)$$

$$= \int \mathbf{F}\bullet d\mathbf{r} \quad + i\int \mathbf{F}\times d\mathbf{r}\bullet\hat{\mathbf{e}}_Z$$

$$= \int \mathbf{F}\bullet d\mathbf{r} \quad + i\int \mathbf{F}\bullet d\mathbf{r}\times\hat{\mathbf{e}}_Z$$

$$= \boxed{\int \mathbf{F}\bullet d\mathbf{r}} \quad \boxed{+ i\int \mathbf{F}\bullet d\mathbf{S}} \qquad \text{where:} \quad d\mathbf{S} = d\mathbf{r}\times\hat{\mathbf{e}}_Z$$





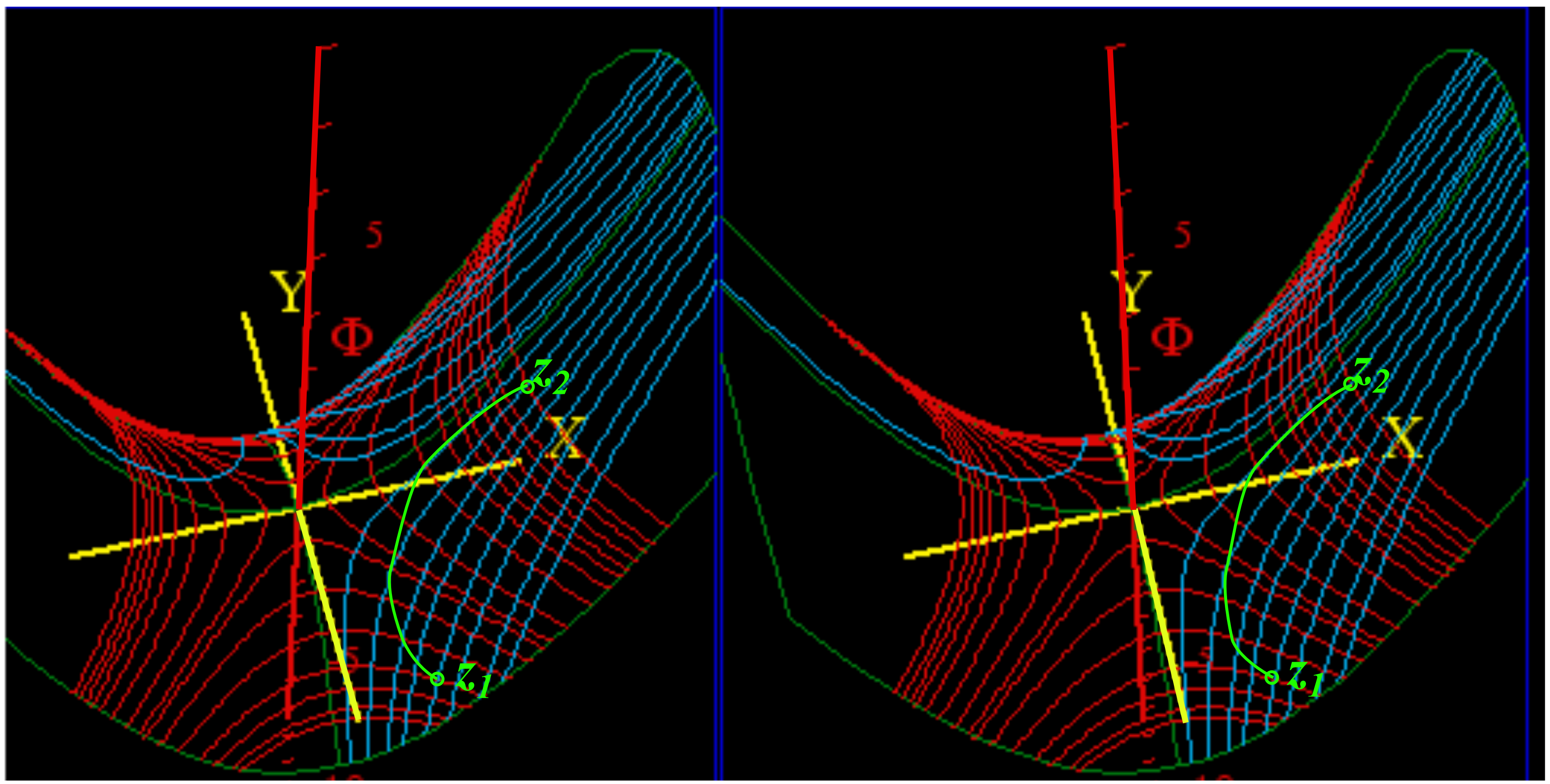
Real part $\int_1^2 \mathbf{F}\bullet d\mathbf{r} = \Delta\Phi$
sums **F** projections *along* path $d\mathbf{r}$ that is, *circulation* on path to get $\Delta\Phi = \Phi(x_2,y_2) - \Phi(x_1,y_1)$

Imaginary part $\int_1^2 \mathbf{F}\bullet d\mathbf{S} = \Delta\mathbf{A}$
sums **F** projection *across* path $d\mathbf{r}$ that is, *flux* thru surface elements $d\mathbf{S} = d\mathbf{r}\times\mathbf{e}_Z$ normal to $d\mathbf{r}$ to get $\Delta\mathbf{A} = \mathbf{A}(x_2,y_2) - \mathbf{A}(x_1,y_1)$

Here the scalar potential $\Phi=(x^2-y^2)/2$ is stereo-plotted vs. (x,y)

The $\Phi=(x^2-y^2)/2=const.$ curves are topography lines

The $A=(xy)=const.$ curves are streamlines normal to topography lines



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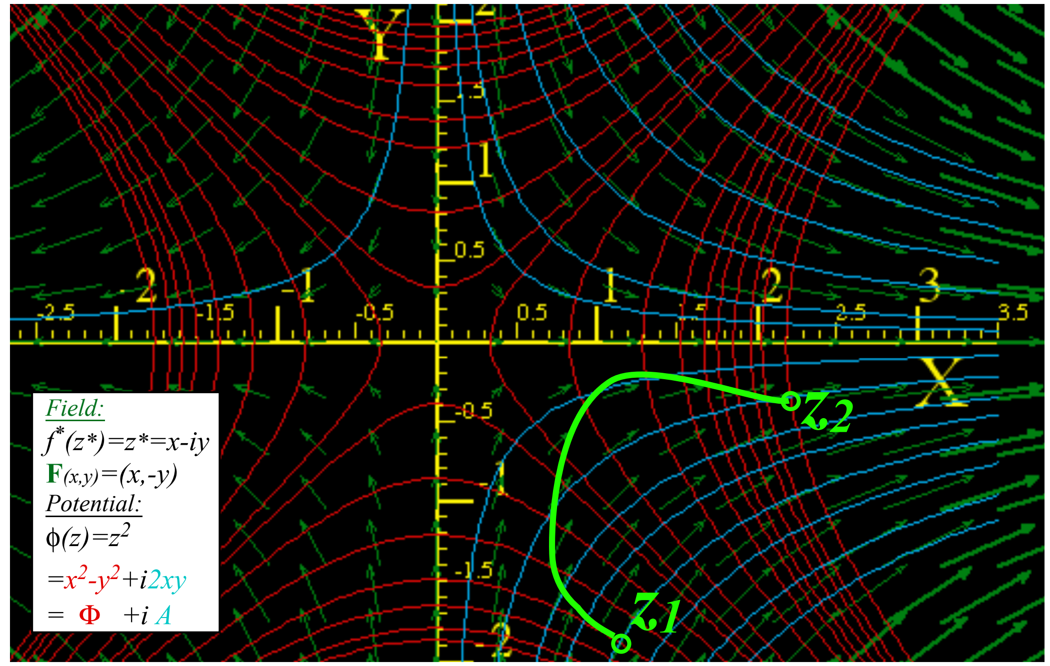
The (Φ, A) grid is a GCC coordinate system*:

$q^1 = \Phi = (x^2 - y^2)/2 = const.$

$q^2 = A = (xy) = const.$

*Actually it's OCC.



$$Kajobian = \begin{pmatrix} \frac{\partial q^1}{\partial x} & \frac{\partial q^1}{\partial y} \\ \frac{\partial q^2}{\partial x} & \frac{\partial q^2}{\partial y} \end{pmatrix} = \begin{pmatrix} \frac{\partial \Phi}{\partial x} & \frac{\partial \Phi}{\partial y} \\ \frac{\partial A}{\partial x} & \frac{\partial A}{\partial y} \end{pmatrix} = \begin{array}{c} \begin{pmatrix} x & -y \\ y & x \end{pmatrix} \end{array} \begin{array}{l} \leftarrow \mathbf{E}^{\Phi} \\ \leftarrow \mathbf{E}^{A} \end{array}$$

$$Jacobian = \begin{pmatrix} \frac{\partial x}{\partial q^1} & \frac{\partial x}{\partial q^2} \\ \frac{\partial y}{\partial q^1} & \frac{\partial y}{\partial q^2} \end{pmatrix} = \begin{array}{cc} \begin{pmatrix} \frac{\partial x}{\partial \Phi} & \frac{\partial x}{\partial A} \\ \frac{\partial y}{\partial \Phi} & \frac{\partial y}{\partial A} \end{pmatrix} \\ \begin{matrix} \uparrow & \uparrow \\ \mathbf{E}_{\Phi} & \mathbf{E}_{A} \end{matrix} \end{array} = \frac{1}{r^2} \begin{array}{c} \begin{pmatrix} x & y \\ -y & x \end{pmatrix} \\ \begin{matrix} \uparrow & \uparrow \\ \mathbf{E}_{\Phi} & \mathbf{E}_{A} \end{matrix} \end{array}$$

$$Metric\,tensor = \begin{pmatrix} g_{\Phi\Phi} & g_{\Phi A} \\ g_{A\Phi} & g_{AA} \end{pmatrix} = \begin{pmatrix} \mathbf{E}_{\Phi}\bullet\mathbf{E}_{\Phi} & \mathbf{E}_{\Phi}\bullet\mathbf{E}_{A} \\ \mathbf{E}_{A}\bullet\mathbf{E}_{\Phi} & \mathbf{E}_{A}\bullet\mathbf{E}_{A} \end{pmatrix} = \begin{pmatrix} r^2 & 0 \\ 0 & r^2 \end{pmatrix} \text{ where: } r^2 = x^2 + y^2$$

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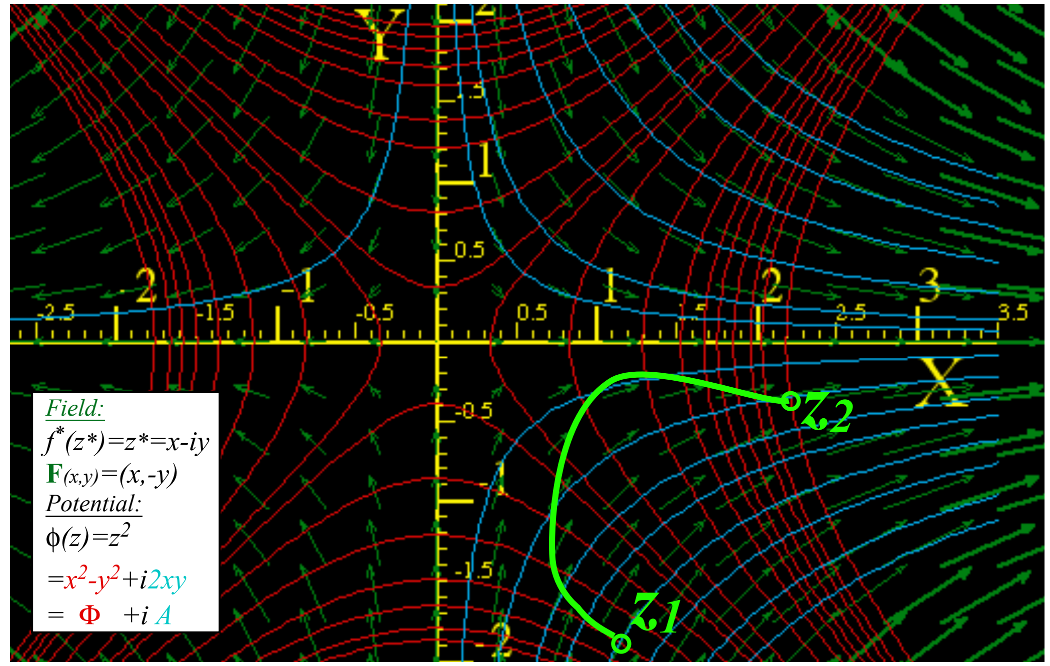

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Riemann-Cauchy Derivative Relations make coordinates orthogonal

$$\nabla\Phi = \begin{pmatrix} \frac{\partial \Phi}{\partial x} \\ \frac{\partial \Phi}{\partial y} \end{pmatrix} = \begin{pmatrix} \frac{\partial}{\partial x}\frac{a}{2}(x^2-y^2) \\ \frac{\partial}{\partial y}\frac{a}{2}(x^2-y^2) \end{pmatrix} = \begin{pmatrix} ax \\ -ay \end{pmatrix} = \mathbf{F}$$

The half-n'-half results assure

$$\mathbf{E}_{\Phi}\bullet\mathbf{E}_{A} = \frac{\partial \Phi}{\partial x}\frac{\partial A}{\partial x} + \frac{\partial \Phi}{\partial y}\frac{\partial A}{\partial y}$$

$$= -\frac{\partial \Phi}{\partial x}\frac{\partial \Phi}{\partial y} + \frac{\partial \Phi}{\partial y}\frac{\partial \Phi}{\partial x} = 0$$

$$\nabla\times\mathbf{A} = \begin{pmatrix} \frac{\partial \mathbf{A}}{\partial y} \\ -\frac{\partial \mathbf{A}}{\partial x} \end{pmatrix} = \begin{pmatrix} \frac{\partial}{\partial y} axy \\ -\frac{\partial}{\partial x} axy \end{pmatrix} = \begin{pmatrix} ax \\ -ay \end{pmatrix} = \mathbf{F}$$

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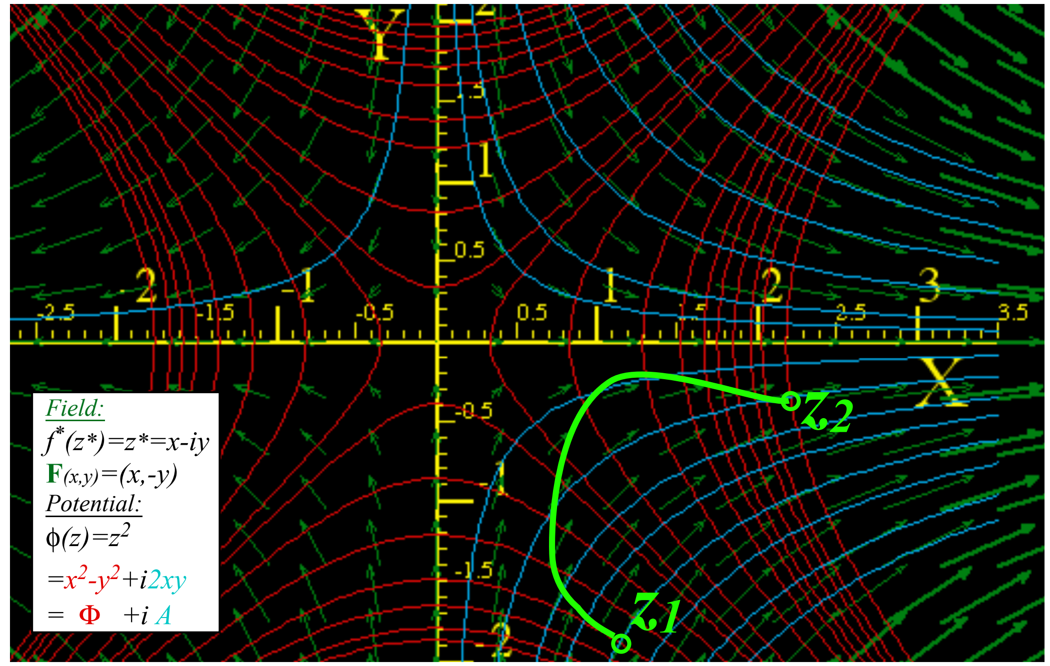

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or Riemann-Cauchy

Zero divergence requirement: $0 = \frac{\partial f_x}{\partial x} + \frac{\partial f_y}{\partial y} = \frac{\partial}{\partial x}\frac{\partial \Phi}{\partial x} + \frac{\partial}{\partial y}\frac{\partial \Phi}{\partial y} = \frac{\partial^2 \Phi}{\partial x^2} + \frac{\partial^2 \Phi}{\partial y^2} = 0$ *potential Φ obeys Laplace equation*

and so does $\mathbf{A}$

4. Riemann-Cauchy conditions *What's analytic? (...and what's <u>not</u>?)*

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11. Complex integrals define 2D monopole fields and potentials

Of all power-law fields $f(z)=az^n$ one lacks a power-law potential $\phi(z)=\frac{a}{n+1}z^{n+1}$. It is the $n = -1$ case.

Unit monopole field: $f(z)=\frac{1}{z}=z^{-1}$ $\qquad$ $f(z)=\frac{a}{z}=az^{-1}$ Source-*a* monopole

It has a *logarithmic potential* $\phi(z)=a\cdot\ln(z)=a\cdot\ln(x+iy)$.

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$$\phi(z) = \quad \Phi \quad + \quad i\mathbf{A} = \int f(z)dz = \int \frac{a}{z}dz = a\ln(z)$$

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What Good Are Complex Exponentials? (contd.)

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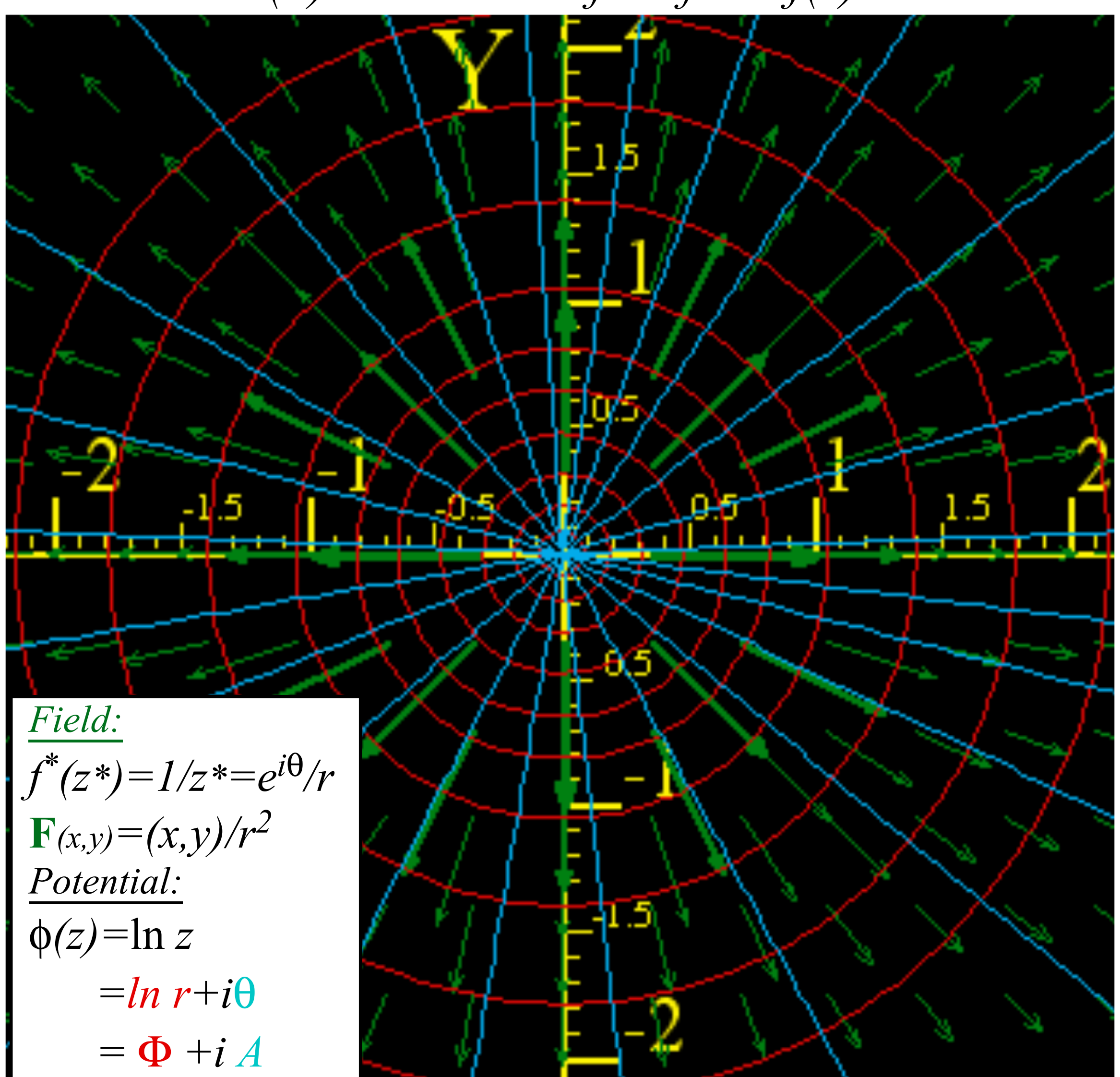
(a) Unit Z-line-flux field f(z)=1/z



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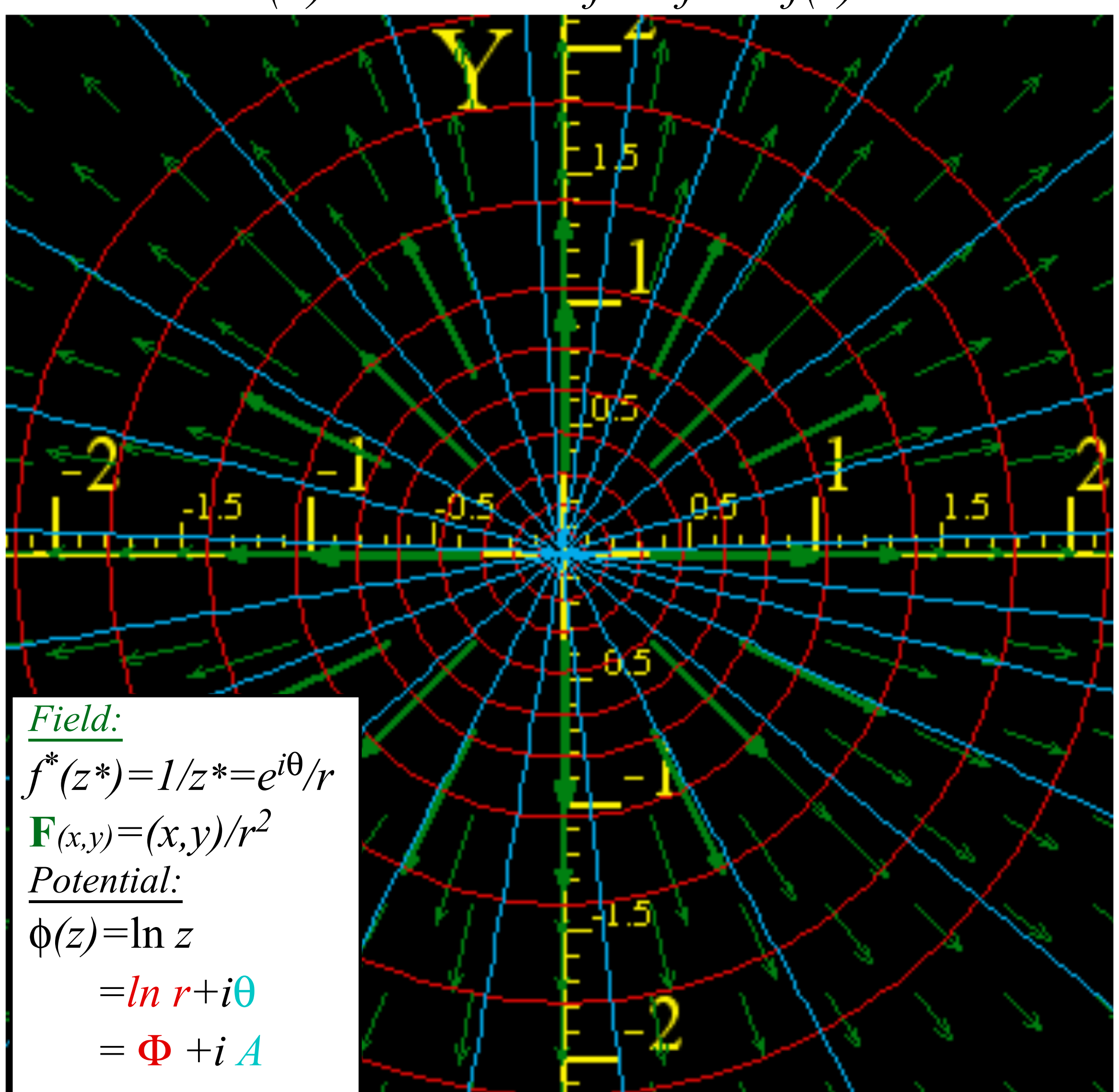
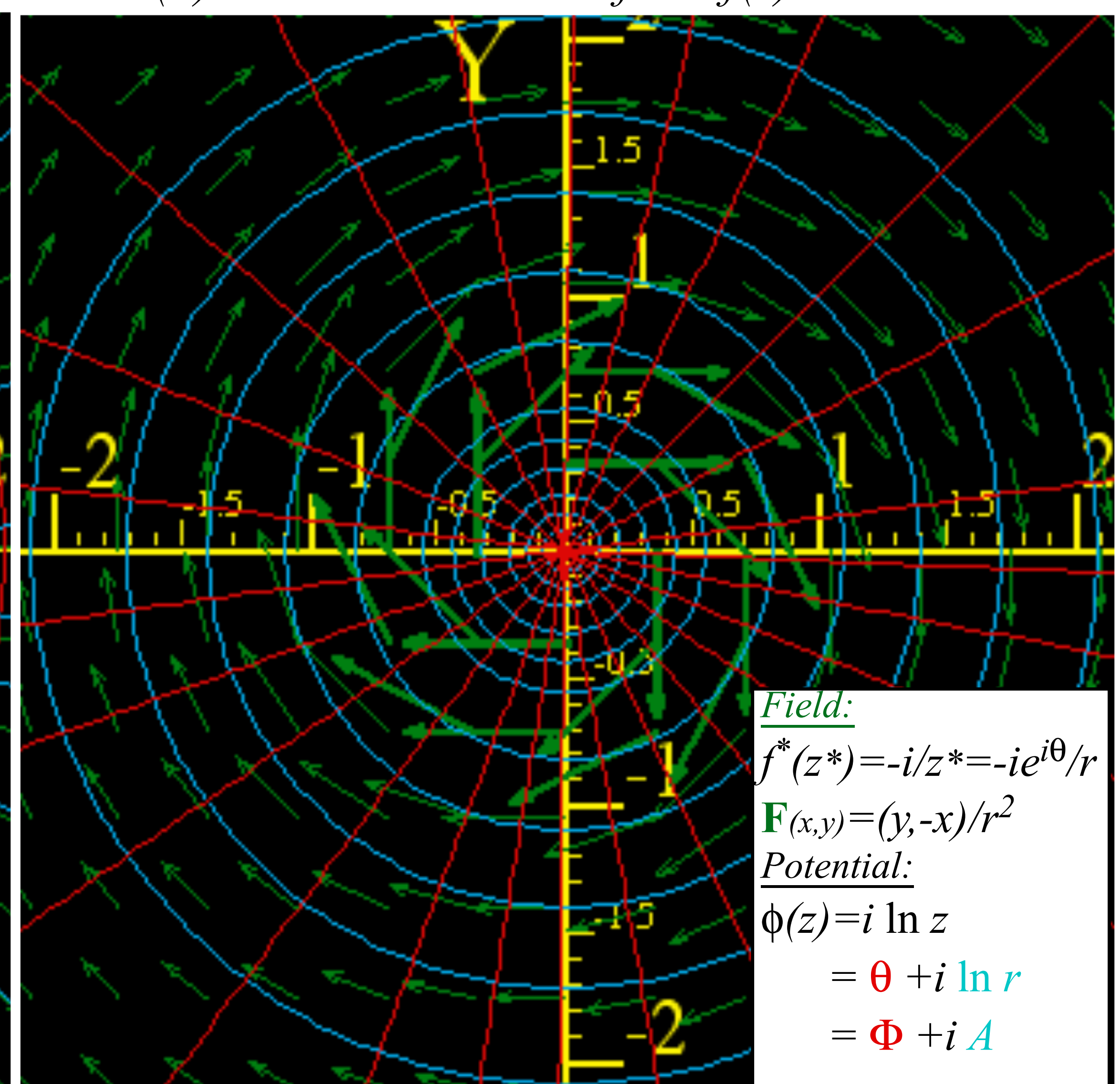


(b) Unit Z-line-vortex field f(z)=i/z



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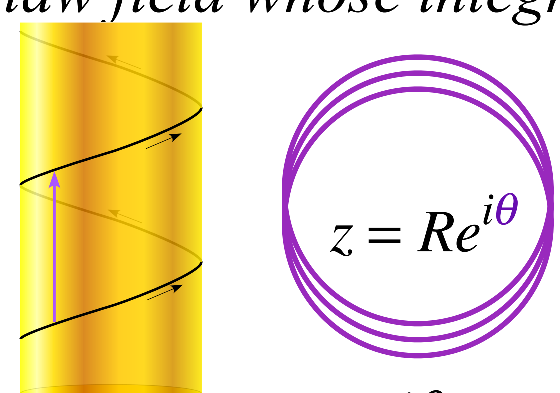
A monopole field is the only power-law field whose integral (potential) depends on path of integration.



path that goes N times around origin (r=0) at constant r = R.

$$\Delta\phi = \oint f(z)dz = a\oint \frac{dz}{z} = a\int_{\theta=0}^{\theta=2\pi N} \frac{d(Re^{i\theta})}{Re^{i\theta}} = a\int_{\theta=0}^{\theta=2\pi N} i d\theta = ai\,\theta\Big|_0^{2\pi N} = 2a\pi iN$$

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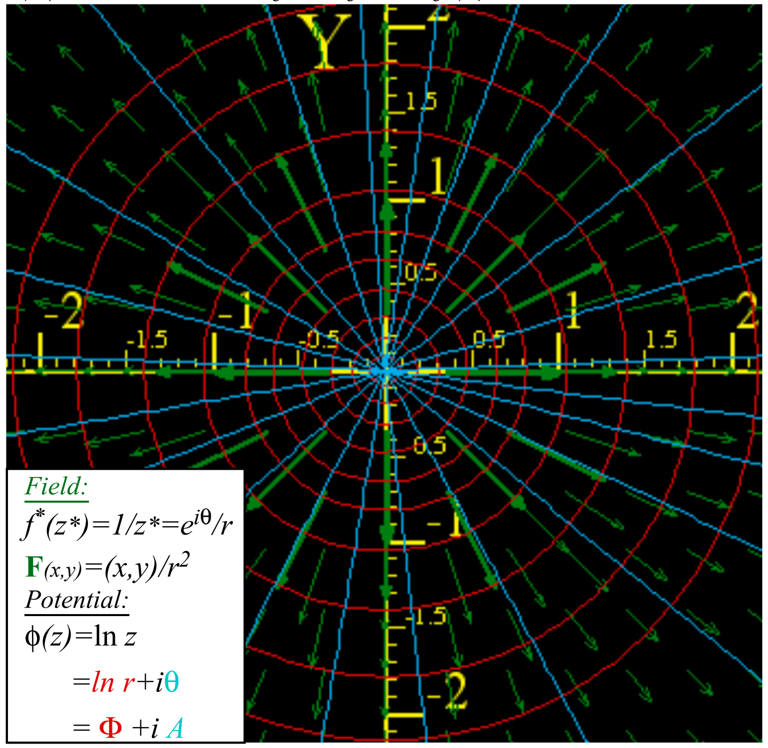
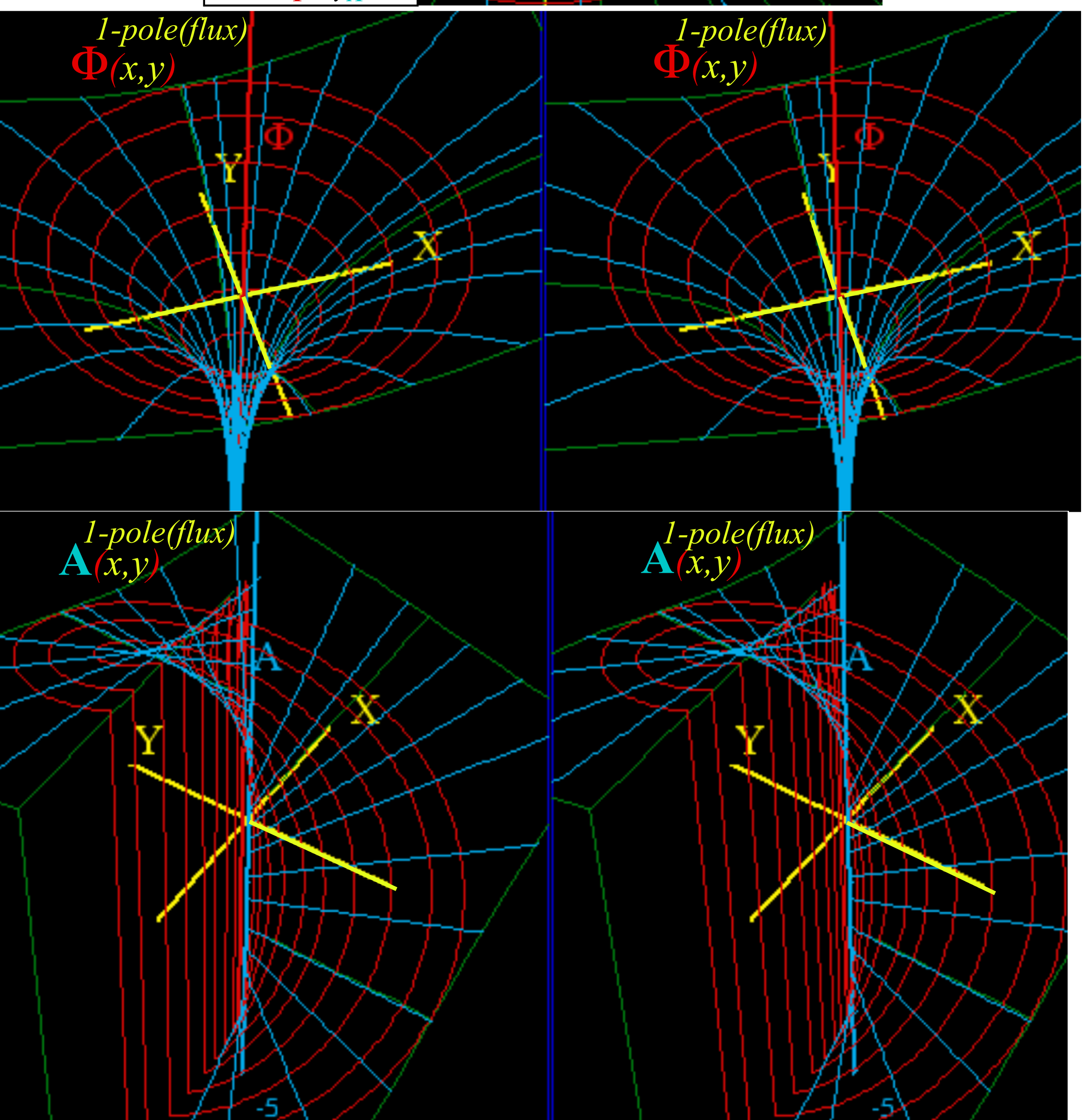


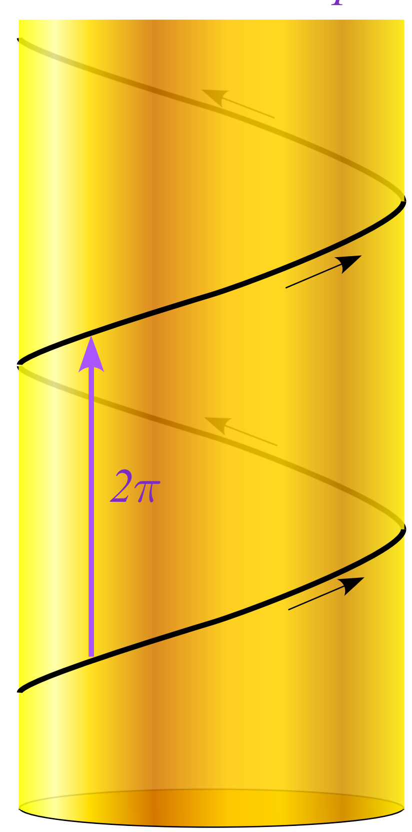
$$\phi(z) = \Phi + i\mathbf{A} = \int f(z)\,dz = \int \frac{a}{z}\,dz = a\ln(re^{i\theta})$$

$$= \ln(r) + i\theta \qquad \textit{(For } a=1\textit{)}$$



Each turn around origin

adds $2\pi i$ to vector potential $i\mathbf{A}$



(For a=1)

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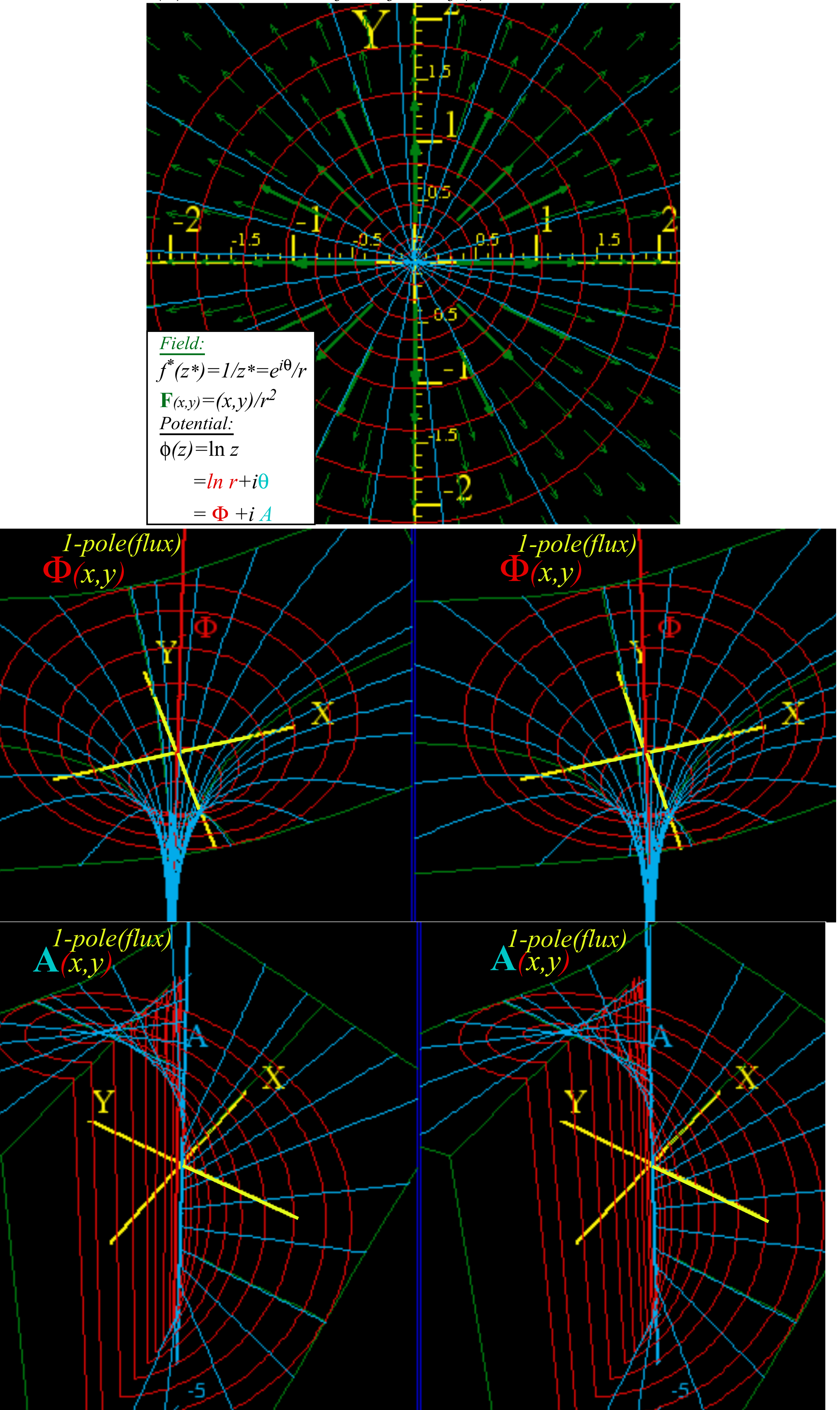
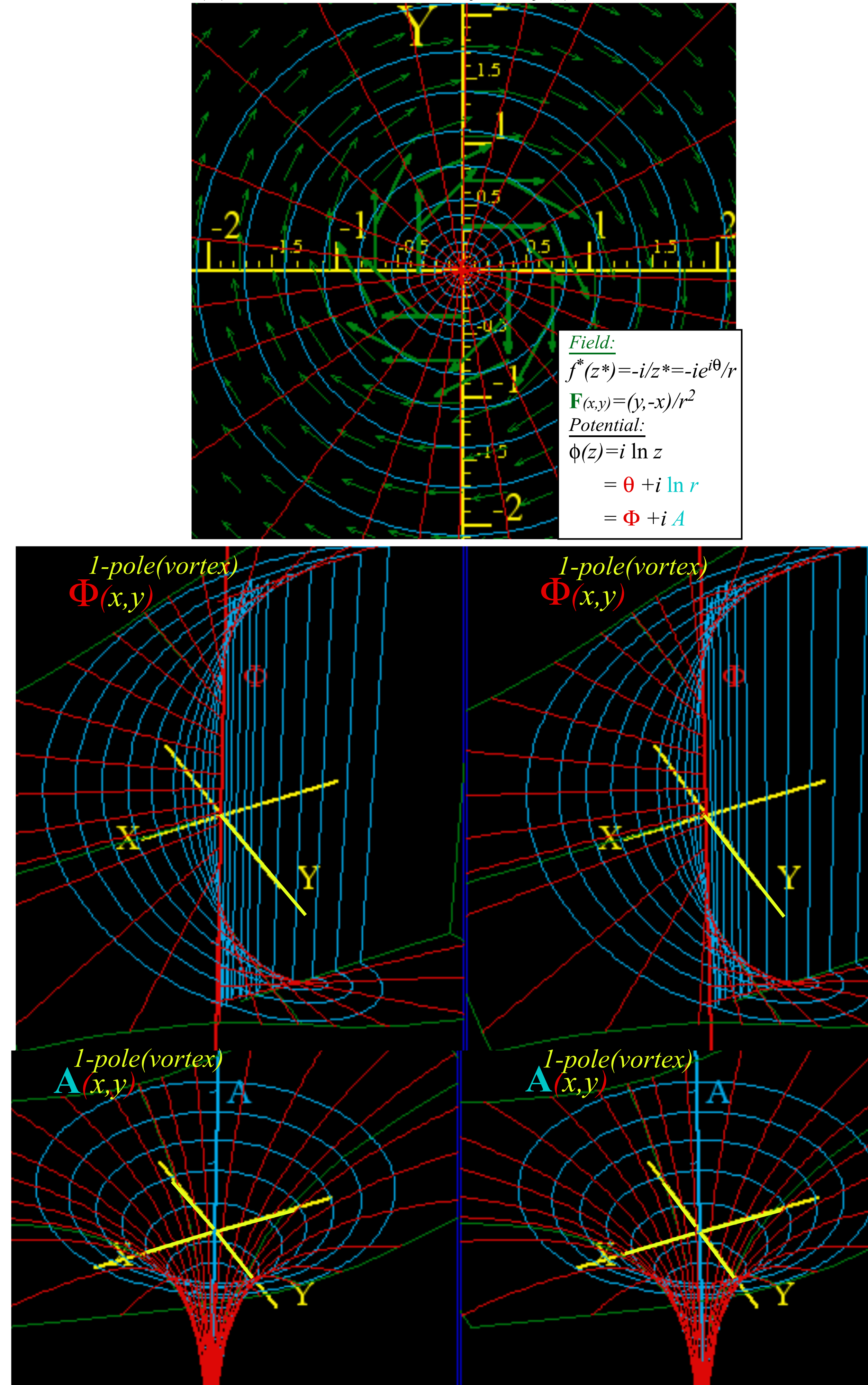


(b) Unit Z-line-vortex field f(z)=i/z

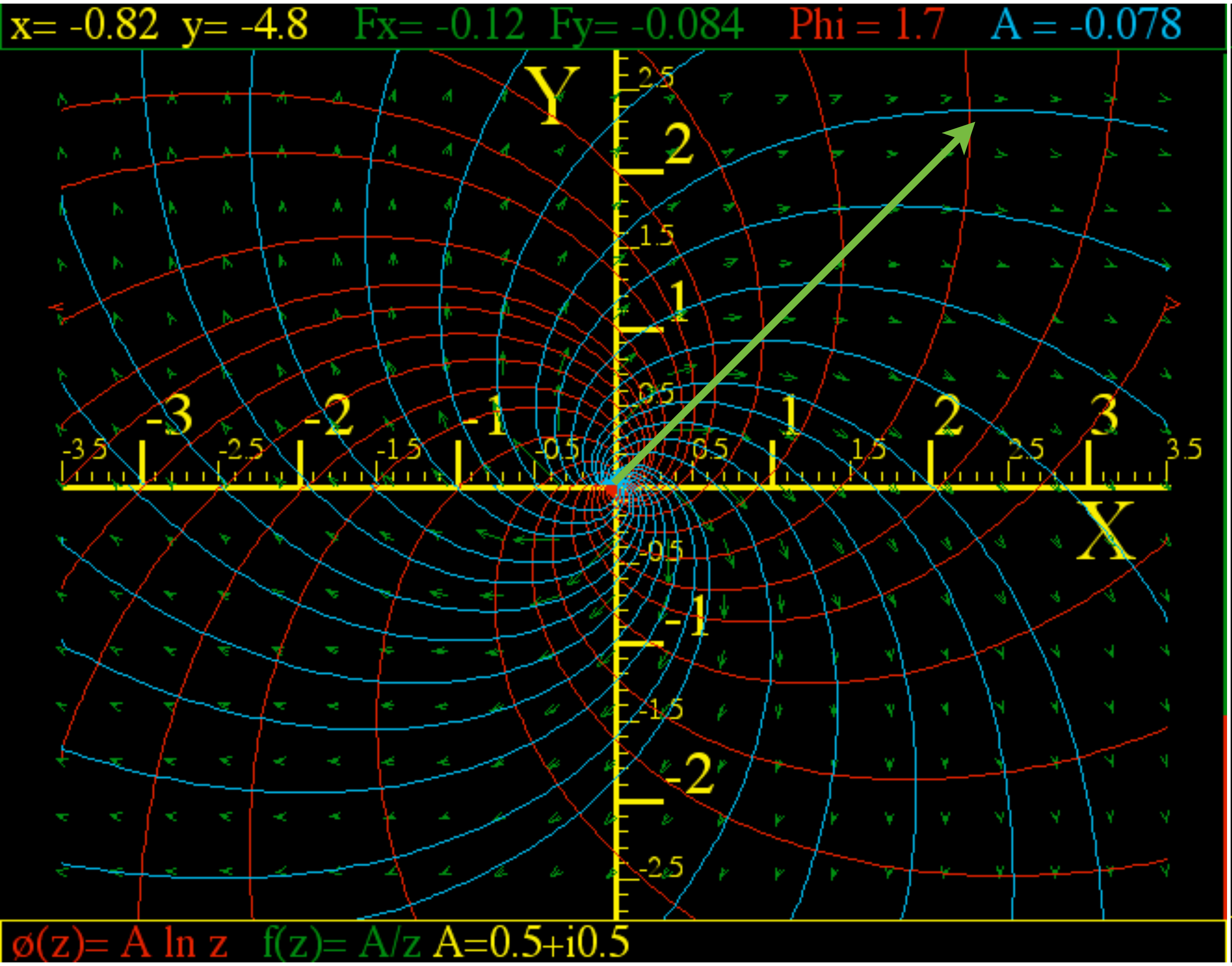


What Good Are Complex Exponentials? (contd.)

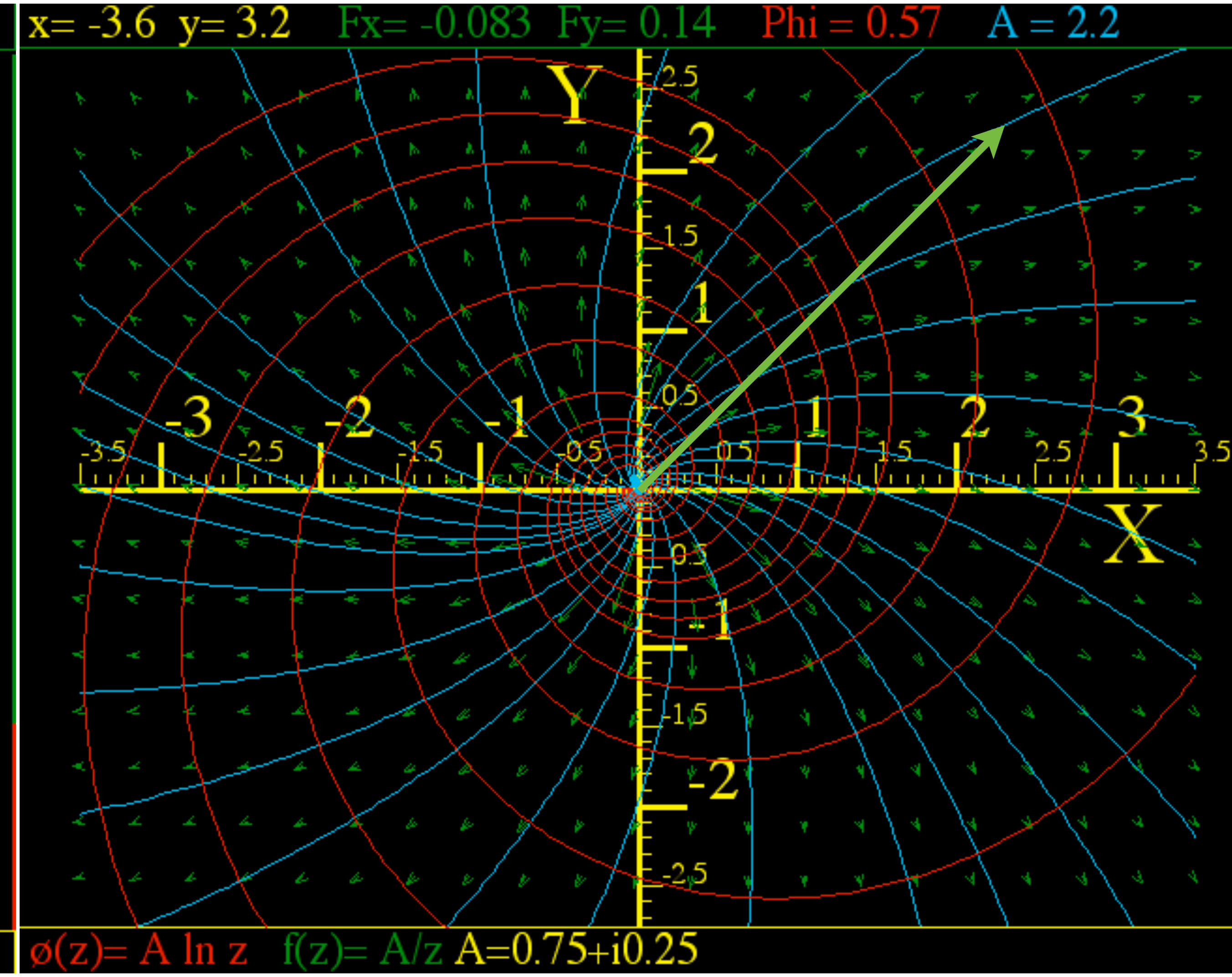
$f(z) = (0.5 + i0.5)/z = e^{i\pi/4}/z\sqrt{2}$

$f(z) = (0.75 + i0.25)/z = e^{i18^\circ}/z\sqrt{n}$

"Vortex"



"Hurricane"



What Good Are Complex Exponentials? (contd.)

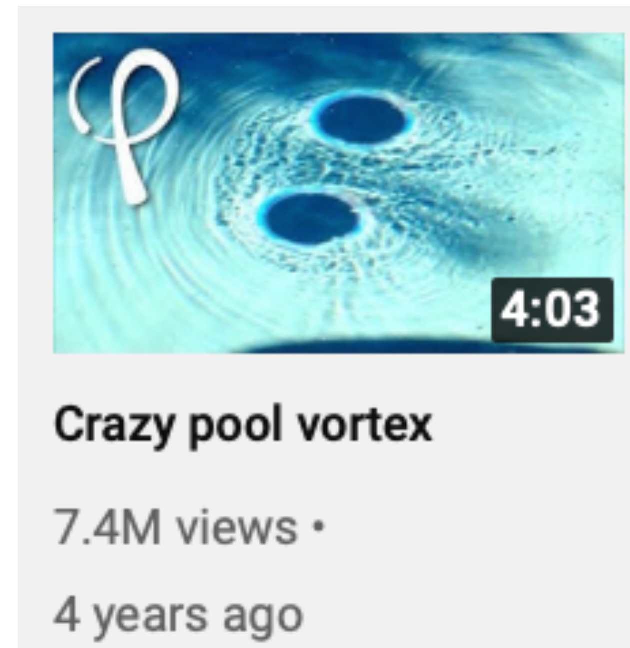
An assist from *Physics Girl* (YouTube Channel):

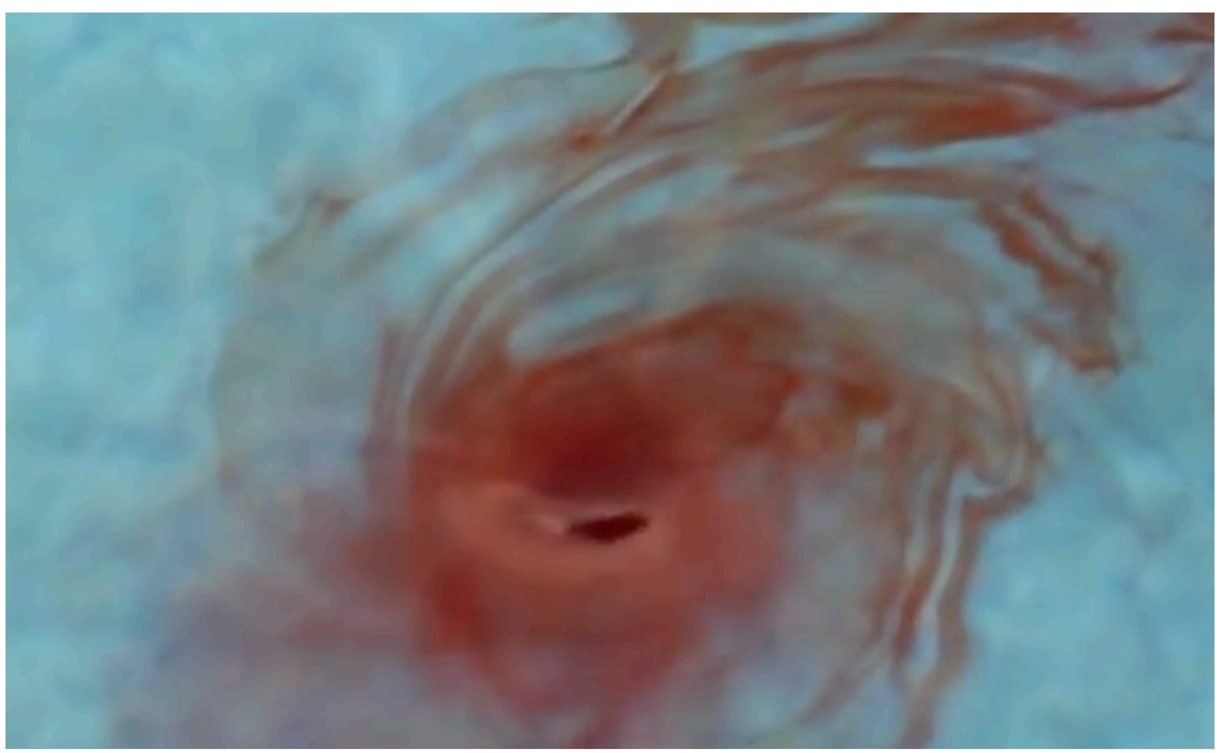
Posted this year:
How to Make VORTEX RINGS in a Pool

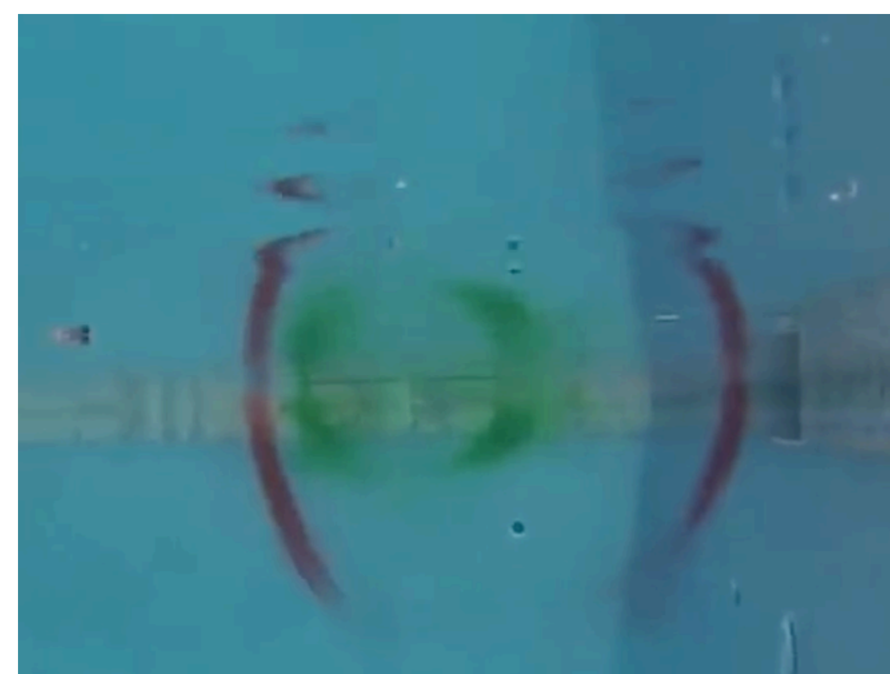
Crazy pool vortex (new inclusion with more background)
Crazy pool vortex - pg-yt-2014

Posting with the best visuals:
Fun with Vortex Rings in the Pool - pg-yt-2014

She covers it beautifully!







4. Riemann-Cauchy conditions *What's analytic? (...and what's <u>not</u>?)*

Easy 2D circulation and flux integrals

Easy 2D curvilinear coordinate discovery

→ *Easy 2D monopole, dipole, and 2^n-pole analysis*

Easy 2^n-multipole field and potential expansion

Easy stereo-projection visualization

Lect 13.5 starts here

12. Complex derivatives give 2D dipole fields

Start with $f(z)=az^{-1}$: 2D line *monopole field* and is its *monopole potential* $\phi(z)=a\ln z$ of source strength *a*.

$$f^{1\text{-}pole}(z)=\frac{a}{z}=\frac{d\phi^{1\text{-}pole}}{dz} \qquad \phi^{1\text{-}pole}(z)=a\ln z$$

Now let these two line-sources of equal but opposite source constants $+a$ and $-a$ be located at $z=\pm\Delta/2$ separated by a small interval Δ. This sum (actually difference) of $f^{1\text{-}pole}$-fields is called a *dipole field*.

$$f^{dipole}(z)=\frac{a}{z+\frac{\Delta}{2}}-\frac{a}{z-\frac{\Delta}{2}}=\frac{-a\cdot\Delta}{z^2-\frac{\Delta^2}{4}} \qquad \phi^{dipole}(z)=a\ln(z-\tfrac{\Delta}{2})-a\ln(z+\tfrac{\Delta}{2})=a\ln\frac{z-\frac{\Delta}{2}}{z+\frac{\Delta}{2}}$$

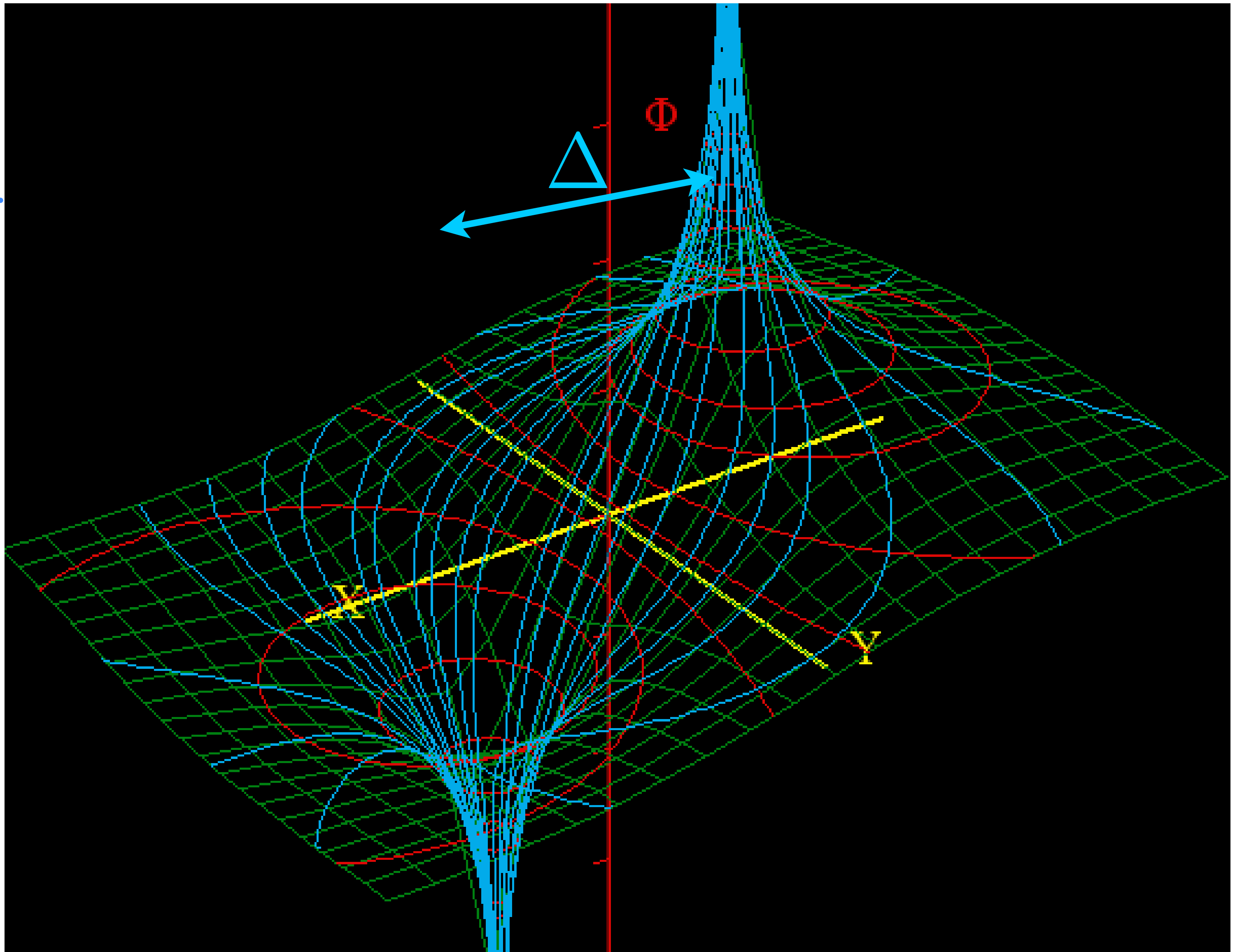
This is like the derivative definition:

$$\frac{df}{dz}=\frac{f(z+\Delta)-f(z)}{\Delta}$$

or:

$$\frac{df}{dz}=\frac{f(z+\frac{\Delta}{2})-f(z-\frac{\Delta}{2})}{\Delta}$$

if Δ is infinitesimal ($\Delta\rightarrow 0$)



So-called "physical dipole" has finite Δ

(+)(−) separation

12. Complex derivatives give 2D dipole fields

Start with $f(z)=az^{-1}$: 2D line *monopole field* and is its *monopole potential* $\phi(z)=a\ln z$ of source strength a.

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If interval Δ is *tiny* and is divided out we get a *point-dipole field* $f^{2\text{-}pole}$ that is the z-derivative of $f^{1\text{-}pole}$.

$$f^{2\text{-}pole}=\frac{-a}{z^2}=\frac{df^{1\text{-}pole}}{dz}=\frac{d\phi^{2\text{-}pole}}{dz} \qquad \phi^{2\text{-}pole}=\frac{a}{z}=\frac{d\phi^{1\text{-}pole}}{dz}$$

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Start with $f(z)=az^{-1}$: 2D line *monopole field* and is its *monopole potential* $\phi(z)=a\ln z$ of source strength a.

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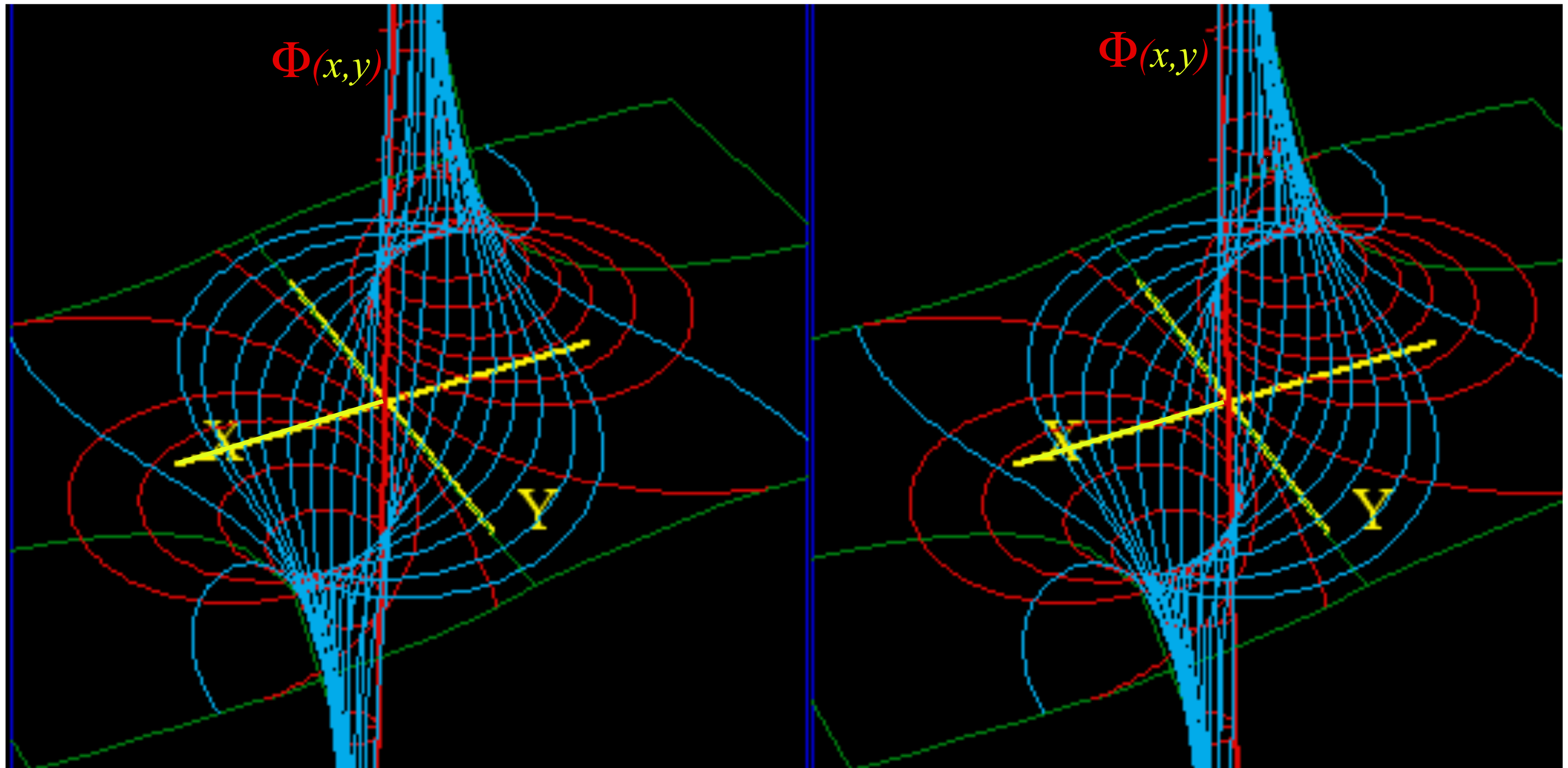

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Start with $f(z)=az^{-1}$: 2D line *monopole field* and is its *monopole potential* $\phi(z)=a\ln z$ of source strength a.

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If interval Δ is *tiny* and is divided out we get a *point-dipole field* $f^{2\text{-}pole}$ that is the z-derivative of $f^{1\text{-}pole}$.

$$f^{2\text{-}pole}=\frac{-a}{z^2}=\frac{df^{1\text{-}pole}}{dz}=\frac{d\phi^{2\text{-}pole}}{dz} \qquad \phi^{2\text{-}pole}=\frac{a}{z}=\frac{d\phi^{1\text{-}pole}}{dz}$$

A *point-dipole potential* $\phi^{2\text{-}pole}$ (whose z-derivative is $f^{2\text{-}pole}$) is a z-derivative of $\phi^{1\text{-}pole}$.

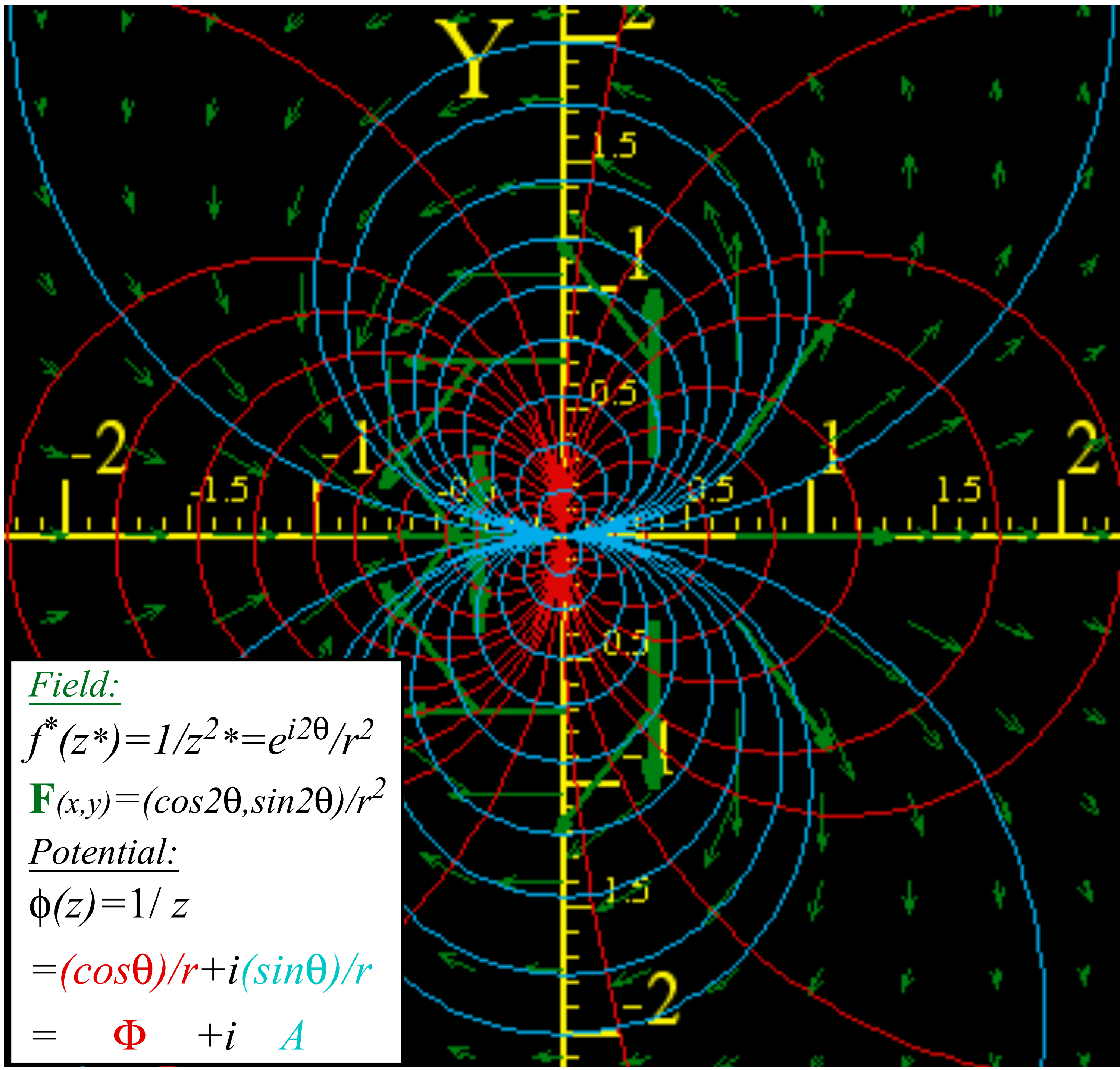
$$\phi^{2\text{-}pole}=\frac{a}{z}=\frac{a}{x+iy}=\frac{a}{x+iy}\frac{x-iy}{x-iy}=\frac{ax}{x^2+y^2}+i\frac{-ay}{x^2+y^2}=\frac{a}{r}\cos\theta-i\frac{a}{r}\sin\theta$$

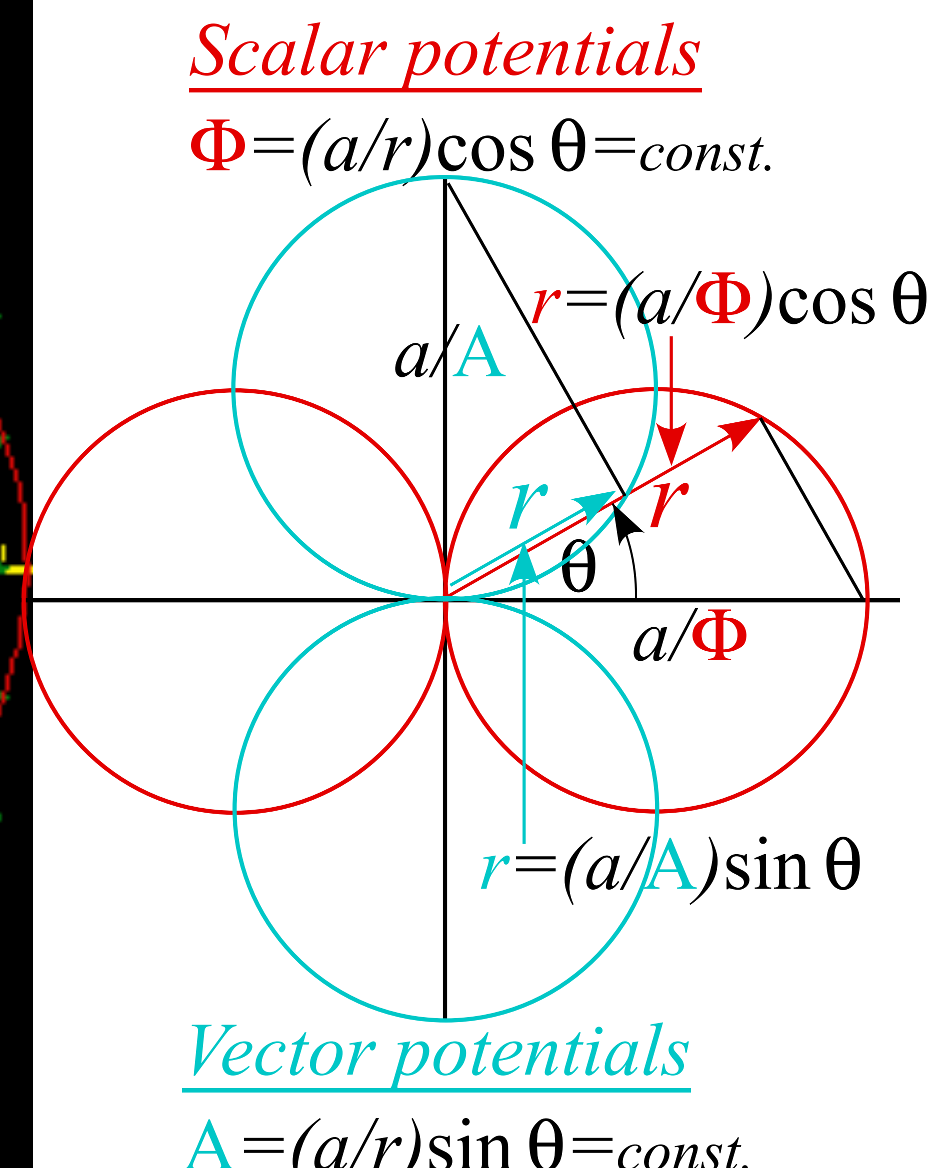
$$=\Phi^{2\text{-}pole}+i\,\mathbf{A}^{2\text{-}pole}$$

A *point-dipole potential* $\phi^{2\text{-}pole}$ (whose z-derivative is $f^{2\text{-}pole}$) is a z-derivative of $\phi^{1\text{-}pole}$.

$$\phi^{2\text{-}pole} = \frac{a}{z} = \frac{a}{x+iy} = \frac{a}{x+iy}\frac{x-iy}{x-iy} = \frac{ax}{x^2+y^2} + i\frac{-ay}{x^2+y^2} = \frac{a}{r}\cos\theta - i\frac{a}{r}\sin\theta$$

$$= \Phi^{2\text{-}pole} + i\,A^{2\text{-}pole}$$





Approximate Lorentz-Green's Function for high quality FDHO (Quantum propagator)

(Preview of Unit 4.)

$$G_{\omega_0}(\omega_s) = \frac{1}{\omega_0^2 - \omega_s^2 - i2\Gamma\omega_s} \xrightarrow[\omega_s \to \omega_0]{} \frac{1}{2\omega_s}\frac{1}{\omega_0 - \omega_s - i\Gamma} \approx \frac{1}{2\omega_0}\frac{1}{\Delta - i\Gamma} = \frac{1}{2\omega_0}L(\Delta - i\Gamma)$$

Complex detuning-decay $\delta = \Delta - i\Gamma$ variable δ is defined with the *real detuning* $\Delta = \omega_0 - \omega_s$

$$L(\Delta - i\Gamma) = \frac{1}{\Delta - i\Gamma} = \mathrm{Re}\,L \quad + i\,\mathrm{Im}\,L \quad = \quad \frac{\Delta}{\Delta^2 + \Gamma^2} \quad + i\frac{\Gamma}{\Delta^2 + \Gamma^2} = \quad |L|^2\,\Delta \quad + i\,|L|^2\,\Gamma$$

$$= |L|\,e^{i\rho} = |L|\cos\rho + i\,|L|\sin\rho = \frac{\cos\rho}{\sqrt{\Delta^2 + \Gamma^2}} + i\frac{\sin\rho}{\sqrt{\Delta^2 + \Gamma^2}} \quad \text{where: } |L| = \frac{1}{\sqrt{\Delta^2 + \Gamma^2}}$$

$$|L| = \frac{1}{\Gamma}\sin\rho$$

$$|L| = \frac{1}{\Delta}\cos\rho$$

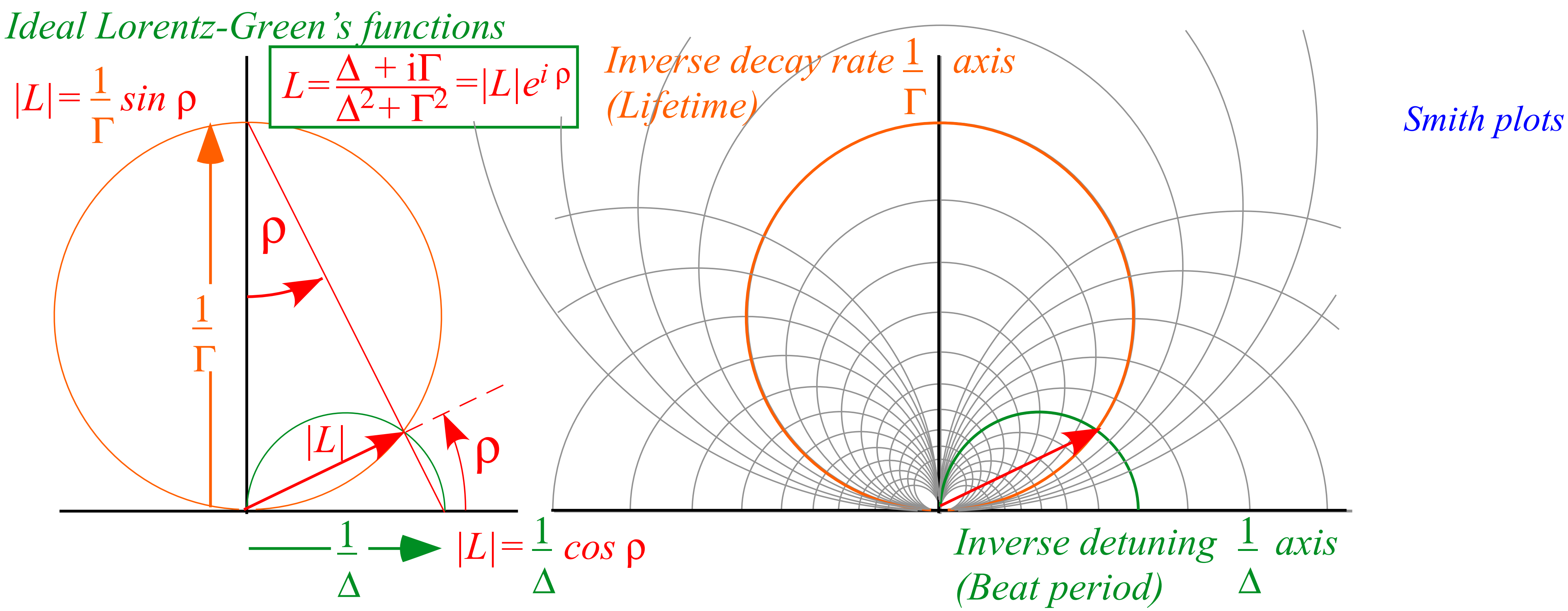


Fig. 4.2.13 Ideal Lorentzian in inverse rate space. (Smith life-time $1/\Gamma$ *vs. beat-period* $1/\Delta$ *coordinates)*

Constant Δ and Γ curves in Fig. 4.2.13 are orthogonal circles of *1/z- dipolar coordinates*. Recall Fig. 1.10.11.

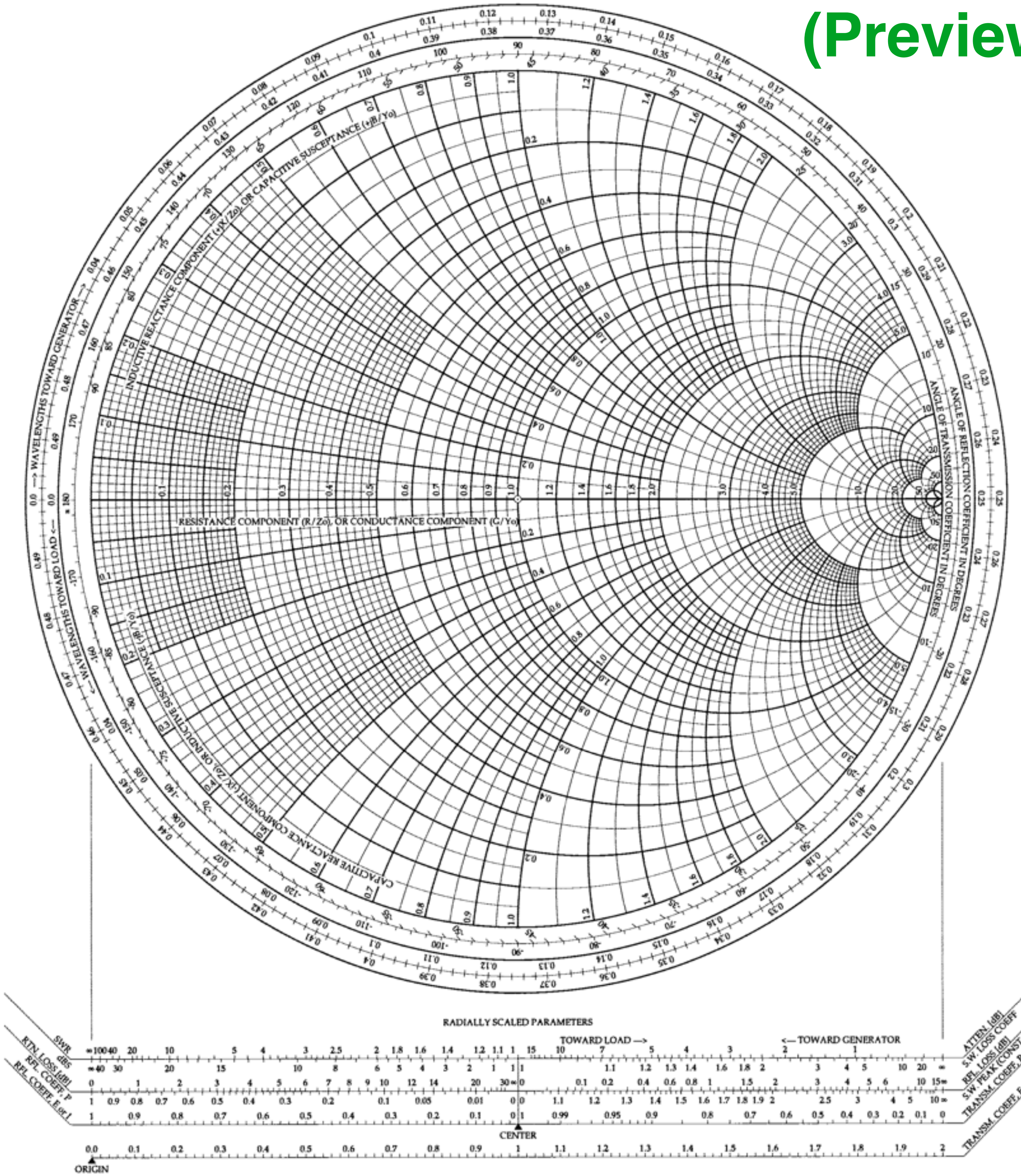
SMITH CHART (Invented by Phillip H. Smith 1905-1987)

(Preview of Unit 4.)

An FDHO Green's Function Slide rule

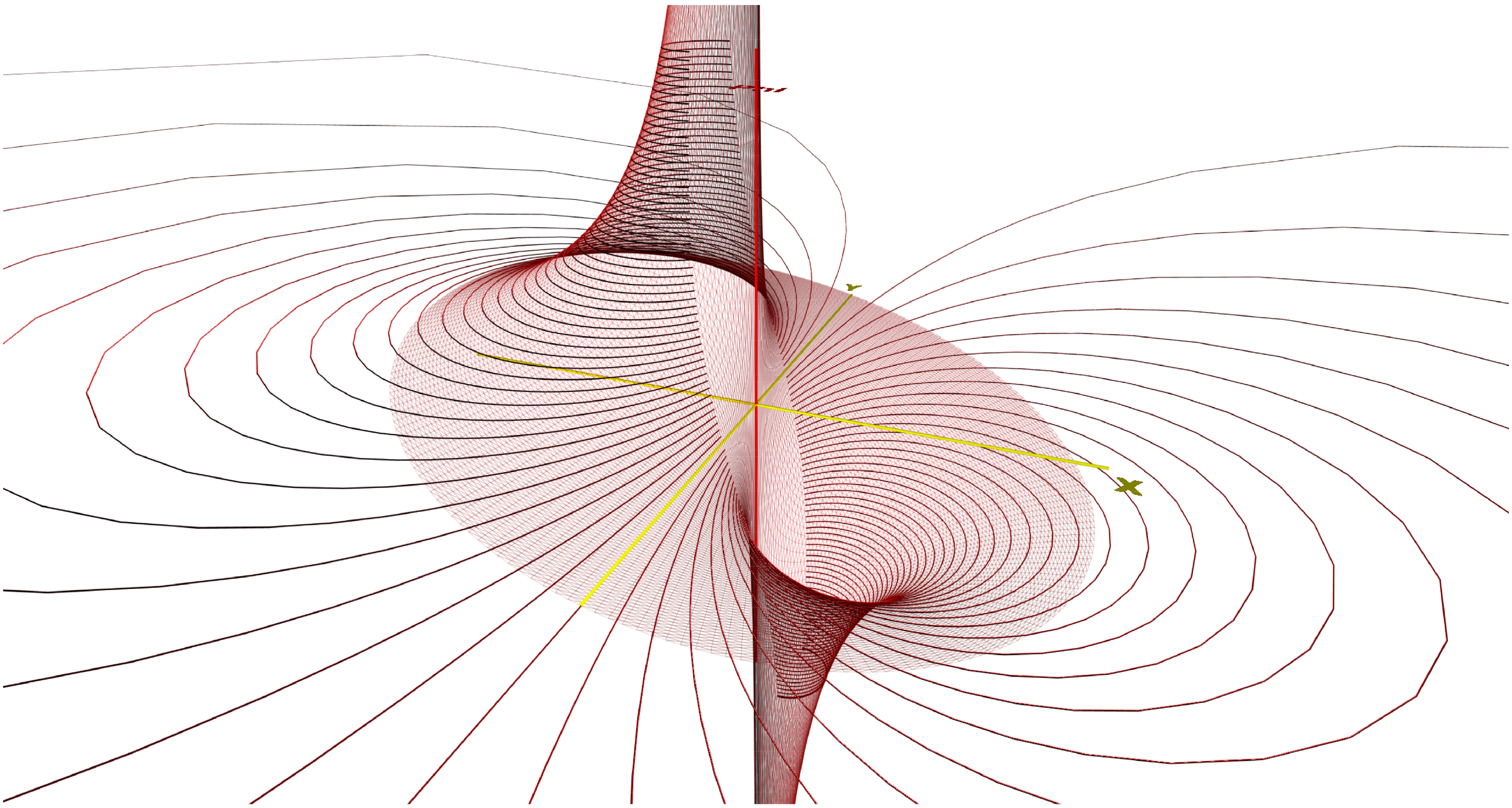
A plot of $f(z) = 1/z$

For wavy "Ohm's Laws"
$V = I \cdot Z$
$I = V/Z$

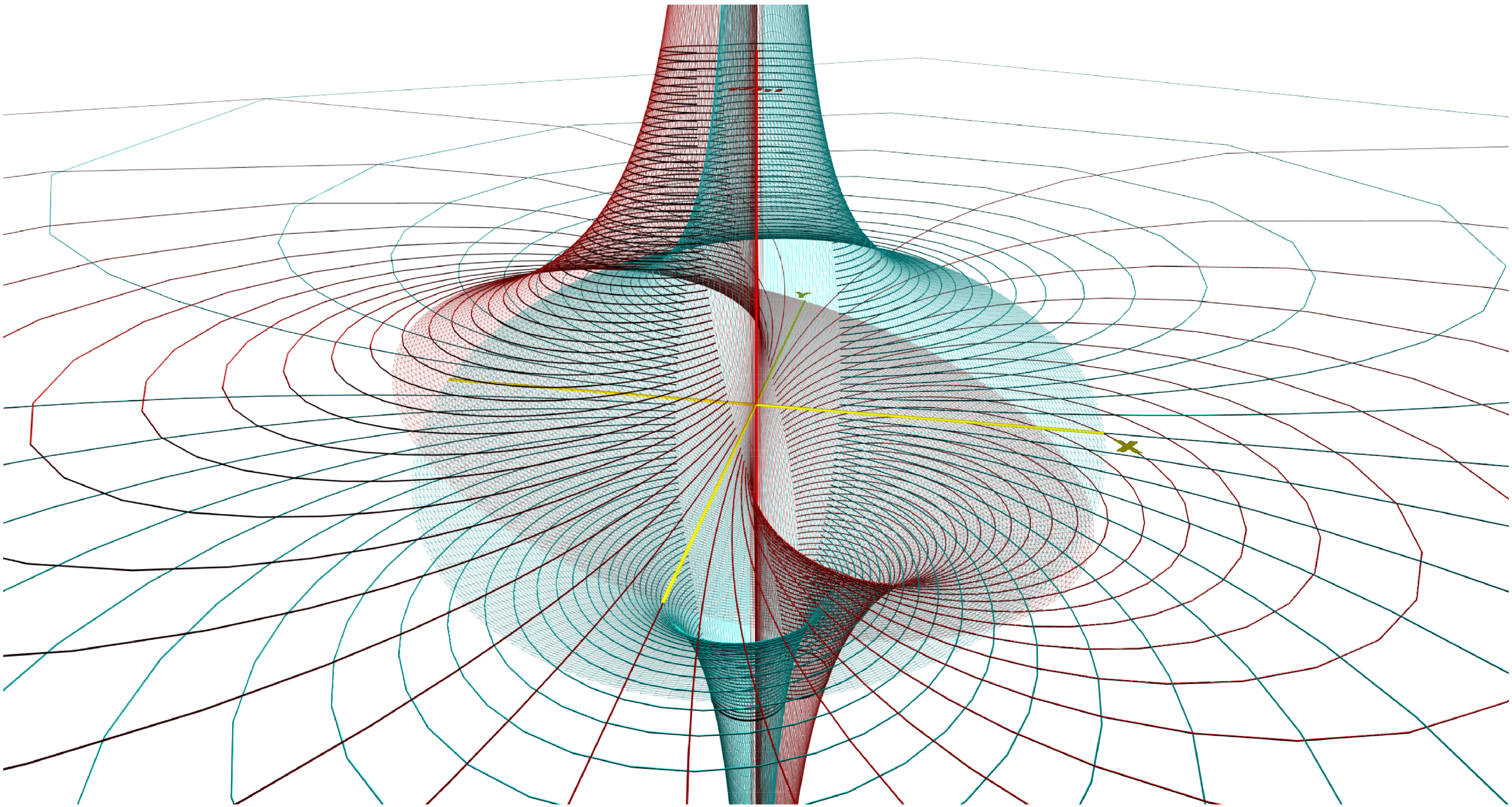


[Smith plot: Graph paper]

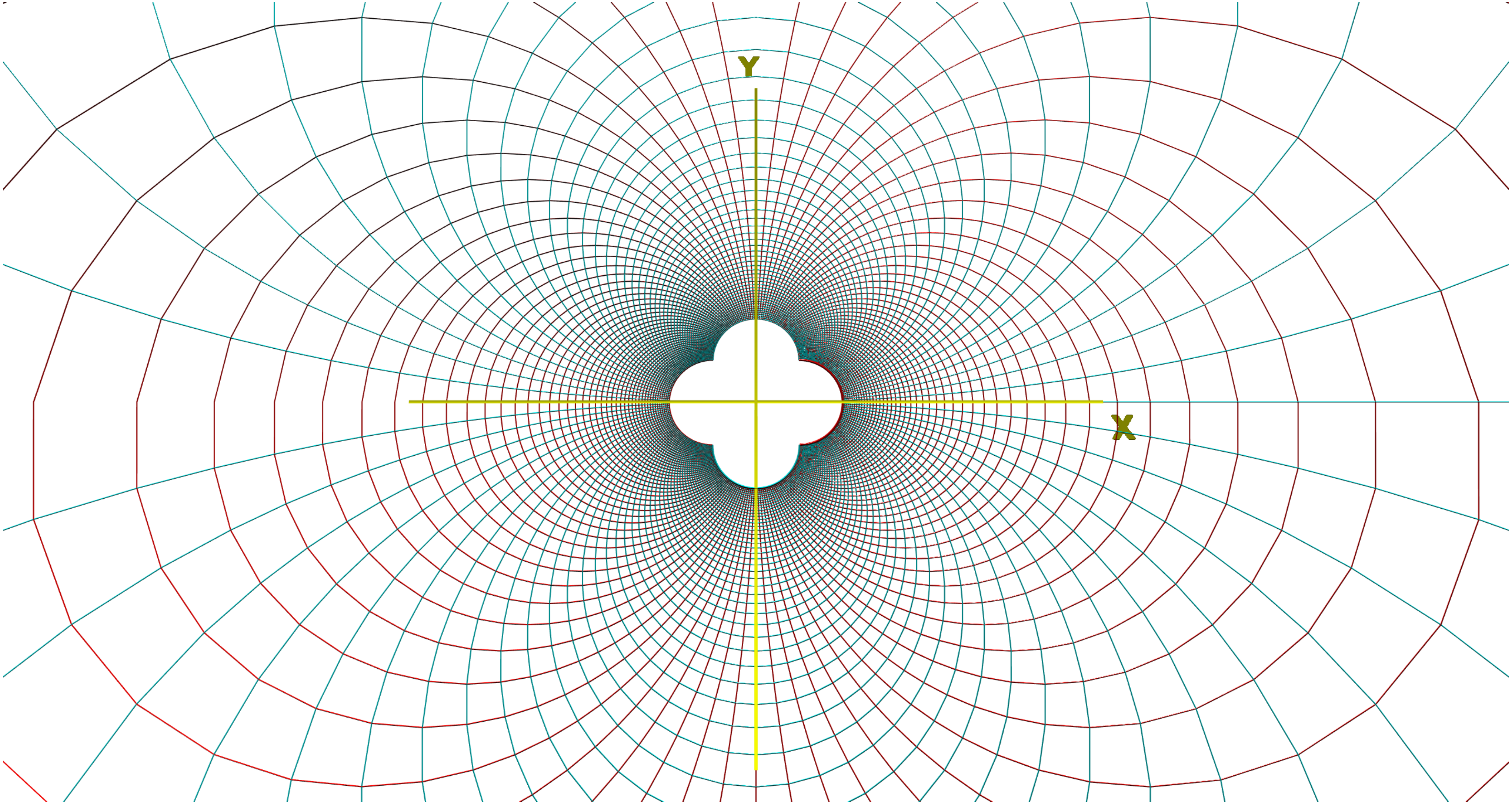
(Preview of Unit 4.)



(Preview of Unit 4.)



(Preview of Unit 4.)



https://modphys.hosted.uark.edu/video/AnalyIt_0-3.webm

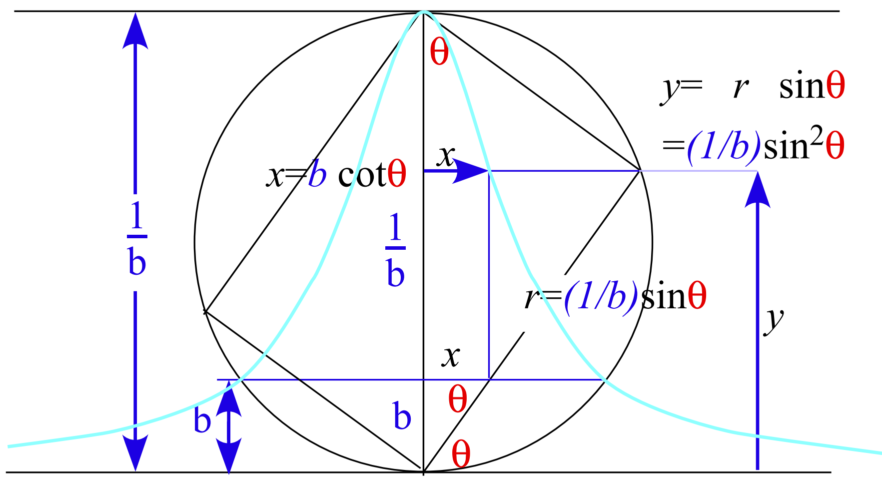
https://modphys.hosted.uark.edu/video/AnalyIt_4-1.webm

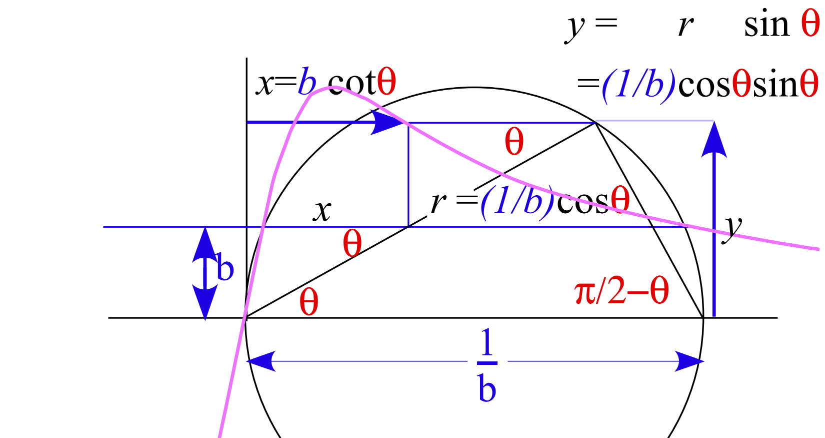
https://modphys.hosted.uark.edu/video/AnalyIt_0-2.webm

https://modphys.hosted.uark.edu/video/AnalyIt_3-1.webm

https://modphys.hosted.uark.edu/video/AnalyIt_1-1.webm

The Common Lorentzian (a.k.a. *The Witch of Agnesi*)





Maria Gaetana Agnesi

Born	May 16, 1718
Died	January 9, 1799 (aged 80)
Residence	Italy
Nationality	Italy
Fields	Mathematics

$$x^2=b^2\cot^2\theta=b^2\frac{\cos^2\theta}{\sin^2\theta}=b^2\frac{1-\sin^2\theta}{\sin^2\theta}=\frac{b^2}{\sin^2\theta}\,b^2$$

$$x^2+b^2=\frac{b^2}{\sin^2\theta}=\frac{b}{y}$$

$$y=\frac{b}{x^2+b^2}$$

Common Lorentzian function I.
(imaginary "absorbtive" part)

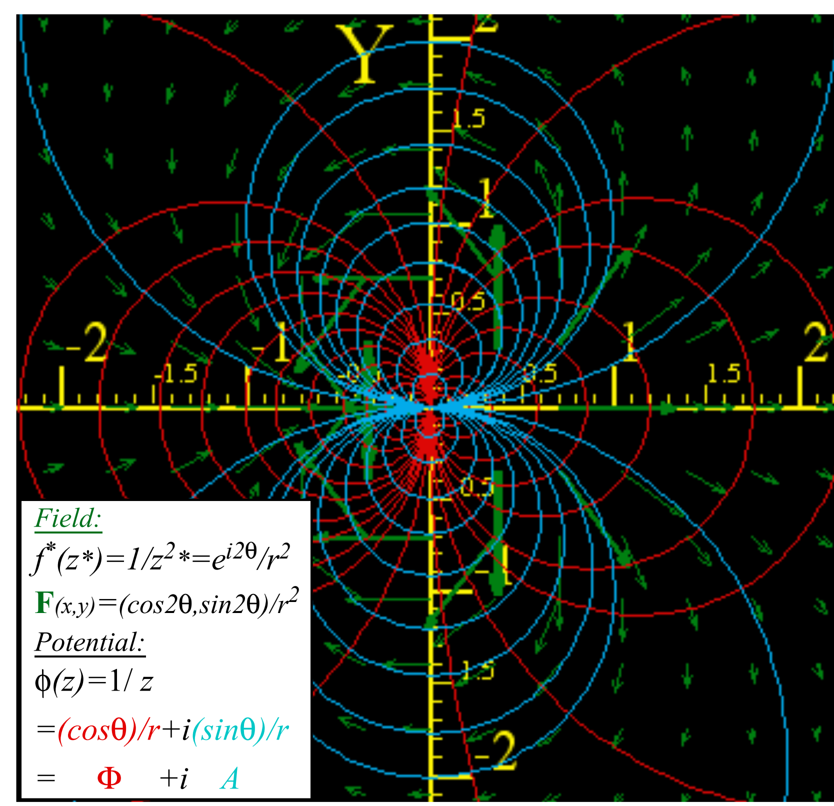
$$\frac{x}{y}=\frac{b\cot\theta}{(1/b)\cos\theta\sin\theta}=\frac{b^2\cos\theta}{\cos\theta\sin^2\theta}=\frac{b^2}{\sin^2\theta}$$

$$x^2+b^2=\frac{b^2}{\sin^2\theta}=\frac{x}{y}$$

$$y=\frac{x}{x^2+b^2}$$

Common Lorentzian function II.
(real "refractory" part)

(Preview of Unit 4.)



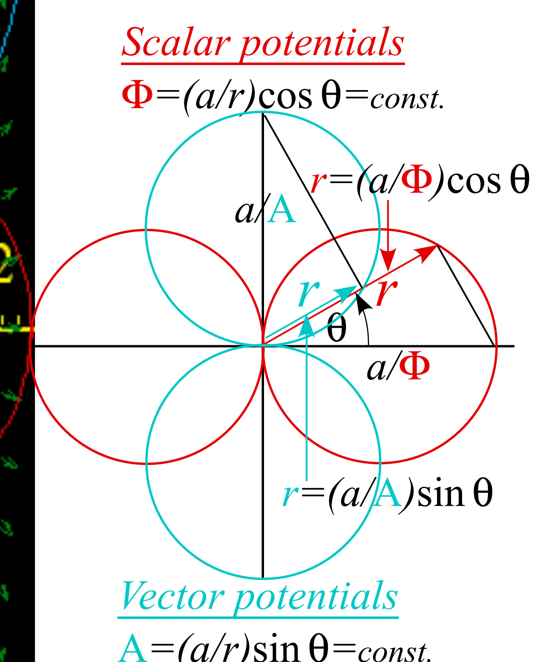
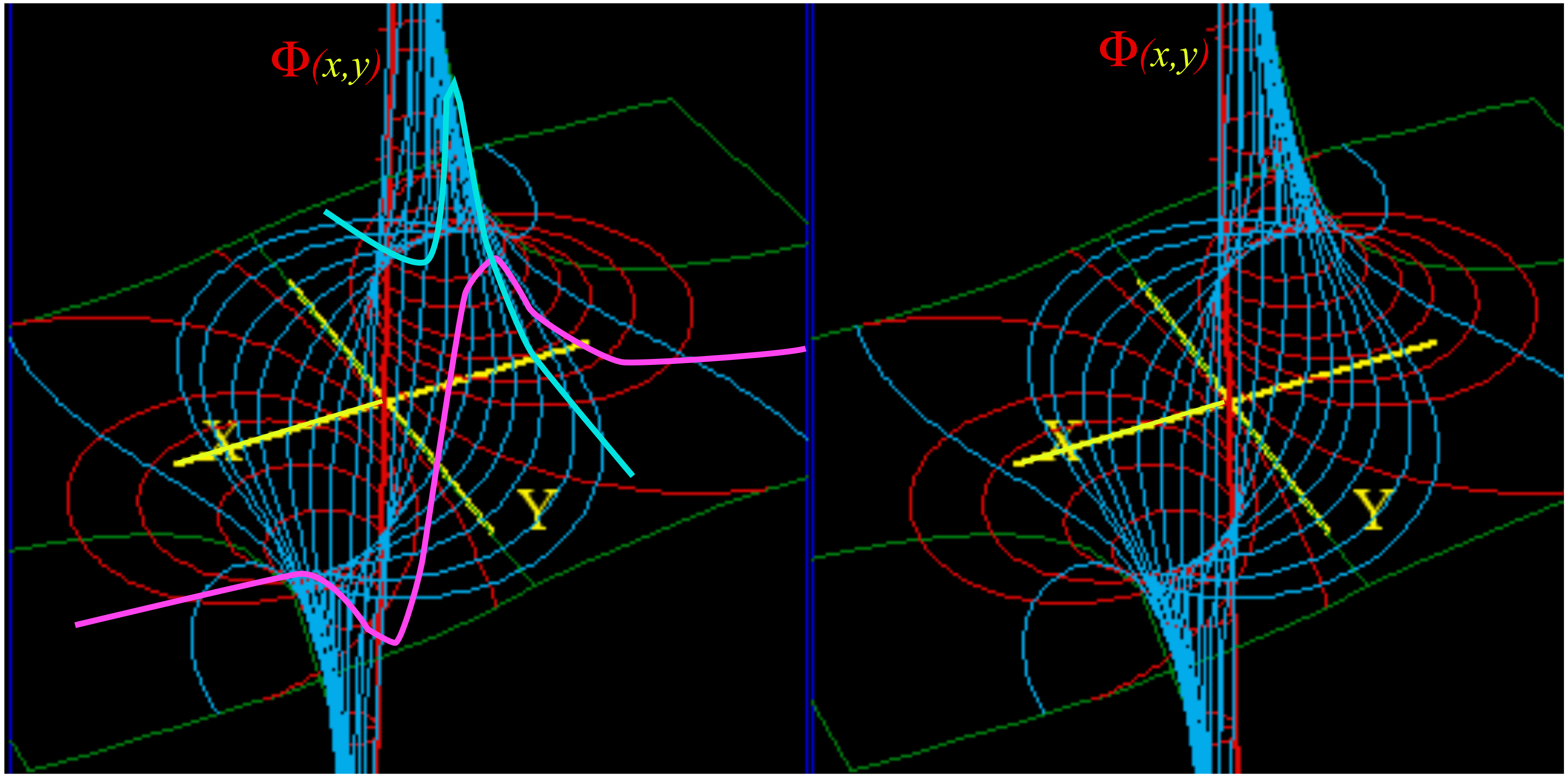


Fig. 10.11 Dipole **F***-field f(z)=1/z² and scalar potential (*Φ*=const.)-circles orthogonal to (*A*=const.)-circles.*

(Preview of Unit 4.)



From: Fig. 1.10.12

2^n*-pole analysis* *(quadrupole: 2^2=4-pole, octapole: 2^3=8-pole, …, pole dancer,*

What if we put a *(-)*copy of a *2-pole* near its original?

Well, the result is *4-pole* or *quadrupole* field $f^{\text{4-pole}}$ and potential $\phi^{\text{4-pole}}$.

Each a z-derivative of $f^{\text{2-pole}}$ and $\phi^{\text{2-pole}}$.

$$f^{\text{4-pole}} = \frac{a}{z^3} = \frac{1}{2}\frac{df^{\text{2-pole}}}{dz} = \frac{d\phi^{\text{4-pole}}}{dz} \qquad\qquad \phi^{\text{4-pole}} = -\frac{a}{2z^2} = \frac{1}{2}\frac{d\phi^{\text{2-pole}}}{dz}$$

2^n-pole analysis *(quadrupole:2^2=4-pole, octapole:2^3=8-pole, ..., pole dancer,*

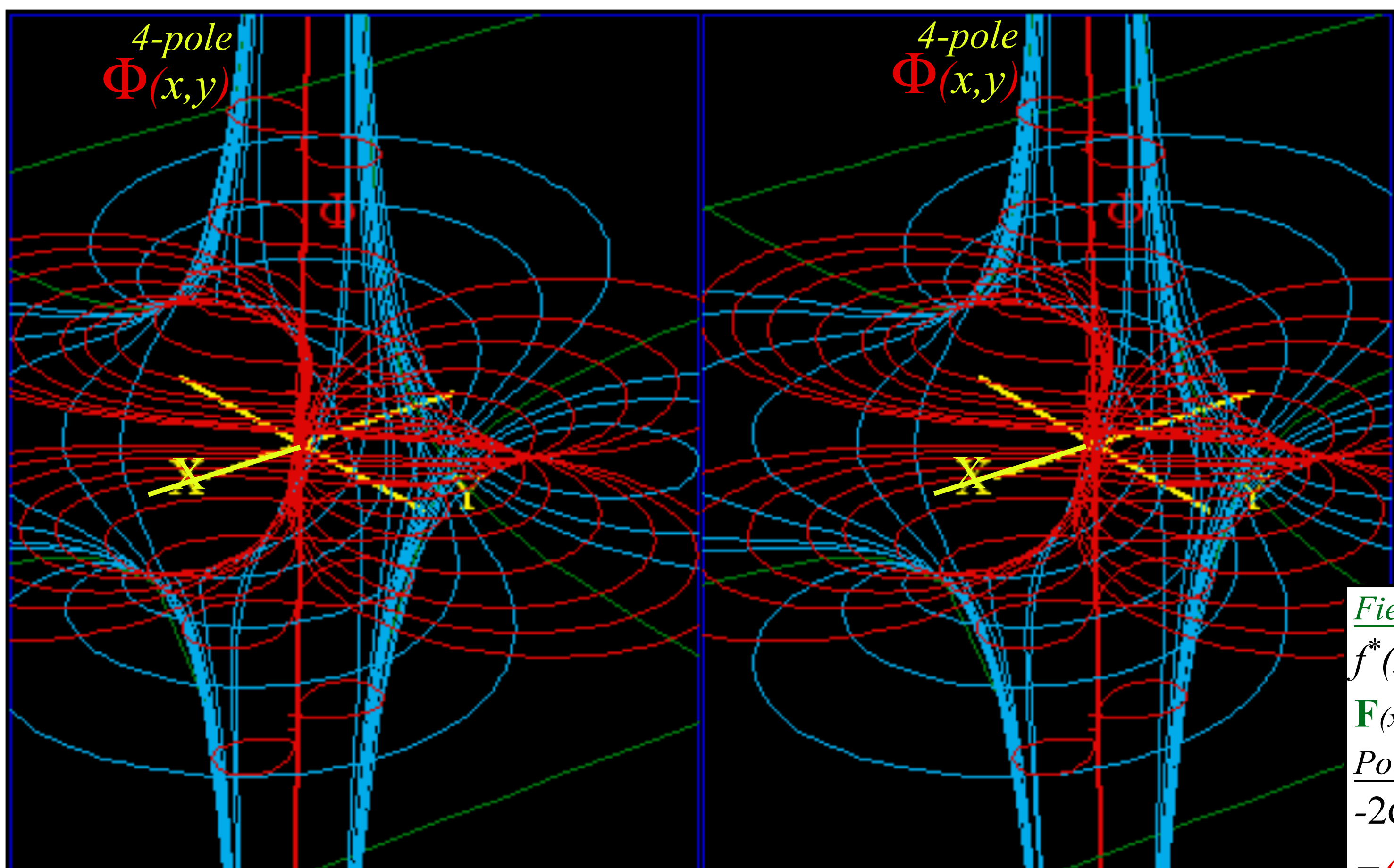
What if we put a *(-)*copy of a *2-pole* near its original?

Well, the result is *4-pole* or *quadrupole* field $f^{\text{4-pole}}$ and potential $\phi^{\text{4-pole}}$.

Each a *z*-derivative of $f^{\text{2-pole}}$ and $\phi^{\text{2-pole}}$, respectively.

$$f^{4\text{-}pole} = \frac{a}{z^3} = \frac{1}{2}\frac{df^{2\text{-}pole}}{dz} = \frac{d\phi^{4\text{-}pole}}{dz}$$

$$\phi^{4\text{-}pole} = -\frac{a}{2z^2} = \frac{1}{2}\frac{d\phi^{2\text{-}pole}}{dz}$$



Field:

$f^*(z^*)=1/z^3{}^*=e^{i3\theta}/r^3$

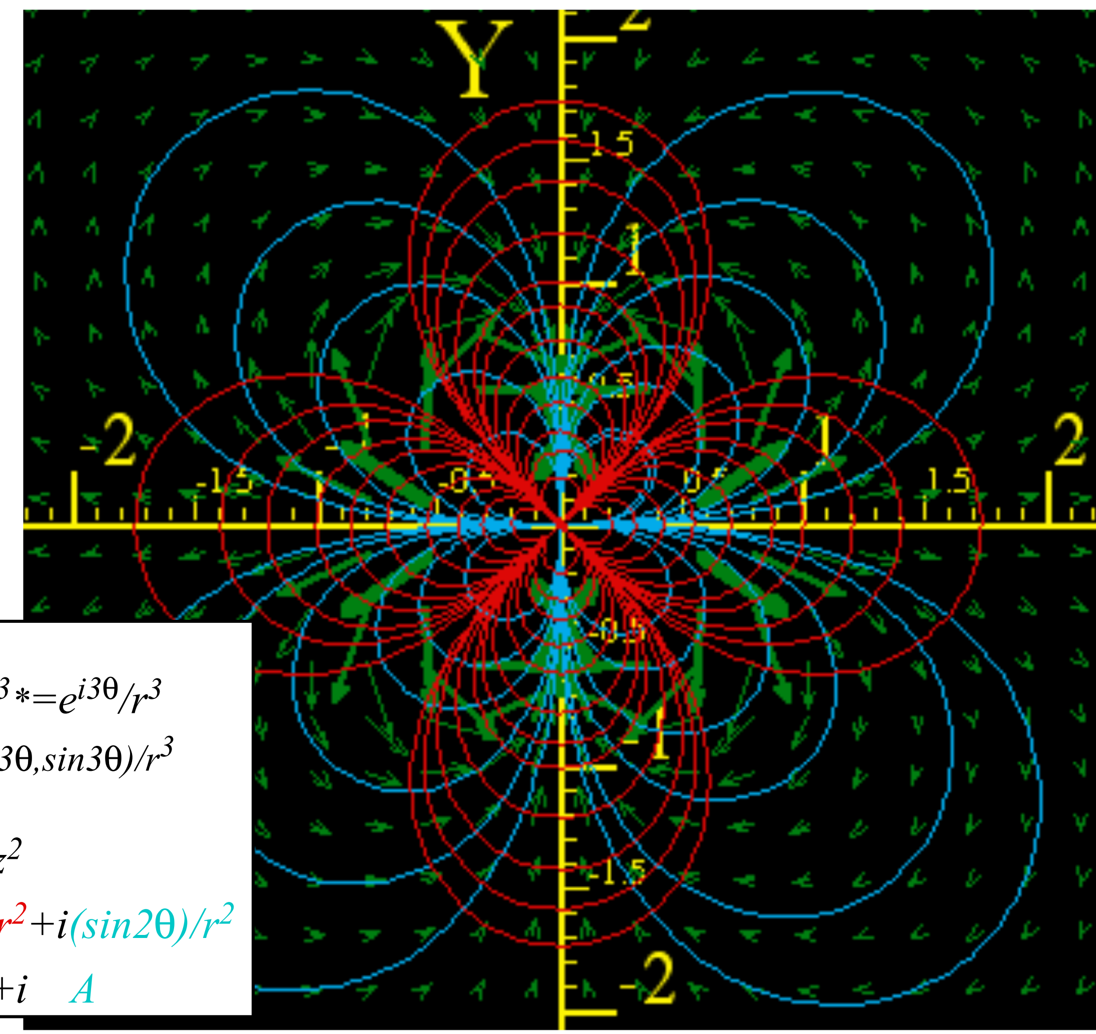
$\mathbf{F}_{(x,y)}=(\cos 3\theta,\sin 3\theta)/r^3$

Potential:

$-2\phi(z)=1/z^2$

$=(\cos 2\theta)/r^2+i(\sin 2\theta)/r^2$

$= \quad \Phi \quad +i \quad A$



4. Riemann-Cauchy conditions *What's analytic? (...and what's <u>not</u>?)*

Easy 2D circulation and flux integrals

Easy 2D curvilinear coordinate discovery

Easy 2D monopole, dipole, and 2^n-pole analysis

Easy 2^n-multipole field and potential expansion

Easy stereo-projection visualization

2^n-pole analysis: Laurent series (Generalization of Maclaurin-Taylor series)

Laurent series or *multipole expansion* of a given complex field function $f(z)$ around $z=0$.

$$\frac{d\phi}{dz} = f(z) = \ldots a_{-3}z^{-3} + a_{-2}z^{-2} + a_{-1}z^{-1} + a_0 + a_1 z + a_2 z^2 + a_3 z^3 + a_4 z^4 + a_5 z^5 + \ldots$$

$\cdots$	2^2-pole (quadrupole) at $z=0$	2^1-pole (dipole) at $z=0$	2^0-pole (monopole) at $z=0$	2^1-pole (dipole) at $z=\infty$	2^2-pole (quadrupole) at $z=\infty$	2^3-pole (octapole) at $z=\infty$	2^4-pole (hexadecapole) at $z=\infty$	2^5-pole at $z=\infty$	2^6-pole at $z=\infty$	$\cdots$

$$\int f\,dz = $$

$$\phi(z) = \ldots \frac{a_{-3}}{-2}z^{-2} + \frac{a_{-2}}{-1}z^{-1} + a_{-1}\ln z + a_0 z + \frac{a_1}{2}z^2 + \frac{a_2}{3}z^3 + \frac{a_3}{4}z^4 + \frac{a_4}{5}z^5 + \frac{a_5}{6}z^6 + \ldots$$

All field terms $a_{m-1}z^{m-1}$ except *1-pole* $\frac{a_{-1}}{z}$ have potential term $a_{m-1}z^m/m$ of a 2^m*-pole.*

These are located at $z=0$ for $m<0$ and at $z=\infty$ for $m>0$.

$$\phi(z) = \ldots \underbrace{\frac{a_{-4}}{-3}z^{-3}}_{(\text{octapole})_0} + \underbrace{\frac{a_{-3}}{-2}z^{-2}}_{(\text{quadrupole})_0} + \underbrace{\frac{a_{-2}}{-1}z^{-1}}_{(\text{dipole})_0} + \underbrace{a_{-1}\ln z}_{(\text{monopole})} + \underbrace{a_0 z}_{(\text{dipole})_\infty} + \underbrace{\frac{a_1}{2}z^2}_{(\text{quadrupole})_\infty} + \underbrace{\frac{a_2}{3}z^3}_{(\text{octapole})_\infty} + \ldots$$

Introducing the concepts of Inverse $w=1/z$ Map

2^n-pole analysis: Laurent series *(Generalization of Maclaurin-Taylor series)*

Laurent series or *multipole expansion* of a given complex field function $f(z)$ around $z=0$.

$$\frac{d\phi}{dz}=f(z)=\ldots a_{-3}z^{-3}+a_{-2}z^{-2}+a_{-1}z^{-1}+a_0+a_1 z+a_2 z^2+a_3 z^3+a_4 z^4+a_5 z^5+\ldots$$

$\cdots 2^2$-pole	2^1-pole	2^0-pole	2^1-pole	2^2-pole	2^3-pole	2^4-pole	2^5-pole	2^6-pole $\cdots$
(quadrupole)	(dipole)	(monopole)	(dipole)	(quadrupole)	(octapole)	(hexadecapole)		
at $z=0$	at $z=0$	at $z=0$	at $z=\infty$	at $z=\infty$	at $z=\infty$	at $z=\infty$	at $z=\infty$	at $z=\infty$

$\int f\,dz=$

$$\phi(z)=\ldots\frac{a_{-3}}{-2}z^{-2}+\frac{a_{-2}}{-1}z^{-1}+a_{-1}\ln z+a_0 z+\frac{a_1}{2}z^2+\frac{a_2}{3}z^3+\frac{a_3}{4}z^4+\frac{a_4}{5}z^5+\frac{a_5}{6}z^6+\ldots$$

All field terms $a_{m-1}z^{m-1}$ except *1-pole* $\frac{a_{-1}}{z}$ have potential term $a_{m-1}z^m/m$ of a *2^m-pole*.

These are located at $z=0$ for $m<0$ and at $z=\infty$ for $m>0$.

(octapole)$_0$ (quadrupole)$_0$ (dipole)$_0$ (monopole) (dipole)$_\infty$ (quadrupole)$_\infty$ (octapole)$_\infty$

$$\phi(z)=\ldots\frac{a_{-4}}{-3}z^{-3}+\frac{a_{-3}}{-2}z^{-2}+\frac{a_{-2}}{-1}z^{-1}+a_{-1}\ln z+a_0 z+\frac{a_1}{2}z^2+\frac{a_2}{3}z^3+\ldots$$

$$\phi(w)=\ldots\frac{a_{-4}}{-3}w^{-3}+\frac{a_{-3}}{-2}w^{-2}+\frac{a_{-2}}{-1}w^{-1}+a_{-1}\ln w+a_0 w+\frac{a_1}{2}w^2+\frac{a_2}{3}w^3+\ldots$$

(*with* $z=w^{-1}$)

Introducing the concepts of Inverse $w=1/z$ Map

2^n-pole analysis: Laurent series *(Generalization of Maclaurin-Taylor series)*

Laurent series or *multipole expansion* of a given complex field function $f(z)$ around $z=0$.

$$\frac{d\phi}{dz}=f(z)=\ldots a_{-3}z^{-3}+a_{-2}z^{-2}+a_{-1}z^{-1}+a_0+a_1 z+a_2 z^2+a_3 z^3+a_4 z^4+a_5 z^5+\ldots$$

$\cdots$ 2^2-pole	2^1-pole	2^0-pole	2^1-pole	2^2-pole	2^3-pole	2^4-pole	2^5-pole	2^6-pole $\cdots$
(quadrupole)	*(dipole)*	*(monopole)*	*(dipole)*	*(quadrupole)*	*(octapole)*	*(hexadecapole)*		
at $z=0$	at $z=0$	at $z=0$	at $z=\infty$	at $z=\infty$	at $z=\infty$	at $z=\infty$	at $z=\infty$	at $z=\infty$

$$\int f\,dz=$$

$$\phi(z)=\ldots\frac{a_{-3}}{-2}z^{-2}+\frac{a_{-2}}{-1}z^{-1}+a_{-1}\ln z+a_0 z+\frac{a_1}{2}z^2+\frac{a_2}{3}z^3+\frac{a_3}{4}z^4+\frac{a_4}{5}z^5+\frac{a_5}{6}z^6+\ldots$$

All field terms $a_{m-1}z^{m-1}$ except *1-pole* $\frac{a_{-1}}{z}$ have potential term $a_{m-1}z^m/m$ of a 2^m*-pole.*

These are located at $z=0$ for $m<0$ and at $z=\infty$ for $m>0$.

(octapole)$_0$	*(quadrupole)*$_0$	*(dipole)*$_0$	*(monopole)*	*(dipole)*$_\infty$	*(quadrupole)*$_\infty$	*(octapole)*$_\infty$

$$\phi(z)=\ldots\frac{a_{-4}}{-3}z^{-3}+\frac{a_{-3}}{-2}z^{-2}+\frac{a_{-2}}{-1}z^{-1}+a_{-1}\ln z+a_0 z+\frac{a_1}{2}z^2+\frac{a_2}{3}z^3+\ldots$$

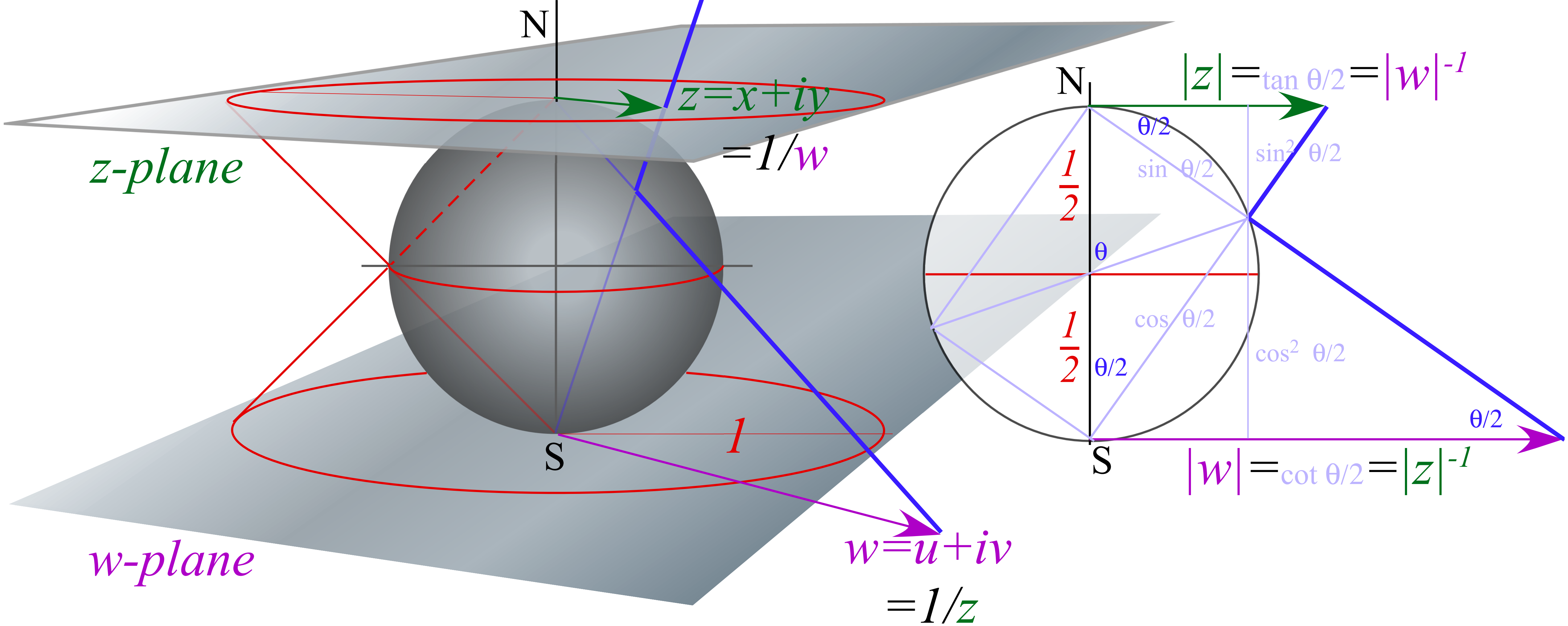
$$\phi(w)=\ldots\frac{a_{-4}}{-3}w^{-3}+\frac{a_{-3}}{-2}w^{-2}+\frac{a_{-2}}{-1}w^{-1}+a_{-1}\ln w+a_0 w+\frac{a_1}{2}w^2+\frac{a_2}{3}w^3+\ldots$$

(*with* $z\rightarrow w$)

$$=\ldots\frac{a_2}{3}z^{-3}+\frac{a_1}{2}z^{-2}+a_0 z^{-1}-a_{-1}\ln z+\frac{a_{-2}}{-1}z+\frac{a_{-3}}{-2}z^2+\frac{a_{-4}}{-3}z^3+\ldots$$

(*with* $w=z^{-1}$)

Introducing the concepts of Inverse $w=1/z$ Map



$$\phi(z) = \ldots \frac{a_{-4}}{-3} z^{-3} + \frac{a_{-3}}{-2} z^{-2} + \frac{a_{-2}}{-1} z^{-1} + a_{-1} \ln z + a_0 z + \frac{a_1}{2} z^2 + \frac{a_2}{3} z^3 + \ldots$$

(octapole)₀ *(quadrupole)₀* *(dipole)₀* *(monopole)* *(dipole)∞* *(quadrupole)∞* *(octapole)∞*

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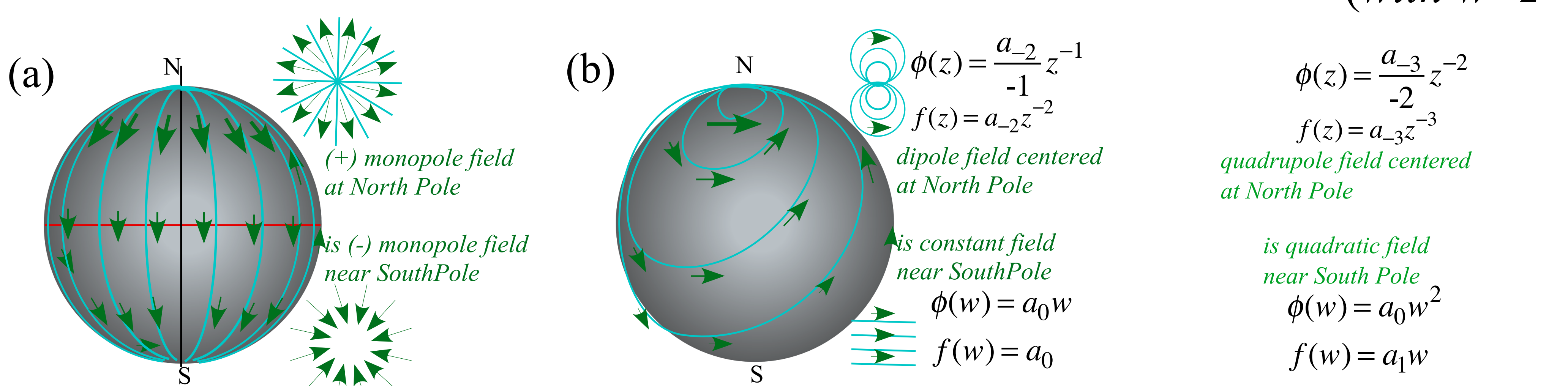
(with $z \to w$)

$$= \ldots \frac{a_2}{3} z^{-2} + \frac{a_1}{2} z^{-2} + a_0 z^{-1} - a_{-1} \ln z + \frac{a_{-2}}{-1} z + \frac{a_{-3}}{-2} z^2 + \frac{a_{-4}}{-3} z^3 + \ldots$$

(with $w=z^{-1}$)



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$$f(z)=\sum_{n=-\infty}^{\infty}a_n(z-a)^n \qquad where:\;a_n=\frac{1}{2\pi i}\oint\frac{f(z)}{(z-a)^{n+1}}dz\left(=\frac{1}{n!}\frac{d^nf(a)}{da^n}\quad for:\;n\geq 0\right)$$

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Source analysis starts with 1-pole loop integrals $\oint z^{-1}dz = 2\pi i$ or, with origin shifted $\oint (z-a)^{-1}dz = 2\pi i$.

They hold for any loop about point-a. Function $f(z)$ is just $f(a)$ on a *tiny* circle around point-a.

(assume *tiny* circle around $z=a$)

$$\oint \frac{f(z)}{z-a}dz = \oint \frac{f(a)}{z-a}dz = f(a)\oint \frac{1}{z-a}dz = 2\pi i f(a)$$

(but any contour that doesn't "touch a gives same answer)

$$\boxed{f(a) = \frac{1}{2\pi i}\oint \frac{f(z)}{z-a}dz}$$

The $f(a)$ result is called a *Cauchy integral*. Then repeated a-derivatives gives a sequence of them.

$$\frac{df(a)}{da} = \frac{1}{2\pi i}\oint \frac{f(z)}{(z-a)^2}dz \ , \ \frac{d^2 f(a)}{da^2} = \frac{2}{2\pi i}\oint \frac{f(z)}{(z-a)^3}dz \ , \ \frac{d^3 f(a)}{da^3} = \frac{3!}{2\pi i}\oint \frac{f(z)}{(z-a)^4}dz, \ \cdots, \frac{d^n f(a)}{da^n} = \frac{n!}{2\pi i}\oint \frac{f(z)}{(z-a)^{n+1}}dz$$

This leads to a general *Taylor-Laurent* power series expansion of function $f(z)$ around point-a.

$$f(z) = \sum_{n=-\infty}^{\infty} a_n (z-a)^n \qquad where:\ a_n = \frac{1}{2\pi i}\oint \frac{f(z)}{(z-a)^{n+1}}dz \left(= \frac{1}{n!}\frac{d^n f(a)}{da^n} \quad for:\ n \geq 0 \right)$$

(quadrupole)₀ (dipole)₀ (monopole) (dipole)∞ (quadrupole)∞ (octapole)∞ (hexadecapole)∞ ...

$$f(z) = \ldots a_{-3}z^{-3} + \underset{\text{dipole moment}}{a_{-2}z^{-2}} + \underset{\text{monopole moment}}{a_{-1}z^{-1}} + a_0 + a_1 z + a_2 z^2 + a_3 z^3 + a_4 z^4 + a_5 z^5 + \ldots$$

5. Mapping and Non-analytic 2D source field analysis

The half-n'-half results

are called

Riemann-Cauchy

Derivative Relations

$$\boxed{\begin{aligned}\frac{\partial \Phi}{\partial x} &= \frac{\partial \mathbf{A}}{\partial y} \\ \frac{\partial \Phi}{\partial y} &= -\frac{\partial \mathbf{A}}{\partial x}\end{aligned}} \begin{aligned}\text{is:} \\ \text{is:}\end{aligned} \boxed{\begin{aligned}\frac{\partial \mathrm{Re}\phi(z)}{\partial x} &= \frac{\partial \mathrm{Im}\phi(z)}{\partial y} \\ \frac{\partial \mathrm{Re}\phi(z)}{\partial y} &= -\frac{\partial \mathrm{Im}\phi(z)}{\partial x}\end{aligned}} \begin{aligned}\text{or:} \\ \text{or:}\end{aligned} \boxed{\begin{aligned}\frac{\partial \mathrm{Re}f(z)}{\partial x} &= \frac{\partial \mathrm{Im}f(z)}{\partial y} \\ \frac{\partial \mathrm{Re}f(z)}{\partial y} &= -\frac{\partial \mathrm{Im}f(z)}{\partial x}\end{aligned}} \begin{aligned}\text{is:} \\ \text{is:}\end{aligned} \boxed{\begin{aligned}\frac{\partial f_x(z)}{\partial x} &= \frac{\partial f_y(z)}{\partial y} \\ \frac{\partial f_x(z)}{\partial y} &= -\frac{\partial f_y(z)}{\partial x}\end{aligned}}$$

RC applies to analytic potential $\phi(z) = \Phi + i\mathbf{A}$ *and analytic field* $f(z) = f_x + if_y$ *and any analytic function*

The *half*-n'-*half* results

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Derivative Relations

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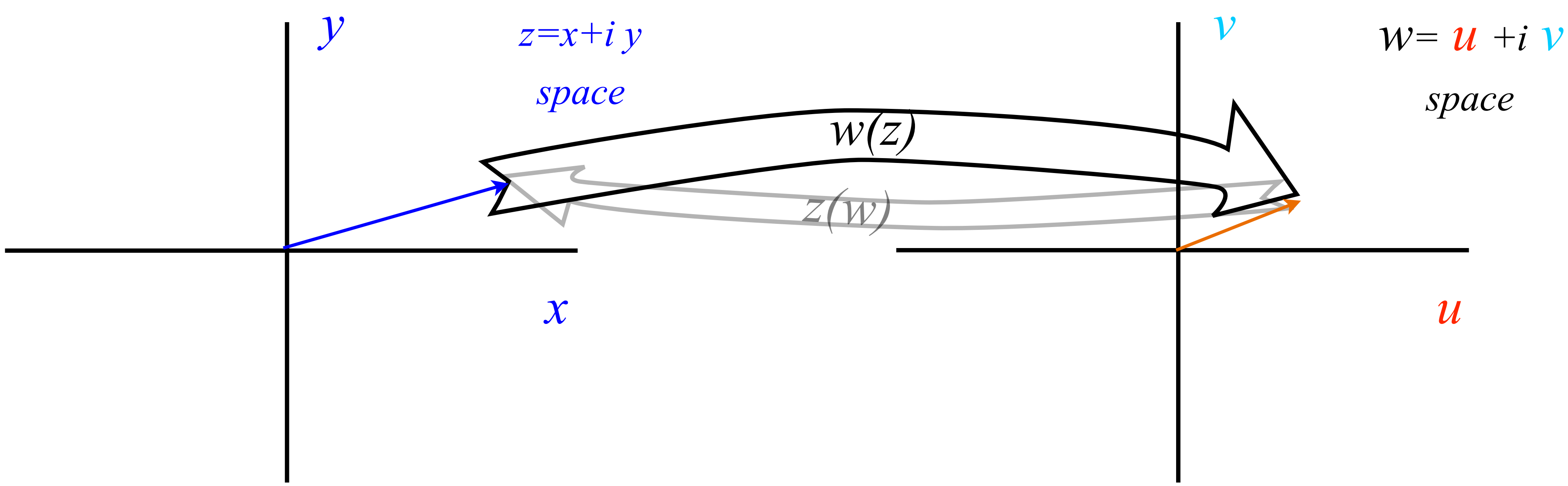
RC applies to analytic potential $\phi(z)=\Phi+iA$ *and analytic field* $f(z)=f_x+if_y$ *and any analytic function*

Common notation for mapping: $w(z)=u+iv$



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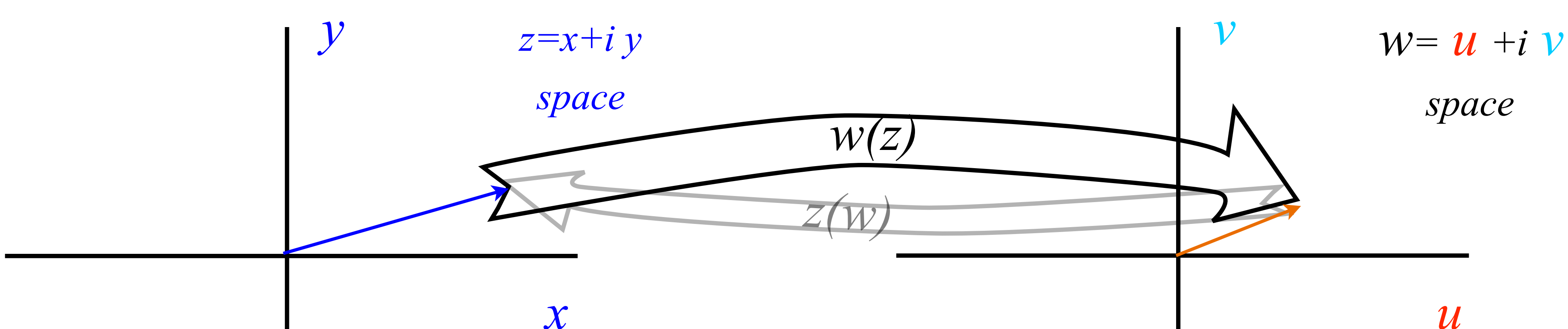

$$du=\frac{\partial u}{\partial x}dx+\frac{\partial u}{\partial y}dy$$

$$dv=\frac{\partial v}{\partial x}dx+\frac{\partial v}{\partial y}dy$$

Jacobian for mapping:

$$\begin{pmatrix}du\\ dv\end{pmatrix}=\begin{pmatrix}\frac{\partial u}{\partial x} & \frac{\partial u}{\partial y}\\ \frac{\partial v}{\partial x} & \frac{\partial v}{\partial y}\end{pmatrix}\begin{pmatrix}dx\\ dy\end{pmatrix}$$

Complex derivative for mapping:

$$\frac{dw}{dz}=\frac{1}{2}\left(\frac{\partial}{\partial x}-i\frac{\partial}{\partial y}\right)(u+iv)=\frac{1}{2}\left(\frac{\partial u}{\partial x}+\frac{\partial v}{\partial y}\right)+\frac{i}{2}\left(\frac{\partial v}{\partial x}-\frac{\partial u}{\partial y}\right)$$

The half-n'-half results are called Riemann-Cauchy Derivative Relations

$$\boxed{\begin{array}{c}\frac{\partial \Phi}{\partial x}=\frac{\partial A}{\partial y}\\[1em] \frac{\partial \Phi}{\partial y}=-\frac{\partial A}{\partial x}\end{array}}\;\begin{array}{c}\text{is:}\\[1em]\text{is:}\end{array}\;\boxed{\begin{array}{c}\frac{\partial \mathrm{Re}\,\phi(z)}{\partial x}=\frac{\partial \mathrm{Im}\,\phi(z)}{\partial y}\\[1em] \frac{\partial \mathrm{Re}\,\phi(z)}{\partial y}=-\frac{\partial \mathrm{Im}\,\phi(z)}{\partial x}\end{array}}\;\begin{array}{c}\text{or:}\\[1em]\text{or:}\end{array}\;\boxed{\begin{array}{c}\frac{\partial \mathrm{Re}\,f(z)}{\partial x}=\frac{\partial \mathrm{Im}\,f(z)}{\partial y}\\[1em] \frac{\partial \mathrm{Re}\,f(z)}{\partial y}=-\frac{\partial \mathrm{Im}\,f(z)}{\partial x}\end{array}}\;\begin{array}{c}\text{is:}\\[1em]\text{is:}\end{array}\;\boxed{\begin{array}{c}\frac{\partial f_x(z)}{\partial x}=\frac{\partial f_y(z)}{\partial y}\\[1em] \frac{\partial f_x(z)}{\partial y}=-\frac{\partial f_y(z)}{\partial x}\end{array}}$$

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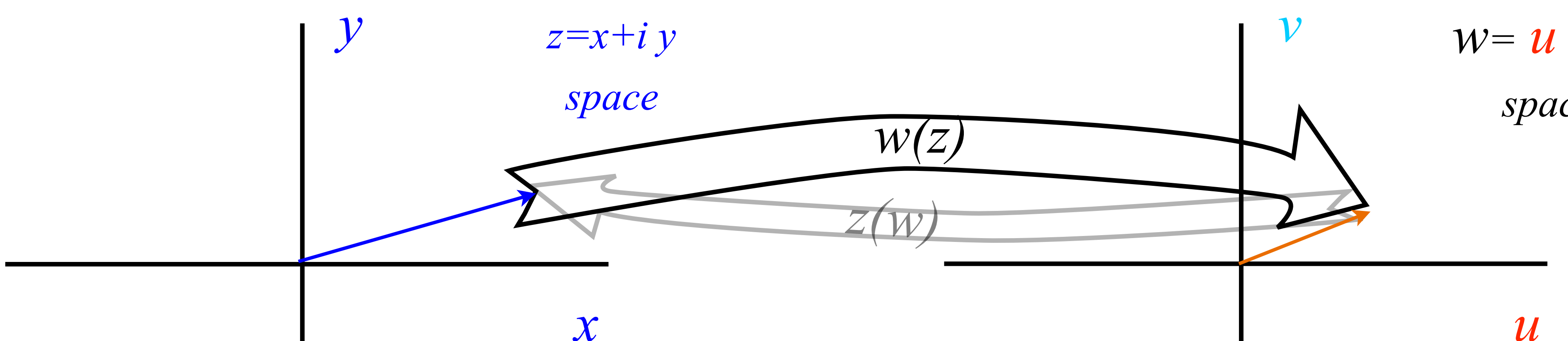

$$du=\frac{\partial u}{\partial x}dx+\frac{\partial u}{\partial y}dy$$

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$$=\begin{pmatrix} \frac{\partial u}{\partial x} & \frac{\partial u}{\partial y} \\ -\frac{\partial u}{\partial y} & \frac{\partial u}{\partial x} \end{pmatrix}\begin{pmatrix} dx \\ dy \end{pmatrix}=\begin{pmatrix} \frac{\partial v}{\partial y} & -\frac{\partial v}{\partial x} \\ \frac{\partial v}{\partial x} & \frac{\partial v}{\partial y} \end{pmatrix}\begin{pmatrix} dx \\ dy \end{pmatrix}$$

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$$=\frac{\partial u}{\partial x}-i\frac{\partial u}{\partial y} \qquad =\frac{\partial v}{\partial y}+i\frac{\partial v}{\partial x}$$

Complex derivative abs-square:

$$\left|\frac{dw}{dz}\right|^2=\left(\frac{\partial u}{\partial x}\right)^2+\left(\frac{\partial u}{\partial y}\right)^2=\left(\frac{\partial v}{\partial y}\right)^2+\left(\frac{\partial v}{\partial x}\right)^2$$

The half-n'-half results

are called

Riemann-Cauchy

Derivative Relations

$$\boxed{\frac{\partial \Phi}{\partial x} = \frac{\partial A}{\partial y}} \text{ is: } \boxed{\frac{\partial \mathrm{Re}\,\phi(z)}{\partial x} = \frac{\partial \mathrm{Im}\,\phi(z)}{\partial y}} \text{ or: } \boxed{\frac{\partial \mathrm{Re}\,f(z)}{\partial x} = \frac{\partial \mathrm{Im}\,f(z)}{\partial y}} \text{ is: } \boxed{\frac{\partial f_x(z)}{\partial x} = \frac{\partial f_y(z)}{\partial y}}$$

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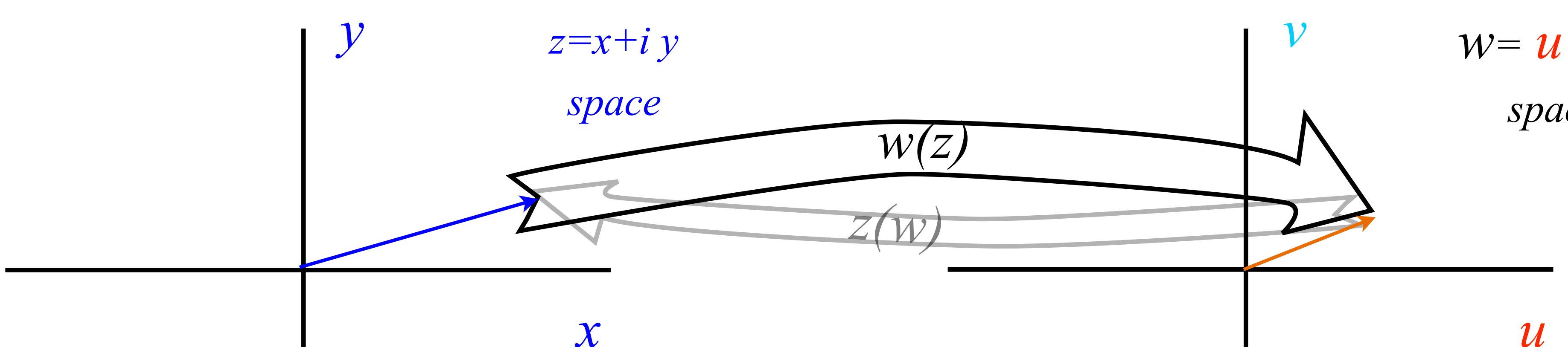

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$$\begin{pmatrix} du \\ dv \end{pmatrix} = \begin{pmatrix} \frac{\partial u}{\partial x} & \frac{\partial u}{\partial y} \\ \frac{\partial v}{\partial x} & \frac{\partial v}{\partial y} \end{pmatrix} \begin{pmatrix} dx \\ dy \end{pmatrix}$$

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Complex derivative for mapping:

$$\frac{dw}{dz} = \frac{1}{2}\left(\frac{\partial}{\partial x} - i\frac{\partial}{\partial y}\right)(u+iv) = \frac{1}{2}\left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y}\right) + \frac{i}{2}\left(\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}\right)$$

$$= \frac{\partial u}{\partial x} - i\frac{\partial u}{\partial y} \qquad = \frac{\partial v}{\partial y} + i\frac{\partial v}{\partial x}$$

Complex derivative abs-square:

$$\left|\frac{dw}{dz}\right|^2 = \left(\frac{\partial u}{\partial x}\right)^2 + \left(\frac{\partial u}{\partial y}\right)^2 = \left(\frac{\partial v}{\partial y}\right)^2 + \left(\frac{\partial v}{\partial x}\right)^2 = \det|J|$$

...equals Jacobian Determinant

The half-n'-half results are called Riemann-Cauchy Derivative Relations

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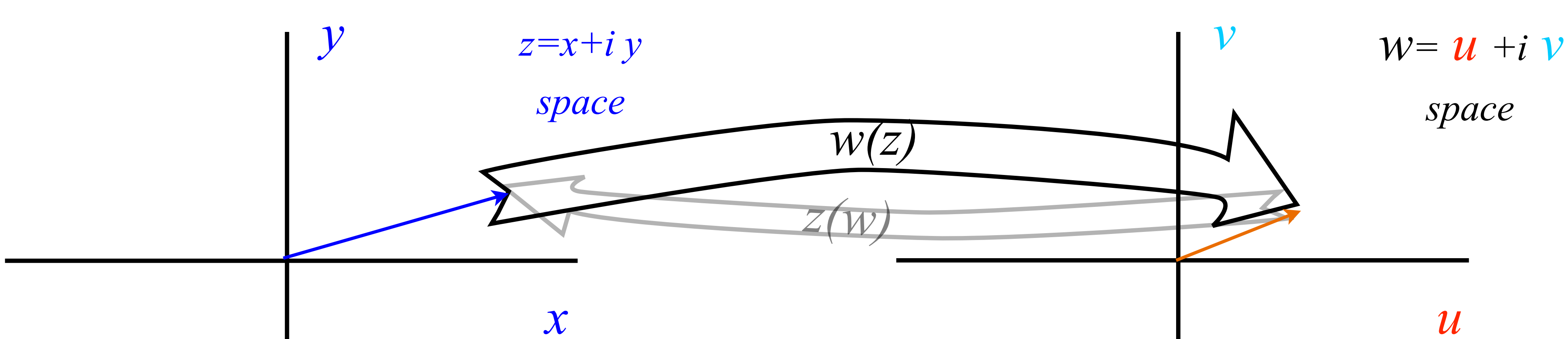


Important result:

$$dw=\sqrt{J}\cdot e^{i\theta}\cdot dz$$

is scaled rotation of dz.

Jacobian for mapping is scaled rotation:

$$du=\frac{\partial u}{\partial x}dx+\frac{\partial u}{\partial y}dy$$
$$dv=\frac{\partial v}{\partial x}dx+\frac{\partial v}{\partial y}dy$$

$$\begin{pmatrix}du\\dv\end{pmatrix}=\begin{pmatrix}\frac{\partial u}{\partial x}&\frac{\partial u}{\partial y}\\ \frac{\partial v}{\partial x}&\frac{\partial v}{\partial y}\end{pmatrix}\begin{pmatrix}dx\\dy\end{pmatrix}=\sqrt{\det J}\begin{pmatrix}\cos\theta&-\sin\theta\\ \sin\theta&\cos\theta\end{pmatrix}\begin{pmatrix}dx\\dy\end{pmatrix}$$

$$=\begin{pmatrix}\frac{\partial u}{\partial x}&\frac{\partial u}{\partial y}\\ -\frac{\partial u}{\partial y}&\frac{\partial u}{\partial x}\end{pmatrix}\begin{pmatrix}dx\\dy\end{pmatrix}=\begin{pmatrix}\frac{\partial v}{\partial y}&-\frac{\partial v}{\partial x}\\ \frac{\partial v}{\partial x}&\frac{\partial v}{\partial y}\end{pmatrix}\begin{pmatrix}dx\\dy\end{pmatrix}$$

Complex derivative for mapping:

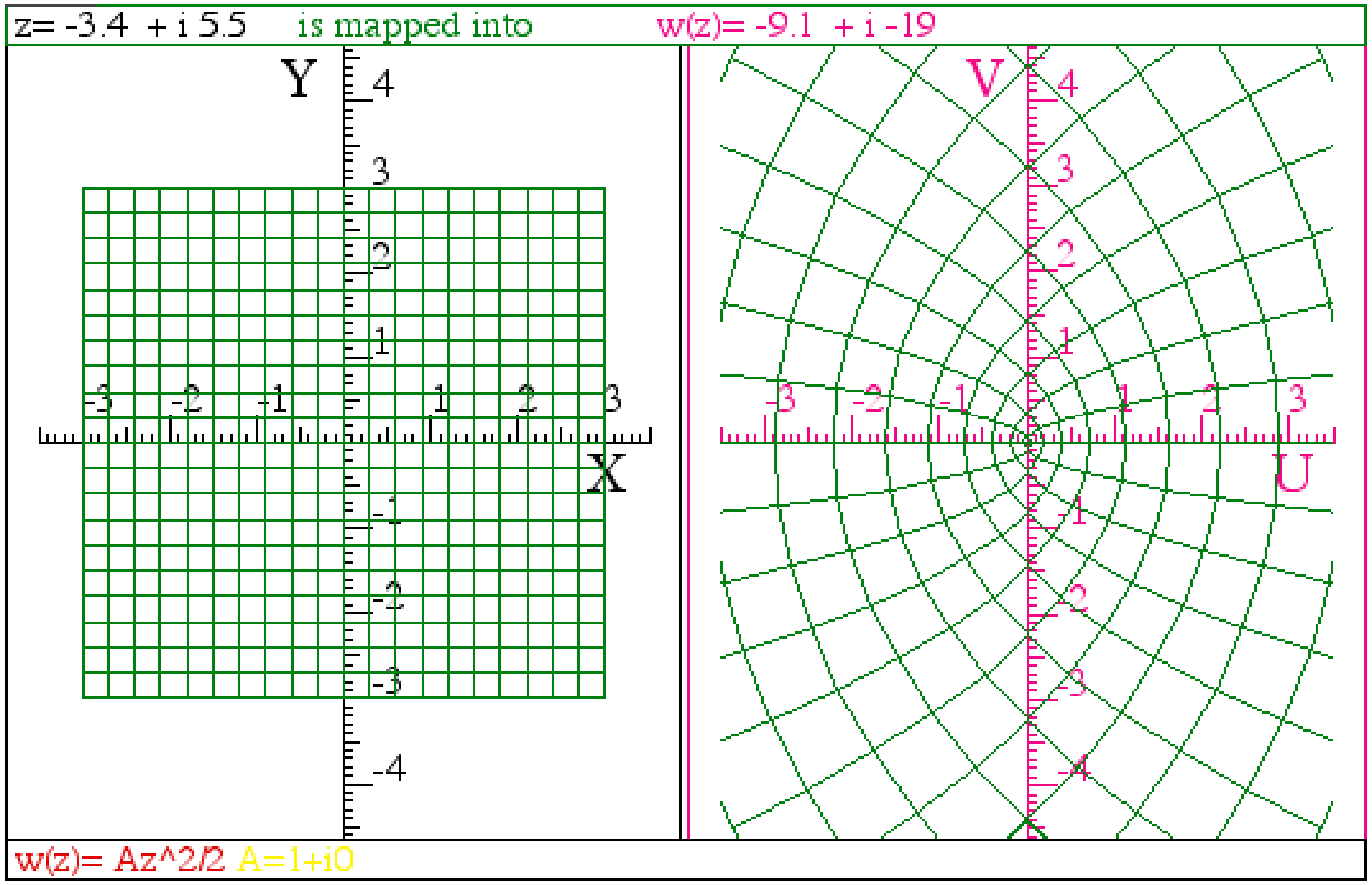
$$\frac{dw}{dz}=\frac{1}{2}\left(\frac{\partial}{\partial x}-i\frac{\partial}{\partial y}\right)(u+iv)=\frac{1}{2}\left(\frac{\partial u}{\partial x}+\frac{\partial v}{\partial y}\right)+\frac{i}{2}\left(\frac{\partial v}{\partial x}-\frac{\partial u}{\partial y}\right)$$
$$=\frac{\partial u}{\partial x}-i\frac{\partial u}{\partial y}\qquad=\frac{\partial v}{\partial y}+i\frac{\partial v}{\partial x}$$

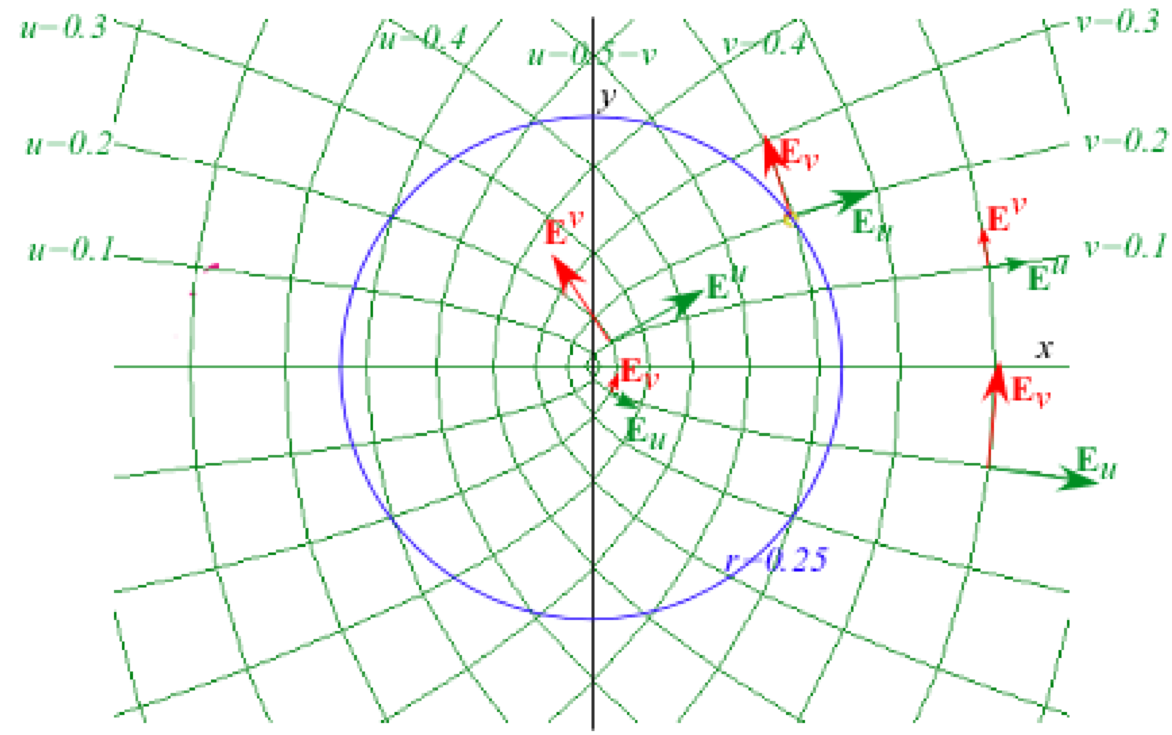
Complex derivative abs-square:

$$\left|\frac{dw}{dz}\right|^2=\left(\frac{\partial u}{\partial x}\right)^2+\left(\frac{\partial u}{\partial y}\right)^2=\left(\frac{\partial v}{\partial y}\right)^2+\left(\frac{\partial v}{\partial x}\right)^2=\det|J|$$

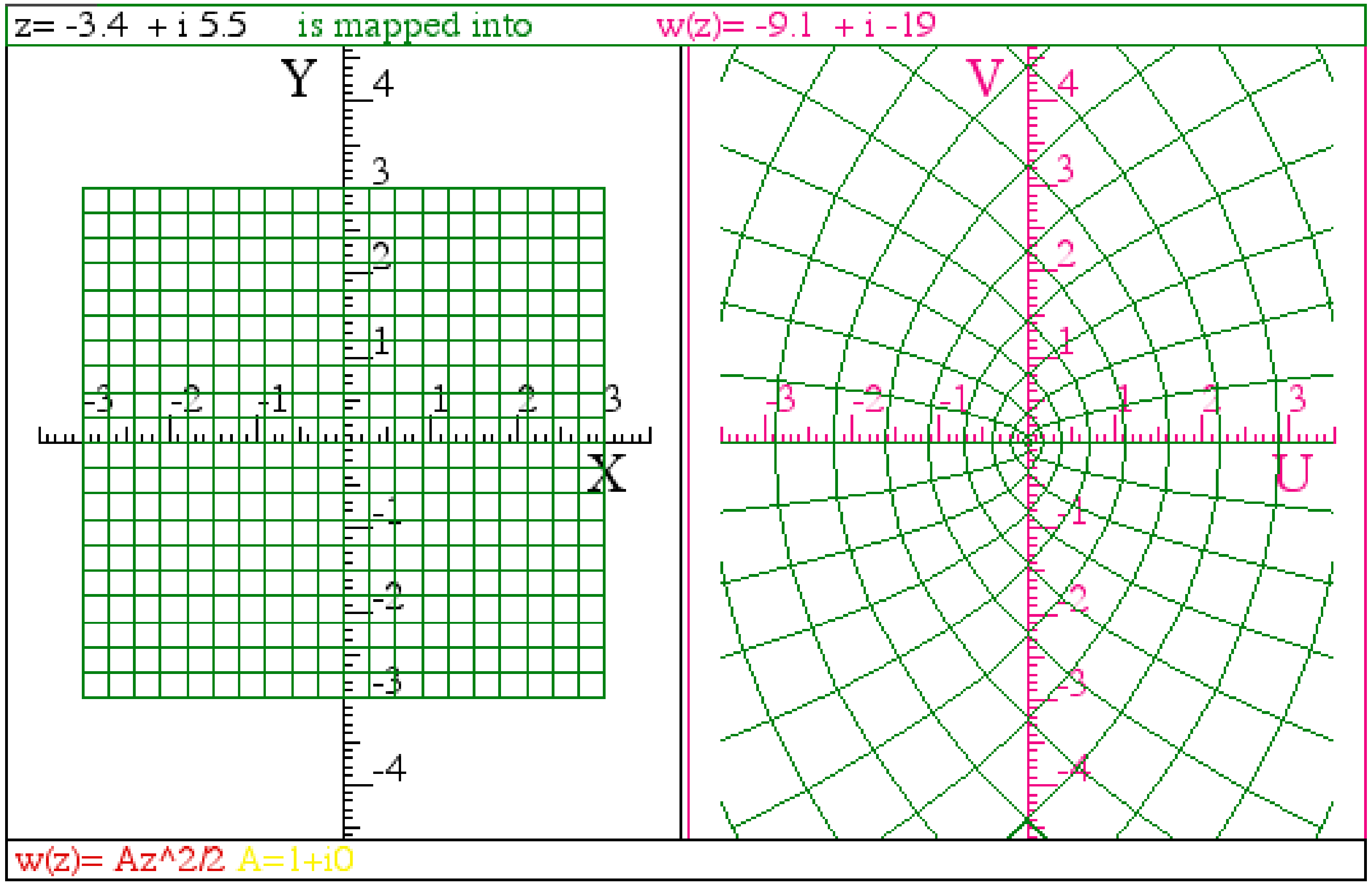
...equals Jacobian Determinant

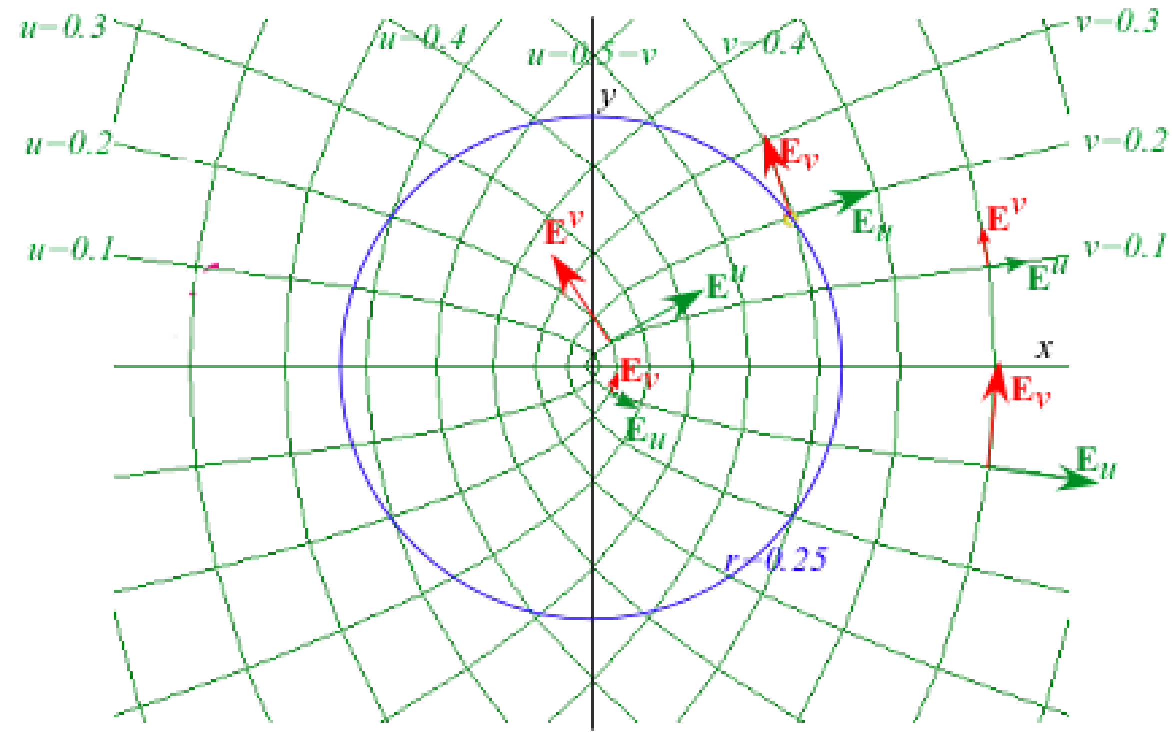
$w(z)=z^2$ gives parabolic OCC



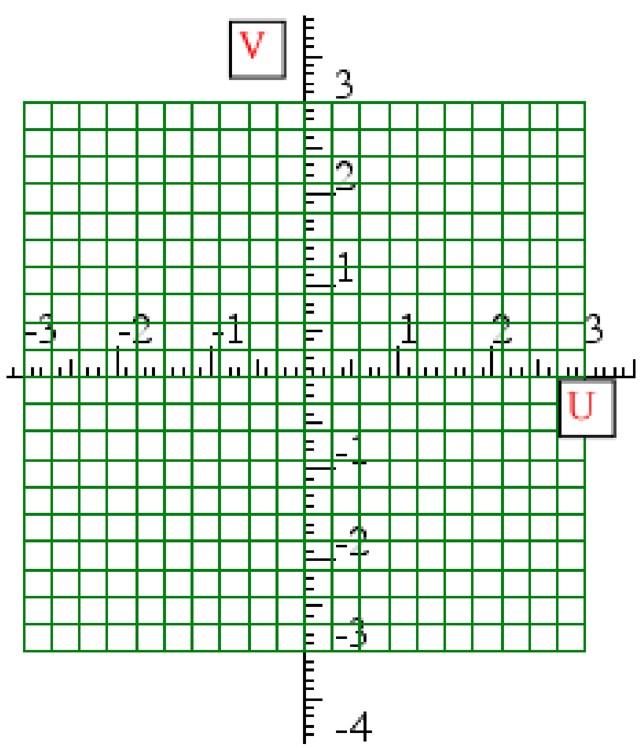


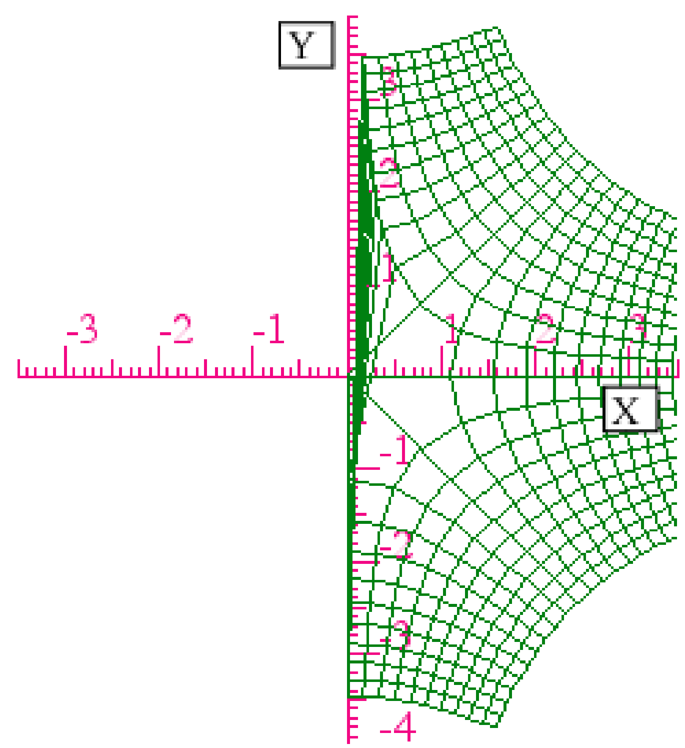
$w(z) = z^2$ gives parabolic OCC





Inverse: $z(w) = w^{1/2}$ gives hyperbolic OCC





5. Mapping and Non-analytic 2D source field analysis

Non-analytic potential, force, and source field functions (Excerpts of Unit 1-Ch.10 and AnalyIt)

A general 2D complex field may have:

1. non-analytic *potential field function* $\phi(z,z^*)=\Phi(x,y)+iA(x,y)$,
2. non-analytic *force field function* $f(z,z^*) = fx(x,y) + ify(x,y)$,
3. non-analytic *source distribution function* $s(z,z^*) = \rho(x,y) + i\, I(x,y)$.

Non-analytic potential, force, and source field functions (Excerpts of Unit 1-Ch.10 and AnalyIt)

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Source definitions generalize source-free fields ($\frac{df(z^*)}{dz}=0=\frac{df(z)}{dz^*}$) based on relations. $\frac{df}{dz}=\frac{1}{2}(\frac{\partial f_x}{\partial x}+\frac{\partial f_y}{\partial y})+\frac{i}{2}(\frac{\partial f_y}{\partial x}-\frac{\partial f_x}{\partial y})=\frac{1}{2}\nabla\bullet\mathbf{F}+\frac{i}{2}|\nabla\times\mathbf{F}|$

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$$2\frac{df^*}{dz}=s^*(z,z^*)$$

$$2\frac{df}{dz^*}=s(z,z^*)$$

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$$2\frac{df^*}{dz}=s^*(z,z^*) \qquad\qquad 2\frac{df}{dz^*}=s(z,z^*)$$

Field-f-from-potential-ϕ equations are like the older ($f(z)=\frac{d\phi}{dz}$ or $f^*(z^*)=\frac{d\phi^*}{dz^*}$) but with an extra factor of *2*.

$$2\frac{d\phi}{dz}=f(z,z^*) \qquad\qquad 2\frac{d\phi^*}{dz^*}=f^*(z,z^*)$$

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$$2\frac{d\phi}{dz}=f(z,z^*) \qquad 2\frac{d\phi^*}{dz^*}=f^*(z,z^*)$$

The new source equations expand into a real and imaginary parts that are divergence and curl terms, respectively.

$$s^*(z,z^*)=2\frac{df^*}{dz}=\left[\frac{\partial}{\partial x}-i\frac{\partial}{\partial y}\right]\left[f_x^*(x,y)+if_y^*(x,y)\right]=\rho-i\,I\ , \quad \text{where: } f_x^*=f_x\ , \text{ and: } f_y^*=-f_y$$

Non-analytic potential, force, and source field functions (Excerpts of Unit 1-Ch.10 and AnalyIt)

A general 2D complex field may have:

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Source definitions generalize source-free fields ($\frac{df(z^*)}{dz}=0=\frac{df(z)}{dz^*}$) based on relations. $\frac{df}{dz}=\frac{1}{2}(\frac{\partial f_x}{\partial x}+\frac{\partial f_y}{\partial y})+\frac{i}{2}(\frac{\partial f_y}{\partial x}-\frac{\partial f_x}{\partial y})=\frac{1}{2}\nabla\bullet\mathbf{F}+\frac{i}{2}|\nabla\times\mathbf{F}|$

$$2\frac{df^*}{dz}=s^*(z,z^*) \qquad\qquad 2\frac{df}{dz^*}=s(z,z^*)$$

Field-f-from-potential-ϕ equations are like the older ($f(z)=\frac{d\phi}{dz}$ or $f^*(z*)=\frac{d\phi*}{dz*}$) but with an extra factor of *2*.

$$2\frac{d\phi}{dz}=f(z,z^*) \qquad\qquad 2\frac{d\phi^*}{dz^*}=f^*(z,z^*)$$

The new source equations expand into a real and imaginary parts that are divergence and curl terms, respectively.

$$s^*(z,z^*)=2\frac{df^*}{dz}=\left[\frac{\partial}{\partial x}-i\frac{\partial}{\partial y}\right]\left[f_x^*(x,y)+if_y^*(x,y)\right] \quad =\rho-i\,I\ , \quad \text{where: } f_x^*=f_x\ , \text{ and: } f_y^*=-f_y$$

$$=\left[\frac{\partial f_x^*}{\partial x}+\frac{\partial f_y^*}{\partial y}\right]+i\left[\frac{\partial f_y^*}{\partial x}-\frac{\partial f_x^*}{\partial y}\right] \quad =\left[\nabla\bullet\mathbf{f}^*\right]+i\left[\nabla\times\mathbf{f}^*\right]_Z$$

Non-analytic potential, force, and source field functions (Excerpts of Unit 1-Ch.10 and AnalyIt)

A general 2D complex field may have:

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Real part: *Poisson scalar source equation* (*charge density* ρ). Imaginary part: *Biot-Savart vector source equation*(*current density I*)

$$\nabla\bullet\mathbf{f}^*=\rho \qquad\qquad \nabla\times\mathbf{f}^*=-I$$

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Field-*vs*-potential has Re and Im parts that are x and y components of grad Φ and curlA_Z from potential $\phi=\Phi+iA$ or $\phi^*=\Phi - iA$.

$$f^*(z,z^*)=2\frac{d\phi^*}{dz^*}=\left[\frac{\partial}{\partial x}+i\frac{\partial}{\partial y}\right]\left(\Phi-iA\right)=f_x^*+if_y^*$$

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$$=\left[\frac{\partial\Phi}{\partial x}+i\frac{\partial\Phi}{\partial y}\right]+\left[\frac{\partial A}{\partial y}-i\frac{\partial A}{\partial x}\right]=\left[\nabla\Phi\right]+\left[\nabla\times\mathbf{A}_Z\right]$$

(For source-free analytic functions these two fields are identical.)

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$$=\left[\frac{\partial\Phi}{\partial x}+i\frac{\partial\Phi}{\partial y}\right]+\left[\frac{\partial A}{\partial y}-i\frac{\partial A}{\partial x}\right]=\left[\nabla\Phi\right]+\left[\nabla\times\mathbf{A}_Z\right]$$

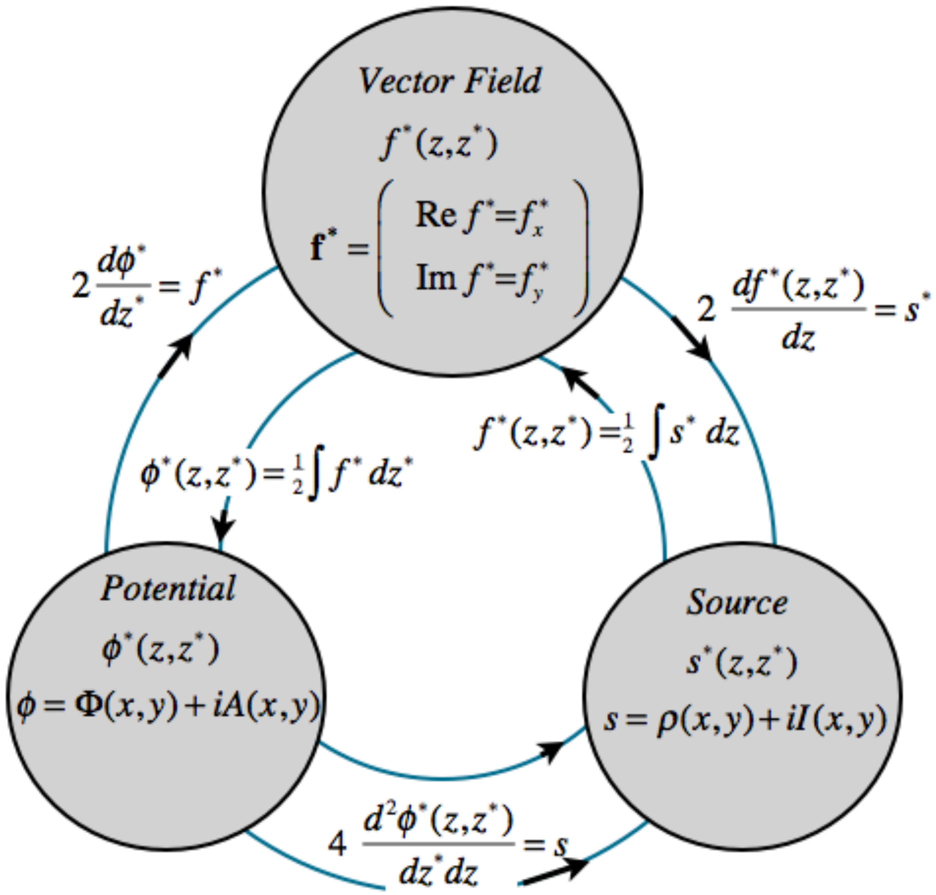
(For source-free analytic functions these two fields are identical.)

Gradient of scalar potential is the *longitudinal field* $\mathbf{f}_\mathbf{L}^*$ and curl of a vector potential is the *transverse field* $\mathbf{f}_\mathbf{T}^*$.

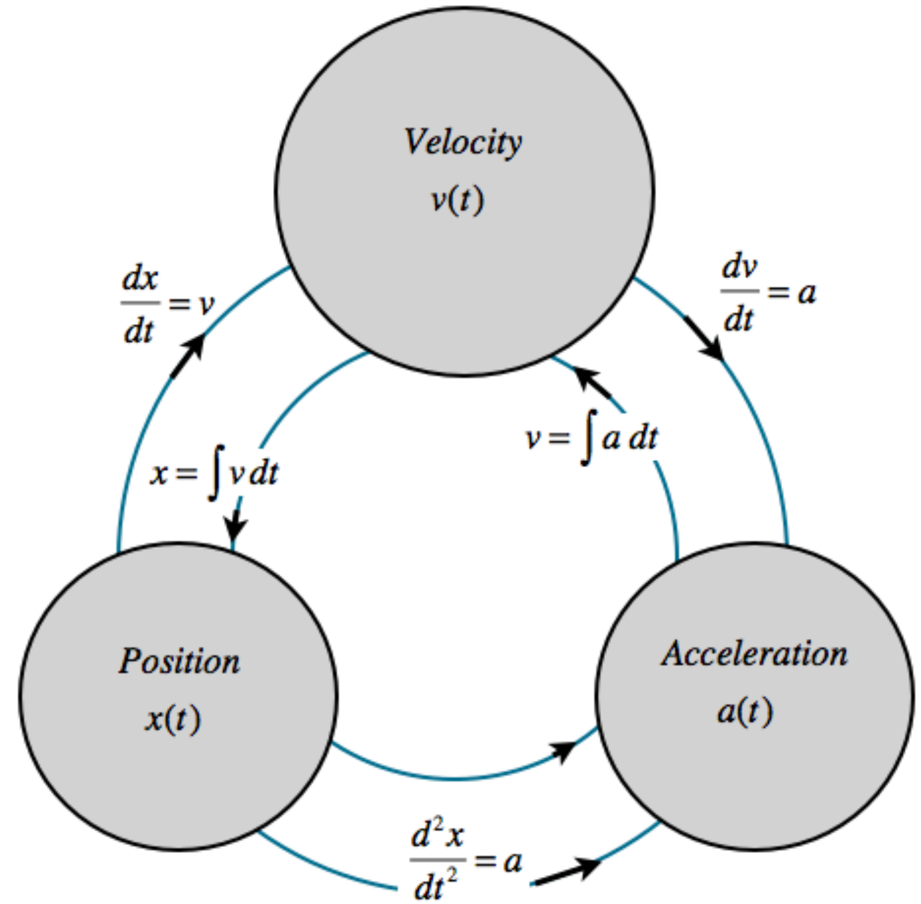
Total field is: $\mathbf{f}^*=\mathbf{f}_\mathbf{L}^*+\mathbf{f}_\mathbf{T}^*$ $\qquad \mathbf{f}_\mathbf{L}^*=\nabla\Phi$ $\qquad \mathbf{f}_\mathbf{T}^*=\nabla\times\mathbf{A}$

Potential, force, and source field equations vs. position, velocity, and acceleration equations

Field equations



Newton equations



Potential and source field theory reduced to sophomore mechanics of 1D-motion!

Example 1 Consider a non-analytic field $f(z) = (z^*)^2$ or $f^*(z) = z^2$.

Non-analytic source s^* is derivative of field f^*

$$s^*(z,z^*) = 2\frac{df^*}{dz} = 4z = 4x + i4y,$$

$$or: \quad \rho = 4x, \quad and: \quad I = -4y.$$

Non-analytic potential ϕ is integral of field f^*

$$\phi(z,z^*) = \frac{1}{2}\int f(z)\,dz = \frac{1}{2}\int (z^*)^2\,dz = \frac{z(z^*)^2}{2} = \frac{(x+iy)(x^2-y^2-i2xy)}{2},$$

$$or: \quad \Phi = \frac{x^3+xy^2}{2}, \quad and: \quad A = \frac{-y^3-yx^2}{2}.$$

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The longitudinal field $\mathbf{f_L}^*$ is quite different from the transverse field $\mathbf{f_T}^*$

$$\mathbf{f}_\mathbf{L}^* = \nabla\Phi = \nabla\left(\frac{x^3+xy^2}{2}\right) = \begin{pmatrix} \frac{\partial\Phi}{\partial x} \\ \frac{\partial\Phi}{\partial y} \end{pmatrix} = \begin{pmatrix} \frac{3x^2+y^2}{2} \\ xy \end{pmatrix},$$

$$\mathbf{f}_\mathbf{T}^* = \nabla\times\mathbf{A} = \nabla\times\left(\frac{-y^3-yx^2}{2}\mathbf{e_z}\right) = \begin{pmatrix} \frac{\partial A}{\partial y} \\ -\frac{\partial A}{\partial x} \end{pmatrix} = \begin{pmatrix} \frac{-3y^2-x^2}{2} \\ xy \end{pmatrix}.$$

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Longitudinal field $\mathbf{f_L}^*$ has no curl and the transverse field $\mathbf{f_T}^*$ has no divergence. Sum field $\mathbf{f}$ has both.

$$\mathbf{f}^* = \mathbf{f}_\mathbf{L}^* + \mathbf{f}_\mathbf{T}^* = \begin{pmatrix} \frac{3x^2+y^2}{2} \\ xy \end{pmatrix} + \begin{pmatrix} \frac{-3y^2-x^2}{2} \\ xy \end{pmatrix} = \begin{pmatrix} x^2-y^2 \\ 2xy \end{pmatrix}, \quad \nabla\bullet\mathbf{f}^* = \nabla\bullet\mathbf{f}_\mathbf{L}^* = 4x = \rho, \quad \nabla\times\mathbf{f}^* = \nabla\times\mathbf{f}_\mathbf{T}^* = 4y = -I.$$

Example 1 Consider a non-analytic field $f(z) = (z^*)^2$ or $f^*(z) = z^2$.

Non-analytic source s^* is derivative of field f^*

$$s^*(z,z^*) = 2\frac{df^*}{dz} = 4z = 4x + i4y,$$

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Non-analytic potential ϕ is integral of field f^*

$$\phi(z,z^*) = \frac{1}{2}\int f(z)\,dz = \frac{1}{2}\int (z^*)^2\,dz = \frac{z(z^*)^2}{2} = \frac{(x+iy)(x^2-y^2-i2xy)}{2},$$

$$or: \quad \Phi = \frac{x^3+xy^2}{2}, \quad and: \quad A = \frac{-y^3-yx^2}{2}.$$

The longitudinal field $\mathbf{f_L}^*$ is quite different from the transverse field $\mathbf{f_T}^*$

$$\mathbf{f}_\mathbf{L}^* = \nabla\Phi = \nabla\left(\frac{x^3+xy^2}{2}\right) = \begin{pmatrix} \frac{\partial\Phi}{\partial x} \\ \frac{\partial\Phi}{\partial y} \end{pmatrix} = \begin{pmatrix} \frac{3x^2+y^2}{2} \\ xy \end{pmatrix}, \qquad \mathbf{f}_\mathbf{T}^* = \nabla\times\mathbf{A} = \nabla\times\left(\frac{-y^3-yx^2}{2}\mathbf{e_z}\right) = \begin{pmatrix} \frac{\partial A}{\partial y} \\ -\frac{\partial A}{\partial x} \end{pmatrix} = \begin{pmatrix} \frac{-3y^2-x^2}{2} \\ xy \end{pmatrix}.$$

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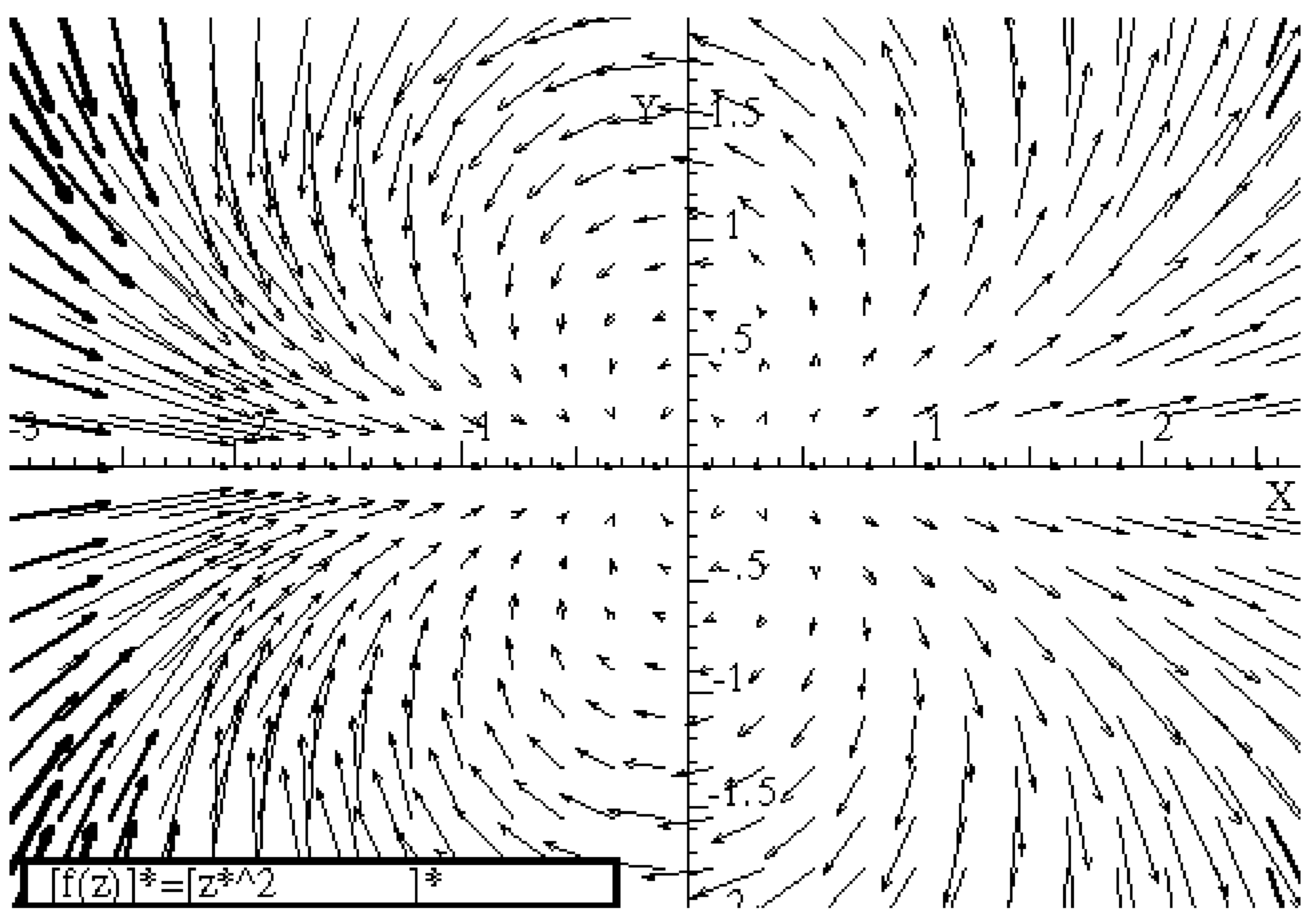
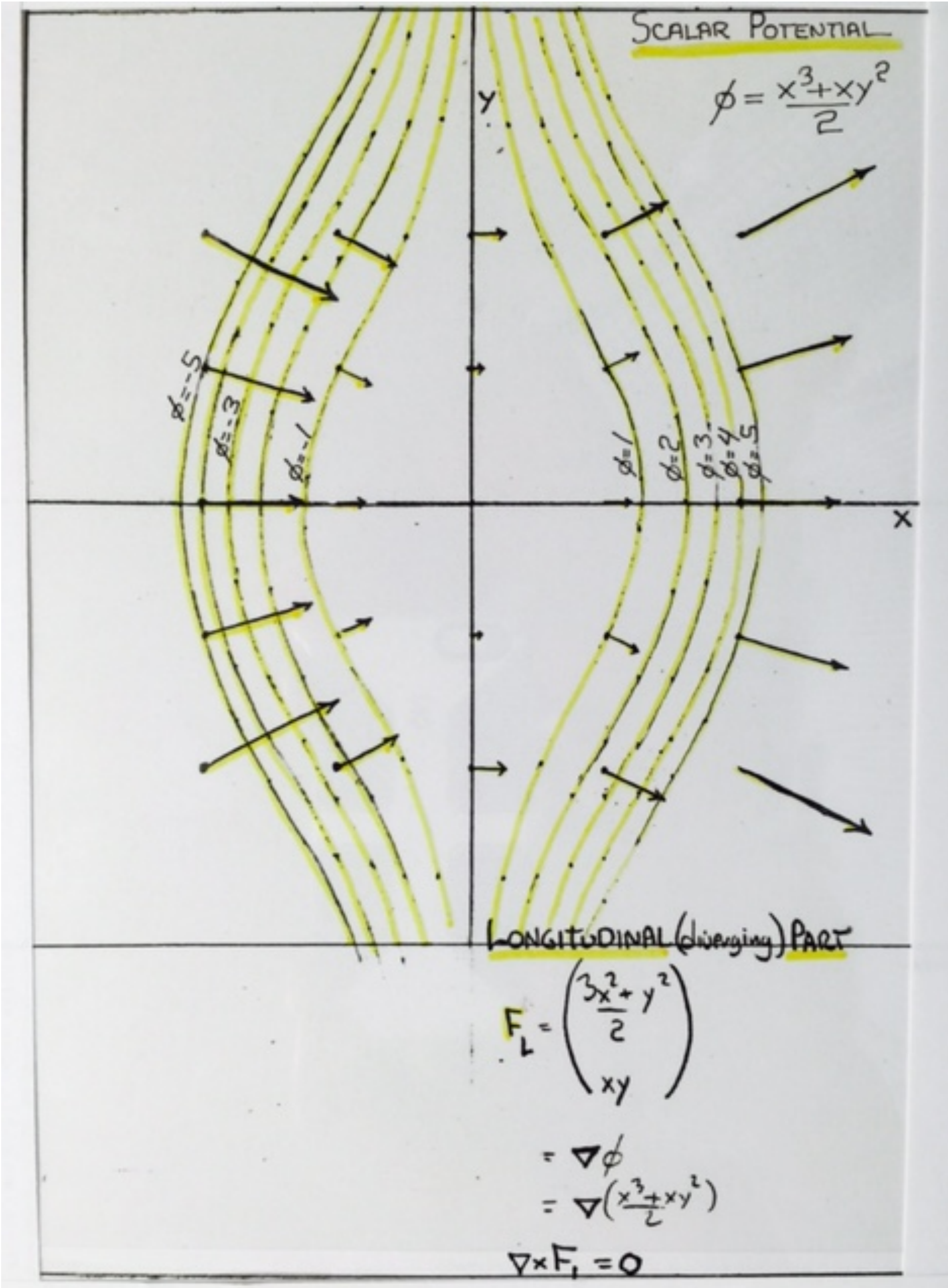
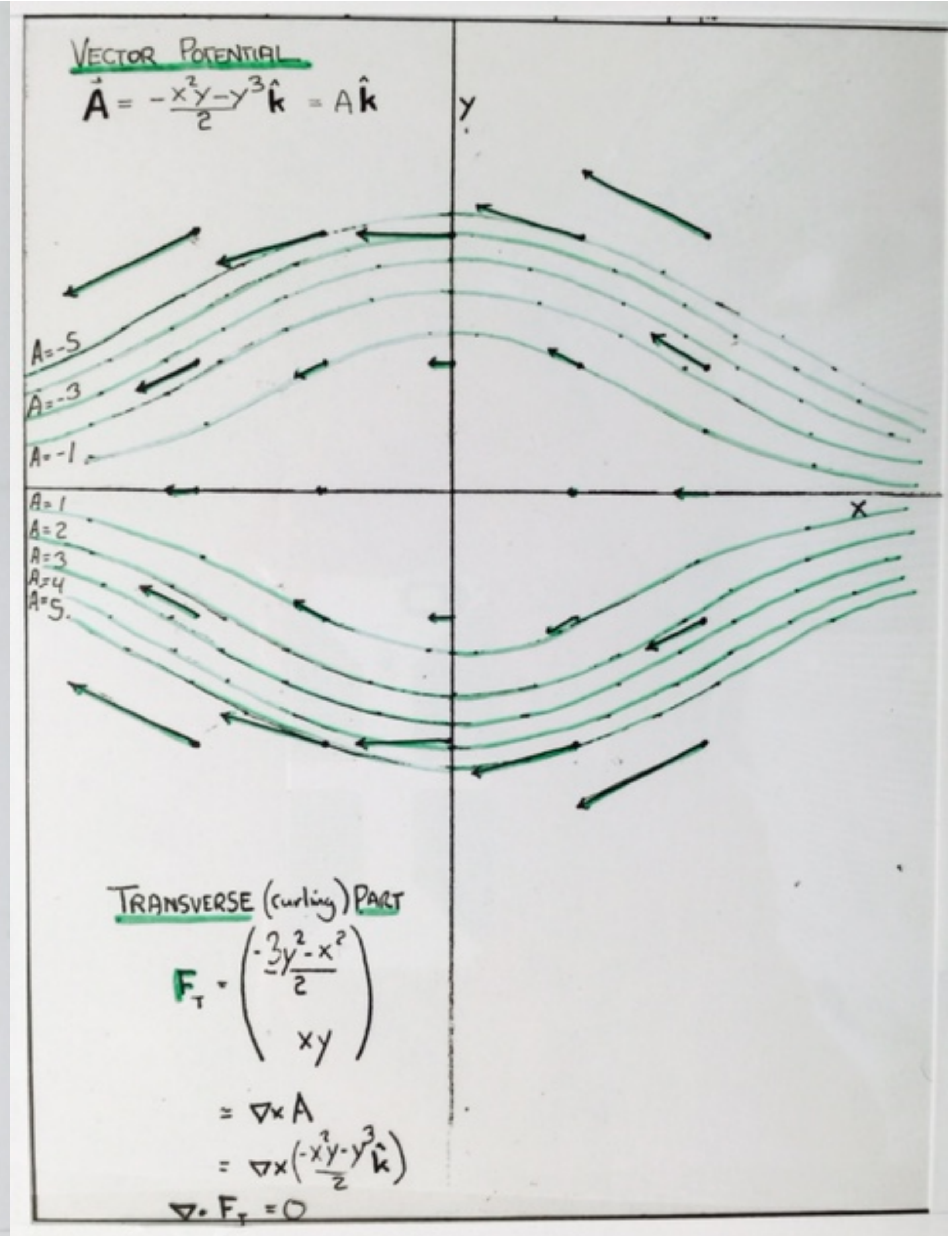
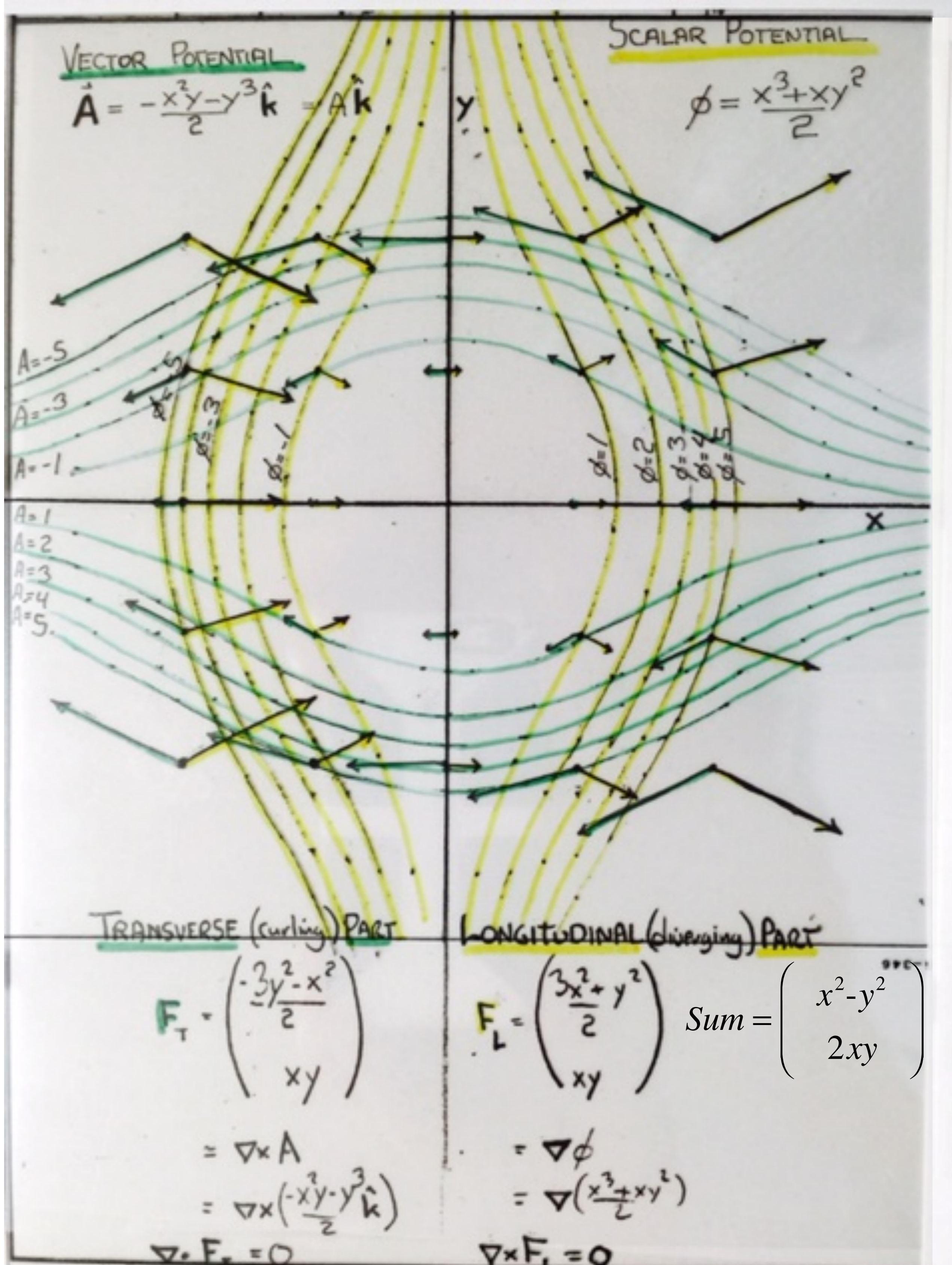


$\mathbf{f}_\mathbf{L}^* + \mathbf{f}_\mathbf{T}^*$

Fig.10.17 Force field vectors for non-analytic function $f(z) = (z^*)^2$


SCALAR POTENTIAL
$\phi = \frac{x^3+xy^2}{2}$
y
x
$\phi=-5$
$\phi=-3$
$\phi=-1$
$\phi=1$
$\phi=2$
$\phi=3$
$\phi=4$
$\phi=5$
LONGITUDINAL (diverging) PART
$F_L = \begin{pmatrix} \frac{3x^2+y^2}{2} \\ xy \end{pmatrix}$
$= \nabla\phi$
$= \nabla\left(\frac{x^3+xy^2}{2}\right)$
$\nabla\times F_L = 0$


VECTOR POTENTIAL
$\vec{A} = -\frac{x^2y-y^3}{2}\hat{k} = A\hat{k}$
y
x
$A=-5$
$A=-3$
$A=-1$
$A=1$
$A=2$
$A=3$
$A=4$
$A=5$
TRANSVERSE (curling) PART
$F_T = \begin{pmatrix} \frac{-3y^2-x^2}{2} \\ xy \end{pmatrix}$
$= \nabla\times A$
$= \nabla\times\left(\frac{-x^2y-y^3}{2}\hat{k}\right)$
$\nabla\cdot F_T = 0$


VECTOR POTENTIAL
$\vec{A} = -\frac{x^2y - y^3}{2}\hat{k} = A\hat{k}$
SCALAR POTENTIAL
$\phi = \frac{x^3 + xy^2}{2}$
y
x
A=-5
A=-3
A=-1
A=1
A=2
A=3
A=4
A=5
$\phi=-3$
$\phi=-1$
$\phi=1$
$\phi=2$
$\phi=3$
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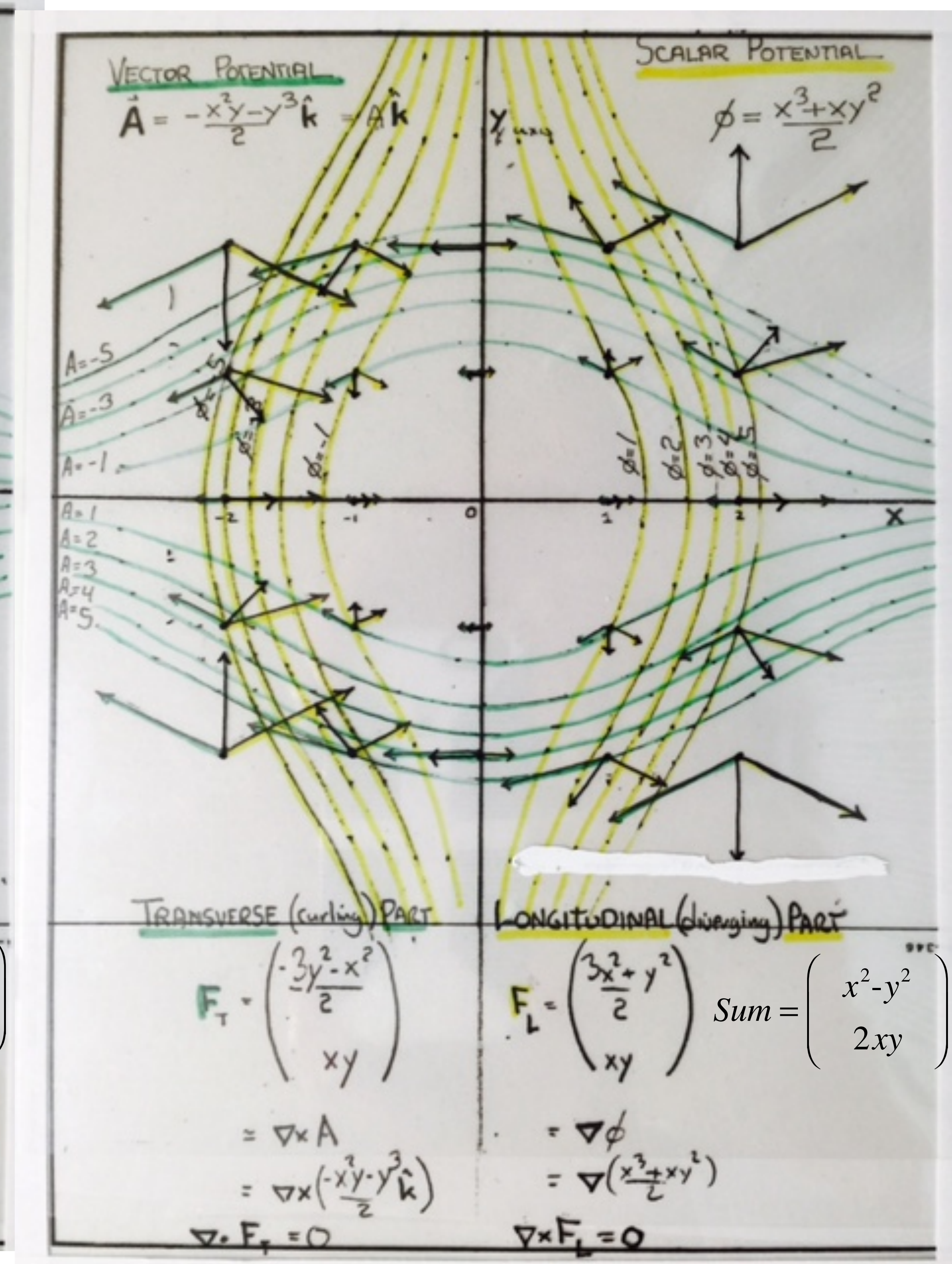
TRANSVERSE (curling) PART
LONGITUDINAL (diverging) PART
$F_T = \begin{pmatrix} \frac{-3y^2 - x^2}{2} \\ xy \end{pmatrix}$
$F_L = \begin{pmatrix} \frac{3x^2 + y^2}{2} \\ xy \end{pmatrix}$
$Sum = \begin{pmatrix} x^2 - y^2 \\ 2xy \end{pmatrix}$
$= \nabla \times A$
$= \nabla \times \left(\frac{-x^2y - y^3}{2}\hat{k}\right)$
$= \nabla\phi$
$= \nabla\left(\frac{x^3 + xy^2}{2}\right)$
$\nabla \cdot F_T = 0$
$\nabla \times F_L = 0$

VECTOR POTENTIAL
$\vec{A} = -\frac{x^2y - y^3}{2}\hat{k} = A\hat{k}$
SCALAR POTENTIAL
$\phi = \frac{x^3 + xy^2}{2}$
y
x
-2
-1
0
1
2
A=-5
A=-3
A=-1
A=1
A=2
A=3
A=4
A=5
$\phi=-1$
$\phi=1$
$\phi=2$
$\phi=3$
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TRANSVERSE (curling) PART
LONGITUDINAL (diverging) PART
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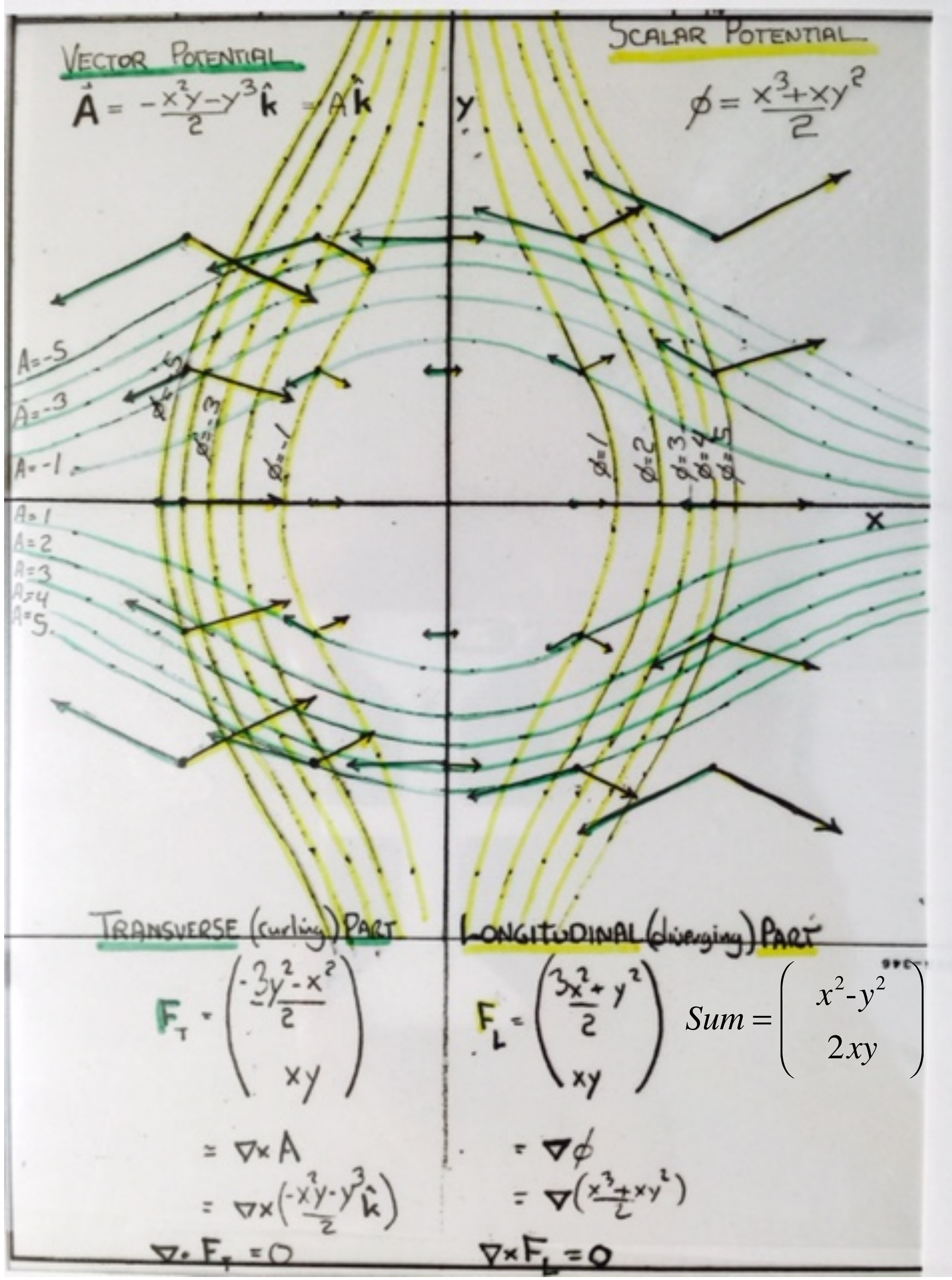
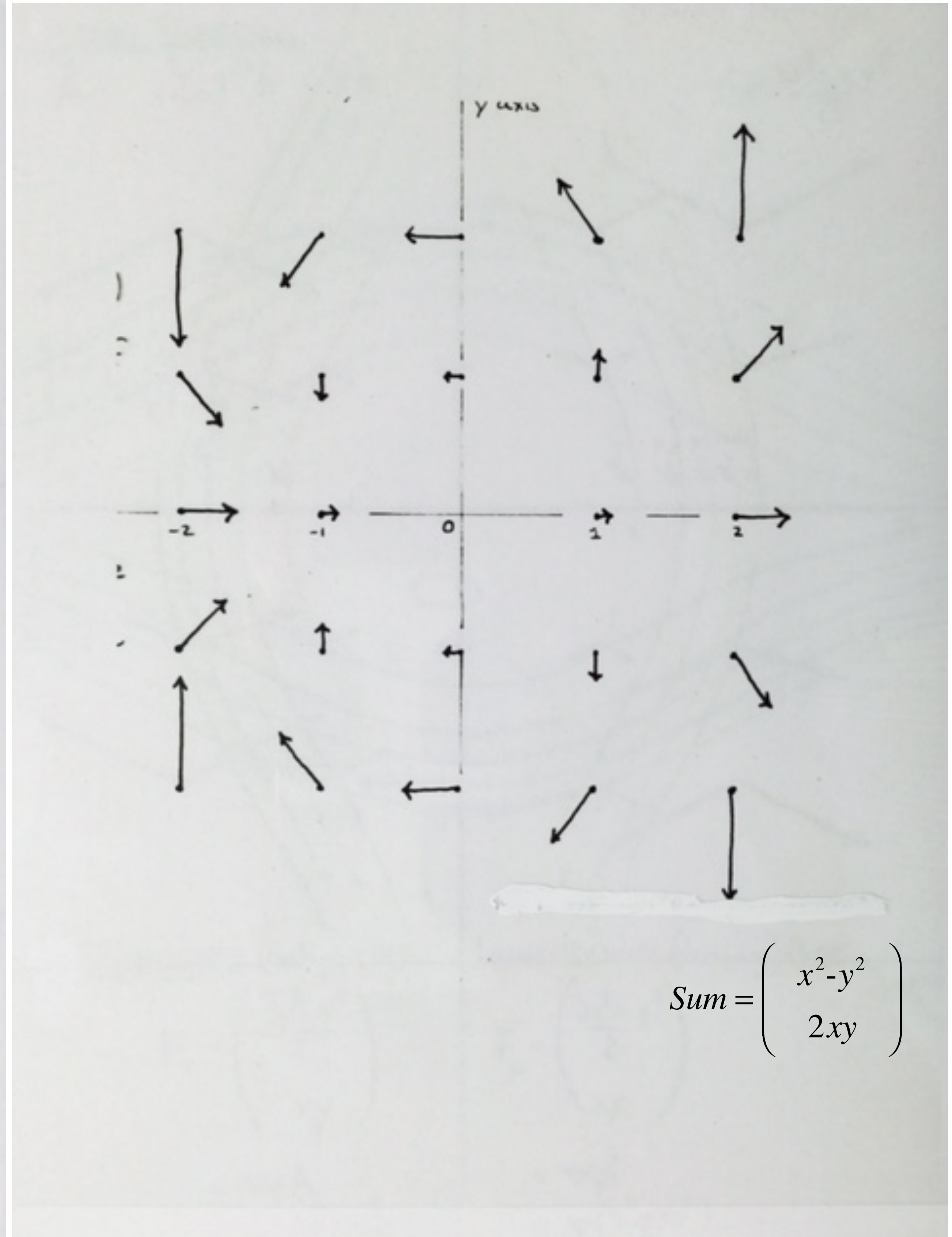
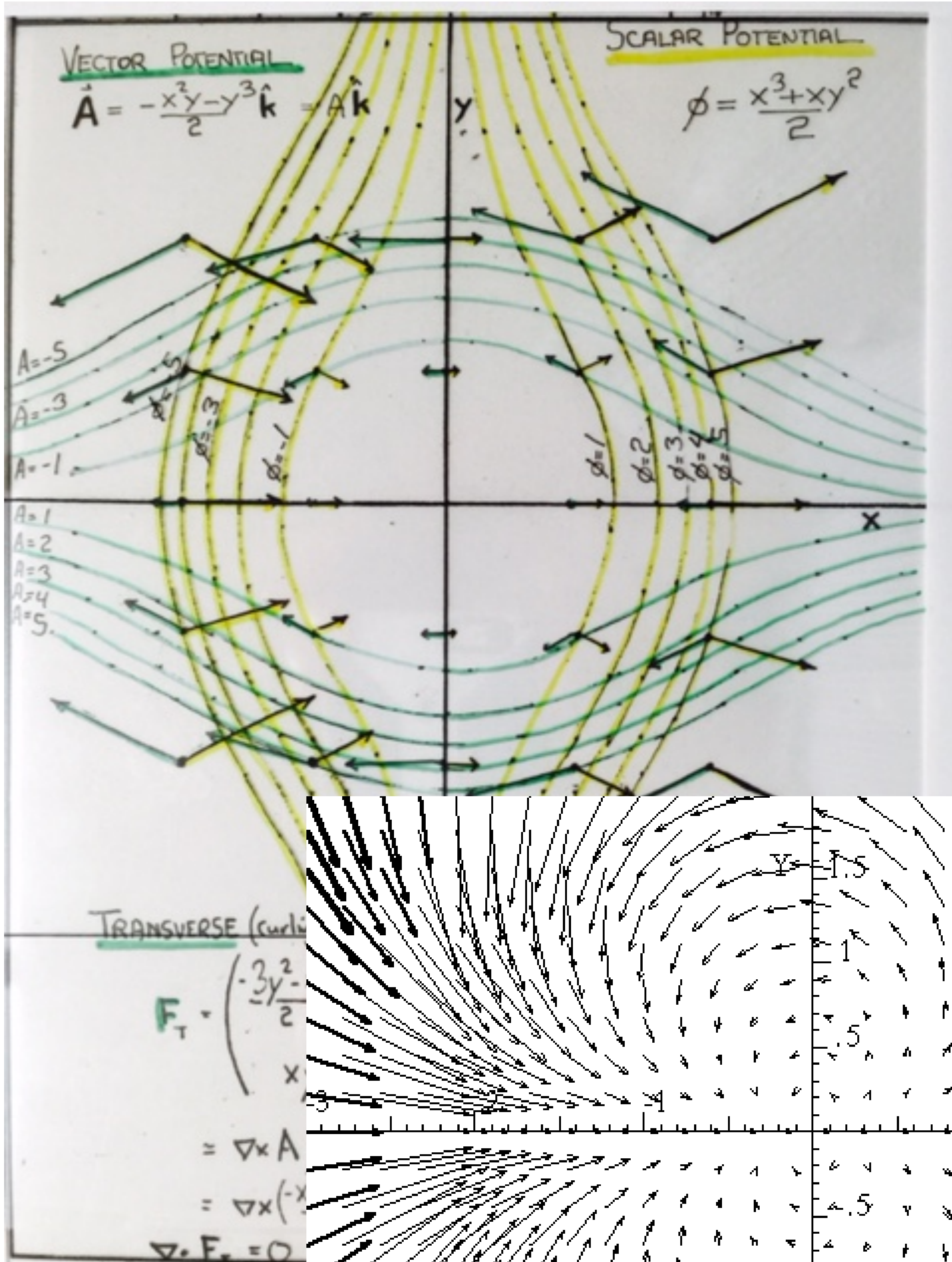
$Sum = \begin{pmatrix} x^2 - y^2 \\ 2xy \end{pmatrix}$
$= \nabla \times A$
$= \nabla \times \left(\frac{-x^2y - y^3}{2}\hat{k}\right)$
$= \nabla\phi$
$= \nabla\left(\frac{x^3 + xy^2}{2}\right)$
$\nabla \cdot F_T = 0$
$\nabla \times F_L = 0$

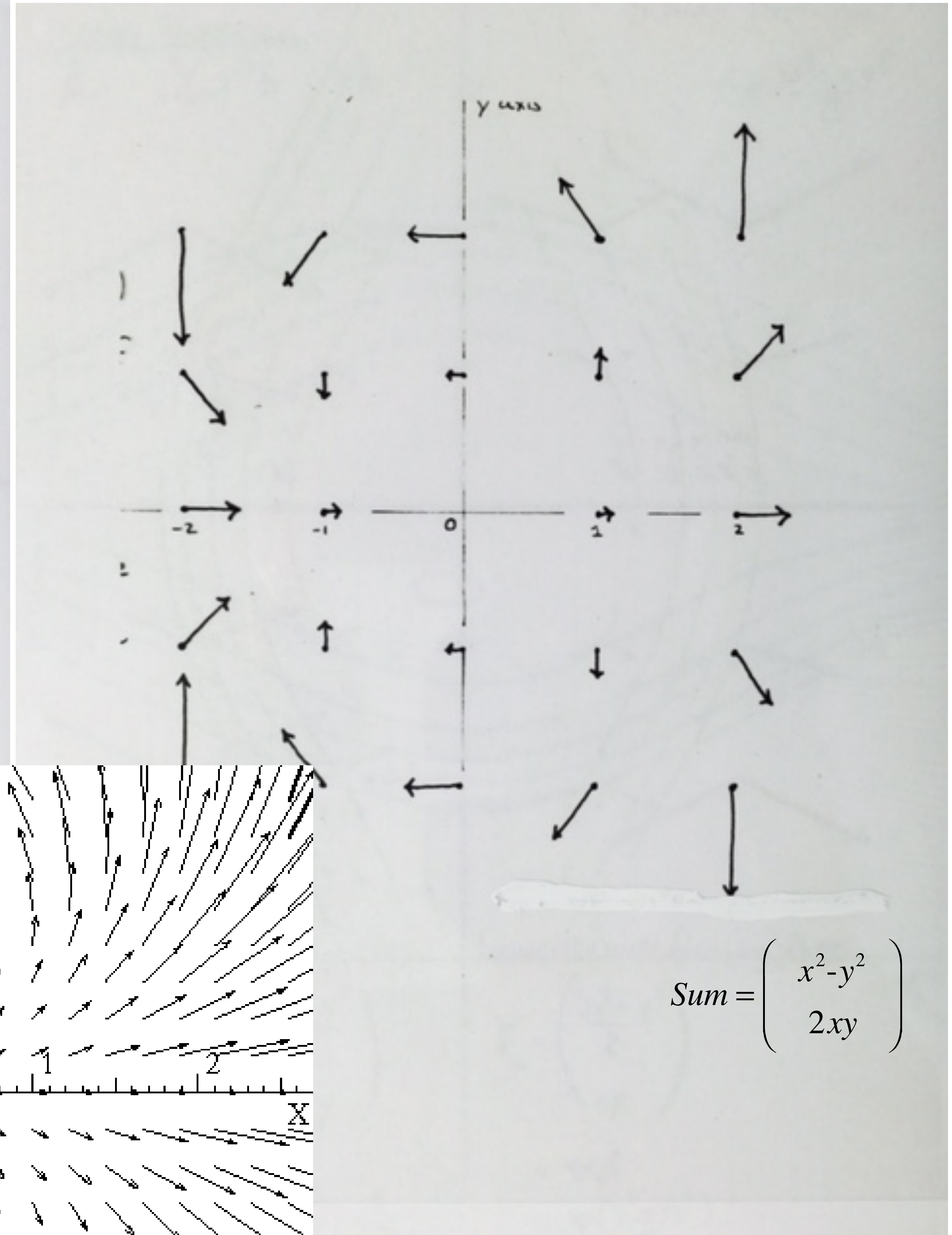


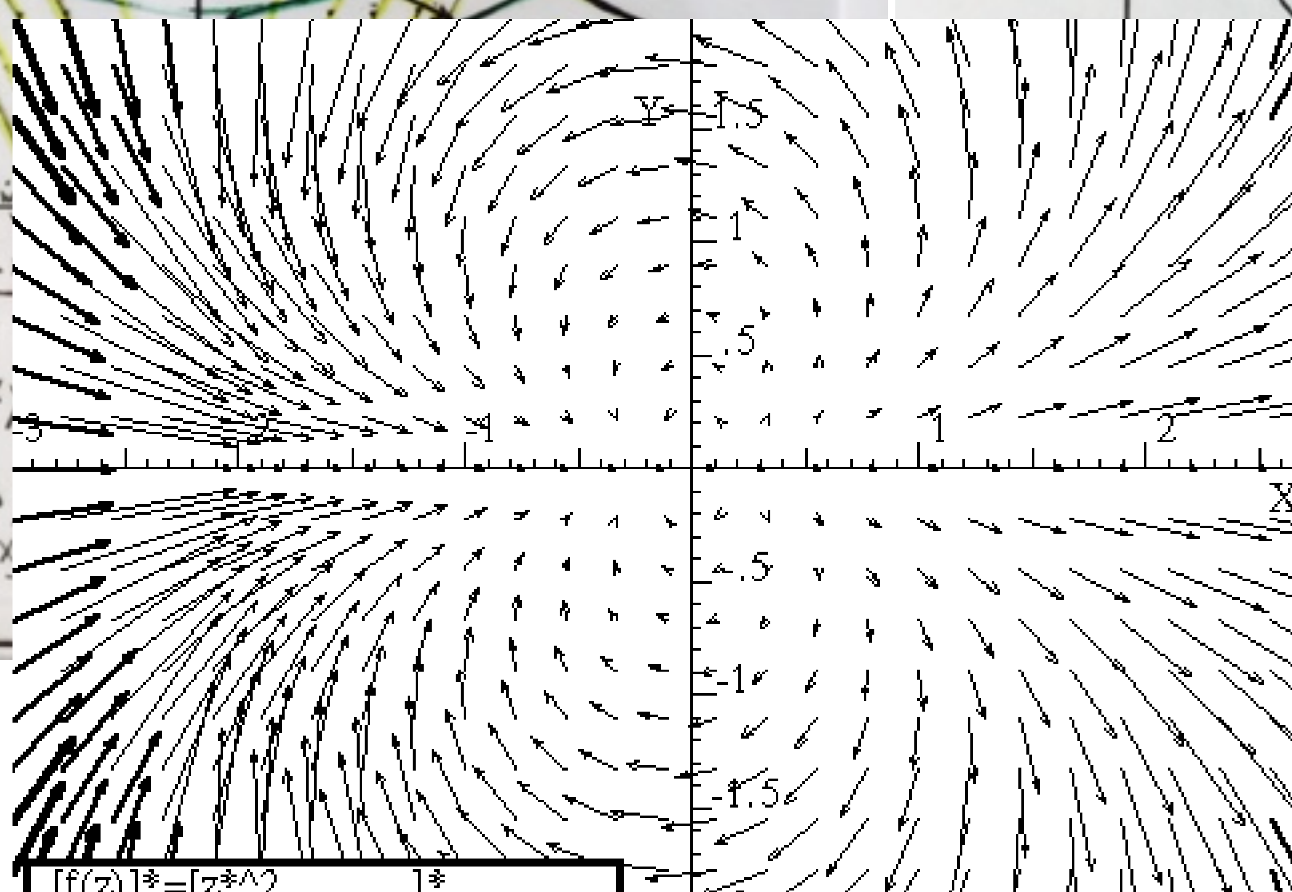
$$Sum = \begin{pmatrix} x^2 - y^2 \\ 2xy \end{pmatrix}$$



$$Sum = \begin{pmatrix} x^2 - y^2 \\ 2xy \end{pmatrix}$$







Lect 13.5 ends here