

Exercise Set 7 Due Wednesday Oct. 16: Based on Unit 1 Chapter 10 and 12 and Lectures 12-13 (2018).

"Professional" Parabolic and Hyperbolic Coordinates (Relates to Fig. 1.10.7)

1. Consider GCC definition: $q^1 = \Phi = x^2 - y^2$, $q^2 = A = 2xy$. Both $(x^1 = x, x^2 = y)$ and $(q^1 = u = \Phi, q^2 = v = A)$ are Orthogonal Curvilinear Coordinates (OCC) related by an analytic function $w = z^2$ or $(u+iv) = (x+iy)^2$. You can treat either one as Cartesian. (This is based on the analytic function $f(z) = 2z$ whose complex potential is $\phi = \underline{\hspace{2cm}}$)

- (a) Plot $(q^1 = u, q^2 = v)$ coordinate curves in a Cartesian $(x^1 = x, x^2 = y)$ graph. Derive the Jacobian, Kajobian, unitary vectors \mathbf{E}_k and \mathbf{E}^k and metric tensors g_{mn} and g^{mn} for this GCC.
- (b) Plot $(x^1 = x, x^2 = y)$ coordinate curves in a Cartesian $(q^1 = u, q^2 = v)$ graph. Derive the Jacobian, Kajobian, unitary vectors and metric tensors for this GCC.

Galaxy Grids

2. Consider the monopole field function $f(z) = e^{i\alpha}/z$ with complex source $e^{i\alpha}$ discussed in Lectures 12-13.

- (a) Derive its $(q^1 = \Phi, q^2 = A)$ scalar and vector potential coordinate functions.
- (b) Plot examples for angle $\alpha = 30^\circ$ and $\alpha = 45^\circ$.

Fun with Exponentials & more from The Story of e

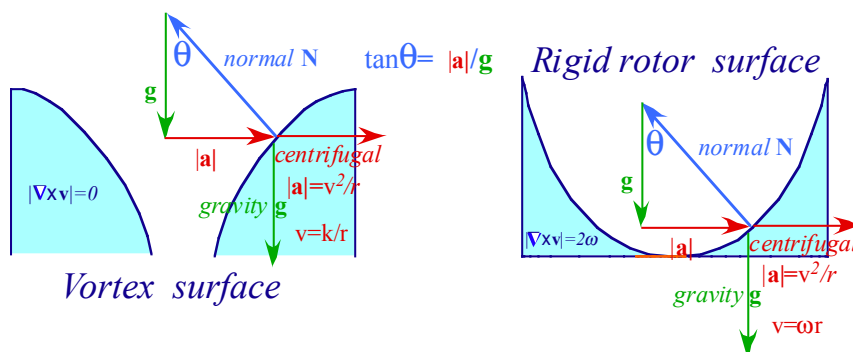
3. Consider a sequence of functions, $f_1(z) = z^z, f_2(z) = z^{f_1(z)} = z^{z^z}, f_3(z) = z^{f_2(z)} = z^{z^{z^z}}, \dots$. The function $f_N(z)$ has a finite limit $f_\infty(z)$ for N approaching infinity if argument z is small enough. ($z=1$ works! But, so does $z=\sqrt{2}$.)

- (a) Find $f_\infty(\sqrt{2}) = \underline{\hspace{2cm}}?$
- (b) Find analytic expression for limiting real z_{max} to give finite $f_\infty(z)$. It involves Euler constant. $e = 2.718281828\dots$

Fun in the bathtub (This has a peculiar connection to "Sophomore-Physics-Earth" potential.)

4. Derive surface shape of rotating fluid subject to constraints on curl function $\nabla \times \mathbf{v}$ for velocity field. From this you should be able to derive surface altitude $S = S(r)$ as a function of radius r by relating balanced forces to differential slope. (Objects floating on these surfaces would not move up or down their $S(r)$ surface.)

- (a) $\nabla \times \mathbf{v} = 0$ (Whirlpool or Vortex) Complex vortex field $f(z^*) = v_x(x,y) + i v_y(x,y) = i/z^*$ has zero z-derivative and zero divergence (flux derivative $\nabla \cdot \mathbf{v} = 0$) and zero curl (circulation derivative $\nabla \times \mathbf{v} = \mathbf{0}$).
- (b) $\nabla \times \mathbf{v} = \text{const.}$ (Rigid rotation) Complex vortex field $f(z) = v_x(x,y) + i v_y(x,y) = i\omega z$ has constant imaginary z-derivative and therefore zero divergence (flux derivative $\nabla \cdot \mathbf{v} = 0$) and constant curl (circulation derivative $\nabla \times \mathbf{v} = \boldsymbol{\omega}$).



- (c) How might the "Sophomore-Physics-Earth" potential be related to a surface whirlpool in deep water

Solutions: Assignment 6 "Professional" Parabolic and Hyperbolic Coordinates

1 Consider the GCC(Cartesian) definition: $q^1 = (x^1)^2 - (x^2)^2$, $q^2 = 2(x^1)(x^2)$. Both $(x^1=x, x^2=y)$ and $(q^1=u, q^2=v)$ are Orthogonal Curvilinear Coordinates (OCC) related by an analytic function $w=z^2$ or $(u+iv)=(x+iy)^2$.

(a) Plot $(q^1=u, q^2=v)$ coordinate curves in a Cartesian $(x^1=x, x^2=y)$ graph. Derive the Jacobian, Kajibian, unitary vectors and metric tensors for this GCC. See below.

(b) Plot $(x^1=x, x^2=y)$ coordinate curves in a Cartesian $(q^1=u, q^2=v)$ graph. Derive the Jacobian, Kajibian, unitary vectors and metric tensors for this GCC.

$w = (u + iv) = z^2 = (x + iy)^2$ is analytic function of z and w

Expansion: $u = x^2 - y^2$ and $v = 2xy$ may be solved using $|w| = |z^2| = |z|^2$

Expansion: $|w| = \sqrt{u^2 + v^2} = x^2 + y^2 = |z|^2$

Solution: $x^2 = \frac{u + \sqrt{u^2 + v^2}}{2}$ $y^2 = \frac{-u + \sqrt{u^2 + v^2}}{2}$

Jacobian follows by inversion:

Easier to get Kajibian first:

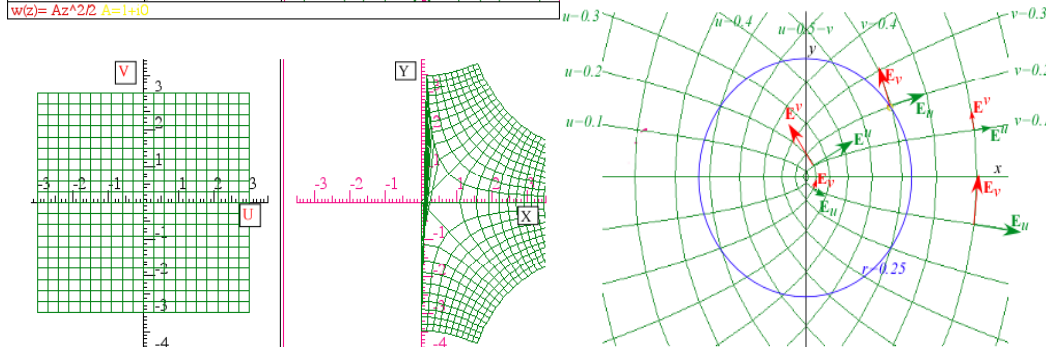
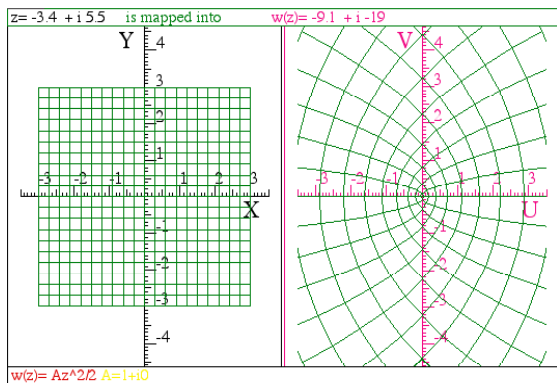
$$\begin{pmatrix} \frac{\partial u}{\partial x} & \frac{\partial u}{\partial y} \\ \frac{\partial v}{\partial x} & \frac{\partial v}{\partial y} \end{pmatrix} = \begin{pmatrix} \bar{\mathbf{E}}^u \\ \bar{\mathbf{E}}^v \end{pmatrix} = \begin{pmatrix} 2x & -2y \\ +2y & 2x \end{pmatrix}$$

$$\begin{pmatrix} \frac{\partial x}{\partial u} & \frac{\partial x}{\partial v} \\ \frac{\partial y}{\partial u} & \frac{\partial y}{\partial v} \end{pmatrix} = \begin{pmatrix} \bar{\mathbf{E}}_u & \bar{\mathbf{E}}_v \end{pmatrix} = \frac{1}{4(x^2 + y^2)} \begin{pmatrix} 2x & +2y \\ -2y & 2x \end{pmatrix}$$

$$\sqrt{2} \begin{pmatrix} \sqrt{u + \sqrt{u^2 + v^2}} & \sqrt{\sqrt{u^2 + v^2} - u} \\ -\sqrt{\sqrt{u^2 + v^2} - u} & \sqrt{u + \sqrt{u^2 + v^2}} \end{pmatrix}$$

$$= \frac{1}{4\sqrt{u^2 + v^2}} \begin{pmatrix} \sqrt{u + \sqrt{u^2 + v^2}} & \sqrt{\sqrt{u^2 + v^2} - u} \\ -\sqrt{\sqrt{u^2 + v^2} - u} & \sqrt{u + \sqrt{u^2 + v^2}} \end{pmatrix}$$

$\det J = 4(x^2 + y^2) = 0$ when: $x = y = 0$
 $g^{11} = \bar{\mathbf{E}}^u \cdot \bar{\mathbf{E}}^u = (2x \ -2y) \cdot (2x \ -2y) = 4x^2 + 4y^2$
 $g^{12} = \bar{\mathbf{E}}^u \cdot \bar{\mathbf{E}}^v = (2x \ -2y) \cdot (2y \ 2x) = 0$
 $g^{22} = \bar{\mathbf{E}}^v \cdot \bar{\mathbf{E}}^v = (2y \ 2x) \cdot (2y \ 2x) = 4x^2 + 4y^2$
 $\det g = g_{11}g_{22} - g_{12}g_{12} = (\det J)^2 = (4x^2 + 4y^2)^2$
 OCC everywhere: $g^{12} = 0$



Note how parabolic covariant vectors grow with distance from origin while contravariant vectors shrink.

Galaxy Grids-Solutions:

2. Consider the monopole field function $f(z)=e^{i\alpha}/z$ with complex source $e^{i\alpha}$ discussed in Lect. 12-13.

(a) Derive its ($q^1=\Phi, q^2=A$) scalar and vector potential coordinate functions.

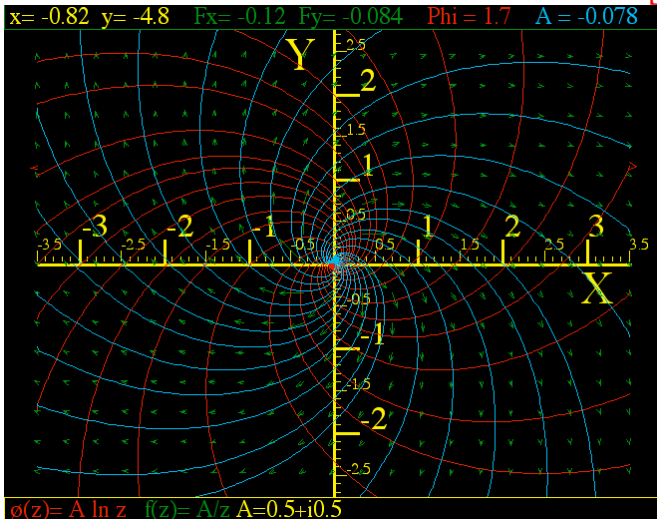
(b) Plot examples for angle $\alpha=30^\circ$ and $\alpha=45^\circ$.

$$\phi = \int f(z) dz = \int (A + iB) / z dz = (A + iB) \ln(z) = (A + iB)(\ln(r) + i\theta) = [A \ln(r) - B\theta] + i[A\theta + B \ln(r)]$$

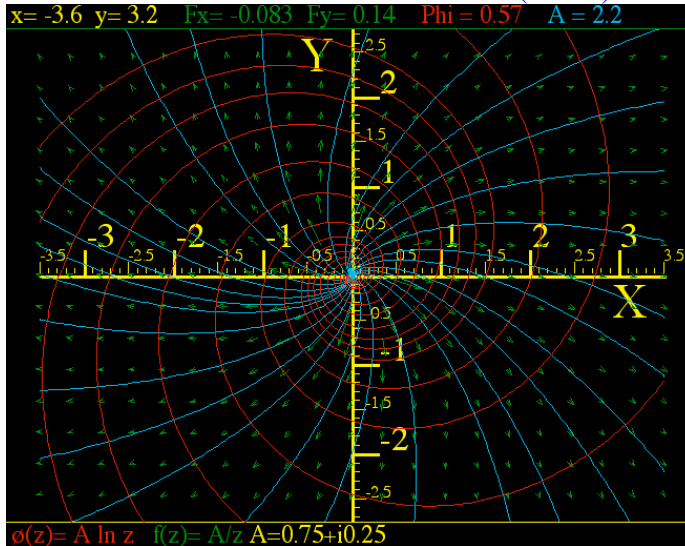
$A = \cos \alpha$ and: $B = \sin \alpha$

$\Phi = [A \ln(r) - B\theta] = const.$

$|A| = [A\theta + B \ln(r)] = const.$



$A=0.5$ and $B=0.5$ $e^{i\alpha} = 0.5 + i 0.5$ has $\alpha = \text{atan}(0.5/0.5) = 45^\circ$.



$A=0.75$ and $B=0.25$ $A=e^{i\alpha} = 0.75 + i 0.25$ has $\alpha = \text{atan}(0.25/0.75) = 18^\circ$.

Fun with Exponents & more of the Story of e

3. Consider a sequence of functions, $f_1(z) = z^z, f_2(z) = z^{f_1(z)} = z^{z^z}, f_3(z) = z^{f_2(z)} = z^{z^{z^z}}, \dots$. The function $f_N(z)$ has a limit for N approaching infinity if argument z is small enough. ($z=1$ works! But, so does $z=\sqrt{2}$.) Find an analytic expression for the limiting real z that involves the Euler constant. $e=2.718281828\dots$

Solution: Assume limiting case can be reached. Then: $f(z) = z^{f(z)}$ or: $z = f^{\frac{1}{f}}$. Let's plot this function: $y(x) = x^{\frac{1}{x}}$
 $1^{\frac{1}{1}} = 1.00, 1.5^{\frac{1}{1.5}} = 1.30, 2^{\frac{1}{2}} = 1.414, 3^{\frac{1}{3}} = 1.442, 4^{\frac{1}{4}} = 1.414, \dots$ Has max when: $\frac{dy}{dx} = 0$ with: $\ln y = \ln x^{\frac{1}{x}} = \frac{\ln x}{x}$

Take implicit derivative of $x \ln y = \ln x$:

then set $\frac{dy}{dx} = 0$:

$$\frac{d}{dx} x \ln y = \frac{d}{dx} \ln x \quad \text{or:} \quad \ln y + \frac{x}{y} \frac{dy}{dx} = \frac{1}{x}, \quad \text{so: } x^{\frac{1}{x}} \equiv y = e^{\frac{1}{x}} = e^{\frac{1}{e}} = 1.44466786$$

Fun in the bathtub (This has a peculiar connection to “Sophomore-Physics-Earth” potential.)
 Solution

4. Derive surface shape of rotating fluid subject to constraints on curl function $\nabla \times \mathbf{v}$ on velocity field.

(a) $\nabla \times \mathbf{v} = 0$ (Whirlpool or Vortex) Complex vortex field $f(z^*) = v_x(x,y) + i v_y(x,y) = i/z^*$ has zero z-derivative and therefore zero divergence (flux derivative $\nabla \cdot \mathbf{v} = 0$) and zero curl (circulation derivative $\nabla \times \mathbf{v} = \mathbf{0}$).

$$0 = \frac{dv(z^*)}{dz} = \frac{1}{2} \left(\frac{\partial}{\partial x} - i \frac{\partial}{\partial y} \right) (v_x + i v_y) = \frac{1}{2} \left(\frac{\partial v_x}{\partial x} + \frac{\partial v_y}{\partial y} \right) + \frac{i}{2} \left(\frac{\partial v_x}{\partial y} - \frac{\partial v_y}{\partial x} \right) = \frac{1}{2} (\nabla \cdot \mathbf{v}) + \frac{i}{2} |\nabla \times \mathbf{v}|_{\perp x,y}$$

for $v(z^*) = \frac{i}{z^*} = \frac{i}{x-iy} = \frac{i}{x-iy} \frac{x+iy}{x+iy} = i \left(\frac{x}{x^2+y^2} + i \frac{y}{x^2+y^2} \right) = v_x + i v_y$ where:

$$v_x = -\frac{y}{r^2} = -\frac{\sin \theta}{r}$$

$$v_y = +\frac{x}{r^2} = +\frac{\cos \theta}{r}$$

(b) $\nabla \times \mathbf{v} = \text{const.}$ (Rigid rotation) Complex vortex field $f(z) = v_x(x,y) + i v_y(x,y) = i\omega z$ has constant imaginary z-derivative and therefore zero divergence (flux derivative $\nabla \cdot \mathbf{v} = 0$) and constant curl (circulation derivative $\nabla \times \mathbf{v} = \boldsymbol{\omega}$).

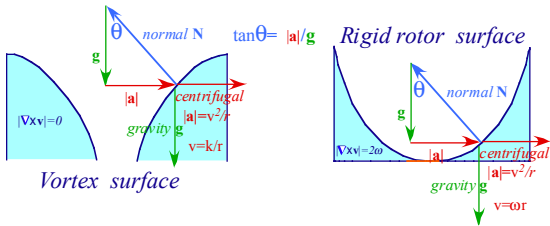
$$i\omega = \frac{d(i\omega z)}{dz} = \frac{1}{2} \left(\frac{\partial}{\partial x} - i \frac{\partial}{\partial y} \right) (v_x + i v_y) = \frac{1}{2} (\nabla \cdot \mathbf{v}) + \frac{i}{2} |\nabla \times \mathbf{v}|_{\perp x,y}$$

So: $\nabla \cdot \mathbf{v} = 0$ and: $|\nabla \times \mathbf{v}|_{\perp x,y} = 2\omega$

for $v(z) = i\omega z = i\omega(x+iy) = v_x + i v_y$ where:

$$v_x = -\omega y = -\omega r \sin \theta$$

$$v_y = +\omega x = +\omega r \cos \theta$$



surface:

$$\tan \theta = \frac{dz}{dr} = \frac{|a|}{g} = \frac{v^2/r}{g} = \begin{cases} (k^2/g)r^{-3} & \text{for vortex } (v = k/r) \\ (\omega^2/g)r^{+1} & \text{for rotor } (v = \omega r) \end{cases}$$

$$z(r) = \begin{cases} -(k^2/2g)r^{-2} & \text{for vortex } (v = k/r) \\ (\omega^2/2g)r^{+2} & \text{for rotor } (v = \omega r) \end{cases}$$

(c) Ideal whirlpools have parabolic core-bottoms starting wherever viscosity wins over curl-free flow.

